

Dark Matter and Leptoquarks

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Dark Matter vs B-physics?

- Experimental fact: Dark Matter exists
- In principle, its mass and couplings can span tens of orders of magnitude!
- If it has (weak) interactions with thermal bath of SM particles, elegant way to predict its cosmological abundance
- Coupling with the SM Dark Sector chessboard Collider physics Direct detection freeze-out Astro-particle physics and Cosmology Gravitational-wave physics WIMP MeV GeV keV SIDM,mDM,... only scalars freeze-in **PBHs** misalignment

- → M_{DM} ~ TeV with ~ EW coupling
- Interplay with popular models of New Physics at the EW scale (e.g. SUSY, composite models, ...)

Post-LHC: -



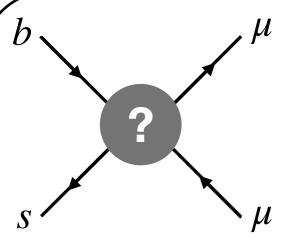
no light on-shell resonances



very interesting anomalies in flavour observables

Leptoquarks as mediators

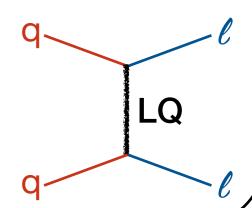
B-anomalies:



Evidence for new interaction between quarks & leptons

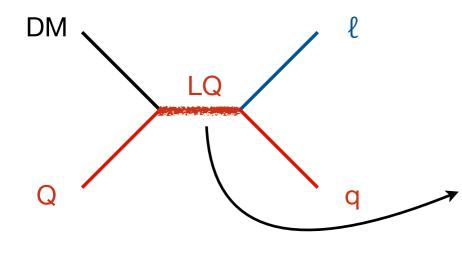
$$\frac{1}{\Lambda^2} (\bar{b}_L \gamma_\nu s_L) (\bar{\mu}_L \gamma^\nu \mu_L) \qquad \Lambda \approx 30 \,\text{TeV} \approx 6 \,\text{TeV} \cdot V_{cb}^{-1/2}$$

- Yukawa-like couplings, larger for heavy generations
- Effect only in semi-leptonic processes VS.
 strong bounds from pure-quark and -lepton observables
 - hint for leptoquark mediator



see talks by Luca, Claudia, David

Can Leptoquarks be the mediators of SM-DM interactions?

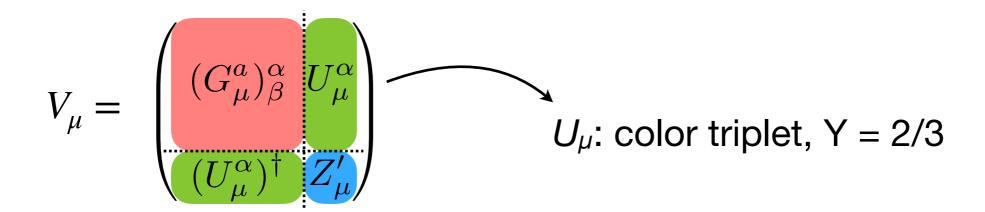


colored: must couple to a quark

A (too) simple picture

Example: Pati-Salam SU(4) unification

extension of SM gauge group $SU(4)_{PS} \longrightarrow SU(3)_c \times U(1)_{B-L}$



Fermions are 4-plets of SU(4): quark-lepton unification

Possible mediators

Consider a fermion SM singlet χ (sterile neutrino).

Mediators that can couple (at tree-level) to χ and a SM fermion:

+ Vector Leptoquark: $U_1 \sim (3, 1, -2/3)$

$$\mathcal{L}_{U} = \left[g_{L}(\bar{Q}_{L}\gamma_{\mu}L_{L}) + g_{R}(\bar{d}_{R}\gamma_{\mu}\ell_{R}) + g_{\chi}(\bar{u}_{R}\gamma_{\mu}\chi) \right] U^{\mu}$$

