

- Photoproduction off nucleons and physics motivations
- The BGOOD Detector
- Overview of analysis and results





BARYON RESONANCES AND MESON PHOTOPRODUCTION OFF NUCLEONS

Physics motivation for the study of photon-induced meson production off nucleons is the investigation of the **excitation spectrum of the nucleon**.

 \rightarrow Static and dynamical properties of baryon resonances are the testing ground for our "understanding" of the quark structure of the matter.

Baryon resonances have been first observed in πN scattering and most states listed in PDG and their properties have been extracted from πN data.

BUT

in the last 30 years large efforts have been done to perform experiments with real and virtual photons at high duty-cycle facilities and dedicated apparata for detection.

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Particle	J^P	Year	Overall status		
				Νγ	Νπ
N(1860)	5/2+	1996	-	-	_
		2018	**	*	**
N(1875)	3/2-	1996	-	_	_
		2018	***	**	**
N(1880)	1/2+	1996	-	-	-
		2018	***	**	*
N(1895)	$1/2^{-}$	1996	-	-	-
		2018	****	****	*
N(1900)	3/2+	1996	**		**
		2018	****	****	**

Strakovski, Progress in Particle and Nuclear Physics 111 (2020) 103752

BARYON RESONANCES AND MESON PHOTOPRODUCTION OFF NUCLEONS

OPEN PROBLEMS



"MISSING RESONANCES"

 \rightarrow effective degrees of freedom inside the nucleon

NATURE OF CERTAIN RESONANCES (not well explained by conventional models) → meson-baryon molecular states



Meson photoproduction allows:

- + access to resonance states coupled to photons
- + polarization observables are accessibles
- + decay amplitudes extraction
- low e.m. cross-sections
- non-resonant contributions are significant
- N.B. isospin filters (η , η ', ω)





W(MeV)

From B.Krusche, Prog. Part. Nucl. Phys. 51 (2003), 399-485

N.B. Photoproduction amplitudes on **proton** and **neutron** of the same meson (in its different charge states) can be decomposed in terms if isospin aamplitudes.

NEED FOR EXPERIMENTS ON NEUTRONS AND PROTONS

AN EXAMPLE: THE $\,\eta\,$ Photoproduction

Let us see the expressions for unpolarized cross section and beam asymmetry in terms of multipoles for η photoproduction case, obtained after truncation at the first orders:

$$\begin{pmatrix} \frac{d\sigma}{d\Omega} \end{pmatrix}_{\text{UNP}} = \frac{q^{\text{CM}}}{2 \, k_{\gamma}^{\text{CM}}} \underbrace{|E_{0+}|^2}_{2 \, k_{\gamma}^{\text{CM}}} Re \left[E_{0+}^* (E_{2-} - 3 \, M_{2-}) \right] + \frac{1}{2 \cos \theta} Re \left[E_{0+}^* (3 E_{1+} + M_{1+} - M_{1-}) \right] + 3 \cos^2 \theta Re \left[E_{0+}^* (E_{2-} - 3 \, M_{2-}) \right]$$

$$\Sigma = - \frac{q^{\text{CM}}}{2 \, k_{\gamma}^{\text{CM}}} \frac{1}{\left(\frac{d\sigma}{d\Omega} \right)_{\text{UNP}}} 3 \, \sin^2 \theta \, \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (2 - 3 \, M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (E_{2-} + M_{2-})} \underbrace{Re \left[E_{0+}^* (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2 \, e^{-2} (E_{2-} + M_{2-})} \underbrace{Re \left[E_{0+} (E_{2-} + M_{2-}) \right]}_{\text{CM}} + \frac{1}{2$$

Since the E_{0+} multipole dominates at threshold, in the differential cross-section it is hard to distinguish the contribution of other resonances.

In beam asymmetry, on the contrary, this contribution is amplified in the interference with the dominating multipole E_{0+} .

NEED FOR POLARIZED BEAMS AND/OR TARGETS EXPERIMENTS

STRANGENESS PHOTOPRODUCTION

In the last years focus has been put on the investigation of unconventional multi-duark states.

