# Electron cloud incoherent effects 

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## Outline

1. Introduction
2. Observations in LHC during Run 2
3. Progress in simulations
a) Description of e-cloud interaction
b) Simulation results

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## Buildup example



- E-cloud buildup is a multi-bunch effect.
- Incoherent e-cloud effect concerns the motion of single beam particles as they encounter these electrons.
- Single particles stay inside the same bunch


## Pinch



- Motion of electrons is very complex
- Complex electron densities $\rightarrow$ complex induced forces.
- Betatron oscillations: up-down, left-right
- Synchrotron oscillations: back-forth in "time"
- Non-linear forces + betatron/synchrotron oscillations can lead to oscillation amplitude increase $\rightarrow$ losses + emittance growth.


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## Motivation



Losses come from:

- Luminosity burn-off that decreases gradually.
- Continuous rate of additional losses.


## Filling scheme



Standard 2018 Physics filling scheme (2556 bunches) [lpc.web.cern.ch]

Beam is composed of repeating patterns (trains):

- $2 \times 48$ bunches,
- $3 \times 48$ bunches.

Magnification:


## All bunch-by-bunch losses



Global picture: Fairly constant loss rate (Corrected for burn-off).

- Grows from head to tail of each train


## Number of BBLR interactions

Number of Beam-Beam Long-Range interactions changes for each bunch in the filling scheme.


"Train"

- Group 1: Few BBLR, reduced e-cloud effects
- Group 2: Max BBLR, reduced e-cloud effects
- Group 3: Max BBLR, stronger e-cloud effects
- Group 4: Few BBLR, stronger e-cloud effects


## Example \#1: Physics fills




- Bunch-by-bunch pattern emerges
- Reminds of e-cloud buildup behaviour.
- Beam-beam effects alone cannot explain behaviour


## Example \#2: Crossing angle

Typical physics fill:



Special test:




- A reduced crossing angle typically enhances BBLR interactions.
- In this case, it enhances the e-cloud pattern losses.


## Example \#3: Buildup simulations in Inner Triplet quadrupoles

One beam:


Two beams:


- One beam: In the small 200 ns between batches, the electron cloud decays significantly.
- Two beams: Beams are not synchronized and the e-cloud does not decay.


## Example \#3: Buildup simulations in Inner Triplet quadrupoles

One beam:



Two beams:



The bunch-by-bunch pattern of the losses resembles the e-cloud buildup simulations of the Inner Triplet quadrupoles.

## Example \#4: Measurements with different betatron functions

$\beta^{*}=65 \mathrm{~cm}, \varphi=120 \mu \mathrm{rad}$
Large ATS telescope ${ }^{1} \rightarrow$
$\rightarrow$ enhancement of arc beta functions

$$
\beta^{*}=30 \mathrm{~cm}, \varphi=150 \mu \mathrm{rad}
$$

Moderate ATS telescope

- Decreasing $\beta$ in the inner triplet quadrupoles should reduce effect of the e-cloud in the inner triplet.
- Increasing $\beta$ in arcs should enhances e-cloud effect:
no significant losses.


## Summary - Observations

Electron cloud related losses are enhanced when:

1. reducing $\beta^{*}$ (increasing $\beta$ in IT)
2. reducing crossing angle (changes closed orbit in IT)
3. Two beams are present (enhanced buildup in IT) but not when:
4. Increasing $\beta$ in arcs


All observations point to the Inner Triplet Quadrupoles.

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## Introduction to simulations

[G. Iadarola, CERN-ACC-NOTE-2019-0033]


- Complex time-dependent e-cloud density $\rightarrow$ complex time-dependent forces
- Slow incoherent effects $\rightarrow$ e-cloud can be re-used = weak-strong approximaton (no self-consistency)
- But: e-cloud potential (PIC) is defined on a 3D grid. Needs to be interpolated.


## Symplecticity

- Numerical methods in solving Hamiltonian systems can break the symplectic condition, making them less accurate at long timescales. (Millions of turns)
- Typically important to preserve symplecticity, even at the expense of accuracy.