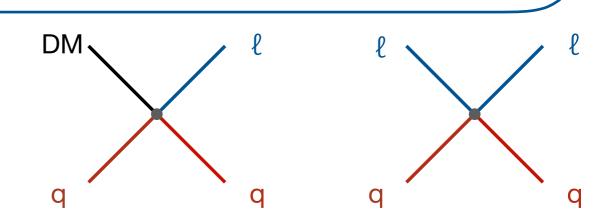
Scalar Leptoquarks:

$$S_1 \sim (\mathbf{3}, \mathbf{1}, 1/3) \qquad \mathcal{L}_{S_1} = \left[g_L(\bar{Q}_L \varepsilon L_L^c) + g_R(\bar{u}_R \ell_R^c) + g_{\chi}(\bar{d}_R \chi) \right] S_1$$

$$R_2 \sim (\mathbf{3}, \mathbf{2}, 1/6) \qquad \mathcal{L}_{\tilde{R}_2} = \left[g_L(\bar{d}_R L_L \varepsilon) + g_{\chi}(\bar{Q}_L \chi) \right] \tilde{R}_2$$

- + Right-handed vector: $W_R \sim (\mathbf{1},\mathbf{1},1)$ $\left[g_R(\bar{u}_R\gamma_\mu d_R) + g_\chi(\bar{\ell}_R\gamma_\mu\chi)\right]W_R^\mu$
- + (Higgs doublet: $Y_{\chi}(\bar{L}_L\chi)H^c$)

Integrating out the heavy mediator: effective interactions



Effective interactions

Consider a fermion SM singlet χ (sterile neutrino).

* Most general effective Lagrangian for χ interacting with SM:

$$\begin{split} \mathcal{L}_{\mathrm{DM}}^{\mathrm{eff}} &= -\,C_{\chi R}(\bar{\chi}\gamma^{\mu}\ell_R)(\bar{d}_R\gamma_{\mu}u_R) + 2C_{\chi L}(\bar{\chi}L_L)(\bar{Q}_Lu_R) \\ &\quad -\,C_{\chi L}'(\bar{L}_L\chi)\varepsilon(\bar{Q}_Ld_R) - C_{\chi L}''(\bar{Q}_L\chi)\varepsilon(\bar{L}_Ld_R). \end{split}$$
 with $C_{IJ} = g_Ig_J/M_{\mathrm{LO}}^2$

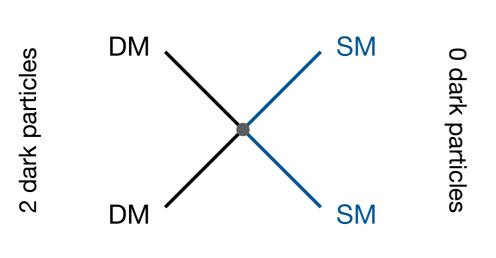
Below the QCD scale: eff. coupling to pions and hadrons

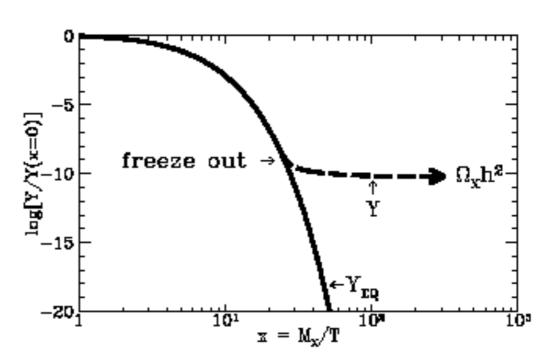
$$\mathcal{L}_{\mathrm{DM}}^{\pi} = \frac{C_{\chi R}}{\sqrt{2}} (\bar{\chi} \gamma^{\mu} \ell_R) \Big[f_{\pi} \, D_{\mu} \pi^+ + i (\pi^0 \partial_{\mu} \pi^+ - \pi^+ \partial_{\mu} \pi^0) + \cdots \Big] + \\ - i C_{\chi L} f_{\pi} B_0 \Big[\sqrt{2} (\bar{\chi} \ell_L) \pi^+ + (\bar{\chi} \nu_L) \pi^0 + \cdots \Big] + \mathrm{h.c.}$$
 (chiral Lagrangian)

