



MSCA

Marie Skłodowska-Curie Actions

*Developing talents,
advancing research*



Coherent and Incoherent Structures in Fuzzy Dark Matter Halos

Gary (I-Kang) Liu

Gerasimos Rigopoulos & Nick Proukakis



G. Rigopoulos

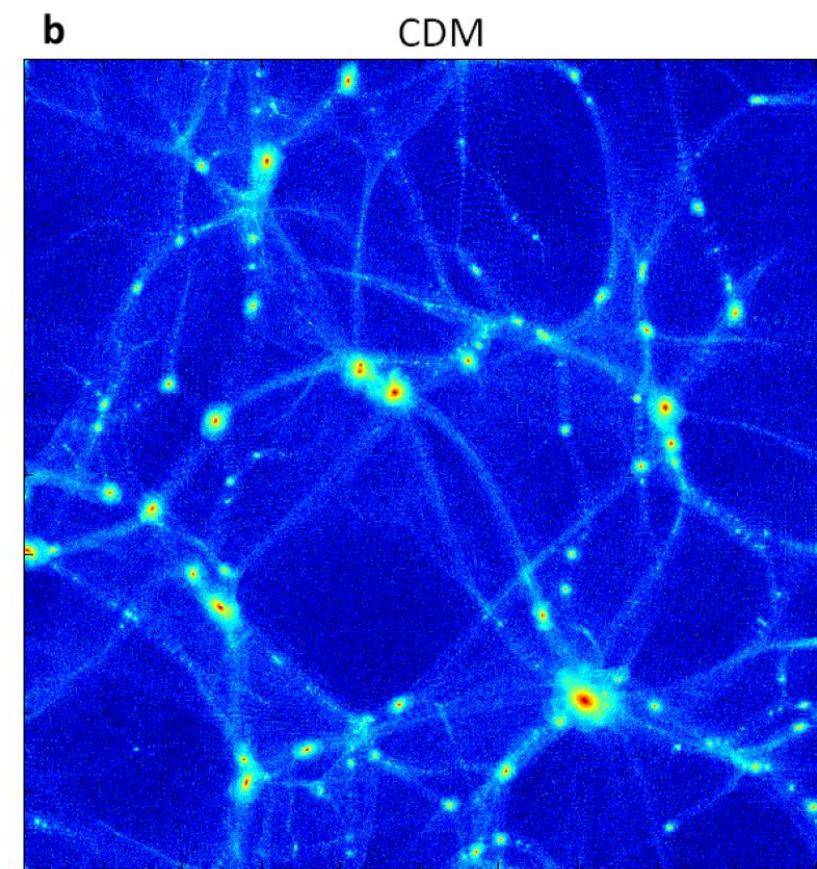
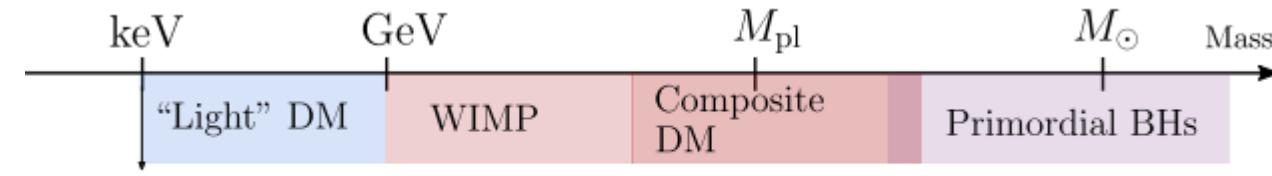


N. Proukakis



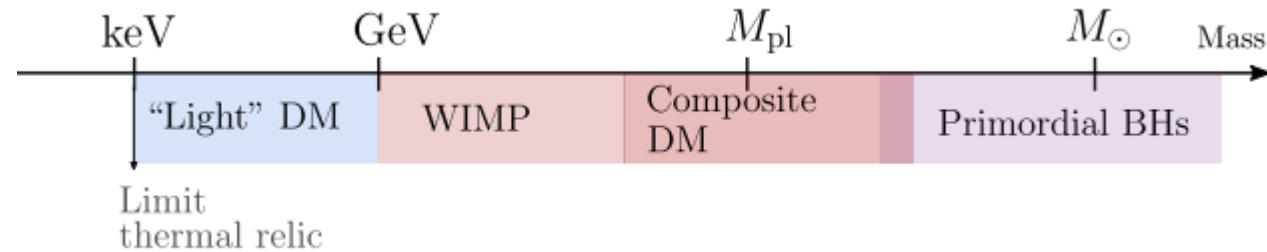
M. Indjin

Fuzzy Dark Matter



Schive et al., Nat.Phys. 10, 496 (2014)

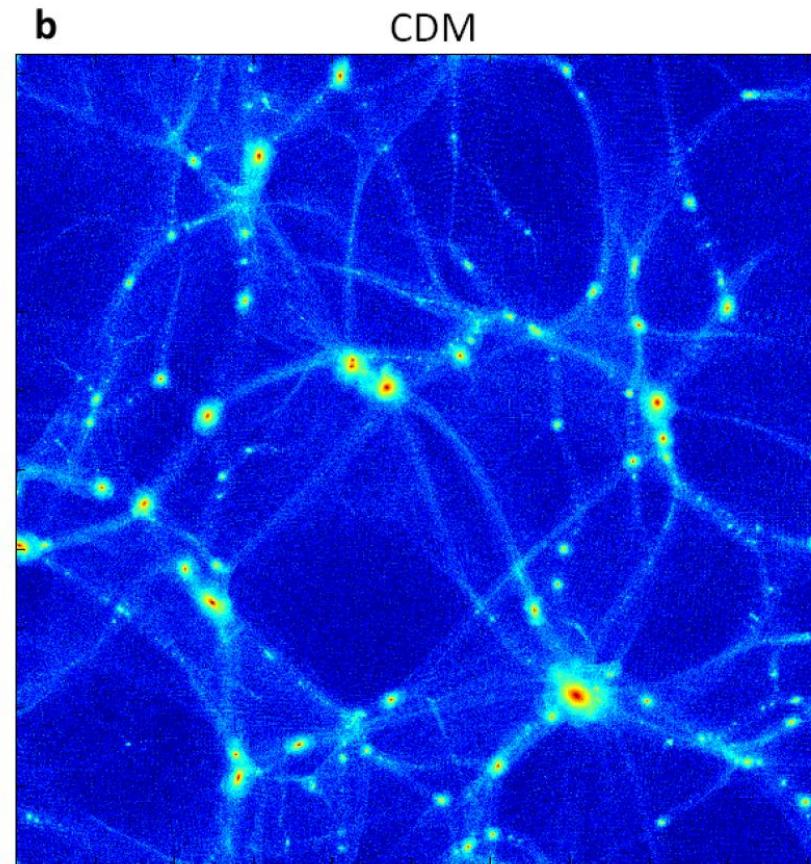
Fuzzy Dark Matter



Small scale crises

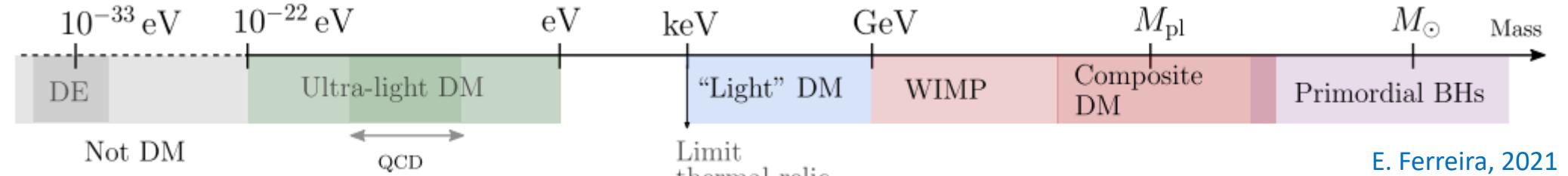
- Missing satellites problem
 - Too-big-to-fail problem
 - Cusp-core problem

$$\rho_{\text{NFW}}(r) = \rho_h \left(\frac{r}{r_h}\right)^{-1} \left(1 + \frac{r}{r_h}\right)^{-2}$$



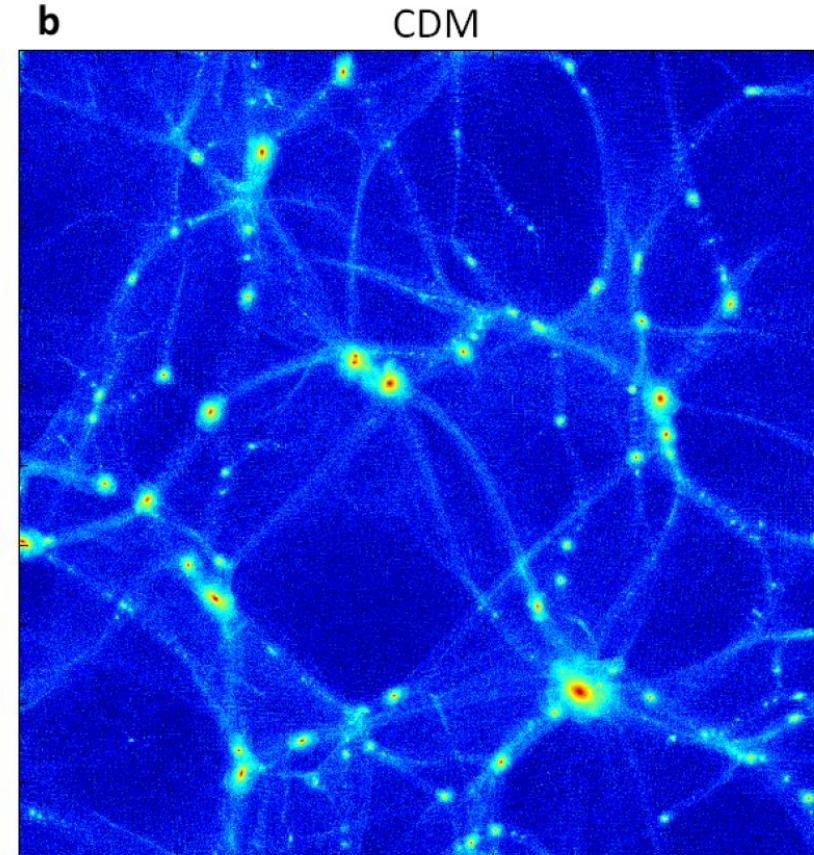
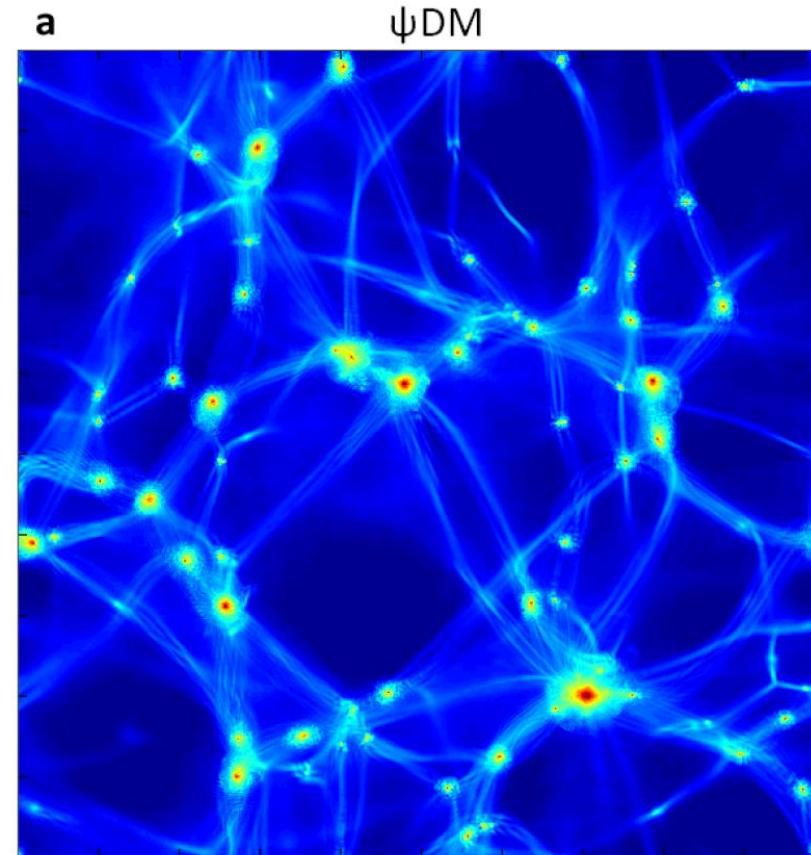
Schive et al., Nat.Phys. 10, 496 (2014)

Fuzzy Dark Matter



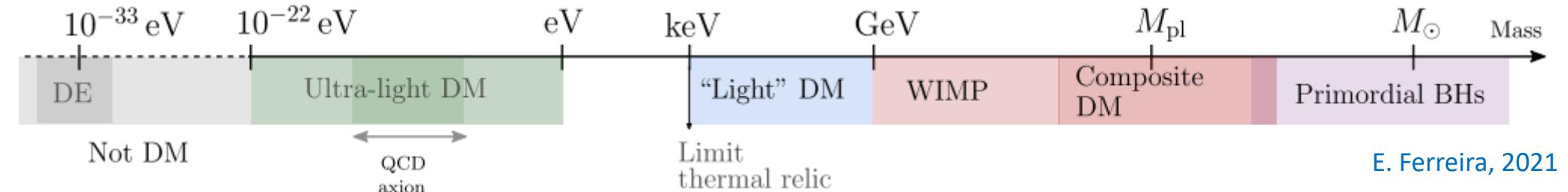
E. Ferreira, 2021

$$\lambda_{dB} = \frac{\hbar}{mv}$$

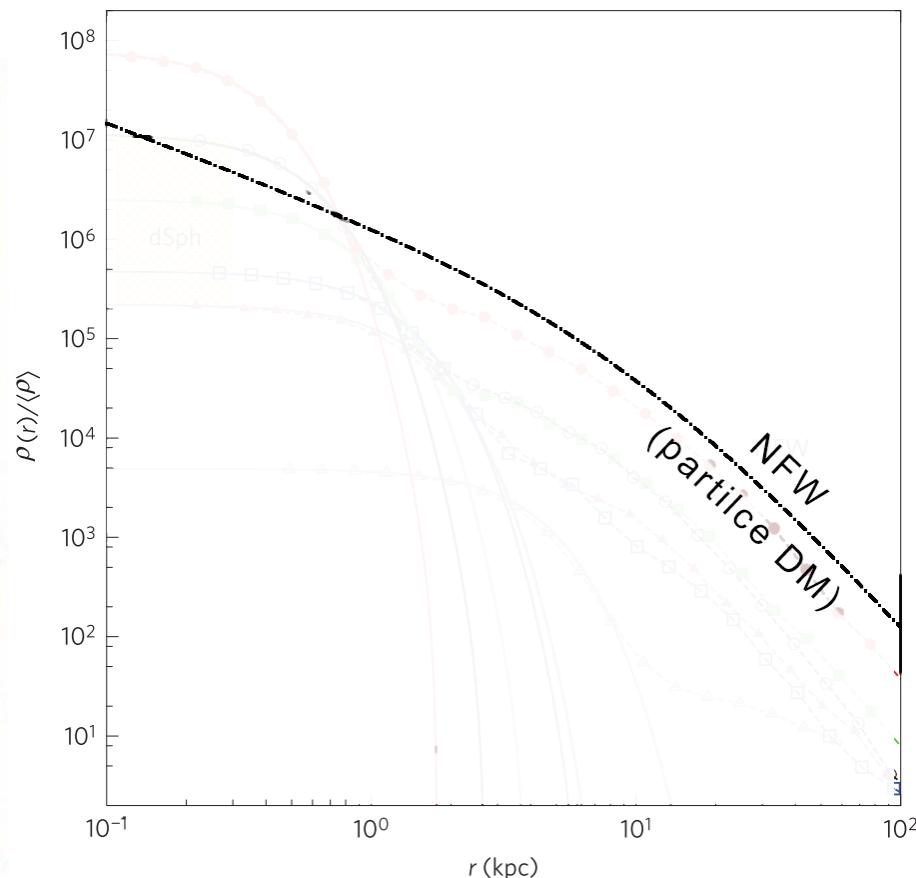
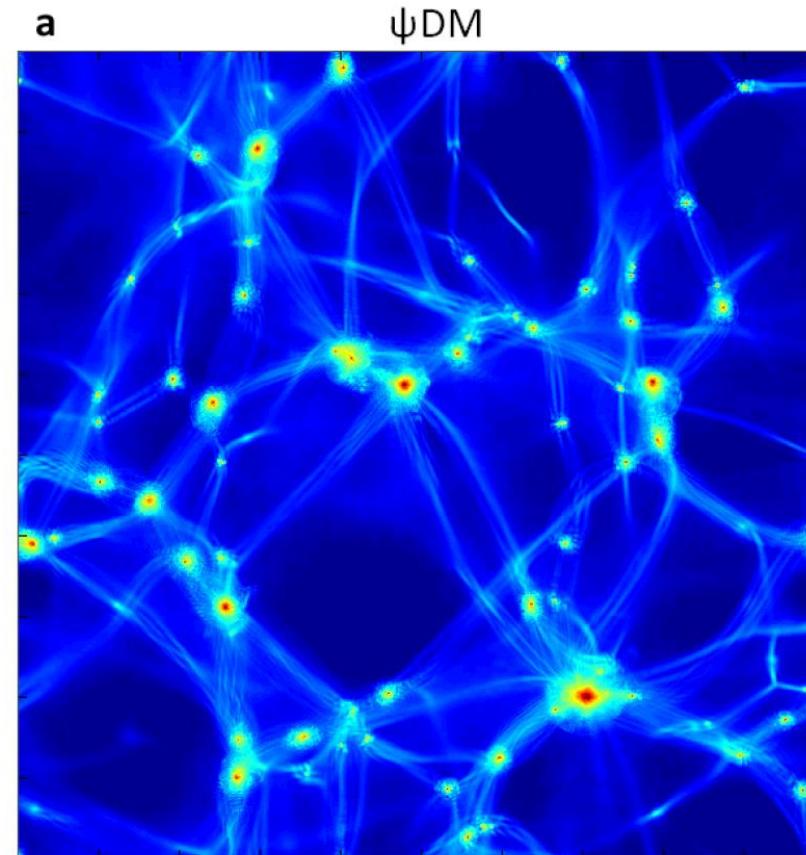


Schive et al., Nat.Phys. 10, 496 (2014)

Fuzzy Dark Matter

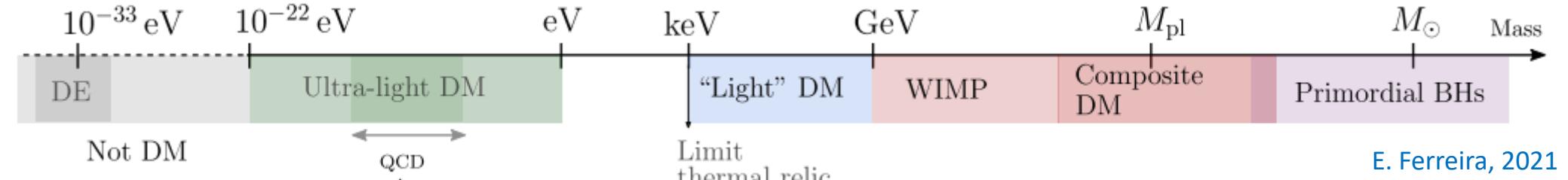


$$\lambda_{dB} = \frac{h}{mv}$$

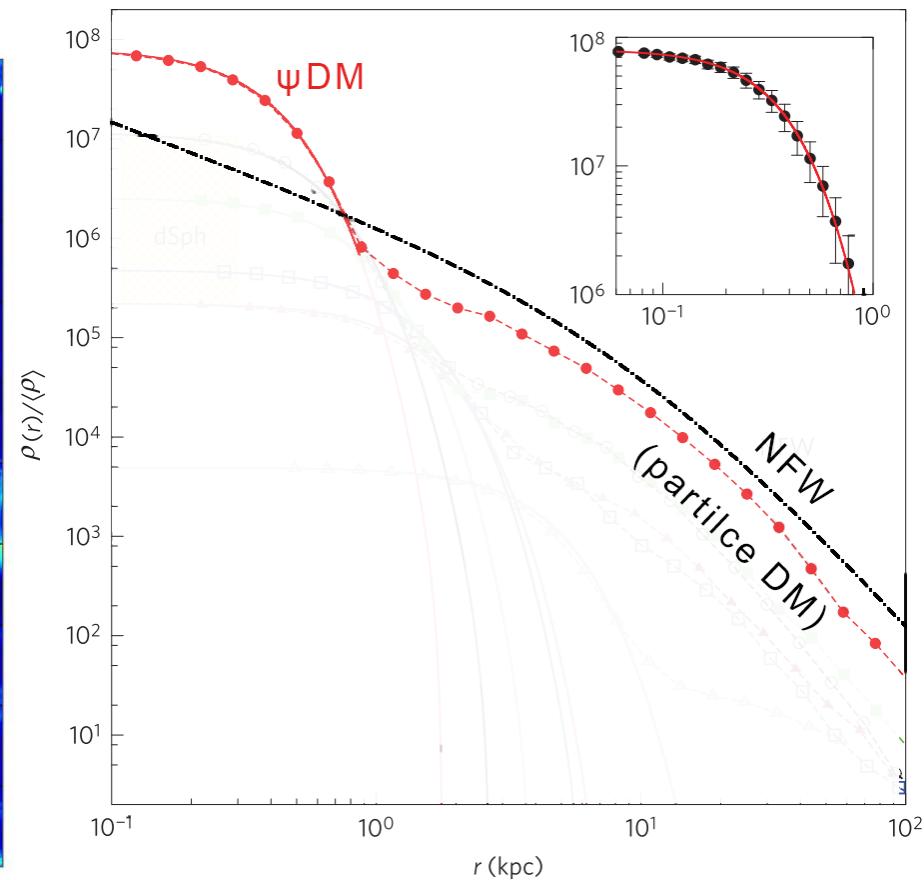
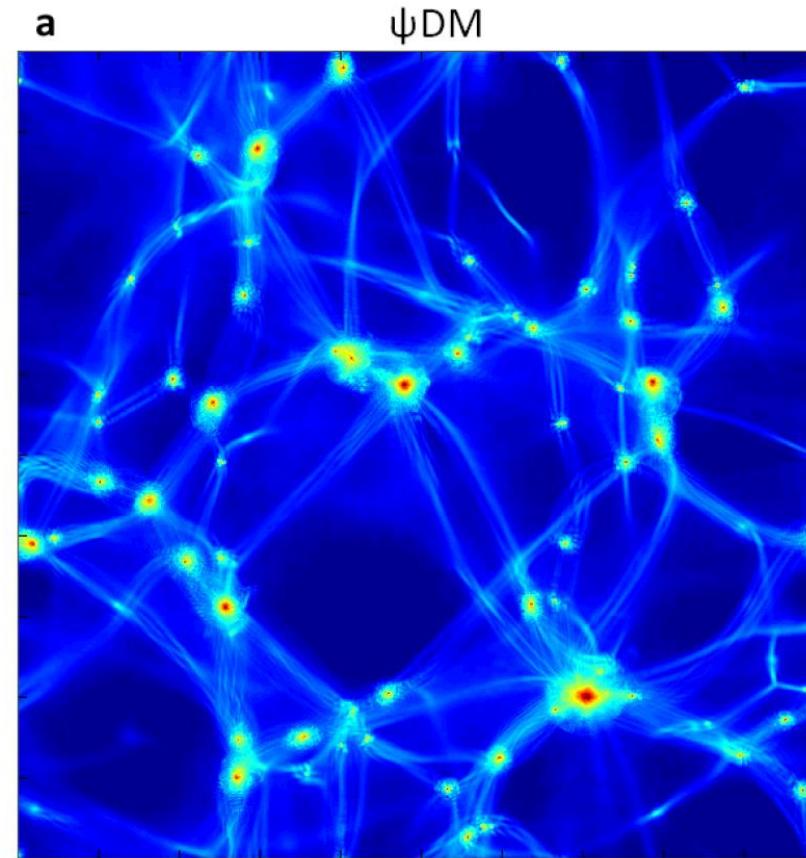


