# Tensions in cosmology and implications for the standard model

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# The **ACDM** model

Among a number of cosmological models introduced in the literature, the Lambda Cold Dark Matter (ACDM) cosmological model is the mathematically simplest model, and has now practically been selected as the "standard" cosmological scenario, because it provides a remarkable description of a wide range of astrophysical and cosmological probes.

However, despite its marvelous fit to the available observations, <u>ACDM harbours large areas of phenomenology and ignorance</u>. For example, it still cannot explain key pillars in our understanding of the structure and evolution of the Universe, namely, <u>Dark Energy, Dark Matter and Inflation</u>.

# The $\Lambda CDM$ model

In the ACDM paradigm these three pillars are our simplest guesses.

- DE assumes its simplest form, that is the cosmological constant, without any strong physical basis.
- The nature of DM is still a mystery except for its gravitational interaction, as suggested by the observational evidence. We know, however, that DM is essential for structure formation in the late Universe, so most of it must be pressure-less, cold, and stable on cosmological time scales. Moreover, despite the significant efforts in the last decades to investigate DM and the physics beyond the SM of particle physics, in laboratory experiments and from devised astrophysical observations, no evidence pointing to the dark matter particle has been found.
- Finally, even though the theory of inflation has solved a number of crucial puzzles related to the early evolution of the Universe, in the standard model this is given by a single, minimally coupled, slow-rolling scalar field.

# The **ACDM** model

Therefore, the 6 parameter ACDM model lacks the deep underpinnings a model requires to approach fundamental physics laws. It can be rightly considered, at best, as an approximation of an underlying physical theory, yet to be discovered. In this situation, we must be careful not to cling to the model too tightly or to risk missing the appearance of departures from the paradigm.

With the improvement of the number and the accuracy of the observations, deviations from ACDM may be expected. And, actually, discrepancies among key cosmological parameters of the models have emerged with different statistical significance.

While some proportion of these discrepancies may have a systematic origin, their persistence across probes should require multiple and unrelated errors, strongly hinting at cracks in the standard cosmological scenario and the necessity of new physics.

These tensions can indicate a failure of the canonical ACDM model.

Submitted to the Proceedings of the US Community Study on the Future of Particle Physics (Snowmass 2021)

### Cosmology Intertwined: A Review of the Particle Physics, Astrophysics, and Cosmology Associated with the Cosmological Tensions and Anomalies

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The current cosmological tensions and anomalies are the argument of Review Paper we submitted for the SNOWMASS call, that includes contributions from more than 200 people, who participated in brainstorming sessions from August 2020, and provided feedback via regular Zoom seminars and meetings.

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## **Cosmology Intertwined I: Perspectives for the Next Decade**

Eleonora Di Valentino, Luis A. Anchordoqui, Ozgur Akarsu, Yacine Ali-Haimoud, Luca Amendola, Nikki Arendse, Marika Asgari, Mario Ballardini, Spyros Basilakos, Elia Battistelli, Micol Benetti, Simon Birrer, François R. Bouchet, Marco Bruni, Erminia Calabrese, David Camarena, Salvatore Capozziello, Angela Chen, Jens Chluba, Anton Chudaykin, Eoin Ó Colgáin, Francis-Yan Cyr-Racine, Paolo de Bernardis, Javier de Cruz Pérez, Jacques Delabrouille, Jo Dunkley, Celia Escamilla-Rivera, Agnès Ferté, Fabio Finelli, Wendy Freedman, Noemi Frusciante, Elena Giusarma, Adrià Gómez-Valent, Will Handley, Ian Harrison, Luke Hart, Alan Heavens, Hendrik Hildebrandt, Daniel Holz, Dragan Huterer, Mikhail M. Ivanov, Shahab Joudaki, Marc Kamionkowski, Tanvi Karwal, Lloyd Knox, Suresh Kumar, Luca Lamagna, Julien Lesgourgues, Matteo Lucca, Valerio Marra, Silvia Masi, Sabino Matarrese, Arindam Mazumdar, Alessandro Melchiorri, Olga Mena, Laura Mersini-Houghton, Vivian Miranda, Cristian Moreno-Pulido, David F. Mota, Jessica Muir, Ankan Mukherjee, Florian Niedermann, Alessio Notari, Rafael C. Nunes, Francesco Pace, Andronikos Paliathanasis, Antonella Palmese, Supriya Pan, Daniela Paoletti, Valeria Pettorino, Francesco Piacentini, Vivian Poulin, Marco Raveri, Adam G. Riess, Vincenzo Salzano, Emmanuel N. Saridakis, Anjan A. Sen, Arman Shafieloo, Anowar J. Shajib, Joseph Silk, Alessandra Silvestri, Martin S. Sloth, Tristan L. Smith, Joan Solà, Carsten van de Bruck, Licia Verde, Luca Visinelli, Benjamin D. Wandelt, Deng Wang, Jian-Min Wang, Anil K. Yadav, Weiqiang Yang

The standard  $\Lambda$  Cold Dark Matter cosmological model provides an amazing description of a wide range of astrophysical and astronomical data. However, there are a few big open questions, that make the standard model look like a first-order approximation to a more realistic scenario that still needs to be fully understood. In this Letter of Interest we will list a few important goals that need to be addressed in the next decade, also taking into account the current discordances present between the different cosmological probes, as the Hubble constant  $H_0$  value, the  $\sigma_8 - S_8$  tension, and the anomalies present in the Planck results. Finally, we will give an overview of upgraded experiments and next-generation space-missions and facilities on Earth, that will be of crucial importance to address all these questions.

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## **Cosmology Intertwined II: The Hubble Constant Tension**

Eleonora Di Valentino, Luis A. Anchordoqui, Ozgur Akarsu, Yacine Ali-Haimoud, Luca Amendola, Nikki Arendse, Marika Asgari, Mario Ballardini, Spyros Basilakos, Elia Battistelli, Micol Benetti, Simon Birrer, François R. Bouchet, Marco Bruni, Erminia Calabrese, David Camarena, Salvatore Capozziello, Angela Chen, Jens Chluba, Anton Chudaykin, Eoin Ó Colgáin, Francis-Yan Cyr-Racine, Paolo de Bernardis, Javier de Cruz Pérez, Jacques Delabrouille, Jo Dunkley, Celia Escamilla-Rivera, Agnès Ferté, Fabio Finelli, Wendy Freedman, Noemi Frusciante, Elena Giusarma, Adrià Gómez-Valent, Julien Guy, Will Handley, Ian Harrison, Luke Hart, Alan Heavens, Hendrik Hildebrandt, Daniel Holz, Dragan Huterer, Mikhail M. Ivanov, Shahab Joudaki, Marc Kamionkowski, Tanvi Karwal, Lloyd Knox, Suresh Kumar, Luca Lamagna, Julien Lesgourgues, Matteo Lucca, Valerio Marra, Silvia Masi, Sabino Matarrese, Arindam Mazumdar, Alessandro Melchiorri, Olga Mena, Laura Mersini-Houghton, Vivian Miranda, Cristian Moreno-Pulido, David F. Mota, Jessica Muir, Ankan Mukherjee, Florian Niedermann, Alessio Notari, Rafael C. Nunes, Francesco Pace, Andronikos Paliathanasis, Antonella Palmese, Supriya Pan, Daniela Paoletti, Valeria Pettorino, Francesco Piacentini, Vivian Poulin, Marco Raveri, Adam G. Riess, Vincenzo Salzano, Emmanuel N. Saridakis, Anjan A. Sen, Arman Shafieloo, Anowar J. Shajib, Joseph Silk, Alessandra Silvestri, Martin S. Sloth, Tristan L. Smith, Joan Solà, Carsten van de Bruck, Licia Verde, Luca Visinelli, Benjamin D. Wandelt, Deng Wang, Jian-Min Wang, Anil K. Yadav, Weiqiang Yang

The current cosmological probes have provided a fantastic confirmation of the standard  $\Lambda$  Cold Dark Matter cosmological model, that has been constrained with unprecedented accuracy. However, with the increase of the experimental sensitivity a few statistically significant tensions between different independent cosmological datasets emerged. While these tensions can be in portion the result of systematic errors, the persistence after several years of accurate analysis strongly hints at cracks in the standard cosmological scenario and the need for new physics. In this Letter of Interest we will focus on the 4.4 $\sigma$  tension between the Planck estimate of the Hubble constant  $H_0$  and the SH0ES collaboration measurements. After showing the  $H_0$  evaluations made from different teams using different methods and geometric calibrations, we will list a few interesting new physics models that could solve this tension and discuss how the next decade experiments will be crucial.

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Astrophysics > Cosmology and Nongalactic Astrophysics [Submitted on 25 Aug 2020 (v1), last revised 13 Oct 2020 (this version, v4)]	
<b>Cosmology Intertwined III:</b> $f\sigma_8$ and $S_8$ Eleonora Di Valentino, Luis A. Anchordogui, Ozgur Akarsu, Yacine Ali-Haimoud, Luca Amendola, Nikki Arendse, Marika Asgari, Mario Ballardini, Sp	vros
Basilakos, Elia Battistelli, Micol Benetti, Simon Birrer, François R. Bouchet, Marco Bruni, Erminia Calabrese, David Camarena, Salvatore Capozziello, J Jens Chluba, Anton Chudaykin, Eoin Ó Colgáin, Francis-Yan Cyr-Racine, Paolo de Bernardis, Javier de Cruz Pérez, Jacques Delabrouille, Jo Dunkley, Escamilla-Rivera, Agnès Ferté, Fabio Finelli, Wendy Freedman, Noemi Frusciante, Elena Giusarma, Adrià Gómez-Valent, Will Handley, Ian Harrison, Alan Heavens, Hendrik Hildebrandt, Daniel Holz, Dragan Huterer, Mikhail M. Ivanov, Shahab Joudaki, Marc Kamionkowski, Tanvi Karwal, Lloyd Know	Angela Chen, Celia Luke Hart, x, Suresh

Kumar, Luca Lamagna, Julien Lesgourgues, Matteo Lucca, Valerio Marra, Silvia Masi, Sabino Matarrese, Arindam Mazumdar, Alessandro Melchiorri, Olga Mena, Laura Mersini–Houghton, Vivian Miranda, Cristian Moreno–Pulido, David F. Mota, Jessica Muir, Ankan Mukherjee, Florian Niedermann, Alessio Notari, Rafael C. Nunes, Francesco Pace, Andronikos Paliathanasis, Antonella Palmese, Supriya Pan, Daniela Paoletti, Valeria Pettorino, Francesco Piacentini, Vivian Poulin, Marco Raveri, Adam G. Riess, Vincenzo Salzano, Emmanuel N. Saridakis, Anjan A. Sen, Arman Shafieloo, Anowar J. Shajib, Joseph Silk, Alessandra Silvestri, Martin S. Sloth, Tristan L. Smith, Joan Solà, Carsten van de Bruck, Licia Verde, Luca Visinelli, Benjamin D. Wandelt, Deng Wang, Jian–Min Wang, Anil K. Yadav, Weiqiang Yang

The standard  $\Lambda$  Cold Dark Matter cosmological model provides a wonderful fit to current cosmological data, but a few tensions and anomalies became statistically significant with the latest data analyses. While these anomalies could be due to the presence of systematic errors in the experiments, they could also indicate the need for new physics beyond the standard model. In this Letter of Interest we focus on the tension of the Planck data with weak lensing measurements and redshift surveys, about the value of the matter energy density  $\Omega_m$ , and the amplitude or rate of the growth of structure ( $\sigma_8, f\sigma_8$ ). We list a few interesting models for solving this tension, and we discuss the importance of trying to fit with a single model a full array of data and not just one parameter at a time.

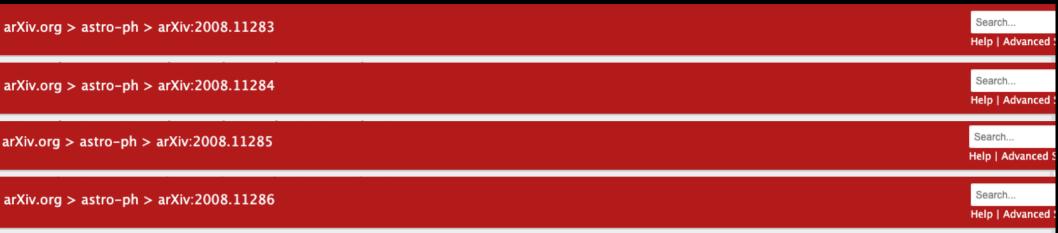
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## Cosmology Intertwined IV: The Age of the Universe and its Curvature

Eleonora Di Valentino, Luis A. Anchordoqui, Ozgur Akarsu, Yacine Ali-Haimoud, Luca Amendola, Nikki Arendse, Marika Asgari, Mario Ballardini, Spyros Basilakos, Elia Battistelli, Micol Benetti, Simon Birrer, François R. Bouchet, <u>Marco Bruni</u>, Erminia Calabrese, David Camarena, Salvatore Capozziello, Angela Chen, Jens Chluba, Anton Chudaykin, Eoin Ó Colgáin, Francis-Yan Cyr-Racine, Paolo de Bernardis, Javier de Cruz Pérez, Jacques Delabrouille, Celia Escamilla-Rivera, Agnès Ferté, Fabio Finelli, Wendy Freedman, Noemi Frusciante, Elena Giusarma, Adrià Gómez-Valent, Will Handley, Ian Harrison, Luke Hart, Alan Heavens, Hendrik Hildebrandt, Daniel Holz, Dragan Huterer, Mikhail M. Ivanov, Shahab Joudaki, Marc Kamionkowski, Tanvi Karwal, Lloyd Knox, Suresh Kumar, Luca Lamagna, Julien Lesgourgues, Matteo Lucca, Valerio Marra, Silvia Masi, Sabino Matarrese, Arindam Mazumdar, Alessandro Melchiorri, Olga Mena, Laura Mersini-Houghton, Vivian Miranda, Cristian Moreno-Pulido, David F. Mota, Jessica Muir, Ankan Mukherjee, Florian Niedermann, Alessio Notari, Rafael C. Nunes, Francesco Pace, Andronikos Paliathanasis, Antonella Palmese, Supriya Pan, Daniela Paoletti, Valeria Pettorino, Francesco Piacentini, Vivian Poulin, Marco Raveri, Adam G. Riess, Vincenzo Salzano, Emmanuel N. Saridakis, Anjan A. Sen, Arman Shafieloo, Anowar J. Shajib, Joseph Silk, Alessandra Silvestri, Martin S. Sloth, Tristan L. Smith, Joan Solà, Carsten van de Bruck, Licia Verde, Luca Visinelli, Benjamin D. Wandelt, Deng Wang, Jian-Min Wang, Anil K. Yadav, Weigiang Yang

A precise measurement of the curvature of the Universe is of primeval importance for cosmology since it could not only confirm the paradigm of primordial inflation but also help in discriminating between different early Universe scenarios. The recent observations, while broadly consistent with a spatially flat standard  $\Lambda$  Cold Dark Matter ( $\Lambda$ CDM) model, are showing tensions that still allow (and, in some cases, even suggest) a few percent deviations from a flat universe. In particular, the Planck Cosmic Microwave Background power spectra, assuming the nominal likelihood, prefer a closed universe at more than 99\% confidence level. While new physics could be in action, this anomaly may be the result of an unresolved systematic error or just a statistical fluctuation. However, since a positive curvature allows a larger age of the Universe, an accurate determination of the age of the oldest objects provides a smoking gun in confirming or falsifying the current flat  $\Lambda$ CDM model.

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### Cosmology Intertwined: A Review of the Particle Physics, Astrophysics, and Cosmology Associated with the Cosmological Tensions and Anomalies

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Jaffe,<sup>7</sup> In Sung Jang,<sup>52</sup> Karsten Jedamzik,<sup>108</sup> Raul Jimenez,<sup>109,110</sup> Melissa Joseph,<sup>11</sup> Shahab Joudaki,<sup>111,112</sup> Marc Kamionkowski,<sup>37</sup> Tanvi Karwal,<sup>113</sup> Lavrentios Kazantzidis,<sup>10</sup> Ryan E. Keeley,<sup>114</sup> Michael Klasen,<sup>3</sup> Eiichiro Komatsu,<sup>115,116</sup> Léon V.E. Koopmans,<sup>117</sup> Suresh Kumar,<sup>118</sup> Luca Lamagna,<sup>28,29</sup> Ruth Lazkoz,<sup>119</sup> Chung-Chi Lee,<sup>120</sup> Julien Lesgourgues,<sup>121</sup> Jackson Levi Said,<sup>122,123</sup> Tiffany R. Lewis,<sup>124</sup> Benjamin L'Huillier,<sup>125</sup> Matteo Lucca,<sup>126</sup> Roy Maartens,<sup>23,127,128</sup> Lucas M. Macri,<sup>129</sup> Danny Marfatia,<sup>130</sup> Valerio Marra,<sup>131, 132, 133</sup> Carlos J. A. P. Martins,<sup>134, 135</sup> Silvia Masi,<sup>28, 29</sup> Sabino Matarrese, 136, 137, 138, 139 Arindam Mazumdar, 140 Alessandro Melchiorri, 28, 29 Olga Mena, 141 Laura Mersini-Houghton,<sup>142</sup> James Mertens,<sup>143</sup> Dinko Milaković,<sup>133,132,144</sup> Yuto Minami,<sup>145</sup> Vivian Miranda,<sup>146</sup> Cristian Moreno-Pulido,<sup>147</sup> Michele Moresco,<sup>46,47</sup> David F. Mota,<sup>148</sup> Emil Mottola,<sup>63</sup> Simone Mozzon,<sup>149</sup> Jessica Muir,<sup>150</sup> Ankan Mukherjee,<sup>151</sup> Suvodip Mukherjee,<sup>150</sup> Pavel Naselsky,<sup>152</sup> Pran Nath,<sup>153</sup> Savvas Nesseris,<sup>99</sup> Florian Niedermann,<sup>154</sup> Alessio Notari,<sup>155</sup> Rafael C. Nunes,<sup>156</sup> Eoin Ó Colgáin,<sup>157, 158</sup> Kayla A. Owens,<sup>52</sup> Emre Özülker,<sup>5</sup> Francesco Pace, <sup>159, 160</sup> Andronikos Paliathanasis, <sup>161, 162</sup> Antonella Palmese, <sup>163</sup> Supriya Pan, <sup>164</sup> Daniela Paoletti, <sup>85, 22</sup> Santiago E. Perez Bergliaffa,<sup>165</sup> Leandros Perivolaropoulos,<sup>10</sup> Dominic W. Pesce,<sup>166,167</sup> Valeria Pettorino,<sup>168</sup> Oliver H. E. Philcox,<sup>169,107</sup> Levon Pogosian,<sup>170</sup> Vivian Poulin,<sup>2</sup> Gaspard Poulot,<sup>80</sup> Marco Raveri,<sup>171</sup> Mark J. Reid,<sup>172</sup> Fabrizio Renzi,<sup>173</sup> Adam G. Riess,<sup>37</sup> Vivian I. Sabla,<sup>54</sup> Paolo Salucci,<sup>174,175</sup> Vincenzo Salzano,<sup>176</sup> Emmanuel N. Saridakis,<sup>26,75,177</sup> Bangalore S. Sathyaprakash,<sup>178,179,94</sup> Martin Schmaltz,<sup>11</sup> Nils Schöneberg,<sup>180</sup> Dan Scolnic,<sup>181</sup> Anjan A. Sen,<sup>182,183</sup> Neelima Sehgal,<sup>184</sup> Arman Shafieloo,<sup>185</sup> M.M. Sheikh-Jabbari,<sup>186</sup> Joseph Silk,<sup>187</sup> Alessandra Silvestri,<sup>173</sup> Foteini Skara,<sup>10</sup> Martin S. Sloth,<sup>188</sup> Marcelle Soares-Santos,<sup>58</sup> Joan Solà Peracaula,<sup>147</sup> Yu-Yang Songsheng,<sup>81</sup> Jorge F. Soriano,<sup>13,14</sup> Denitsa Staicova,<sup>189</sup> Glenn D. Starkman,<sup>6,7</sup> István Szapudi,<sup>190</sup> Elsa M. Teixeira,<sup>80</sup> Brooks Thomas,<sup>191</sup> Tommaso Treu,<sup>60</sup> Emery Trott,<sup>58</sup> Carsten van de Bruck,<sup>80</sup> J. Alberto Vazquez,<sup>192</sup> Licia Verde,<sup>193,194</sup> Luca Visinelli,<sup>195</sup> Deng Wang,<sup>196</sup> Jian-Min Wang,<sup>81</sup> Shao-Jiang Wang,<sup>197</sup> Richard Watkins,<sup>198</sup> Scott Watson,<sup>199</sup> John K. Webb,<sup>120</sup> Neal Weiner,<sup>200</sup> Amanda Weltman,<sup>201</sup> Samuel J. Witte,<sup>202</sup> Radosław Wojtak,<sup>4</sup> Anil Kumar Yadav,<sup>203</sup> Weiqiang Yang,<sup>204</sup> Gong-Bo Zhao,<sup>205, 206</sup> and Miguel Zumalacárregui<sup>207</sup>

Finally, you can also find a section where we discuss in a unified manner many less discussed less-standard existing signals in cosmological and astrophysical data that appear to be in some tension (20 or larger) with the standard ACDM model as defined by the Planck 2018 parameter values.

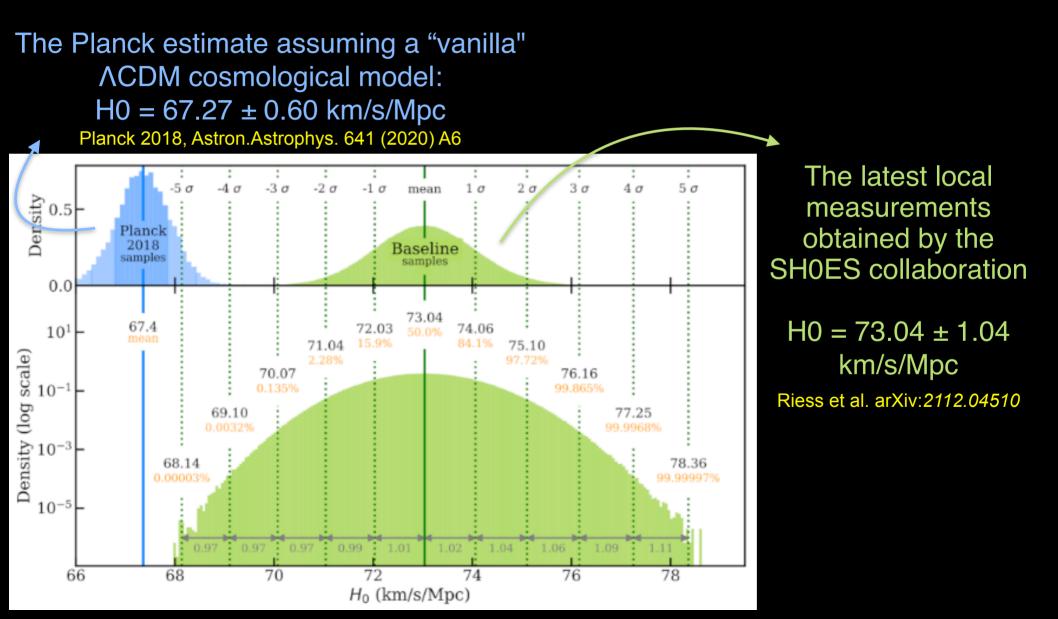
In many cases the signals are controversial and there is currently debate in the literature on the possible systematics origin of some of these signals. I encourage you to have a look at the paper if you are interested in learning more.