Interest has been fueled by discoveries, in the (hidden) charm sector, of: **X(3872)** @ BELLE (**PRL91, 262001, 2003**) **X(3872)** @ Babar (**PRD77 091101, 2008**) $B^+ \rightarrow K^+\pi^+\pi J/\psi$ very close to $D^0\overline{D}^{*0}$ threshold **interpreted as a tetraquark state**

 P_{c} (4380/4450) states @ LHCb (PRL 115, 072001, 2015) observed in M(J/ ψ p) distributions predicted in Wu et al., PRL 105, 232001 (2010) interpreted as pentaquark states





Parallels between the charm and strange sectors are possible and predictions for lighter multi-quark states can be done (exchanging a $c \rightarrow s$), that can be interpreted as meson-baryon molecular-type structures.

→See Katrin Kohl's talk(Oct. 18^{th} @ 15:00)→See Johannes Groβ' talk(Oct 19^{th} @ 15:40)

The same model that predicts *P*_C states (**A. Ramos and E. Oset, Phys. Lett. B 727, (2013) 287**):

- explains the cusp observed in $K^0\Sigma^+$ (CBELSA/TAPS Phys. Lett. B 713 (2012) 180)

- predicts an enhancement in $K^0\Sigma^0$ on the neutron

due to destructive/constructive interference between $K^*\Lambda$ and $K^*\Sigma$ states magnified by a same resonance $N^*(2030)$, a loosely bound $K^*\Sigma$ state.



Same resonance $N^*(2030)$ could be responsible for $\Lambda(1405)$ photoproduction via the triangle singularity (**E. Wang, et al., Phys. Rev. C 95 (2017) 015205**)



R. Di Salvo - NSTAR22 - Oct. 20th, 2022

BGOball**O**pen**D**ipole magnet **Collaboration**

⇒ A collaboration has been established in 2009 (~ 50 collaborators, Germany, Italy, Great Britain, Russia, Switzerland, Ukraine)
- LoI signed and accepted at the PAC Mainz-Bonn 2009, June 25th-26th

- MoU signed on 2010, March 9th

Location: S-Beamline ELSA Accelerator in Bonn

Beam:

polarized and tagged photon beam in the energy range 0.2-3. GeV

Spokespersons: H. Schmieden, P. Levi Sandri











GENERAL FEATURES OF BGOOD

Setup is a **unique combination of:**

- central region: large solid angle calorimeter with excellent energy resolution for photons and good detection efficiency for neutrons; charged particle tracking and identification; neutron/photon discrimination

- high momentum resolution forward tracking of charged particles

and has trigger capabilities (trigger on the released energy in BGO)

Ideally suited for:

- **mixed charged and neutral final state detection** (reconstruction of the complete kinematics of complex final states. both charged and neutrals)

- investigation of kinematical regimes with minimal momentum transfer

- open trigger



BGOOD Detector Performances



EXPERIMENTAL PROGRAM AT BGOOD

Extensive program for:

- strangeness photoproduction → CROSS SECTIONS
 - ${}^{\bullet} \ \gamma \ p \ \rightarrow \ K^{+} \Lambda$
 - $\gamma p \rightarrow K^+ \Sigma^0$
 - $\gamma p \rightarrow K^+ \Lambda(1405) \rightarrow K^+ \pi^0 \Sigma^0$
 - $\gamma n \rightarrow K^0 \Sigma^0$
 - $\gamma n \rightarrow K^+ \Sigma^-$ (preliminary)

- baryon-baryon dynamics in the ud sector

• $\gamma d \rightarrow \pi^0 \pi^0 d$

See Johannes Groβ' talk (Oct 19th@ 15:40)

See Katrin Kohl's talk (Oct 18th @ 15:00)

See Johannes Groβ' talk (Oct 19th@ 15:40)

See Tom Jude's talk (Oct 18th@ 15:40)

- pseudoscalar meson photoproduction → BEAM ASYMMETRIES

- $\gamma p \rightarrow \pi^0 p$ $\gamma n \rightarrow \pi^0 n$
- $\gamma p \rightarrow \eta p$ $\gamma n \rightarrow \eta n$
- $\gamma p \rightarrow \eta' p$



K⁺ detected in forward spectrometer and $\Lambda \rightarrow \pi^0$ n with 2γ in BGO **cosθ**^{K+}_{CM} > **0.9**

• VERY HIGH STATISTICS

→ Diff x-section vs. $\cos\theta_{CM}$ in W bins is flat at threshold and more forward peaked for higher W, consisting with increasing t-channel *K* and *K** exchange processes.