Schive et al., Nat.Phys. 10, 496 (2014)
 Navarro, Frenk & White, ApJ, 462, 563 (1996)

Fuzzy Dark Matter

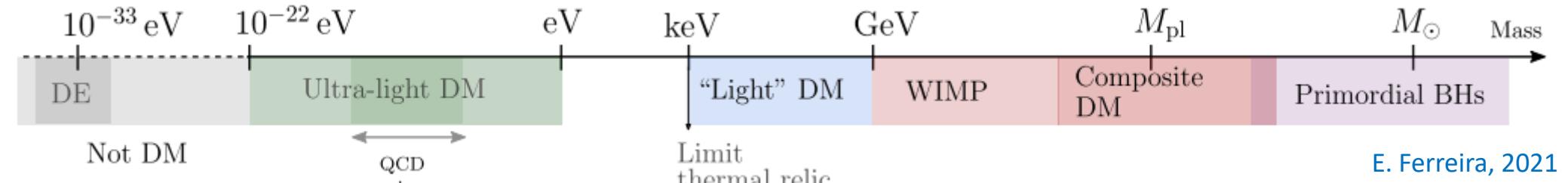


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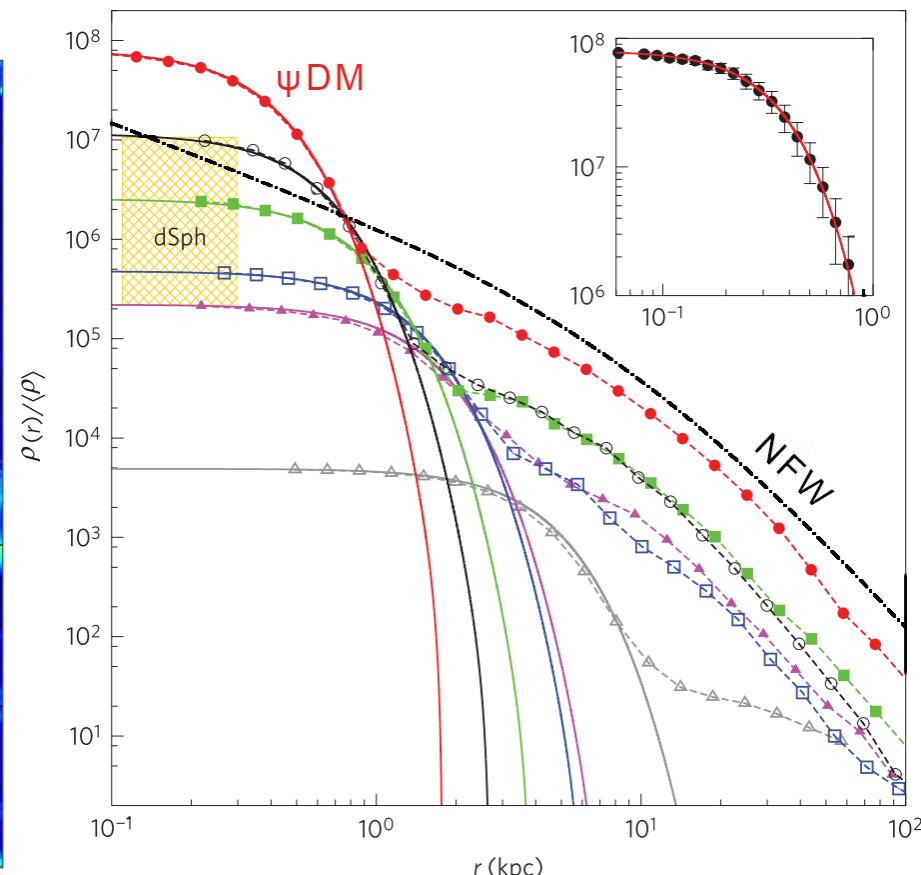
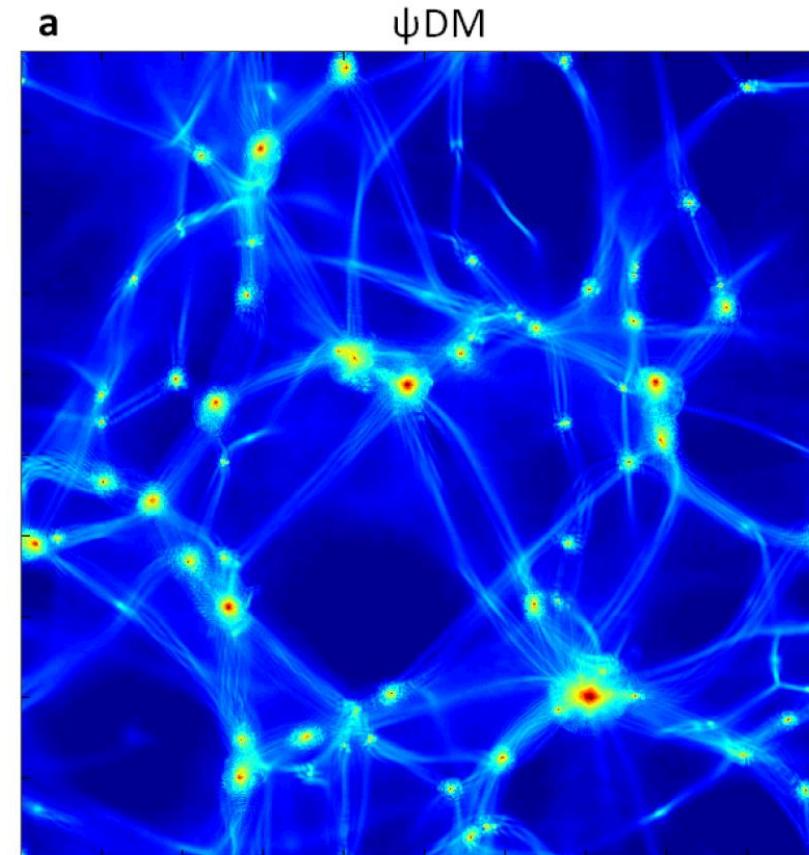


Schive et al., Nat.Phys. 10, 496 (2014)

Fuzzy Dark Matter



$$\lambda_{dB} = \frac{\hbar}{mv}$$



Schive et al., Nat.Phys. 10, 496 (2014)

PRL 113, 261302 (2014)

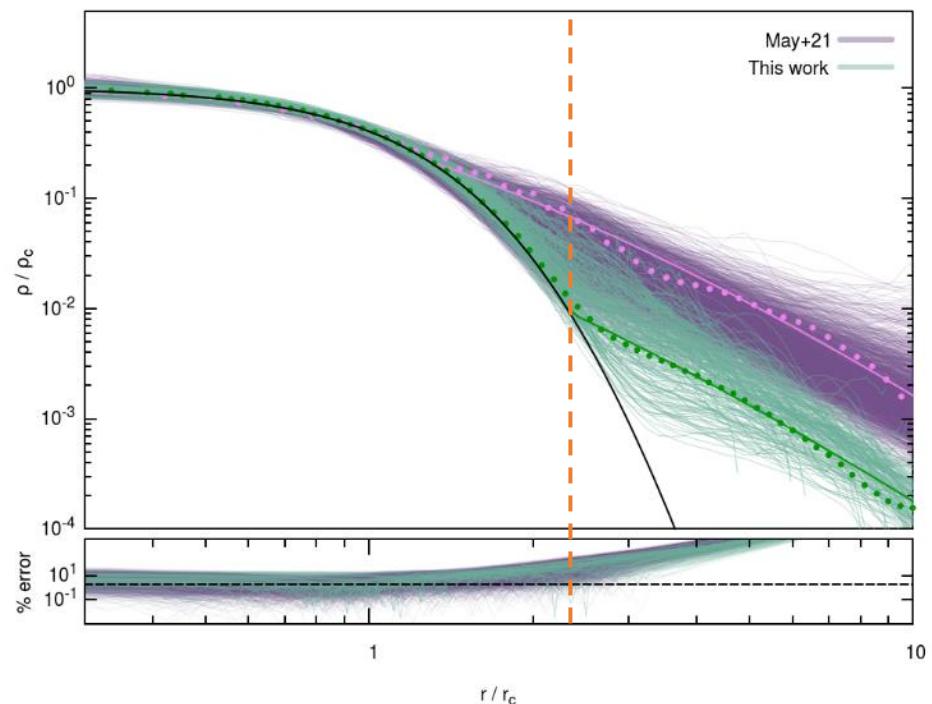
Halo structures

- Cored-halo profile (Bimodal fit)

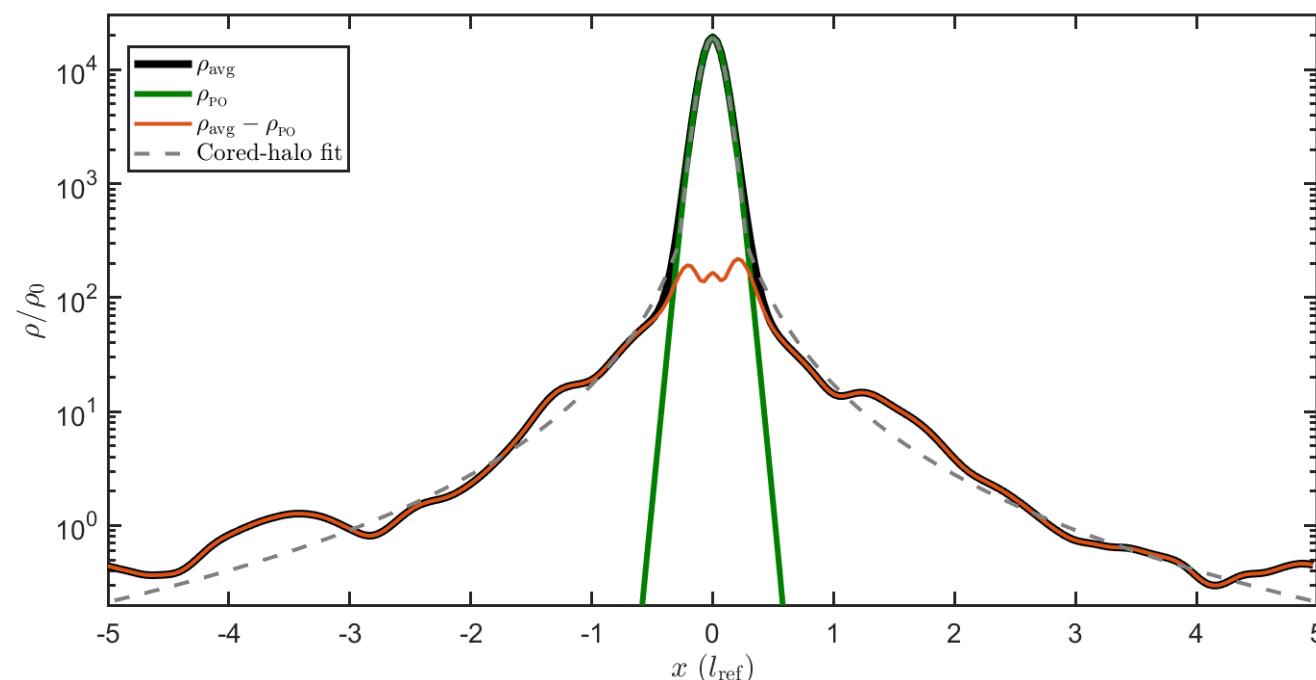
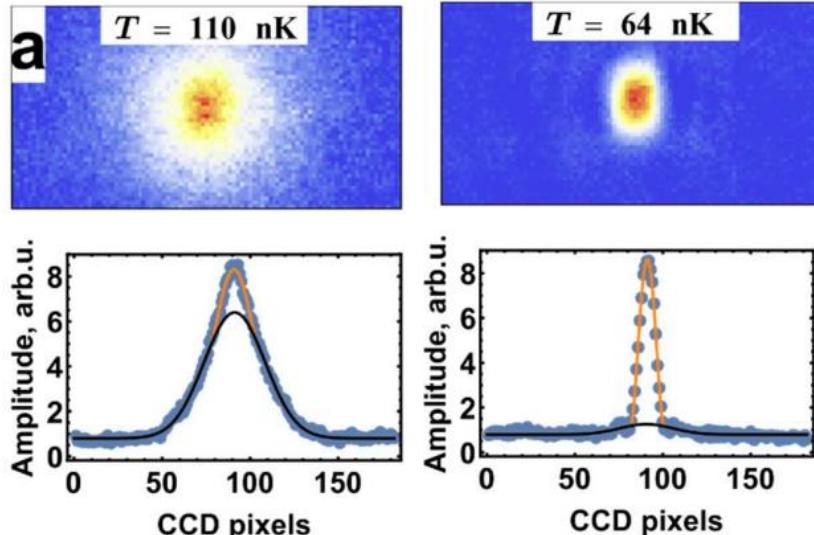
$$\rho_{\text{cNFW}}(r) = \begin{cases} \rho_{\text{soliton}}(r) & , r \leq r_t \\ \rho_{\text{NFW}}(r) & , r > r_t \end{cases}$$

$$\rho_{\text{soliton}}(r) = \rho_c \left[1 - \lambda \left(\frac{r}{r_c} \right)^2 \right]^{-8} \quad \rho_{\text{NFW}}(r) = \rho_h \left(\frac{r}{r_h} \right)^{-1} \left(1 + \frac{r}{r_h} \right)^{-2}$$

Chan et al., arXiv:2110.11882



E. T. Davletov et al., PRA 102, 011302(R) (2020)



Fuzzy Dark Matter

- Ultralight mass, $m \approx [10^{-22}, 10]$ eV

L. Hui, arXiv:2101.11735

Dentler et al., arXiv:2111.01199

Schrödinger-Poisson equations (SPE)

$$i\hbar \frac{\partial}{\partial t} \Psi(\mathbf{r}, t) = \left[-\frac{\hbar^2 \nabla^2}{2m} + m\Phi(\mathbf{r}, t) \right] \Psi(\mathbf{r}, t)$$

$$\nabla^2 \Phi(\mathbf{r}, t) = 4\pi G [\rho(\mathbf{r}, t) - \bar{\rho}]$$

$$\Psi(\mathbf{r}, t) = \sqrt{\rho(\mathbf{r}, t)} e^{i\varphi(\mathbf{r}, t)}$$

$$\mathbf{v}(\mathbf{r}, t) = \frac{\hbar}{m} \nabla \varphi(\mathbf{r}, t)$$

Hydrodynamic description

$$\frac{\partial}{\partial t} \rho + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial}{\partial t} \mathbf{v} + \nabla \left[\frac{|\mathbf{v}|^2}{2} + \Phi - \frac{\hbar^2}{2m^2} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}} \right] = 0$$

Scaling invariant/symmetry

$$\{t, x, \Phi, \Psi, \rho\} \rightarrow \{\Lambda^2 t', \Lambda x', \Lambda^{-2} \Phi', \Lambda^{-2} \Psi', \Lambda^{-4} \rho'\}$$

$$\{t, x, \Phi, \Psi, \rho\} \rightarrow \{\alpha^{-1} t', x', \alpha^2 \Phi', \alpha \Psi', \alpha^2 \rho'\}$$

Ji & Sin, PRD, 50, 3655 (1994); Mocz et al., MNRAS 471, 4559 (2017)

Wavy/quantum features

- Wave interferences (fuzzy DM/ψDM/BECDM)
- Quantum pressure

Superfluidity

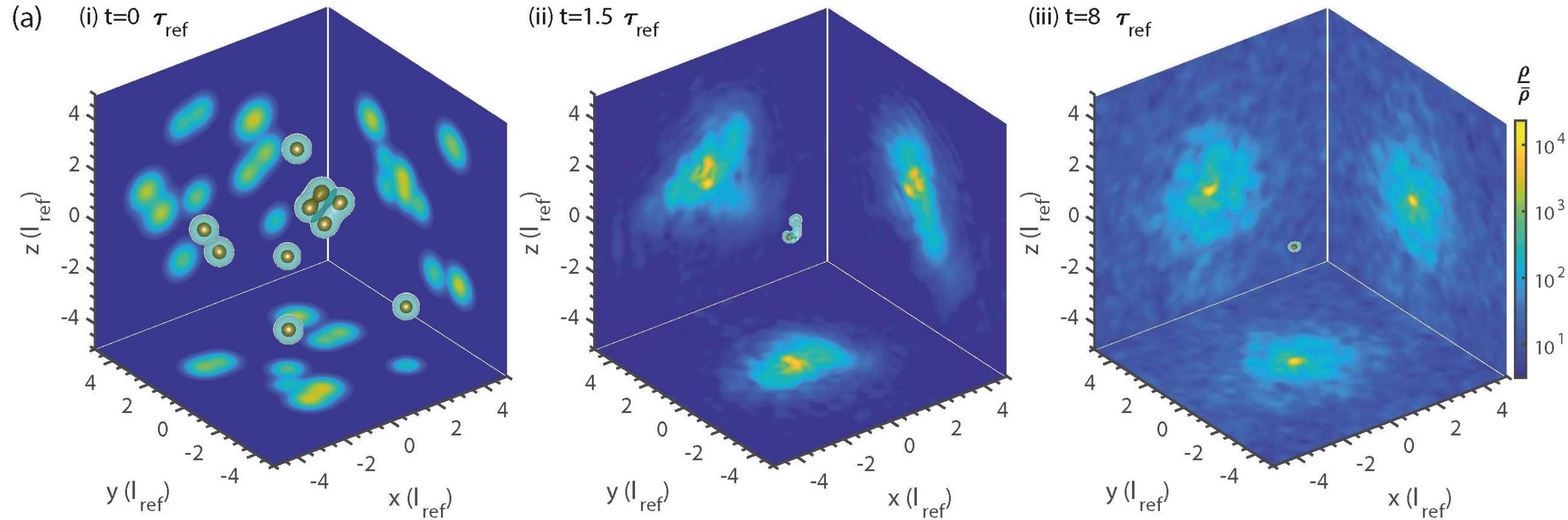
- Viscosity free
- Critical velocity (m -dependent Jeans scale)
- Irrational

Halo structure

Soliton merger simulation

Mocz et al., MNRAS 471, 4559 (2017)
 Chan et al., arXiv:2110.11882

Primary sample ($M = 100M_{\text{ref}}$)



$$E_{\text{ref}} = \hbar \sqrt{G \rho_{\text{ref}}} \quad \tau_{\text{ref}} = (G \rho_{\text{ref}})^{-1/2}$$

$$l_{\text{ref}} = \left(\frac{\hbar^2}{m^2 G \rho_{\text{ref}}} \right)^{1/4}$$

$$\rho_{\text{ref}} = 10^3 M_{\odot} \text{kpc}^{-3}$$

$$m_{\text{ref}} = 2.5 \times 10^{-23} \text{ eV}$$

$$M_{\text{ref}} \approx 1.26 \times 10^6 M_{\odot}$$

$$\tau_{\text{ref}} \approx 14.91 \text{ Gyr}$$

$$l_{\text{ref}} \approx 10.81 \text{ kpc}$$

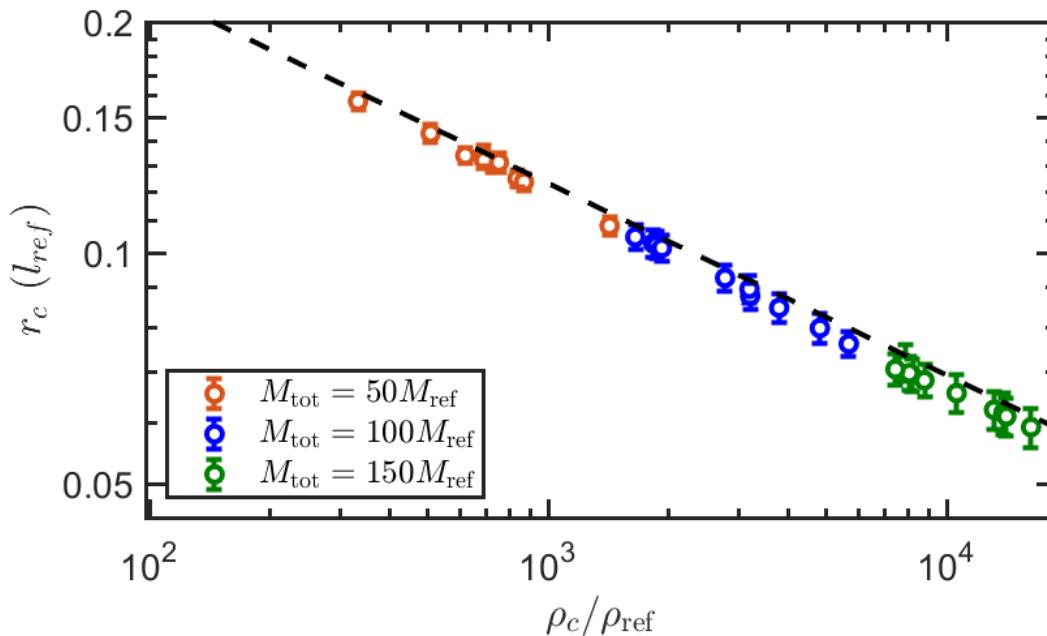
Halo structure

Energy functional of the SP system

$$E = E[\Psi^*, \Psi] = \frac{1}{m} \int d\mathbf{r} \Psi^* \left[-\frac{\hbar^2 \nabla^2}{2m} + \frac{1}{2} m\Phi \right] \Psi$$

Energetically preferred soliton solution, $\Psi(r) = \sqrt{\rho_{\text{soliton}}(r)}$, for a fixed M_{soliton}

$$r_c = \left(\frac{11303\lambda^2}{128} \frac{\hbar^2}{m^2 G \pi} \right)^{\frac{1}{4}} \rho_c^{-\frac{1}{4}}$$



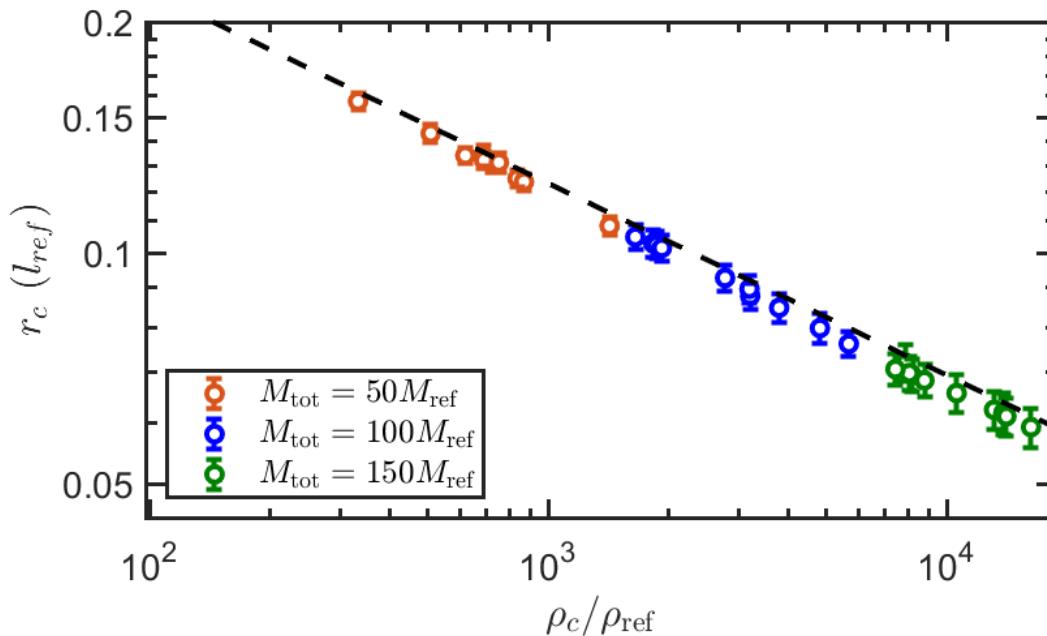
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$$r_c = \left(\frac{11303\lambda^2}{128} \frac{\hbar^2}{m^2 G \pi} \right)^{\frac{1}{4}} \rho_c^{-\frac{1}{4}} \approx 23.686 \left(\frac{2.5 \times 10^{-23} \text{ eV}}{m} \right)^{\frac{1}{2}} \left(\frac{10^3 M_\odot \text{kpc}^{-3}}{\rho_c} \right)^{\frac{1}{4}} \text{kpc}$$