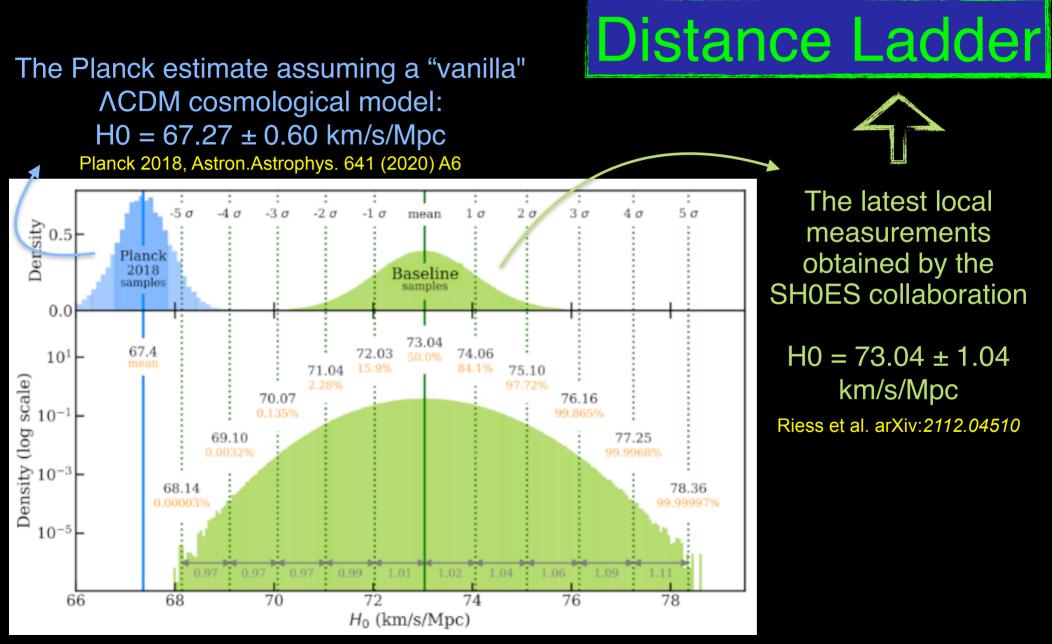
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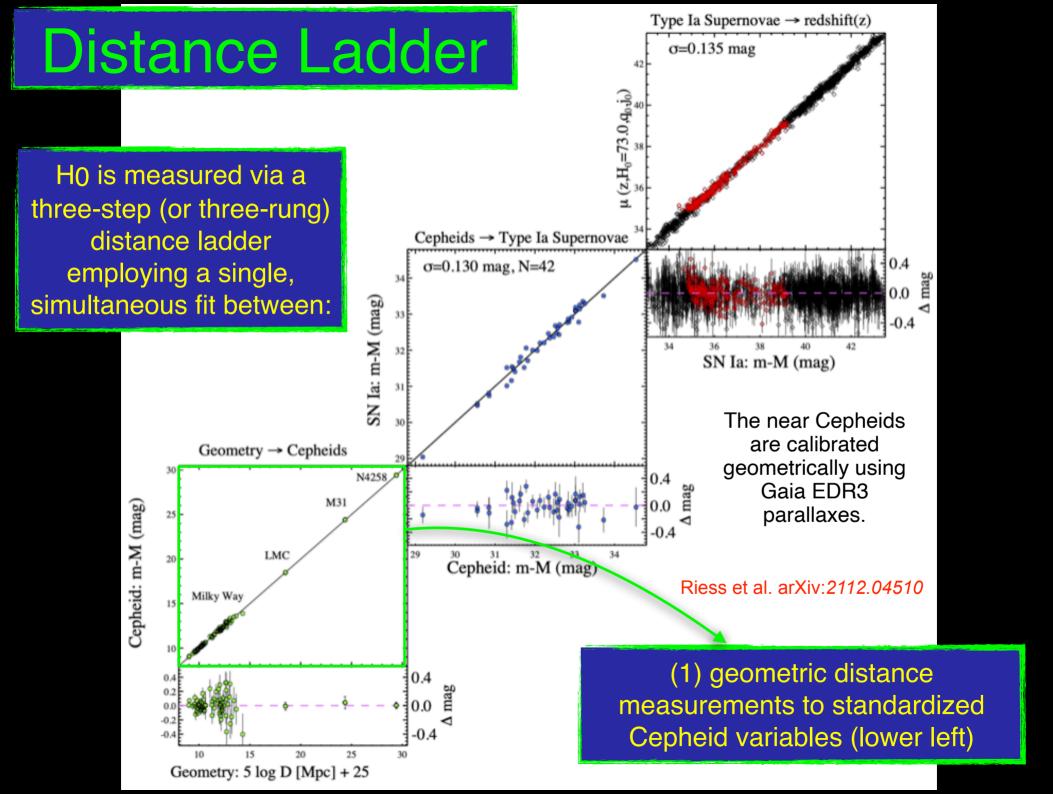
Corresponding author (e.divalentino@sheffield.ac.uk)

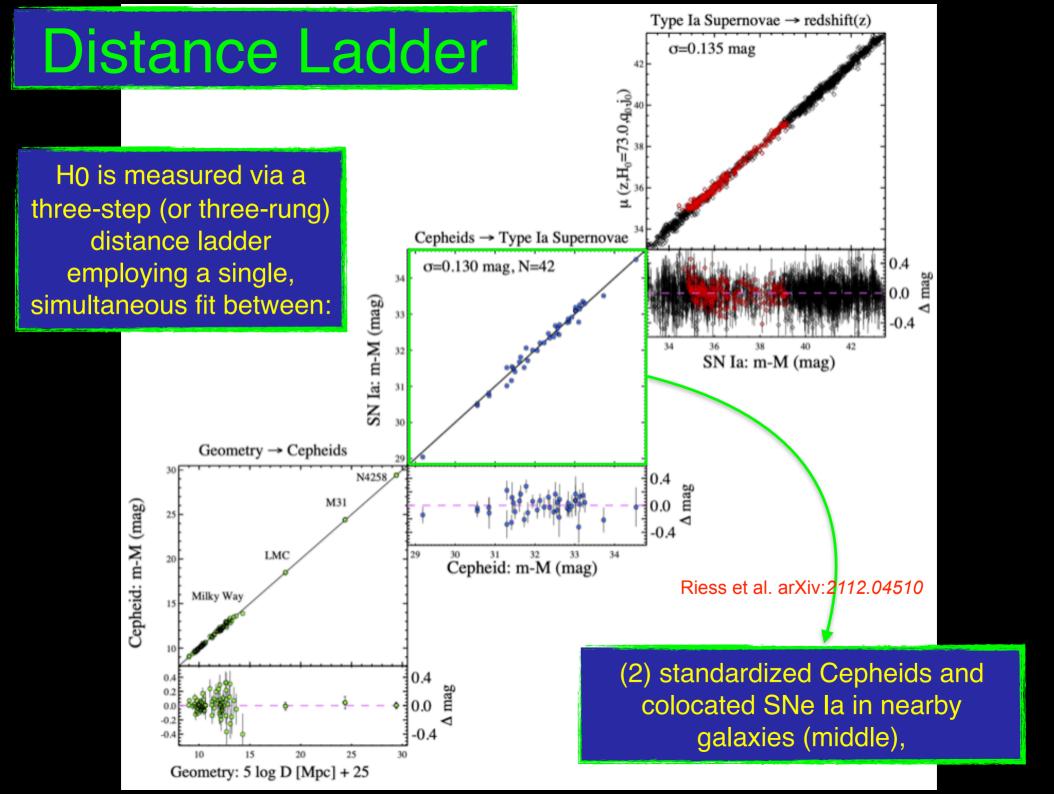
# The H0 tension exceeds $5\sigma!!$

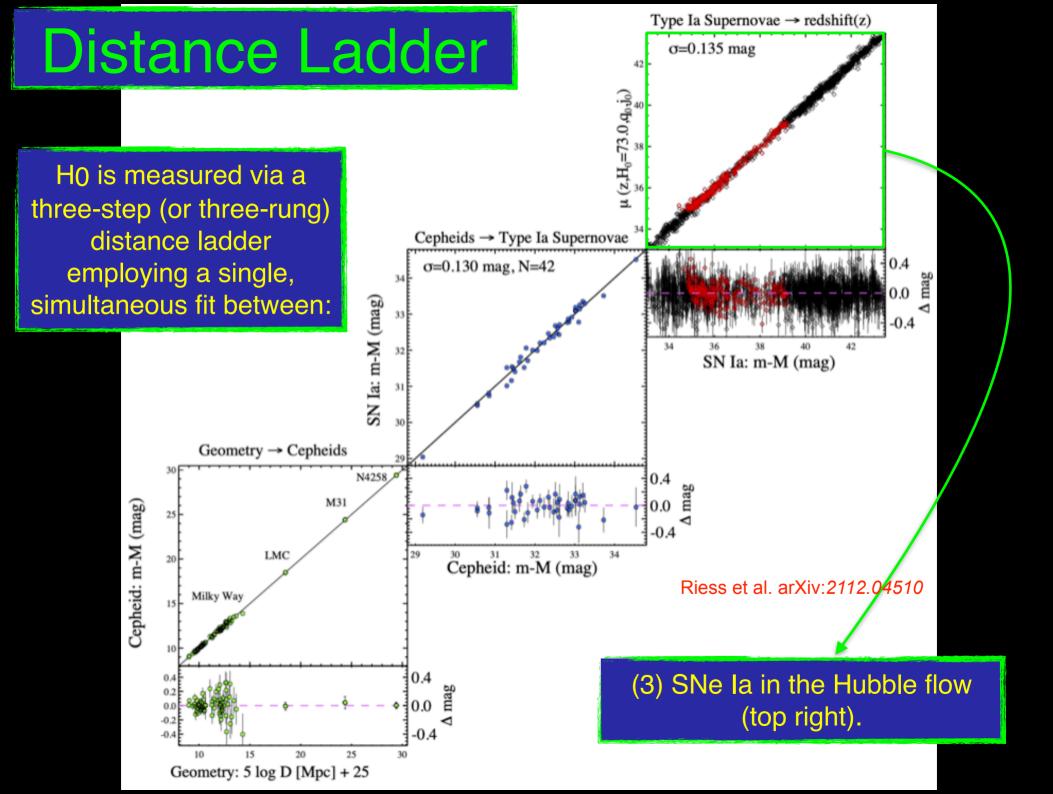
The H0 tension is the most statistically significant, long-lasting and widely persisting disagreement we have currently in cosmology.

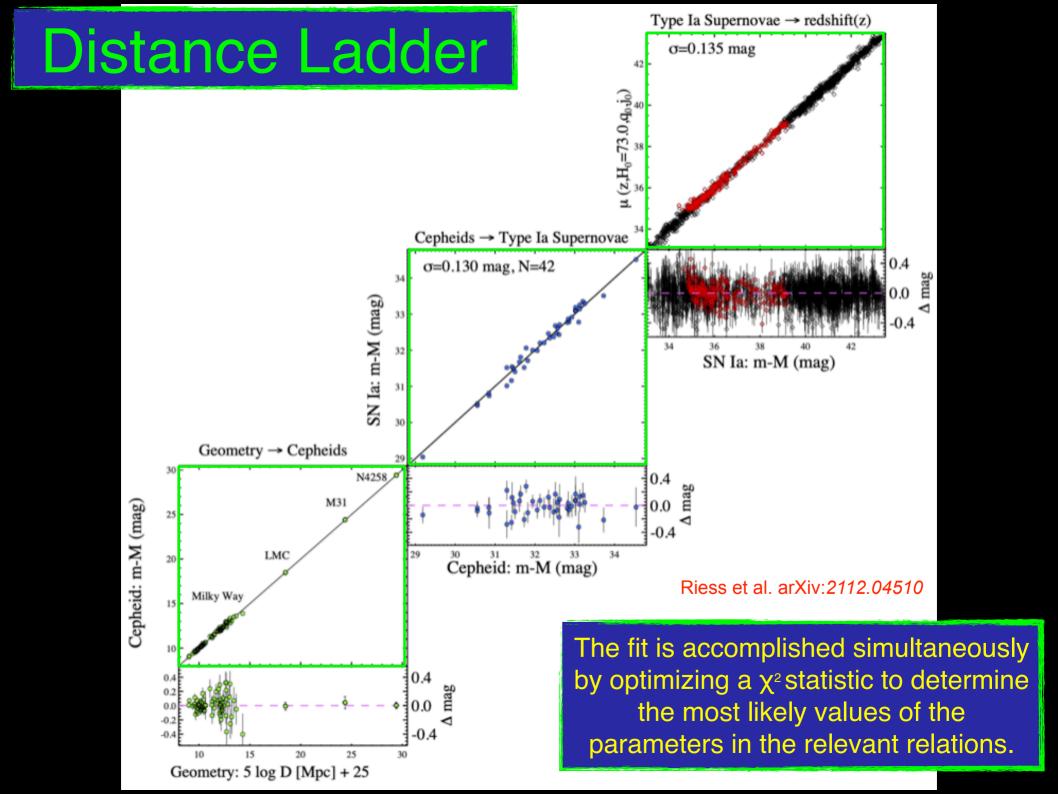






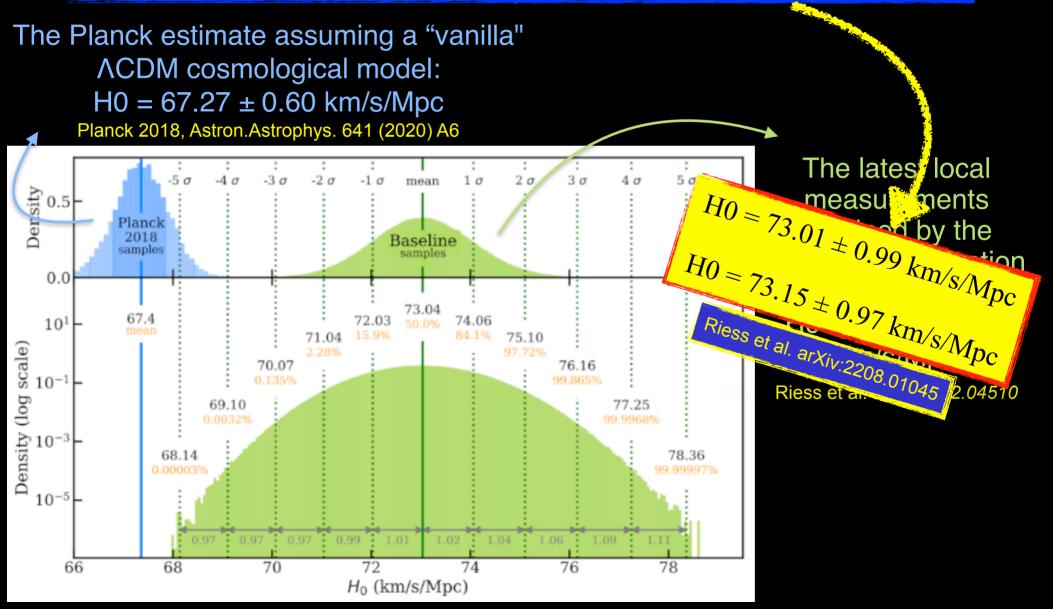


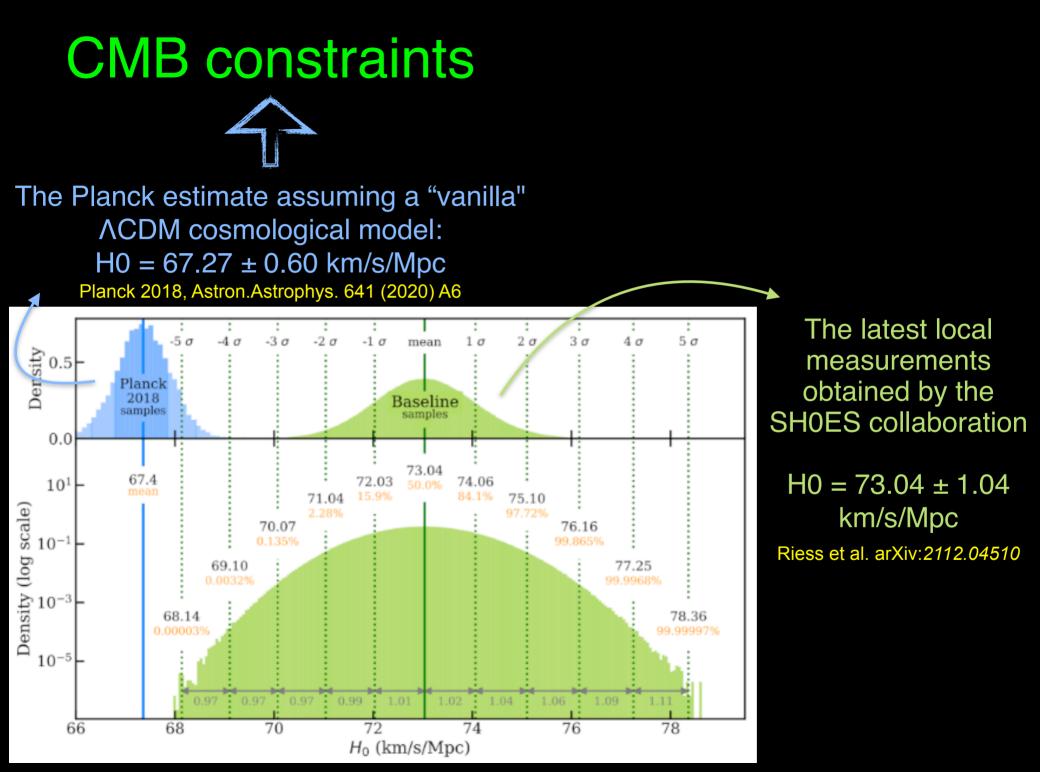




# The H0 tension exceeds $5\sigma!!$

# The H0 tension at 5.3 or



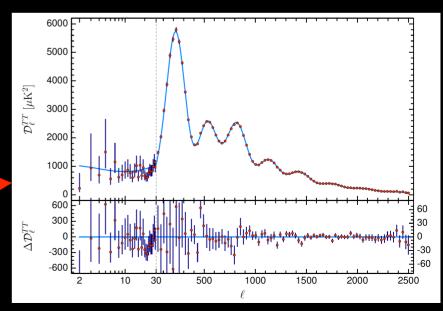


## **CMB** constraints

$$\left\langle \frac{\Delta T}{T} \left( \vec{\gamma}_1 \right) \frac{\Delta T}{T} \left( \vec{\gamma}_2 \right) \right\rangle = \frac{1}{2\pi} \sum_{\ell} (2\ell + 1) C_{\ell} P_{\ell} \left( \vec{\gamma}_1 \cdot \vec{\gamma}_2 \right)$$

From the map of the CMB anisotropies we can extract the temperature angular power spectrum.





# Cosmological parameters: $(\Omega_b h^2, \Omega_m h^2, H0, n_s, \tau, As)$



We choose a set of cosmological parameters that describes our theoretical model and compute the angular power spectra.

Because of the correlations present between the parameters, variation of different quantities can produce similar effects on the CMB.

## 100 (a) Curvature (b) Dark Energy 80 $\Delta_T (\mu K)$ 4020 $\Omega_{tot}$ $\Omega$ (d) Matter (c) Baryons 100 80 $\Delta_T (\mu \mathbf{K})$ 60 40 20 $\Omega_{\rm b}h^2$ $\Omega_m h^2$ 0.04 10 100 1000 10 100 1000

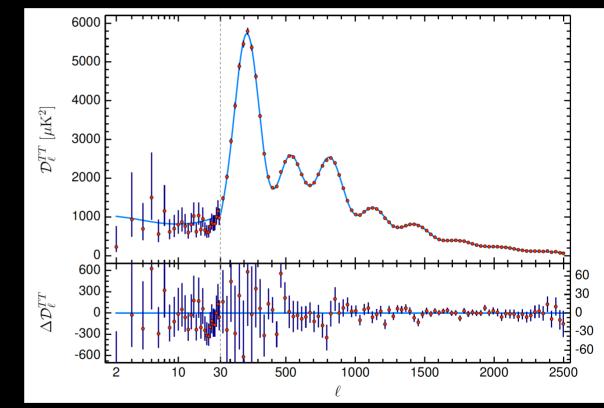
Wayne Hu's tutorial

## Cosmological parameters: $(\Omega_b h^2, \Omega_m h^2, H0, n_s, \tau, As)$

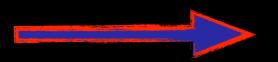


**Theoretical model** 

We compare the angular power spectra we computed with the data and, using a bayesian analysis, we get a combination of cosmological parameter values in agreement with these.



Planck 2018, Astron.Astrophys. 641 (2020) A6



## Parameter constraints

# **CMB** constraints

	TT+lowE	TE+lowE	EE+lowE	TT,TE,EE+lowE	TT,TE,EE+lowE+lensing	TT,TE,EE+lowE+lensing+BAO
Parameter	68% limits	68% limits	68% limits	68% limits	68% limits	68% limits
$\Omega_{\rm b}h^2$	$0.02212 \pm 0.00022$	$0.02249 \pm 0.00025$	$0.0240 \pm 0.0012$	$0.02236 \pm 0.00015$	$0.02237 \pm 0.00015$	$0.02242 \pm 0.00014$
$\Omega_{\rm c} h^2$	$0.1206 \pm 0.0021$	$0.1177 \pm 0.0020$	$0.1158 \pm 0.0046$	$0.1202 \pm 0.0014$	$0.1200 \pm 0.0012$	$0.11933 \pm 0.00091$
100θ <sub>MC</sub>	$1.04077 \pm 0.00047$	$1.04139 \pm 0.00049$	$1.03999 \pm 0.00089$	$1.04090 \pm 0.00031$	$1.04092 \pm 0.00031$	$1.04101 \pm 0.00029$
τ	$0.0522 \pm 0.0080$	$0.0496 \pm 0.0085$	$0.0527 \pm 0.0090$	$0.0544^{+0.0070}_{-0.0081}$	$0.0544 \pm 0.0073$	$0.0561 \pm 0.0071$
$\ln(10^{10}A_{\rm s})$	$3.040 \pm 0.016$	$3.018^{+0.020}_{-0.018}$	$3.052\pm0.022$	$3.045 \pm 0.016$	$3.044 \pm 0.014$	$3.047 \pm 0.014$
<i>n</i> <sub>s</sub>	$0.9626 \pm 0.0057$	$0.967 \pm 0.011$	$0.980 \pm 0.015$	$0.9649 \pm 0.0044$	$0.9649 \pm 0.0042$	$0.9665 \pm 0.0038$
$H_0 [\mathrm{kms^{-1}Mpc^{-1}}]$	$66.88 \pm 0.92$	$68.44 \pm 0.91$	69.9 ± 2.7	$67.27 \pm 0.60$	67.36 ± 0.54	$67.66 \pm 0.42$
$\Omega_{\Lambda}$	$0.679 \pm 0.013$	$0.699 \pm 0.012$	$0.711^{+0.033}_{-0.026}$	$0.6834 \pm 0.0084$	$0.6847 \pm 0.0073$	$0.6889 \pm 0.0056$
$\Omega_m \ldots \ldots \ldots \ldots \ldots$	$0.321 \pm 0.013$	$0.301 \pm 0.012$	$0.289^{+0.026}_{-0.033}$	$0.3166 \pm 0.0084$	$0.3153 \pm 0.0073$	$0.3111 \pm 0.0056$
$\Omega_{\rm m} h^2$	$0.1434 \pm 0.0020$	$0.1408 \pm 0.0019$	$0.1404^{+0.0034}_{-0.0039}$	$0.1432 \pm 0.0013$	$0.1430 \pm 0.0011$	$0.14240 \pm 0.00087$
$\Omega_{\rm m} h^3$	$0.09589 \pm 0.00046$	$0.09635 \pm 0.00051$	$0.0981^{+0.0016}_{-0.0018}$	$0.09633 \pm 0.00029$	$0.09633 \pm 0.00030$	$0.09635 \pm 0.00030$
<i>σ</i> <sub>8</sub>	$0.8118 \pm 0.0089$	$0.793 \pm 0.011$	$0.796 \pm 0.018$	$0.8120 \pm 0.0073$	$0.8111 \pm 0.0060$	$0.8102 \pm 0.0060$
$S_8\equiv\sigma_8(\Omega_{\rm m}/0.3)^{0.5}$ .	$0.840 \pm 0.024$	$0.794 \pm 0.024$	$0.781^{+0.052}_{-0.060}$	$0.834 \pm 0.016$	$0.832 \pm 0.013$	$0.825 \pm 0.011$

## Planck 2018, Astron.Astrophys. 641 (2020) A6

2018 Planck results are a wonderful confirmation of the flat standard ΛCDM cosmological model, but are **model dependent**!

- The cosmological constraints are obtained assuming a cosmological model.
- The results are affected by the degeneracy between the parameters that induce similar effects on the observables.

