# A plasma-based acceleration method suitable for non-relativistic muons

### Chiara Badiali, Bernardo Malaca, Thales Silva, Ricardo Fonseca, Jorge Vieira

GoLP / Instituto de Plasmas e Fusão Nuclear Instituto Superior Técnico, Lisbon, Portugal

**epp**.tecnico.ulisboa.pt || **golp**.tecnico.ulisboa.pt

EUROPEAN NETWORK FOR NOVEL ACCELERATORS



NPACT supported by EU via I-FAST















What are the best particles to collide?





## Why muons?

 $p^+$ 

What are the best particles to collide?









## Why muons?

### What are the best particles to collide?



 $p^+$ 





Usually the energy available is less than 10% of this value



 $p^+$ 



## Why muons?



Furthermore,  $m_{\mu} = 200 \ m_e$  so that:

Muons have a finite life-time  $(2.2 \ \mu s) \rightarrow$  We have to accelerate them quickly to mitigate muon decay losses, exploiting the time-life dilation in the lab reference frame given by the Lorentz boost ( $\tau' = \gamma \tau$ ) → possible advantage in using plasma based accelerators.



Muons, like electrons, are **fundamental particles**, for this reason their full energy is available in collisions

 $\rightarrow$  a **I4 TeV** muon collider with sufficient enough luminosity would provide similar discovery reach as a **I00 TeV** proton-proton collider.

$$\frac{e^{-}/\mu^{-}}{\Delta E_{e}} = \left(\frac{m_{e}}{m_{\mu}}\right)^{4} \sim 10^{-12}$$







## Muon production: proton beam against a dense target 6 GEANT4

### $\mu^-$ production in the earth's atmosphere



\* J.Allison et al., Nucl. Instrum. Meth.A 835 (2016) 186-225.



### **Muons production on Earth**



Ν<sub>u</sub>

High energetic proton beam against a dense target



A GEANT4\* simulation with a proton beam of  $5 \cdot 10^6$  monoenergetic protons ( $E_p = 450 \text{ GeV}$ ) was performed (in reality ~ $10^{12}$  protons).



## Subluminal spatio-temporal pulses as drivers

### A first step toward the acceleration of non-relativistic particles is having slower drivers ( $v_g < c$ ).

In free space, we can sculpt optical pulses with a modulation of the spatio-temporal degrees of freedom \*.





This allows us to have pulses propagating with a group velocity  $v_g \neq c$ : dω  $= c \tan \theta$ 





## Subluminal spatio-temporal pulses as drivers

### A first step toward the acceleration of non-relativistic particles is having slower drivers ( $v_g < c$ ).









This allows us to have pulses propagating with a group velocity  $v_g \neq c$ :  $d\omega$  $= c \tan \theta$ 





## Subluminal spatio-temporal pulses as drivers

### A first step toward the acceleration of non-relativistic particles is having slower drivers ( $v_g < c$ ).









This allows us to have pulses propagating with a group velocity  $v_g \neq c$ : dω  $- = c \tan \theta$ 





### Toward accelerating drivers

If we assign a finite spatial spectrum instead of a singular frequency of the stationary case, we can make these pulses accelerate with an axially encoded changing velocity\*.



\*M. Yessenov and Y.F.Abouraddy, Phys. Rev. Lett. 125, 244901 (2020).

## Osiris



### **Analytical model**

We model the energy gain of non-relativistic particles using an external field with a time dependent phase velocity. Energy gain:

$$\frac{dp}{dt} = \frac{d}{dt} \left( \frac{\beta_z(t)}{\sqrt{1 - \beta_z^2(t)}} \right) = eE_0 cos[k_p(z(t) - \int \beta_\phi(t) dt)]$$

Imposing the phase-locking condition ( $\beta_z(t) = \beta_\phi(t)$ ), we find:







## Toward accelerating drivers

If we assign a finite spatial spectrum instead of a singular frequency of the stationary case, we can make these pulses accelerate with an axially encoded changing velocity\*.



\*M. Yessenov and Y.F.Abouraddy, Phys. Rev. Lett. 125, 244901 (2020).