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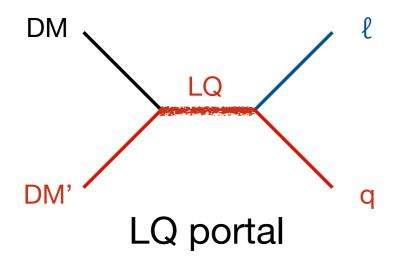
DM production

- ♦ Usually: DM is stable thanks to an exact (\mathbb{Z}_2) global symmetry $\chi \to -\chi$
- * Production in early Universe: thermal freeze-out of $2 \rightarrow 0$ scatterings





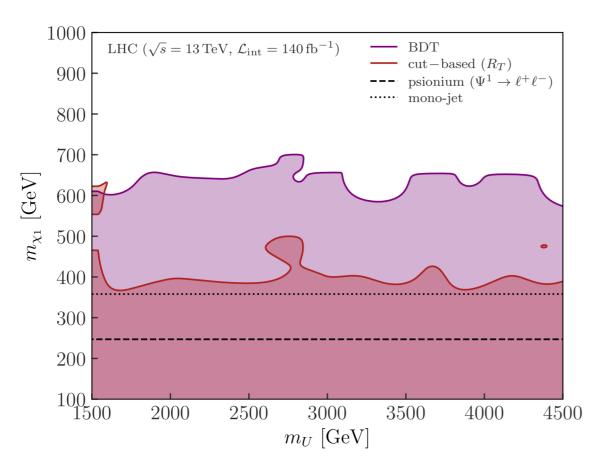
→ Introduce a Dark Sector with new colored states odd under Z₂

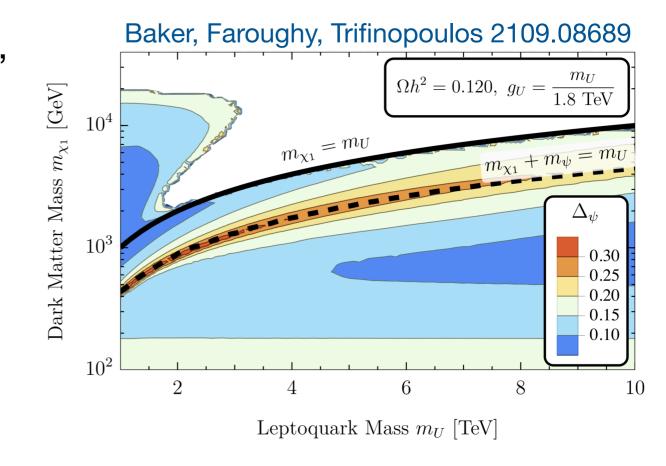


- Baker, Faroughy, Trifinopoulos 2109.08689
- Belanger et al. 2111.08027
- Guadagnoli, Reboud, Stangl 2005.10117

DM production: annihilation

- Relic abundance can be reproduced, by coannihilation with colored DM'
 (DM-DM' splitting can be adjusted)
- Consequences for phenomenology:





can search for colored states at LHC: Leptoquark & DM partner

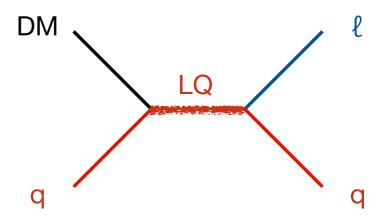
See also: Belanger et al. 2111.08027

Guadagnoli, Reboud, Stangl 2005.10117

Second possibility: allow only one LQ-DM-SM coupling (LQ is charged under Z₂ and part of dark sector) Mandal 1808.07844

Thermal decays

- → We want to explore a more minimal setup: no separate Dark Sector
 - $1 \rightarrow 0$ processes are allowed
 - Does induce DM decay! DM → SM

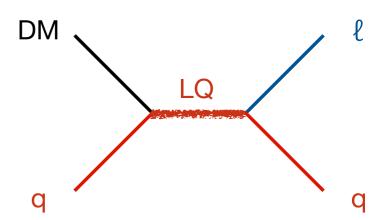


Contribute to DM annihilation in early universe, DM SM → SM SM similar to DM decays in a thermal bath: "thermal decays"

