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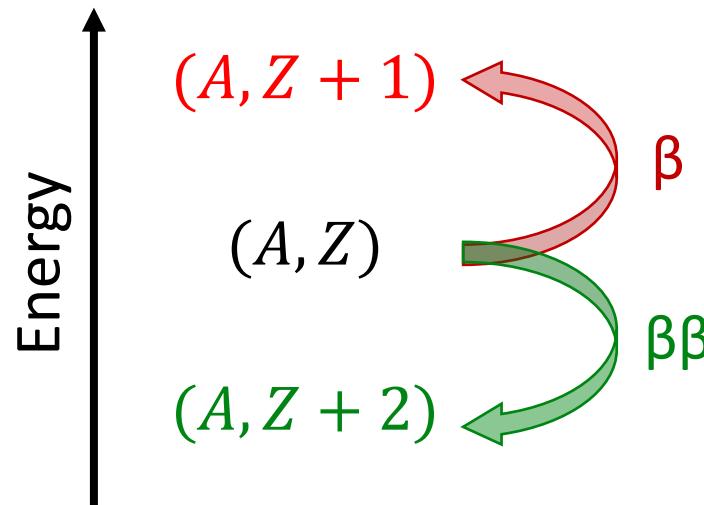


# Final results from GERDA: a neutrinoless double beta decay search

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on behalf of the GERDA collaboration

La Thuile 6-12 March 2022  
Les Rencontres de Physique de la Vallée d'Aoste  
08.03.2022

# Double Beta Decay



- measured in several isotopes:  $^{76}\text{Ge}$ ,  $^{82}\text{Se}$ ,  $^{100}\text{Mo}$ ,  $^{130}\text{Te}$ ,  $^{136}\text{Xe}$ , ...

## OBSERVED

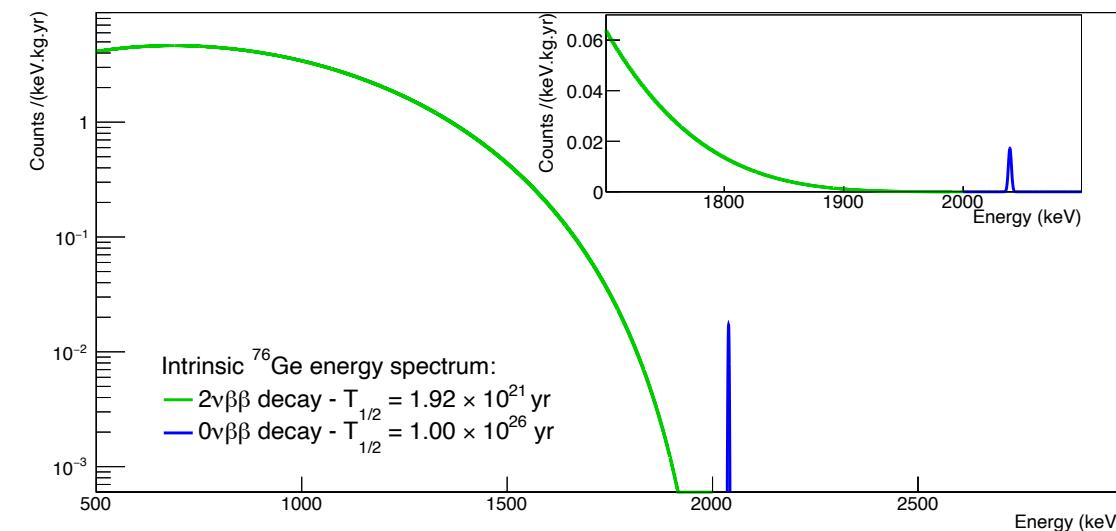
- $2\nu\beta\beta$ :  $^{76}\text{Ge} \rightarrow ^{76}\text{Se} + 2e^- + 2\bar{\nu}_e$
- broad continuous spectrum
- $T_{1/2}^{0\nu} \sim \mathcal{O}(10^{21}) \text{ yr}$
- $\Delta L = 0$

## NOT YET OBSERVED

- $0\nu\beta\beta$ :  $^{76}\text{Ge} \rightarrow ^{76}\text{Se} + 2e^-$
- peak at  $Q_{\beta\beta} = 2039 \text{ keV}$
- $T_{1/2}^{0\nu} > \mathcal{O}(10^{25}) \text{ yr}$
- $\Delta L = 2$

### Search for $0\nu\beta\beta$ informs:

- violation of lepton number conservation
- nature of neutrinos  
(Dirac vs Majorana)
- Neutrino mass scale and ordering  
(normal vs inverted)
- matter-antimatter asymmetry in the universe



# Experimental sensitivities

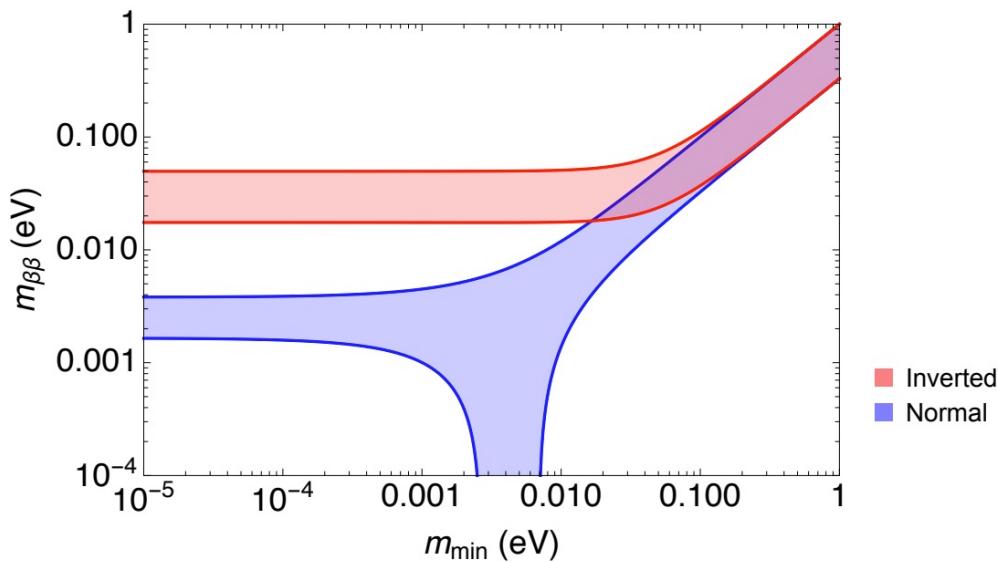


$$n_b \gg 1 \quad T_{1/2}^{0\nu} \propto f e \sqrt{\frac{Mt}{B\sigma_E}}$$

Half-life sensitivity:

(background free)  $n_b < 1$

$$T_{1/2}^{0\nu} \propto f e M t$$



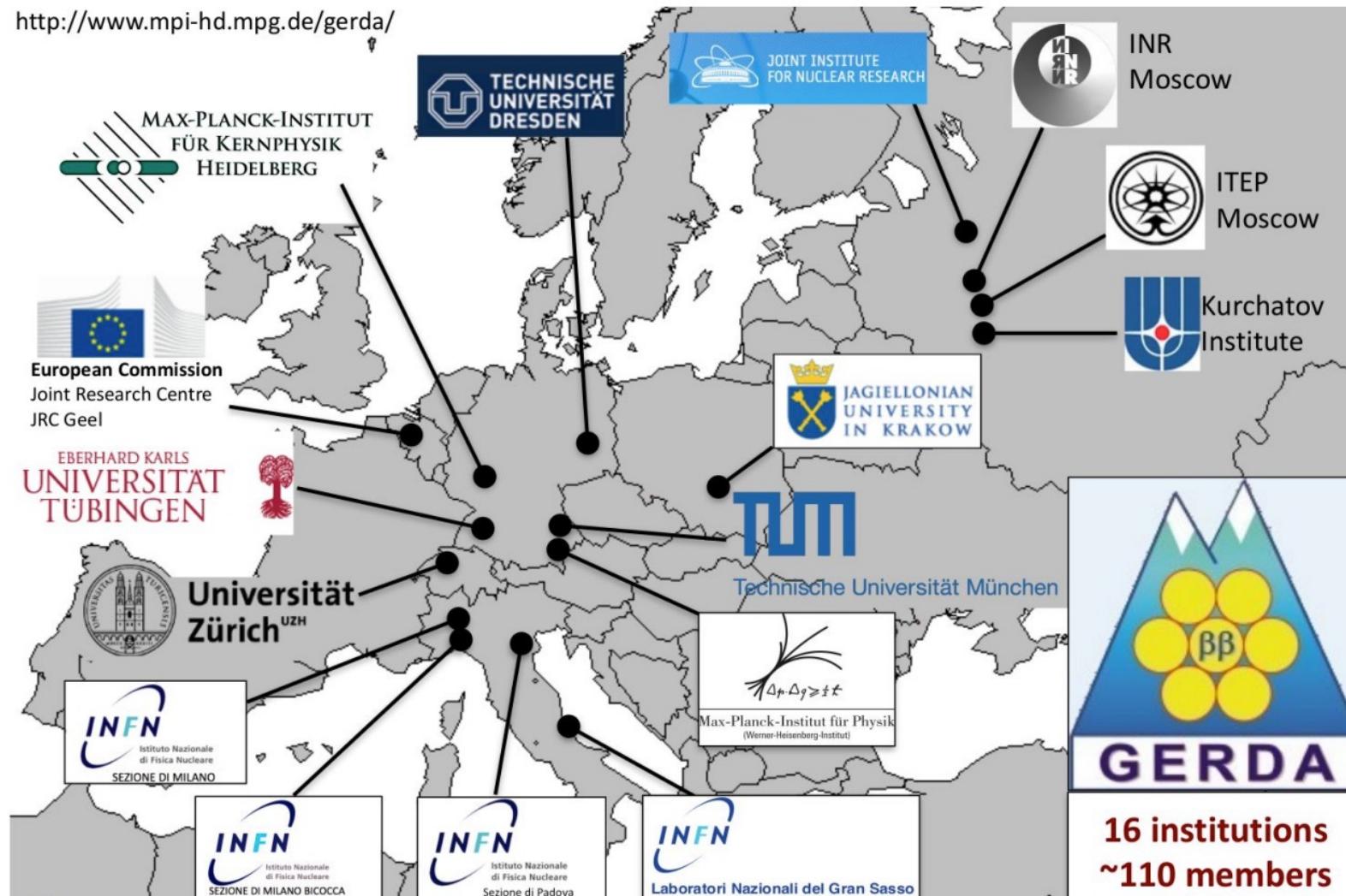
- $f$ : enrichment fraction
  - $e$ : efficiency
  - $M$ : mass
  - $t$ : measurement time
  - $B$ : background index
  - $\sigma_E$ : energy resolution at  $Q_{\beta\beta}$
  - $n_b$  : expected background events
- $\mathcal{E} = Mt$ : exposure

