

# Universität Zürich<sup>uzH</sup>

## Is the $(g - 2)_{\mu}$ anomaly a threat to Lepton Flavor Conservation?

## Julie Pagès University of Zurich

G. Isidori, J. Pagès and F. Wilsch, arXiv:<u>2111.13724</u> Based on



La Thuile, 9 March 2022 - Les Rencontres de Physique de la Vallée d'Aoste







## in anomalous magnetic dipole moment of the muon $a_{\mu} \equiv \frac{(g-2)_{\mu}}{2}$





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## Anomalies



## see talk by Massimo







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## Anomalies

see talk by Massimo

see talks by Vitalii, Luca, Claudia





## both suggest

Lepton Flavor Universality Violation (LFUV)

see talk by Andreas

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## Anomalies

What about Lepton Flavor Violation (LFV)?



## $(g-2)_{\mu}$ and LFV





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describe heavy New Physics (NP) effects on low energy observables i.e.  $m_{\rm NP} \gtrsim O(100 \text{ GeV})$ 





## Same operator mediate $(g - 2)_{\mu}$ and $\mu \rightarrow e\gamma$



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 $(g-2)_{\mu}$  and LFV

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$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}} + \sum_{i} \mathscr{C}_{i} \mathscr{O}_{i}$$





Same operator mediate  $(g - 2)_{\mu}$  and  $\mu \rightarrow e\gamma$ 

Tree-level contribution to  $\Delta a_{\mu}$ 

$$\Delta a_{\mu} = \frac{4m_{\mu}v}{e\sqrt{2}} \operatorname{Re} \mathscr{C}'_{e\gamma}$$
22



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$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}} + \sum_{i} \mathscr{C}_{i} \mathscr{O}_{i}$$

Tree-level contribution to  $\mu \rightarrow e\gamma$ 

$$\mathscr{B}(\mu \to e\gamma) = \frac{m_{\mu}^{3}v^{2}}{8\pi\Gamma_{\mu}} \left( \left| \mathscr{C}_{e\gamma}^{\prime} \right|^{2} + \left| \mathscr{C}_{\mu}^{\prime} \right|^{2} \right) \right)$$





## Observables and Constraints on NP

 $(g - 2)_{\mu}$ 

Combined FNAL and BNL result: [hep-ex/0602035, 2104.03281, 2006.04822]  $\Delta a_{\mu} \equiv a_{\mu}^{\exp} - a_{\mu}^{SM} = (251 \pm 59) \times 10^{-11}$ 

Evidence of

$$\operatorname{Re} \mathscr{C}'_{e\gamma} = 1 \times 10^{-5} \operatorname{TeV}^{-2}_{22}$$

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$$\mu \to e \gamma$$



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$$\mu \to e \gamma$$

Branching ratio measured by MEG: [1605.05081]

 $\mathscr{B}(\mu^+ \to e^+ \gamma) < 4.2 \times 10^{-13} \ [90 \% \text{ C} . \text{L}.]$ 

Upper bound on

$$|\mathscr{C}'_{e\gamma}| < 2 \times 10^{-10} \text{ TeV}^{-2}$$



## **Observables and Constraints on NP**

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Upper bound on

$$|\mathscr{C}'_{e\gamma}| < 2 \times 10^{-10} \text{ TeV}^{-2}$$

with strong flavor alignment

$$\epsilon_{12}^{L} \equiv \frac{12}{\mathscr{C}'_{e\gamma}} < 2 \times 10^{-5}$$

$$\frac{12}{22}$$





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Julie Pagès — Is the  $(g - 2)_u$  anomaly a threat to Lepton Flavor Conservation? — La Thuile 09/03/22

From high-scale, 2 sources can spoil alignment at low-scale:

Operators mix through Renormalization Group Evolution Jenkins et al. [1308.2627, 1310.4838, 1312.2014]

$$\mu \frac{d}{d\mu} C_i = \frac{1}{16\pi^2} \beta_i \quad \text{where} \quad \beta_i = \sum_j \gamma_{ij} C_j$$
  
lution 
$$C_i \left(\mu_L\right) = C_i \left(\mu_H\right) + \frac{1}{16\pi^2} \log\left(\frac{\mu_L}{\mu_H}\right) \beta_i$$
$$\underbrace{-\hat{L}}_{-\hat{L}}$$





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Rotation to the mass basis:

$$\Theta_{L(R)}^{\mathscr{Y}} = -\frac{\left[\mathscr{Y}_{e}\right]_{12(21)}}{\left[\mathscr{Y}_{e}\right]_{22}}\Big|_{\mu_{L}}$$