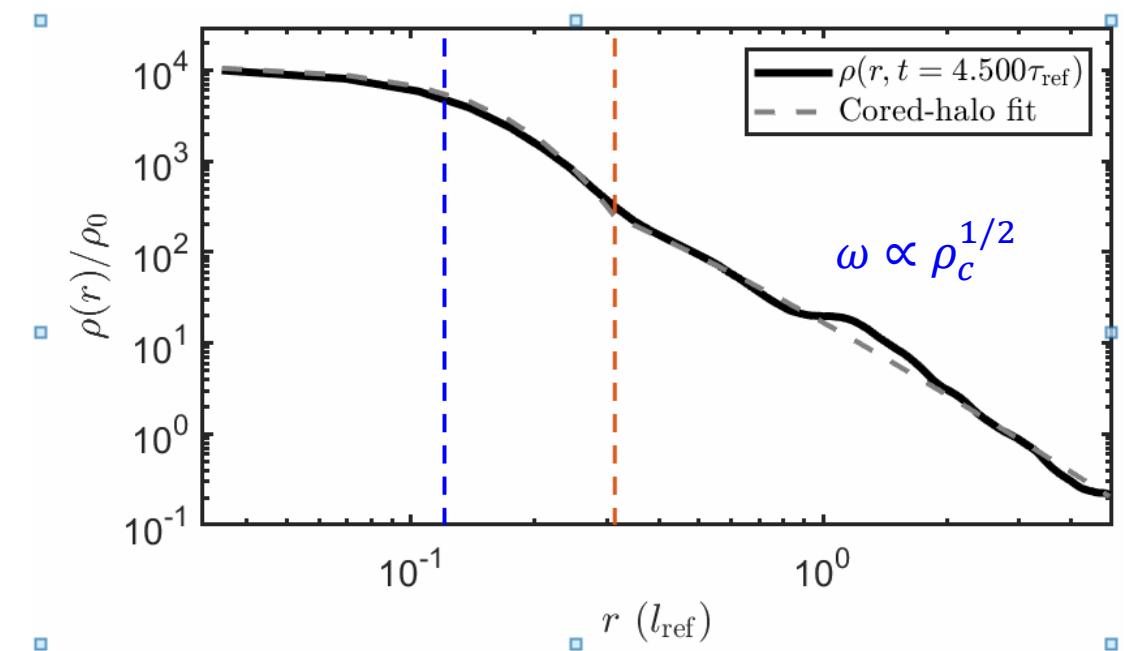
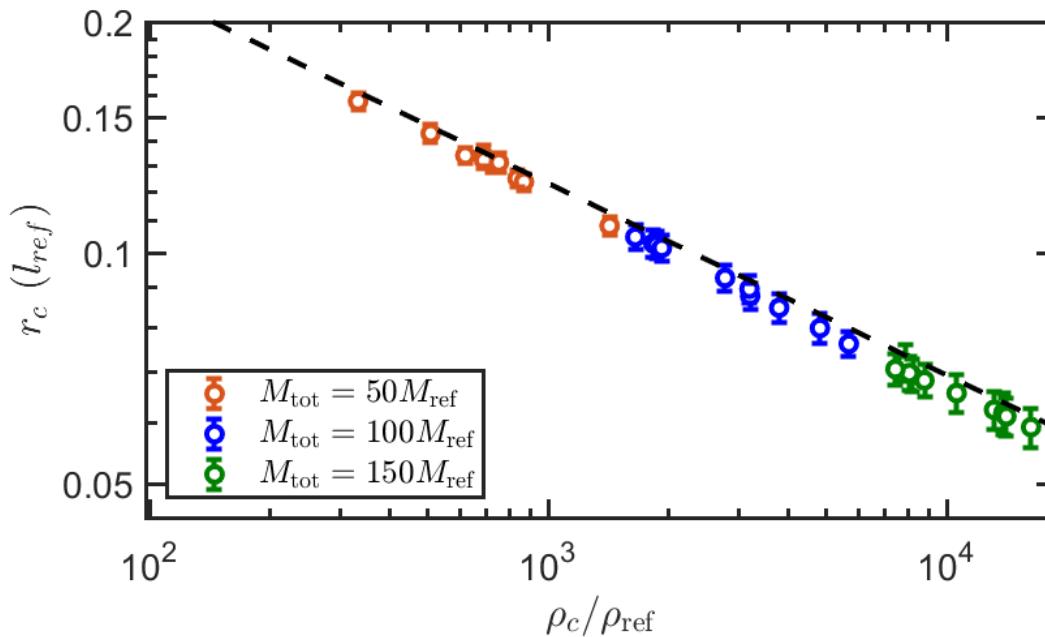
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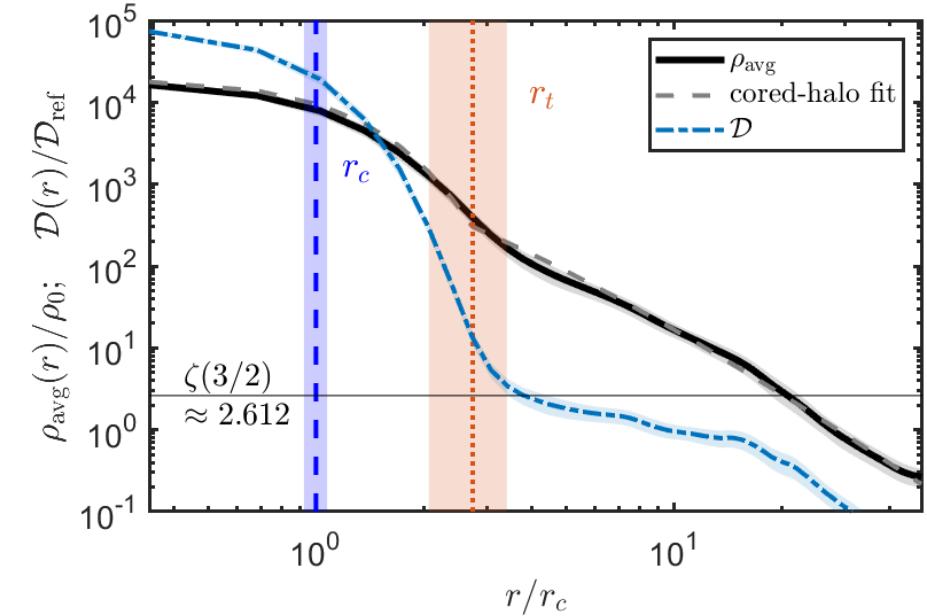


Halo structures

Coherence measures

Phase space density

$$\mathcal{D}(r) = \lambda_{\text{dB}}^3 n = \frac{\langle \rho'(r) \rangle}{\langle v'(r) \rangle^3} \left(\frac{\hbar^3 \rho_{\text{ref}}}{m^4 v_{\text{ref}}^3} \right)$$



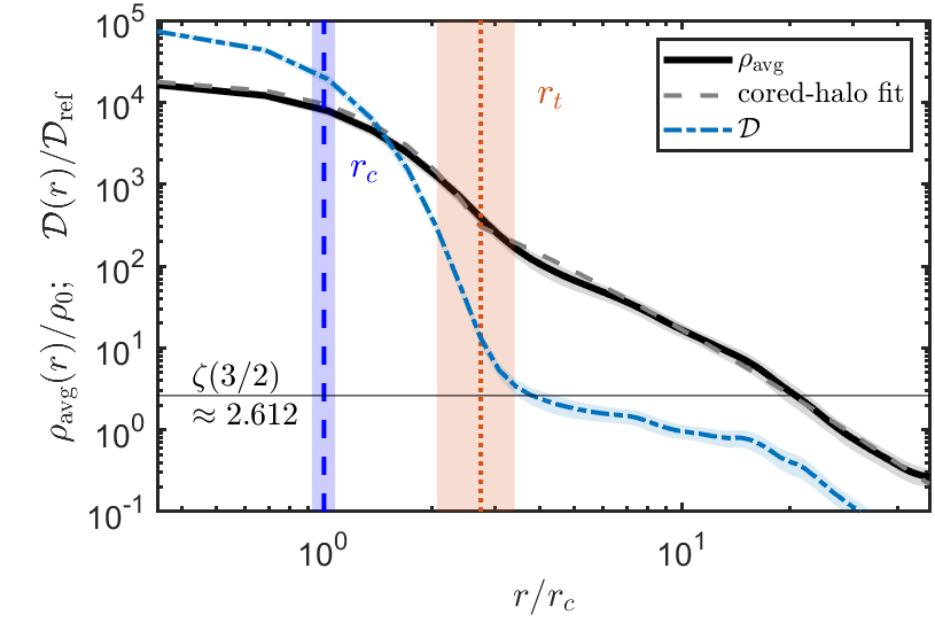
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High phase density!



Halo structures

Coherence measures

Phase space density

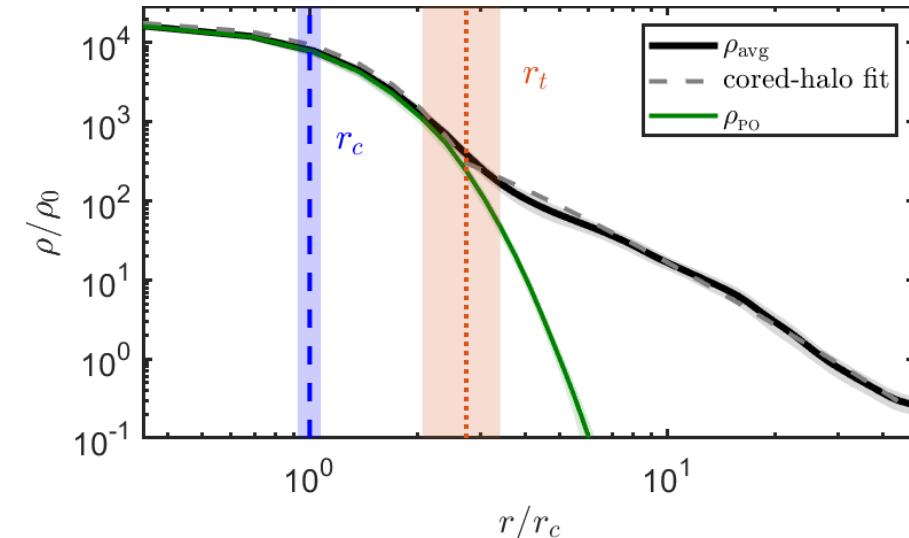
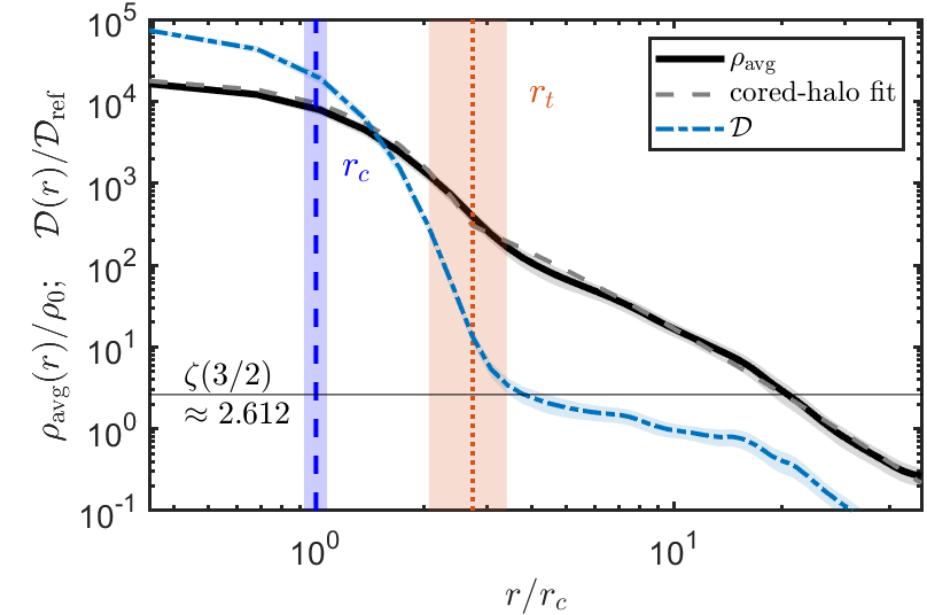
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High phase density!

Penrose-Onsager condensate mode

$$\hat{\varrho}(\mathbf{r}, \mathbf{r}') = \langle \Psi^*(\mathbf{r}') \Psi(\mathbf{r}) \rangle$$

$$\int \hat{\varrho}(\mathbf{r}, \mathbf{r}') \Psi_{\text{PO}}(\mathbf{r}') = N_{\text{PO}} \Psi_{\text{PO}}(\mathbf{r})$$



Halo structures

Coherence measures

Phase space density

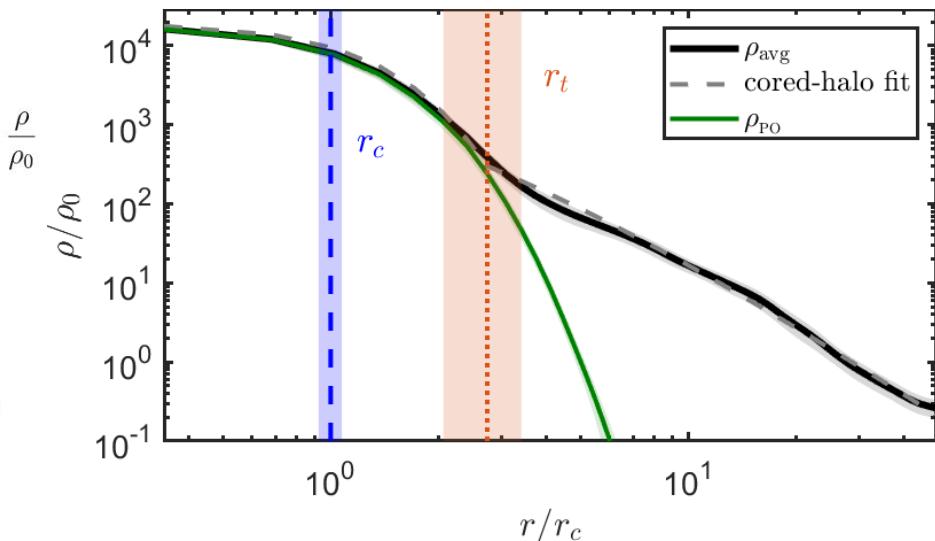
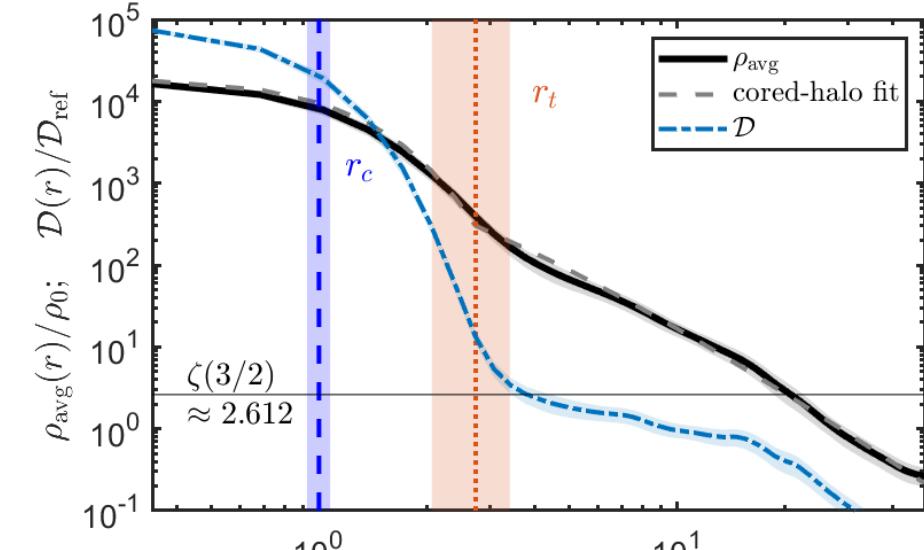
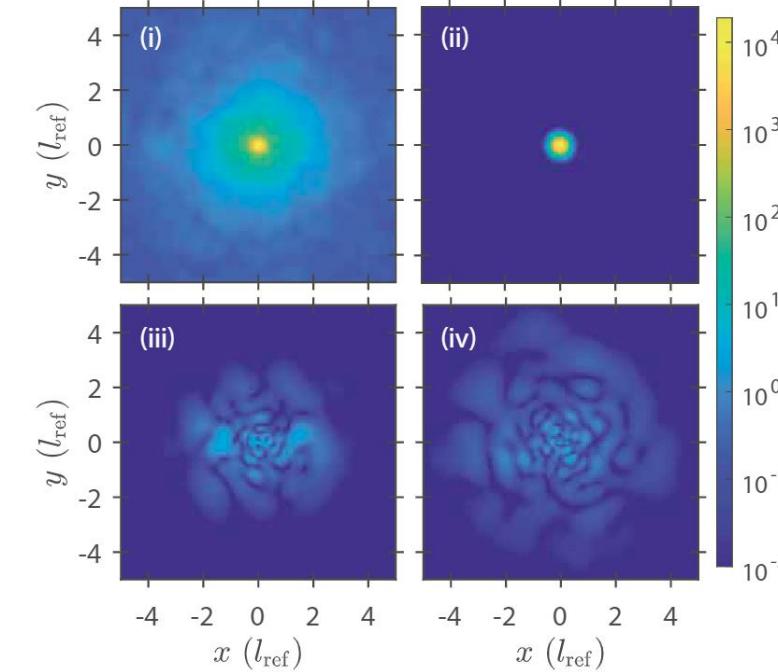
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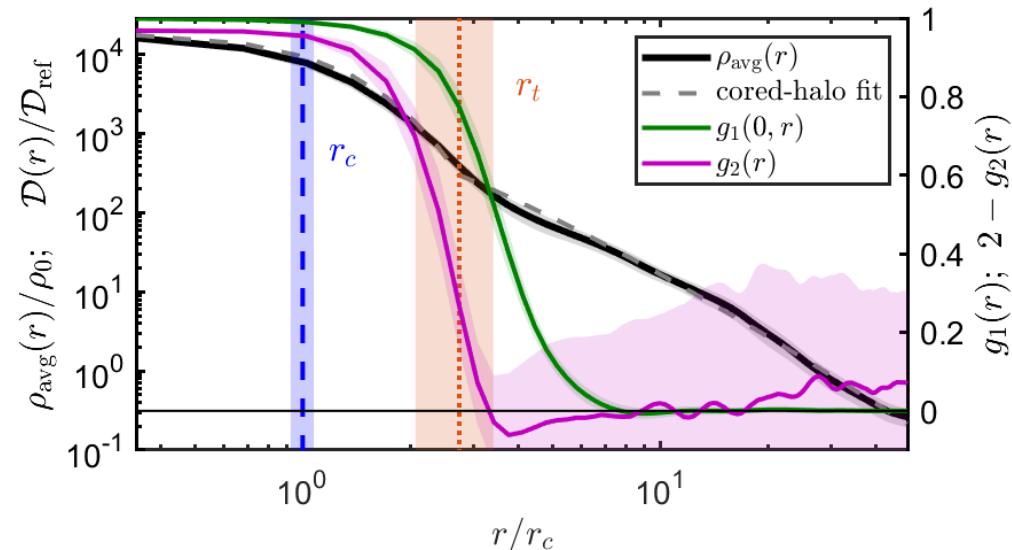
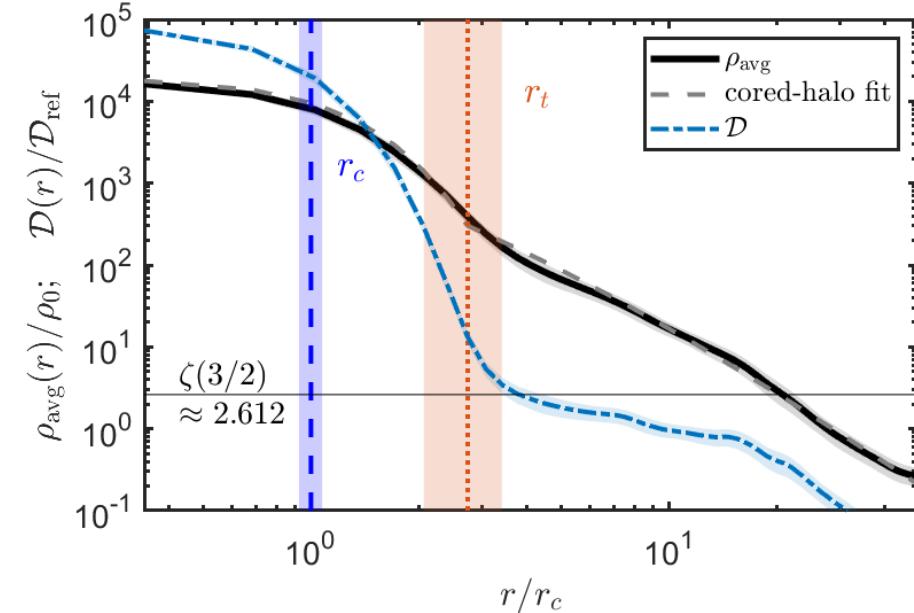
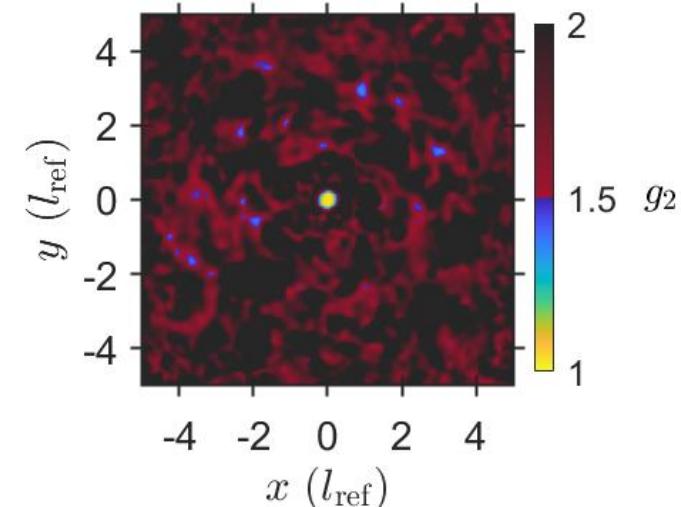
First-order spatial correlation function

$$g_1(\mathbf{r}, \mathbf{r}') = \frac{\langle \Psi^*(\mathbf{r}') \Psi(\mathbf{r}) \rangle}{\sqrt{\langle |\Psi(\mathbf{r}')|^2 \rangle \langle |\Psi(\mathbf{r})|^2 \rangle}} \quad \mathbf{r}' = \mathbf{0}$$

Second-order spatial auto correlation function

$$g_2(\mathbf{r}) = \frac{\langle \rho^2(\mathbf{r}) \rangle}{\langle \rho(\mathbf{r}) \rangle^2}$$