2111.14808 with Belfatto, Gross, Panci, Strumia, Vignaroli, Vittorio, Watanabe

Thermal decays

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Contribute to DM annihilation in early universe, DM SM → SM SM similar to DM decays in a thermal bath: "thermal decays"

Naively:

- + $\tau_{\rm DM} = 1/\Gamma_{\rm DM} > 10^{18}\,{\rm s}$ for DM stability (bound 10²⁵ s in many cases!)
- $+ \Gamma_{\rm ann} \sim H \sim T^2/M_{\rm Pl} \approx \frac{(T/{\rm GeV})^2}{10^{-5}\,s} \ {\rm for\ freeze-out\ at\ T\sim M}$
- $\rightarrow \Gamma_{\rm DM}/\Gamma_{\rm ann} \lesssim 10^{-(23 \div 30)}$

However, different phase space and energy scale!

Freeze-out of relativistic DM

Decay and "thermal decay" rates not equal in general!

$$\Gamma_{\rm DM} pprox G_{\rm NP}^2 M_{\rm DM}^5$$
 must be small

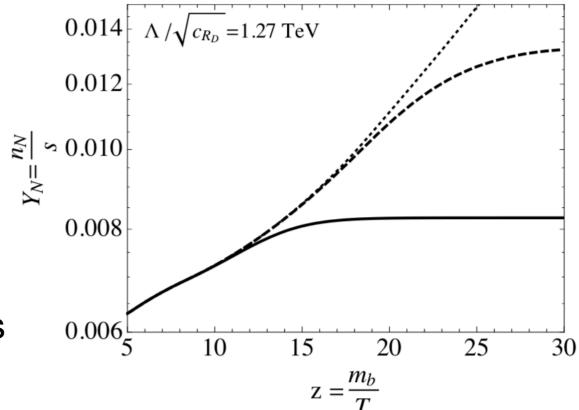
$$\Gamma_{\rm ann} \approx G_{\rm NP}^2 T^5$$
 must be large

Freeze-out at T ≫ M: relativistic DM. Decay suppressed by small mass

Azatov, Barducci, Ghosh, Marzocca, Ubaldi 1807.10745

$$\Omega_{\rm DM} h^2 = 0.12 \frac{50}{g_{\star}(T_d)} \frac{M_{\chi}}{50 \, {\rm eV}}$$

 $(\bar{\chi}\gamma^{\mu}\tau_R)(\bar{b}_R\gamma_{\mu}c_R)$ effective interaction with couplings that fit flavor anomalies



Sterile neutrino of 50 eV: hot DM, ruled out by structure formation

Cold Dark Matter

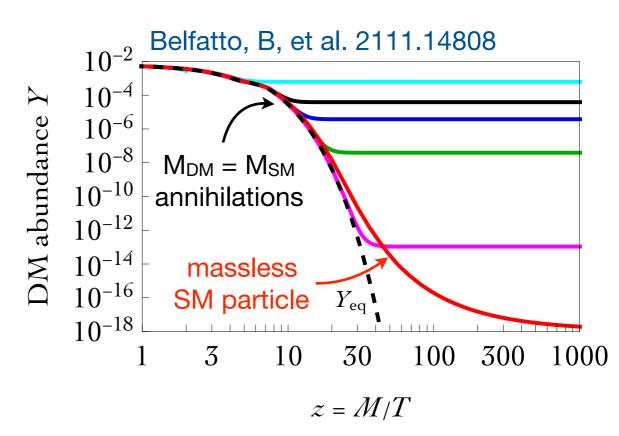
Boltzmann equation:

$$z \frac{dY}{dz} \approx -2\gamma_{\rm ann} \left(\frac{Y^2}{Y_{\rm eq}^2} - 1\right) - \gamma_{\rm dec} \left(\frac{Y}{Y_{\rm eq}} - 1\right)$$

usual DM DM → SM SM

$$\gamma_{AB\to X} \propto \langle \sigma v \rangle n_A n_B$$

 $\begin{array}{c} \mathsf{DM} \; \mathsf{SM} \to \mathsf{SM} \; \mathsf{SM} \\ \text{``thermal decay''} \end{array}$