• VERY GOOD ANGULAR RESOLUTION IN THE FORWARD DIRECTION

- \rightarrow sensitivity to high spin states (3rd resonance region, only access to *N**)
- Structure visible at 1720 MeV then data are more flat for energies above 1800 MeV.

Helpful in constraining elementary reaction mechanism for hypernuclei electroproduction¹ at very low Q², an important prerequisite to study hypernuclei.



K⁺ detected in forward spectrometer and $\Sigma \rightarrow \Lambda \gamma$ with γ 's in BGO

cosθ^{K+}_{CM}>0.9

• VERY GOOD ANGULAR RESOLUTION IN THE FORWARD DIRECTION

- **Cusp at W=1900 MeV where the x-sections drops by 1/3 over a 20 MeV range.** Structure nicely resolved in bins of θ_{CM} , (visible for 0.94< $\cos\theta_{CM}$ < 0.96, pronounced for $\cos\theta$ >0.98).
- ALMOST NO EXTRAPOLATION NEEDED TO *t*=*t*_{min}



K⁺Σ⁰(1385) (mass 1873) K⁺Λ(1405) φ p K⁺K⁻p (threshold) *K*Kp (mass 1920MeV) φ p (mass 1950MeV)

- In *dσ/dt* drop very pronounced at W=1900 MeV.
- Drop due to pentaquark X(2000)? (proposed by the Sphinx Collaboration (for $p_T < 141$ MeV) in diffractive $p+C \rightarrow [K^+\Sigma^0] + C$, *Z. Phys. C* 68 (1995) 585))
- But also compatible with re-scattering effects close to open and hidden strange thresholds in a region with various hadronic bound states.



K⁺ detected in forward spectrometer and $\Lambda(1405) \rightarrow \Sigma^0 \pi^0 \rightarrow (\gamma \Lambda) (\gamma \gamma) \rightarrow (\gamma \pi^- p) (\gamma \gamma)$ with 3γ's in BGO 3 charged p.

- **HERMETIC COVERAGE** \rightarrow allows detection of all final state particles
- EXTENSION TO VERY FORWARD ANGLES WITH UNPRECEDENTED ENERGY RESOLUTION

 $\rightarrow \Lambda(1405)$ is assumed to mainly proceed via the kaon t-channel exchange dominated by small momentum transfer, *t*, in particular if it has a relatively loosely bound molecular structure.

→ Another proposed mechanism (**PRC 95 (2017) 015205**) is a triangle singularity driven by the *N**(2030) (bound $K^*\Sigma$) which is expected to vanish once the threshold of free $K^*\Sigma$ production is exceeded.

Cross-section data support the triangle mechanism with $N^*(2030)$ (bound $K^*\Sigma$)

• Measured line-shape of $\Lambda(1405) \rightarrow \Sigma^0 \pi^0$



- **HERMETIC COVERAGE** → allows detection of all final state particles
- Structure may become visible @ 0.2<cosθ_{CM}<0.5,
- → could be interpreted as due to $N^*(2030)$ (dinamically generated $K^*\Sigma$ state) (**Phys. Lett. B 727, (2013) 287**) causing constructive interference between $K^*\Sigma$ and $K^*\Lambda$ states.

N.B. $N^*(2030)$ could also explain the **cusp observed in** $K^0\Sigma^+$ and be responsible for **A(1405)** photoproduction **(supported by BGOOD data**)

In principle other interpretations in terms of conventional states cannot be ruled out.