- $g_2 = 1$ for coherent density
- $g_2 = 2$ for strong density fluctuation



Halo structures

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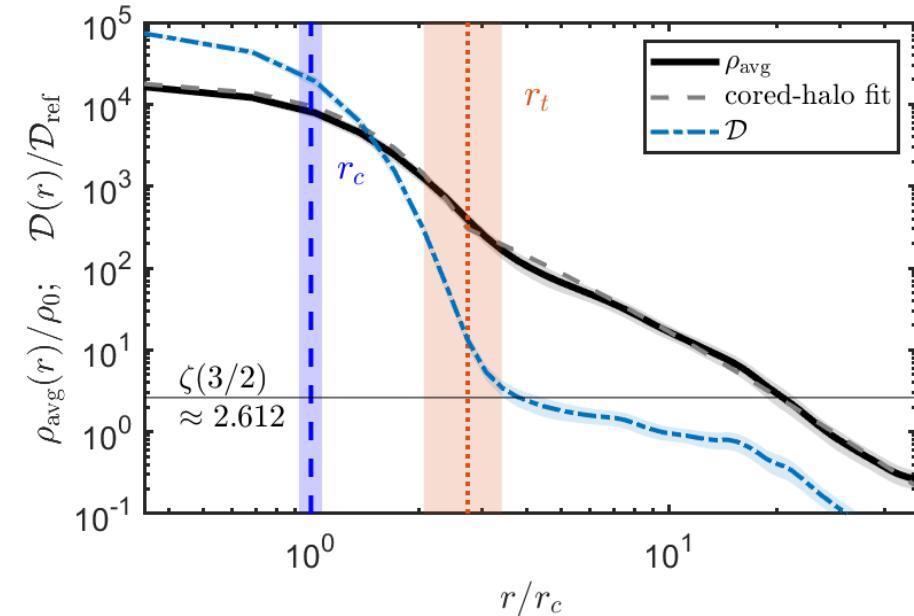
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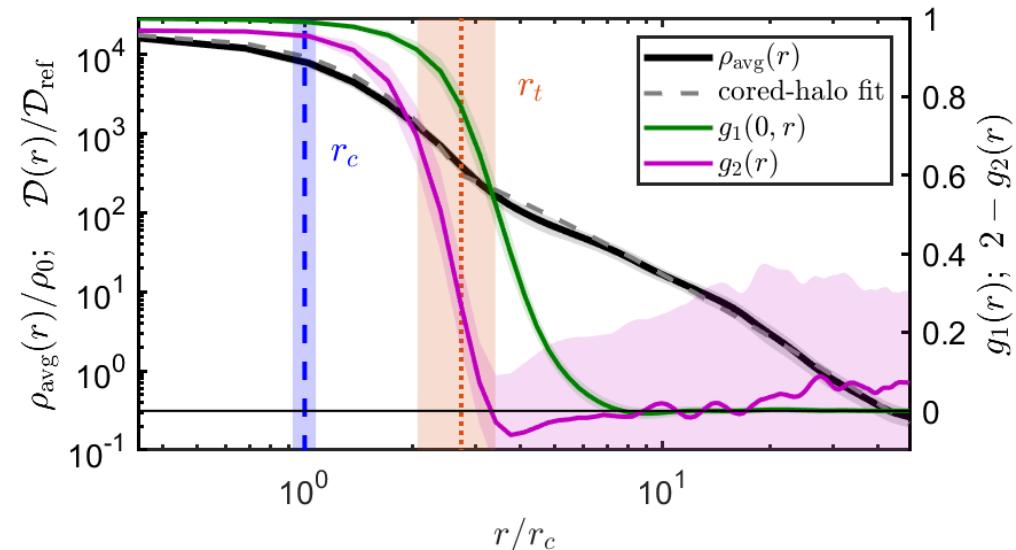
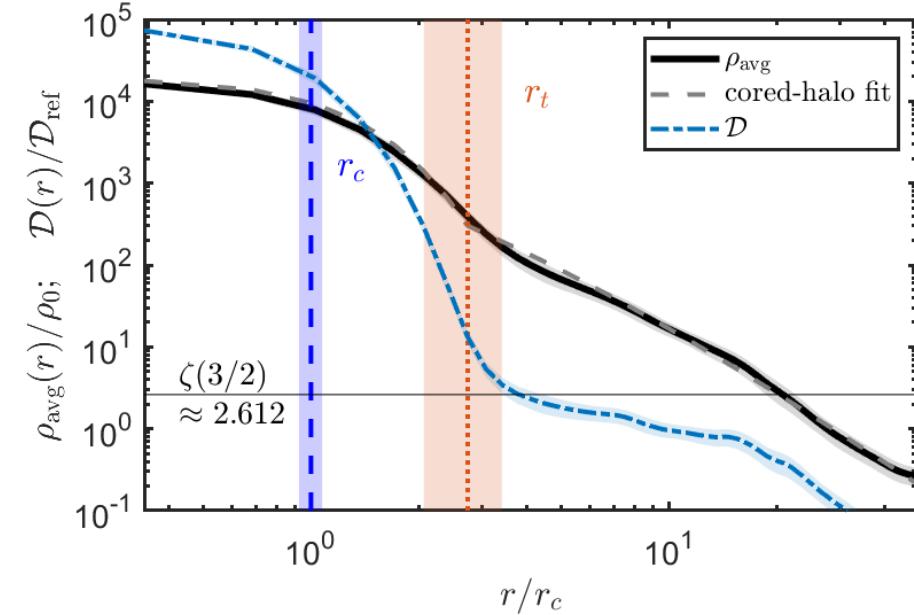
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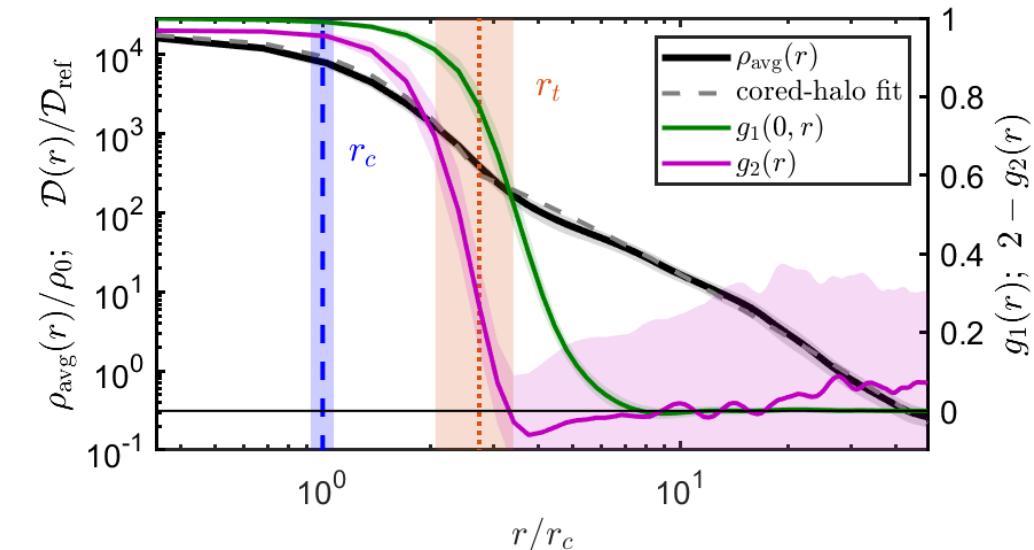
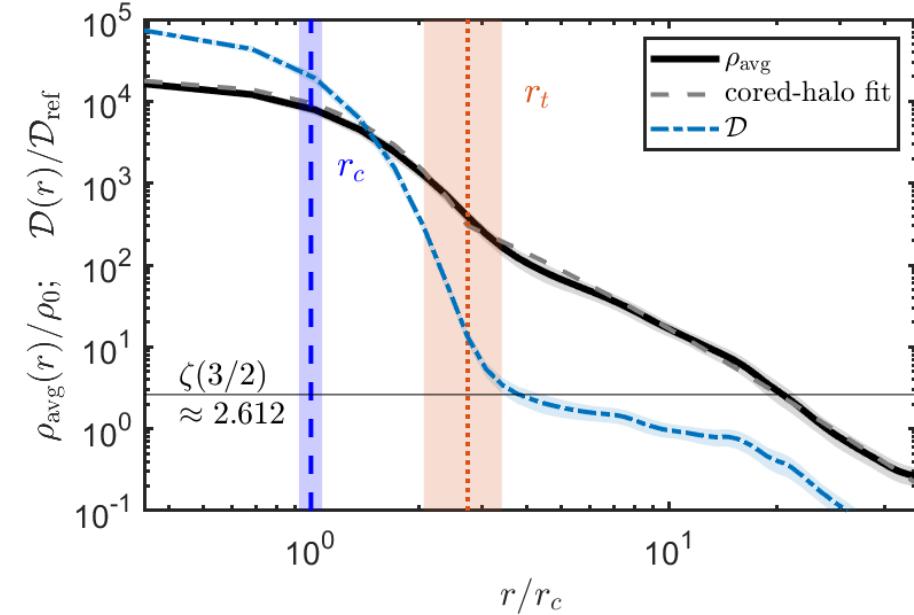
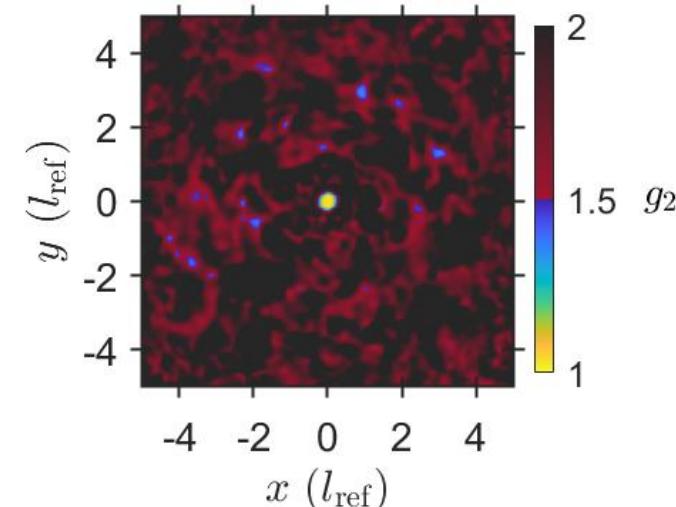
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Halo structures

Energy profiles

$$E = \int d\mathbf{r} \left\{ \frac{-\hbar^2 \nabla^2 / 2m^2}{[\varepsilon_{ke}(\mathbf{r}) + \varepsilon_{qp}(\mathbf{r})]} + \Phi(\mathbf{r})\rho(\mathbf{r})/2 \right\} = E_{ke} + E_{qp}$$
$$\varepsilon_{ke}(\mathbf{r}) = \frac{1}{2}\rho(\mathbf{r})|\mathbf{v}(\mathbf{r})|^2 \quad \varepsilon_{qp}(\mathbf{r}) = \frac{\hbar^2}{2m^2} |\nabla\sqrt{\rho(\mathbf{r})}|^2$$

$$E_{qp} \gtrsim 2E_{ke}$$

Halo structures

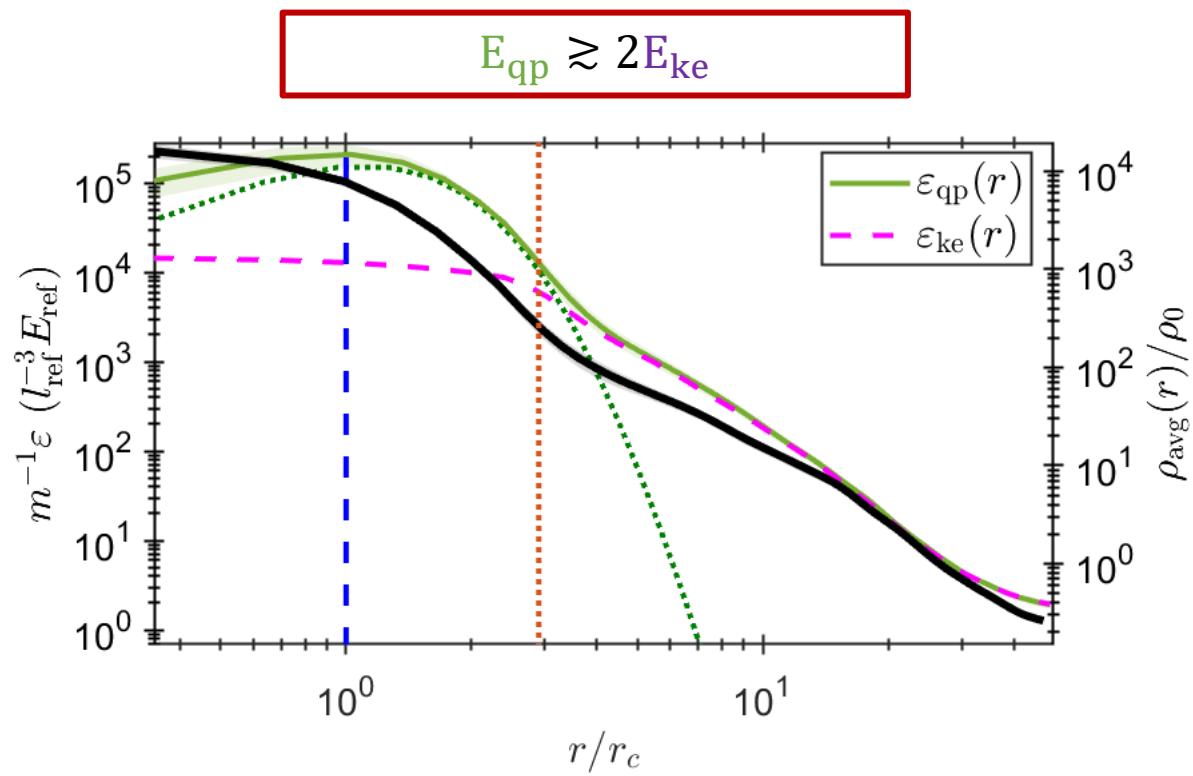
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$$\varepsilon_{qp}(\mathbf{r}) = \frac{\hbar^2}{2m^2} \left| \nabla \sqrt{\rho(\mathbf{r})} \right|^2$$

$$E_{qp} \approx 2E_{ke}$$



Halo structures

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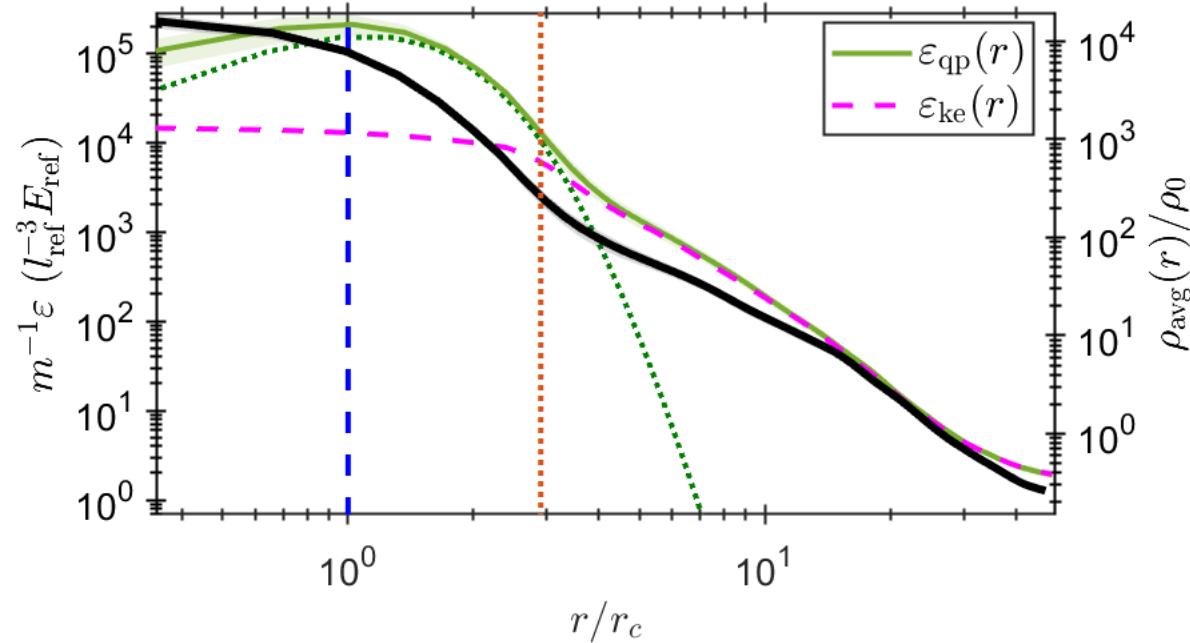
$$E_{ke} = E_{ke}^c + E_{ke}^i$$

$$\varepsilon_{ke}(\mathbf{r}) = \frac{1}{2} \left| \mathbf{F}^c(\mathbf{r}) + \mathbf{F}^i(\mathbf{r}) \right|^2$$

$$\nabla \times \mathbf{F}^c(\mathbf{r}) = 0$$

$$\nabla \cdot \mathbf{F}^i(\mathbf{r}) = 0$$

$$E_{qp} \approx 2E_{ke}$$



Halo structures

Energy profiles

$$E = \int d\mathbf{r} \left\{ \left[\varepsilon_{ke}(\mathbf{r}) + \varepsilon_{qp}(\mathbf{r}) \right] + \varepsilon_\Phi(\mathbf{r}) \right\} = E_{ke} + E_{qp}$$

$$\varepsilon_{ke}(\mathbf{r}) = \frac{1}{2} \rho(\mathbf{r}) |\mathbf{v}(\mathbf{r})|^2 \quad \varepsilon_{qp}(\mathbf{r}) = \frac{\hbar^2}{2m^2} \left| \nabla \sqrt{\rho(\mathbf{r})} \right|^2$$

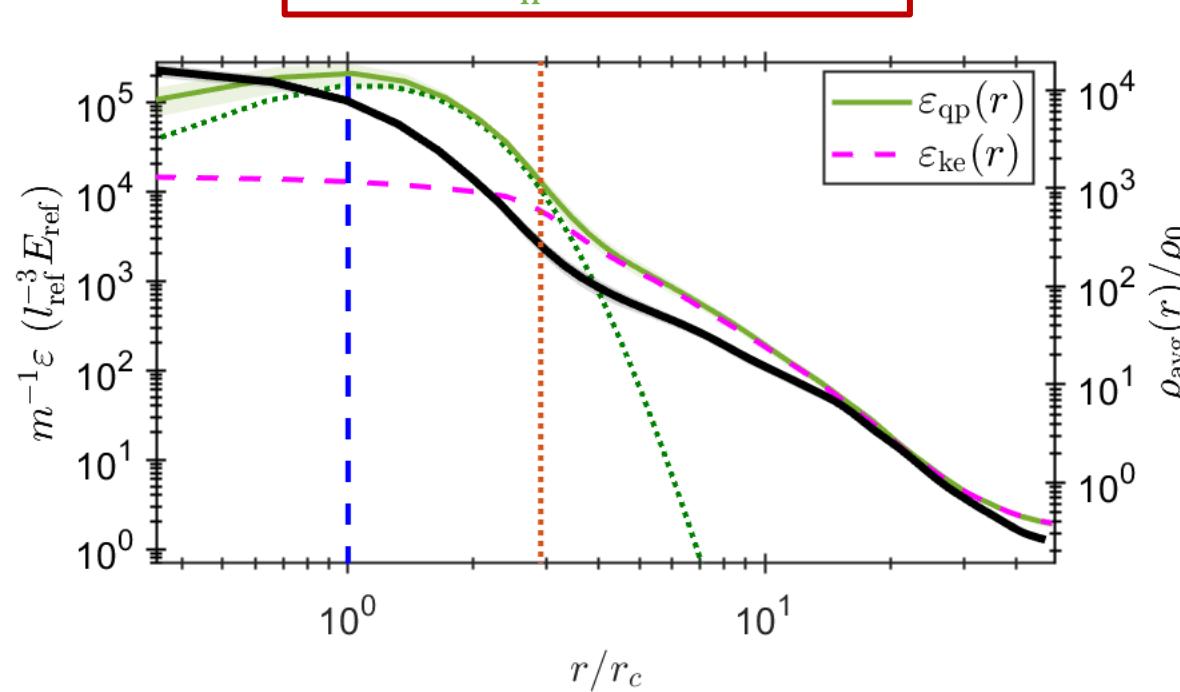
$$E_{ke} = E_{ke}^c + E_{ke}^i$$

$$\varepsilon_{ke}(\mathbf{r}) = \frac{1}{2} \left| \mathbf{F}^c(\mathbf{r}) + \mathbf{F}^i(\mathbf{r}) \right|^2$$

$$\nabla \times \mathbf{F}^c(\mathbf{r}) = 0$$

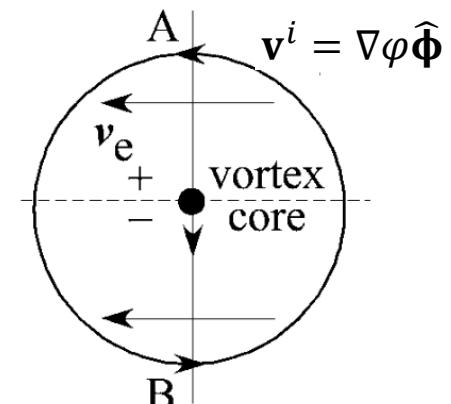
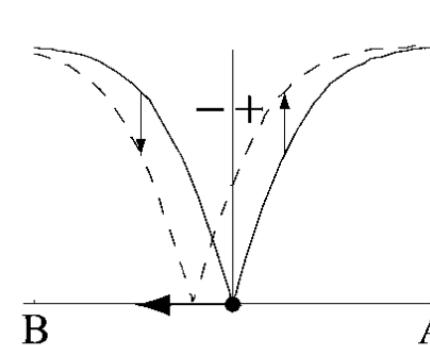
$$\nabla \cdot \mathbf{F}^i(\mathbf{r}) = 0$$

$$E_{qp} \approx 2E_{ke}$$



Quantized vortex

The incompressible/rotational component corresponds to the topological defect



Halo structures

Energy profiles

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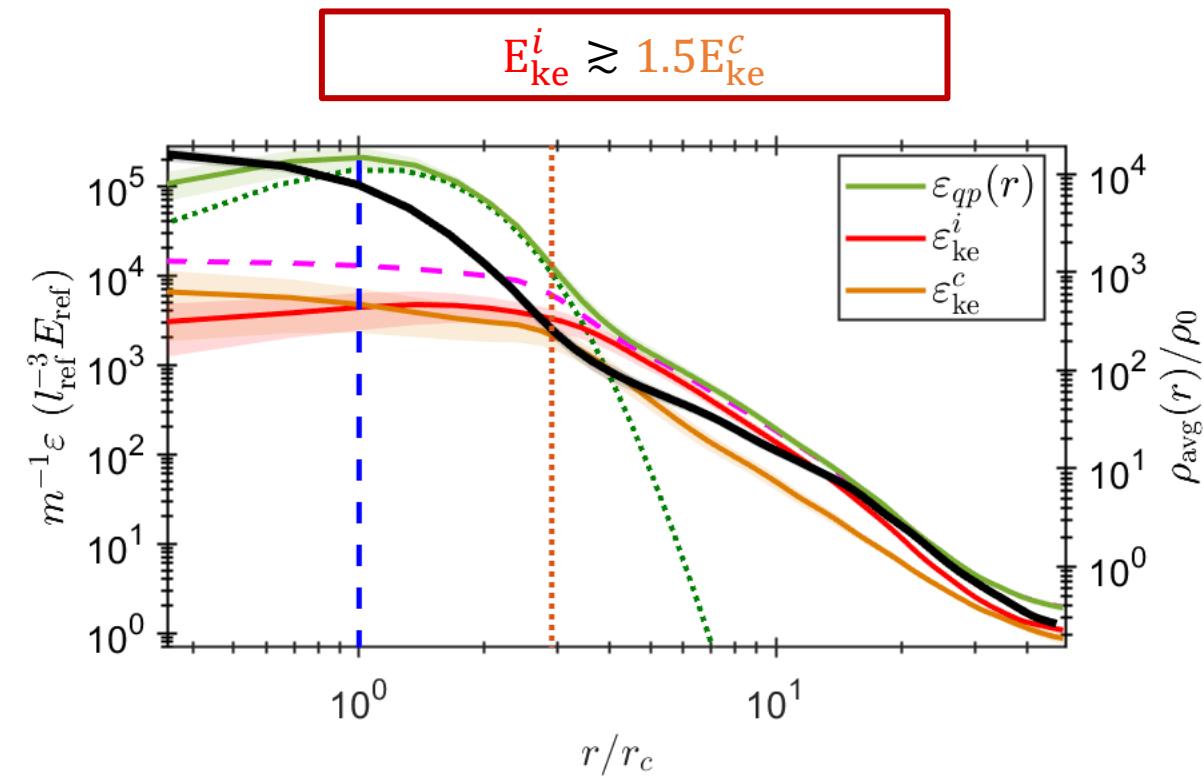
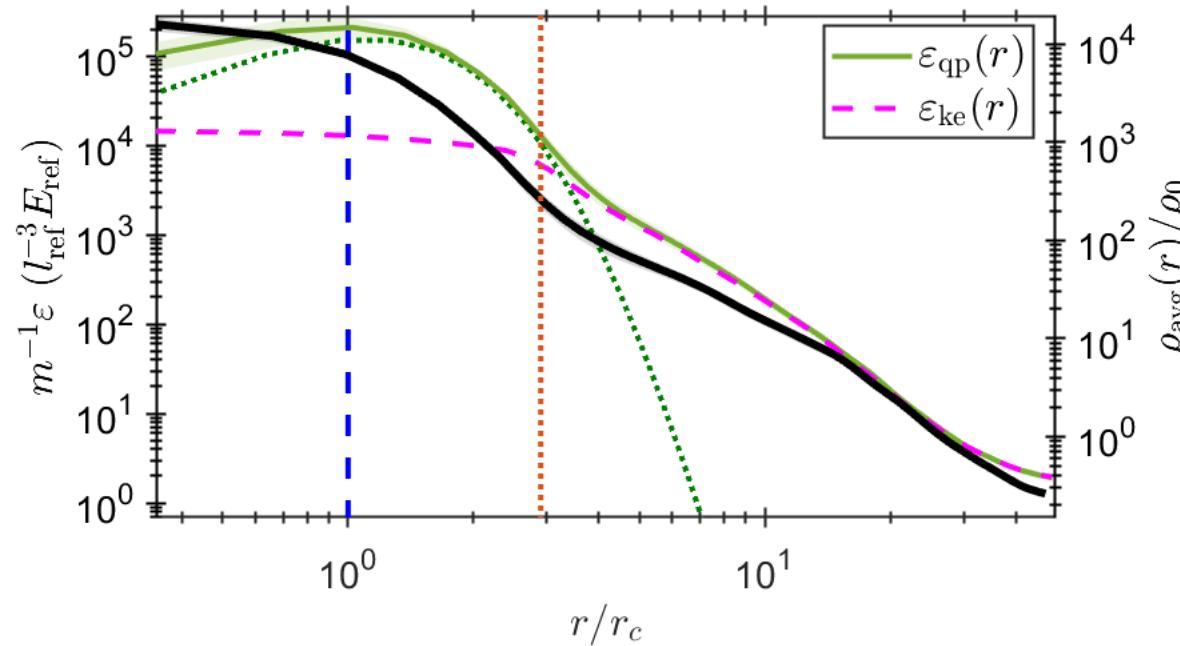
$$E_{ke} = E_{ke}^c + E_{ke}^i$$