# Are there other H0 estimates?

# The H0 tension

CMB Polarization Measurements with SPTpol

On the same side of Planck, i.e. preferring smaller values of H0 we have:

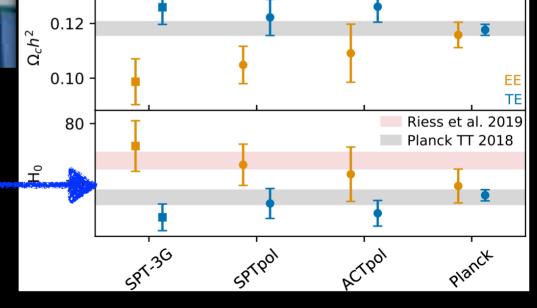
Ground based CMB telescope

Nicholas Harrington UC Berkeley

 $\frac{\text{SPT-3G}}{\text{H0} = 68.8 \pm 1.5 \text{ km/s/Mpc} \text{ in } \Lambda \text{CDM}}$ 

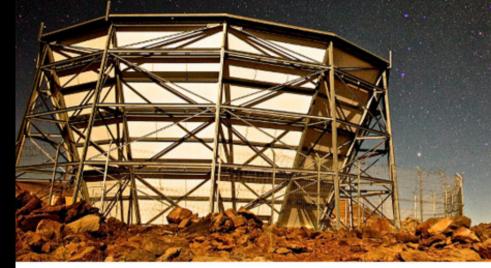
LCD/M - dependent

SPT-3G, Dutcher et al., Phys.Rev.D 104 (2021) 2, 022003

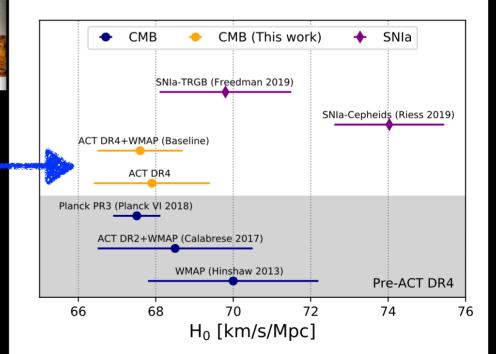


# The H0 tension

On the same side of Planck, i.e. preferring smaller values of H<sub>0</sub> we have:



Ground based CMB telescope



ACT collaboration, Aiola et al., JCAP 12 (2020) 047

ACT-DR4: H0 =  $67.9 \pm 1.5$  km/s/Mpc in  $\Lambda$ CDM

 $\frac{\text{ACT-DR4} + \text{WMAP}}{\text{H0} = 67.6 \pm 1.1 \text{ km/s/Mpc} \text{ in } \text{ACDM}}$ 

LCDM - dependent

# The H0 tension

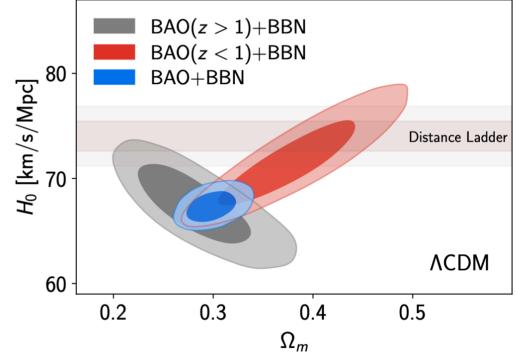
On the same side of Planck, i.e. preferring smaller values of H<sub>0</sub> we have:

## BAO+Pantheon+BBN+ $\theta_{MC, Planck}$ : H0 = 67.9 ± 0.8 km/s/Mpc

Planck 2018, Aghanim et al., Astron.Astrophys. 641 (2020) A6

## BAO+BBN from BOSS and eBOSS: $H_0 = 67.35 \pm 0.97 \text{ km/s/Mpc}$

eBOSS, Alam et al., Phys.Rev.D 103 (2021) 8, 083533



eBOSS, Alam et al., Phys.Rev.D 103 (2021) 8, 083533

LCD/M - dependent

#### CMR with Planel

Balkenhol et al. (2021), Planck 2018+SPT+ACT : 67.49 ± 0.5 Pogosian et al. (2020), eBOSS+Planck mH2: 69.6 ± 1.8 Aghanim et al. (2020), Planck 2018: 67.27 ± 0.60 ghanim et al. (2020), Planck 2018+CMB lensing: 67.36 ± 0.54 Ade et al. (2016), Planck 2015, H0 = 67.27 ± 0.60

Dutcher et al. (2021), SPT: 68.8 ± 1.5 Aiola et al. (2020). ACT: 67.9 + 1 Aiola et al. (2020), WMAP9+ACT: 67.6 + 1 Zhang Huang (2019) WMAP9+BAO: 68 36+9-53 Henning et al. (2018), SPT: 71.3 ± 2. Hinshaw et al. (2013), WMAP9: 70.0 ± 2.2

#### No CMB, with BBN

Zhang et al. (2021), BOSS correlation function+BAO+BBN: 68.19±0.99 Chen et al. (2021). P+BAO+BBN: 69 23+0 7 Philcox et al. (2021), P+Bispectrum+BAO+BBN: 68.31+0.8 D' Amico et al. (2020). BOSS DR12+BBN: 68.5 + 2 Colas et al. (2020), BOSS DR12+BBN: 68.7 ± 1.5 Ivanov et al. (2020), BOSS+BBN: 67.9 ± 1. Alam et al. (2020). BOSS+cBOSS+BBN: 67.35 + 0.9

#### CMB lensin Baxter et al. (2020): 73.5 ± 5

Philcox et al. (2020), P1(k)+CMB lensing: 70.6-3 LSS teg standard rule Farren et al. (2021): 69 5+3

#### SNIa-Cenheid

Riess et al. (2022), R22: 73.04 ± 1.04 Camarena Marra (2021): 74 30 + 1 45 Riess et al. (2020), R20: 73.2 ± 1. Breuval et al. (2020): 72.8 ± 2.7 Riess et al. (2019), R19: 74.03 ± 1.4 Camarena, Marra (2019): 75.4 + 1

### SNIa-TRGB

Dhawan et al. (2022): 76.94 ± 6.4 Iones et al. (2022): 72.4 + 3.3 Anand, Tully, Rizzi, Riess, Yuan (2021): 71.5 ± 1. Freedman (2021): 69.8 ± 1. Kim, Kang, Lee, Jang (2021): 69.5 ± 4.2 Soltis Casertano Riess (2020): 72.1 + 2.0 Freedman et al. (2020): 69.6 ± 1.9 Reid Pesce Riess (2019) SH0ES: 71.1 + 1.99 Yuan et al. (2019): 72.4 + 2.0

> SNIa-Miras Huang et al. (2019): 73.3 ± 4.0

Blakeslee et al. (2021) IR-SBF w/ HST: 73.3 ± 2.5 Khetan et al. (2020) w/ LMC DEB: 71.1 + 4.1 Cantiello et al. (2018): 71.9 + 7.

> de Jacquer et al. (2022): 75.4+3de Jaeger et al. (2020): 75.8-3 c

> > Masers Pesce et al. (2020): 73.9 + 3.0

Tully Fisher Kourkchi et al. (2020): 76.0 ± 2.6

Schombert, McGaugh, Lelli (2020): 75.1 ± 2.8 HII galaxy Fernandez Arenas et al. (2018): 71.0 ± 3 Wang, Meng (2017): 76.12+3-4

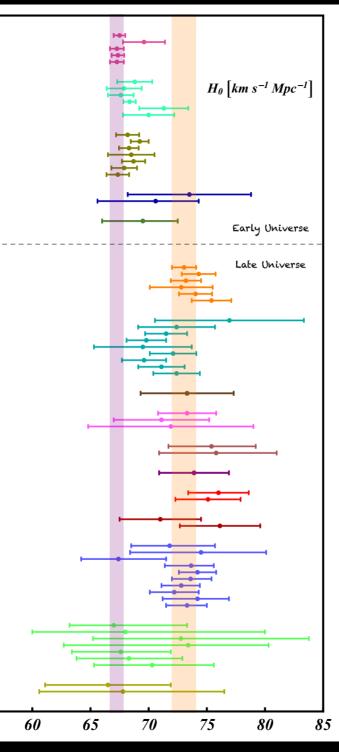
## Lensing related,mass model depende Denzel et al. (2021): 71.8<sup>±</sup>

Birrer et al. (2020). TDCOSMO: 74.5+ Birrer et al. (2020), TDCOSMO+SLACS: 67.4 Yang, Birrer, Hu (2020): 73.65+1-3 Millon et al. (2020). TDCOSMO: 74.2  $\pm$  1. Oi et al (2020): 73.6+ Liao et al. (2020): 72.8+ Liao et al. (2019): 72.2 ± Shajib et al. (2019), STRIDES: 74.2 Wong et al. (2019), H0LiCOW 2019: 73.3+

### CW rol

GW related Mukherjee et al. (2022), GW170817+GWTC-3: 67<sup>+9</sup>, Abbott et al. (2021), GWTC-3: 68<sup>+1</sup> Palmese et al. (2021), GW170817: 72.77<sup>+1</sup> -2.55 Gavathri et al. (2020), GW190521+GW170817: 73.4+6 Mukherjee et al. (2020), GW170817+ZTF: 67.6 Mukherice et al. (2019), GW170817+VLBI: 68.3+2 Hotokezaka et al. (2019): 70.3+5

### Cosmic chronometer Moresco et al. (2022), flat ACDM with systematics: 66.5 ± 5.4 Moresco et al. (2022), open wCDM with systematics: 67.8±8-3



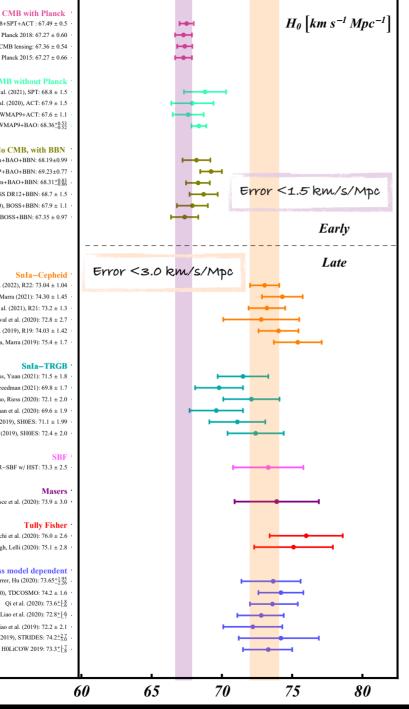
Hubble constant measurements made by different astronomical missions and groups over the years.

The orange vertical band corresponds to the H0 value from **SHOES Team** and the light pink vertical band corresponds to the H0 value as reported by Planck 2018 team within a ACDM scenario.

Abdalla et al., JHEAp 34 (2022) 49-211

## High precision measurements of HO

The high precision and consistency of the data at both ends present strong challenges to the possible solution space and demands a hypothesis with enough rigor to explain multiple observations whether these invoke new physics, unexpected largescale structures or multiple, unrelated errors.



Balkenhol et al. (2021), Planck 2018+SPT+ACT : 67.49 ± 0.5

Aghanim et al. (2020), Planck 2018: 67.27 ± 0.60 Aghanim et al. (2020). Planck 2018+CMB lensing: 67.36 + 0.54 Ade et al. (2016), Planck 2015: 67.27 ± 0.66

Dutcher et al. (2021). SPT: 68.8 + 1.5 Aiola et al. (2020), ACT: 67.9 ± 1.5 Aiola et al. (2020). WMAP9+ACT: 67.6 + 1.1 Zhang, Huang (2019), WMAP9+BAO: 68.36+0.53

### No CMB, with BBN

Zhang et al. (2022). BOSS correlation function+BAO+BBN: 68.19+0.99 Chen et al. (2022), P+BAO+BBN: 69.23±0.77 Philcox, Ivanov (2022), P+Bispectrum+BAO+BBN: 68.31+0.83 Colas et al. (2020), BOSS DR12+BBN: 68.7 ± 1.5 Ivanov et al. (2020), BOSS+BBN: 67.9 + 1.1 Alam et al. (2020), BOSS+eBOSS+BBN: 67.35 ± 0.97

#### SnIa-Cepheid

- Riess et al. (2022), R22: 73.04 ± 1.04 Camarena, Marra (2021): 74.30 ± 1.45 Riess et al. (2021), R21: 73.2 ± 1.3 Breuval et al. (2020): 72.8 ± 2.7
- Riess et al. (2019), R19: 74.03 ± 1.42
- Camarena, Marra (2019): 75.4 ± 1.7

### SnIa-TRGB

- Anand, Tully, Rizzi, Riess, Yuan (2021): 71.5 ± 1.8
  - Freedman (2021): 69.8 + 1.7
  - Soltis, Casertano, Riess (2020): 72.1 + 2.0
  - Freedman et al. (2020): 69.6 + 1.9
  - Reid, Pesce, Riess (2019), SH0ES: 71.1 ± 1.99
  - Yuan et al. (2019). SH0ES: 72.4 + 2.0

Blakeslee et al. (2021) IR-SBF w/ HST: 73.3 + 2.5

#### Masers Pesce et al. (2020): 73.9 + 3.0

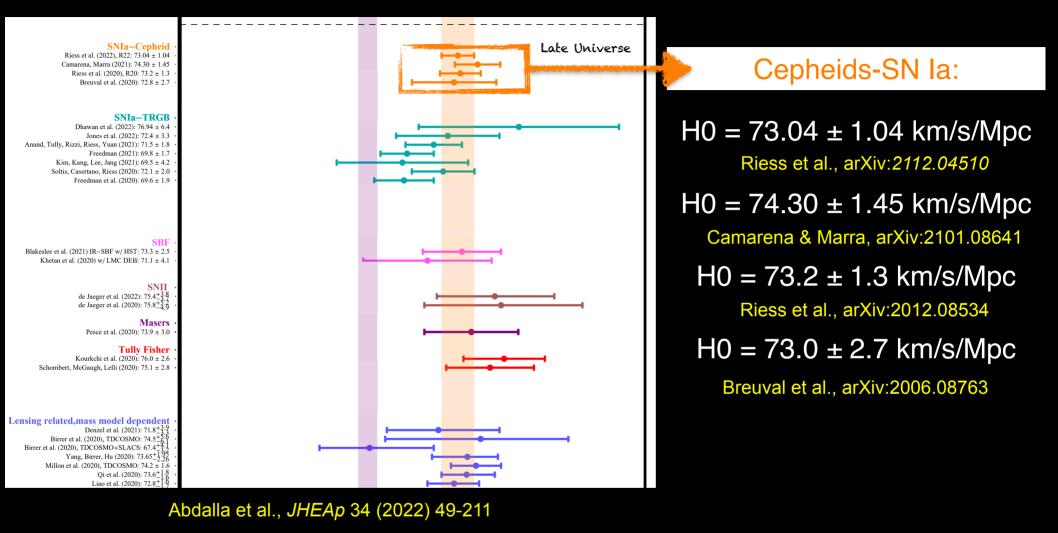
### Tully Fisher

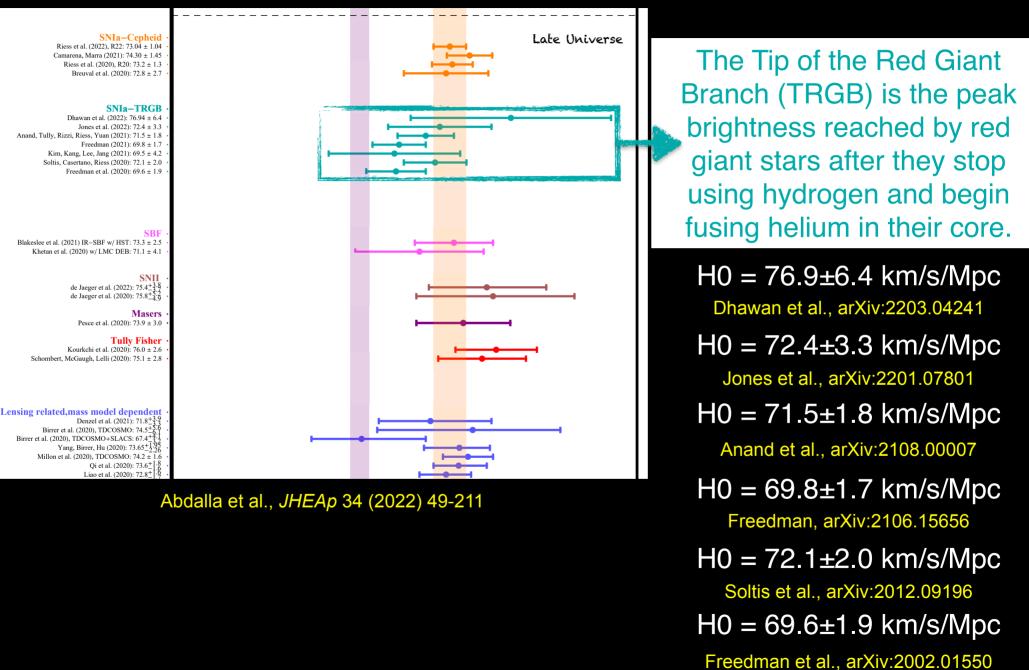
Kourkchi et al. (2020): 76.0 ± 2.6 Schombert, McGaugh, Lelli (2020): 75.1 ± 2.8

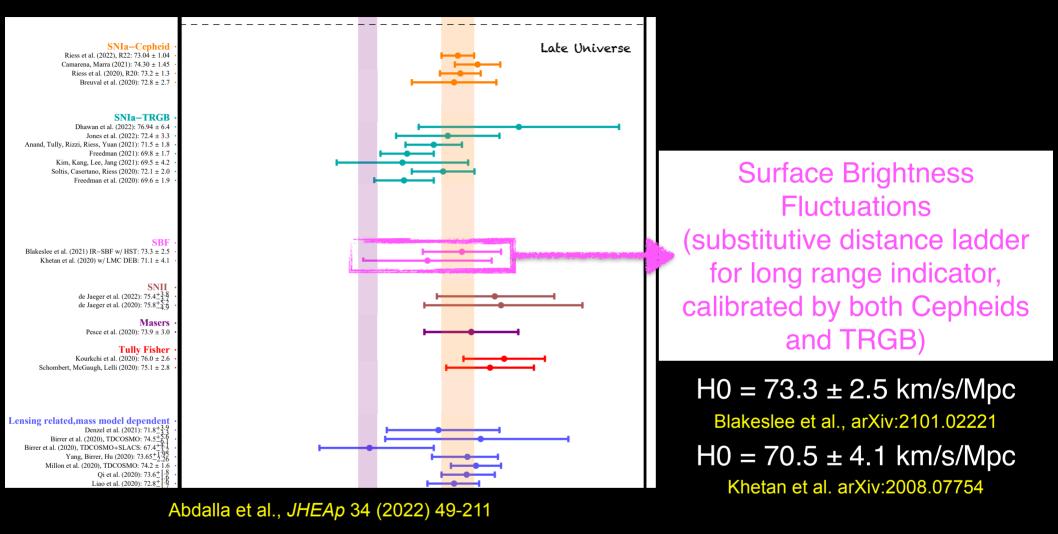
### TD lensing related, mass model dependent

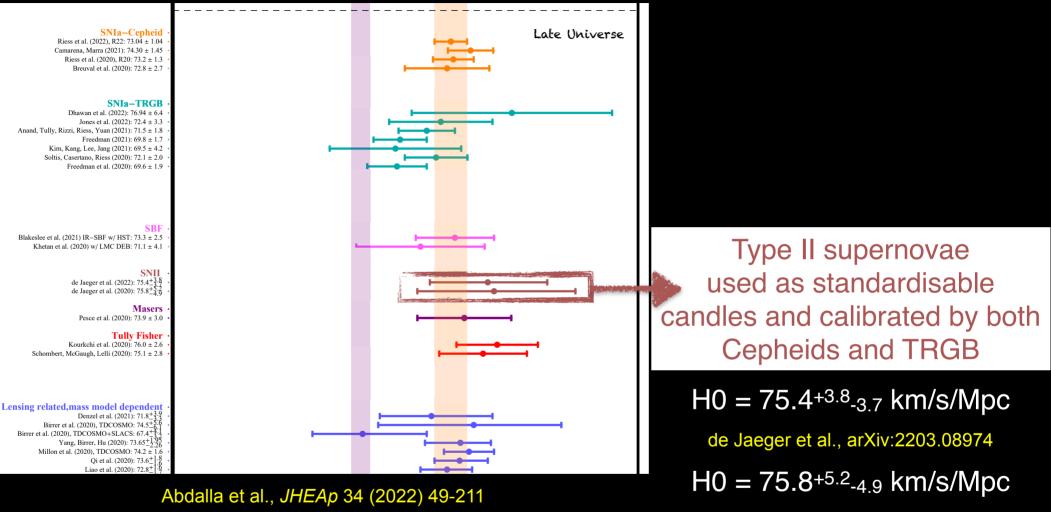
Yang, Birrer, Hu (2020): 73.65+1.95 Millon et al. (2020), TDCOSMO: 74.2 ± 1.6 Qi et al. (2020): 73.6+1-8 Liao et al. (2020): 72.8+1-9 Liao et al. (2019): 72.2 ± 2.1 Shajib et al. (2019), STRIDES: 74.2+27 Wong et al. (2019), H0LiCOW 2019: 73.3+1-7

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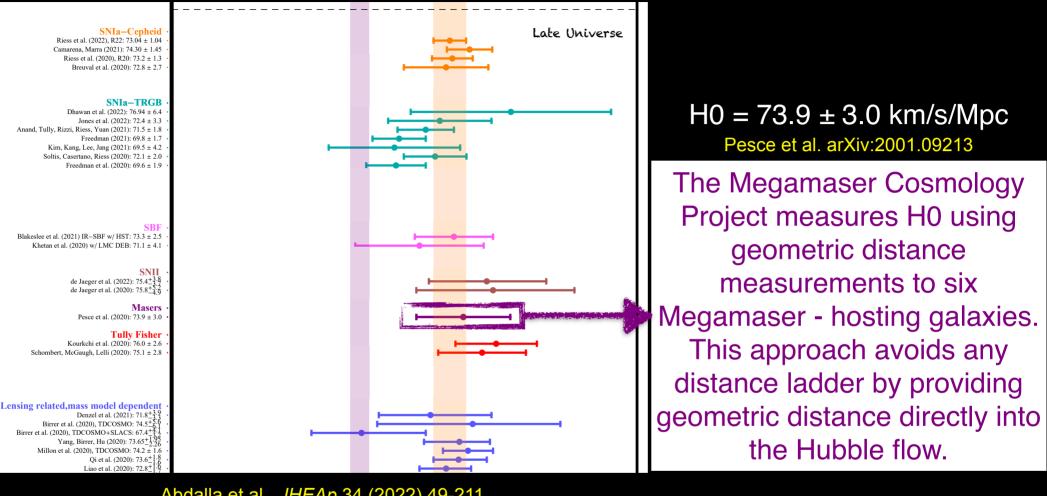




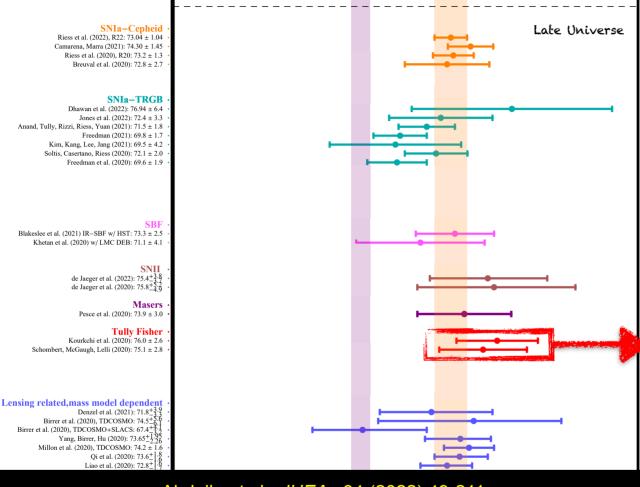




de Jaeger et al., arXiv:2006.03412



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H0 = 76.00 ± 2.55 km/s/Mpc Kourkchi et al. arXiv:2004.14499

H0 = 75.10 ± 2.75 km/s/Mpc Schombert et al. arXiv:2006.08615

Tully-Fisher Relation (based on the correlation between the rotation rate of spiral galaxies and their absolute luminosity, and using as calibrators Cepheids and TRGB)

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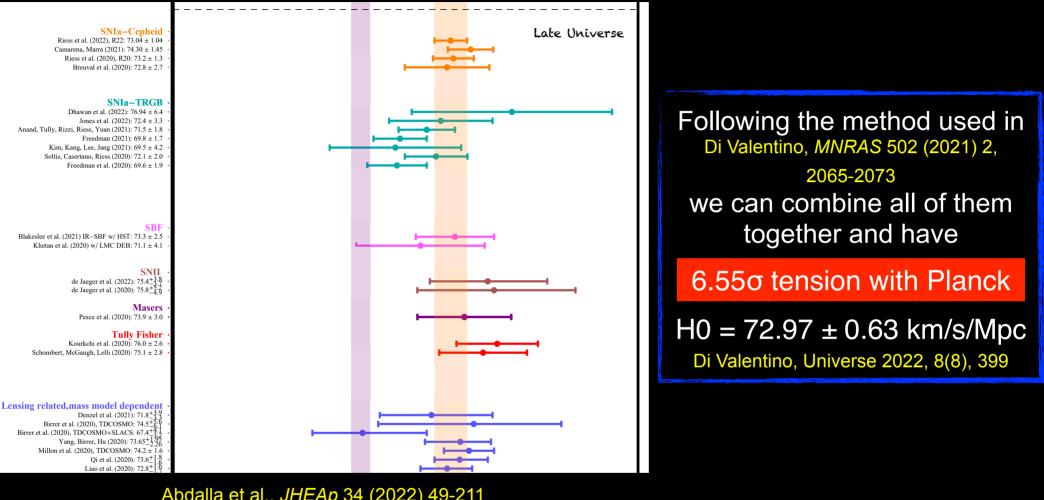
SNIa-Cepheid Late Universe Riess et al. (2022), R22: 73.04 ± 1.04 Camarena, Marra (2021): 74.30 ± 1.45 Riess et al. (2020) R20: 73.2 + 1.3 Breuval et al. (2020): 72.8 ± 2.7 SNIa-TRGB Dhawan et al. (2022): 76.94 + 6.4 Jones et al. (2022): 72.4 ± 3.3 Anand, Tully, Rizzi, Riess, Yuan (2021); 71.5 ± 1.8 Freedman (2021): 69.8 ± 1.7 Kim, Kang, Lee, Jang (2021): 69.5 ± 4.2 Soltis, Casertano, Riess (2020): 72.1 ± 2.0 Freedman et al. (2020): 69.6 ± 1.9 Blakeslee et al. (2021) IR-SBF w/ HST: 73.3 ± 2.5 Khetan et al. (2020) w/ LMC DEB: 71.1 ± 4.1 SNII de Jaeger et al. (2022): 75.4+3.5 de Jaeger et al. (2020): 75.8+5.2 Masers Pesce et al. (2020): 73.9 ± 3.0 Tully Fisher Kourkchi et al. (2020): 76.0 ± 2.6 Schombert, McGaugh, Lelli (2020): 75.1 ± 2.8 Lensing related, mass model dependen Denzel et al. (2021): 71.8+3.9 Birrer et al. (2020), TDCOSMO: 74.5+2 Birrer et al. (2020), TDCOSMO+SLACS: 67.4<sup>42</sup>/<sub>-3</sub> Yang, Birrer, Hu (2020): 73.65-1-9-2 Millon et al. (2020). TDCOSMO: 74.2 + 1.0 Qi et al. (2020): 73.6+ Liao et al. (2020): 72.8+

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Model Dependent

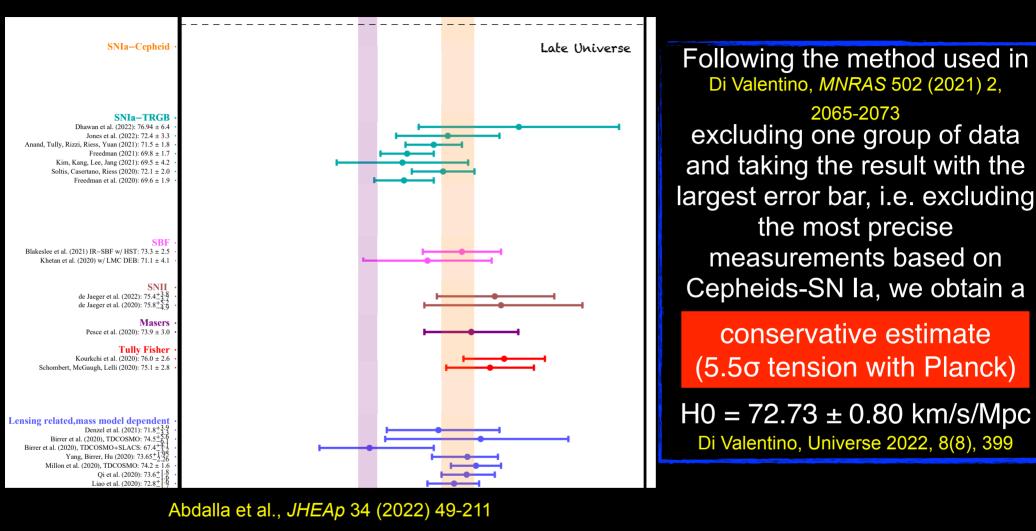
H0 = 72.8 + 1.6 - 1.7 km/s/MpcLiao et al. arXiv:2002.10605 H0 = 73.6 + 1.8 - 1.6 km/s/MpcQi et al. arXiv:2011.00713 H0 = 73.65 + 1.95 - 2.26 km/s/MpcYang et al. arXiv:2003.03277 **TDCOSMO** H0 = 74.5 + 5.6 - 6.1 km/s/Mpc**TDCOSMO+SLACS** H0 = 67.4 + 4.1 - 3.2 km/s/MpcBirrer et al. arXiv:2007.02941 H0 = 71.8 + 3.9 - 3.3 km/s/MpcDenzel et al. arXiv:2007.14398

 Strong Lensing measurements of the time delays of multiple images of quasar systems caused by the strong gravitational lensing from a foreground galaxy.
 Uncertainties coming from the lens mass profile.