Majorana mass sensitivity:

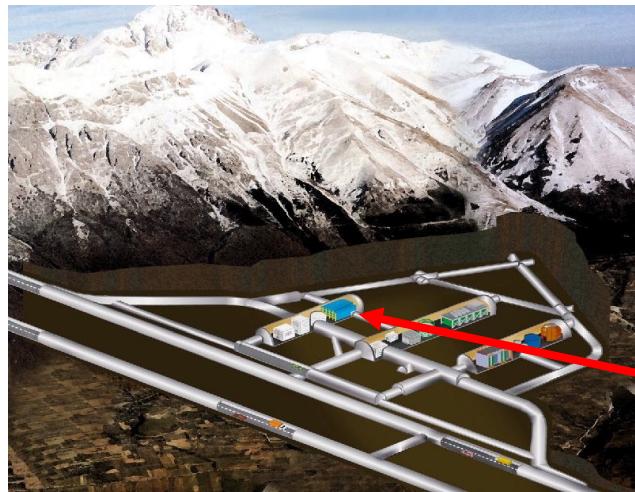
$$\frac{1}{\sqrt{T_{1/2}^{0\nu}}} \propto |\mathcal{M}| \sqrt{G} \frac{\langle m_{\beta\beta} \rangle}{\sqrt{m_e}}$$

- $\langle m_{\beta\beta} \rangle = m_1 |U_{e1}|^2 e^{i\rho} + m_2 |U_{e2}|^2 + m_3 |U_{e3}|^2 e^{i\sigma}$  effective neutrino mass (Majorana)
- $\mathcal{M}$  : matrix nuclear element
- $G$  : phase space factor
- $\phi$  : Majorana phases

# GERmanium Detector Array - GERDA Collaboration

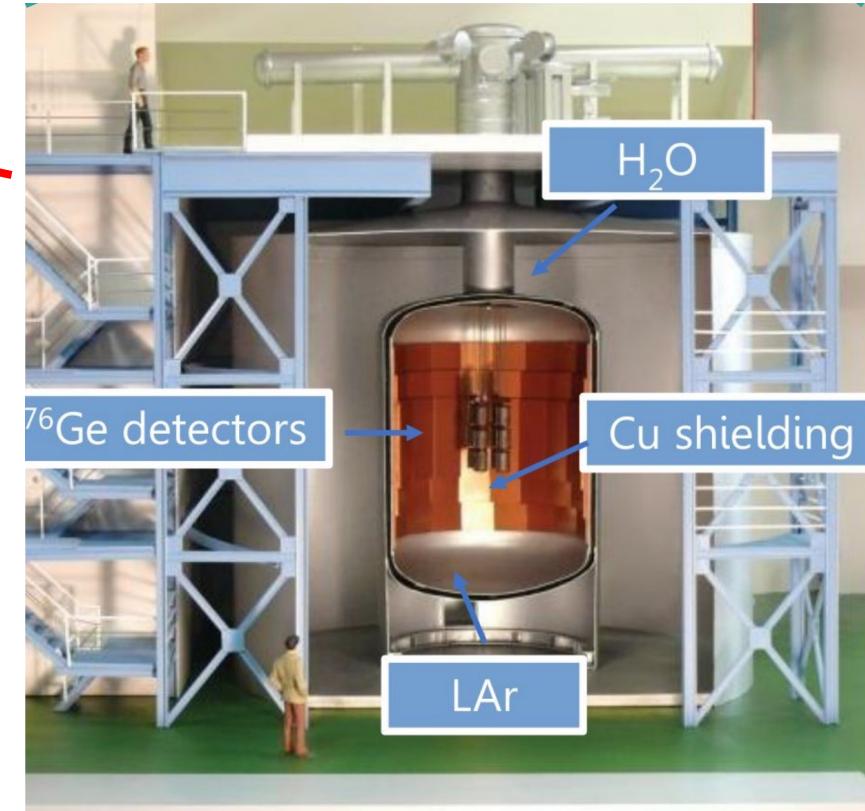


# GERDA Experiment: site and infrastructure



- Laboratori Nazionali del Gran Sasso (LNGS), Italy
- 1400 m rock overburden (3500 mwe)
- Cosmic muon reduction  $\mathcal{O}(10^6)$

- 590 m<sup>3</sup> pure water tank equipped with PMTs:
  - Cherenkov light detection
- cryostat filled with Liquid Argon (LAr):
  - shielding
  - cooling
  - active veto
- Detector Array

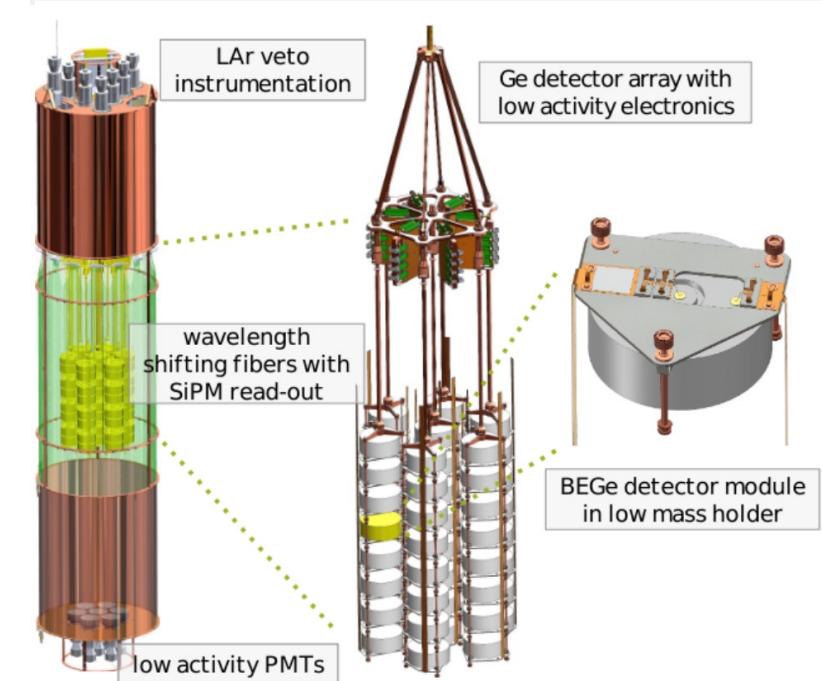


Eur. Phys. J. C 78 388 (2018)

# GERDA Setup



- Germanium is a promising candidate since 1967  
E. Fiorini et al., *Phys Lett B*, 25 (1967), no. 10, 602–603
- Up to 41 detectors in 6 to 7 strings covered by nylon cylinders
- high-purity bare detectors (HPGe) with enriched  $^{76}\text{Ge}$  fraction  
( up to ~87% )  
GERDA, *Astropart.Phys.* 91 (2017) 15-21
  - maximizes detection efficiency: source = detector
  - Excellent energy resolution: ~0.1% FWHM at  $Q_{\beta\beta}$
- lowest background per FWHM energy resolution in the field
- surrounded by fibers coated with the wavelength-shifter TPB (tetraphenyl butadiene)



# Germanium detectors



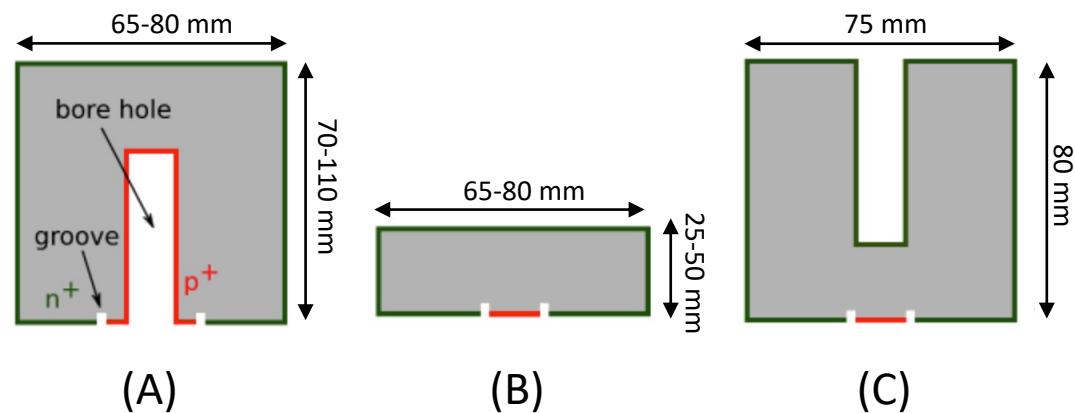
## A. Semi-coaxial (Coax): 6-7

- typical mass 2-3 kg



## B. Broad Energy Germanium (BEGe): 30

- average mass 670 g
- small p+ contact at bottom: good for PSD (Pulse Shape Discrimination)
- excellent energy resolution