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$$\nabla \times \mathbf{F}^c(\mathbf{r}) = 0$$

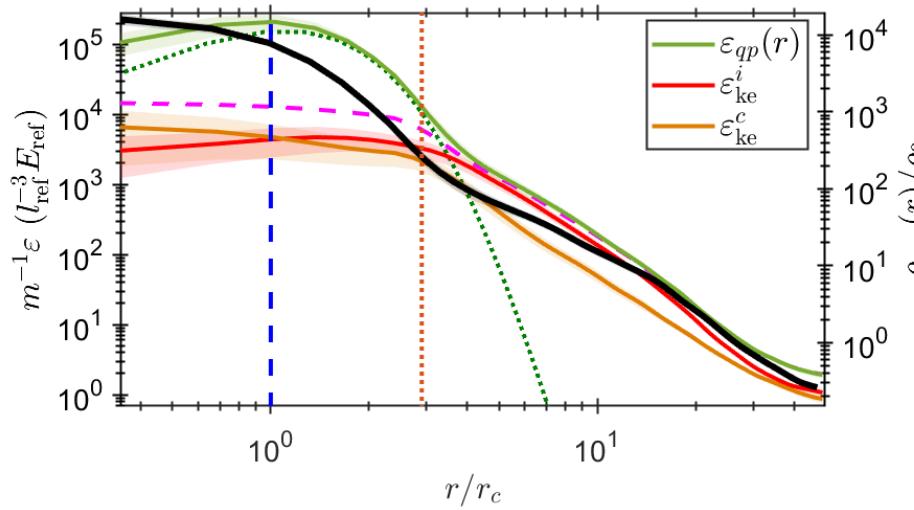
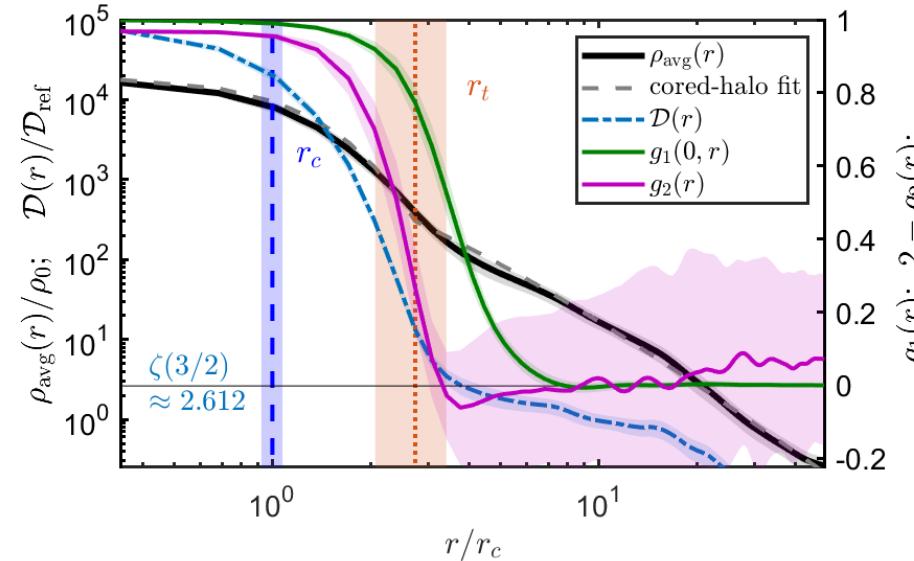
$$\nabla \cdot \mathbf{F}^i(\mathbf{r}) = 0$$

$$E_{qp} \approx 2E_{ke}$$



Vortices in FDM halos and granule size

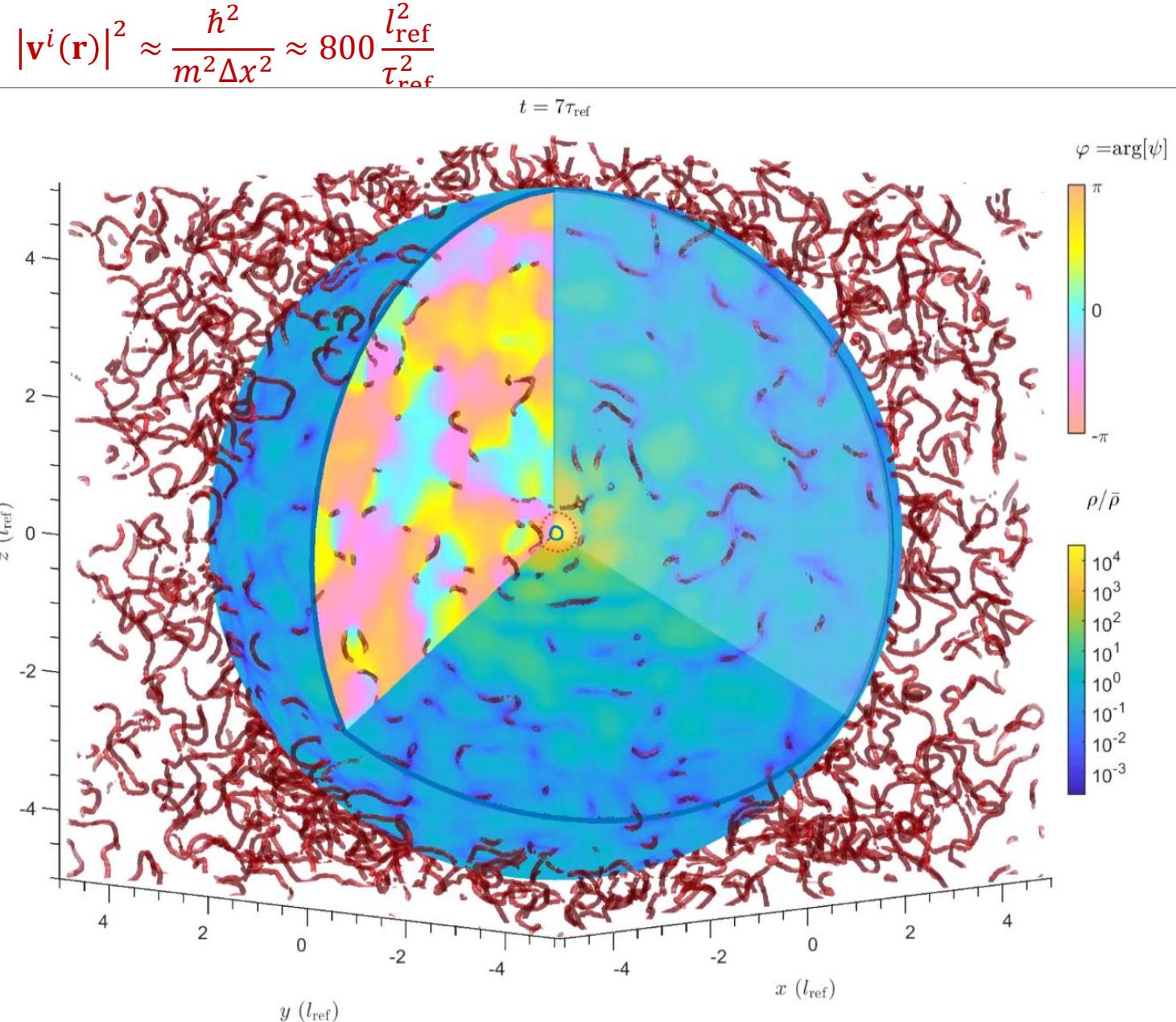
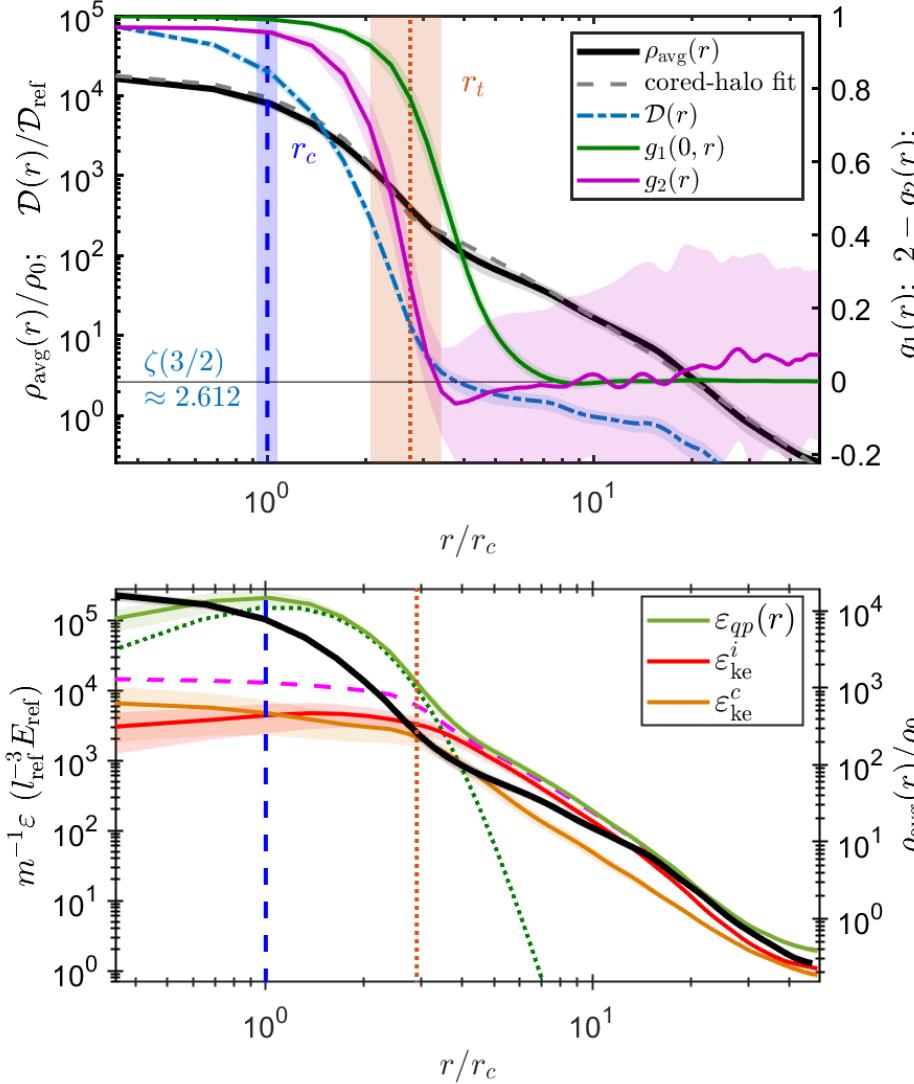
Vortical structure visualization



* Vortices in FDM have also been investigated in L. Hui et al., 2021 & Mocz et al., MNRAS 471, 4559 (2017)

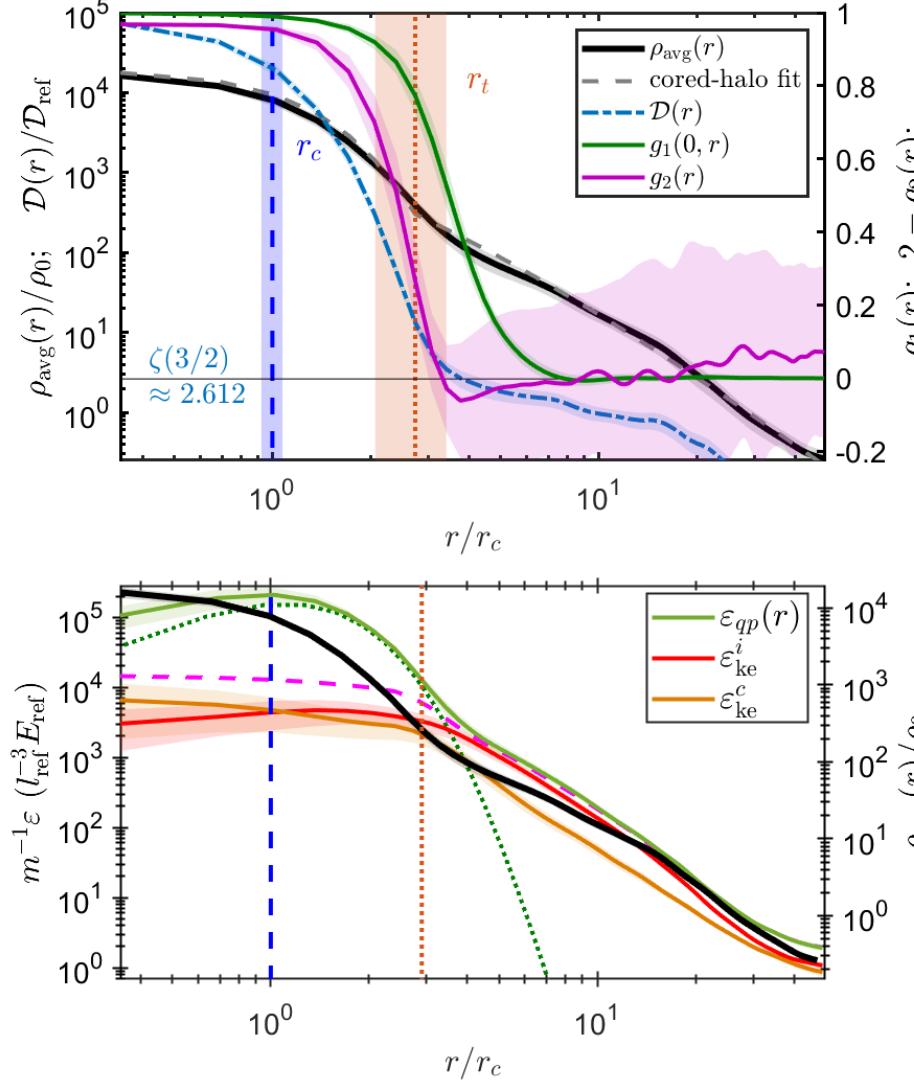
Vortices in FDM halos and granule size

Vortical structure visualization

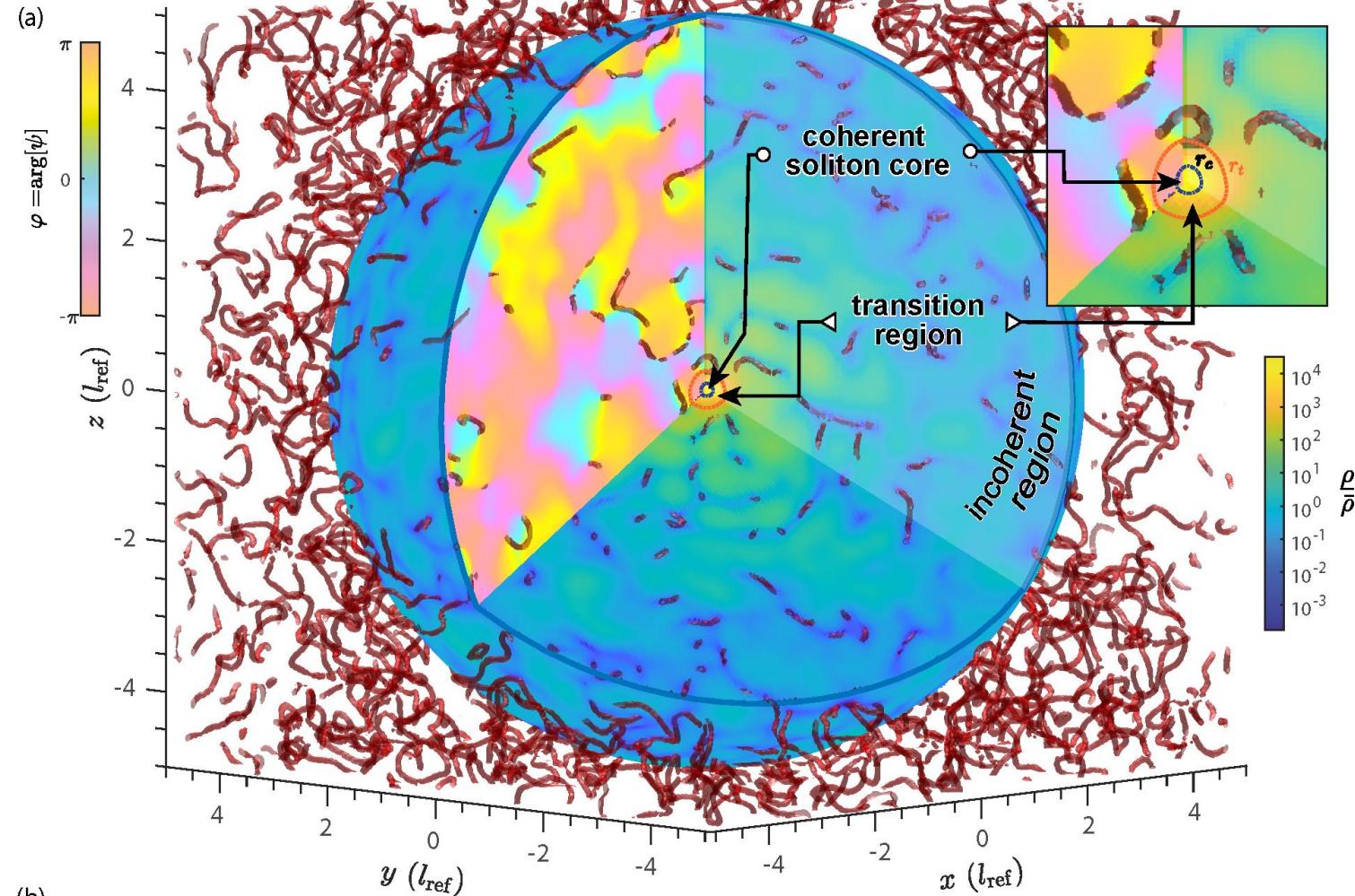


Vortices in FDM halos and granule size

Vortical structure visualization

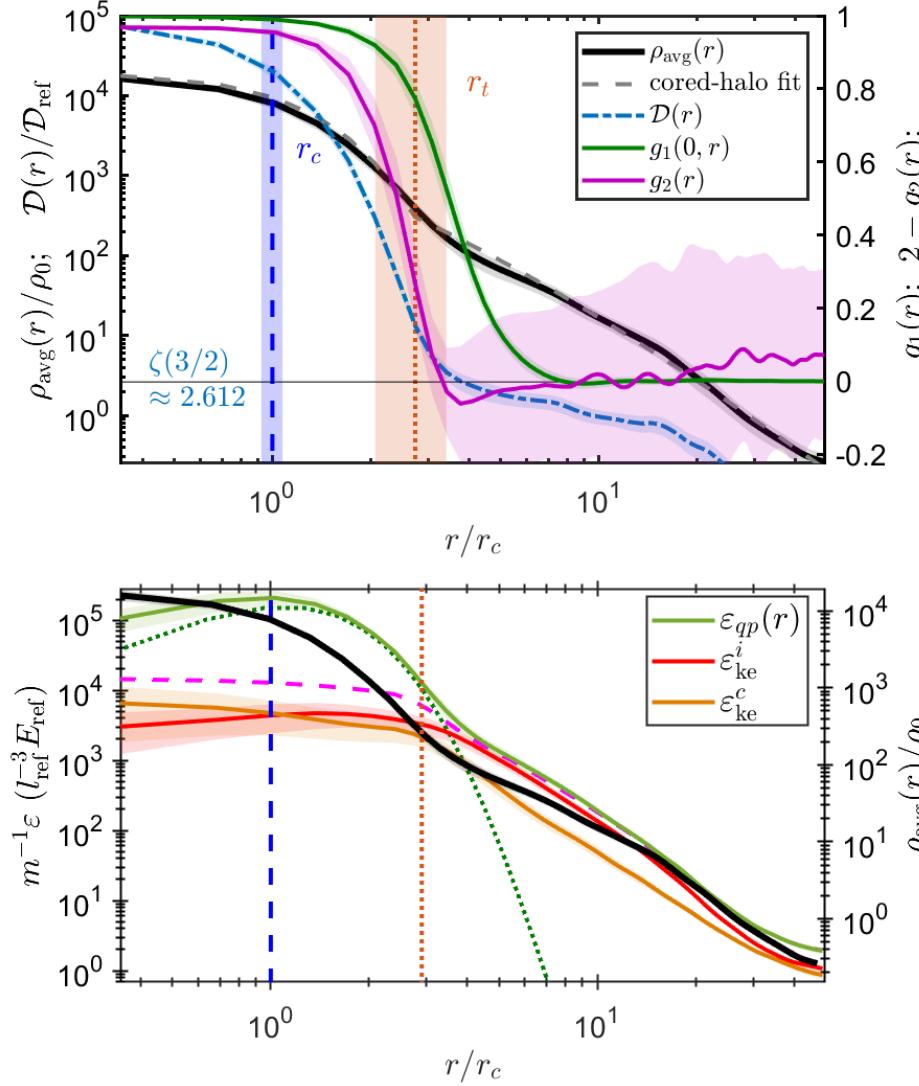


$$|\mathbf{v}^i(\mathbf{r})|^2 \approx \frac{\hbar^2}{m^2 \Delta x^2} \approx 800 \frac{l_{\text{ref}}^2}{\tau_{\text{ref}}^2}$$