 Non-relativistic limit: freeze-out below M_{DM}.
 "Decay" term is less Boltzmann-suppressed (analytic results in 2111.14808)

Cold Dark Matter

Boltzmann equation:

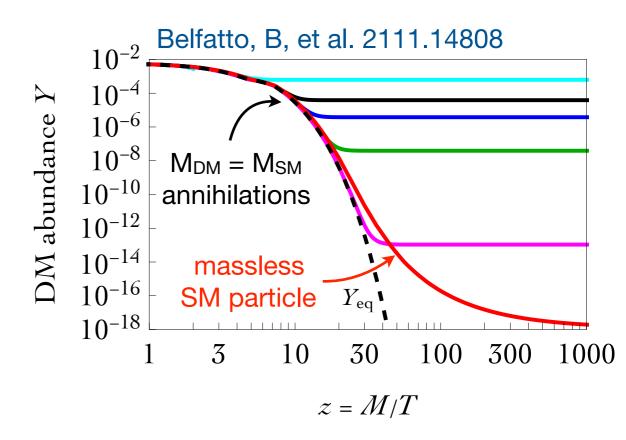
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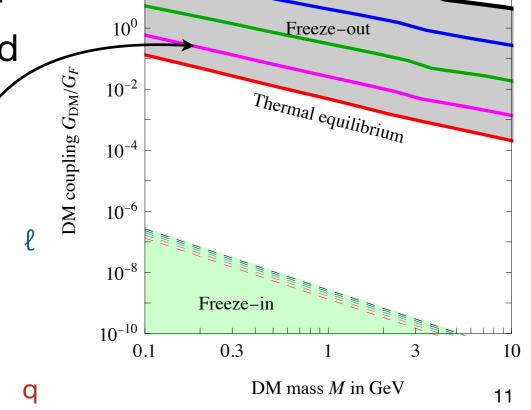
DM SM → SM SM "thermal decay"

DM.



- Non-relativistic limit: freeze-out below M_{DM}.
 "Decay" term is less Boltzmann-suppressed (analytic results in 2111.14808)
- Freeze-out of dim.6 interaction:

$$G_{\rm DM} \gtrsim 10^{-4} G_F \left({\rm GeV}/M_{\rm DM}\right)^{3/2}$$



Cold Dark Matter

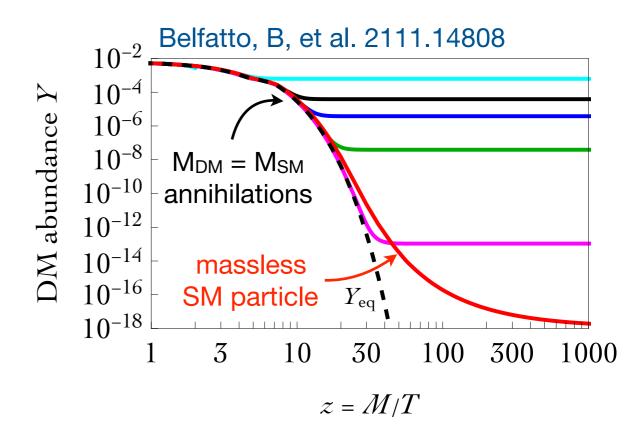
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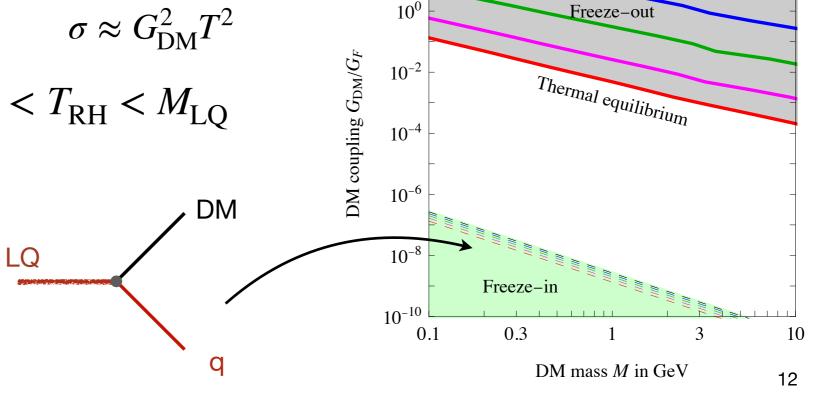