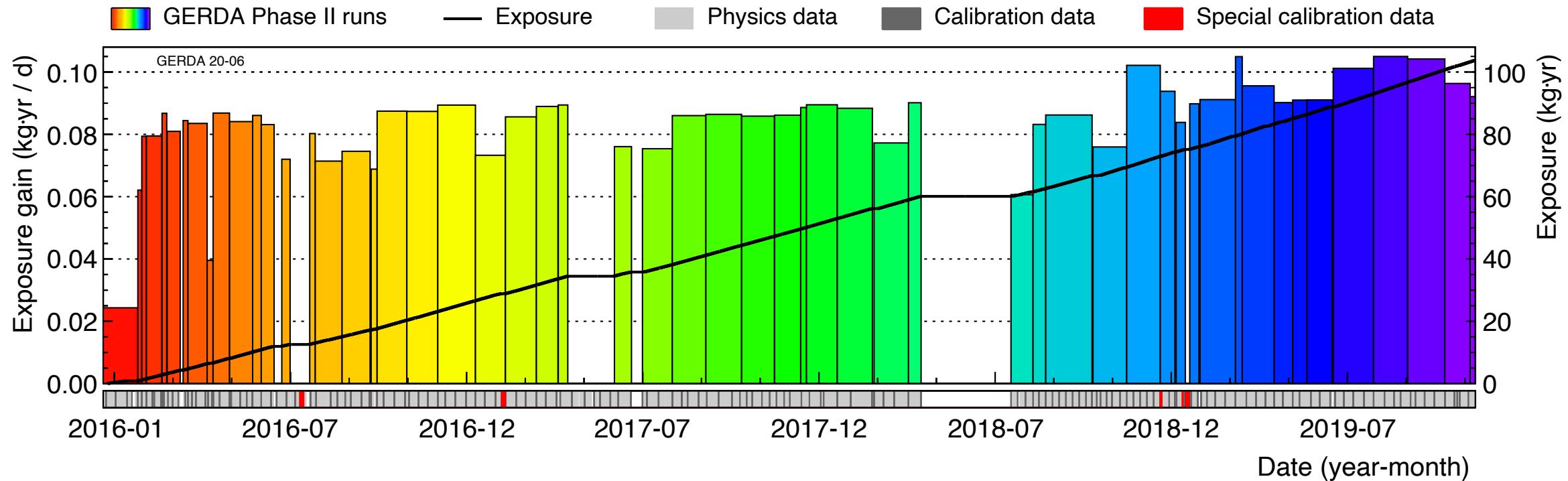
## C. Inverted Coaxial (IC): 5

- Average mass 2 kg
- excellent energy resolution & PSD (like BEGe)

Eur. Phys. J. C. 79 11 978 (2019)

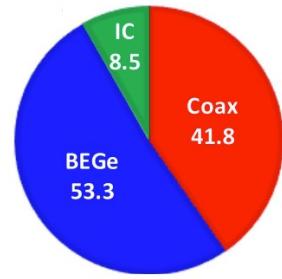
Eur. Phys. J. C, 81 6 505 (2021)

# Data taking

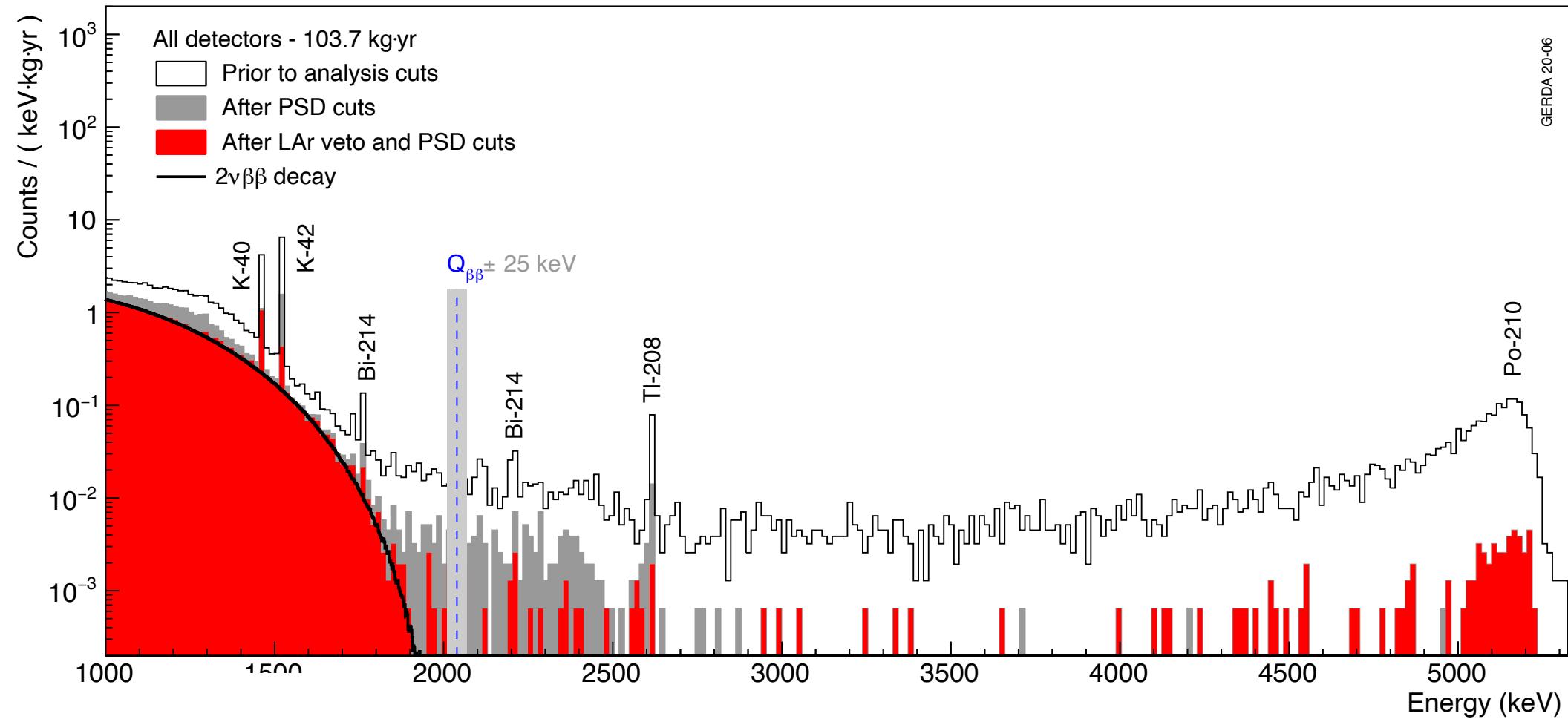


- 2011 – 2013: Phase I, 23.5 kg yr exposure
- 2015 – 2019: Phase II, 103.7 kg yr exposure

- Installation of LAr veto
- 2018: upgrade with 5 IC detectors
- Operation in bkg-free regime



# Background reduction



- < 0.5 MeV:  $^{39}\text{Ar}$

- [0.5 - 2] MeV:  $2\nu\beta\beta$

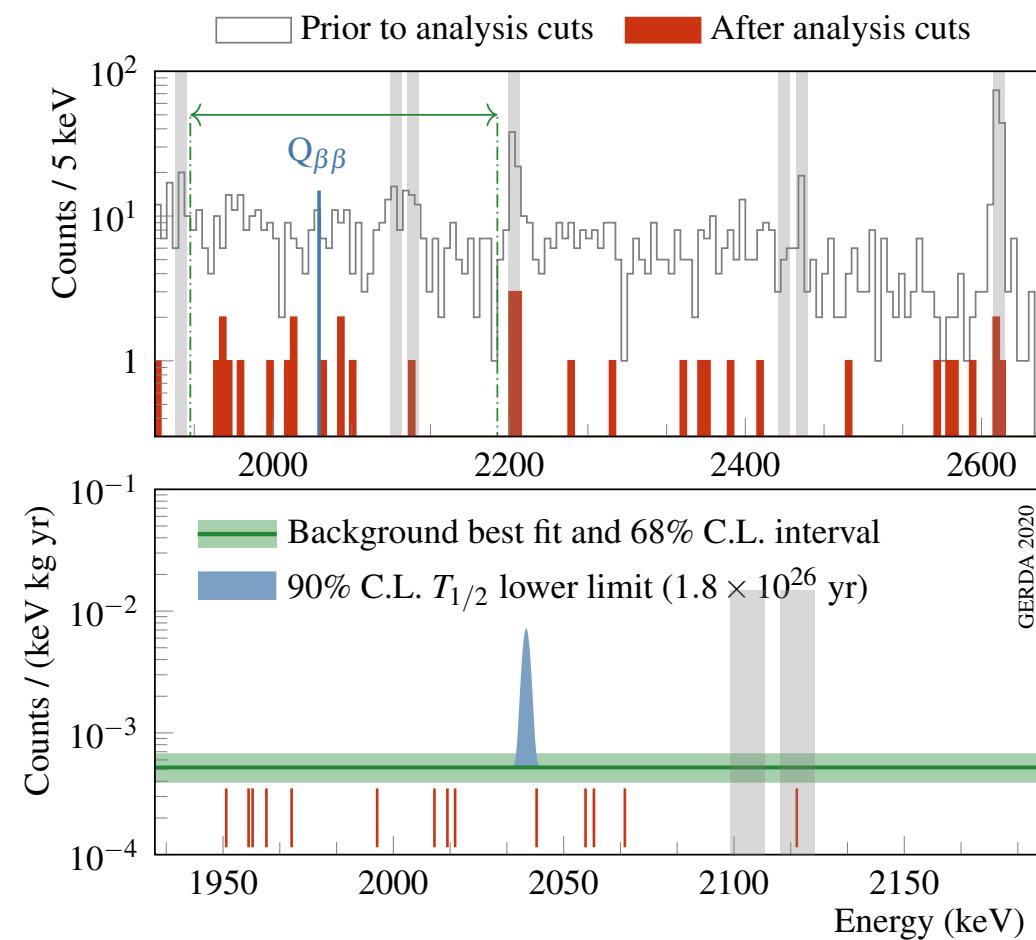
- > 4 MeV:  $\alpha$

# Analysis workflow



The analysis proceeds as follows:

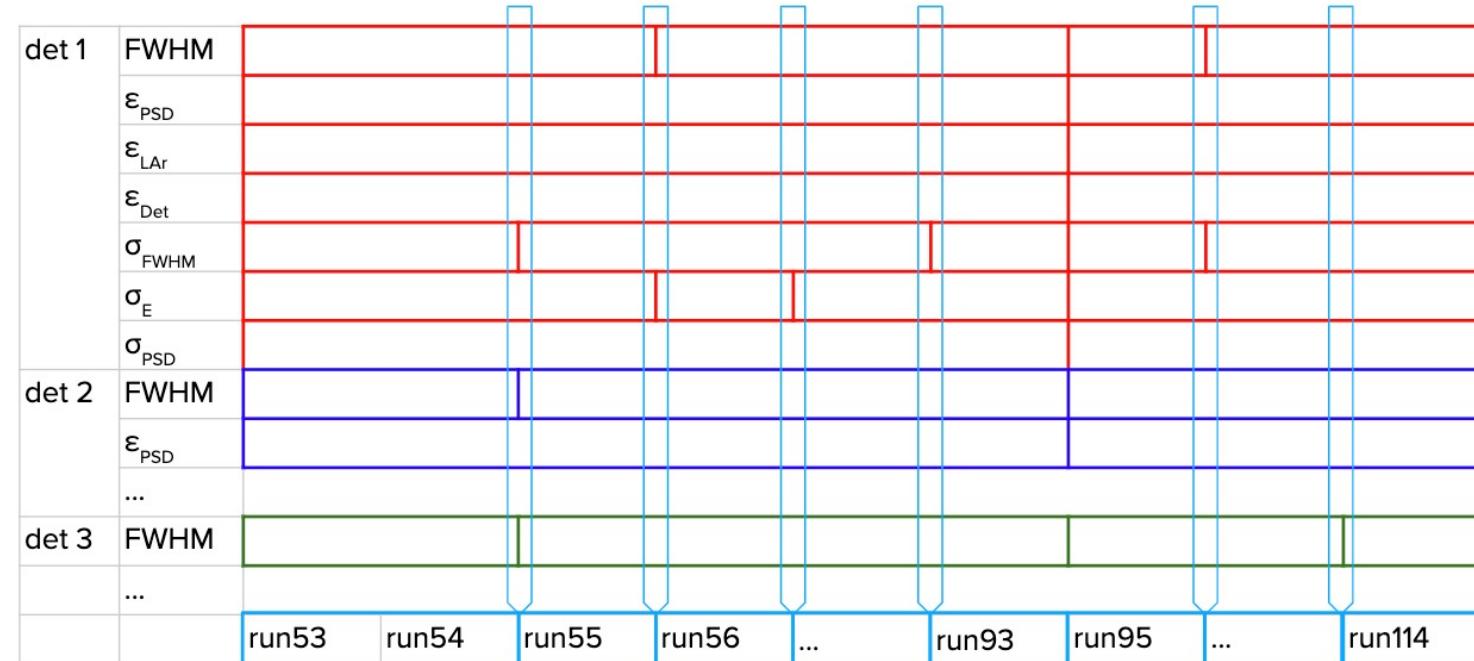
- events in the interval  $Q_{\beta\beta} \pm 25$  keV are not analysed but only stored on disk
- continuous monitoring of detectors
- freezing of analysis procedure and parameters
- blinded events are processed
- data analysis of events detected in the analysis window (1930–2190 keV) excluding the 2 gamma line regions:
  - $2104 \pm 5$  keV :  $^{208}\text{TI}$  (from  $^{232}\text{Th}$  decay chain)
  - $2119 \pm 5$  keV :  $^{214}\text{Bi}$  (from  $^{238}\text{U}$  decay chain)



# Data partitioning



- **Partition:** period of time in which all parameters are constant
- cut data with respect to different detectors
- cut data with respect to time windows that share the same constant parameters
- background indices can be common parameters among partitions
- **Each partition has its own efficiency ( $\epsilon_k$ ), exposure ( $\mathcal{E}_k$ ), energy resolution ( $\sigma_k = \text{FWHM}/2.35$ ) and background index ( $B_k$ )**
- **the result is 383 partitions**

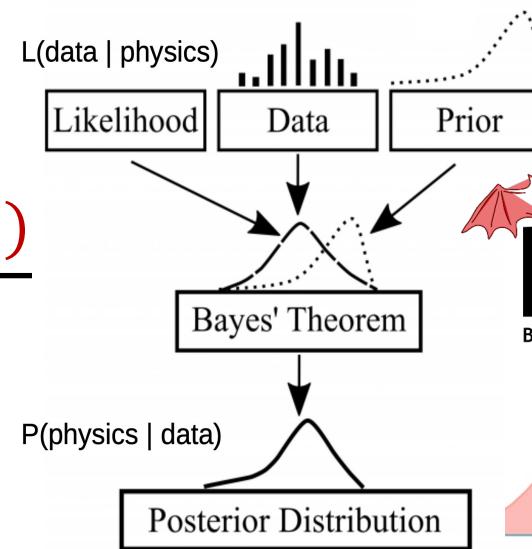


# Bayesian analysis in a nutshell

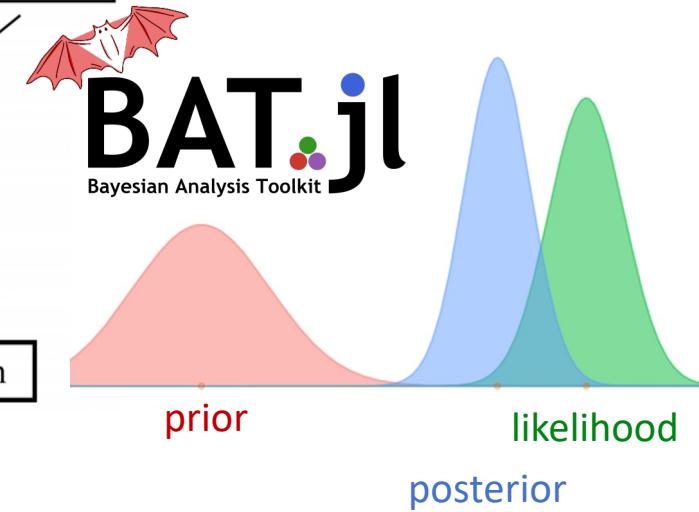


Bayes' theorem:  $P(\theta|D) = \frac{P(D|\theta) \cdot P(\theta)}{P(D)}$

- $D$ : data
- $\theta$ : parameters
- allows to incorporate prior knowledge into computing statistical probabilities
- $p(D) = \int d\theta p(D|\theta) \cdot p(\theta)$ : evidence



<https://github.com/bat/BAT.jl>



## Bayes factor

- For a given model  $M_1$ :  $P(\theta|D, M_1) = \frac{P(D|\theta, M_1) \cdot P(\theta|M_1)}{P(D|M_1)}$

$$BF_{12} = \frac{P(D|M_1)}{P(D|M_2)}$$

# Unbinned Extended Likelihood



$$\mathcal{L} = \prod_k \left[ \frac{(\mu_{s,k} + \mu_{b,k})^{N_k} e^{-(\mu_{s,k} + \mu_{b,k})}}{N_k!} \times \left( \frac{\mu_{b,k}}{\Delta E} + \frac{\mu_{s,k}}{\sqrt{2\pi}\sigma_k} e^{-\frac{(E_i - Q_{\beta\beta})^2}{2\sigma_k^2}} \right) \right]$$

Expected counts Poisson weight  
Flat background Gaussian signal

- $E_i$  is the energy
- GERDA, *Nature* 544 (2017), 47–52
- $\mu_{s,k}$  and  $\mu_{b,k}$  are the expected signal and background counts respectively
- $N_k$  is the number of events in the  $k$ -th partition
- Systematic uncertainties on  $E_i$ ,  $\epsilon_k$  and  $\sigma_k$  are included in the analysis and modelled as normal distributions
- the hypothesis of a flat background is supported by means of a test-statistic derived from Order-Statistic, which models the distribution of spacings between statistical samples arXiv:2008.02048

# Models for the Signal strength



There are 2 different priors on the signal strength  $S = \frac{1}{T_{1/2}^{0\nu}}$  (which ranges from 0 to  $10^{-24}$  1/yr):

- $p(S) \sim \text{Uniform}$ 
  - equiprobable signal strengths
- $p(S) \sim \frac{1}{\sqrt{S}}$ 
  - equiprobable Majorana neutrino masses  $m_{\beta\beta}$
  - $S \propto m_{\beta\beta}^2$