## Osiris



### **Analytical model**

We model the energy gain of non-relativistic particles using an external field with a time dependent phase velocity. Energy gain:

$$\frac{dp}{dt} = \frac{d}{dt} \left( \frac{\beta_z(t)}{\sqrt{1 - \beta_z^2(t)}} \right) = eE_0 cos[k_p(z(t) - \int \beta_\phi(t) dt)]$$

Imposing the phase-locking condition ( $\beta_z(t) = \beta_\phi(t)$ ), we find:







## 2D tests in OSIRIS of the use of these accelerating pulses as drivers

 $x_2 [c / \omega_p]$ 

2D simulations have been performed using OSIRIS, testing the acceleration in the quasi-linear regime, with  $a_0 = 0.8$ .













## 2D tests in OSIRIS of the use of these accelerating pulses as drivers J







## **Conclusions & Future Work**

#### Plasma accelerators so far are only applicable to relativistic particles.

To fill this gap, we propose the possibility of accelerating non-relativistic particles using optical wave packets with a group velocity smaller than the speed of light.

Accelerating space time wave packets have been implemented into OSIRIS, and then tested as drivers for the acceleration of non-relativistic muons.

In the future, we will investigate the **non-linear** regime to see if we can improve the energy gain.

### Thank you for your attention! :)







# BACK UP SLIDES







A pulsed plane wave is split into two paths:

$$E(x, z; t) = e^{i(k_0 z - \omega_0 t)} \int z dk_x \tilde{\psi}(k_x) e^{i(k_x x + [k_z - k_0][z - ct \tan \theta])}$$
  
=  $e^{i(k_0 z - \omega_0 t)} \psi(x, z - v_g t).$ 



the ST wave packet is synthesized using a twodimensional pulse shaper formed of a diffraction grating (G), cylindrical lens (L), and spatial light modulator (SLM)

## Testing of the analytical model in 2D using OSIRIS





Moving window at **0.95c** Initial velocity of the muons of **0.9c** 



## Optical space-time wave packets with arbitrary group velocities

The group velocity of an optical pulse can usually be modified in the propagation in a material.

In free space, we can sculpt optical pulses with a modulation of the spatial and temporal degrees of freedom.

Spatio-temporal wave packets: each spatial frequency is uniquely associated with a specific temporal frequency (or wavelength)



Kondakci, H. Esat; Abouraddy, Ayman F., Nature Communications, 10, Article number: 929 (2019) https://doi.org/10.1038/s41467-019-08735-8









## Optical space-time wave packets with arbitrary group velocities

OSIRIS 2D simulations of subluminal space-time wake packets\*



#### \*B. Malaca et al., in preparation







## Plasma ramp for the "acceleration" of the plasma wake











## Plasma ramp for the "acceleration" of the plasma wake



$$z_B = v_g t$$

 $z_A = v_g t - \lambda_p(n_p)$ 

S. Bulanov et al, Phys. Rev. E 58, R5257(R) – Published 1 November (1998), https://doi.org/10.1103/PhysRevE.58.R5257

$$v_B = v_g$$

$$v_{A} = \frac{dz_{A}}{dt} = v_{g} - \frac{d\lambda_{p}(n_{p})}{dt} = v_{g} - \frac{\partial\lambda_{p}}{\partial n_{p}}\frac{\partial n_{p}}{\partial t}$$







## Plasma ramp for the "acceleration" of the plasma wake



$$z_B = v_g t$$

 $z_A = v_g t - \lambda_p(n_p)$ 

This accordion effect results in an acceleration of the back of the plasma wake  $\rightarrow$  it could help us to extend the acceleration distance.

S. Bulanov et al, Phys. Rev. E 58, R5257(R) – Published 1 November (1998), https://doi.org/10.1103/PhysRevE.58.R5257



$$v_B = v_g$$

$$v_{A} = \frac{dz_{A}}{dt} = v_{g} - \frac{d\lambda_{p}(n_{p})}{dt} = v_{g} - \frac{\partial\lambda_{p}}{\partial n_{p}}\frac{\partial n_{p}}{\partial t}$$







## Acknowledgments

Simulation results obtained at MareNostrum (BSC).

## Thank you for your attention











C Badiali | EPP Weekly Meeting | April 8th, 2022