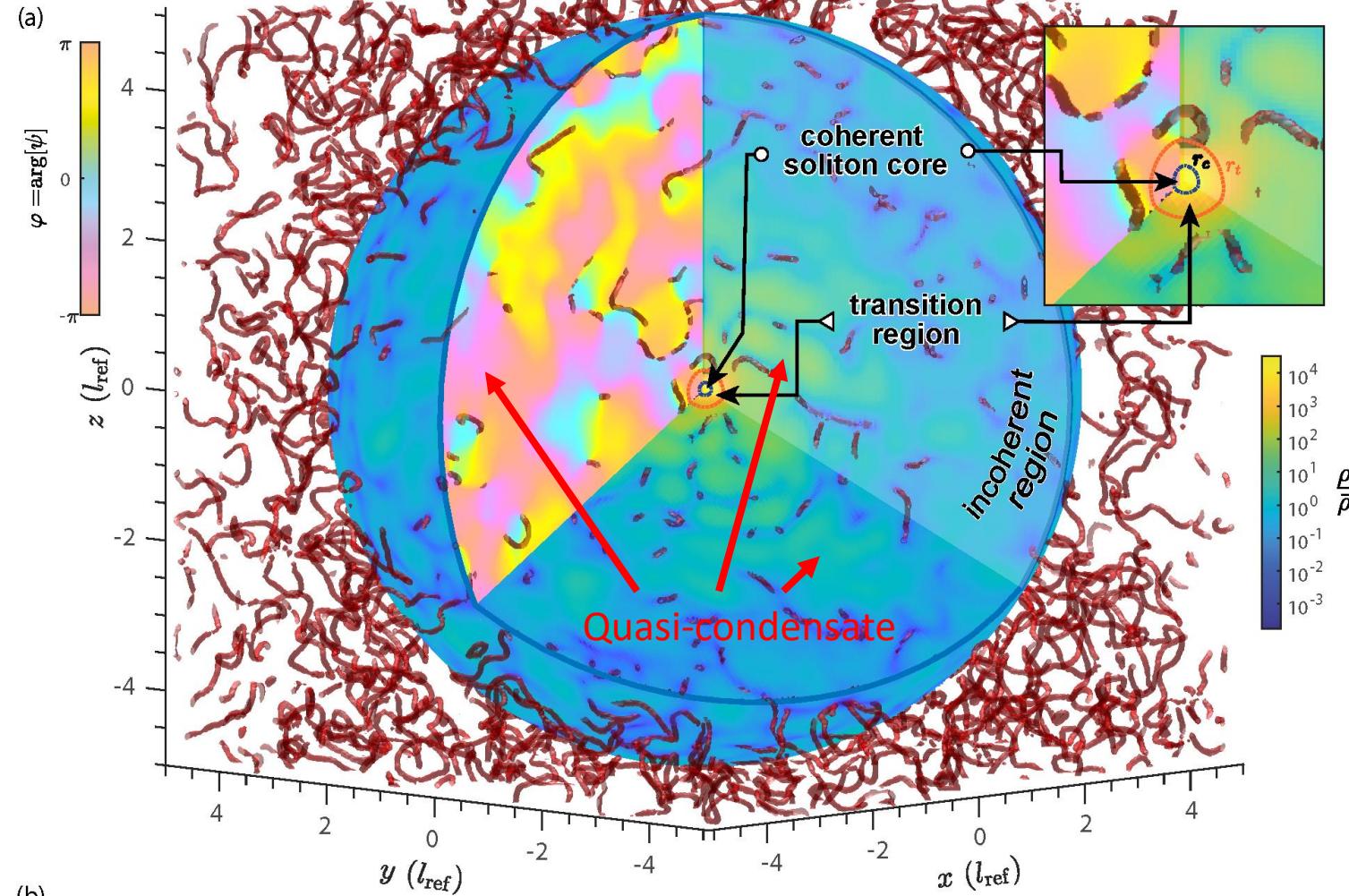


Vortices in FDM halos and granule size

Vortical structure visualization



$$|\mathbf{v}^i(\mathbf{r})|^2 \approx \frac{\hbar^2}{m^2 \Delta x^2} \approx 800 \frac{l_{\text{ref}}^2}{\tau_{\text{ref}}^2}$$



(b)

Vortices in FDM halos and granule size

Vortex energy and granule power spectra

Incompressible energy spectrum

$$\tilde{\varepsilon}_{\text{ke}}^i(k_r) \equiv \int d\Omega_k k^2 \varepsilon_{\text{ke}}^i(\mathbf{r}, t)$$

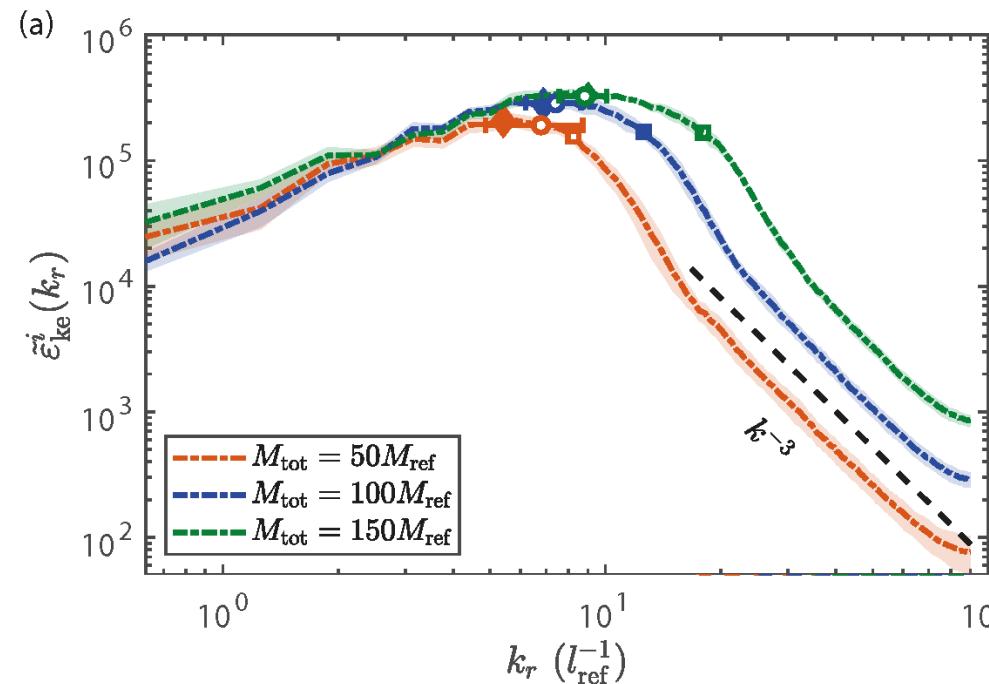
- No quasi-classical turbulence, $k^{-5/3}$
 - required vortex bundles in SF
[A. Baggaley et al., PRL 109, 205304 \(2012\)](#)
 - FDM vortices are more chaotic

[Mocz et al., MNRAS 471, 4559 \(2017\)](#)

- Vortex core structure, k^{-3} : $\rho |\mathbf{v}^i|^2 \approx \text{Const.}$ in large k

[Nore et al., Phys. of Fluids 9, 2644 \(1997\)](#)

[Stagg et al., PRA 94, 053632 \(2016\)](#)



Vortices in FDM halos and granule size

Vortex energy and granule power spectra

Incompressible energy spectrum

$$\tilde{\varepsilon}_{\text{ke}}^i(k_r) \equiv \int d\Omega_k k^2 \varepsilon_{\text{ke}}^i(\mathbf{r}, t)$$

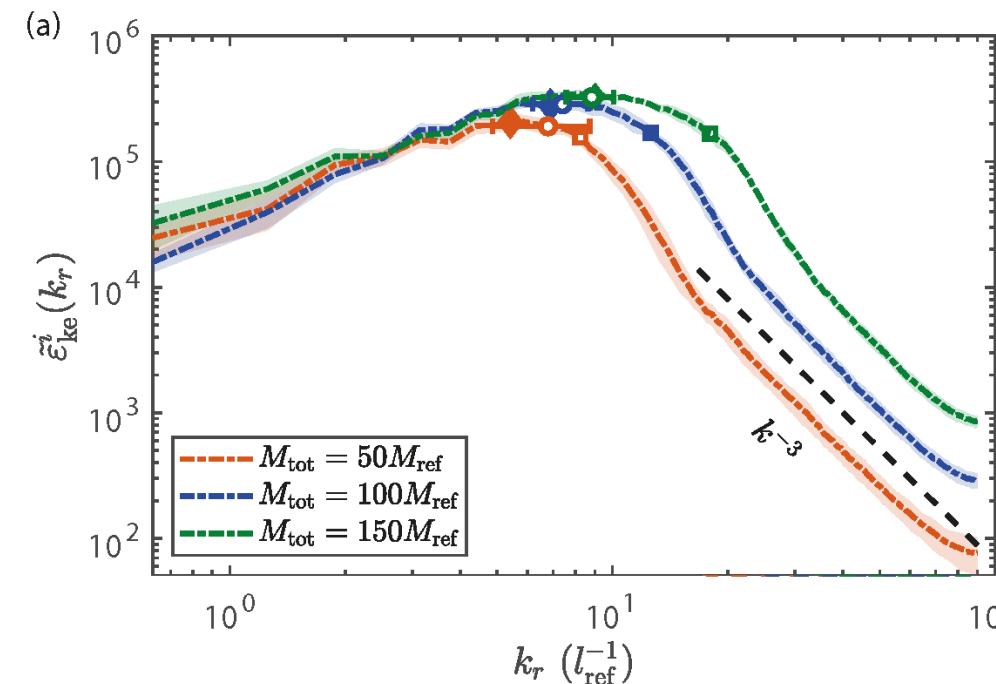
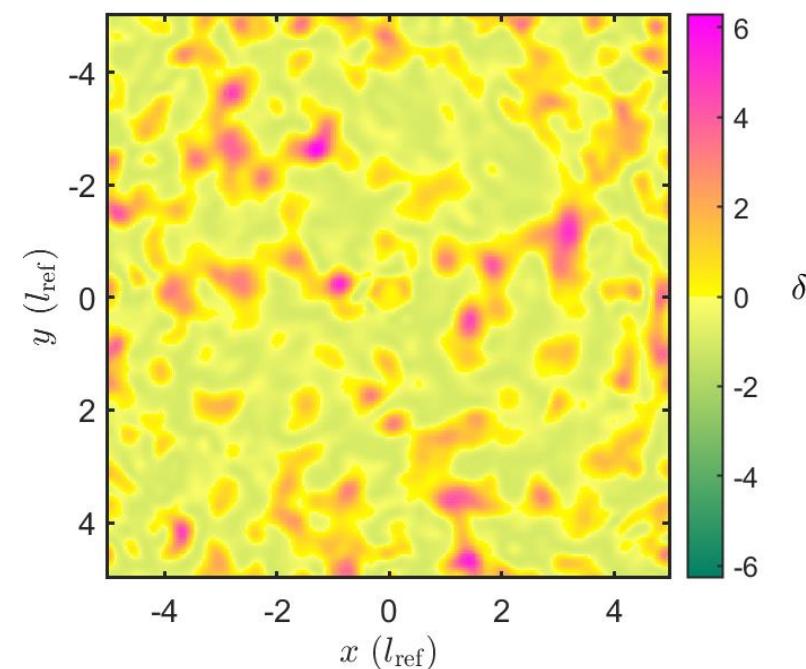
Granule power spectrum

$$\delta(\mathbf{r}, t) = \frac{\rho(\mathbf{r}, t) - \rho_{\text{avg}}(\mathbf{r})}{\rho_{\text{avg}}(\mathbf{r})}$$

Chan et al., MNRAS 478, 2686 (2018)

Lin et al. PRD 97, 103523 (2018)

Dutta Chowdhury et al. (2021)



Vortices in FDM halos and granule size

Vortex energy and granule power spectra

Incompressible energy spectrum

$$\tilde{\varepsilon}_{\text{ke}}^i(k_r) \equiv \int d\Omega_k k^2 \varepsilon_{\text{ke}}^i(\mathbf{r}, t)$$

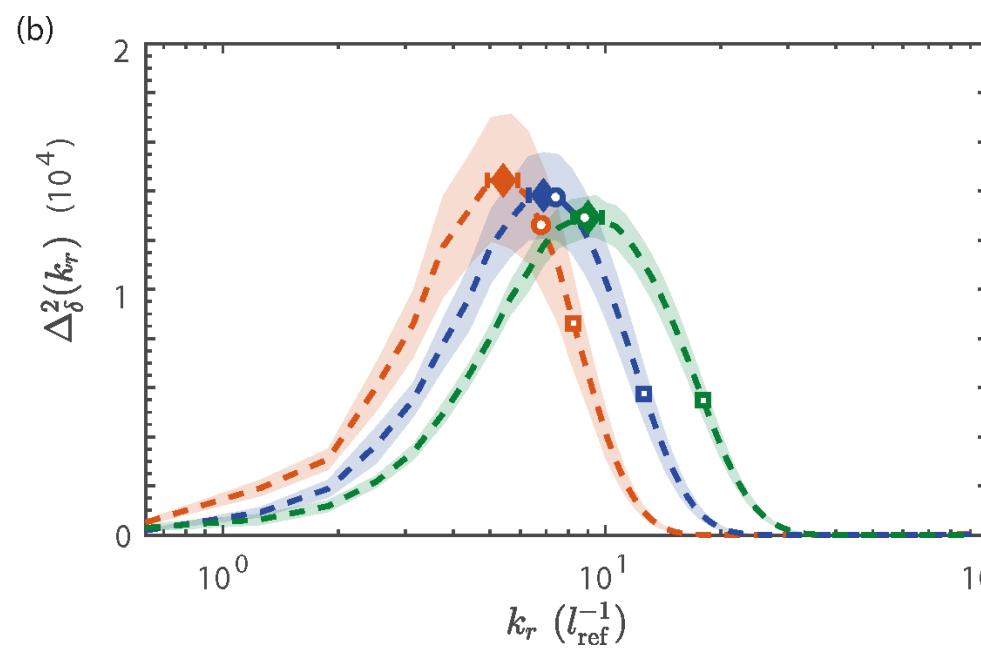
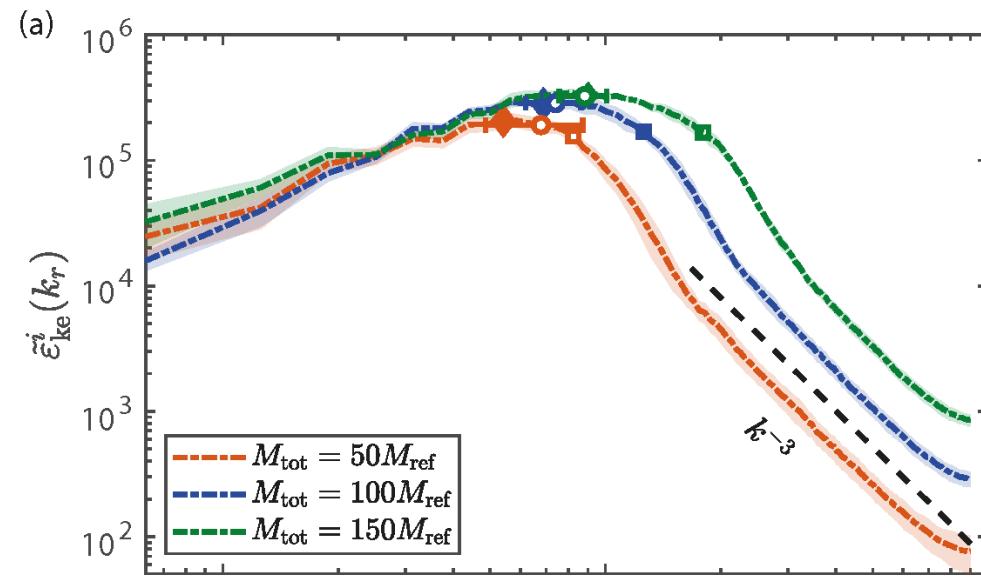
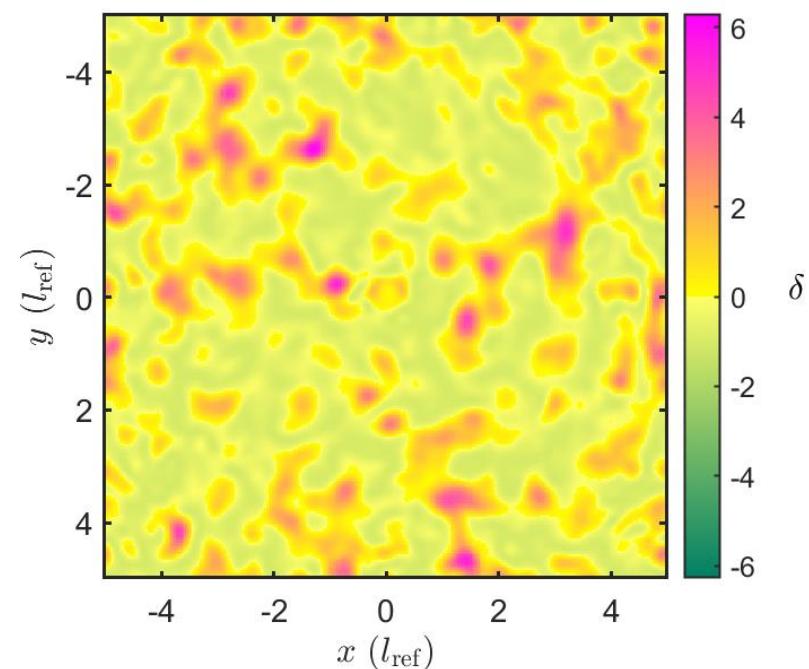
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Chan et al., MNRAS 478, 2686 (2018)

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Dutta Chowdhury et al. (2021)



Vortices in FDM halos and granule size

Vortex energy and granule power spectra

Incompressible energy spectrum

$$\tilde{\epsilon}_{\text{ke}}^i(k_r) \equiv \int d\Omega_k k^2 \epsilon_{\text{ke}}^i(\mathbf{r}, t)$$

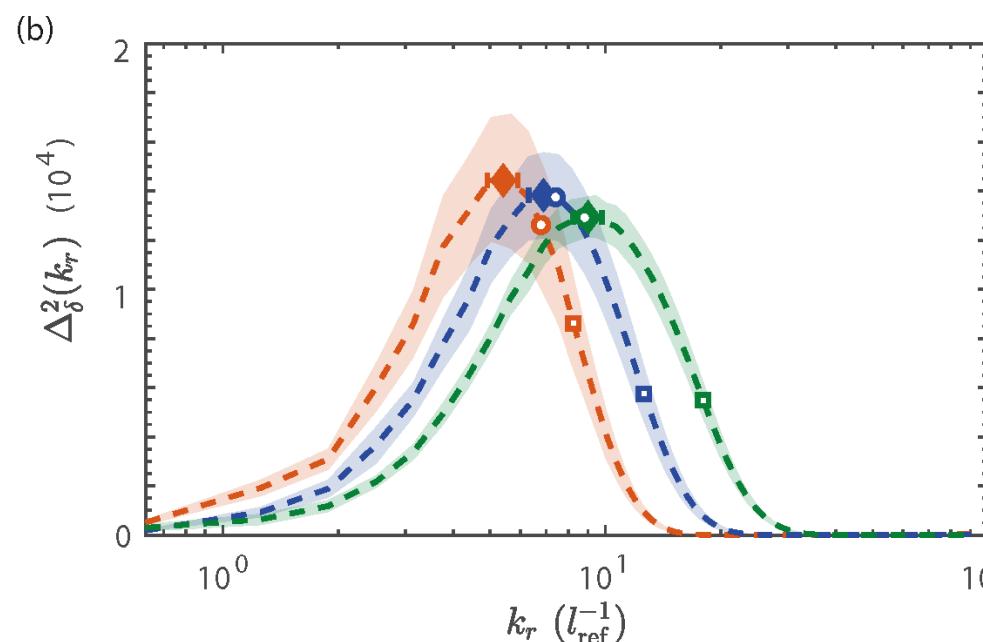
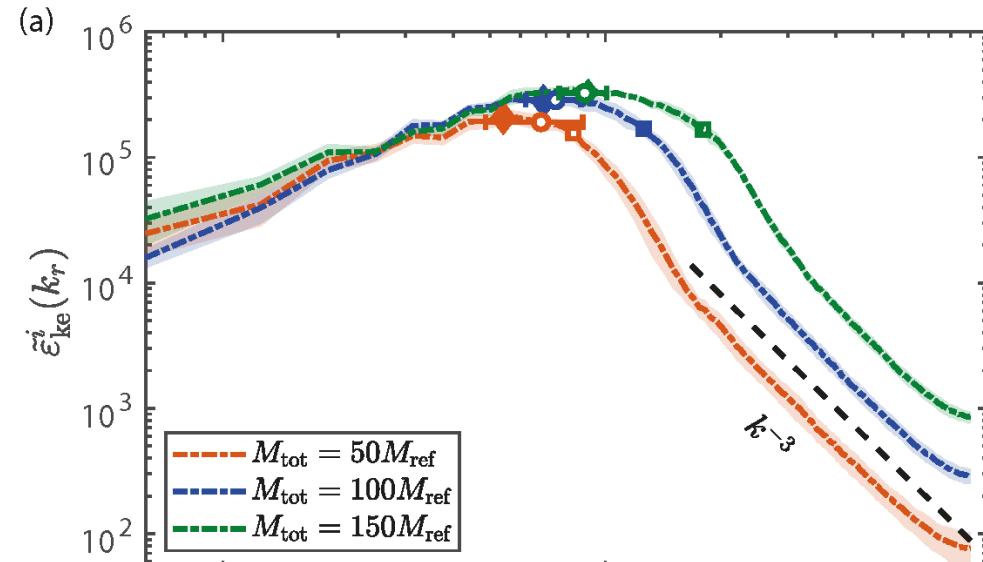
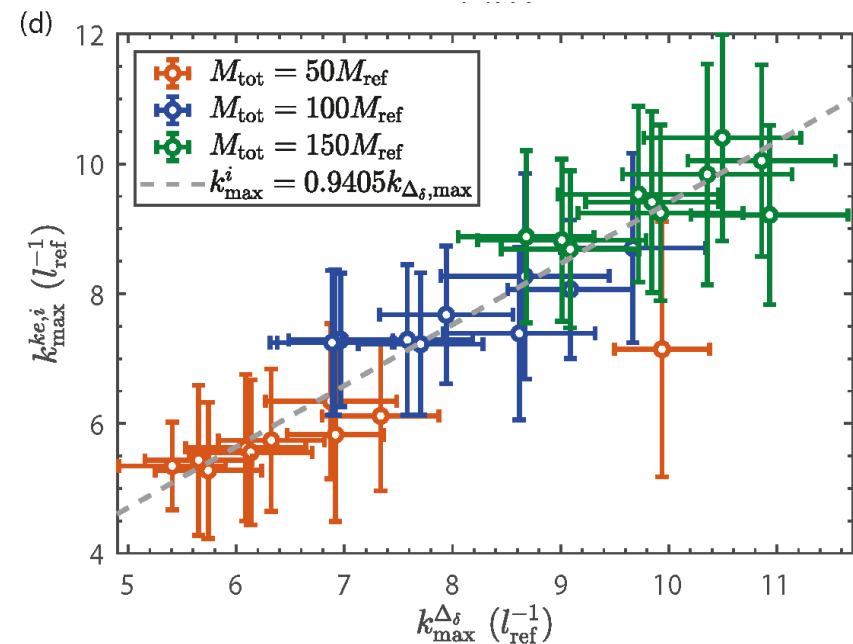
Granule power spectrum

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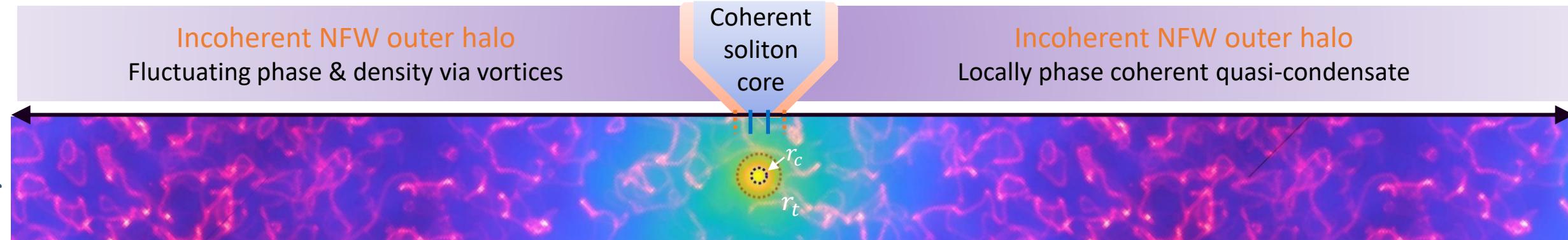
Chan et al., MNRAS 478, 2686 (2018)

Lin et al. PRD 97, 103523 (2018)

Dutta Chowdhury et al. (2021)



Conclusion



- The radial halo profile can be fairly caught by the **cored-halo fit** with the parameters r_c and r_t .
- The FDM mass can roughly be estimated by

$$\mathcal{D}_{\text{ref}} = \left(\frac{\hbar^3 \rho_{\text{ref}}}{m^4 v_{\text{ref}}^3} \right) = 1.25 \times 10^{90} \left(\frac{10^{-22} \text{ eV}}{mc^2} \right)^4 \left(\frac{\rho_{\text{ref}}}{10^3 M_\odot \text{kpc}^{-3}} \right) \left(\frac{250 \text{ km/s}}{v_{\text{ref}}} \right)^3$$

giving $m \approx [10^{-22}, 10] \text{ eV}$ for FDM in our simulation.