- Interpolation scheme should guarantee symplecticity.

$$
\begin{aligned}
& \quad x, y, \tau \mapsto x, y, \tau \\
& \quad \text { In our case, } \quad \frac{\partial^{2} \phi}{\text { symplecticity: }} \frac{\partial^{2} \phi}{\partial x \partial y}=\frac{p_{x}}{\partial y \partial x} \mapsto p_{x}-\frac{q L}{\beta_{0} P_{0} c} \frac{\partial \phi}{\partial x}(x, y, \tau) \\
& p_{y} \mapsto p_{y}-\frac{q L}{\beta_{0} P_{0} c} \frac{\partial \phi}{\partial y}(x, y, \tau) \\
& \text { Linear interpolation is not symplectic. } p_{\tau} \mapsto p_{\tau}-\frac{q L}{\beta_{0} P_{0} c} \frac{\partial \phi}{\partial \tau}(x, y, \tau)
\end{aligned}
$$

## Tricubic interpolation

Given a regular 3D grid of any function $\mathrm{f}^{\mathrm{jjk}}$, we interpolate locally in a way that the following quantities are continuous globally.

$$
\left\{f, \frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z}, \frac{\partial^{2} f}{\partial x \partial y}, \frac{\partial^{2} f}{\partial x \partial z}, \frac{\partial^{2} f}{\partial y \partial z}\right\}
$$



Lekien and Marsden* proved that it is possible to meet this condition by using a tricubic interpolation scheme.

$$
f(x, y, z)=\sum_{i=0}^{3} \sum_{j=0}^{3} \sum_{k=0}^{3} a_{i j k} x^{i} y^{j} z^{k}
$$

The coefficients $\mathrm{a}_{\mathrm{ijk}}$ change from cell to cell but required quantities stay continuous across the cells.

- Analytical derivatives for interaction.
*F. Lekien and J. Marsden, "Tricubic interpolation in three dimensions". https://doi.org/10.1002/nm2.1296


## Issue with PIC potential



- PIC simulation suffers from macroparticle noise.
- Can be reduced by averaging many simulations.

Averaging 4000 reveals the physical structures in the induced forces.

## Issue with cubic interpolator




- Close look reveals irregularities from Tricubic interpolation.
- Inaccuracies are correlated with discontinuity of second derivative accross cells.

$$
\mathcal{E}_{x}^{i j k}=\left.\frac{\partial e_{x}^{\mathrm{int}}}{\partial x}\right|_{x \rightarrow x_{i}^{+}}-\left.\frac{\partial e_{x}^{\mathrm{int}}}{\partial x}\right|_{x \rightarrow x_{i}^{-}}=-2 \frac{\partial^{3} \phi}{\partial x^{3}}\left(x_{i}, y_{j}, \tau_{k}\right) \Delta x+O\left(\Delta x^{3}\right)
$$

## Refinement of potential

We found that we can treat our potential by:

1. Interpolate charge density on an auxilliary finer grid (by factor h).
2. Recalculate $\varphi$ and derivatives in the finer grid.
3. Store recalculated $\varphi$ and derivatives on original grid. $\Delta x_{\text {refined }}=\frac{\Delta x}{h}$

Minimal expense on memory and speed (performed during pre-processing) Proved analytically that error becomes:

$$
\mathcal{E}_{x, \text { refined }}^{i j k}=-2 \frac{\partial^{3} \phi}{\partial x^{3}}\left(x_{i}, y_{j}, \tau_{k}\right) \frac{\Delta x}{h^{2}}+O\left(h^{-4} \Delta x^{3}\right)
$$



Complete mitigation of the irregularities.

## Quick recap

- Analytical form of e-cloud kick.
- Used a high-order interpolation scheme (tri-cubic) to preserve symplecticity everywhere in phase space.
- Averaged multiple Particle-In-Cell e-cloud simulations to reduce macroparticle noise in the interpolated data.
- Solved Poisson's equation in a finer auxiliary grid (done only once) to improve performance of the interpolation scheme.

$$
\begin{aligned}
& x, y, \tau \mapsto x, y, \tau \\
& p_{x} \mapsto p_{x}-\frac{q L}{\beta_{0} P_{0} c} \frac{\partial \phi}{\partial x}(x, y, \tau) \\
& p_{y} \mapsto p_{y}-\frac{q L}{\beta_{0} P_{0} c} \frac{\partial \phi}{\partial y}(x, y, \tau) \\
& p_{\tau} \mapsto p_{\tau}-\frac{q L}{\beta_{0} P_{0} c} \frac{\partial \phi}{\partial \tau}(x, y, \tau)
\end{aligned}
$$

## General procedure for the simulation



## E-cloud setup

E-cloud exists across the full length of the LHC beam pipe.
Different magnetic fields lead to completely different e-clouds.
Most significant contributors:

1. E-cloud in arc dipoles (MB) $(66 \%)$
2. E-cloud in arc quadrupoles (MQ) (7\%)

We place one interaction for each three dipoles and each quadrupole.