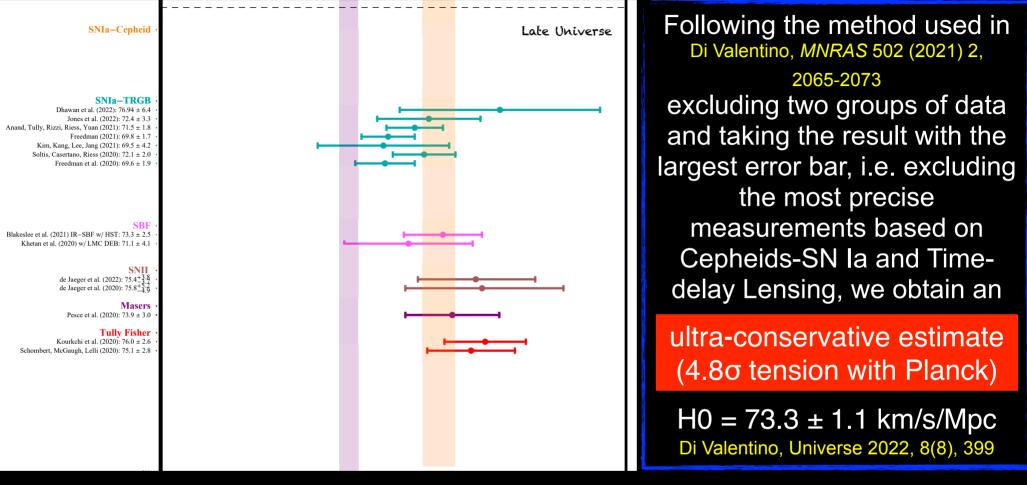


Abdalla et al., JHEAp 34 (2022) 49-211

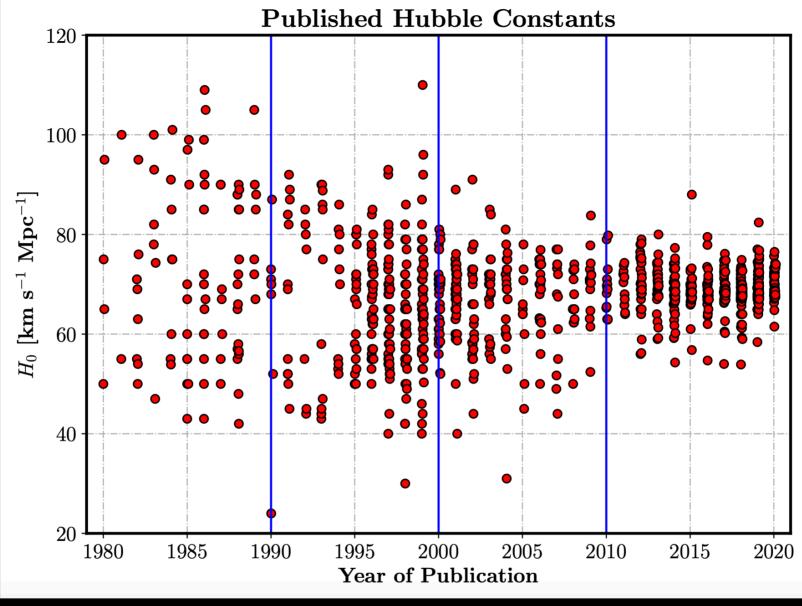
### Late universe measurements since 2020



### Late universe measurements since 2020



Abdalla et al., JHEAp 34 (2022) 49-211

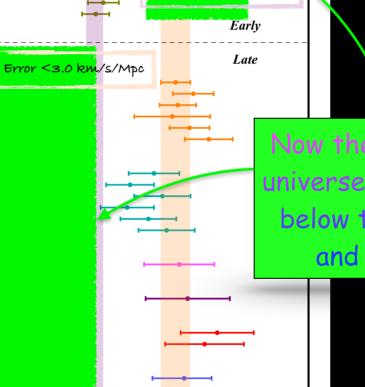


Freedman, Astrophys. J. 919 (2021) 1, 16

In the past the tension was within the same types of measurements and at the same redshifts and thus pointing directly to systematics.

### High precision measurements of Ho

Now there are no late universe measurements below the early ones and vice versa.



Error <1.5 km/s/MDC

 $H_{\theta}$  [km s<sup>-1</sup> Mpc<sup>-1</sup>]

### CMB with Planck

Balkenhol et al. (2021), Planck 2018+SPT+ACT : 67.49 ± 0.5 Aghanim et al. (2020), Planck 2018: 67.27 ± 0.60 Aghanim et al. (2020), Planck 2018+CMB lensing: 67.36 ± 0.54 Ade et al. (2016), Planck 2015: 67.27 ± 0.66

### CMB without Planck

Dutcher et al. (2021), SPT: 68.8 ± 1.5 Aiola et al. (2020), ACT: 67.9 ± 1.5 Aiola et al. (2020), WMAP9+ACT: 67.6 ± 1.1 Zhang, Huang (2019), WMAP9+BAO: 68.364\_0.52

### No CMB, with BBN

Zhang et al. (2022), BOSS correlation function+BAO+BBN: 68.19±0.99 -Chen et al. (2022), P+BAO+BBN: 69.23±0.77 -Philcox, Ivanov (2022), P+Bispectrum+BAO+BBN: 68.31±053 Colas et al. (2020), BOSS DR12+BBN: 67.9 ± 1.1 -Ivanov et al. (2020), BOSS+BBN: 67.9 ± 1.1 -Alam et al. (2020), BOSS+BBN: 67.9 ± 0.97 -

### SnIa-Cepheid

- Riess et al. (2022), R22: 73.04 ± 1.04
- Camarena, Marra (2021): 74.30 ± 1.45
- Riess et al. (2021), R21: 73.2 ± 1.3 ·
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  - Freedman (2021); 69.8 ± 1.7
  - Soltis, Casertano, Riess (2020): 72.1 ± 2.0
  - Freedman et al. (2020): 69.6 ± 1.9
  - Reid, Pesce, Riess (2019), SH0ES: 71.1 ± 1.99
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### SBF

Blakeslee et al. (2021) IR–SBF w/ HST: 73.3  $\pm$  2.5

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Pesce et al. (2020): 73.9 ± 3.0

### Tully Fisher

Kourkchi et al. (2020): 76.0 ± 2.6 Schombert, McGaugh, Lelli (2020): 75.1 ± 2.8

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  - Liao et al. (2020): 72.8+1.6 Liao et al. (2019): 72.2 ± 2.1
  - Shajib et al. (2019), STRIDES: 74.2+2.7
- Wong et al. (2019), H0LiCOW 2019: 73.3+1-7

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60

**65** 

70

75

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It is hard to conceive of a single type of systematic error that would apply to the measurements of the disparate phenomena we saw before as to effectively resolve the Hubble constant tension. Because the tension remains with the removal of the measurements of any single type of object, mode or calibration, it is challenging to devise a single error that would suffice. While multiple, unrelated systematic errors have a great deal more flexibility to resolve the tension but become less likely by

their inherent independence.

Since the early universe (indirect) constraints are model dependent, we can try to expand the cosmological scenario and see which extensions work in solving the tensions between the cosmological probes.

## Let's modify the $\Lambda CDM$ model...

### The Neutrino effective number

We can consider modifications in the dark matter sector.

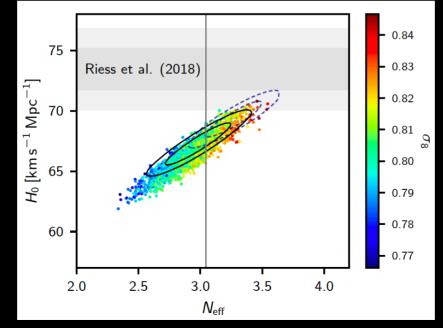
A classical extension is the effective number of relativistic degrees of freedom, i.e. additional relativistic matter at recombination, corresponding to a modification of the expansion history of the universe at early times.

### The Neutrino effective number

The expected value is Neff = 3.044, if we assume standard electroweak interactions and three active massless neutrinos. If we measure a Neff > 3.044, we are in presence of extra radiation.

If we vary Neff at 68% cl H0 passes is equal to  $66.4 \pm 1.4$  km/s/Mpc, and the tension with SH0ES increases from  $1.7\sigma$  to  $3.9\sigma$  also varying Neff.

 $N_{\rm eff} = 2.92^{+0.36}_{-0.37}$  (95%, *Planck* TT, TE, EE+lowE),



Planck 2018, Astron. Astrophys. 641 (2020) A6

### The Dark energy equation of state

# For example, we can consider modifications in the dark energy sector.

### A classical extension is a varying dark energy equation of state, that is a modification of the expansion history of the universe at late times.

## The Dark energy equation of state

If we change the cosmological constant with a Dark Energy with equation of state w, we are changing the expansion rate of the Universe:

$$H^{2} = \left(\frac{\dot{a}}{a}\right)^{2} = H_{0}^{2} \left(\frac{\Omega_{r}}{a^{4}} + \frac{\Omega_{m}}{a^{3}} + \frac{\Omega_{k}}{a^{2}} + \Omega_{\Lambda}\right)$$

$$H^{2} = H_{0}^{2} \left[ \Omega_{m} (1+z)^{3} + \Omega_{r} (1+z)^{4} + \Omega_{de} (1+z)^{3(1+w)} + \Omega_{k} (1+z)^{2} \right]$$

w introduces a geometrical degeneracy with the Hubble constant that is almost unconstrained using the CMB data only, resulting in agreement with SH0ES.

We have in 2018 w =  $-1.58^{+0.52}_{-0.41}$  with H0 > 69.9 km/s/Mpc at 95% c.l.

Planck data prefer a phantom dark energy, with an energy component with w < -1, for which the density increases with time in an expanding universe that will end in a Big Rip. A phantom dark energy violates the energy condition  $\rho \ge |p|$ , that means that the matter could move faster than light and a comoving observer measure a negative energy density, and the Hamiltonian could have vacuum instabilities due to a negative kinetic energy.

### Formally successful models in solving H0

tension $\leq 1\sigma$ "Excellent models"	tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"
Dark energy in extended parameter spaces [289]	Early Dark Energy [235]	Early Dark Energy [229]
Dynamical Dark Energy [309]	Phantom Dark Energy [11]	Decaying Warm DM [474]
Metastable Dark Energy [314]	Dynamical Dark Energy [11,281,309]	Neutrino-DM Interaction [506]
PEDE [392, 394]	GEDE [397]	Interacting dark radiation [517]
Elaborated Vacuum Metamorphosis [400–402]	Vacuum Metamorphosis [402]	Self-Interacting Neutrinos [700, 701]
IDE $[314, 636, 637, 639, 652, 657, 661-663]$	IDE [314, 653, 656, 661, 663, 670]	IDE [656]
Self-interacting sterile neutrinos [711]	Critically Emergent Dark Energy [997]	Unified Cosmologies [747]
Generalized Chaplygin gas model [744]	$f(\mathcal{T})$ gravity [814]	Scalar-tensor gravity [856]
Galileon gravity [876, 882]	Über-gravity [59]	Modified recombination [986]
Power Law Inflation [966]	Reconstructed PPS [978]	Super $\Lambda CDM$ [1007]
$f(\mathcal{T})$ [818]		Coupled Dark Energy [650]

Table B1. Models solving the  $H_0$  tension with R20 within the  $1\sigma$ ,  $2\sigma$  and  $3\sigma$ Planck only confidence levels considering the *Planck* dataset only.

Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

## Let's see an example...

There is a model considered in the early days of dark energy investigations that possesses the phenomenological properties needed to solve the H0 tension, but is based on a sound theoretical foundation: the vacuum metamorphosis model of Parker and Raval, Phys. Rev. D 62, 083503 (2000), Parker and Vanzella, Phys. Rev. D 69, 104009 (2004),

Caldwell, Komp, Parker and Vanzella, Phys. Rev. D 73, 023513 (2006), which has a phase transition in the nature of the vacuum.

Vacuum metamorphosis arises from a nonperturbative summation of quantum gravity loop corrections due to a massive scalar field.

We found that the Parker vacuum metamorphosis model, physically motivated by quantum gravitational effects, with the same number of parameters as LCDM, but not nested with it, can remove the H<sub>0</sub> tension, because can mimic a phantom DE behaviour at low redshifts.

First principles theory

When the Ricci scalar evolves during cosmic history to reach the scalar field mass squared m<sup>2</sup>, then a phase transition occurs and R freezes with

$$R = 6(\dot{H} + 2H^2 + ka^{-2}) = m^2$$
 and defining  $M = m^2/(12H_0^2)$ 

The expansion behaviour above and below the phase transition is

$$H^{2}/H_{0}^{2} = \Omega_{m}(1+z)^{3} + \Omega_{r}(1+z)^{4} + \Omega_{k}(1+z)^{2} + M \left\{ 1 - \left[ 3\left(\frac{4}{3\Omega_{m}}\right)^{4}M(1-M-\Omega_{k}-\Omega_{r})^{3} \right]^{-1} \right\}, \ z > z_{t}$$
$$H^{2}/H_{0}^{2} = (1-M-\Omega_{k})(1+z)^{4} + \Omega_{k}(1+z)^{2} + M, \quad z \le z_{t}$$

# with $z_t = -1 + rac{3\Omega_m}{4(1-M-\Omega_k-\Omega_r)}$

We see that above the phase transition, the universe behaves as one with matter (plus radiation plus spatial curvature) plus a constant, and after the phase transition it effectively has a dark radiation component that rapidly redshifts away leaving a de Sitter phase.

When the Ricci scalar evolves during cosmic history to reach the scalar field mass squared m<sup>2</sup>, then a phase transition occurs and R freezes with

$$R = 6(\dot{H} + 2H^2 + ka^{-2}) = m^2$$
 and defining  $M = m^2/(12H_0^2)$ 

The expansion behaviour above and below the phase transition is

$$\begin{split} H^2/H_0^2 &= \Omega_m (1+z)^3 + \Omega_r (1+z)^4 + \Omega_k (1+z)^2 + M \left\{ 1 - \left[ 3 \left( \frac{4}{3\Omega_m} \right)^4 M (1-M-\Omega_k - \Omega_r)^3 \right]^{-1} \right\}, \ z > z_t \\ H^2/H_0^2 &= (1-M-\Omega_k)(1+z)^4 + \Omega_k (1+z)^2 + M \,, \quad z \le z_t \end{split}$$

### with

$$z_t = -1 + \frac{3\Omega_m}{4(1 - M - \Omega_k - \Omega_r)}$$

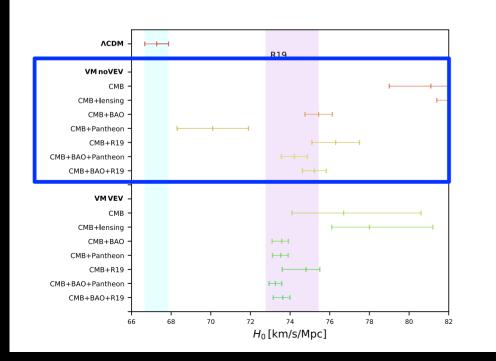
The original model did not include an explicit high redshift cosmological constant; we see that this implies that

$$\Omega_m = \frac{4}{3} \left[ 3M(1 - M - \Omega_k - \Omega_r)^3 \right]^{1/4}$$

i.e. the parameter M is fixed and depends on the matter density, and this model has the same number of degrees of freedom as ACDM.

### Constraints at 68% cl.

Parameters	CMB	CMB+lensing	CMB+BAO	CMB+Pantheon	CMB+R19	CMB+BAO+Pantheon	CMB+BAO+R19
$\Omega_b h^2$	$0.02238 \pm 0.00014$	$0.02242 \pm 0.00013$	$0.02218 \pm 0.00012$	$0.02201 \pm 0.00013$	$0.02221 \pm 0.00012$	$0.02213 \pm 0.00012$	$0.02217 \pm 0.00012$
$100 heta_{MC}$	$1.04091 \pm 0.00030$	$1.04097 \pm 0.00029$	$1.04060 \pm 0.00029$		$1.04063 \pm 0.00029$		$1.04060 \pm 0.00029$
au	$0.0524 \pm 0.0078$	$0.0510 \pm 0.0078$	$0.0458\substack{+0.0083\\-0.0067}$	$0.039\substack{+0.010\\-0.007}$	$0.0469 \pm 0.0075$	$0.0449\substack{+0.0079\\-0.0065}$	$0.0456\substack{+0.0083\\-0.0068}$
M	$0.9363\substack{+0.0055\\-0.0044}$	$0.9406 \pm 0.0034$	$0.9205 \pm 0.0023$	$0.8996\substack{+0.0081\\-0.0073}$	$0.9230\substack{+0.0042\\-0.0036}$	$0.9163 \pm 0.0023$	$0.9198 \pm 0.0020$
$\ln(10^{10}A_s)$	$3.041\pm0.016$	$3.036 \pm 0.015$	$3.035\substack{+0.017\\-0.014}$	$3.027\substack{+0.020\\-0.014}$	$3.036\pm0.016$	$3.035\substack{+0.017\\-0.014}$	$3.035\substack{+0.017\\-0.015}$
$n_s$	$0.9643 \pm 0.0039$	$0.9663 \pm 0.0036$	$0.9572 \pm 0.0031$	$0.9511 \pm 0.0036$	$0.9585 \pm 0.0033$	$0.9560 \pm 0.003$	$0.9571 \pm 0.0031$
$H_0[{ m km/s/Mpc}]$	$81.1 \pm 2.1$	$82.9 \pm 1.5$	$75.44 \pm 0.69$	$70.1 \pm 1.8$	$76.3\pm1.2$	$74.21 \pm 0.66$	$75.22\pm0.60$
$\sigma_8$	$0.9440 \pm 0.0077$	$0.9392 \pm 0.0067$	$0.9430_{-0.0070}^{+0.0082}$	$0.9419_{-0.0069}^{+0.0008}$	$0.9437 \pm 0.0073$	$0.9401_{-0.0068}^{+0.0080}$	$0.9457^{+0.0082}_{-0.0073}$
$S_8$	$0.805\pm0.022$	$0.783 \pm 0.014$	$0.865 \pm 0.010$	$0.927 \pm 0.023$	$0.856 \pm 0.015$	$0.880 \pm 0.010$	$0.8675 \pm 0.0098$
$\Omega_m$	$0.218\substack{+0.010 \\ -0.012}$	$0.2085 \pm 0.0076$	$0.2510 \pm 0.0046$	$0.291 \pm 0.015$	$0.2458\substack{+0.0074\\-0.0084}$	$0.2593 \pm 0.0046$	$0.2525 \pm 0.0040$
$\overline{\chi^2_{ m bf}}$	2767.74	2776.23	2806-22	$3874\ 13$	2777 04	3910.01	280.34
$\Delta \chi^2_{ m bf}$	-4.91	-5.81	+26.51	+66.63	-14.80	+95.83	+1.29



For the full dataset combinations H0~74 km/s/Mpc !!

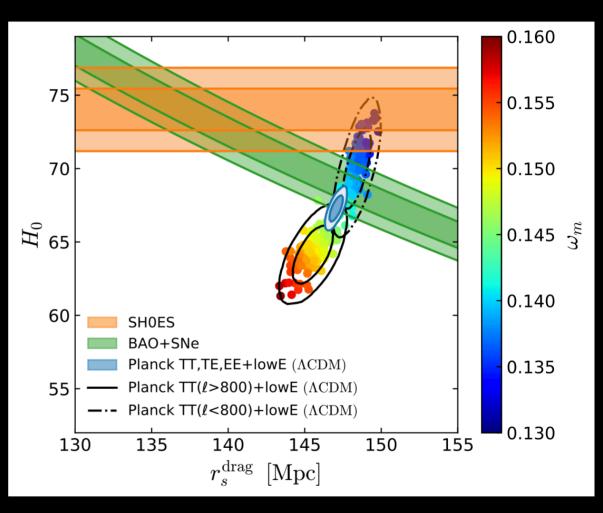
H0 is exactly in agreement with SH0ES even if BAO and Pantheon are included. However, this worsen considerably the fit of the data because the model fails in recover the shape of H(z) at low redshifts.

Di Valentino et al., Phys.Dark Univ. 30 (2020) 100733

### What about BAO+Pantheon?

BAO+Pantheon measurements constrain the product of H0 and the sound horizon r<sub>s</sub>.

In order to have a higher H0 value in agreement with SH0ES, we need r<sub>s</sub> near 137 Mpc. However, Planck by assuming  $\Lambda$ CDM, prefers r<sub>s</sub> near 147 Mpc. Therefore, a cosmological solution that can increase H0 and at the same time can lower the sound horizon inferred from CMB data is the most promising way to put in agreement all the measurements.

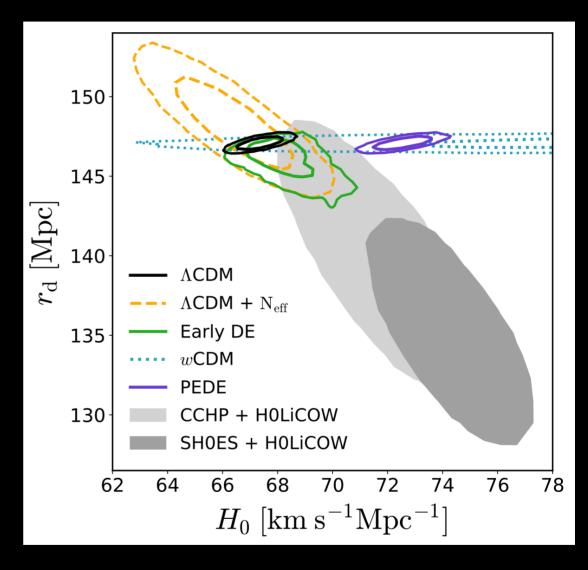


Knox and Millea, Phys. Rev. D 101 (2020) 4, 043533

## Early vs late time solutions

Here we can see the comparison of the 2 $\sigma$  credibility regions of the CMB constraints and the measurements from late-time observations (SN + BAO + H0LiCOW + SH0ES).

We see that the late time solutions, as wCDM, increase H0 because they decrease the expansion history at intermediate redshift, but leave r<sub>s</sub> unaltered.