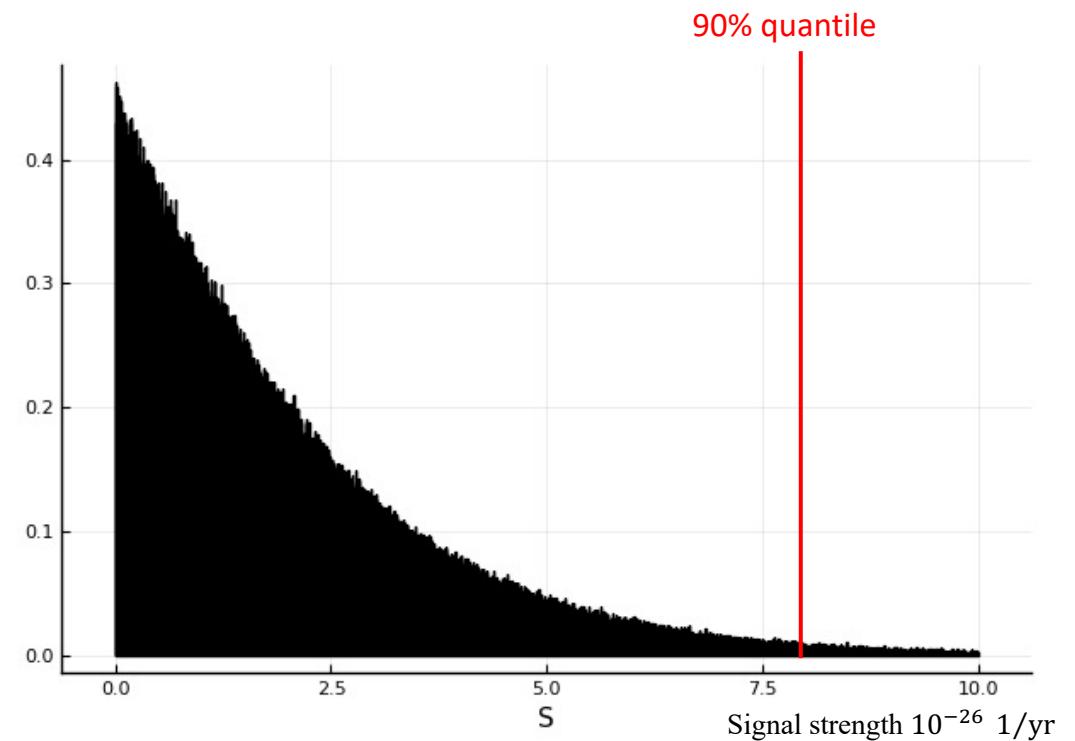
1) Perform fit to Phase I data

- 61 events, 23.5 kg · yr exposure

2) Feed posterior from Phase-I to Phase-II analysis

- 13 events, 103.7 kg · yr exposure

3) Get limit at 90% C. I. from posterior distribution of  $S$



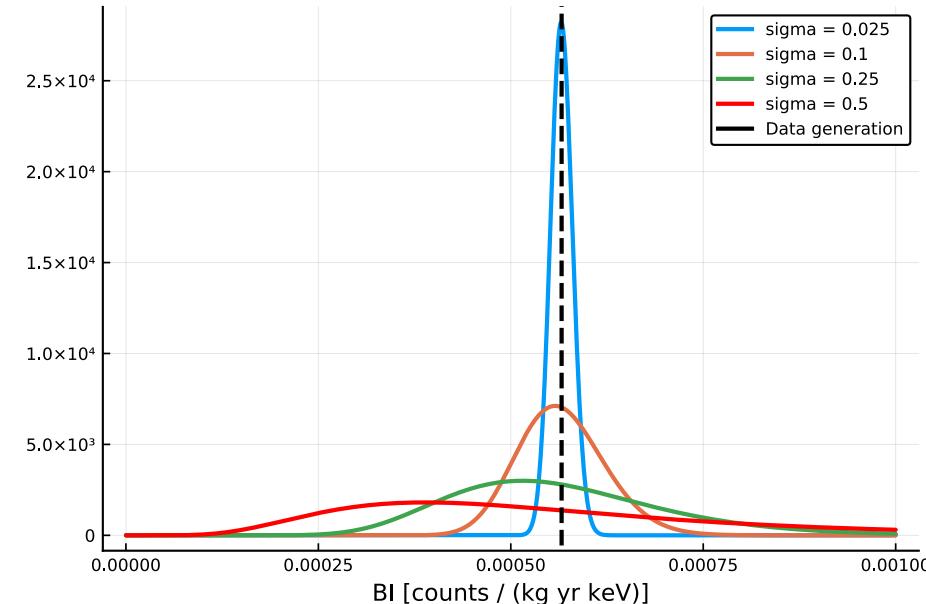
# Models for the Background Indices



There are 3 types of detectors: BEGe (Broad Energy Germanium), Coax (Coaxial), Inv-Coax (Inverted Coaxial)

The **background index ( $B$ )** can be treated in 3 different ways, which gives rise to 3 different models:

- **Single background index:** there is only one background index for all detector types:  $B \sim Uniform$
- **Uncorrelated background indices:** each detector type has its own independent  $B_i$ :  $B_i \sim Uniform$
- **Correlated background indices:** each detector type has a different  $B_i$  but they are all correlated. This implies a hierarchical model
  - $\sigma_B \sim Uniform$
  - $m_B \sim Uniform$
  - $B_i \sim LogNormal \left( \ln(m_B) - \frac{\sigma_B^2}{2}, \sigma_B \right)$
- Changing the range of  $\sigma_B$  allows the correlated model to replicate the previous two models: **smooth change**
  - **Small sigma** ---> Single BI
  - **Large sigma** ---> Uncorrelated BI

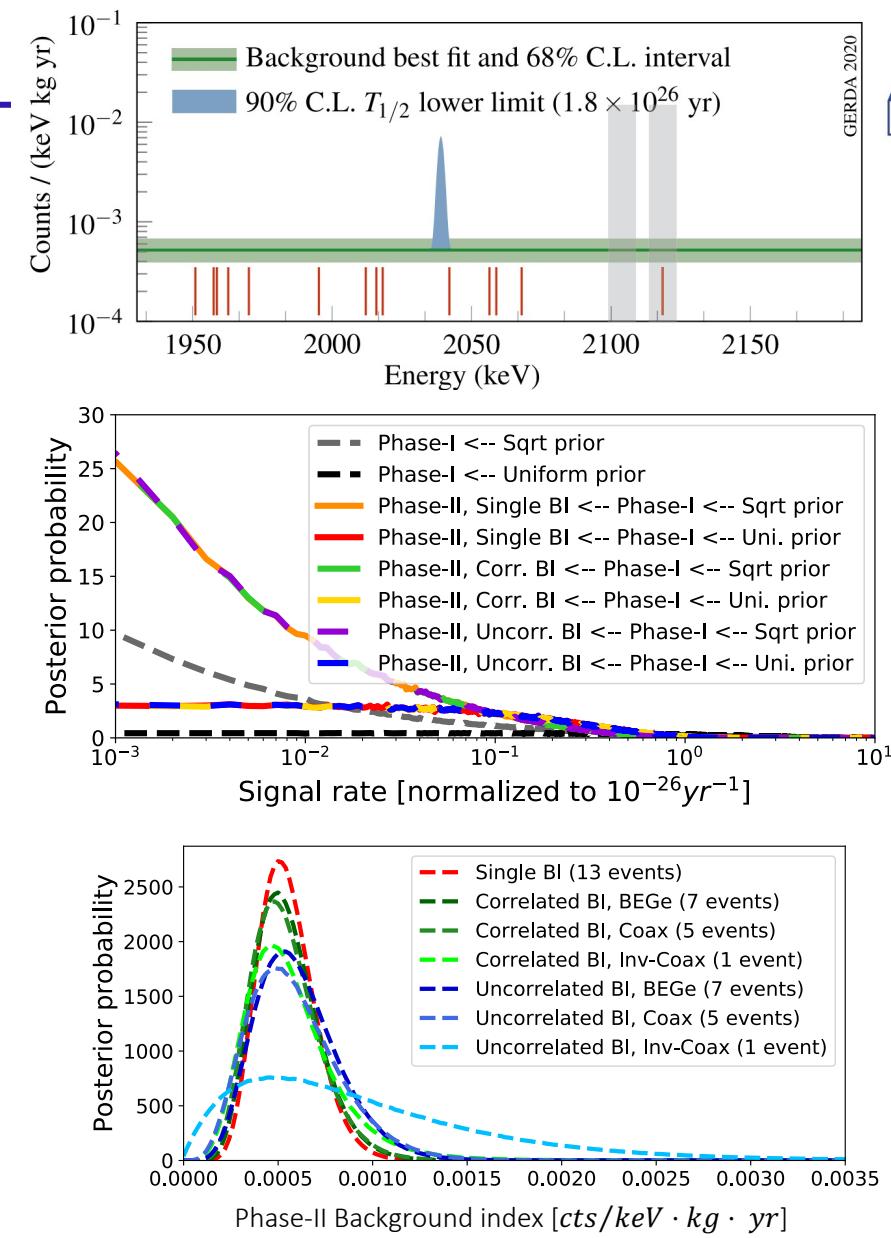


# Data and Results.

PRL 125, 252502 (2020)



- In the analysis window we detect 13 events
- After the analysis we cannot claim a signal
- Half-life limit (and sensitivities) extracted is the same for all bkg models (Bayesian): SN COMPUT. SCI. 2, 210 (2021)
  - Uniform prior:  $T_{1/2}^{0\nu} > 1.4 \cdot 10^{26} \text{ yr}$  (90% C.I.)
  - $1/\sqrt{S}$  prior:  $T_{1/2}^{0\nu} > 2.3 \cdot 10^{26} \text{ yr}$  (90% C.I.)
- Frequentist limit (and sensitivity):
  - $T_{1/2}^{0\nu} > 1.8 \cdot 10^{26} \text{ yr}$  (90% C.L.)
- Limits on effective neutrino mass:
  - $|m_{\beta\beta}| < [79 - 180] \text{ meV}$
- Background index for Phase-II analysis of Single B model is:
 
$$(\text{Phase-II}) \quad B = 5.2_{-1.3}^{+1.6} \cdot 10^{-4} \left[ \frac{\text{cts}}{\text{keV} \cdot \text{kg} \cdot \text{yr}} \right] \text{ (68\% SI)}$$