- Betatron and dispersion functions stay the same between each cell.
- Approximate SEY as uniform everywhere. Large fluctuations in reality.
- Effect from saturated e-cloud.


## E-cloud setup



- One MB e-cloud per half-cell
$\rightarrow 46$ interactions per arc
$\rightarrow 368$ interactions.
- One MQ e-cloud per half-cell
$\rightarrow 45$ interactions per arc
$\rightarrow 360$ interactions.

Tracking time per e-cloud type ( $\sim 360$ interactions) is about as much as rest of the lattice (11k tracking elements).

## E-cloud setup



- Dipoles: Reduced bunch intensity leads to larger e- density close to the beam.
- Quadrupoles: Small dependence on bunch intensity, large e- densities close to beam.


## E-cloud setup

Nominal intensity (1.2 $\left.10^{11} \mathrm{p} / \mathrm{bunch}\right)$


Reduced intensity ( $0.610^{11} \mathrm{p} / \mathrm{bunch}$ )




- Dipoles: Reduced bunch intensity leads to larger e- density close to the beam.
- Quadrupoles: Small dependence on bunch intensity, large e- densities close to beam.


## Simulation Parameters

Typical LHC at injection, 2018

> Bunch intensity : $1.2010^{11}$ protons
> Energy : $450 \mathbf{G e V}$

Chromaticity: 15/15
Octupole magnet's current : 40 A
Bunch spacing : 25 ns
Transverse norm emittances : $2 \mu \mathrm{~m} / 2 \mu \mathrm{~m}$
R.M.S. bunch length : 0.09 m

Betatron tunes : 62.270/60.295
RF voltage : 6 MV

The three primary collimators (TCP) in IR7 (as black absorbers) are included in the lattice at their typical configuration (5.7 "collimation" $\sigma \rightarrow 7.5$ beam $\sigma$ ).

There is no uncorrected linear coupling, magnet field imperfections, magnet misalignments or beam-beam interactions in the lattice.

## Secondary Emission Yield (SEY) - Intensity scan




- Larger Secondary electron Emission Yield (of beam pipe) $\rightarrow$ $\rightarrow$ stronger e-cloud $\rightarrow$ less DA
- Dipoles (MB): strong dependence with bunch intensity, correlated to $\mathrm{e}^{-}$ density close to the beam.
- Quadrupoles (MQ): weak dependence with bunch intensity



## Long simulations (10M turns $\rightarrow$ 15min beam time)



Incoherent effects in the LHC are typically very slow processes. Need to simulate long timescales. Recent advances (SixTrackLib/XSuite) allow the direct simulation of particle distributions with GPUs for such times.

Simulation using a V100 GPU took 1 week / 20000 particles / $\mathbf{1 0}$ M turns. Specific study used 6 GPUs at the same time to simulated more particles.

In long term simulations we observe:

- small increase of losses
- horizontal emittance growth,
- vertical emittance growth, when e-clouds are included.
Experimental observations show emittance growth in the same order of magnitude. For quantitative comparisons we have planned dedicated MDs in Run 3.


## Long simulations ( 10 M turns $\rightarrow \mathbf{1 5 m i n}$ beam time)



MB (Dipoles):

- Losses stronger at reduced intensity.
- Emittance growth only at reduced intensity.
- Vertical growth larger than horizontal.

MQ (Quadrupoles):

- Losses across all intensities.
- Emittance growth at all intensities.
- Similar growths in both horizontal and vertical.

Effects strongly correlated with the $\mathrm{e}^{-}$ density close to the beam.

Reminder:

- MB show large densities around the beam for reduced intensities,
- MQ for all intensities.


## Conclusion and Remarks

## Observations:

- Electron cloud in the insertion region quadrupoles is significant. Reduces integrated luminosity.