* Freeze-in: DM production most efficient at the highest energy $\sigma \approx G_{\rm DM}^2 T^2$

Reheating temperature $M_{\rm DM} < T_{\rm RH} < M_{\rm LQ}$

$$Y_{\infty} \approx G_{\rm DM}^2 M_{\rm Pl} T_{\rm RH}^3$$

(LQ in thermal eq. that decays into DM)

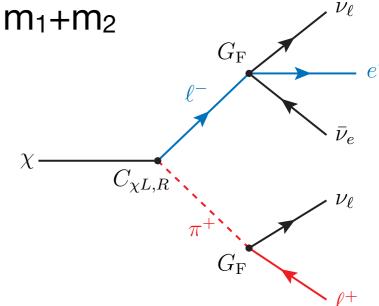


Dark Matter decays

- **→** Tree-level DM \rightarrow SM₁ SM₂ must be closed: M_{DM} < m₁+m₂
- Off-shell decays:

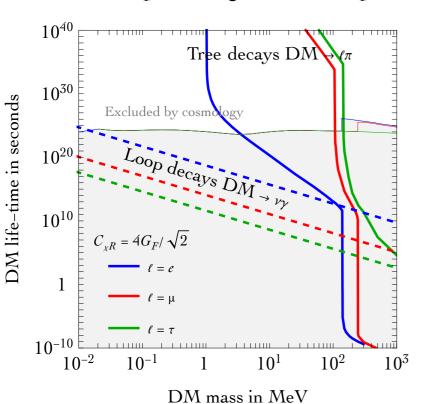
$$\Gamma(\mathrm{DM} \to \mathrm{SM_1\,X}) \approx \frac{\Gamma(\mathrm{DM} \to \mathrm{SM_1SM_2^*})}{1 - p_2^2/m_2^2} \cdot \frac{\Gamma(\mathrm{SM_2} \to \mathrm{X})}{4\pi m_2}$$

need long-lived SM₂ particles

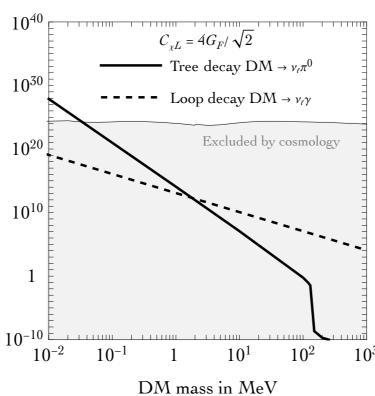


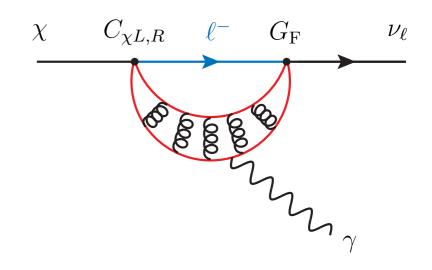
Loop decays: $\Gamma(\mathrm{DM} \to \nu \gamma) \propto M_{\mathrm{DM}}^3 C_{\chi L, R}^2 \times \left(\frac{G_F \Lambda_{\mathrm{QCD}}^2}{4\pi}\right)^2$