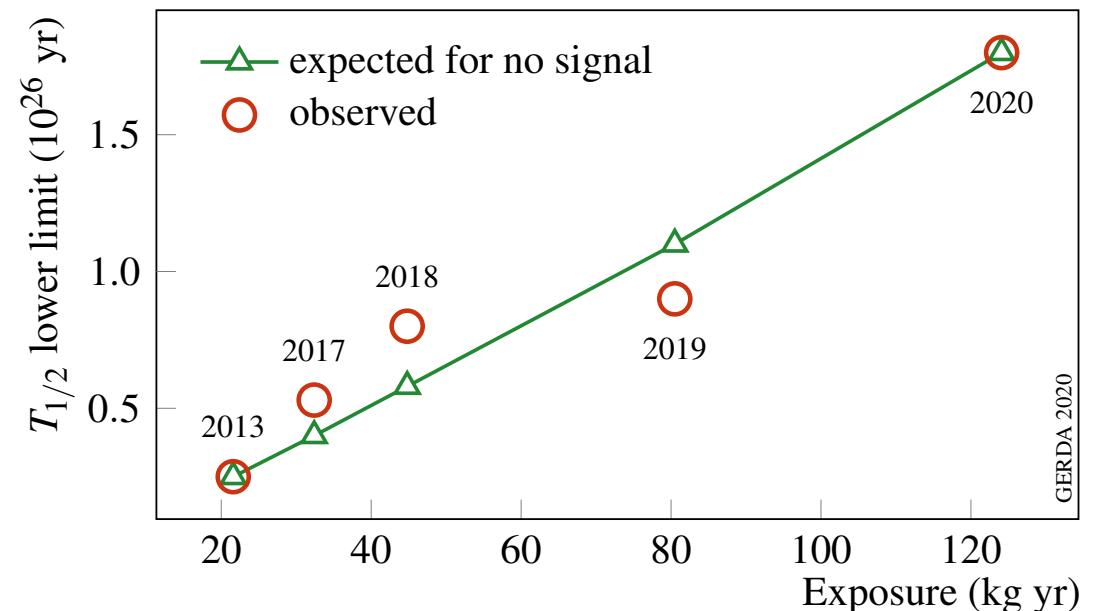


# Summary



- GERDA employed an array of HPGe detectors enriched in  $^{76}\text{Ge}$  to search for  $0\nu\beta\beta$
- GERDA ran in background-free regime for the entire duration of its data taking
- Provides the most stringent constraints on the half-life of  $0\nu\beta\beta$  decay
- $T_{1/2}^{0\nu} > 1.8 \cdot 10^{26} \text{ yr} (90\% \text{ C.L.})$
- $|m_{\beta\beta}| < [79 - 180] \text{ meV}$
- $B = 5.2^{+1.6}_{-1.3} \cdot 10^{-4} \left[ \frac{\text{cts}}{\text{keV}\cdot\text{kg}\cdot\text{yr}} \right] (68\% \text{ SI})$
- Next step ...

LEGEND



*Thank you for your attention !*

*BACKUP*

# Background reduction



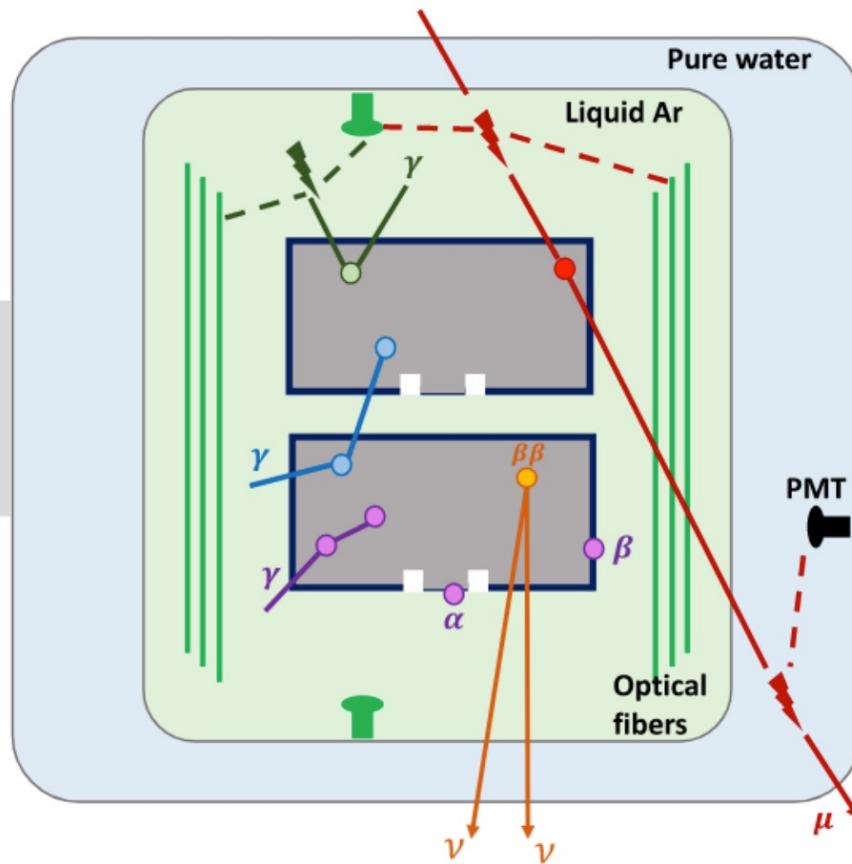
Analysis cuts:

- PSD
- multiplicity/coincidence
- Lar veto
- muon veto

Signal efficiencies after cuts:

- Coax 46%
- BEGe 61%
- IC 66%

$\beta\beta$  decay signal:  
single-site event  
energy deposition  
in a  $1 \text{ mm}^3$  volume



Pulse shape  
discrimination (PSD)  
for multi-site and  
surface  $\alpha, \beta$  events

Ge detector  
anti-coincidence

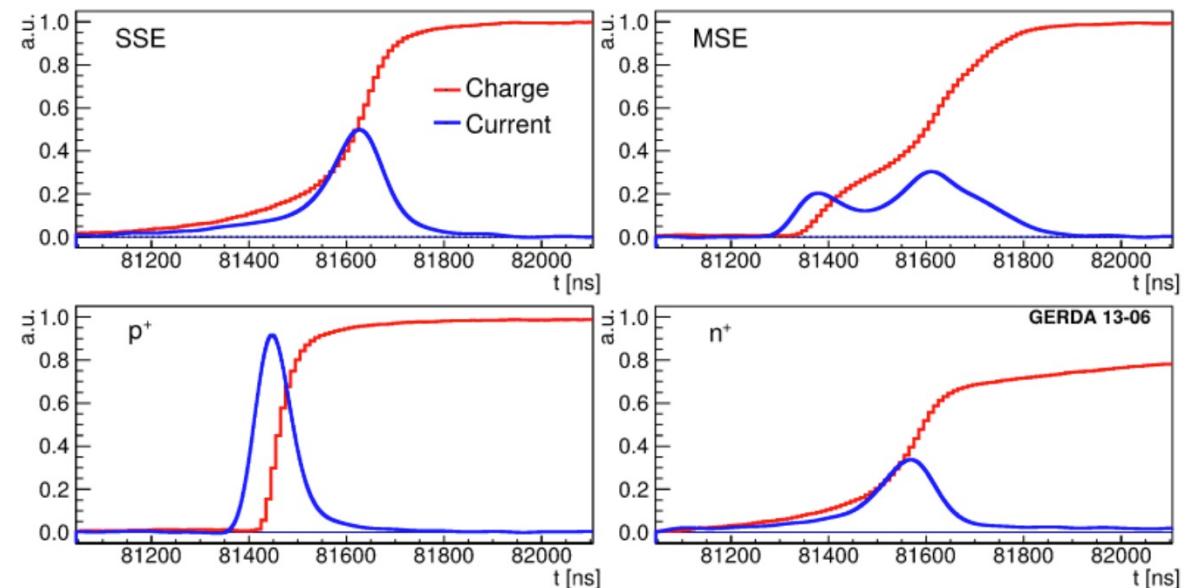
LAr veto based on Ar  
scintillation light read  
by fibers and PMT

Muon veto based on  
Cherenkov light and  
plastic scintillator

# Pulse Shape Discrimination



- **single-site events**: signal-like
- **multi-site events**: induce double-peak structure
- **surface  $\alpha$  events**: fast risetime, high current
- **surface  $\beta$  events**: incomplete charge collection
- rejection based on current amplitude over energy (A/E) for BEGe, IC & on artificial neural network comparing pulse shape for Coax

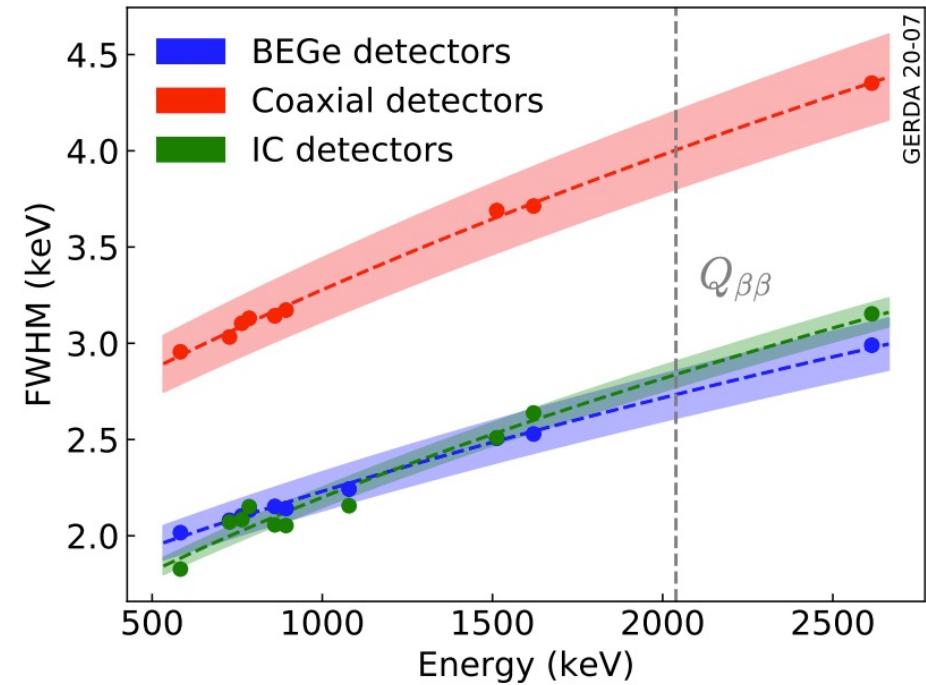


Eur. Phys. J. C 73 2583 (2013)