- The turbulent-like vortex dynamics results in the large density fluctuations in the outer halo.
- The inter-vortex distances and the granule length scale are comparable

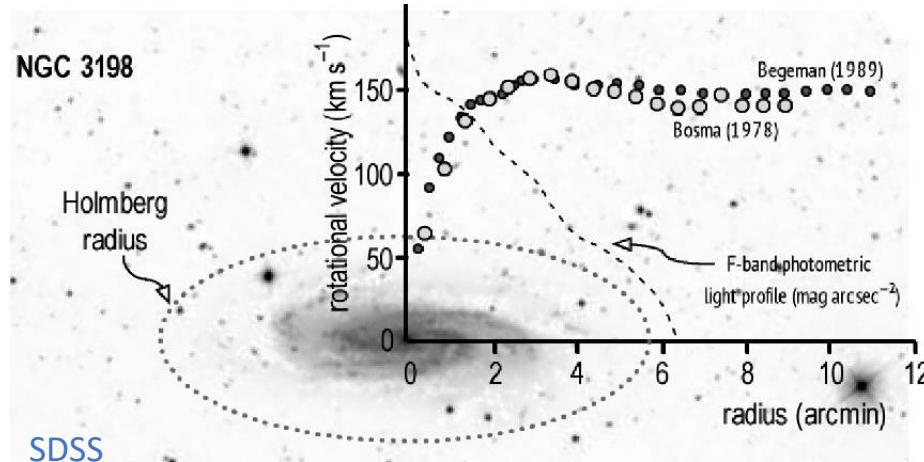
Thank you for your attention

Outline

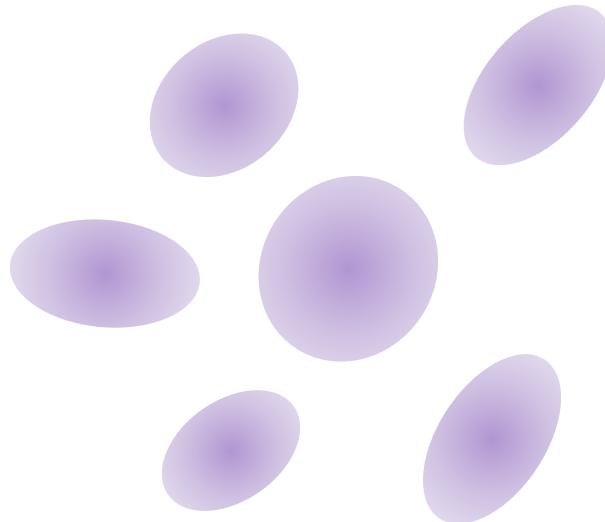
- Fuzzy dark matter
- Halo structures
- Vortices in FDM halos and granule size
- Conclusion

Fuzzy Dark Matter

- Rotational curve



- Galaxy alignment



- Lensing



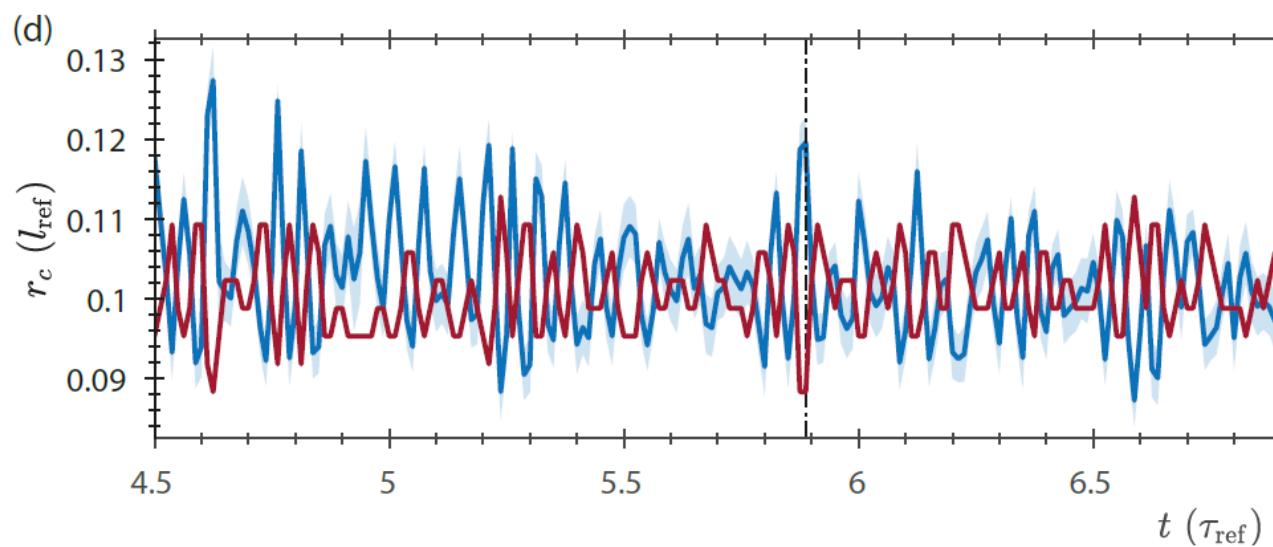
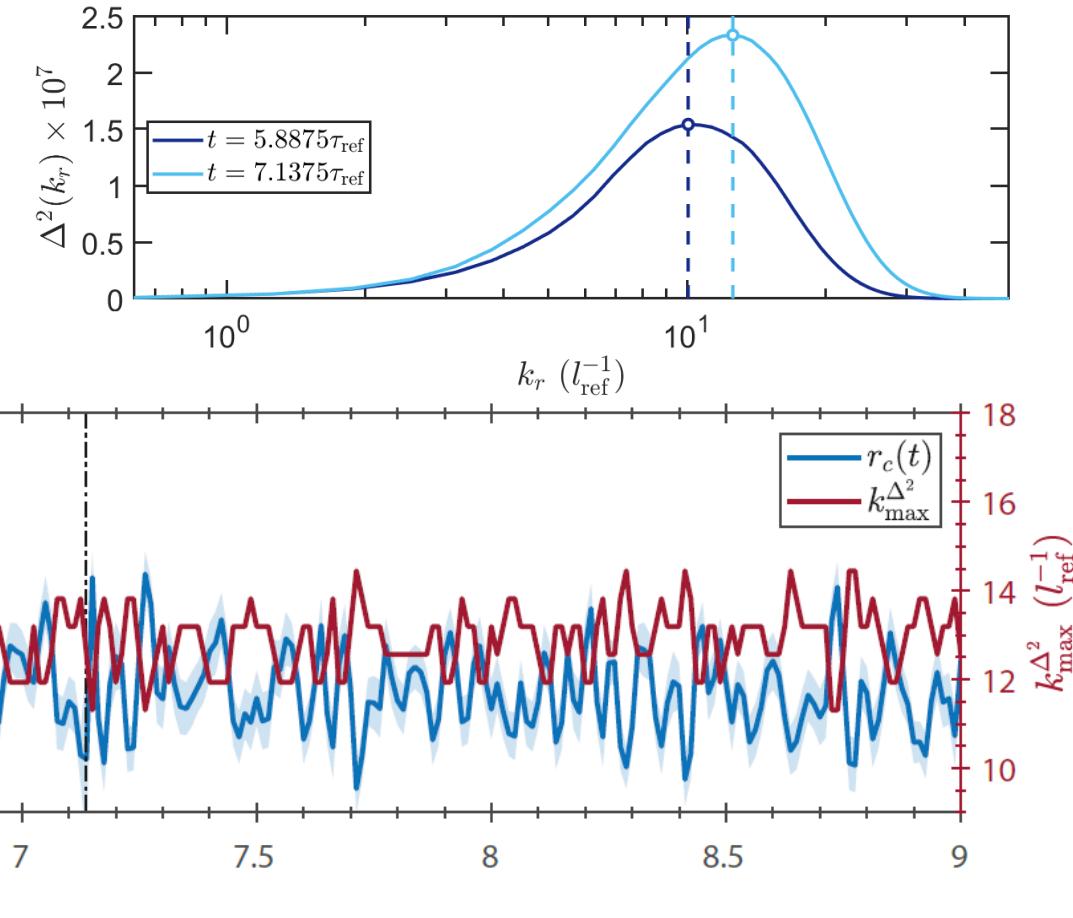
Halo structure: soliton core oscillations

Power spectrum

$$\Delta^2(k_r, t) = \frac{k_r^3}{(2\pi)^2} \int \frac{d\Omega_{\mathbf{k}}}{4\pi k_r^2} P(\mathbf{k}, t)$$

$$P(\mathbf{k}, t) = |\tilde{\eta}(\mathbf{k}, t)|^2$$

$$\eta(\mathbf{r}, t) = \frac{\rho(\mathbf{r}, t) - \bar{\rho}}{\bar{\rho}}$$



$r_c - k_{\max}^{\Delta^2}$ is anti correlated

Halo structures

Mode-mixing and nonlinear dynamic

Lin et al. PRD 97, 103523 (2018)

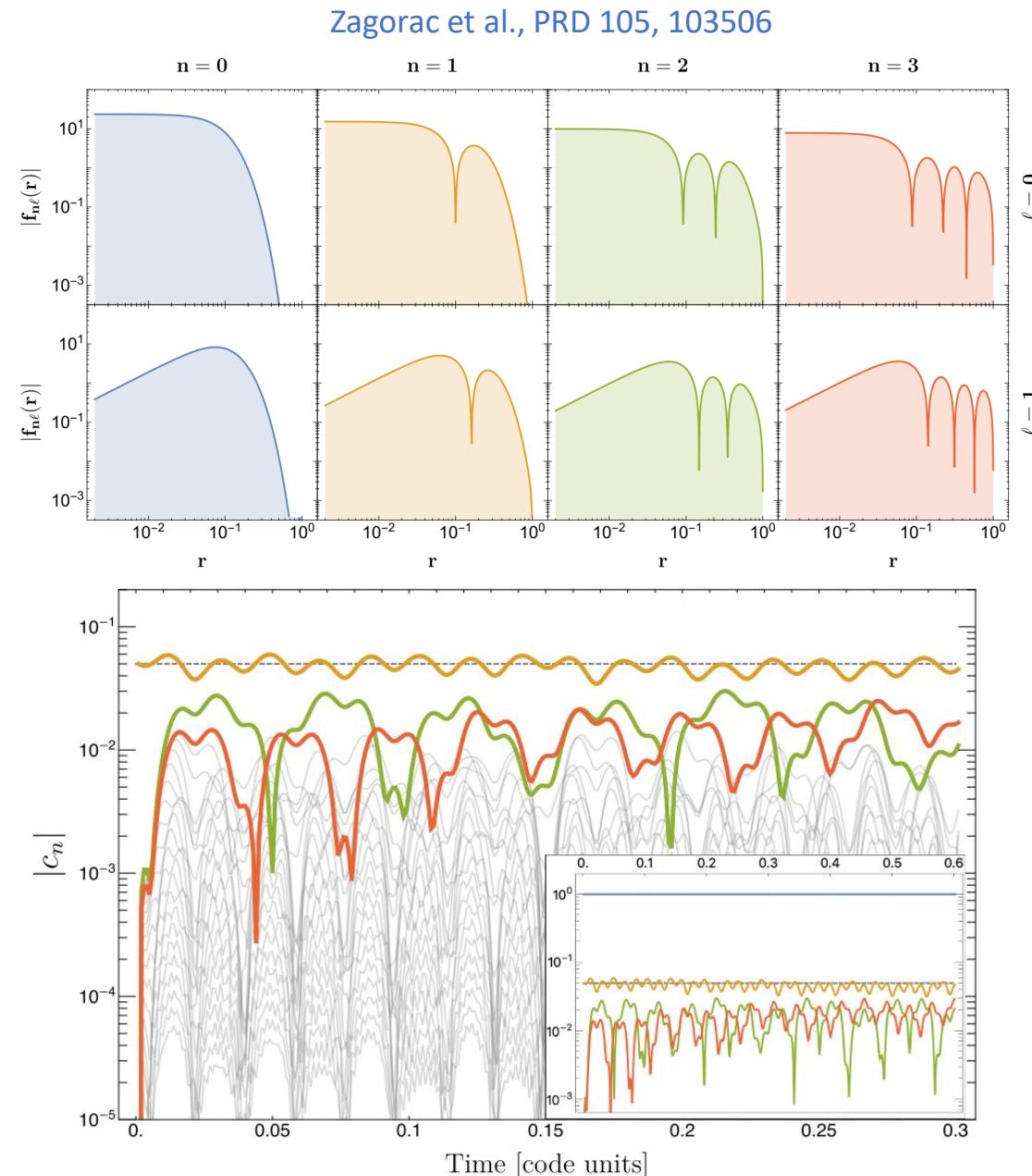
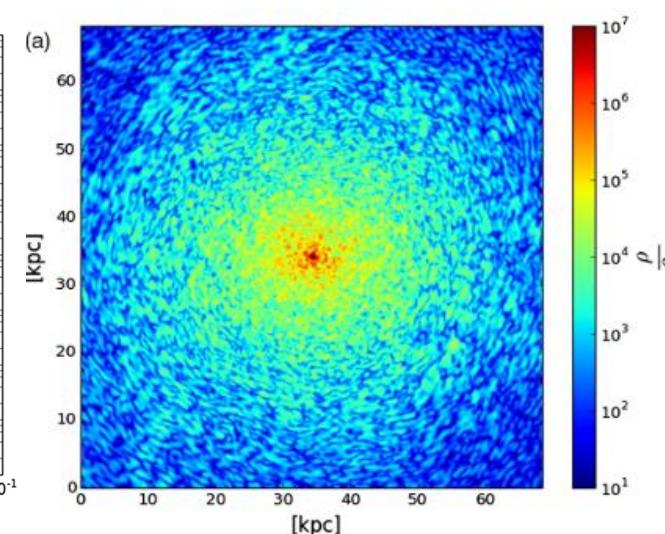
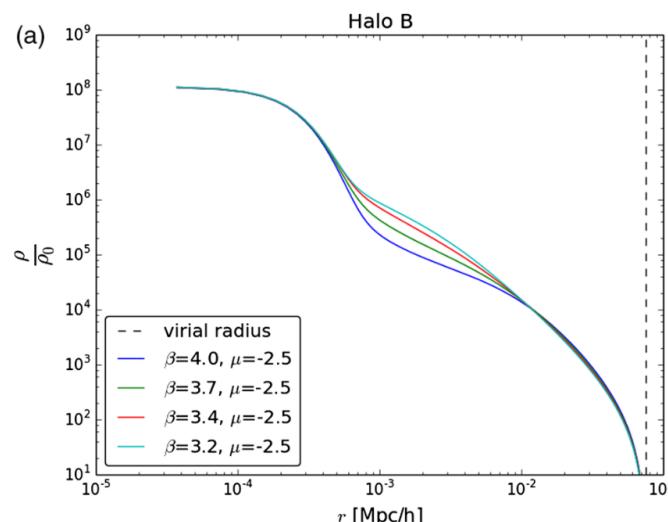
$$\Psi(\mathbf{r}, t) = \sum_{nlm} a_{nlm} e^{iE_{nl}t/\hbar} \phi_{nlm}(\mathbf{r})$$

$$E_{nl} \phi_{nlm}(\mathbf{r}) = \left[-\frac{\hbar^2 \nabla^2}{2m} + m\bar{\Phi}(\mathbf{r}) \right] \phi_{nlm}(\mathbf{r})$$

$|\phi_{000}(\mathbf{r})|^2$ is the soliton profile

Random phase + Spherical-symmetry approximations

$$\nabla^2 \bar{\Phi}(\mathbf{r}) = 4\pi G \bar{\rho}(r) = G \int d\Omega \sum_{nlm} |a_{nlm} \phi_{nlm}(\mathbf{r})|^2$$



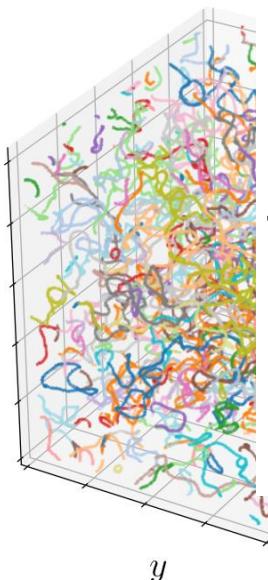
Vortices in FDM halos and granule size

Vortical structure visualization

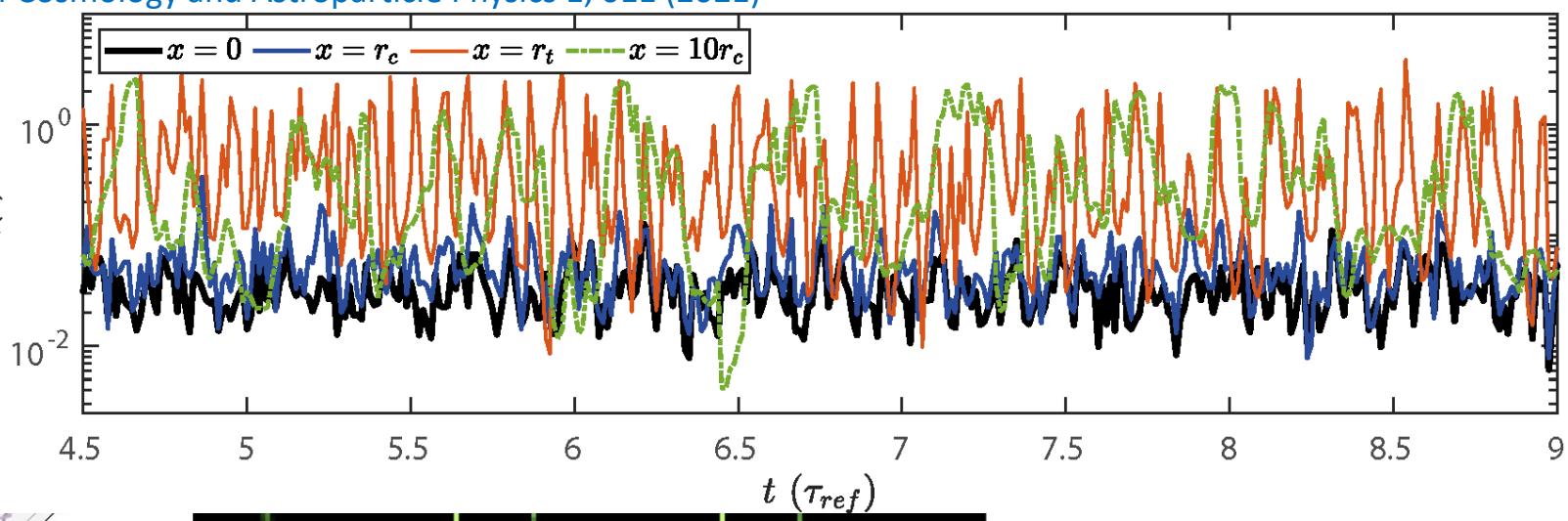
$$|\mathbf{v}^i(\mathbf{r})|^2 \approx \frac{\hbar^2}{m^2 \Delta x^2} \approx 800 \frac{l_{\text{ref}}^2}{\tau_{\text{ref}}^2}$$

- Vortices are in closed loops.
- They are stretching, shrinking, reconnecting and tangling together

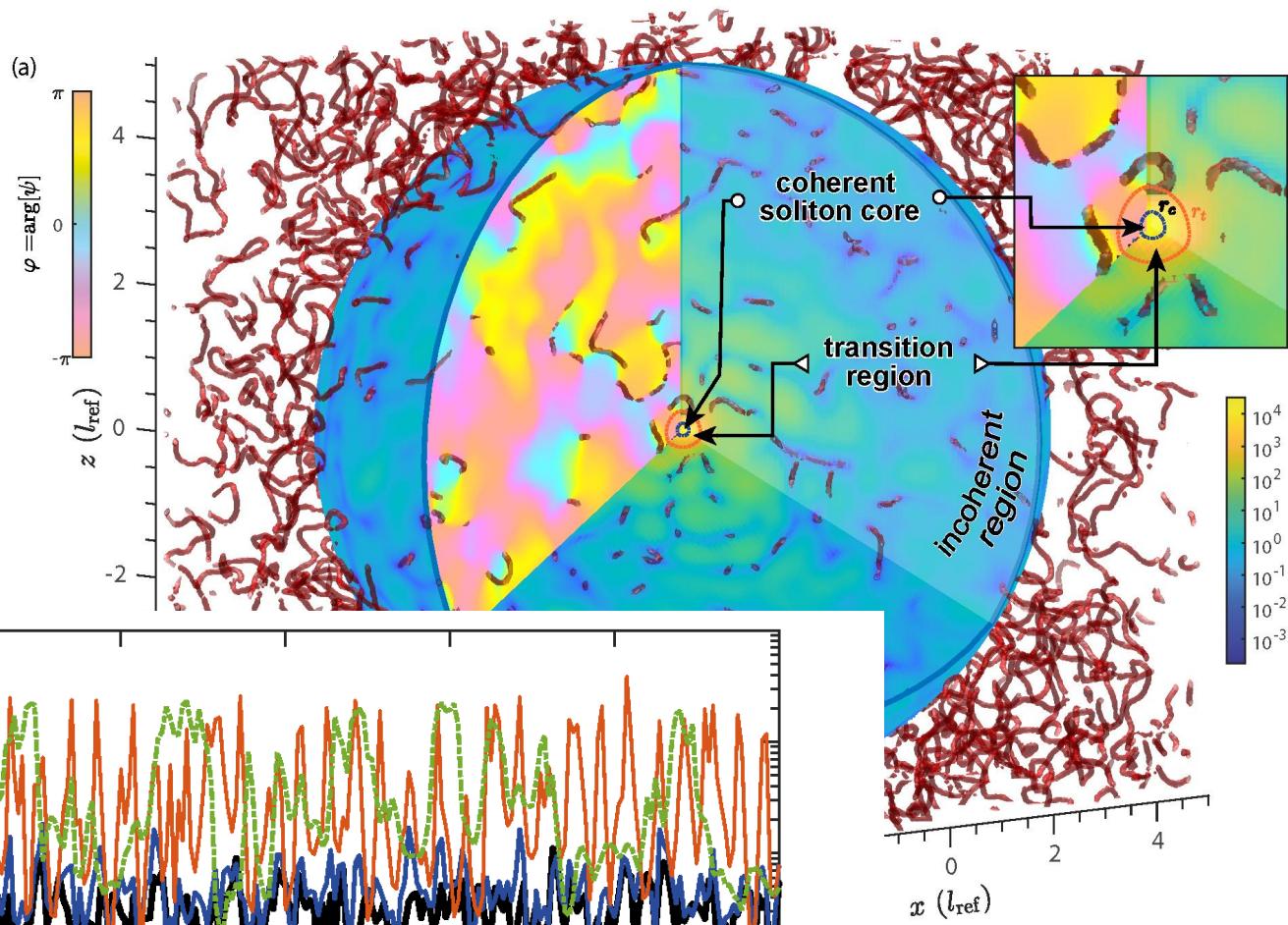
L. Hui et al., Journal of Cosmology and Astroparticle Physics 1, 011 (2021)



$$\zeta_x(t) = \int_{V_c(x)} d\mathbf{r} |\mathbf{v}^i(\mathbf{r})|^2$$



Mocz et al., MNRAS 471, 4559 (2017)