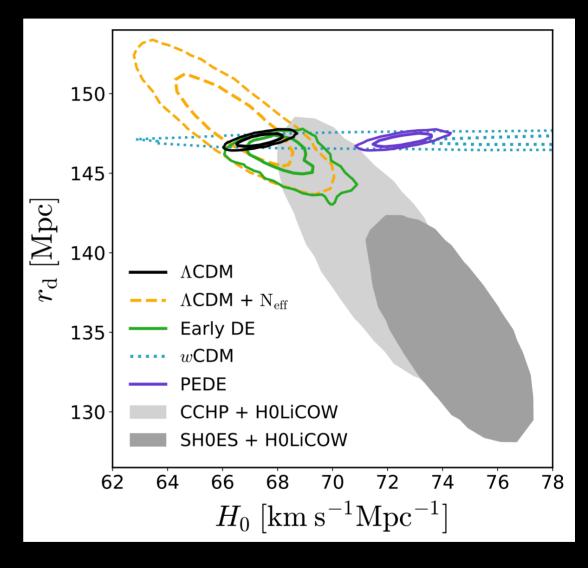


Arendse et al., Astron.Astrophys. 639 (2020) A57

## Early vs late time solutions

Here we can see the comparison of the 2 $\sigma$  credibility regions of the CMB constraints and the measurements from late-time observations (SN + BAO + H0LiCOW + SH0ES).

However, the early time solutions, as Neff or Early Dark Energy, move in the right direction both the parameters, but can't solve completely the H0 tension between Planck and SH0ES.



Arendse et al., Astron.Astrophys. 639 (2020) A57

### Formally successful models in solving H0

tension $\leq 1\sigma$ "Excellent models"	tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"
Early Dark Energy [228, 235, 240, 250]	Early Dark Energy [212, 229, 236, 263]	DE in extended parameter spaces [289]
Exponential Acoustic Dark Energy [259]	Rock 'n' Roll [242]	Dynamical Dark Energy [281, 309]
Phantom Crossing [315]	New Early Dark Energy [247]	Holographic Dark Energy [350]
Late Dark Energy Transition [317]	Acoustic Dark Energy [257]	Swampland Conjectures [370]
Metastable Dark Energy [314]	Dynamical Dark Energy [309]	MEDE [399]
PEDE [394]	Running vacuum model [332]	Coupled DM - Dark radiation [534]
Vacuum Metamorphosis [402]	Bulk viscous models [340, 341]	Decaying Ultralight Scalar [538]
Elaborated Vacuum Metamorphosis [401, 402]	Holographic Dark Energy [350]	$BD-\Lambda CDM$ [852]
Sterile Neutrinos [433]	Phantom Braneworld DE [378]	Metastable Dark Energy [314]
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IDE $[637, 639, 657, 661]$	IDE $[659, 670]$	IDE $[634-636, 653, 656, 663, 669]$
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combination of datasets **Table B2.** Models solving the  $H_0$  tension with R20 within  $1\sigma$ ,  $2\sigma$  and  $2\sigma$ *Planck* in combination with additional cosmological probes. datasets are discussed in the main text.

Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

### Let's see another example...

In the standard cosmological framework, DM and DE are described as separate fluids not sharing interactions beyond gravitational ones. At the background level, the conservation equations for the pressureless DM and DE components can be decoupled into two separate equations with an inclusion

of an arbitrary function, Q, known as the coupling or interacting function:

$$\dot{\rho}_c + 3\mathcal{H}\rho_c = Q,$$

$$\dot{\rho}_x + 3\mathcal{H}(1+w)\rho_x = -Q,$$

and we assume the phenomenological form for the interaction rate:

$$Q = \xi \mathcal{H} \rho_x$$

proportional to the dark energy density  $\rho_x$  and the conformal Hubble rate  $\mathcal{H}$ , via a negative dimensionless parameter  $\xi$  quantifying the strength of the coupling, to avoid early-time instabilities.

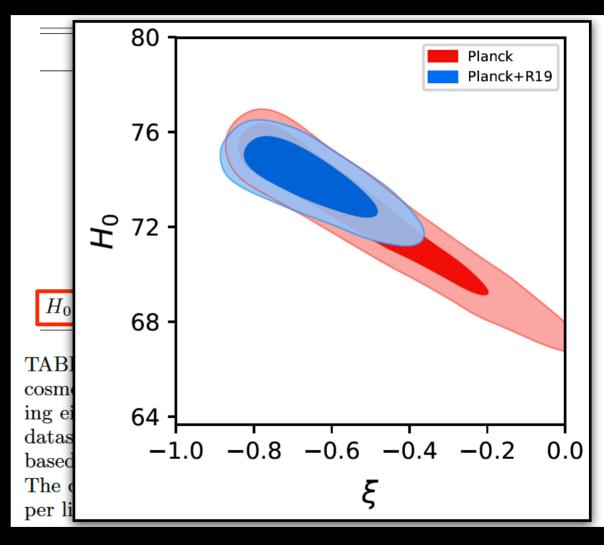
In this scenario of IDE the tension on H0 between the Planck satellite and SH0ES is completely solved. The coupling could affect the value of the present matter energy density  $\Omega_m$ . Therefore, if within an interacting model  $\Omega_m$  is smaller (because for negative  $\xi$  the dark matter density will decay into the dark energy one), a larger value of H0 would be required in order to satisfy the peaks structure of CMB observations, which accurately determine the value of  $\Omega_m h^2$ .

]	Parameter	Planck	Planck+R19
	$\Omega_{ m b}h^2$	$0.02239 \pm 0.00015$	$0.02239 \pm 0.00015$
[	$\Omega_{\rm c} h^2$	< 0.105	< 0.0615
	$n_s$	$0.9655 \pm 0.0043$	$0.9656 \pm 0.0044$
	$100\theta_{s}$	$1.0458\substack{+0.0033\\-0.0021}$	$1.0470 \pm 0.0015$
	au	$0.0541 \pm 0.0076$	$0.0534 \pm 0.0080$
	ξ	$-0.54^{+0.12}_{-0.28}$	$-0.66\substack{+0.09\\-0.13}$
$H_0$ []	${\rm kms^{-1}Mpc^{-1}}]$	$72.8^{+3.0}_{-1.5}$	$74.0^{+1.2}_{-1.0}$

TABLE I. Mean values with their 68% C.L. errors on selected cosmological parameters within the  $\xi \Lambda \text{CDM}$  model, considering either the *Planck* 2018 legacy dataset alone, or the same dataset in combination with the *R19* Gaussian prior on  $H_0$  based on the latest local distance measurement from *HST*. The quantity quoted in the case of  $\Omega_c h^2$  is the 95% C.L. upper limit.

Di Valentino et al., Phys.Dark Univ. 30 (2020) 100666

Therefore we can safely combine the two datasets together, and we obtain a nonzero dark matter-dark energy coupling  $\xi$  at more than FIVE standard deviations.



Di Valentino et al., Phys.Dark Univ. 30 (2020) 100666

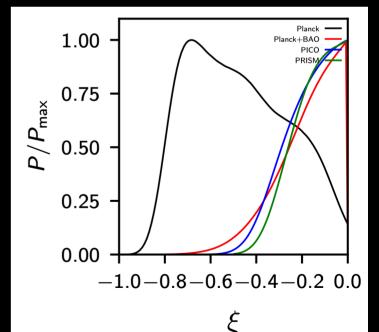
## fake IDE detection

Parameters	Fiducial model	Planck	Planck+BAO	PICO	PRISM
$\Omega_b h^2$	0.02236	$0.02238 \pm 0.00015$	$0.02230 \pm 0.00014$	$0.022364 \pm 0.000029$	$0.022361 \pm 0.000019$
$\Omega_c h^2$	0.1202	$0.056^{+0.025}_{-0.047}$	$0.101^{+0.019}_{-0.006}$	$0.100^{+0.019}_{-0.008}$	$0.103^{+0.016}_{-0.007}$
$100 \theta_{MC}$	1.04090	$1.0451_{-0.0032}^{+0.0021}$	$1.0419^{+0.0005}_{-0.0011}$	$1.04206^{+0.0005}_{-0.0011}$	$1.04191\substack{+0.00042\\-0.00094}$
au	0.0544	$0.0528^{+0.010}_{-0.009}$	$0.0517 \pm 0.0098$	$0.0543^{+0.0016}_{-0.0019}$	$0.0542^{+0.0017}_{-0.0019}$
$n_s$	0.9649	$0.9652 \pm 0.0041$	$0.9624 \pm 0.0036$	$0.9571 \pm 0.0014$	$0.9657 \pm 0.0012$
$\ln(10^{10}A_s)$	3.045	$3.041^{+0.020}_{-0.018}$	$3.042\pm0.019$	$3.0436^{+0.0030}_{-0.0034}$	$3.0435 \pm 0.0032$
ξ		$-0.48^{+0.16}_{-0.30}$	> -0.223	> -0.220	> -0.195

Di Valentino & Mena, Mon.Not.Roy.Astron.Soc. 500 (2020) 1, L22-L26, arXiv:2009.12620

For a simulated Planck-like experiment, due to the strong correlation present between the standard and the exotic physics parameters, there is a dangerous detection at more than  $3\sigma$  for a coupling between dark matter and dark energy different from zero, even if the fiducial model has  $\xi = 0$ :

 $-0.85 < \xi < -0.02$  at 99% CL



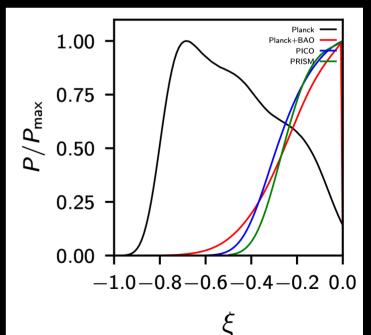
simulated experiments

## fake IDE detection

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au	0.0544	$0.0528_{-0.009}^{+0.010}$	$0.0517 \pm 0.0098$	$0.0543^{+0.0016}_{-0.0019}$	$0.0542^{+0.0017}_{-0.0019}$
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ξ	0	$-0.43^{+0.16}_{-0.30}$	> -0.223	> -0.220	> -0.195

Di Valentino & Mena, Mon.Not.Roy.Astron.Soc. 500 (2020) 1, L22-L26, arXiv:2009.12620

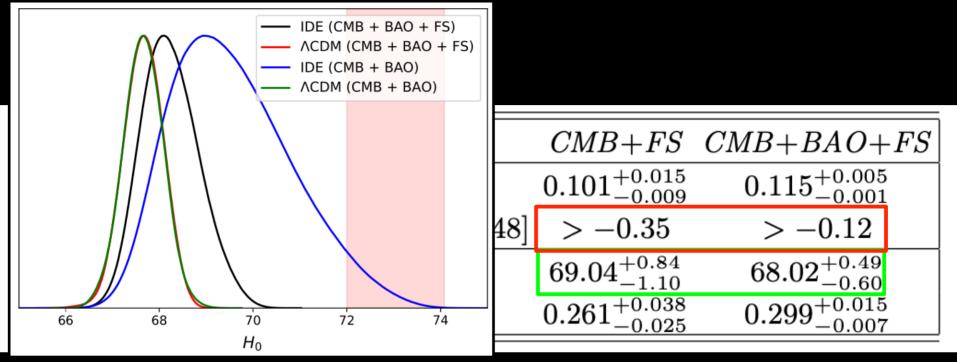
The inclusion of simulated BAO data, a mock dataset built using the same fiducial cosmological model than that of the CMB, helps in breaking the degeneracy, providing a lower limit for the coupling  $\xi$ in perfect agreement with zero.



simulated experiments

Parameter	CMB+BAO	CMB+FS	CMB + BAO + FS
$\omega_c$	$0.094\substack{+0.022\\-0.010}$	$0.101\substack{+0.015\\-0.009}$	$0.115\substack{+0.005\\-0.001}$
ξ	$-0.22_{-0.09}^{+0.18} [> -0.$	[48] > -0.35	> -0.12
$H_0  [{ m km/s/Mpc}]$	$69.55\substack{+0.98\\-1.60}$	$69.04\substack{+0.84 \\ -1.10}$	$68.02\substack{+0.49\\-0.60}$
$\Omega_m$	$0.243\substack{+0.054\\-0.030}$	$0.261\substack{+0.038\\-0.025}$	$0.299\substack{+0.015\\-0.007}$

The addition of low-redshift measurements, as BAO data, still hints to the presence of a coupling, albeit at a lower statistical significance. Also for this data sets the Hubble constant values is larger than that obtained in the case of a pure LCDM scenario, enough to bring the H0 tension at 2.1σ with SH0ES.



Nunes, Vagnozzi, Kumar, Di Valentino, and Mena, arXiv:2203.08093 [astro-ph.CO]

The addition of low-redshift measurements, as BAO data, still hints to the presence of a coupling, albeit at a lower statistical significance. Also for this data sets the Hubble constant values is larger than that obtained in the case of a pure LCDM scenario, enough to bring the H0 tension at 2.1σ with SH0ES.

However, the IDE model does not survive to the additional information coming from the full shape (FS) power spectrum of the BOSS DR12 galaxies.

### **Baryon Acoustic Oscillations**

BAO is formed in the early universe, when baryons are strongly coupled to photons, and the gravitational collapse due to the CDM is counterbalanced by the radiation pressure. Sound waves that propagate in the early universe imprint a characteristic scale on the CMB. Since the scale of these oscillations can be measured at recombination, BAO is considered a "standard ruler". These fluctuations have evolved and we can observe BAO at low redshifts in the distribution of galaxies. Since the data reduction process leading to these measurements requires assumptions about the fiducial cosmology, BAO is model dependent.

In other words, the tension between Planck+BAO or Planck+FS and SH0ES could be due to a statistical fluctuation in this case.

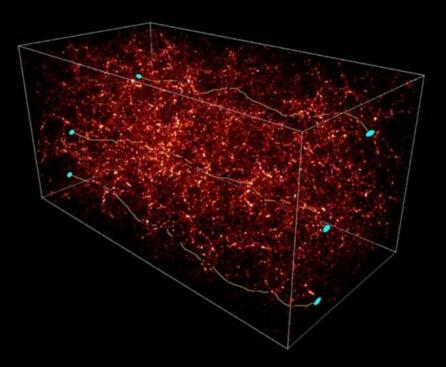
Actually, BAO and FS data are extracted under the assumption of LCDM, and the modified scenario of interacting dark energy could affect the result.

In fact, the full procedure which leads to the BAO and FS datasets carried out by the different collaborations might be not necessarily valid in extended DE models with important perturbations in the non-linear scales.

BAO and FS datasets (both the pre- and post- reconstruction measurements) might need to be revised in a non-trivial manner when applied to constrain more exotic dark energy cosmologies.

Additional complication: the models proposed to alleviate the H0 tension increase the S8 tension!

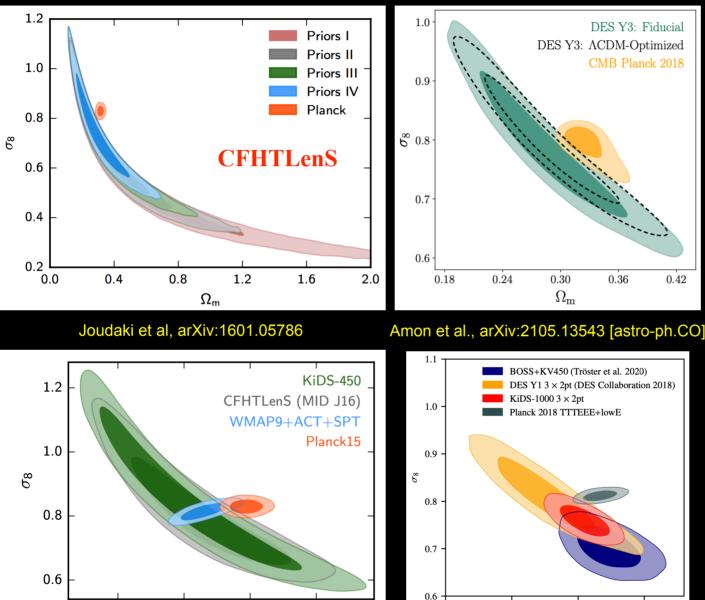
### The S8 tension



$$S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$$

A tension on S8 is present between the Planck data in the ΛCDM scenario and the cosmic shear data.

## The S8 tension



The S8 tension is now at 3.4σ between Planck assuming ΛCDM and KiDS+VIKING-450 and BOSS combined together, or 3.1σ with KiDS-1000, or 2.5σ with DES-Y3.

Hildebrandt et al., arXiv:1606.05338

0.24

0.32

 $\Omega_{m}$ 

0.40

0.16

Heymans et al., arXiv:2007.15632

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0.20

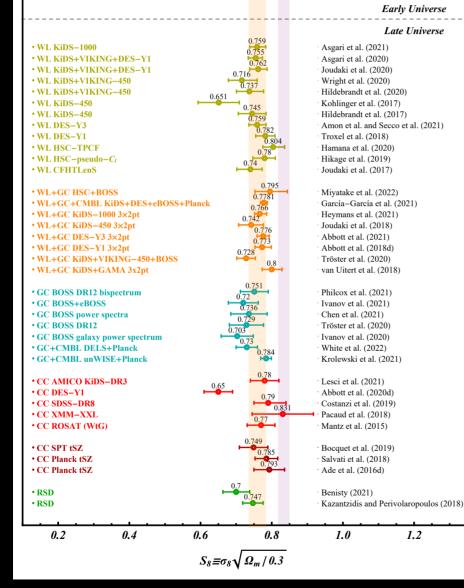
### The S8 tension

Aghanim et al. (2020d)

Aghanim et al. (2020d)

Aiola et al. (2020)

• CMB Planck TT,TE,EE+lowE • CMB Planck TT,TE,EE+lowE+lensing • CMB ACT+WMAP



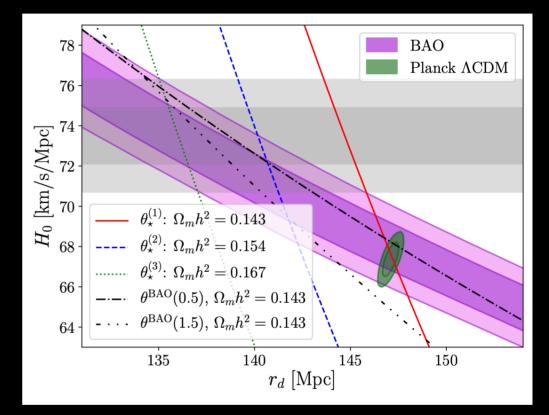
See Di Valentino et al. *Astropart.Phys.* 131 (2021) 102604 and Abdalla et al., arXiv:2203.06142 [astro-ph.CO] for a summary of the possible candidates proposed to solve the S8 tension.

Abdalla et al., *JHEAp* 34 (2022) 49-211

## Early solutions to the H0 tension

Actually, a dark energy model that merely changes the value of rd would not completely resolve the tension, since it will affect the inferred value of  $\Omega$ m and transfer the tension to it.

This is a plot illustrating that achieving a full agreement between CMB, BAO and SH0ES through a reduction of rd requires a higher value of  $\Omega_m h^2$ .



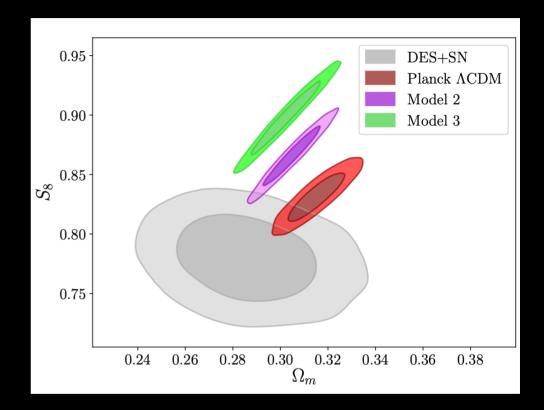
Jedamzik et al., Commun.in Phys. 4 (2021) 123

## Early solutions to the H0 tension

Model 2 is defined by the simultaneous fit to BAO and CMB acoustic peaks at  $\Omega_m h^2 = 0.155$ , while model 3 has  $\Omega_m h^2 = 0.167$ 

The sound horizon problem should be considered not only in the plane H0–rd, but it should be extended to the parameters triplet H0–rd– $\Omega$ m.

The figure shows that when attempting to find a full resolution of the Hubble tension, with CMB, BAO and SH0ES in agreement with each other, one exacerbates the tension with DES and KiDS.

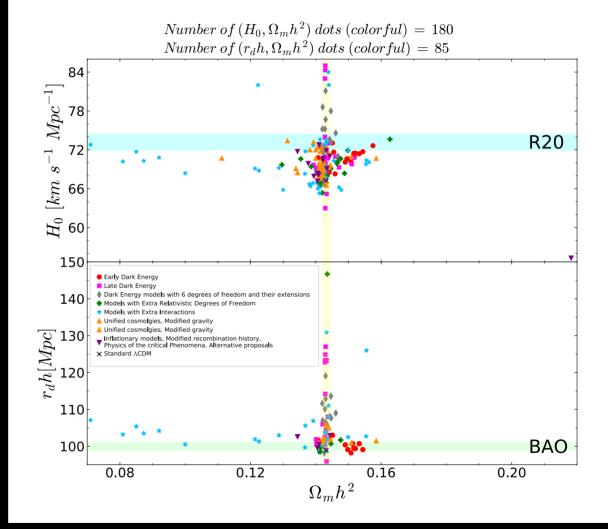


Jedamzik et al., Commun.in Phys. 4 (2021) 123

### Successful models?

This is the density of the proposed cosmological models:

At the moment no specific proposal makes a strong case for being highly likely or far better than all others !!!



Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

...but the excess of lensing in Planck could explain S8...

### A<sub>L</sub> internal anomaly

The gravitational effects of intervening dark matter fluctuations bend the path of CMB light on its way from the early universe to the CMB telescope. This "gravitational lensing" distorts our image of the CMB.

The lensing amplitude AL parameterizes the rescaling of the lensing potential  $\phi(n)$ , then the power spectrum of the lensing field:

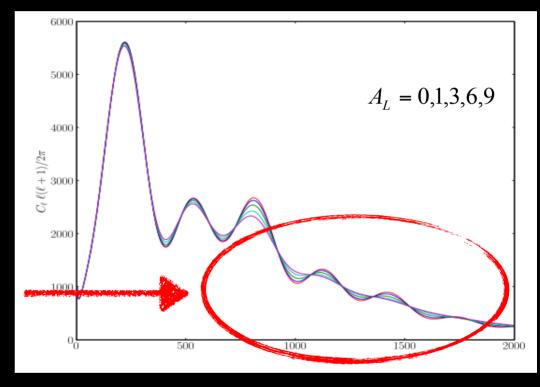
$$C_{\ell}^{\phi\phi} \to A_{\rm L} C_{\ell}^{\phi\phi}$$

### A<sub>L</sub> internal anomaly

The lensing effect on the power spectrum is the smoothing of the acoustic peaks, increasing AL.

Interesting consistency checks is if the amplitude of the smoothing effect in the CMB power spectra matches the theoretical expectation AL = 1 and whether the amplitude of the smoothing is consistent with that measured by the lensing reconstruction.

If AL =1 then the theory is correct, otherwise we have a new physics or systematics.



Calabrese et al., Phys. Rev. D, 77, 123531

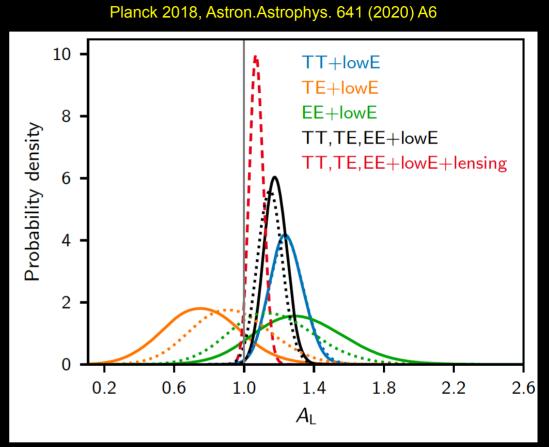
### A<sub>L</sub> : a failed consistency check

The Planck lensing-reconstruction power spectrum is consistent with the amplitude expected for LCDM models that fit the CMB spectra, so the Planck lensing measurement is compatible with AL = 1.

However, the distributions of AL inferred from the CMB power spectra alone indicate a preference for AL > 1.