DM coupled to right-handed leptons

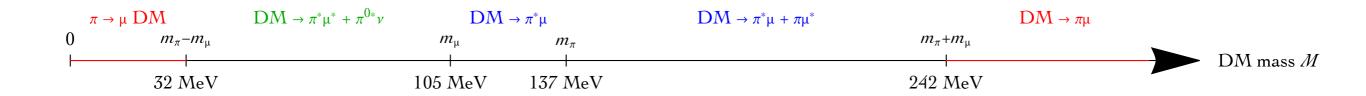


DM coupled to left-handed leptons





Example: DM coupled to muons and pions



Thermal freeze-out is excluded by too fast DM decays

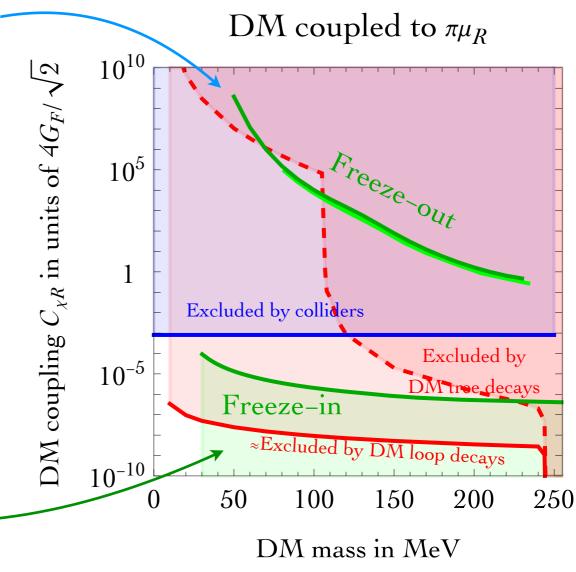
(& collider bounds on lepton + MET)

Thermal freeze-in is viable:

If
$$T_{RH} > M_{LQ}$$
, $Y_{\infty} \approx g_{\chi}^2 \frac{M_{Pl}}{M_{LQ}}$

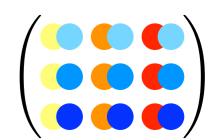
Matching the cosmological DM abundance gives

$$g_{\chi} \approx 10^{-8} \sqrt{\frac{\text{MeV}}{M_{\text{DM}}} \frac{M_{\text{LQ}}}{\text{TeV}}}$$



Connection with flavor

- Is there an explanation for the tiny coupling required by freeze-in?
- What is the flavor structure of the LQ couplings?
 - + Simple flavor ansatz: $Y_u^{ij} \approx \epsilon_{Q_i} \epsilon_{u_j}$ (and similar for d, e) (e.g. Froggatt-Nielsen models, partial compositeness, ...)



B-anomalies:

- Yukawa-like couplings, larger for heavy generations
- CKM-like suppression of flavor-changing effects

Assume LQ couplings follow the same pattern:

$$g_L^{ij} pprox \epsilon_{Q_i} \epsilon_{L_j}$$
, $g_R^{ij} pprox \epsilon_{d_i} \epsilon_{e_j}$, $g_\chi^i pprox \epsilon_{u} \epsilon_\chi$ new parameter

Connection with flavor

+ $\epsilon_{O,u,d}$ can be estimated from V_{CKM} and quark masses

$$V_{\mathrm{CKM}}^{ij} \approx \epsilon_{Q_i}/\epsilon_{Q_j}$$

• ϵ_L can be estimated from b \rightarrow sµµ anomalies:

$$\Delta C_9 \approx V_{ts} V_{tb}^* \frac{(\epsilon_{L_2})^2}{M_{LO}^2} \sim V_{ts} V_{tb}^* \frac{1}{(6 \text{ TeV})^2}$$

 $(\epsilon_{L_2} \approx 0.1 \text{ for } M_{LO} \sim \text{TeV})$

+ No sizable effect in electrons: $\epsilon_{L_1} \ll \epsilon_{L_2}$

$$\begin{split} \epsilon_Q &\sim (\lambda^3, \lambda^2, \lambda), \\ \epsilon_u &\sim (\lambda^4, \lambda, 1), \\ \epsilon_d &\sim (\lambda^4, \lambda^3, \lambda^2). \end{split}$$

+ DM freeze-in:
$$g_\chi \approx \epsilon_\chi \epsilon_u \sim 10^{-8} \sqrt{{\rm MeV}/M_\chi}$$

If dominant coupling is with 1st generation, $\epsilon_{\gamma} \approx 10^{-3} \approx \epsilon_{e_1}$