 $-\hat{L}$ 

 $\Rightarrow$  Dipole in the mass basis:

$$\begin{aligned} \mathscr{C}_{e\gamma}'\left(\mu_{L}\right) &= \mathscr{C}_{e\gamma}\left(\mu_{L}\right) + \Theta_{L}^{\mathscr{Y}} \mathscr{C}_{e\gamma}\left(\mu_{L}\right) \\ & \begin{array}{c} 12 \\ \mathscr{C}_{e\gamma}'\left(\mu_{L}\right) \approx \mathscr{C}_{e\gamma}\left(\mu_{L}\right) \\ & \begin{array}{c} 22 \end{array} \end{aligned}$$





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 $\mu \frac{d}{d\mu} C_i = \frac{1}{16\pi^2} \beta_i \qquad \text{where} \quad \beta_i = \sum_j \gamma_{ij} C_j$ with solution  $C_i(\mu_L) = C_i(\mu_H) + \frac{1}{16\pi^2} \log\left(\frac{\mu_L}{\mu_H}\right) \beta_i$ 

Rotation to the mass basis:

$$\Theta_{L(R)}^{\mathscr{Y}} = -\frac{\left[\mathscr{Y}_{e}\right]_{12(21)}}{\left[\mathscr{Y}_{e}\right]_{22}}\Big|_{\mu_{L}}$$

 $\Rightarrow$  Dipole in the mass basis:

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## Operators in the broken phase

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Operators in the broken phase Operators in the unbroken phase

$$\bigcirc e_{rs} = \frac{v}{\sqrt{2}} \bar{e}_{Lr} \sigma^{\mu\nu} e_{Rs} F_{\mu\nu}$$

$$\bigcirc e_{rs} = \frac{v}{\sqrt{2}} \bar{e}_{Lr} \sigma^{\mu\nu} e_{Rs} Z_{\mu\nu}$$

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$$\bigcirc g_{rs} = \frac{v}{\sqrt{2}} \bar{e}_{Lr} e_{Rs}$$

$$\bigcirc g_{he} = \frac{h}{\sqrt{2}} \bar{e}_{Lr} e_{Rs}$$

$$\bigcirc e_{rs} = \frac{v}{\sqrt{2}} \bar{e}_{Lr} e_{Rs}$$

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μν

 $W^{I}_{\mu
u}$ 

<u>(</u>)



Operators in the broken phase Operators in the unbroken phase

$$\bigcirc e_{rs} = \frac{v}{\sqrt{2}} \bar{e}_{Lr} \sigma^{\mu\nu} e_{Rs} F_{\mu\nu}$$

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 $W^{I}_{\mu
u}$ 

weak mixing angle

 $\begin{pmatrix} \mathscr{C}_{e\gamma} \\ \mathscr{C}_{eZ} \end{pmatrix} = \begin{pmatrix} c_{\theta} & -s_{\theta} \\ -s_{\theta} & -c_{\theta} \end{pmatrix} \begin{pmatrix} C_{eB} \\ C_{eW} \end{pmatrix}$ 

 $\begin{pmatrix} \mathscr{Y}_e \\ \mathscr{Y}_{he} \end{pmatrix} = \begin{pmatrix} 1 & -\frac{1}{2} \\ 1 & -\frac{3}{2} \end{pmatrix} \begin{pmatrix} Y_e \\ v^2 C_{eH} \end{pmatrix}$ 



Operators in the broken phase Operators in the unbroken phase

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$$\bigcirc \varphi_{he} = \bar{e}_{Lr} e_{Rs} H(H^{\dagger}H)$$

4-fermions operators for RGE mixing

$$O_{lequ}^{(3)} = \left(\overline{\ell}_{Lp}^{j} \sigma^{\mu\nu} e_{Rr}\right) \epsilon_{prst}$$

$$O_{lequ}^{(1)} = \left(\overline{\ell}_{Lp}^{j} e_{Rr}\right) \epsilon_{jk} \left(\frac{1}{p}\right)$$

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 $\begin{pmatrix} \mathscr{C}_{e\gamma} \\ \mathscr{C}_{e7} \end{pmatrix} = \begin{pmatrix} c_{\theta} & -s_{\theta} \\ -s_{\theta} & -c_{\theta} \end{pmatrix} \begin{pmatrix} C_{eB} \\ C_{eW} \end{pmatrix}$ 

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 $\epsilon_{jk} \left( \bar{q}_{Ls}^k \, \sigma^{\mu\nu} u_{Rt} \right)$ 

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Assumptions:

•  $g_i^2, \lambda \to 0$ 

• 
$$y_{i\neq t} \to 0$$

• 
$$\theta_{eH} = \theta_Y$$

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Operators in the broken phase | Operators in the unbroken phase

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 $\epsilon_{jk} \left( \bar{q}_{Ls}^k \, \sigma^{\mu\nu} u_{Rt} \right)$ 