• Peak at W = 1900 MeV @ $-0.10 < \cos \theta_{CM} < 0.20 \rightarrow \Delta(1900)1/2^{-1}$







d detected in forward spectrometer and $\pi^0 \pi^0 \rightarrow 4\gamma$'s in BGO **cos\theta^d_{CM} >0.8**

X-section peaks at 2650MeV around 4nb/sr

 \rightarrow not compatible with coherent photoproduction (green curve), "Toy pick-up" models (red curve)

→ **compatible** with a 3-isoscalar dibaryon scenario, $d^*(2380)$, 2470 and 2630 MeV/c² reported by ELPH collaboration (PLB 789 (2019) 413) (red curve, right picture), but statistics not sufficient to draw conclusions.

• Evidence of $d^*(2380)$ is also supported by the $2\pi^0$ inv. mass confirming the so-called "ABC" effect.

Extracted:



 $\gamma d \rightarrow d^*(2380) \rightarrow \pi^0 \pi^0 d$

cross section = $(11.3\pm3.2stat\pm2.7 sys)$ nb

 Double peak structure for W>2500 MeV for π⁰d invariant mass, compatible with observations at ELPH of an isovector dibaryon with:

 $M = 2140 MeV/c^2$ Width = 91 MeV/c².

BGOOD: M = 2117 MeV/c² Width = 20MeV/c

The isovector is produced via a sequential decay mechanism isoscalar-isovetor.



BEAM ASYMMETRY IN MESON PHOTOPRODUCTION

With two polarization states \rightarrow systematical errors not depending on the polarization cancel:

 $N_{1,2}(E_v, \theta_{C.M.}) = \#$ events in pol 1/2 = flux of photons in pol 1/2 $F_{1,2}(E_{v})$



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100

150 Phi Meson (dea)

Extraction of the asymmetry

In the hypothesis that the two polarization states have different degrees of polarization, we have:

$$N_{Pol+}^{P} = F_{Pol+}^{P} \left(\frac{d \sigma}{d \Omega} \right)_{UNP} \varepsilon \left(\varphi \right) N_{SC} \left(1 + P_{Pol+}^{P} \Sigma \cdot \cos \left(2 \varphi \right) \right)$$
(1)
$$P = period$$
$$N_{Pol-}^{P} = F_{Pol-}^{P} \left(\frac{d \sigma}{d \Omega} \right)_{UNP} \varepsilon \left(\varphi \right) N_{SC} \left(1 - P_{Pol-}^{P} \Sigma \cdot \cos \left(2 \varphi \right) \right)$$
(2)

If we extract the asymmetry from the usual ratio:

$$\frac{\frac{N_{Pol+}^{P}}{F_{Pol+}^{P}}}{\frac{N_{Pol+}^{P}}{F_{Pol+}^{P}} + \frac{N_{Pol-}^{P}}{F_{Pol-}^{P}}} = \frac{1 + P_{Pol+}^{P} \Sigma \cdot \cos(2\varphi)}{2 + \left(P_{Pol+}^{P} - P_{Pol-}^{P}\right) \Sigma \cdot \cos(2\varphi)}$$
(3)

We get a behaviour which depends on phi also in the denominator. This depends on the fact that the denominator is not proportional to the unpolarized cross section. The unpolarized can be defined as:

$$\frac{N_{UNP}^{P}}{F_{UNP}^{P}} = \frac{1}{2} \frac{1}{P_{Pol+}^{P} + P_{Pol-}^{P}} \left(P_{Pol-}^{P} \frac{N_{Pol+}^{P}}{F_{Pol+}^{P}} + P_{Pol+}^{P} \frac{N_{Pol-}^{P}}{F_{Pol-}^{P}} \right)$$
(4)



$\begin{array}{l} \gamma \, p \ \rightarrow \ \eta \ p \\ Preliminary \\ R. \ Di \ Salvo, \ A. \ Fantini, \ P. \ Levi \ Sandri \end{array}$