Connection with flavor

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Leptons
$$\epsilon_{L_1} < 0.1, \qquad \epsilon_{L_2} \sim 0.1, \qquad \epsilon_{L_3} \sim 1, \\ \epsilon_{e_1} > 10^{-5}) \qquad \epsilon_{e_2} \sim 10^{-3}, \qquad \epsilon_{e_3} \sim 10^{-2}.$$

$$\epsilon_{u} \sim (\lambda^4, \lambda, 1), \\ \epsilon_{d} \sim (\lambda^4, \lambda^3, \lambda^2).$$

$$\begin{split} \epsilon_Q &\sim (\lambda^3, \lambda^2, \lambda), \\ \epsilon_u &\sim (\lambda^4, \lambda, 1), \\ \epsilon_d &\sim (\lambda^4, \lambda^3, \lambda^2). \end{split}$$

+ DM freeze-in: $g_{\chi} \approx \epsilon_{\chi} \epsilon_{u} \sim 10^{-8} \sqrt{\text{MeV}/M_{\chi}}$

If dominant coupling is with 1st generation, $(\epsilon_\chi \approx 10^{-5}) \approx \epsilon_{e_1}$

More than one DM flavor

- * DM could also carry a flavor index, χ^i .
- With more than one state, thermal decay freeze-out is possible:

- 1. $C_{2,12}$ large enough to achieve thermal freeze-out $\chi_1 \leftrightarrow \chi_2 \leftrightarrow SM$
- 2. χ_2 decays dominantly into χ_1 if $M_1 < M_2 < m_{\ell} + m_{\pi}$
- 3. χ_1 is long-lived if C_1 is small enough
 - DM is not part of a separate Dark Sector

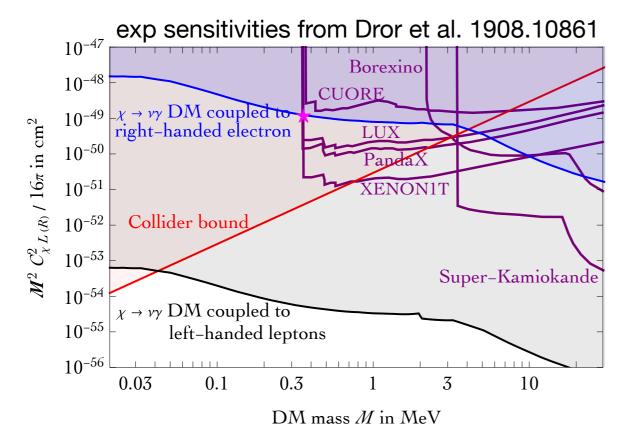
 χ_2

SM

Direct and indirect detection

+ Charged-current interaction: $C_{\chi R}(\bar{\chi}\gamma_{\mu}e_R)\left(\bar{n}\gamma^{\mu}\frac{1+g_A\gamma_5}{2}p\right)+C_{\chi L}(\bar{\chi}e_L)(\bar{n}p)$

DM-induced
$$\beta$$
- transition DM + $^{A}_{Z}N \rightarrow e^{-} + ^{A}_{Z+1}N$



$$\langle \sigma v \rangle = \frac{C_{\chi L(R)}^2 M_{\chi}^2}{16\pi} f(E)$$
Nucleus

monochromatic electron with $E=M_{\rm DM}-m_e-Q$

- Depending on nuclear thresholds Q, sensitive to DM heavier than few 100 keV
- Can also look for signal from decaying $_{Z+1}^{A}N$
- ▶ Neutral-current interaction: nuclear recoil in $DM + N \rightarrow \nu + N$

DM

Summary

- Dark Matter could interact with the SM through Leptoquarks.
- DM abundance from freeze-in of DM SM → SM SM scatterings
 - small coupling, mass ~ MeV
 consistent with expectation for 1st generation fermions
 - Exotic signatures in Direct Detection: DM-induced β decay
- DM abundance through thermal freeze-out
 - → DM SM \rightarrow SM SM thermal decays, if more than one generation of DM
 - DM DM → SM SM annihilations, if Dark Sector contains colored states