 $(\bar{q}_{Ls}^k u_{Rt})$ 

Flavor phases defined as  $\theta_X = \frac{12}{C_X}$ Δ 22  $\mu_H$ 



## Contributions to dipole operator



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## Contributions to dipole operator



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## Alignment master formula:

$$\left| \begin{array}{c} \mathcal{C}'_{e\gamma} \\ \epsilon_{12}^{L} \equiv \frac{12}{\mathcal{C}'_{e\gamma}} \\ 22 \end{array} \right|_{\mu_{L}} =$$

$$(\theta_{e\gamma} - \theta_Y) + (\theta_{lequ^{(3)}})$$

with 
$$\Delta_{3} = \frac{-16\hat{L}ey_{t}C_{lequ}^{(3)}(\mu_{H})}{\sum_{22}^{2233}} \text{ and } \Delta_{1} = \frac{-6\hat{L}y_{t}^{3}v^{2}C_{lequ}^{(1)}(\mu_{H})}{[\mathscr{Y}_{e}]_{22}(\mu_{L})}$$

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## Alignment of New Physics







## Alignment master formula:

How can we reach this alignment?

0

0

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## Alignment of New Physics



Dynamical alignments

Flavor symmetries





 $\epsilon_{12}^{L} = (\theta_{e\gamma} - \theta_{Y}) + (\theta_{lequ^{(3)}} - \theta_{e\gamma}) \Delta_{3} + (\theta_{lequ^{(1)}} - \theta_{Y}) \Delta_{1}$ 







 $\epsilon_{12}^{L} = (\theta_{e\gamma} - \theta_{Y}) + (\theta_{lequ^{(3)}} - \theta_{e\gamma}) \Delta_{3} + (\theta_{lequ^{(1)}} - \theta_{Y}) \Delta_{1}$ 







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$$\theta_i - \theta_j = \square^2 \left( \frac{b_j}{a_j} - \frac{b_i}{a_i} \right) < 2 \times 10^{-5} \Rightarrow \text{unnatural turn}$$





$$\theta_i - \theta_j = \square^2 \left( \frac{b_j}{a_j} - \frac{b_i}{a_i} \right) < 2 \times 10^{-5} \Rightarrow \text{unnatural tu}$$



UV theory:

scalar leptoquark  $S_1 \sim (\bar{3}, 1)_{1/3}$ 

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additional Higgs  $\Phi \sim (1,2)_{1/2}$ 













 $U(2)^2$  for Yukawa coupling?

$$\rightarrow \quad \theta_{lequ^{(1)}} = \theta_{lequ^{(3)}}$$

ded top dominant 
$$\rightarrow \theta_{e\gamma} =$$

$$\rightarrow \quad \theta_Y - \theta_{e\gamma} \approx \quad \blacksquare \quad < 2 \times 10^{-5}$$







 $U(2)^2$  for Yukawa co

## Tension in aligning:

 $\theta_{Y} \leftarrow$ 

$$\rightarrow \quad \theta_{lequ^{(1)}} = \theta_{lequ^{(3)}}$$

ded top dominant 
$$\rightarrow \theta_{e\gamma} = \theta_{e\gamma}$$

$$\rightarrow \quad \Theta_{e\gamma} = \Theta_{lequ^{(3)}}$$

$$\rightarrow \quad \theta_Y - \theta_{e\gamma} \approx \quad \blacksquare \quad < 2 \times 10^{-5}$$

$$\rightarrow \theta_{e\gamma} \longleftrightarrow \theta_{lequ^{(3)}}$$





- \* Tight bound on flavor alignment in dipole operator in SMEFT from  $(g 2)_{\mu}$  and  $\mu \rightarrow e\gamma$
- Misaligned NP at high-scale can spoil alignment at low-scale
  - through direct RGE contribution to the dipole  $C_{lease}^{(3)}$ • through rotation to the mass basis  $Y_e$ ,  $C_{leau}^{(1)}$  provided  $(g-2)_{\mu}$  is confirmed

  - $\rightarrow$  impose constraints on some 4-fermion operators
- Flavor symmetries and/or Dynamical mechanism can help explain flavor alignment  $U(2)_{\ell_L} \times U(2)_{e_R}$ •  $\theta_{lequ^{(3)}} = \theta_{lequ^{(1)}}$  from matching & RG dipole and Yukawa •  $\theta_{e\gamma} = \theta_Y$  from matching &  $C_{leau}^{(1)}$ ,  $C_{leau}^{(3)}$  not generated •  $U(1)_{aL_u+bL_\tau}$ •  $\theta_{e\gamma} = \theta_Y = \theta_{lequ^{(3)}} = \theta_{lequ^{(1)}}$  from matching

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## Conclusion

- $\Rightarrow$  If  $(g 2)_{\mu}$  anomaly is a sign of NP, quark sector  $\neq$  lepton sector beyond the SM
  - Thank you for your attention!