# Calibration



- Detectors calibrated weekly with 3  $^{228}\text{Th}$  sources
- Energy shifts between calibration  $< 1 \text{ keV}$
- peak fitting algorithm to determine each detector's resolution
- Gaussian mixture models to determine resolutions per detector type
- digital shaping with “zero area cusp” (ZAC) filter



Eur. Phys. J. C 75 255 (2015)

Eur. Phys. J. C 81 682 (2021)

# Detector specifications



TABLE I. Summary of the GERDA Phase II parameters for different detector types and before and after the upgrade. The components of the total efficiency  $\epsilon$  for  $0\nu\beta\beta$  decays are reported individually. The efficiencies of muon veto and quality cuts are above 99.9% and are not shown. Energy resolutions and all  $0\nu\beta\beta$  decay detection efficiencies are reported as exposure-weighted averages for each detector type and their uncertainties are given as standard deviations.

	Dec 2015–May 2018		July 2018–Nov 2019		
	Coaxial	BEGe	Coaxial	BEGe	Inverted coaxial
Number of detectors	7	30	6	30	5
Total mass	15.6 kg	20 kg	14.6 kg	20 kg	9.6 kg
Exposure $\mathcal{E}$	28.6 kg yr	31.5 kg yr	13.2 kg yr	21.9 kg yr	8.5 kg yr
Energy resolution at $Q_{\beta\beta}$ (FWHM)	$(3.6 \pm 0.2)$ keV	$(2.9 \pm 0.3)$ keV	$(4.9 \pm 1.4)$ keV	$(2.6 \pm 0.2)$ keV	$(2.9 \pm 0.1)$ keV
$0\nu\beta\beta$ decay detection efficiency $\epsilon$ :	$(46.2 \pm 5.2)\%$	$(60.5 \pm 3.3)\%$	$(47.2 \pm 5.1)\%$	$(61.1 \pm 3.9)\%$	$(66.0 \pm 1.8)\%$
Electron containment	$(91.4 \pm 1.9)\%$	$(89.7 \pm 0.5)\%$	$(92.0 \pm 0.3)\%$	$(89.3 \pm 0.6)\%$	$(91.8 \pm 0.5)\%$
$^{76}\text{Ge}$ enrichment	$(86.6 \pm 2.1)\%$	$(88.0 \pm 1.3)\%$	$(86.8 \pm 2.1)\%$	$(88.0 \pm 1.3)\%$	$(87.8 \pm 0.4)\%$
Active volume	$(86.1 \pm 5.8)\%$	$(88.7 \pm 2.2)\%$	$(87.1 \pm 5.8)\%$	$(88.7 \pm 2.1)\%$	$(92.7 \pm 1.2)\%$
Liquid argon veto		$(97.7 \pm 0.1)\%$		$(98.2 \pm 0.1)\%$	
Pulse shape discrimination	$(69.1 \pm 5.6)\%$	$(88.2 \pm 3.4)\%$	$(68.8 \pm 4.1)\%$	$(89.0 \pm 4.1)\%$	$(90.0 \pm 1.8)\%$

- closest event at  $2.4\sigma$

PRL 125, 252502 (2020)

# Likelihood

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- Signal rate  $S = 1 / T_{1/2}^{0\nu}$
- Expected number of signal events in partition  $k$  ( $\varepsilon_k$  exposure,  $\epsilon_k$  efficiency,  $m_{76}$  molar mass):

$$\mu_{s,k} = \frac{\ln 2 \mathcal{N}_A}{m_{76}} \epsilon_k \varepsilon_k S$$

- Expected number of background events in partition  $k$  ( $B_k$  bkg index,  $\Delta E$  analysis window width):

$$\mu_{s,b} = B_k \Delta E \varepsilon_k$$

- Gaussian distribution for the signal, centered at  $Q_{\beta\beta}$  with a width corresponding to the energy resolution ( $\sigma_k$ ), and a flat distribution for the background
  - the hypothesis of a flat background is supported by means of a test-statistic derived from Order-Statistic, which models the distribution of spacings between statistical samples

# Sensitivity vs. Exposure



- Number of signal and background counts are Poisson distributed
- Expected number of signal counts ( $f$  enrichment fraction,  $e$  efficiency,  $\epsilon = f \cdot e$ ):

$$n_s = \frac{\ln(2) N_A}{m_{76}} fe \epsilon \frac{1}{T_{1/2}^{0\nu}}$$

- Expected number of background counts (around  $Q_{\beta\beta}$ , thus  $\Delta E = \sigma_E$ ):

$$n_b = \epsilon B \Delta E$$

- If  $n_b = 0$  (background free), then  $T_{1/2}^{0\nu}$  is the time needed to observe 1 signal event:

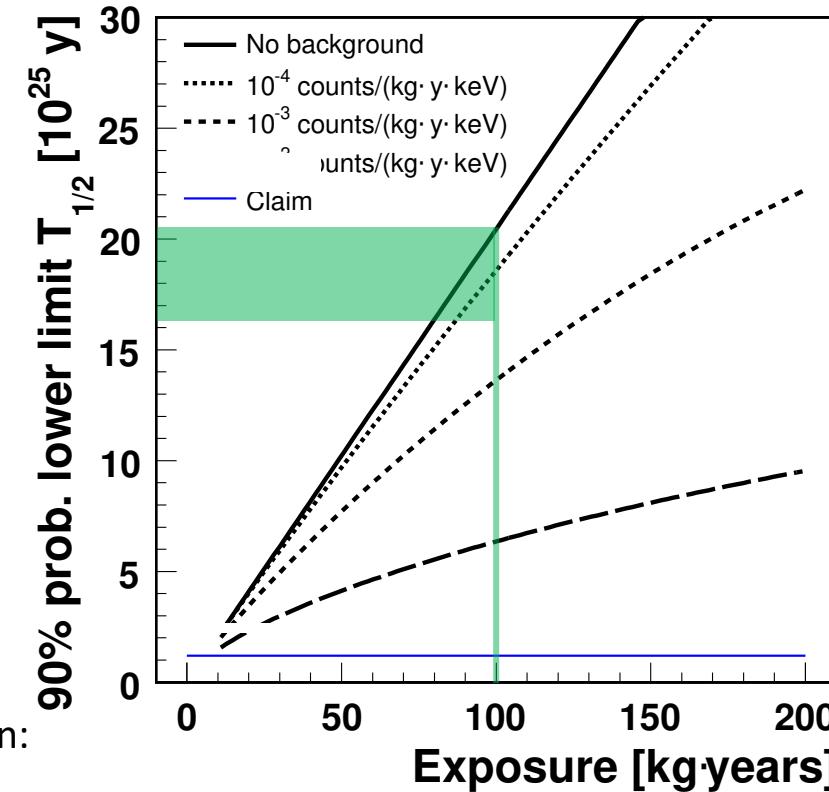
$$T_{1/2}^{0\nu} \propto fe \epsilon$$

- If  $n_b \gg 1$ , then we can approximate the Poisson statistics with a Gaussian distribution:

$$n_b \sim \mathcal{N}(\epsilon B \Delta E, \sqrt{\epsilon B \Delta E})$$

- the minimal number of signal counts that can be distinguished from the background is approximately  $\sqrt{n_b}$ , thus:

$$T_{1/2}^{0\nu} \propto fe \sqrt{\frac{\epsilon}{B \Delta E}}$$



# Frequentist analysis (I)

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- Two sided test statistic based on the profile likelihood  $\lambda(S)$ :