Analysed reactions

1) $\eta \rightarrow 2 \gamma$ with: 2γ in BGO + 1 proton in all detector

- 2) $\eta \rightarrow 3\pi^0 \rightarrow 6 \gamma$ with: 6γ in BGO + 1 proton in all detector
- **3)** $\eta \rightarrow \pi^+\pi^-\pi^0$ with 2γ in BGO + 1 proton in all detector + $\pi^+\pi^-$ in BGO
 - → proton and charged pions momenta reconstructed from momentum conservation between initial and final state particles with no hypothesis on the decaying meson





$$\gamma p \rightarrow \eta p$$



$\gamma p \rightarrow \eta p$

GrAAL Ey 1277 MeV BGOOD Eg 1278 MeV

GrAAL Ey 1330 MeV BGOOD Eg 1338 MeV



GrAAL Ey 1381 MeV BGOOD Eg 1400 MeV





GrAAL Ey 1472 MeV BGOOD Eg 1455 MeV







SUMMARY AND CONCLUSIONS

- BGOOD is a unique detector for forward acceptance and hermetic coverage working in the 3rd resonance region.

An extensive program in progress at BGOOD on:
 strangeness photoproduction focused on forward angles and low momentum
 transfer → potential to investigate unconventional states

 $\gamma p \rightarrow K^{+}\Lambda, K^{+}\Sigma^{0}, K^{+}\Lambda(1405)$ $\gamma n \rightarrow K^{0}\Sigma^{0}, K^{+}\Sigma^{-}$

baryon-baryon dynamics in the ud sector $\gamma p \rightarrow d \pi^0 \pi^0$, $d \pi^0 \pi^0 \pi^0$, $d \eta \pi^0$

pseudoscalar meson photoproduction on proton and neutron (and also vector mesons $\boldsymbol{\omega}$ in the future)

- Improving the statistics of some channels (also already published) is possible. Statistics to be analyzed is already available on tapes.

THANK YOU!!

BACKUP SLIDES

EXPERIMENTAL SITUATION FOR η PHOTOPRODUCTION





Measurement of the $\gamma n \rightarrow K^0 \Sigma^0$ differential cross section over the K^{*} threshold K. Kohl *et al.*

arXiv: 2108.13319 and Submitted to EPJA



 $\gamma n \ \rightarrow \ K^0 \Sigma^0 \ \rightarrow \ (\pi^0 \pi^0) \ (\gamma \Lambda) \ \rightarrow (\gamma \gamma \ \gamma \gamma) \ (\gamma \pi^- p)$

with all y's in BGO and most of charged part in BGO/SciRi

HERMETIC COVERAGE

 \rightarrow allows detection of all final state particles

Structure may become visible at $0.2 < \cos\theta_{CM} < 0.5$,

→ could be interpreted as due to a vector meson-baryon dynamically generated $K^*\Sigma$ resonance, the $N^*(2030)$ (**Phys. Lett. B 727, (2013) 287**) causing the constructive interference between $K^*\Sigma$ and $K^*\Lambda$ states.

N.B.

- In the same model $N^*(2030)$ could explain a **cusp observed in** $K^0\Sigma^+$ (destructive interference between $K^*\Lambda$ and $K^*\Sigma$ states amplified by $N^*(2030)$).

- $N^*(2030)$ is supported by BGOOD data in $\Lambda(1405)$ photoproduction

In principle other interpretations in terms of conventional states cannot be ruled out.

La Rivista del Nuovo Cimento (2022) 45:189–276 https://doi.org/10.1007/s40766-021-00028-5



REVIEW PAPER



Trends in particle and nuclei identification techniques in nuclear physics experiments

7.3 BGOOD (MAMBO): neutral particles discrimination with BGO

Fig. 35 MAMBO. a Ratio between the total cluster energy and the cluster multiplicity for neutrons and photons (in black and red the more and less energetic π^0 decay photons, respectively). b Ratio between the maximum energy released in a crystal and the total cluster energy for neutrons with multiplicity > 2 and photons (sum of red and black distributions from a). The applied graphical cut is the red contour. c photon TOF versus neutron TOF distribution for selected events, the applied cut is presented in red. Ambiguity between two neutral signals with the same multiplicity is solved assigning the neutron tag to the signal with larger TOF (see 45° black line in the picture)

NEUTRON/PHOTON DISCRIMINATION IN BGO

GOAL OF THE JOB

Find criteria for identification and separation neutron/photon in BGO

N.B. Sampling ADC used for BGO crystals readout allow also time measurement with 1.5ns resolution.