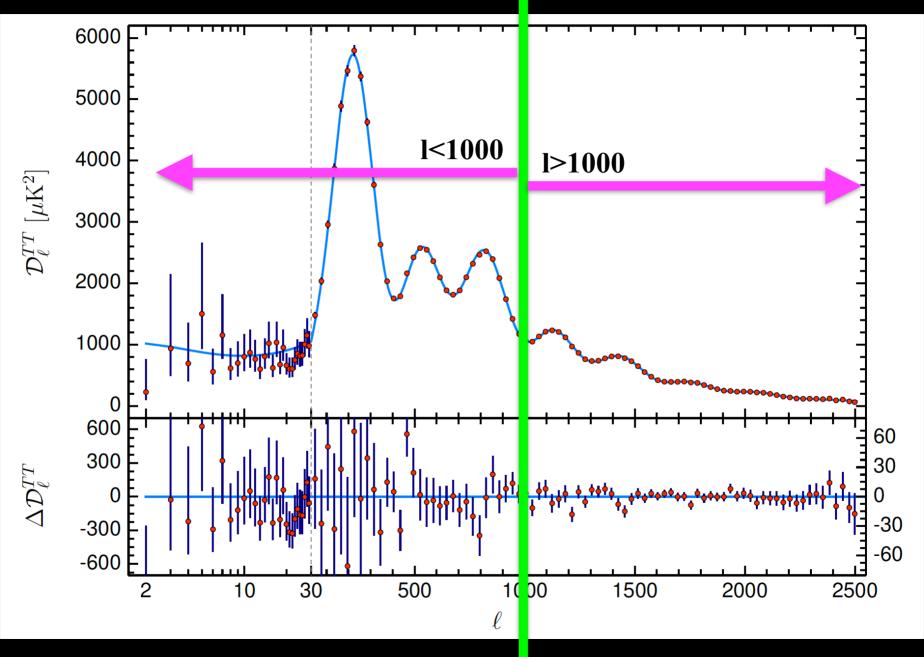
The joint combined likelihood shifts the value preferred by the TT data downwards towards AL = 1, but the error also shrinks, increasing the significance of AL > 1 to  $2.8\sigma$ .

The preference for high AL is not just a volume effect in the full parameter space, with the best fit improved by  $\Delta\chi^2 \sim 9$  when adding AL for TT+lowE and 10 for TTTEEE+lowE.

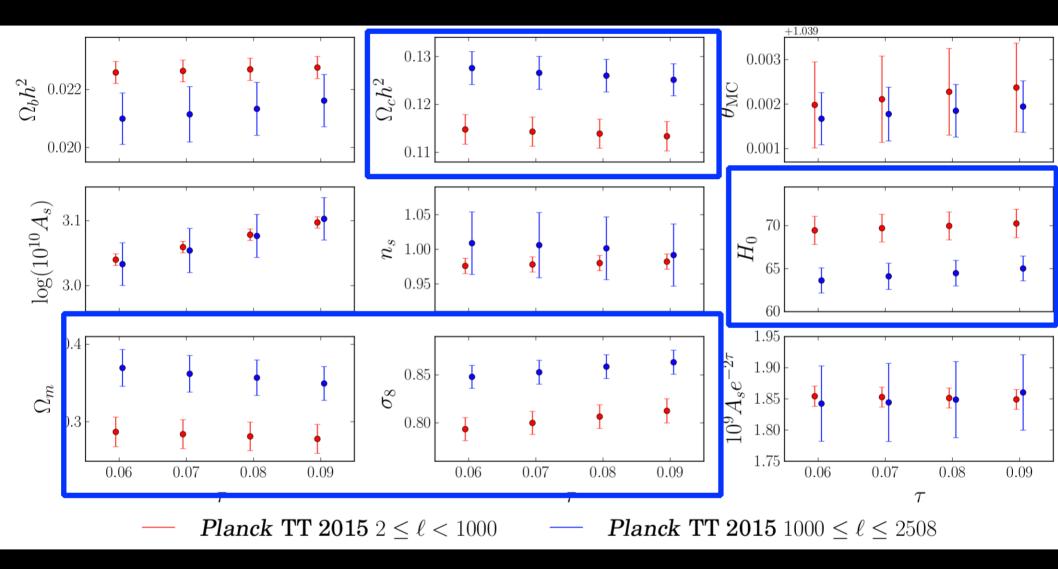


 $A_{\rm L} = 1.243 \pm 0.096$  (68 %, *Planck* TT+lowE),  $A_{\rm L} = 1.180 \pm 0.065$  (68 %, *Planck* TT,TE,EE+lowE),

### A<sub>L</sub> can explain internal tension



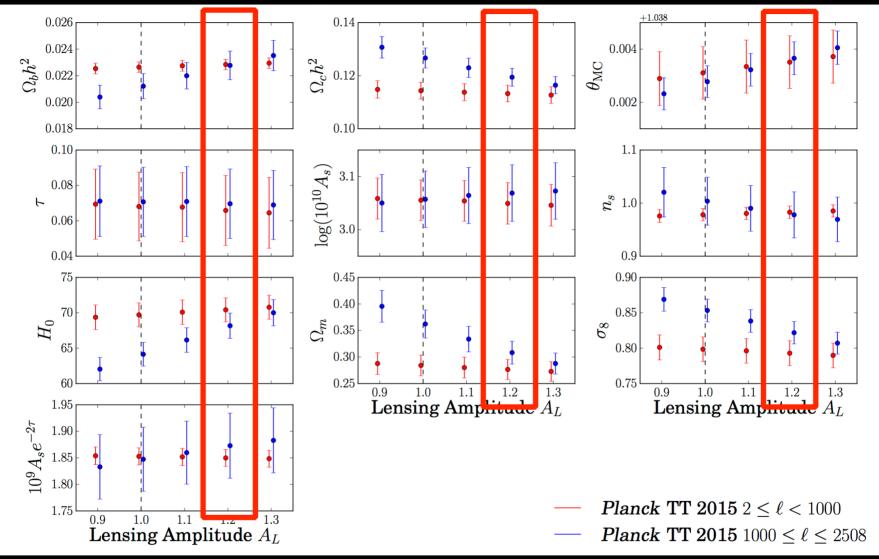
## A<sub>L</sub> can explain internal tension



Marginalized 68.3% confidence  $\Lambda$ CDM parameter constraints from fits to the I < 1000 and I ≥ 1000 Planck TT 2015 spectra. Tension at more than 2 $\sigma$  level appears in  $\Omega_c h^2$ and derived parameters, including H0,  $\Omega$ m, and  $\sigma$ 8.

Addison et al., Astrophys.J. 818 (2016) no.2, 132

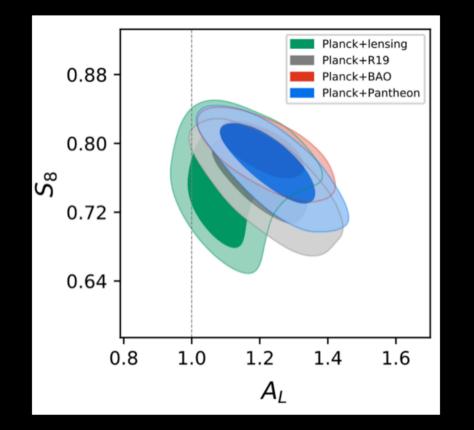
## A<sub>L</sub> can explain internal tension



Increasing AL smooths out the high order acoustic peaks, improving the agreement between the two multipole ranges.

Addison et al., Astrophys.J. 818 (2016) no.2, 132

# A<sub>L</sub> can explain the S8 tension



Di Valentino, Melchiorri and Silk, JCAP 2001 (2020) no.01, 013

A<sub>L</sub> that is larger than the expected value at about 3 standard deviations even when combining the Planck data with BAO and supernovae type Ia external datasets.

# $A_{L}$ can be an indication for Modified Gravity models...

# MG could explains AL

Assuming a flat universe, we can write the line element of the Friedmann-Lemaitre-Robertson-Walker (FLRW) metric in the conformal Newtonian gauge as:

$$ds^{2} = a(\tau)^{2} \left[ -(1+2\Psi)d\tau^{2} + (1-2\Phi)dx^{i}dx_{i} \right]$$

where a is the scale factor,  $\tau$  is the conformal time,  $\Psi$  is the Newton's gravitational potential, and  $\Phi$  the space curvature. We can use a phenomenological parametrization of the gravitational potentials  $\Psi$  and  $\Phi$  and their combinations given by:

•  $\mu(k, a)$  modifies the Poisson equation for the Newton's gravitational potential  $\Psi$ 

$$k^2\Psi = -4\pi G a^2 \mu(k,a) \rho \Delta$$

η(k,a) takes into account the presence of a non-zero anisotropic stress, with Φ the space curvature:

$$\eta(k,a) = rac{\Phi}{\Psi}\,.$$

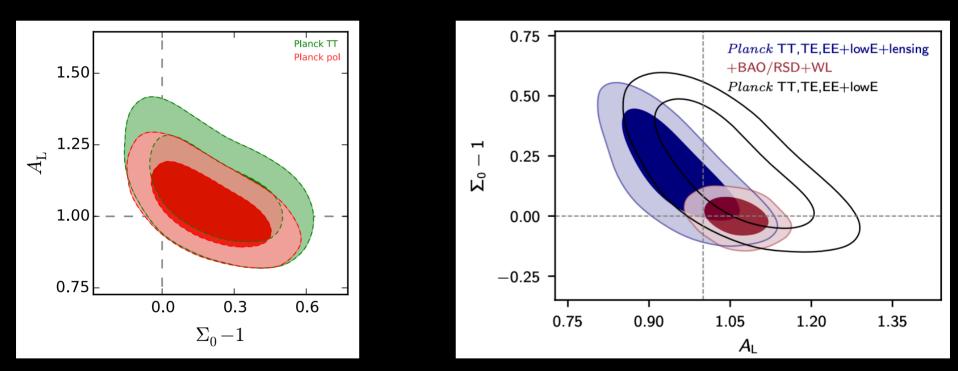
•  $\Sigma(k, a)$  modifies the lensing/Weyl potential  $\Phi+\Psi$ :

$$-k^2(\Phi+\Psi)\equiv 8\pi Ga^2\Sigma(k,a)
ho\Delta$$

$$\Sigma = rac{\mu}{2}(1+\eta)$$
 .

## MG could explains AL

A strong degeneracy is present between  $\Sigma 0$  and AL: if we fix  $\Sigma 0=1$  we have a larger value for AL, but when AL =1 then some indication for  $\Sigma 0>1$  appears.

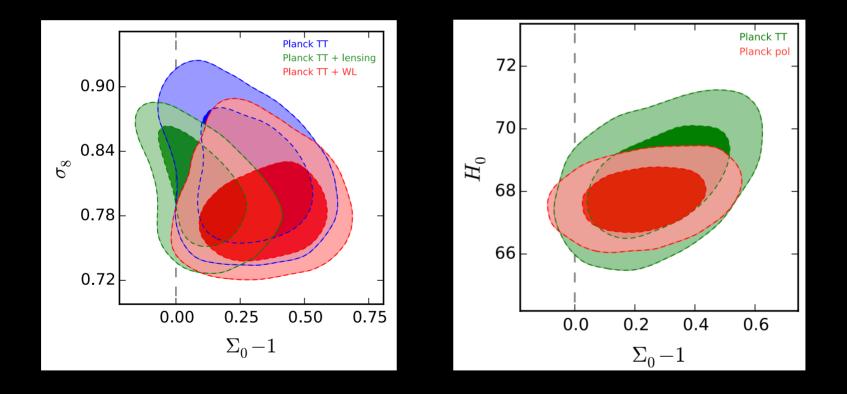


Di Valentino et al., Phys.Rev. D93 (2016) no.2, 023513

Planck 2018, Astron. Astrophys. 641 (2020) A6

# MG could explains S8 and H0

The constraints on the amplitude of matter density fluctuations σ8 are relaxed and in better agreement with weak lensing measurements. Moreover, we have a positive correlation with H0, potentially solving the Hubble constant tension.



# MG could explains H0

tension $\leq 1\sigma$ "Excellent models"	tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"
Early Dark Energy [228, 235, 240, 250]	Early Dark Energy [212, 229, 236, 263]	DE in extended parameter spaces [289]
Exponential Acoustic Dark Energy [259]	Rock 'n' Roll [242]	Dynamical Dark Energy [281, 309]
Phantom Crossing [315]	New Early Dark Energy [247]	Holographic Dark Energy [350]
Late Dark Energy Transition [317]	Acoustic Dark Energy [257]	Swampland Conjectures [370]
Metastable Dark Energy [314]	Dynamical Dark Energy [309]	MEDE [399]
PEDE [394]	Running vacuum model [332]	Coupled DM - Dark radiation [534]
Vacuum Metamorphosis [402]	Bulk viscous models $[340, 341]$	Decaying Ultralight Scalar [538]
Elaborated Vacuum Metamorphosis [401, 402]	Holographic Dark Energy [350]	BD-ACDM [852]
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Ginzburg-Landau theory [996]	Early modified gravity [859]	Primordial Black Holes [545]
Lorentzian Quintessential Inflation [979]	Bianchi type I spacetime [999]	
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**Table B2.** Models solving the  $H_0$  tension with R20 within  $1\sigma$ ,  $2\sigma$  and  $3\sigma$  considering *Planck* in combination with additional cosmological probes. Details of the combined datasets are discussed in the main text.

Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

...or assuming General Relativity, a curved universe can be a physical explanation for AL...

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#### Planck 2018 results. VI. Cosmological parameters

Planck Collaboration: N. Aghanim<sup>54</sup>, Y. Akrami<sup>15,57,59</sup>, M. Ashdown<sup>65,5</sup>, J. Aumont<sup>95</sup>, C. Baccigalupi<sup>78</sup>, M. Ballardini<sup>21,41</sup>, A. J. Banday<sup>95,8</sup>,
 R. B. Barreiro<sup>61</sup>, N. Bartolo<sup>29,62</sup>, S. Basak<sup>85</sup>, R. Battye<sup>64</sup>, K. Benabed<sup>55,90</sup>, J.-P. Bernard<sup>95,8</sup>, M. Bersanelli<sup>32,45</sup>, P. Bielewicz<sup>75,78</sup>, J. J. Bock<sup>63,10</sup>,
 J. R. Bond<sup>7</sup>, J. Borrill<sup>12,93</sup>, F. R. Bouchet<sup>55,90</sup>, F. Bou<sup>1</sup>

J.-F. Cardoso<sup>55,90</sup>, J. Carron<sup>23</sup>, A. Challinor<sup>58,65,11</sup>, H. C. Chia F. Cuttaia41, P. de Bernardis31, G. de Zotti42, J. Delabroui A. Ducout<sup>66</sup>, X. Dupac<sup>35</sup>, S. Dusini<sup>62</sup>, G. Efstathiou<sup>65</sup> J. Fergusson<sup>11</sup>, R. Fernandez-Cobos<sup>61</sup>, F. Finelli<sup>41,47</sup>, F. For S. Galli55,90<sup>†</sup>, K. Ganga<sup>2</sup>, R. T. Génova-Santos<sup>60,16</sup>, M. A. Gruppuso<sup>41,47</sup>, J. E. Gudmundsson<sup>94,25</sup>, J. Hamann<sup>86</sup>, Z. Huang<sup>83</sup>, A. H. Jaffe<sup>53</sup>, W. C. Jones<sup>25</sup>, A. Karakci<sup>59</sup>, N. Krachmalnicoff78, M. Kunz14,54,3, H. Kurki-Suonio24,4 M. Le Jeune<sup>2</sup>, P. Lemos<sup>58,65</sup>, J. Lesgourgues<sup>56</sup>, F. Levri M. López-Caniego35, P. M. Lubin28, Y.-Z. Ma77,80,74, J. A. Marcos-Caballero<sup>61</sup>, M. Maris<sup>43</sup>, P. G. Martin<sup>7</sup>, M. Mar P. R. Meinhold<sup>28</sup>, A. Melchiorri<sup>31,50</sup>, A. Mennella<sup>32,45</sup> D. Molinari<sup>30,41,48</sup>, L. Montier<sup>95,8</sup>, G. Morgante<sup>41</sup>, A. Mo B. Partridge<sup>39</sup>, G. Patanchon<sup>2</sup>, H. V. Peiris<sup>22</sup>, F. Perrott J. P. Rachen<sup>18</sup>, M. Reinecke<sup>72</sup>, M. Remazeilles<sup>64</sup>, A. B. Ruiz-Granados<sup>60,16</sup>, L. Salvati<sup>54</sup>, M. Sandri<sup>41</sup>, M. Savelair R. Sunyaev<sup>72,91</sup>, A.-S. Suur-Uski<sup>24,40</sup>, J. A. Tauber<sup>36</sup>, D. Valenziano41, J. Valiviita24,40, B. Van Tent69, L. Vibert54. S. D. M. W (Affiliation

We present cosmological parameter results from the final isotropies, combining information from the temperature an improved measurements of large-scale polarization allow th cant gains in the precision of other correlated parameters. In many parameters, with residual modelling uncertainties estim spatially-flat 6-parameter ACDM cosmology having a power from polarization, temperature, and lensing, separately and i baryon density  $\Omega_b h^2 = 0.0224 \pm 0.0001$ , scalar spectral inde 68 % confidence regions on measured parameters and 95 %  $100\theta_* = 1.0411 \pm 0.0003$ . These results are only weakly dependent in many commonly considered extensions. Assuming the ba Hubble constant  $H_0 = (67.4 \pm 0.5) \text{ km s}^{-1} \text{Mpc}^{-1}$ ; matter dens We find no compelling evidence for extensions to the base-A considering single-parameter extensions) we constrain the effective the Standard Model prediction  $N_{\text{eff}} = 3.046$ , and find that the to prefer higher lensing amplitudes than predicted in base AC from the ACDM model; however, this is not supported by BAO data. The joint constraint with BAO measurements on s with Type Ia supernovae (SNe), the dark-energy equation of constant. We find no evidence for deviations from a purely Keck Array data, we place a limit on the tensor-to-scalar ra deuterium abundances for the base-ACDM cosmology are in agreement with BAO, SNe, and some galaxy lensing observ including galaxy clustering (which prefers lower fluctuation measurements of the Hubble constant (which prefer a high favoured by the Planck data.

 $\Omega_K = -0.044^{+0.018}_{-0.015}$  (68%, *Planck* TT, TE, EE+lowE), (46b)

Page 40

a detection of curvature at about  $3.4\sigma$ 

an apparent detection of curvature at well over  $2\sigma$ . The 99% probability region for the TT,TE,EE+lowE result is  $-0.095 < \Omega_K < -0.007$ , with only about 1/10000 samples at  $\Omega_K \ge 0$ . This is not entirely a volume effect, since the best-fit  $\chi^2$  changes by  $\Delta \chi^2_{\text{eff}} = -11$  compared to base  $\Lambda$ CDM when adding the one additional curvature parameter. The reasons for the pull towards

\*Corresponding author: G. Efstathiou, gpe@ast.cam.ac.uk <sup>†</sup>Corresponding author: S. Galli, gallis@iap.fr <sup>‡</sup>Corresponding author: A. Lewis, antony@cosmologist.info

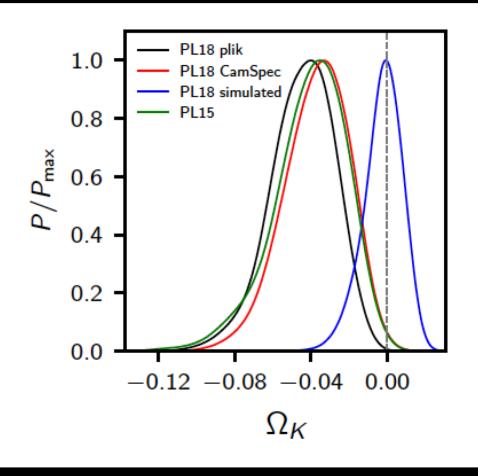
\*Corresponding author: G. Efstathio. gpe@est.can.ac.uk Vature of the Universe <sup>†</sup>Corresponding author: S. Galli, gall.cap... <sup>†</sup>Corresponding author: A. Lewis, antony@cosmologist.info

2019

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### **Curvature of the universe**

Can Planck provide an unbiased and reliable estimate of the curvature of the Universe? This may not be the case since a "geometrical degeneracy" is present with  $\Omega m$ . When precise CMB measurements at arc-minute angular scales are included, since gravitational lensing depends on the matter density, its detection breaks the geometrical degeneracy. The Planck experiment with its improved angular resolution offers the unique opportunity of a precise measurement of curvature from a single CMB experiment. We simulated Planck, finding that such experiment could constrain curvature with a 2% uncertainty, without any significant bias towards closed models.



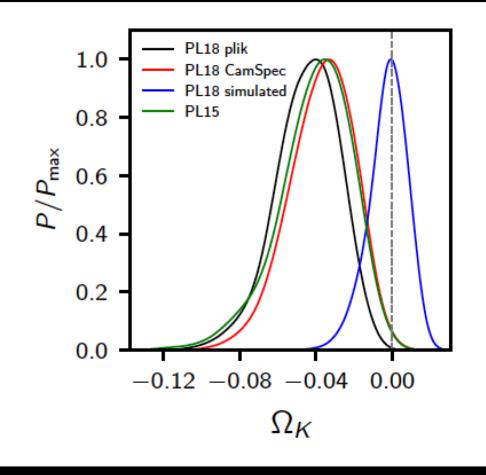
Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

### **Curvature of the universe**

Planck favours a closed Universe  $(\Omega k < 0)$  with 99.985% probability. A closed Universe with  $\Omega K = -0.0438$ provides a better fit to PL18 with respect to a flat model.

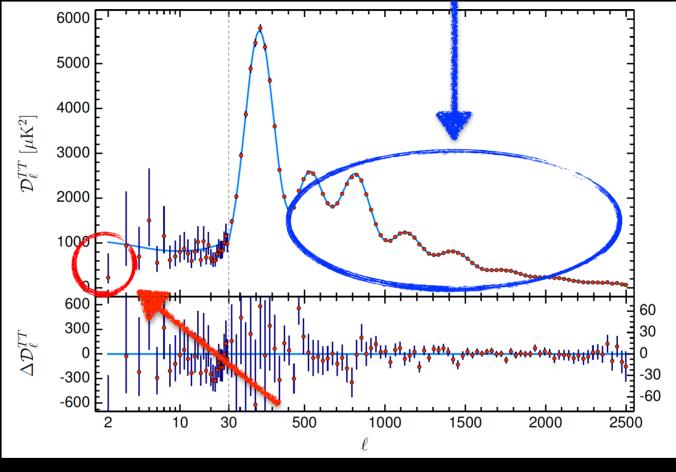
This is not entirely a volume effect, since the best-fit  $\Delta \chi^2$  changes by -11 compared to base ACDM when adding the one additional curvature parameter. The improvement is due also to the fact that closed models could also lead to a large-scale cut-off in the primordial density fluctuations in agreement with the observed low

CMB anisotropy quadrupole.



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

# Low CMB anisotropy quadrupole



Planck 2018, Astron.Astrophys. 641 (2020) A6

A model with  $\Omega \kappa < 0$  is slightly preferred with respect to a flat model with AL > 1, because closed models better fit not only the damping tail, but also the lowmultipole data, especially the quadrupole.

#### Astrophysics

[Submitted on 5 Mar 2003 (v1), last revised 30 Jul 2003 (this version, v2)]

### Is the Low CMB Quadrupole a Signature of Spatial Curvature?

#### G. Efstathiou (University of Cambridge)

The temperature anisotropy power spectrum measured with the Wilkinson Microwave Anisotropy Probe (WMAP) at high multipoles is in spectacular agreement with an inflationary Lambda-dominated cold dark matter cosmology. However, the low order multipoles (especially the quadrupole) have lower amplitudes than expected from this cosmology, indicating a need for new physics. Here we speculate that the low quadrupole amplitude is associated with spatial curvature. We show that positively curved models are consistent with the WMAP data and that the quadrupole amplitude can be reproduced if the primordial spectrum

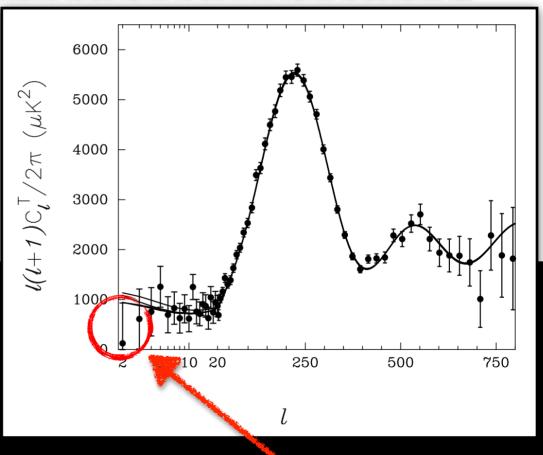
truncates on scales comparable to the curvature scale.

Comments:	4 pages, Latex, 2 figs, revised version accepted by MNRAS
Subjects:	Astrophysics (astro-ph)
Journal reference:	Mon.Not.Roy.Astron.Soc. 343 (2003) L95
DOI:	10.1046/j.1365-8711.2003.06940.x
Cite as:	arXiv:astro-ph/0303127
	(or arXiv:astro-ph/0303127v2 for this version)

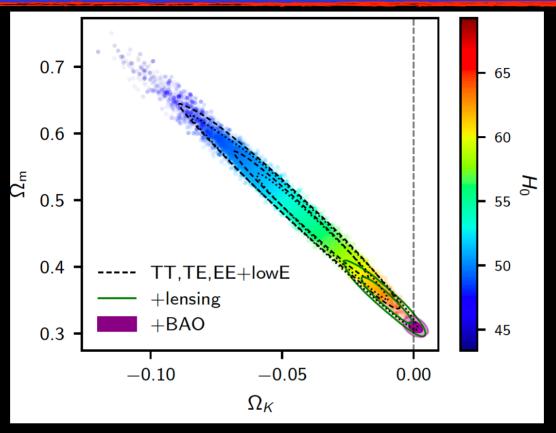
#### **Submission history**

From: George Efstathiou [view email] [v1] Wed, 5 Mar 2003 23:30:33 UTC (21 KB) [v2] Wed, 30 Jul 2003 10:16:45 UTC (22 KB)

A lower quadrupole than predicted by the ACDM was already present in WMAP, and a closed universe to explain this effect was already taken into account.