Liu, Proukakis, Rigopoulos, in preparation

Vortices in FDM halos and granule size

Vortex energy and granule power spectra

Incompressible energy spectrum

$$\tilde{\varepsilon}_{\text{ke}}^i(k_r) \equiv \int d\Omega_k k^2 \varepsilon_{\text{ke}}^i(\mathbf{r}, t)$$

Nore et al., Phys. of Fluids 9, 2644 (1997)
 Stagg et al., PRA 94, 053632 (2016)

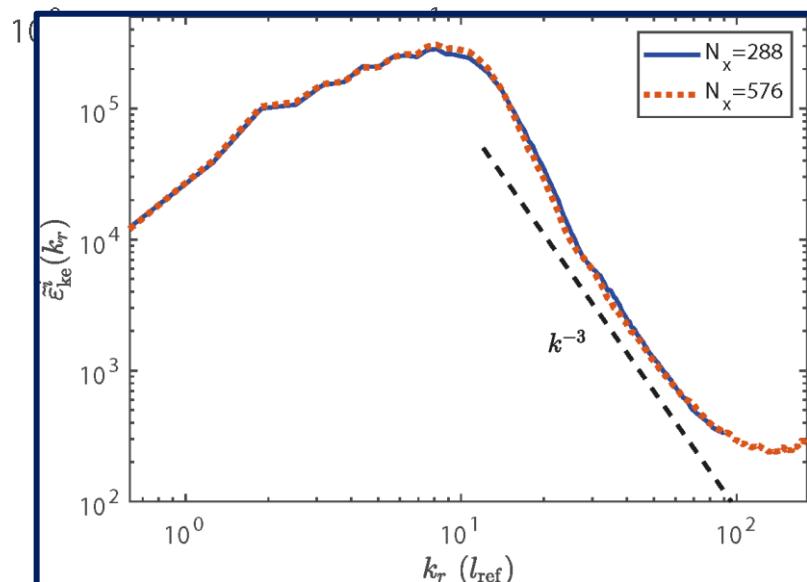
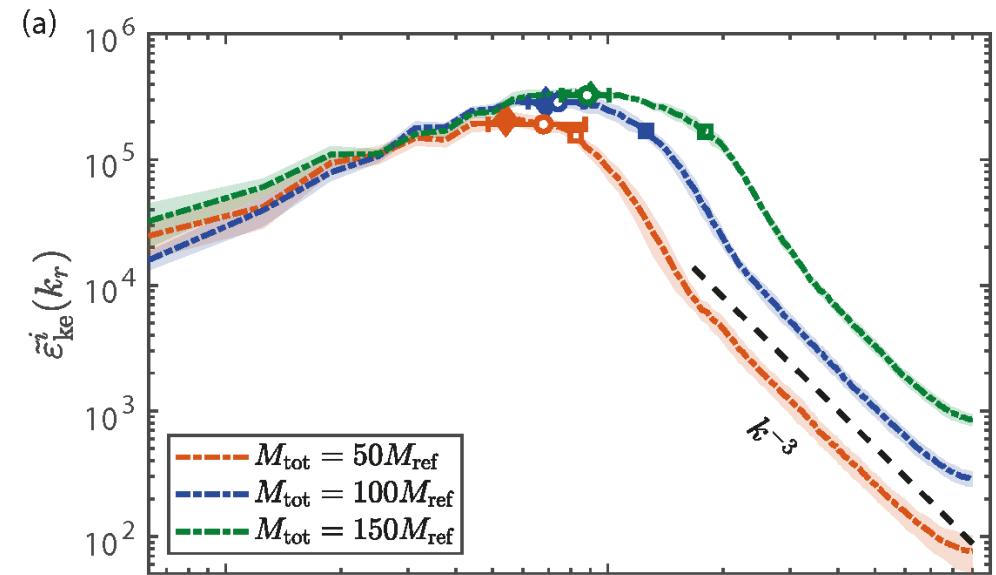
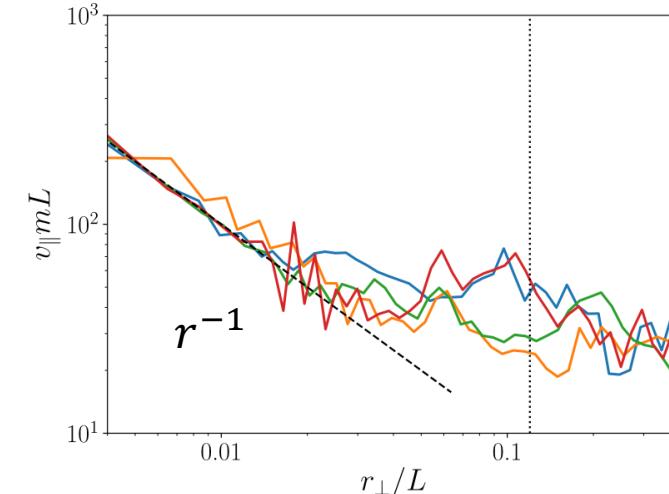
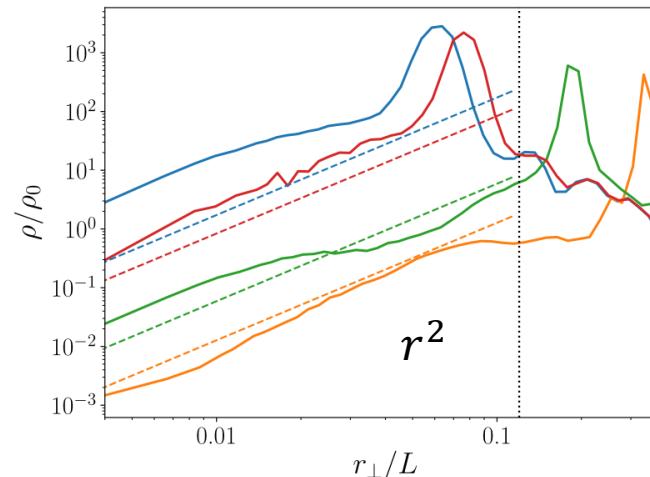
Quasi-classical turbulence, $k^{-5/3}$: Not seen

Mocz et al., MNRAS 471, 4559 (2017)

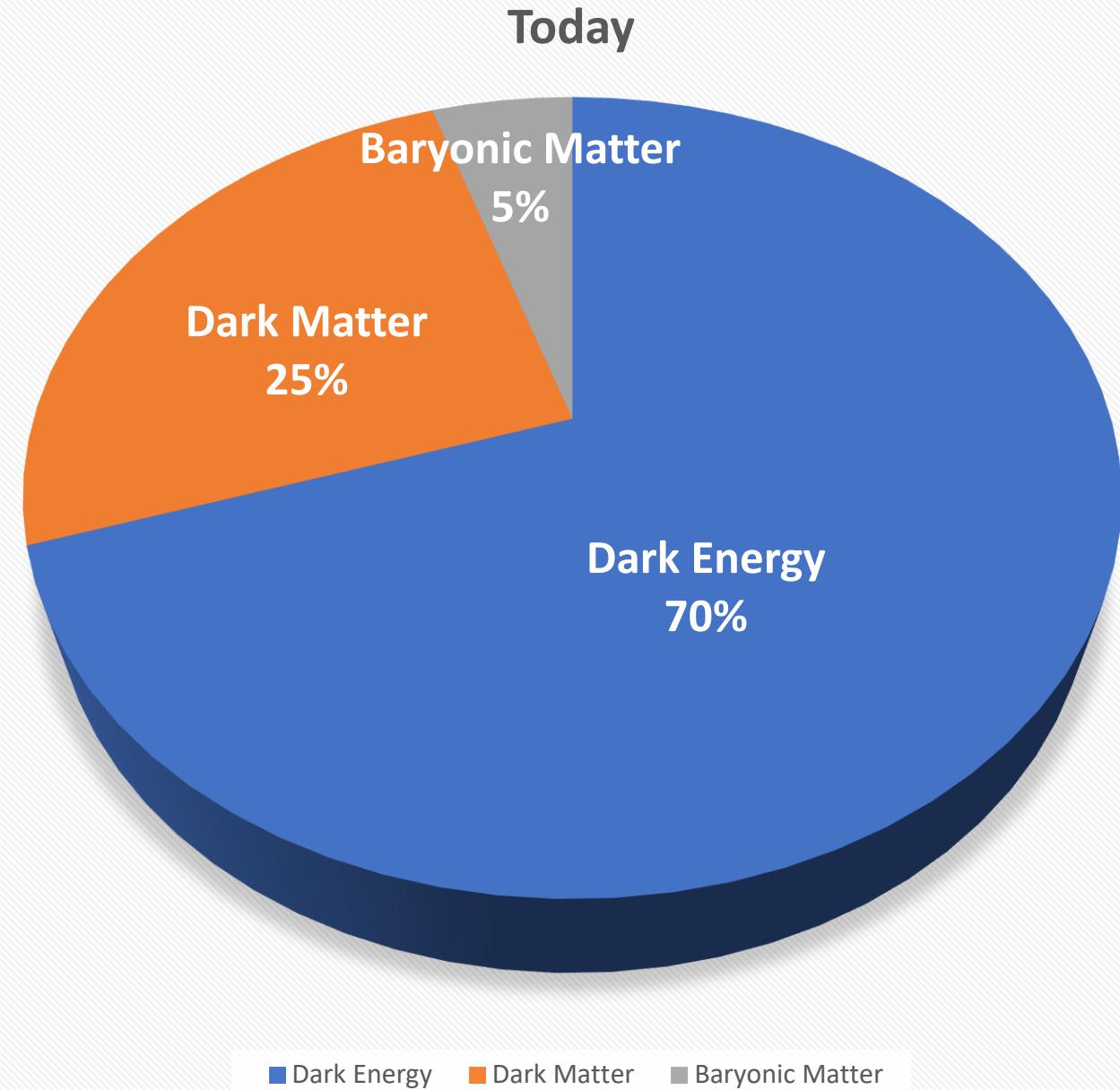
Vortex core structure, k^{-3} : $\rho |\mathbf{v}^i|^2 \approx \text{Const.}$ in large k

L. Hui et al., 2004.01188, JCOPAL (2021)

T. Chiueh et al., J. Phys. B 44, 115101 (2011)

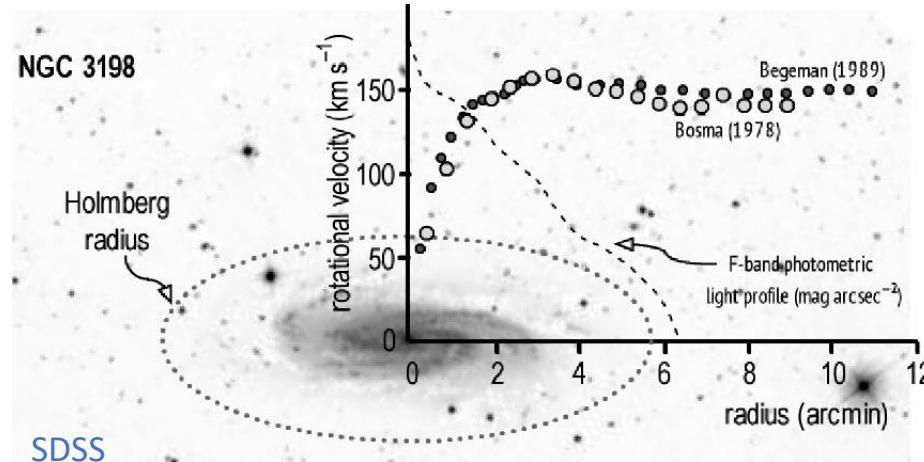


Fuzzy Dark Matter



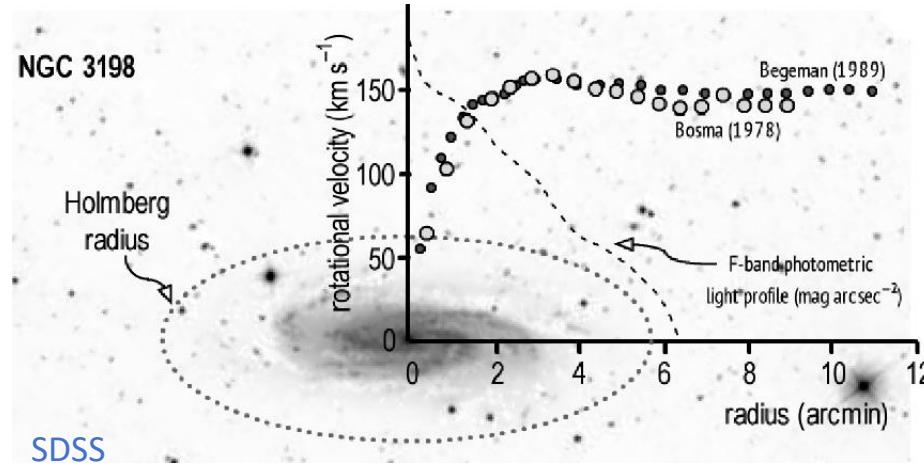
Fuzzy Dark Matter

- Rotational curve

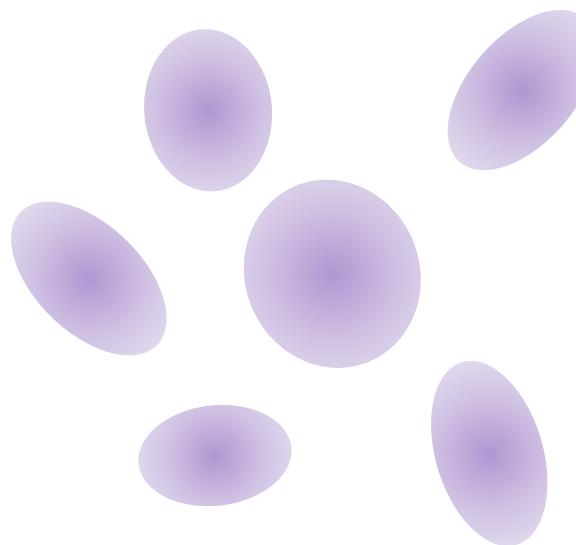


Fuzzy Dark Matter

- Rotational curve

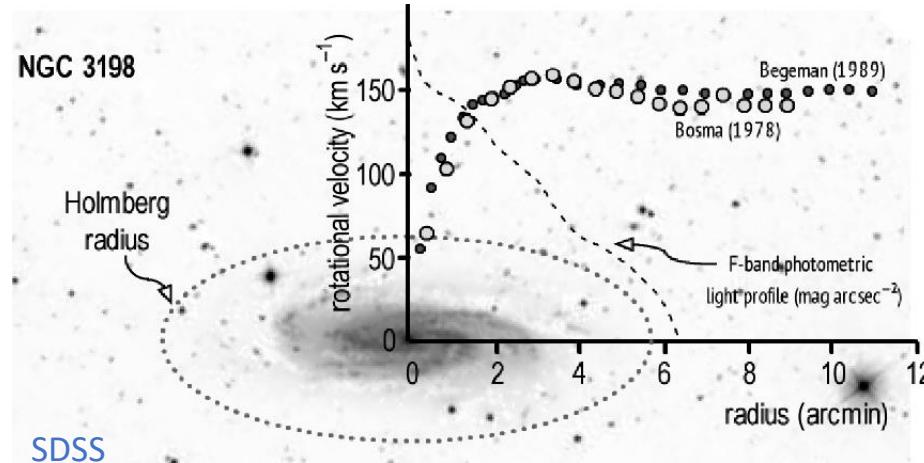


- Galaxy alignment

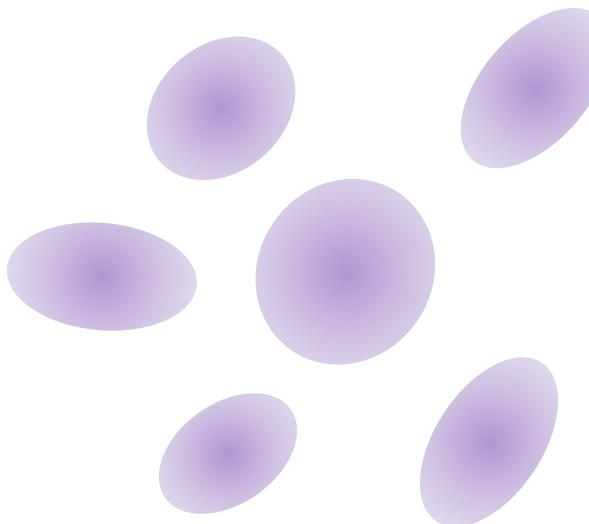


Fuzzy Dark Matter

- Rotational curve

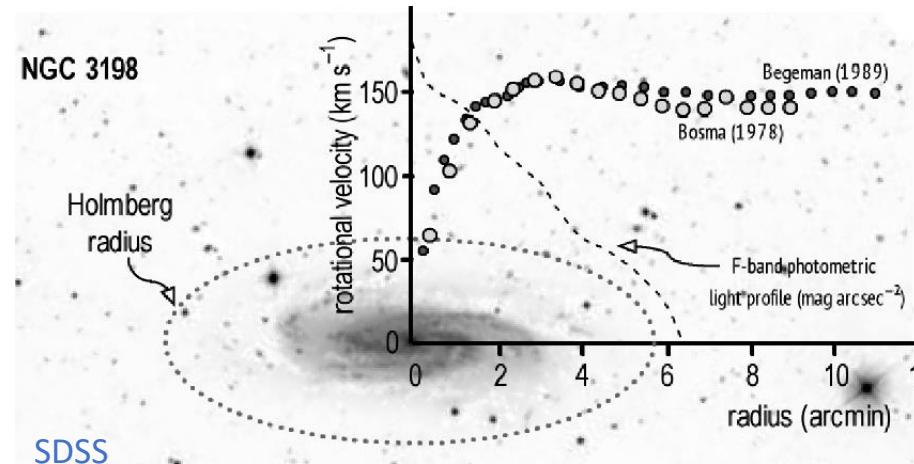


- Galaxy alignment

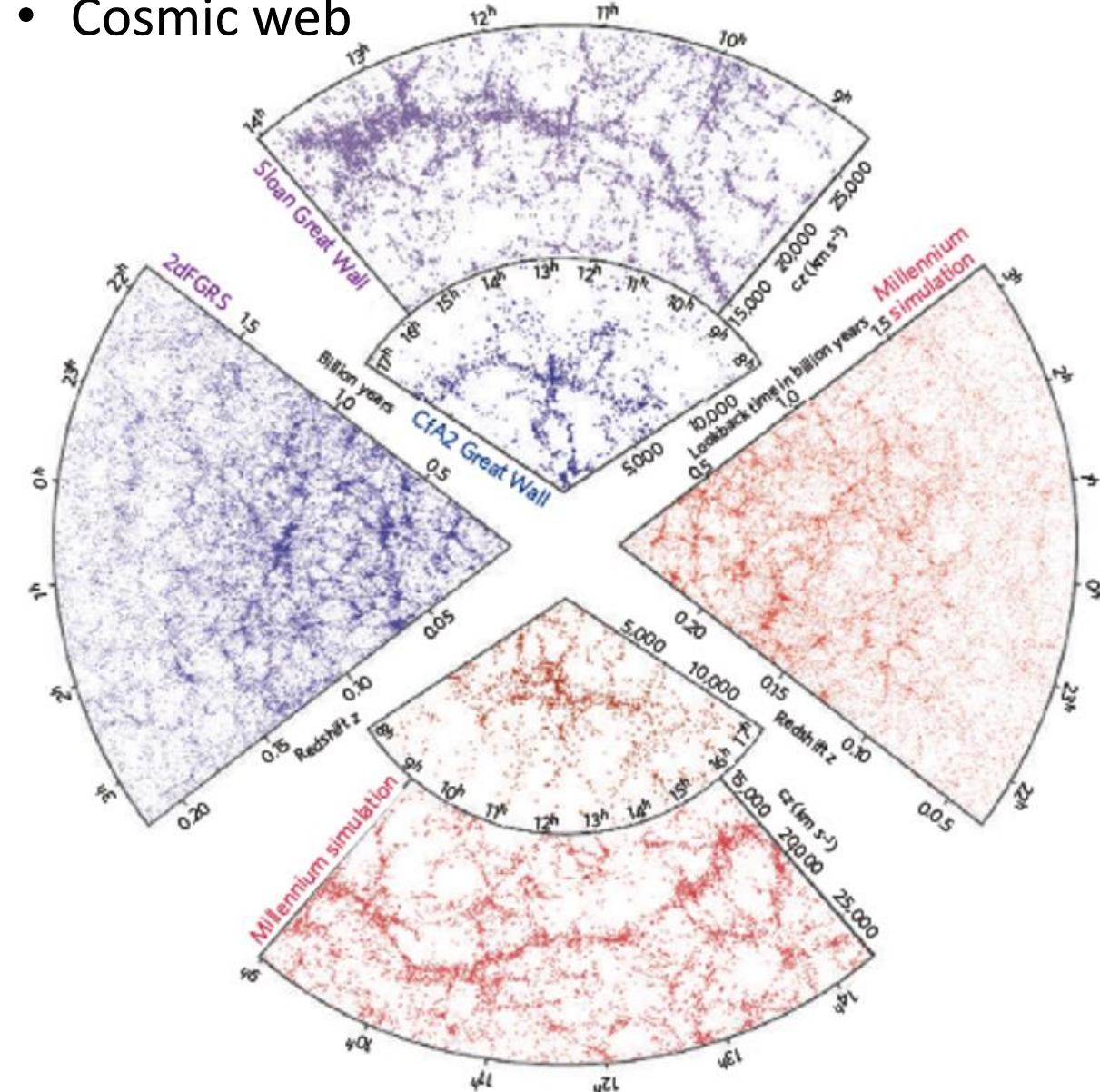


Fuzzy Dark Matter

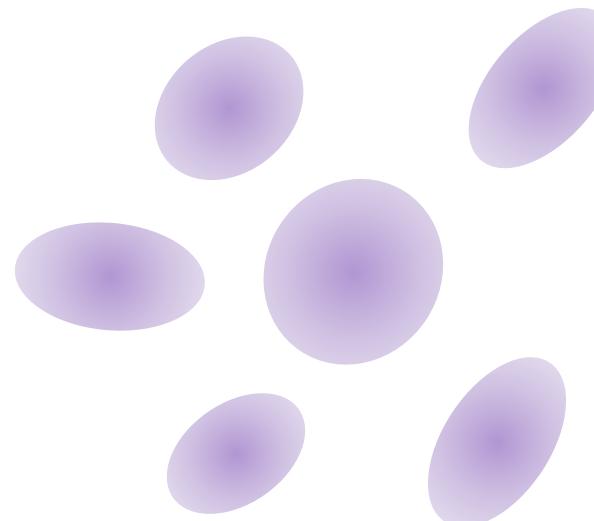
- Rotational curve



- Cosmic web



- Galaxy alignment



Fuzzy Dark Matter

- Cosmic web