TECHNIQUE

Start from a clean sample of selected events with a neutron and a photon in the final state. The selection is based only on kinematical criteria.

Derive criteria for neutron/photon discrimination from this clean sample.

CHOSEN REACTION: $\gamma n \rightarrow \pi^0 n$

SELECTION CRITERIA:

- 3 "neutral" clusters in BGO (i.e. with no signal in barrel/mwpc in geometrical coincidence), arranged into 3 possible combinations







1) Ratio between the cluster total energy and cluster multipl:

$$R = \frac{ETOT_{Clus}}{Mult_{Clus}}$$

R vs. Mult.

Photons: R<100 MeV/crystal Neutrons:R<400 MeV/crystal

$$R' = \frac{E_{Cryst, MAX}}{ETOT_{Clus}}$$

R' vs. ETOT_{Clus}

3) neutron tof (≈ 1 ns) is in average higher than γ (≈ -0.5 ns)



TAGGER DESIGN



- 120 Channels (covering 0.1-0.9 E₀)
- Energy resolution: 10-40 MeV (0.55%E₀-2.15%E₀)
- Beam Intensity: 5 · 10⁷ s⁻¹





Siebke, G.

- High energy e⁻(→ low en. γ): focal plane is not accessible split into 54 horizontal scint. (10-32% E₀) and 66 vertical ones (32-90% E₀)
- trigger on double and triple coincidences

Time Distribution of clusters in BGO

Clusters from a photon

Clusters from a proton



K^+ identification with the BGO ball

- Time delayed, K^+ weak decay within the crystals of the BGO ball
- Technique proven with the Crystal Ball, Mainz



K^+ identification with the BGO ball



Preliminary experimental data (no charged particle ID used):



FORWARD SPECTROMETER: BEFORE DIPOLE

ΜΟΜΟ



Successful commissioning with beam tests in Feb.-March 2012 and June 2012

Scintillation fiber detector (\oslash 44cm) 672 total channels 3 layers of 2x112 parallel fibers (\oslash 2.5mm, Δ x=1.5mm) Each layer rotated of 60° w.r.t. each other Read-out 16 channel PM Central hole (\oslash 5cm) for the passage of beam

Scintillation fiber detector (66cmx51cm) 640 total channels 2 double layers x/y (352 and 288 rot. 90°) Each single layer consists of two fiber arrays (2.0mm track length for crossing particles) Read-out 16 channel PM Central hole (⊘4cm) for the passage of beam





660 mm (x-Layer Hamamatsu PMTs





FORWARD SPECTROMETER: AFTER DIPOLE

DRIFT CHAMBERS

Two sets of 4 double layers X(vert.), Y(90°), U(+9°), V(-9°) Sensitive area 1.2 x 2.4 m² X = horizontal Y = vertical U = 99° w.r.t. X V = 81° w.r.t. X Hexagonally shaped drift cells (inner radius 8.5mm) Distance from target = $3.8 \text{ m} \div 4.5 \text{ m}$ Gas mixture: 70% Ar, 30% Co₂ Central insensitivity spot (5x5cm²)



PNPI- Gatchina

TOF WALLS

- > 2 walls
- > 14 bars vert., 8 bars hor.
- Scintillator dimensions:
- > 3400mm x 210mm x 60mm (horizontal bars)
- 2700mm x 200mm x 45mm (vertical bars)
- > Time resolution ~500ps
- Upgrade using former GRAAL ToF detectors (time resolution ~200ps) JUNE 2014



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O. Jahn – V. Vegna

Missing Mass from Protons Detected in the Forward Spectrometer



FLUX MONITORS

Two flux monitors:

GIM (Gamma Intensity Monitor): efficiency close to 100% Lead Glass Uses the Cerenkov effect to separate charged particles from the e.m. showers generated by photons One single lead block 2" PM

FLUMO (Flux Monitor)

Three scintillators + a Copper foil between the first and second scintillators Low intensity runs to extract the efficiency of the FluMo