### What about Planck+BAO?



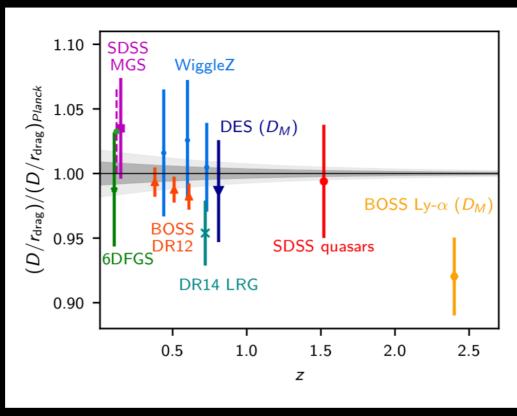
Planck 2018, Astron.Astrophys. 641 (2020) A6

Adding BAO data, a joint constraint is very consistent with a flat universe.

 $\Omega_K = 0.0007 \pm 0.0019$  (68 %, TT, TE, EE+lowE +lensing+BAO).

Given the significant change in the conclusions from Planck alone, it is reasonable to investigate whether they are actually consistent. In fact, a basic assumption for combining complementary datasets is that these ones must be consistent, i.e. they must plausibly arise from the same cosmological model.

# **BAO tension**

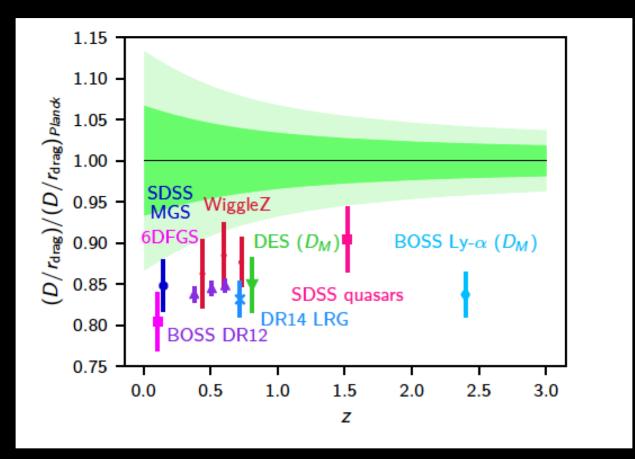


Planck 2018, Astron. Astrophys. 641 (2020) A6

This is a plot of the acoustic-scale distance ratio, DV(z)/rdrag, as a function of redshift, taken from several recent BAO surveys, and divided by the mean acoustic-scale ratio obtained by Planck adopting a model. rdrag is the comoving size of the sound horizon at the baryon drag epoch, and DV, the dilation scale, is a combination of the Hubble parameter H(z) and the comoving angular diameter distance DM(z).

In a ACDM model the BAO data agree really well with the Planck measurements...

## **BAO tension**

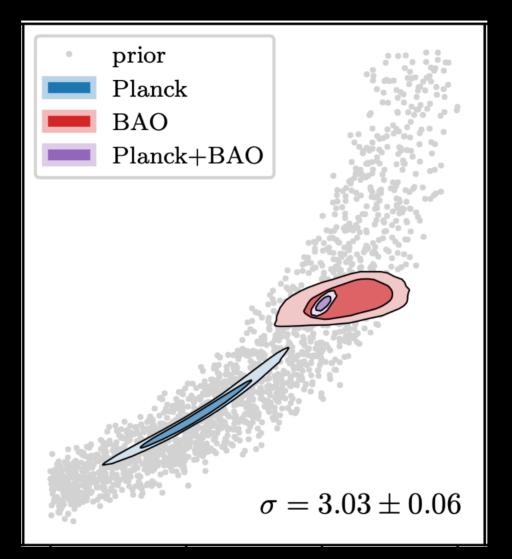


Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

... but when we let curvature to vary

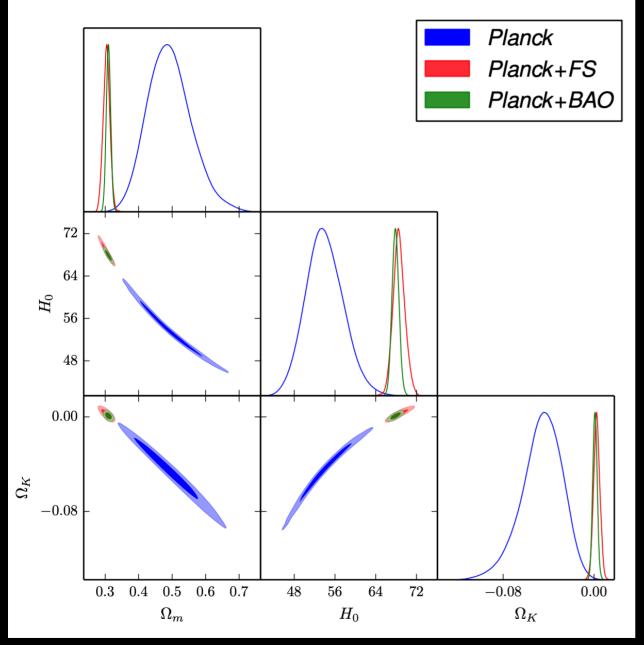
there is a striking disagreement between Planck spectra and BAO measurements!

# **BAO tension**



In agreement with Handley, Phys.Rev.D 103 (2021) 4, L041301

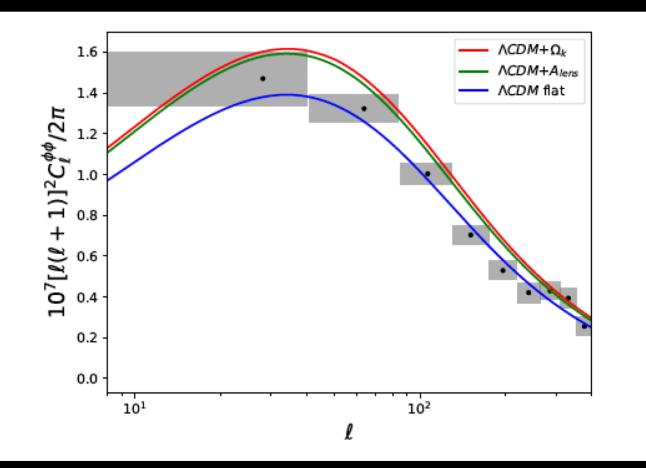
### What about Planck+FS?



The strong disagreement between Planck and BAO it is evident in this triangular plot, as well as that with the full-shape (FS) galaxy power spectrum measurements from the BOSS DR12 CMASS sample, at an effective redshift  $z_{eff} = 0.57$ .

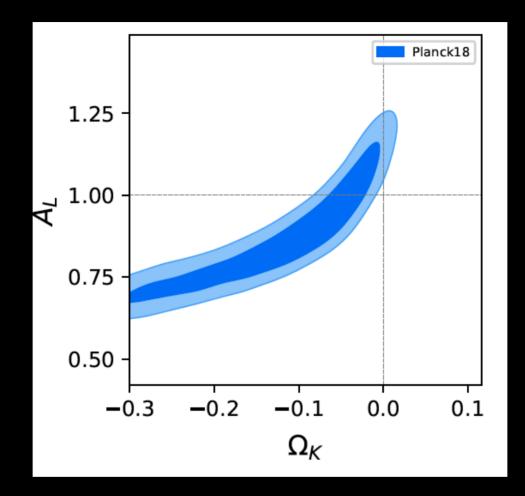
Vagnozzi, Di Valentino, et al., Phys.Dark Univ. 33 (2021) 100851

Closed models predict substantially higher lensing amplitudes than in  $\Lambda$ CDM, because the dark matter content can be greater, leading to a larger lensing signal. The reasons for the pull towards negative values of  $\Omega_K$  are essentially the same as those that lead to the preference for AL > 1.



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

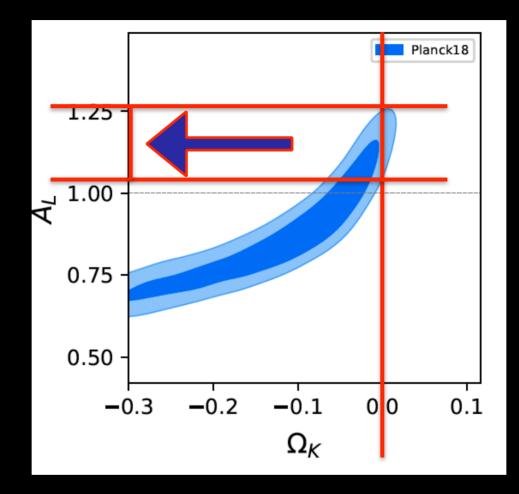
### A closed universe (Friedmann 1922) can explain AL!



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

A degeneracy between curvature and the AL parameter is clearly present. A closed universe can provide a robust physical explanation to the enhancement of the lensing amplitude. In fact, the curvature of the Universe is not new physics beyond the standard model, but it is predicted by the General Relativity, and depends on the energy content of the Universe.

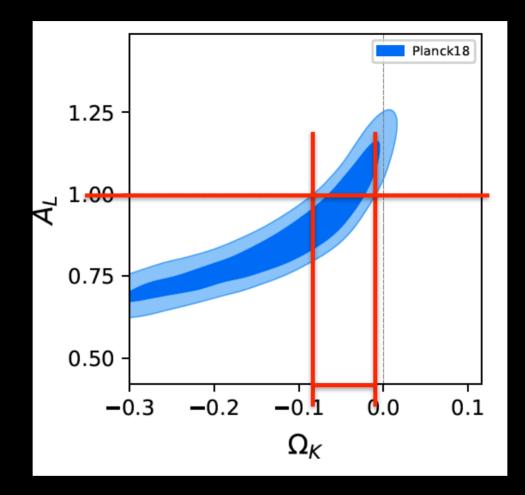
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### The evolution over time of the geometry of the universe is described by Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^2}T_{\mu\nu} + \Lambda g_{\mu\nu}$$

Adopting a 4-dimensional coordinate system for the space-time and the Cosmological Principle, i.e. a universe homogeneous and isotropic at large scales, the resulting metric is the Friedmann-Lemaitre-Robertson-Walker (FLRW), that describes the distance between two events in space-time.

$$ds^{2} = c^{2}dt^{2} - a^{2}(t)\left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}\left(d\theta^{2} + \sin^{2}\theta d\varphi^{2}\right)\right]$$

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The curvature parameter k can be positive, null or negative, depending on the value of the curvature of the universe: positive, flat or negative.

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^2}T_{\mu\nu} + \Lambda g_{\mu\nu}$$

Adopting a 4-dimensional coordinate system for the space-time and the Cosmological Principle, i.e. a universe homogeneous and isotropic at large scales, the resulting metric is the Friedmann-Lemaitre-Robertson-Walker (FLRW), that describes the distance between two events in space-time.

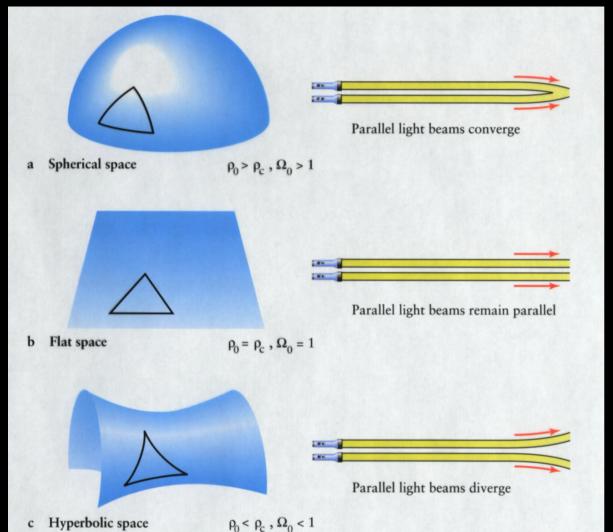
$$ds^{2} = c^{2}dt^{2} - a^{2}(t) \left[ \frac{dr^{2}}{1 - kr^{2}} + r^{2} \left( d\theta^{2} + \sin^{2}\theta d\varphi^{2} \right) \right]$$

Combining together the FLRW metric and Einstein's equations we obtain the Friedmann equations that describe the expansion history of the universe:

1st 
$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} + \frac{\Lambda}{3}$$
  
2nd  $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + 3P\right) + \frac{\Lambda}{3}$ 

The evolution over time of the geometry of the universe is described by Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^2}T_{\mu\nu} + \Lambda g_{\mu\nu}$$



If we divide the 1st Friedmann equation, for the critical density (density of a flat universe), we obtain today:

$$\Omega = \sum_{i} \Omega_{i} = \Omega_{m} + \Omega_{\Lambda} + \Omega_{r} = 1 - \Omega_{k}$$

From this equation it is possible to estimate the curvature of the universe, independently measuring the various contributions to the total density parameter  $\Omega$ .

#### Figure: http://w3.phys.nthu.edu.tw

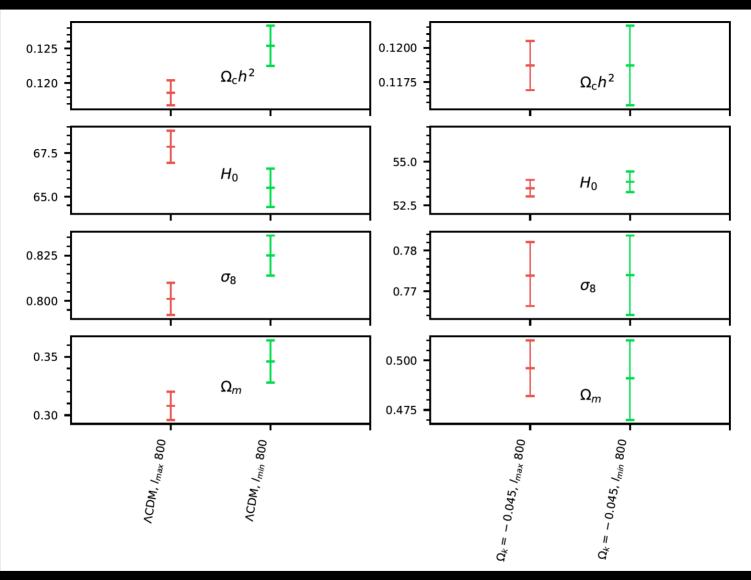
$$\begin{cases} \Omega > 1 \quad \Omega_k < 0 \\ \Omega = 1 \quad \Omega_k = 0 \\ \Omega < 1 \quad \Omega_k > 0 \end{cases} \xrightarrow{k > 0} \xrightarrow{k > 0} \xrightarrow{k > 0} \xrightarrow{k < 0} : \text{ closed Universe} \\ \xrightarrow{k < 0} : \text{ flat Universe} \\ \xrightarrow{k < 0} : \text{ open Universe} \end{cases}$$

### LCDM+ $\Omega_k$ : a 7 parameter standard model

As it has been convincingly pointed out in Anselmi et al., arXiv:2207.06547, absent any theoretical arguments,

we cannot use observations that suggest small Ω<sub>k</sub> to enforce Ω<sub>k</sub>=0.
The common practice to set Ω<sub>k</sub>=0 places the onus on proponents of "curved LCDM" to present sufficient evidence that Ω<sub>k</sub>≠0, and this is needed as an additional parameter.
Given the current tensions in cosmological parameters and CMB anomalies this choice is at least debatable.
So it would be desirable to have the standard cosmological phenomenological model with at least 7 parameters.

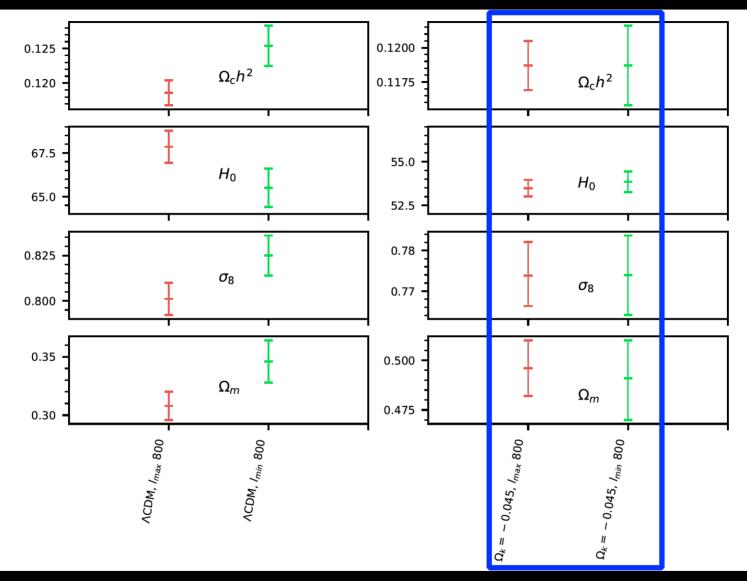
# **Curvature can explain internal tension**



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

In a closed Universe with  $\Omega K = -0.045$ , the cosmological parameters derived in the two different multipole ranges are now fully compatible.

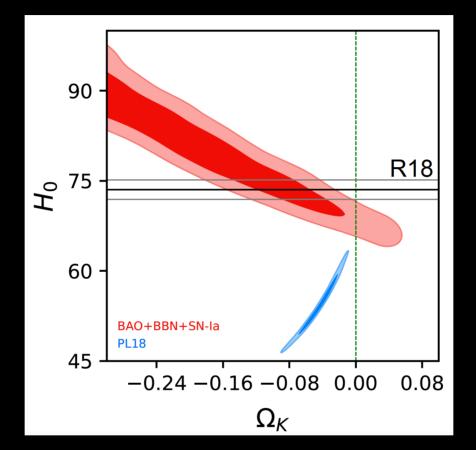
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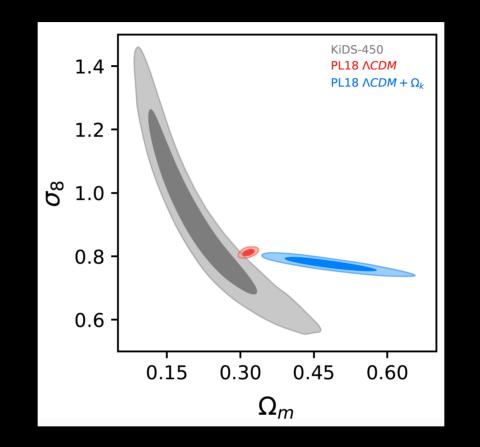
#### Curvature can't explain external tensions



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

Varying  $\Omega \kappa$ , both the well known tensions on H0 and S8 are exacerbated. In a  $\Lambda CDM + \Omega K$  model, Planck gives H0 = 54.4<sup>+3.3</sup>-4.0 km/s/Mpc at 68% cl., increasing the tension with SH0ES at 5.5 $\sigma$ .

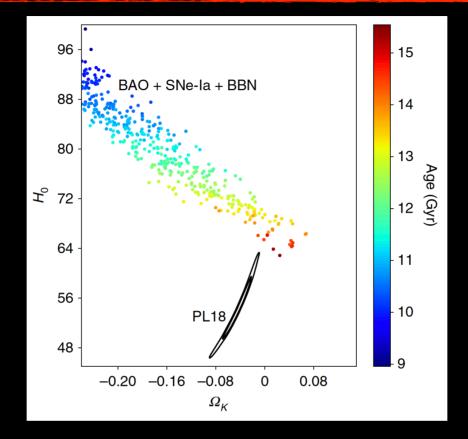
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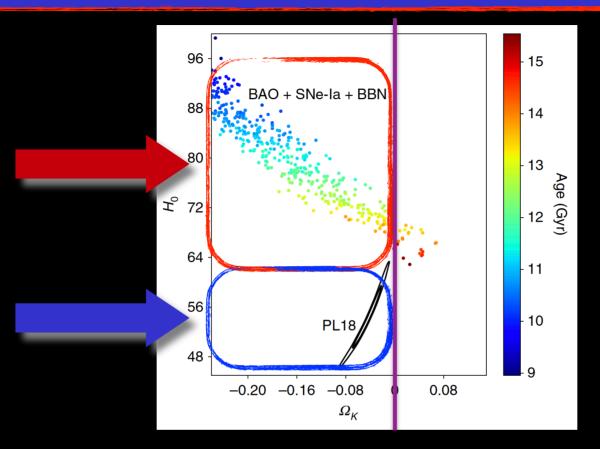
#### What about non-CMB data?



Di Valentino, Melchiorri and Silk, Nature Astron. 4 (2019) 2, 196-203

It is now interesting to address the compatibility of Planck with combined datasets, like BAO + type-la supernovae + big bang nucleosynthesis data. In principle, each dataset prefers a closed universe, but BAO+SN-Ia+BBN gives H0 = 79.6  $\pm$  6.8 km/s/Mpc at 68%cl, perfectly consistent with SH0ES, but at 3.4 $\sigma$  tension with Planck.

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BAO+SNIa+BBN+R18 gives sik = -0.091 ± 0.037 at 68%cl.

#### EFTOFLSS to investigate FS data

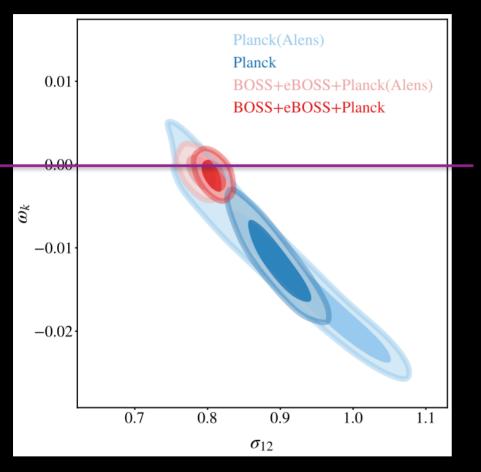
	$\ln(10^{10}A_{s})$	s)	h		$\Omega_{cdm}h^2$		$\Omega_m$		$\Omega_k$	n <sub>s</sub>		$2 * \log(\mathcal{L})$
Flat, fixed $n_s$	$2.85^{+0.11}_{-0.12}$	(3.03)	$0.667^{+0.011}_{-0.011}$	(0.672)	$0.114^{+0.005}_{-0.004}$	(0.115)	$0.307\substack{+0.010\\-0.011}$	(0.304)	-	-		367.2
Curved, fixed $n_s$	$2.55^{+0.21}_{-0.22}$	(2.77)	$0.686^{+0.015}_{-0.016}$	(0.665)	$0.115^{+0.004}_{-0.005}$	(0.111)	$0.291\substack{+0.014\\-0.013}$	(0.302	$-0.089^{+0.049}_{-0.046}$ $-0.042$	) -		366.3
Flat, varying $n_s$	$2.80^{+0.14}_{-0.13}$	(2.97)	$0.669^{+0.012}_{-0.011}$	(0.668)	$0.117^{+0.009}_{-0.008}$	(0.114)	$0.312^{+0.017}_{-0.014}$	(0.304)	-	$0.950\substack{+0.04\\-0.051}$	(0.972)	367.1
Curved, varying <i>n<sub>s</sub></i>	$2.19^{+0.29}_{-0.28}$	(2.62)	$0.707\substack{+0.021\\-0.021}$	(0.686)	$0.127^{+0.011}_{-0.009}$	(0.116)	$0.300\substack{+0.016\\-0.014}$	(0.295	$-0.152^{+0.059}_{-0.053}(-0.089)$	$0.878^{+0.053}_{-0.055}$	(0.932)	364.8

Glanville et al., *arXiv:2205.05892* 

0,10

In this paper they use EFTofLSS to simultaneously constrain measurements from the 6dFGS, BOSS, and eBOSS catalogues, in order to remove some of the assumptions of flatness that enter into other large-scale structure analyses. Fitting the FS data with a BBN prior they measure a >2σ preference for a closed universe.

#### Beyond six parameters: extending $\Lambda CDM + \Omega k$



Semenaite et al., arXiv:2210.07304

A similar result has been obtained by analysing a wKCDM model, and the parameter  $\omega_{K=}\Omega_kh^2$  that gives

 $\omega_{\rm K} = -0.0116^{+0.0029}_{-0.0036}$ 

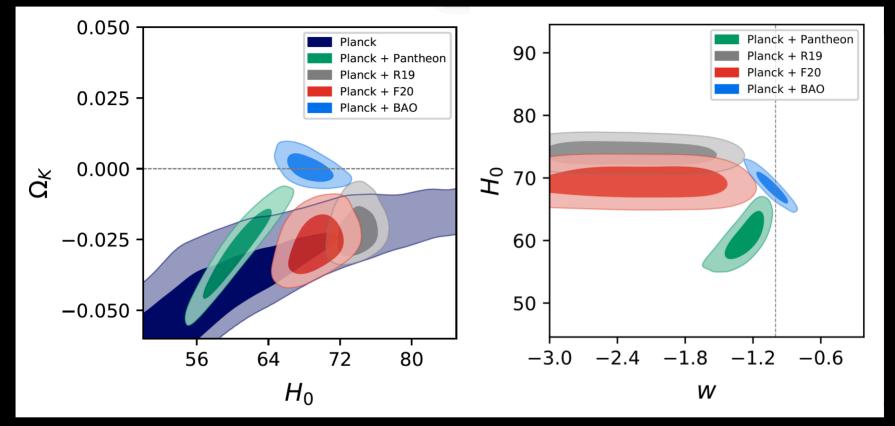
i.e. a  $4\sigma$  preference for a closed universe.