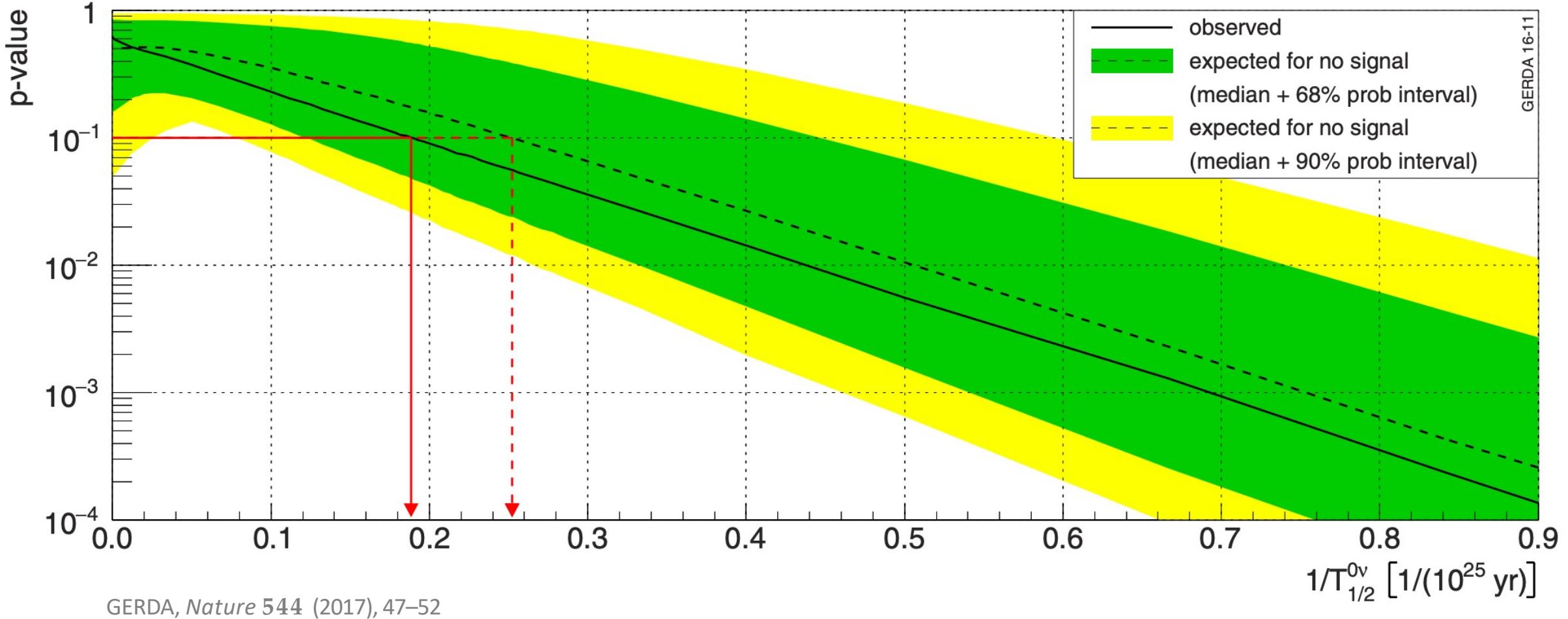
$$t_S = -2 \ln[\lambda(S)] = -2 \ln \left[ \frac{\mathcal{L}(S, \hat{B}, \hat{\theta})}{\mathcal{L}(\hat{S}, \hat{B}, \hat{\theta})} \right]$$

- $\hat{B}$  and  $\hat{\theta}$  denote the value of the parameters that maximize  $\mathcal{L}$  for a fixed  $S$
- $\hat{S}, \hat{B}$  and  $\hat{\theta}$  denote the values corresponding to the absolute maximum likelihood
- Estimate the distribution  $f(t_S | S)$  using MCMC
- The p-value for data at a specific value of  $S$  is:

$$p_S = \int_{t_{obs}}^{\infty} f(t_S | S) d(t_s)$$

- The 90% CL is given by all  $S$  values with  $p_S > 0.1$

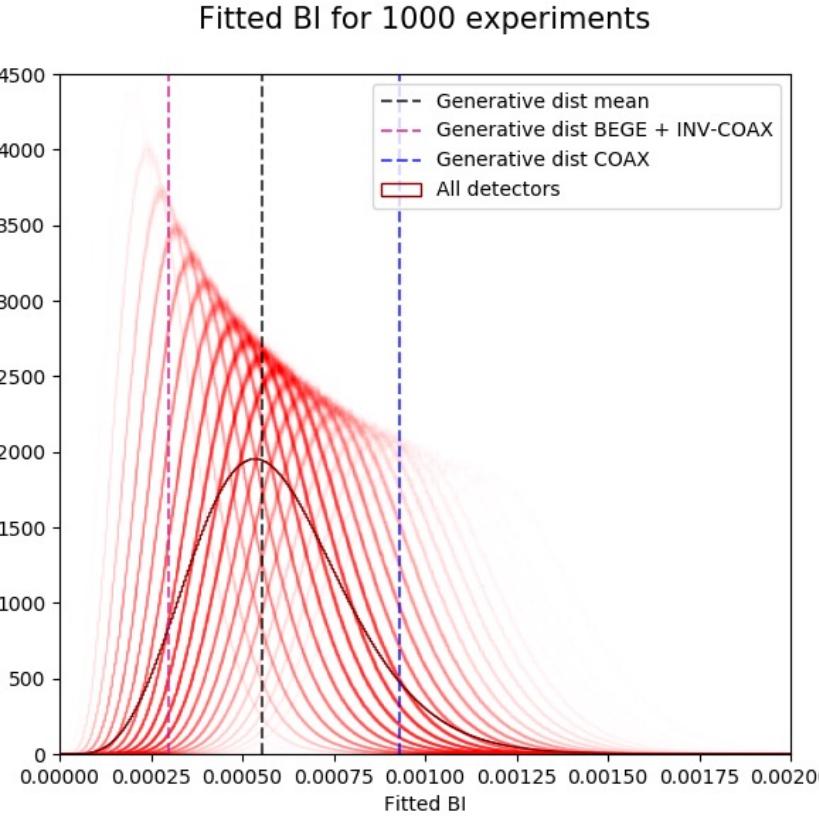
# Frequentist analysis (II)



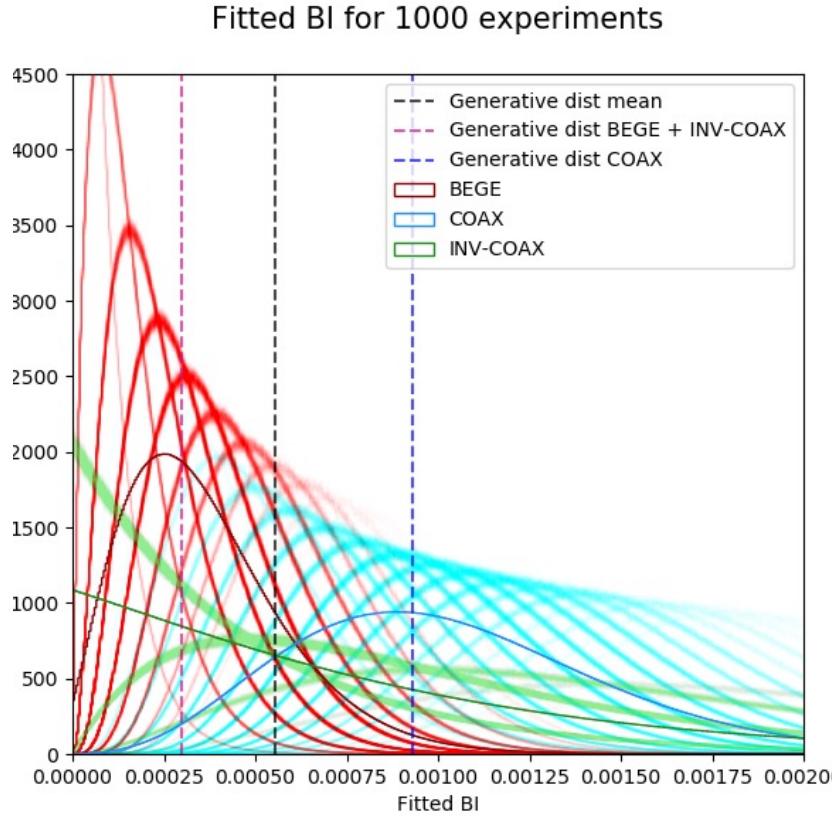
# Comparison with toy experiment



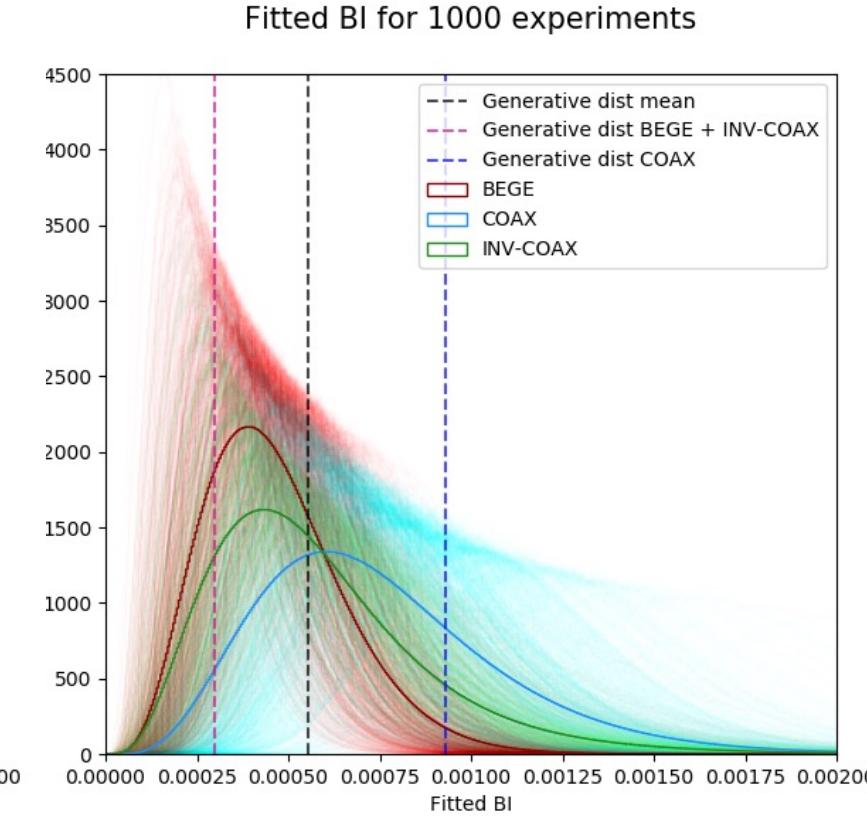
Single BI



Uncorrelated BI



Correlated BI



# Model comparison

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- Single B model reconstructs a larger fraction of signal events on average: **stronger discovery power**
- Uncorrelated B model gives on average a better half-life limit in experiments with only background events: **stronger limit setting capabilities**
- Correlated B model's performance is halfway between the extreme models both in discovery power and limit setting
- The (median) sensitivity of all models assuming no signal and using a uniform prior for  $S$  is

$$T_{1/2}^{0\nu} > 1.4 \cdot 10^{26} \text{ yr (90% C.I.)}$$