## Back-up slides

## Dipole 12 element in mass basis after RGE

RGE for dipole and mass Yukawa

$$\mathcal{C}_{e\gamma}_{ij}(\mu_L) = \left[1 - 3\hat{L}\left(y_t^2 + y_b^2\right)\right]\mathcal{C}_{e\gamma}_{ij}(\mu_H) - \left[16\frac{16}{16}\right]$$
$$[\mathcal{Y}_e]_{ij}(\mu_L) = [Y_e]_{ij}(\mu_H) - \frac{v^2}{2}\mathcal{C}_{eH}_{ij}(\mu_H) + 6v^2\hat{L}$$

LFV Dipole in terms of high-scale quantities

$$\begin{split} \begin{split} \mathcal{C}_{e\gamma}'_{12}(\mu_L) &= (\theta_L^{e\gamma} - \theta_L^Y) \mathcal{C}_{e\gamma}(\mu_L) + (\theta_L^{e\gamma} - \theta_L^{u_3}) (16\hat{L}ey_t) C_{lequ}^{(3)}(\mu_H) \\ &+ \left[ (\theta_L^Y - \theta_L^{u_1}) (6y_t^3) C_{lequ}^{(1)}(\mu_H) + (\theta_L^d - \theta_L^Y) (6y_b^3) C_{ledq}(\mu_H) \right] \frac{1}{[\mathcal{Y}_e]_{22}(\mu_L)} \hat{L}v^2 \mathcal{C}_{e\gamma}(\mu_L) \\ &+ (\theta_L^{eH} - \theta_L^Y) \frac{1 - 9(y_t^2 + y_b^2) \hat{L}}{2} C_{eH}(\mu_H) \frac{1}{[\mathcal{Y}_e]_{22}(\mu_L)} v^2 \mathcal{C}_{e\gamma}(\mu_L) \; . \end{split}$$

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Back-up slides

## Explicit NP model Lagrangian and flavor phases



$$(D^{\mu}S_{1}) - M_{S_{1}}^{2}S_{1}^{\dagger}S_{1} - \left[\lambda_{ilpha}^{L}(ar{q}_{i}^{c}\epsilon\ell_{lpha})S_{1} + \lambda_{ilpha}^{R}(ar{u}_{i}^{c}e_{lpha})S_{1} + ext{h.c.}
ight]^{\mu}\Phi - M_{\Phi}^{2}\Phi^{\dagger}\Phi - \left[\lambda_{lphaeta}^{e}(ar{\ell}_{lpha}e_{eta})\Phi + \lambda_{ij}^{u}(ar{q}_{i}u_{j})ar{\Phi} + ext{h.c.}
ight]^{\mu}$$

$$\frac{+2\lambda_{12}^e\lambda_{33}^uM_{S_1}^2/M_{\Phi}^2}{+2\lambda_{22}^e\lambda_{33}^uM_{S_1}^2/M_{\Phi}^2}$$

$$\frac{\lambda_{\alpha}^{R} \lambda_{i2}^{R} + \lambda_{i1}^{L*} \lambda_{i\alpha}^{L} (Y_{e})_{\alpha 2} - 14 y_{t} \lambda_{31}^{L*} \lambda_{32}^{R}}{\lambda_{\alpha}^{R} \lambda_{i2}^{R} + \lambda_{i2}^{L*} \lambda_{i\alpha}^{L} (Y_{e})_{\alpha 2} - 14 y_{t} \lambda_{32}^{L*} \lambda_{32}^{R}}$$

![](_page_38_Picture_9.jpeg)

Flavor Alignment from  $\tau \rightarrow \mu \gamma$ 

$$\mathcal{B}(\tau^{\pm} \to \mu^{\pm} \gamma) < 4.4 \times 10^{-8} (90\% \text{ CL}) \qquad \Rightarrow \qquad |\mathcal{C}'_{\substack{e\gamma\\23(32)}}$$
  
Flavor alignment in 2-3:  $|\epsilon_{23}^L|, \ |\epsilon_{23}^R| < 1.6 \times 10^{-2} \times \left|\frac{y_{\tau} \, \mathcal{C}'_{e\gamma}}{y_{\mu} \, \mathcal{C}'_{e\gamma}}\right|_{33}$ 

$$\Rightarrow \qquad |\mathcal{C}'_{e\gamma}| < 2.7 \times 10^{-6} \,\mathrm{TeV^{-2}}_{_{23(32)}}$$

![](_page_39_Picture_8.jpeg)