# Constraints at 68% cl. Beyond six parameters: extending $\Lambda CDM + \Omega k$

Parameters	Planck	Planck	Planck	Planck	Planck
1 af affieters	I TAILCK				
		+R19	+F20	+BAO	+ Pantheon
$\Omega_b h^2$	$0.02253 \pm 0.00019$	$0.02253^{+0.00020}_{-0.00016}$	$0.02255^{+0.00019}_{-0.00017}$	$0.02243 \pm 0.00016$	$0.02255 \pm 0.00018$
$\Omega_c h^2$	$0.1183 \pm 0.0016$	$0.1187\substack{+0.0015\\-0.0018}$	$0.1184 \pm 0.0015$	$0.1198 \pm 0.0014$	$0.1186 \pm 0.0015$
$100 heta_{ m MC}$	$1.04099 \pm 0.00035$	$1.04103^{+0.00034}_{-0.00031}$	$1.04105 \pm 0.00034$	$1.04095 \pm 0.00032$	$1.04107 \pm 0.00034$
au	$0.0473 \pm 0.0083$	$0.052^{+0.009}_{-0.011}$	$0.0491 \pm 0.0079$	$0.0563 \pm 0.0081$	$0.0506 \pm 0.0082$
$\Sigma m_{\nu}$ [eV]	$0.43^{+0.16}_{-0.27}$	< 0.513	$0.28^{+0.11}_{-0.22}$	< 0.194	< 0.420
w	$-1.6^{+1.0}_{-0.8}$	$-2.11^{+0.35}_{-0.77}$	$-2.14\pm0.46$	$-1.038\substack{+0.098\\-0.088}$	$-1.27^{+0.14}_{-0.09}$
$\Omega_k$	$\begin{array}{c} -1.6^{+1.0}_{-0.8}\\ -0.074^{+0.058}_{-0.025}\end{array}$	$-0.0192\substack{+0.0036\\-0.0099}$	$-0.0263\substack{+0.0060\\-0.0077}$	$0.0003\substack{+0.0027\\-0.0037}$	$-0.029^{+0.011}_{-0.010}$
$\ln(10^{10}A_s)$	$3.025\pm0.018$	$3.037\substack{+0.016\\-0.026}$	$3.030\pm0.017$	$3.049 \pm 0.017$	$3.034 \pm 0.017$
$n_s$	$0.9689 \pm 0.0054$	$0.9686\substack{+0.0056\\-0.0050}$	$0.9693 \pm 0.0051$	$0.9648 \pm 0.0048$	$0.9685 \pm 0.0051$
$lpha_S$	$-0.0005 \pm 0.0067$	$-0.0012 \pm 0.0066$	$-0.0010 \pm 0.0068$	$-0.0054 \pm 0.0068$	$-0.0023 \pm 0.0065$
$H_0 [{ m km/s/Mpc}]$	$53^{+6}_{-16}$	$73.8 \pm 1.4$	$69.3\pm2.0$	$68.6^{+1.5}_{-1.8}$	$60.5 \pm 2.5$
$\sigma_8$	$0.74^{+0.08}_{-0.16}$	$0.932\pm0.040$	$0.900\pm0.039$	$0.821 \pm 0.027$	$0.812\substack{+0.031\\-0.018}$
$S_8$	$0.989^{+0.095}_{-0.063}$	$0.874 \pm 0.032$	$0.900\substack{+0.034\\-0.031}$	$0.826 \pm 0.016$	$0.927 \pm 0.037$
$Age[{ m Gyr}]$	$16.10^{+0.92}_{-0.80}$	$14.90\substack{+0.72 \\ -0.32}$	$15.22^{+0.054}_{-0.038}$	$13.77\pm0.10$	$14.98 \pm 0.39$
$\Omega_m$	$0.61\substack{+0.21\-0.34}$	$0.264_{-0.013}^{+0.010}$	$0.300\substack{+0.017\\-0.020}$	$0.305\pm0.016$	$0.393\substack{+0.030\\-0.036}$
$\Delta \chi^2_{bestfit}$	0.0	0.62	0.88	14.77	1037.82

We want to check the robustness of these results further increasing the number of parameters, in addition to curvature.

#### **Cosmic Discordance**



#### Evidence for a phantom closed Universe at more than 99% CL!!

It is interesting to note that if a closed universe increases the fine-tuning of the theory, the removal of a cosmological constant reduces it. It is, therefore, difficult to decide whether a phantom closed model is less or more theoretically convoluted than ACDM.

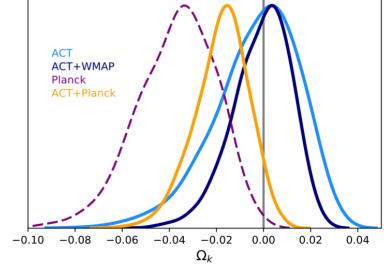
Di Valentino, Melchiorri and Silk, ApJ Letters, 908, L9 (2021), arXiv:2003.04935

#### What about different CMB experiments?



ACT-DR4 + WMAP gives at 68% CL

 $\Omega_k = -0.001 \pm 0.012$ 



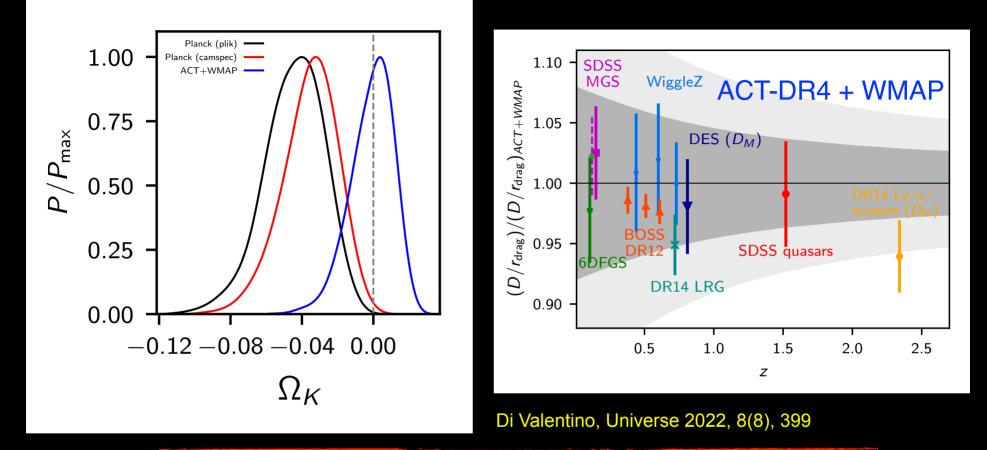
ACT-DR4 2020, Aiola et al., arXiv:2007.07288 [astro-ph.CO]

#### What about different CMB experiments?

CMB Polarization Measurements with SPTpol SPT-3G H<sub>0</sub> [km s<sup>-1</sup> Mpc<sup>-1</sup>] 109 00 108 100Nicholas Harrington Planck UC Berkeley SPT-3G + PlanckSPT-3G + Planck + BAORiess et al. 2020 SPT-3G gives at 68% CL:  $\Omega_K = 0.001^{+0.018}_{-0.019}$ -0.10-0.050.00  $\Omega_K$ 

SPT-3G, arXiv:2103.13618 [astro-ph.CO]

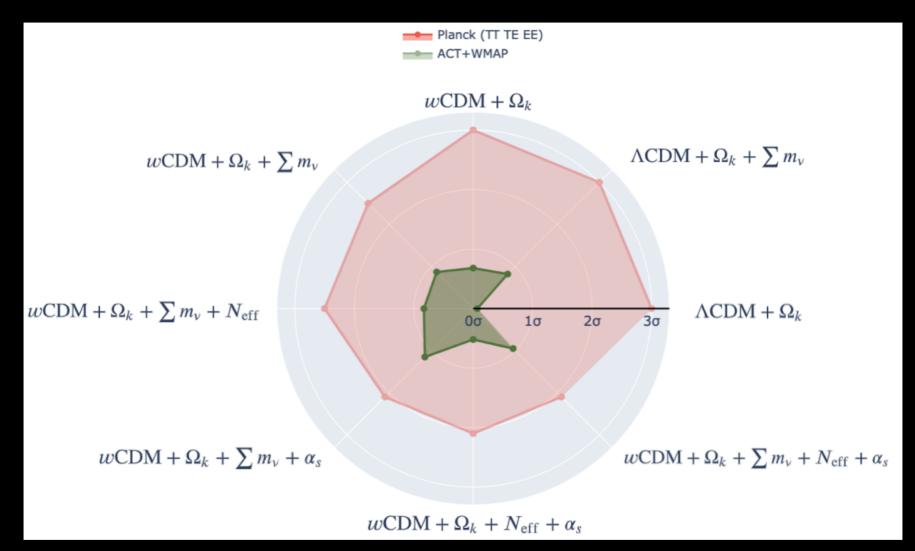
### ACT-DR4



Confirmation that from a CMB experiment you can obtain Omegak!

When precise CMB measurements at arc-minute angular scales are included, since gravitational lensing depends on the matter density, its detection breaks the geometrical degeneracy.

### Tension with $\Omega_k = 0$



Di Valentino et al., arXiv:2209.12872

And the indication we see in the simplest LCDM+  $\Omega_k$  model is robust also in its extensions.

### Inflation: $\Omega_k < 0$ or HZ?

Dataset	Scalar Spectral Index $(n_s)$ $\Lambda \text{CDM}$				
ACT	$1.009\pm0.015$				
ACT+BAO (DR12)	$1.006\pm0.013$				
ACT+BAO (DR16)	$1.006\pm0.014$				
ACT+DESy1	$1.007\pm0.013$				
ACT+SPT+BAO (DR12)	$0.996 \pm 0.012$				
Planck	$0.9649 \pm 0.0044$				
Planck+BAO (DR12)	$0.9668 \pm 0.0038$				
Planck+BAO (DR16)	$0.9677 \pm 0.0037$				
Planck $(2 \le \ell \le 650)$	$0.9655 \pm 0.0043$				
Planck ( $\ell > 650$ )	$0.9634 \pm 0.0085$				

At this point, if Planck seems to disfavour the inflationary prediction for a flat background geometry at more than 3σ, ACT, although in perfect agreement with spatial flatness, shows a preference for a larger spectral index consistent with a Harrison-Zel'dovich scale-invariant spectrum ns=1 of primordial density perturbations introducing a tension with a significance of 2.7σ with the results from the Planck satellite.

Giarè, Renzi, Mena, Di Valentino, and Melchiorri, arXiv:2210.09018

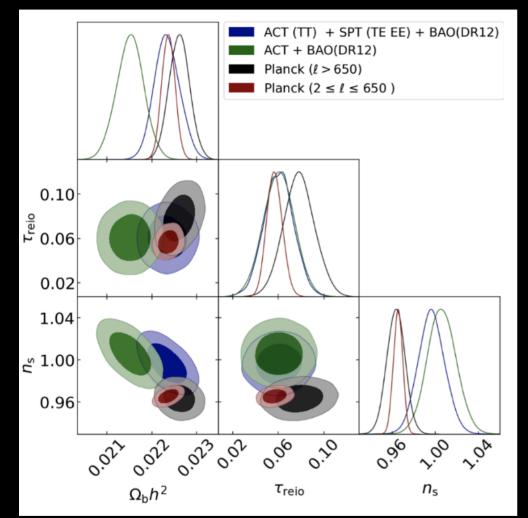
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Giarè, Renzi, Mena, Di Valentino, and Melchiorri, arXiv:2210.09018

n ACT-DR4 2020, arXiv:2007.07288 [astro-ph.CO] this discrepancy was interpreted as a consequence of the lack of information concerning the first acoustic peak of the temperature power spectrum. To verify this origin of the discrepancy in the CMB values of ns, we have performed two separate analyses of the Planck observations, splitting the likelihood into low 2 < I < 650 and high I > 650 multipoles. We find that the discrepancy still persists at the level of  $3\sigma$  ( $2\sigma$ ) for low (high) multiple temperature data. Planck data still prefer a value of the scalar spectral index smaller than unity at  $\sim 4.3\sigma$ when the information about the first acoustic peak is removed.

### Inflation: $\Omega_k < 0$ or HZ?

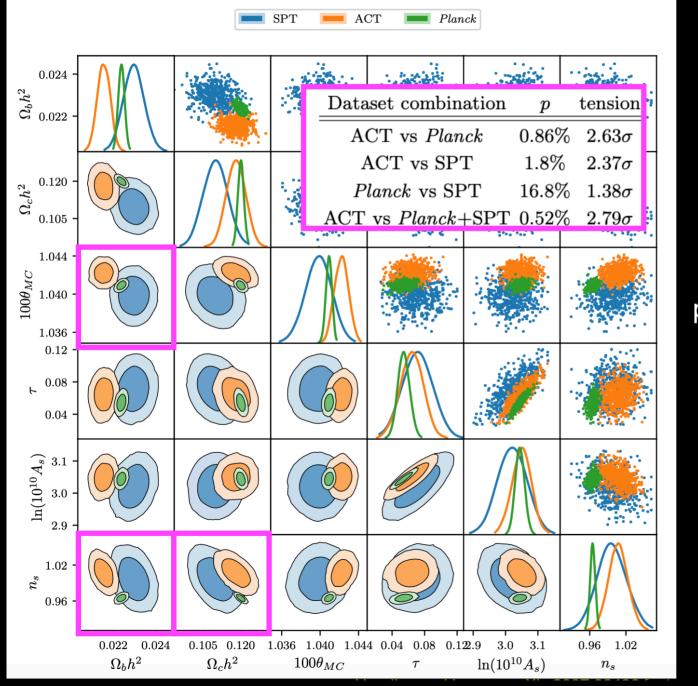


Giarè, Renzi, Mena, Di Valentino, and Melchiorri, arXiv:2210.09018

Such preference remains robust under the addition of large scale structure information, and in the two-dimensional plane it can be definitely noted that the direction of the  $\Omega_bh^2$  - ns degeneracy is opposite for ACT and Planck, and the disagreement here is significantly exceeding 3 $\sigma$ . This tension is partially driven by the ACT polarization data, as we can see replacing it with the SPT polarization measurements, but while the tension is relaxed in the plane  $\Omega_bh^2$  - ns, this combination is still preferring ns=1.

#### Alternative CMB vs Planck: LCDM

Handley and Lemos, arXiv:2007.08496 [astro-ph.CO]



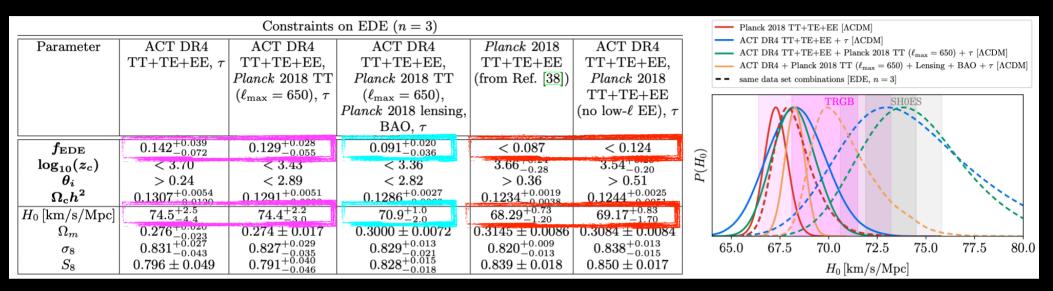
Global tensions between CMB datasets.

For each pairing of datasets this is the tension probability p that such datasets would be this discordant by (Bayesian) chance, as well as a conversion into a Gaussianequivalent tension.

Between Planck and ACT there is a 2.6σ tension.

Assuming LCDM

### ACT-DR4 vs Planck: EDE



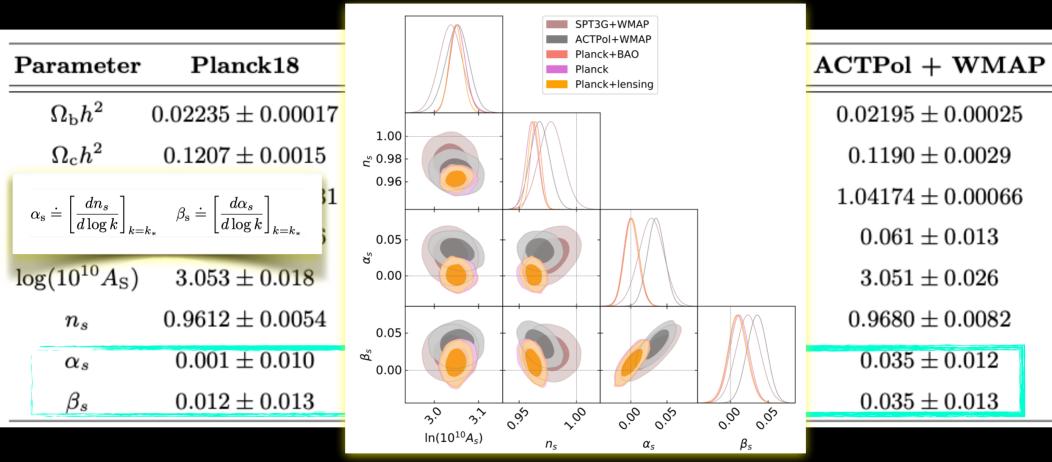
ACT collaboration, Hill et al. arXiv:2109.04451

Considering ACT only data or combined with Planck TT up to multipoles 650, there is an evidence for EDE > 3σ, solving completely the Hubble tension. The evidence for EDE > 3σ persists with the inclusion of Planck lensing + BAO data, but shifting H0 towards a lower value. Once the full Planck data are considered, the evidence for EDE disappears and H0 is again in tension with SH0ES.

The Planck damping tail is in disagreement with EDE different from zero.

### ACT-DR4 vs Planck: $\alpha_s$ and $\beta_s$ .

Forconi, Giarè, Di Valentino and Melchiorri, *Phys.Rev.D* 104 (2021) 10, 103528



ACT-DR4 and SPT-3G are in agreement one with each other, but in disagreement with Planck, for the value of the

running of the scalar spectral index  $\alpha_s$  and of the running of the running  $\beta_s$ . In particular ACT-DR4 + WMAP prefer both a non vanishing running  $\alpha_s$  and running of the running  $\beta_s$  at the level of 2.9 $\sigma$  and 2.7 $\sigma$ , respectively.

### Alternative CMB vs Planck: Σmv

Di Valentino and Melchiorri, 2022 ApJL 931 L18

Constraints at 68% CL	$\Sigma m_{\nu}   [{ m eV}]$
${\rm Planck}(+A_{\rm lens})$	< 0.51
$\mathrm{Planck}{+}\mathrm{BAO}~(+A_{\mathrm{lens}})$	< 0.19
$\mathrm{Planck+Pantheon}~(+A_{\mathrm{lens}})$	< 0.25
$\mathrm{Planck+Lensing}\;(+A_{\mathrm{lens}})$	$0.41\substack{+0.17 \\ -0.25}$
ACT-DR4+WMAP	$0.68\pm0.31$
ACT-DR4+WMAP+BAO	< 0.19
ACT-DR4+WMAP+Pantheon	< 0.25
ACT-DR4+WMAP+Lensing	$0.60\pm0.25$
SPT-3G+WMAP	$0.46\substack{+0.14 \\ -0.36}$
$\mathbf{SPT}\operatorname{-3G}+\mathbf{WMAP}+\mathbf{BAO}$	$0.22\substack{+0.056\\-0.14}$
${ m SPT-3G+WMAP+Pantheon}$	$0.25\substack{+0.052\\-0.19}$
SPT-3G+WMAP+Lensing	< 0.37

Moreover, we have a mildly suggestion from both the ACT-DR4 and SPT-3G data, when combined with WMAP, of a neutrino mass with  $\Sigma m_V = 0.68 \pm 0.31 \text{ eV}$  and  $\Sigma m_V = 0.46^{+0.14} - 0.36 \text{ eV}$  at 68% CL, respectively. A combination of Planck CMB+Lensing constrain  $\Sigma m_V = 0.41^{+0.17} - 0.25 \text{ eV}$  at 68% CL when a variation in the AL parameter is considered.

# Quantifying global CMB tension

Cosmological model	d	$\chi^2$	p	$\log S$	Tension
ΛCDM	6	16.3	0.012	-5.17	$2.51\sigma$
$\Lambda { m CDM} + A_{ m lens}$	7	18.5	0.00977	-5.77	$2.58\sigma$
$\Lambda { m CDM} + N_{ m eff}$	7	13	0.0719	-3	$1.80\sigma$
$\Lambda  ext{CDM} + \Omega_k$	7	16.5	0.0209	-4.75	$2.31\sigma$
$w\mathrm{CDM}$	7	16.8	0.0187	-4.9	$2.35\sigma$
$\Lambda  ext{CDM} + \sum m_{ u}$	7	20.7	0.00421	-6.86	$2.86\sigma$
$\Lambda  ext{CDM} + lpha_s$	7	20.6	0.00448	-6.78	$2.84\sigma$
$w \text{CDM} + \Omega_k$	8	17.6	0.0249	-4.78	$2.24\sigma$
$\Lambda \text{CDM} + \Omega_k + \sum m_{\nu}$	8	21.2	0.00651	-6.62	$2.72\sigma$
$w \text{CDM} + \Omega_k + \sum m_{ii}$	9	19.8	0.0195	-5.38	$2.34\sigma$
$w$ CDM + $\Omega_k$ + $\sum m_{\nu}$ + $N_{\text{eff}}$	10	18.8	0.0434	-4.38	$2.02\sigma$
$w  ext{CDM} + \Omega_k + \sum m_ u + lpha_s$	10	22	0.015	-6.01	$2.43\sigma$
$w  ext{CDM} + \Omega_k + N_{ ext{eff}} + \alpha_s$	10	20.9	0.0218	-5.45	$2.29\sigma$
$w  ext{CDM} + \sum m_{ u} + N_{ ext{eff}} + lpha_s$	10	31.1	0.000575	-10.5	$3.44\sigma$
$w$ CDM + $\Omega_k$ + $\sum m_{\nu}$ + $N_{\text{eff}}$ + $\alpha_s$	11	24.7	0.0102	-6.83	$2.57\sigma$

Di Valentino et al., arXiv:2209.14054

$\Lambda { m CDM} + N_{ m eff}$	Planck	_	$2.92\pm0.19$
	ACT-DR4	_	$2.35\substack{+0.40 \\ -0.47}$

If we now study the global agreement between Planck and ACT in various cosmological models that differ by the inclusion of different combinations of additional parameters, we can use the Suspiciousness statistic, to quantify their global "CMB tension".

We find that the 2.5 $\sigma$  tension within the baseline  $\Lambda$ CDM, is reduced at the level of 1.8 $\sigma$ when Neff is significantly less than 3.044, while it ranges between 2.3 $\sigma$ and 3.5 $\sigma$  in all the other extended models.

## Concluding

Most of the anomalies and tensions are involving the CMB data:

- H0 tension
- S8 tension
- $A_L > 1$  or  $\Omega \kappa < 0$  for Planck
- $\alpha_s$ ,  $\beta_s$  or  $\Sigma m_V$  for ACT and SPT
- EDE for ACT

presenting a serious limitation to the precision cosmology.

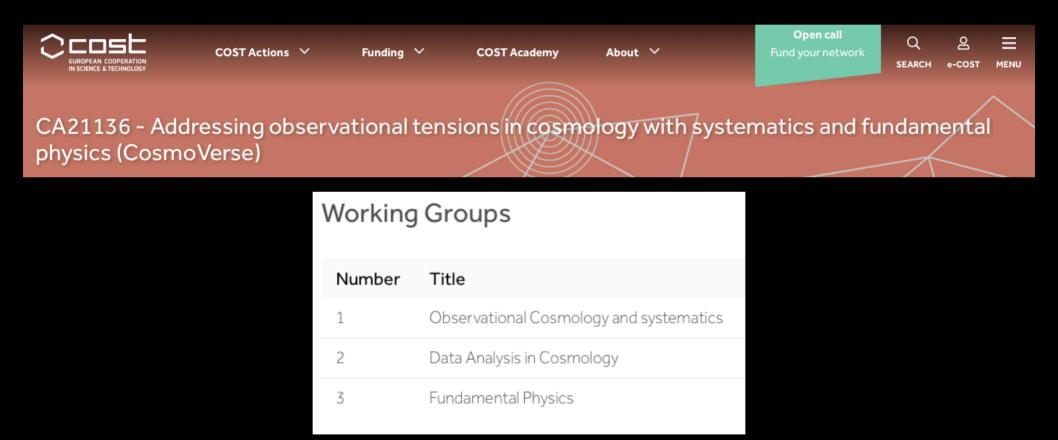
Are we sure that the CMB results are still a confirmation of the flat standard ACDM cosmological model?

At this point, given the quality of all the analyses, probably these discrepancies are indicating a problem with the underlying cosmology and our understanding of the Universe, rather than the presence of systematic effects.

These cosmic discordances

call for new observations and stimulate the investigation of alternative theoretical models and solutions.

# Thank you! e.divalentino@sheffield.ac.uk



#### Action keywords

Cosmological surveys - Observation systematics - Gravitation - Fundamental physics - Dark Energy, Dark Matter