## Simulations:

- We can do particle tracking simulations with arbitrarily complex e-clouds in arbitrarily complex lattices for millions of turns.
- Simulated simplified scenario at injection energy. Interplay with non-linear magnetic imperfections expected.
- Simulations have reproduced the expected qualitative behavior.
- Very long simulation timescales (several minutes) are in reach. (Using GPUs)


## Outlook for the future:

- Comparison with experimental measurements needs specialized tests. $\rightarrow$ Soon to be carried out in the LHC.
- Simulate scenario during collisions: Strong electron clouds in the Insertion Region quadrupoles + strong beam-beam effects.

Thank you for your attention!
Konstantinos Paraschou

# Backup slides 

## Spare slide



## The RF bucket



- DA simulations done for off-momentum particles $\left(p_{\tau}=5.510^{-4}\right)$.
- FMA simulations done for on-momentum particles ( $p_{\tau}=0$ ).
- Long-term tracking simulations with particles across the full bucket.
- Work in progress: FMA with off-momentum particles.


## Frequency Map Analysis - Nominal intensity ( $0.6 \mathbf{1 0}^{11} \mathrm{p} / \mathrm{b}$ )



## Frequency Map Analysis - Reduced intensity ( $0.6 \mathbf{1 0}^{11} \mathrm{p} / \mathrm{b}$ )






Reduced intensity
Dipoles (MB):
$\rightarrow$ larger tune-shift
$\rightarrow$ more resonances

Quadrupoles (MQ):
$\rightarrow$ large tune-shift
$\rightarrow$ more resonances

Particles are on-momentum, picture is not yet complete. Work in progress to try identify synchro-betatron resonances.

## Tune scan



## Why symplectic?

Symplecticity is a property closely related to Hamiltonian mechanics and the associated integrals of motion. If the numerical method for solving Hamilton's differential equations is not symplectic, e.g. $4^{\text {th }}$ order Runge-Kutta method, quantities which would otherwise stay constant will grow in time.

Consider the Hamiltonian: $H=\frac{p_{1}^{2}}{2}+\frac{p_{2}^{2}}{2}+\phi\left(q_{1}, q_{2}\right)$ with $\phi\left(q_{1}, q_{2}\right)=e^{q_{1}-q_{2}}$
These quantities are conserved: $J_{1}=\left(p_{1}-p_{2}\right)^{2}+4 e^{q_{1}-q_{2}}$, (along with the Hamiltonian)

$$
I_{1}=\frac{p_{1}-p_{2}+\sqrt{J_{1}}}{p_{1}-p_{2}-\sqrt{J_{1}}} \exp \left(\sqrt{J_{1}} \frac{q_{1}+q_{2}}{p_{1}+p_{2}}\right)
$$

We can numerically solve the equations of motion with
the method: $\quad q_{1}^{f}=q_{1}^{i}+p_{1} \cdot \Delta t$
$q_{2}^{f}=q_{2}^{i}+p_{2} \cdot \Delta t$
$p_{1}^{f}=p_{1}^{i}-\frac{\partial \phi}{\partial q_{1}}\left(q_{1}^{f}, q_{2}^{f}\right) \cdot \Delta t$
$p_{2}^{f}=p_{2}^{i}-\frac{\partial \phi}{\partial q_{2}}\left(q_{1}^{f}, q_{2}^{f}\right) \cdot \Delta t$

- The potential is discretized on a grid and the two interpolation methods are used.


## Why symplectic?

Non-symplectic method: Use (bi)linear interpolation on the derivatives

$$
\text { of } \phi\left(q_{1}, q_{2}\right)=e^{q_{1}-q_{2}}
$$

Symplectic method: Use (bi)cubic interpolation on $\phi\left(q_{1}, q_{2}\right)=e^{q_{1}-q_{2}}$.




- The relative error on the integrals of motion does not grow with a symplectic method,
- While it grows for non-symplectic methods.

Impact of tricubic interpolation irregularities

- Simple tracking of linear 2D phase space rotation and an e-cloud symplectic kick.
- Very important to minimize irregularities.
- By reducing them, there is significant impact on the particle

 motion.



## Induced forces

Dipole magnet:


Quadrupole magnet:


