

# Analytic Evaluation of the NNLO virtual corrections to Muon-Electron scattering

Manoj Kumar Mandal  
INFN & University of Padova

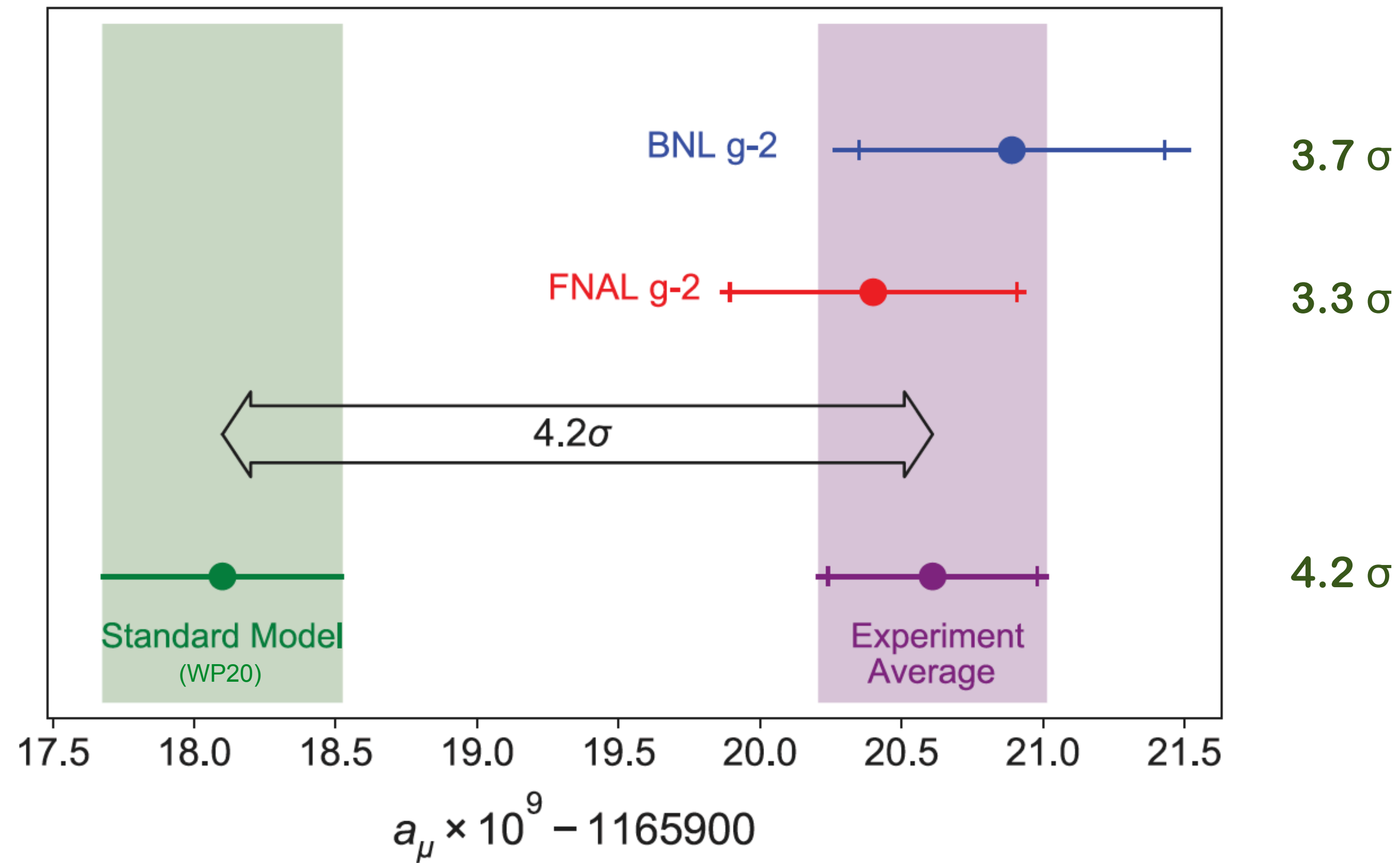
*In collaboration with*

*R. Bonciani, A. Broggio, S. Di Vita, A. Ferroglia, S. Laporta, P. Mastrolia, L. Mattiazzi, M. Passera,  
A. Primo, J. Ronca, U. Schubert, W. J. Torres Bobadilla, and F. Tramontano*

Strong 2020, Remote Seminar  
26th November, 2021



# Motivation



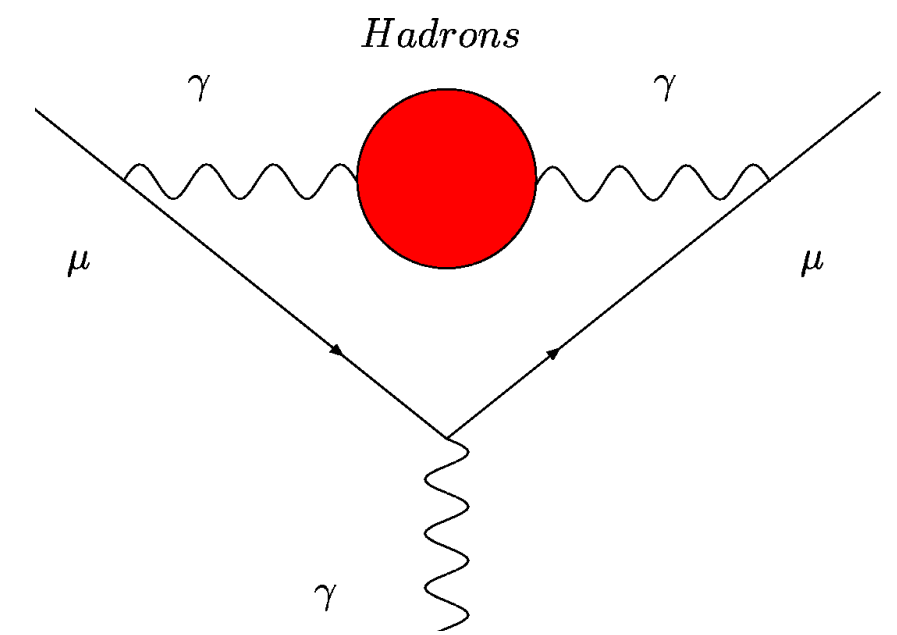
$$a_\mu^{\text{EXP}} = (116592089 \pm 63) \times 10^{-11} [0.54\text{ppm}] \quad \text{BNL E821}$$

$$a_\mu^{\text{EXP}} = (116592040 \pm 54) \times 10^{-11} [0.46\text{ppm}] \quad \text{FNAL E989 Run 1}$$

$$a_\mu^{\text{EXP}} = (116592061 \pm 41) \times 10^{-11} [0.35\text{ppm}] \quad \text{WA}$$

# Motivation:

Leading hadronic contribution computed via the usual dispersive (timelike) formula:



$$a_{\mu}^{\text{HLO}} = \frac{1}{4\pi^3} \int_{m_{\pi}^2}^{\infty} ds K(s) \sigma_{\text{had}}^{(0)}(s)$$

$$K(s) = \int_0^1 dx \frac{x^2 (1-x)}{x^2 + (1-x)(s/m_{\mu}^2)}$$



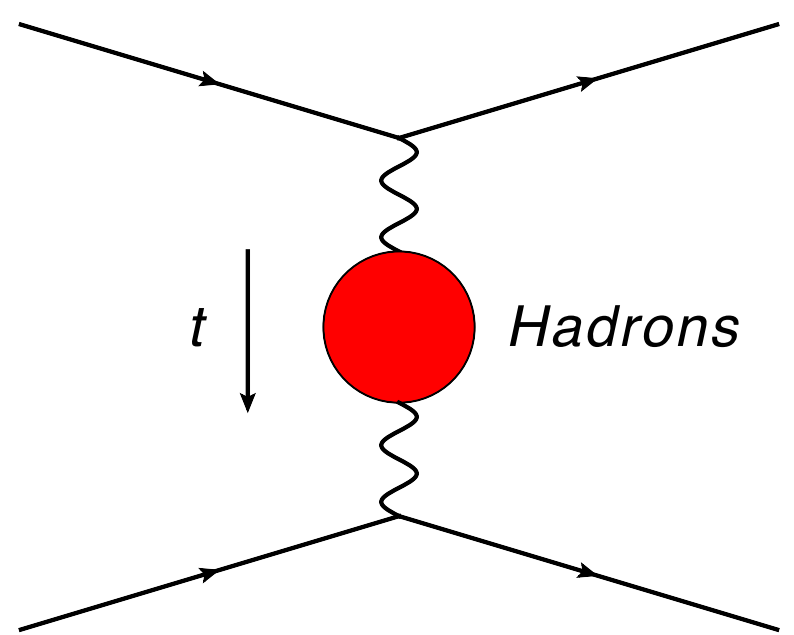
Eur. Phys. J. C (2017) 77:139  
DOI 10.1140/epjc/s10052-017-4633-z

THE EUROPEAN  
PHYSICAL JOURNAL C



Regular Article - Experimental Physics

Alternatively, simply exchanging the x and s integrations:



$$a_{\mu}^{\text{HLO}} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{\text{had}}[t(x)]$$

$$t(x) = \frac{x^2 m_{\mu}^2}{x-1} < 0$$

Lautrup, Peterman, de Rafael, 1972

Measuring the leading hadronic contribution to the muon g-2 via  $\mu e$  scattering

G. Abbiendi<sup>1,a</sup>, C. M. Carloni Calame<sup>2,b</sup>, U. Marconi<sup>3,c</sup>, C. Matteuzzi<sup>4,d</sup>, G. Montagna<sup>2,5,e</sup>, O. Nicrosini<sup>2</sup>, M. Passera<sup>6,g</sup>, F. Piccinini<sup>2,h</sup>, R. Tenchini<sup>7,i</sup>, L. Trentadue<sup>8,4,j</sup>, G. Venanzoni<sup>9,k</sup>

See talk by Stefano Laporta,  
Umberto Marconi, Giovanni Abbiendi

$\Delta\alpha_{\text{had}}(t)$  is the hadronic contribution to the space-like running of  $\alpha$ : **proposal to measure  $a_{\mu}^{\text{HLO}}$  via scattering data!**

# Muon-Electron Scattering @ NNLO

Eur. Phys. J. C (2020) 80:591  
<https://doi.org/10.1140/epjc/s10052-020-8138-9>

THE EUROPEAN  
PHYSICAL JOURNAL C



Review

## Theory for muon-electron scattering @ 10 ppm

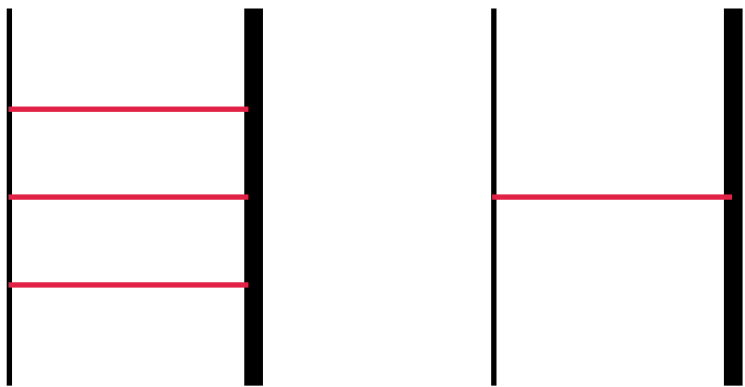
A report of the MUonE theory initiative

P. Banerjee<sup>1</sup>, C. M. Carloni Calame<sup>2</sup>, M. Chiesa<sup>3</sup>, S. Di Vita<sup>4</sup>, T. Engel<sup>1,5</sup>, M. Fael<sup>6</sup>, S. Laporta<sup>7,8</sup>, P. Mastrolia<sup>7,8</sup>, G. Montagna<sup>2,9</sup>, O. Nicrosini<sup>2</sup>, G. Ossola<sup>10</sup>, M. Passera<sup>8</sup>, F. Piccinini<sup>2</sup>, A. Primo<sup>5</sup>, J. Ronca<sup>11</sup>, A. Signer<sup>1,5,a</sup>, W. J. Torres Bobadilla<sup>11</sup>, L. Trentadue<sup>12,13</sup>, Y. Ulrich<sup>1,5</sup>, G. Venanzoni<sup>14</sup>

[Banerjee, Engel, Signer, Ulrich (2020)]  
[Banerjee, Engel, Schalch, Signer, Ulrich (2021)]  
[Budassi, Carloni Calame, Chiesa, Del Pio, Hasan, Montagna, Nicrosini, Piccinini (2021)]  
[Carloni Calame, Chiesa, Hasan, Montagna, Nicrosini, Piccinini (2020)]

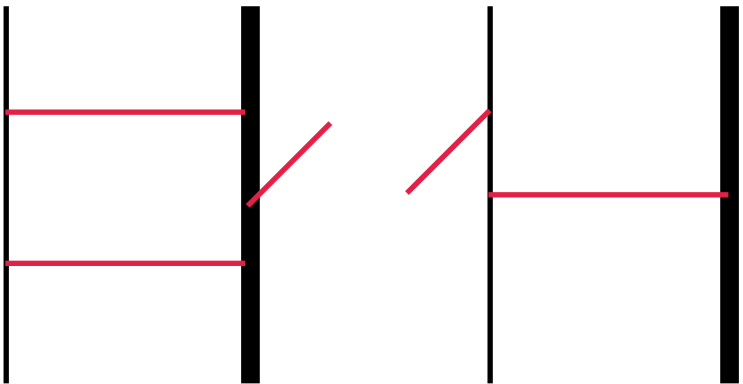
See talk by Ettore Budassi, Tim Engel

$$\int \left[ \frac{VV_4}{\epsilon^4} + \frac{VV_3}{\epsilon^3} + \frac{VV_2}{\epsilon^2} + \frac{VV_1}{\epsilon^1} + VV_0 \right] d\phi_2$$



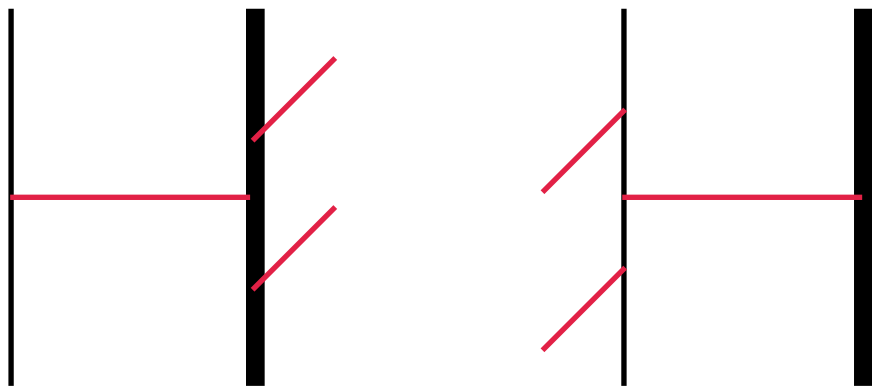
Double Virtual

$$\int \left[ \frac{RV_2}{\epsilon^2} + \frac{RV_1}{\epsilon^1} + RV_0 \right] d\phi_3$$



Real Virtual

$$\int [RR_0] d\phi_4$$



Double Real



# Muon-Electron Scattering @ NNLO

Eur. Phys. J. C (2020) 80:591  
<https://doi.org/10.1140/epjc/s10052-020-8138-9>

THE EUROPEAN  
PHYSICAL JOURNAL C



Review

## Theory for muon-electron scattering @ 10 ppm

A report of the MUonE theory initiative

P. Banerjee<sup>1</sup>, C. M. Carloni Calame<sup>2</sup>, M. Chiesa<sup>3</sup>, S. Di Vita<sup>4</sup>, T. Engel<sup>1,5</sup>, M. Fael<sup>6</sup>, S. Laporta<sup>7,8</sup>, P. Mastrolia<sup>7,8</sup>, G. Montagna<sup>2,9</sup>, O. Nicrosini<sup>2</sup>, G. Ossola<sup>10</sup>, M. Passera<sup>8</sup>, F. Piccinini<sup>2</sup>, A. Primo<sup>5</sup>, J. Ronca<sup>11</sup>, A. Signer<sup>1,5,a</sup>, W. J. Torres Bobadilla<sup>11</sup>, L. Trentadue<sup>12,13</sup>, Y. Ulrich<sup>1,5</sup>, G. Venanzoni<sup>14</sup>

[Banerjee, Engel, Signer, Ulrich (2020)]

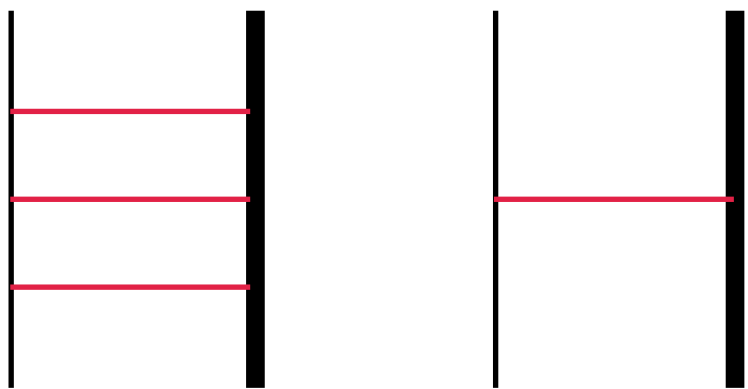
[Banerjee, Engel, Schalch, Signer, Ulrich (2021)]

[Budassi, Carloni Calame, Chiesa, Del Pio, Hasan, Montagna, Nicrosini, Piccinini (2021)]

[Carloni Calame, Chiesa, Hasan, Montagna, Nicrosini, Piccinini (2020)]

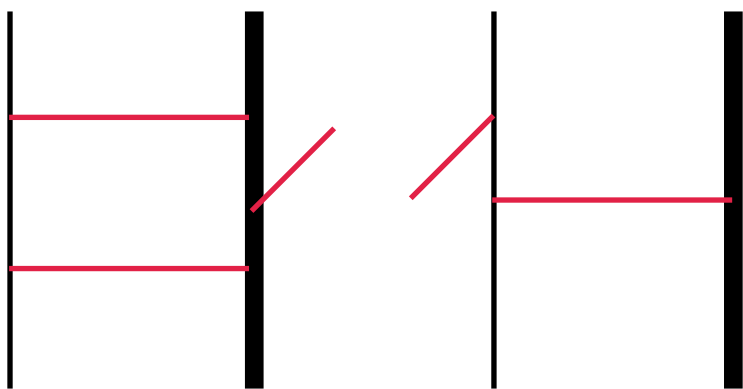
See talk by Ettore Budassi, Tim Engel

$$\int \left[ \frac{VV_4}{\epsilon^4} + \frac{VV_3}{\epsilon^3} + \frac{VV_2}{\epsilon^2} + \frac{VV_1}{\epsilon^1} + VV_0 \right] d\phi_2$$



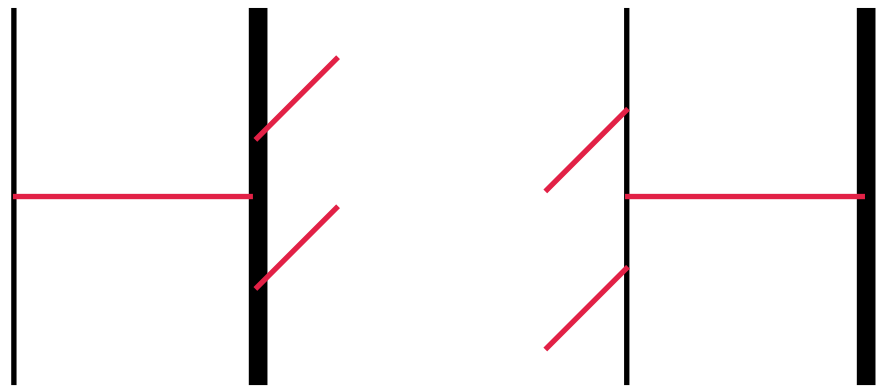
Double Virtual

$$\int \left[ \frac{RV_2}{\epsilon^2} + \frac{RV_1}{\epsilon^1} + RV_0 \right] d\phi_3$$



Real Virtual

$$\int [RR_0] d\phi_4$$



Double Real

This talk is towards the computation of the complete 2-Loop Virtual Amplitude

# Muon-Electron Scattering @ NNLO

Eur. Phys. J. C (2020) 80:591  
<https://doi.org/10.1140/epjc/s10052-020-8138-9>

THE EUROPEAN  
PHYSICAL JOURNAL C



Review

## Theory for muon-electron scattering @ 10 ppm

A report of the MUonE theory initiative

P. Banerjee<sup>1</sup>, C. M. Carloni Calame<sup>2</sup>, M. Chiesa<sup>3</sup>, S. Di Vita<sup>4</sup>, T. Engel<sup>1,5</sup>, M. Fael<sup>6</sup>, S. Laporta<sup>7,8</sup>, P. Mastrolia<sup>7,8</sup>, G. Montagna<sup>2,9</sup>, O. Nicrosini<sup>2</sup>, G. Ossola<sup>10</sup>, M. Passera<sup>8</sup>, F. Piccinini<sup>2</sup>, A. Primo<sup>5</sup>, J. Ronca<sup>11</sup>, A. Signer<sup>1,5,a</sup>, W. J. Torres Bobadilla<sup>11</sup>, L. Trentadue<sup>12,13</sup>, Y. Ulrich<sup>1,5</sup>, G. Venanzoni<sup>14</sup>

[Banerjee, Engel, Signer, Ulrich (2020)]

[Banerjee, Engel, Schalch, Signer, Ulrich (2021)]

[Budassi, Carloni Calame, Chiesa, Del Pio, Hasan, Montagna, Nicrosini, Piccinini (2021)]

[Carloni Calame, Chiesa, Hasan, Montagna, Nicrosini, Piccinini (2020)]

See talk by Ettore Budassi, Tim Engel

$$\int \left[ \frac{VV_4}{\epsilon^4} + \frac{VV_3}{\epsilon^3} + \frac{VV_2}{\epsilon^2} + \frac{VV_1}{\epsilon^1} + VV_0 \right] d\phi_2$$

$$\int \left[ \frac{RV_2}{\epsilon^2} + \frac{RV_1}{\epsilon^1} + RV_0 \right] d\phi_3$$

$$\int [RR_0] d\phi_4$$

## The two-loop four-fermion scattering amplitude in QED

R. Bonciani,<sup>1,\*</sup> A. Broggio,<sup>2,†</sup> S. Di Vita,<sup>3,4</sup> A. Ferroglia,<sup>5,6,‡</sup> M. K. Mandal,<sup>7,8,§</sup> P. Mastrolia,<sup>8,7,¶</sup> L. Mattiazzi,<sup>7,8,||</sup> A. Primo,<sup>9,\*\*</sup> J. Ronca,<sup>10,††</sup> U. Schubert,<sup>11,‡‡</sup> W. J. Torres Bobadilla,<sup>12,§§</sup> and F. Tramontano<sup>10,¶¶</sup>

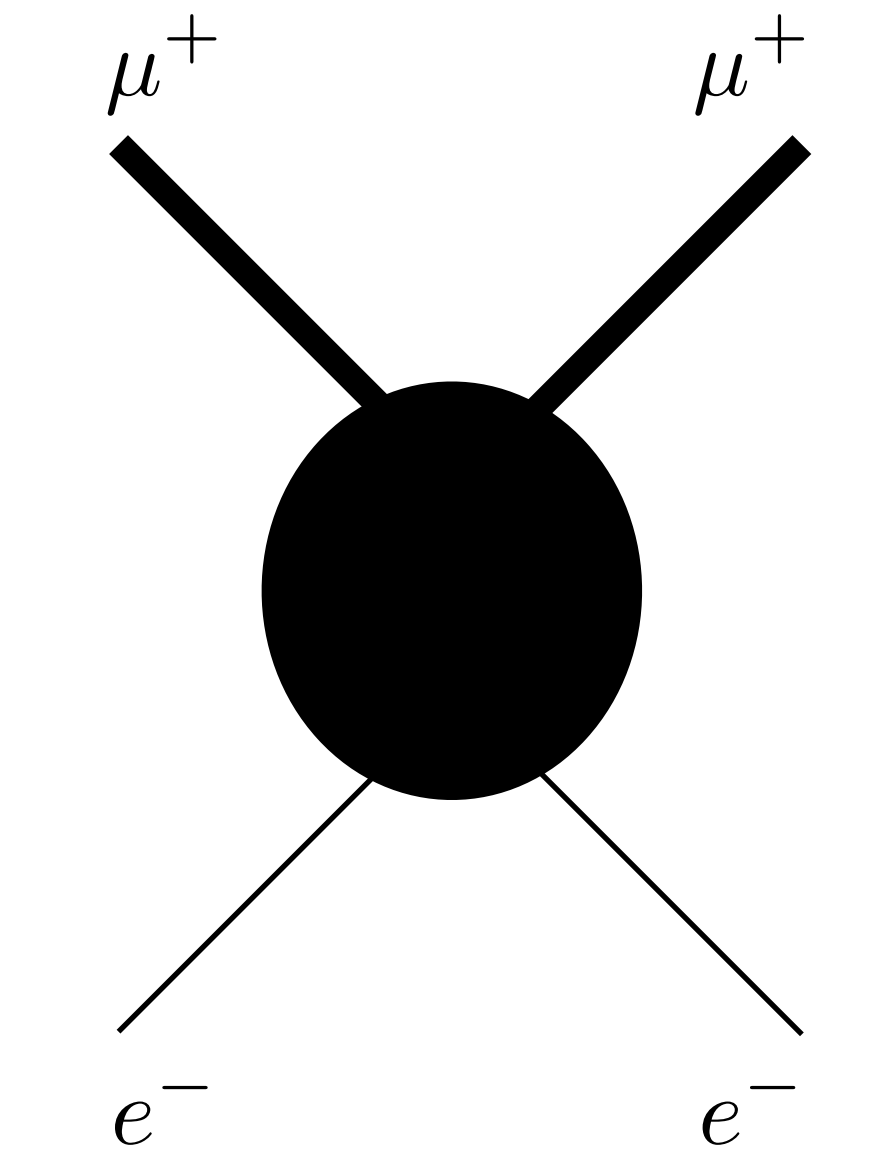
Real Virtual

Double Real

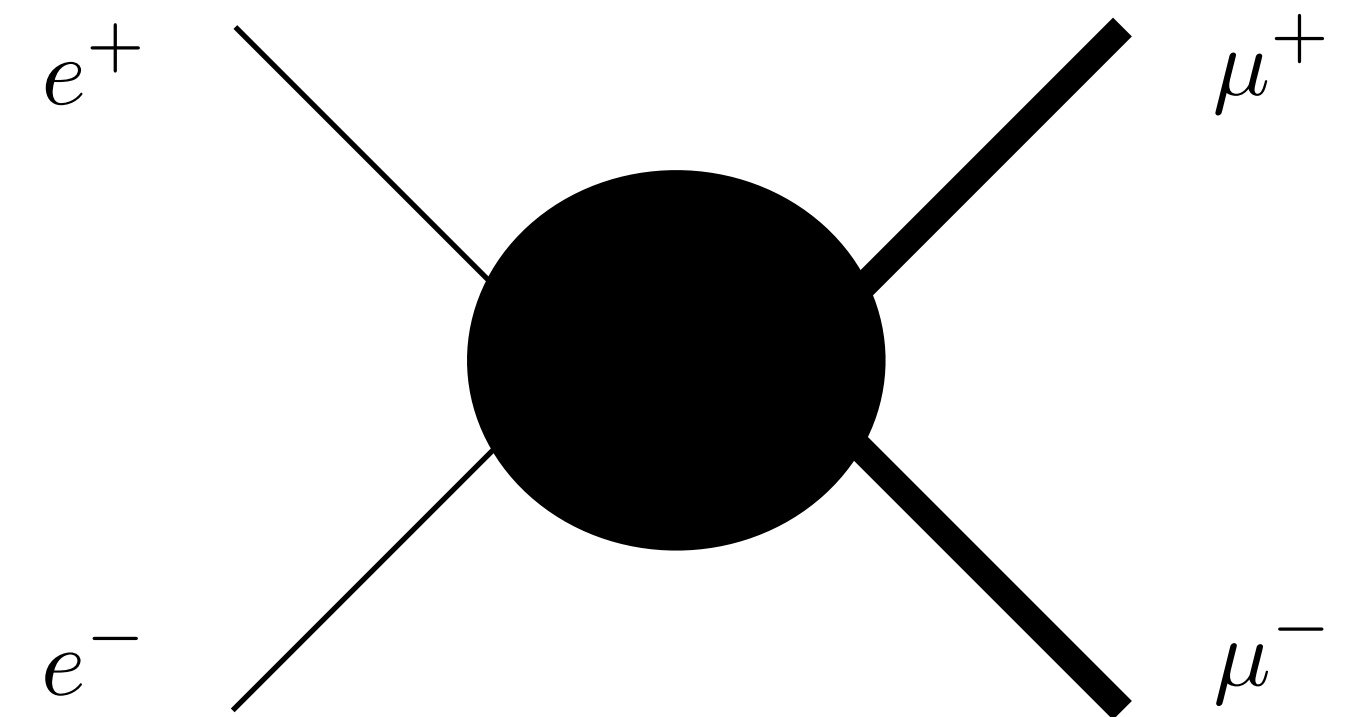
This talk is towards the computation of the complete 2-Loop Virtual Amplitude

# Di-Muon Production

[Bonciani, Broggio, Di Vita, Ferroglia, MKM, Mastrolia, Mattiazzi, Primo, Ronca, Schubert, Torres Bobadilla, Tramontano (2021)]



$$e^- + \mu^+ \rightarrow e^- + \mu^+$$



$$e^- + e^+ \rightarrow \mu^- + \mu^+$$

Crossing

# Amplitude for Di-muon Production

[Bonciani, Broggio, Di Vita, Ferroglia, MKM, Mastrolia, Mattiazzi, Primo, Ronca, Schubert, Torres Bobadilla, Tramontano (2021)]

$$e^-(p_1) + e^+(p_2) \rightarrow \mu^-(p_3) + \mu^+(p_4)$$

$$m_e = 0$$

$$m_\mu = M$$

$$s = (p_1 + p_2)^2$$

$$t = (p_1 - p_3)^2$$

$$u = (p_2 - p_3)^2$$

$$s + t + u = 2M^2$$

$$\mathcal{A}_b(\alpha_b) = 4\pi\alpha_b S_\epsilon \mu^{-2\epsilon} \left[ \mathcal{A}_b^{(0)} + \left(\frac{\alpha_b}{\pi}\right) \mathcal{A}_b^{(1)} + \left(\frac{\alpha_b}{\pi}\right)^2 \mathcal{A}_b^{(2)} \right]$$

Bare Amplitude

LO

NLO

NNLO

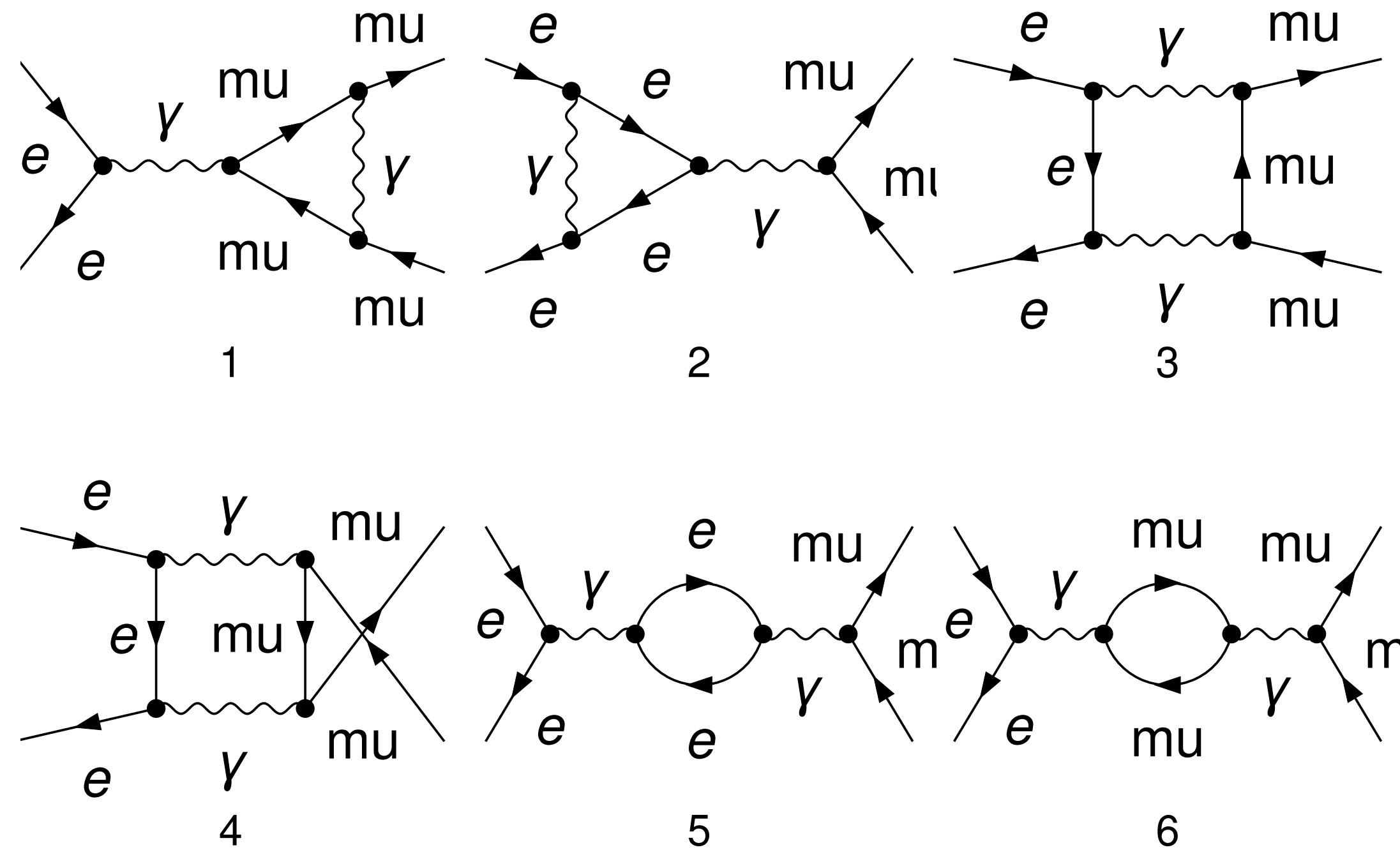
$$\mathcal{M}_b^{(0)} = \frac{1}{4} \sum_{\text{spins}} |\mathcal{A}_b^{(0)}|^2 = \frac{1}{s^2} [2(1 - \epsilon)s^2 + 4(t - M^2)^2 + 4st]$$

$$\mathcal{M}_b^{(n)} = \frac{1}{4} \sum_{\text{spins}} 2 \text{Re}(\mathcal{A}_b^{(0)*} \mathcal{A}_b^{(n)})$$

# 1-Loop Diagrams

6 Diagrams

$$\mathcal{M}^{(1)} = A^{(1)} + n_l B_l^{(1)} + n_h C_h^{(1)}$$



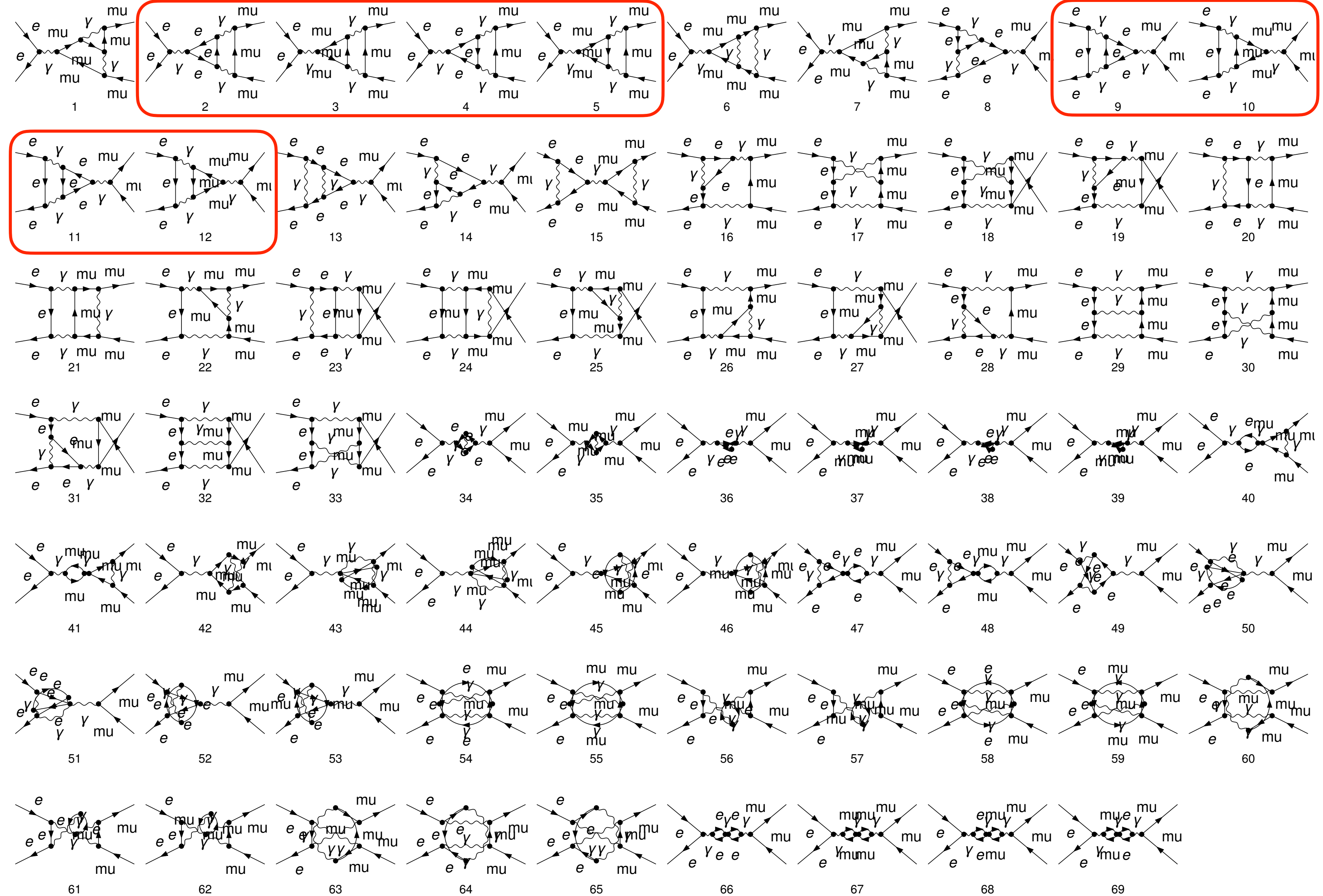


# 2-Loop Diagrams

69 Diagrams

0 Due to Furry's Theorem

$$\mathcal{M}^{(2)} = A^{(2)} + n_l B_l^{(2)} + n_h C_h^{(2)} + n_l^2 D_l^{(2)} \\ + n_h n_l E_{hl}^{(2)} + n_h^2 F_h^{(2)}$$



# Computation of the Loop Amplitude

Mathematica Based Package AIDA

[Mastrolia, Peraro, Primo, Ronca, Torres Bobadilla (To be Published) ]

Generation of Diagram by FeynArts



Spin sums, Dirac Algebra, Trace by FeynCalc



Adaptive Integrand Decomposition



IBP Reduction via Reduze and KIRA



Master Integral evaluation



$$\mathcal{M}_b^{(n)} = (S_\epsilon)^n \int \prod_{i=1}^n \frac{d^d k_i}{(2\pi)^d} \sum_G \frac{N_G}{\prod_{\sigma \in G} D_\sigma}$$

$$\mathcal{M}_b^{(n)} = \mathbb{C}^{(n)} \cdot \mathbf{I}^{(n)}$$

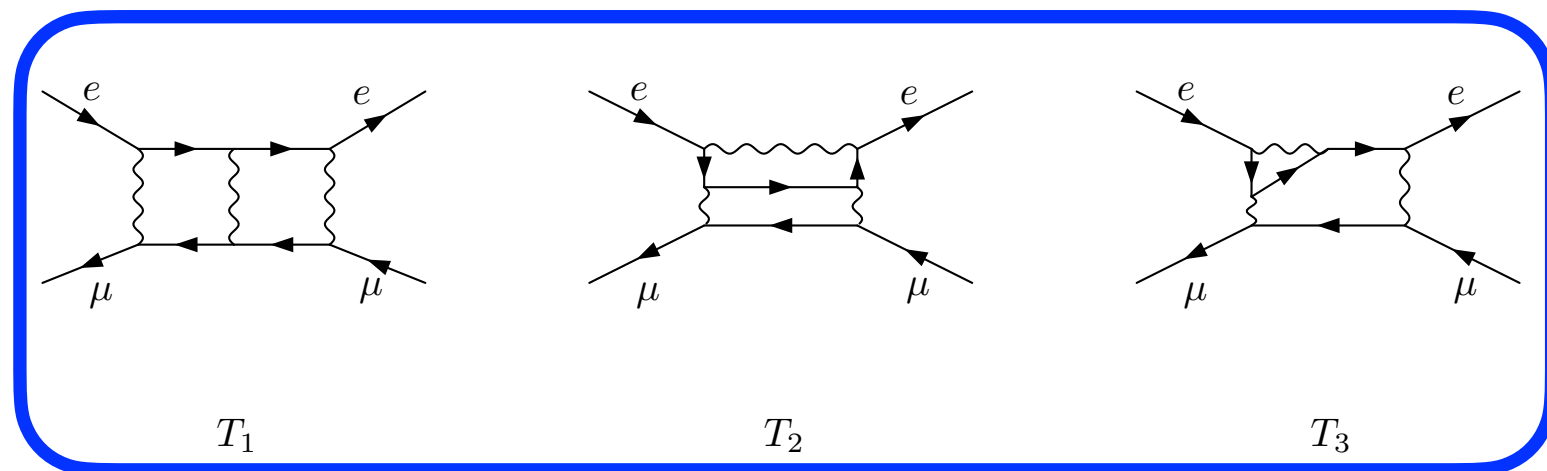
Master Integrals

# Master Integrals

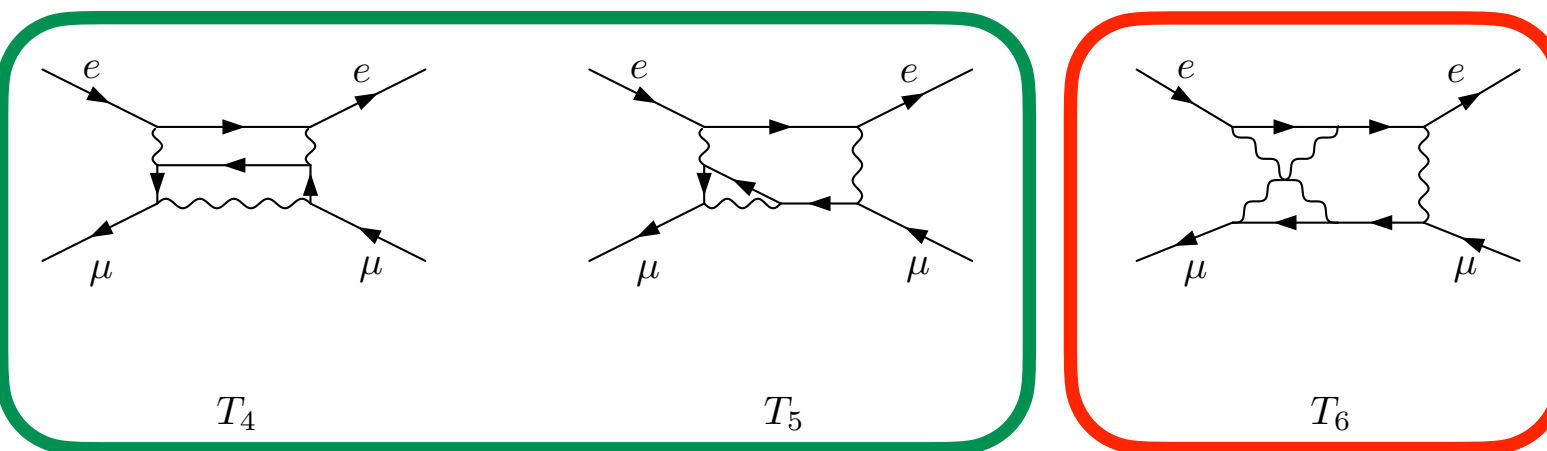
- ◆ The differential equation method has been the most successful in the computation of the MIs

[Kotikov (1990)] [Gehrmann, Remiddi (1999)]

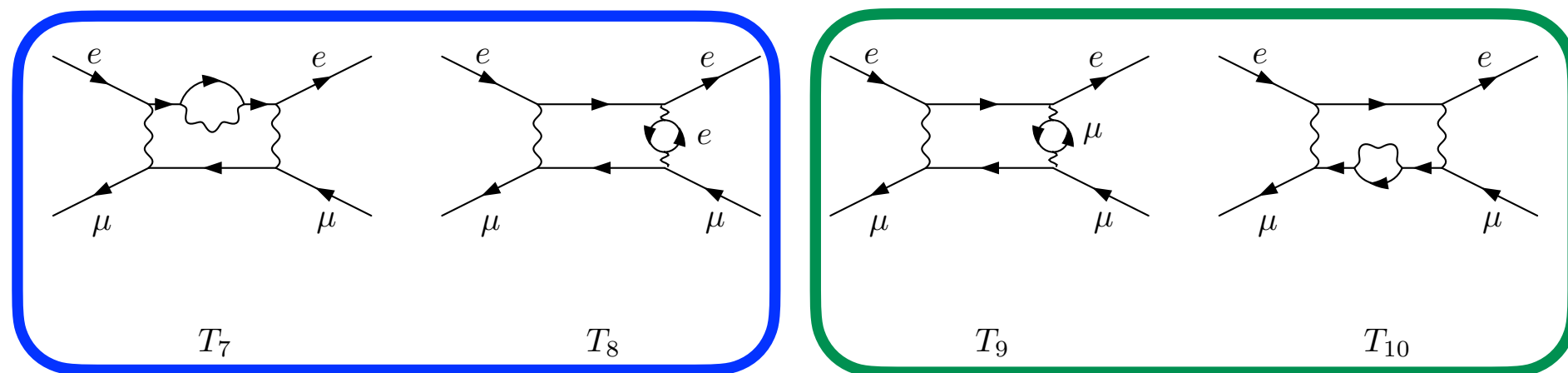
[Henn (2013)] [Argeri, Di Vita, Mastrolia, Mirabella, Schlenk, Schubert, Tancredi (2014)]



33 MIs



44 MIs



42 MIs

[Bonciani, Ferroglia, Gehrmann, von Manteuffel (2008-13)]

[Mastrolia, Passera, Primo, Schubert (2017)]

[Di Vita, Laporta, Mastrolia, Primo, Schubert (2018)]

Generalized Polylogarithms

$$G(w_n, \dots, w_1; \tau) \equiv \int_0^\tau \frac{dt}{t - w_n} G(w_{n-1}, \dots, w_1; t)$$

$$G(w_1; t) \equiv \log(1 - t/w_1)$$

# UV Renormalization

Wave functions of External particles - Onshell

$$\psi_b = \sqrt{Z_2} \psi, \quad A_b^\mu = \sqrt{Z_3} A^\mu, \quad M_b = Z_M M$$

Mass of the muon - Onshell

Renormalized QED Vertex

$$\mathcal{L}_{\text{int}} = e_b \bar{\psi}_b \not{A}_b \psi_b = e Z_1 \bar{\psi} \not{A} \psi$$

$$e Z_1 = e_b Z_2 \sqrt{Z_3}$$

$$Z_1 = Z_2 \quad \text{QED Ward Identity}$$

Coupling constant -  $\overline{\text{MS}}$        $\alpha_b \equiv e_b^2/4\pi$        $\alpha_b S_\epsilon = \alpha(\mu^2) \mu^{2\epsilon} Z_\alpha^{\overline{\text{MS}}}$

$$\mathcal{A} = Z_{2,f} Z_{2,F} \mathcal{A}_b (\alpha_b = \alpha_b(\alpha), M_b = M_b(M))$$

$$Z_j = 1 + \left(\frac{\alpha}{\pi}\right) \delta Z_j^{(1)} + \left(\frac{\alpha}{\pi}\right)^2 \delta Z_j^{(2)} + \mathcal{O}(\alpha^3)$$

# UV Renormalization

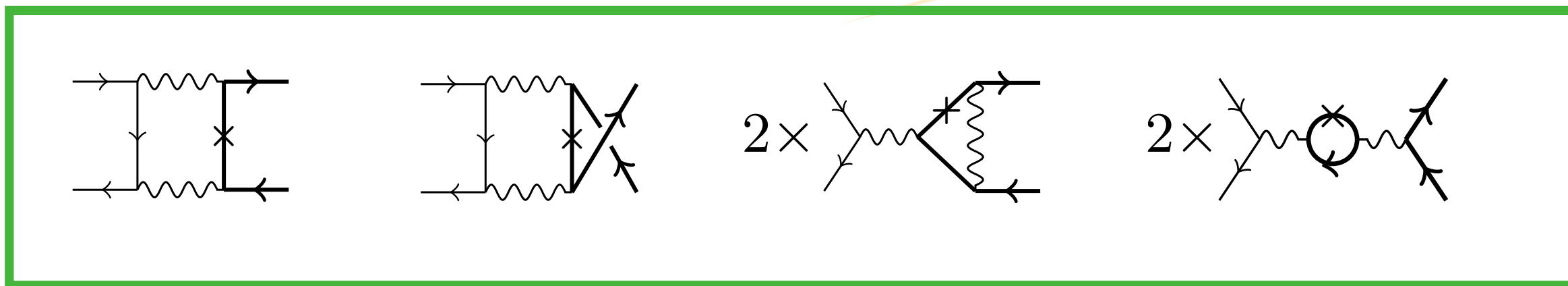
## Renormalized Amplitude

$$\mathcal{A}(\alpha) = 4\pi\alpha \left[ \mathcal{A}^{(0)} + \left(\frac{\alpha}{\pi}\right) \mathcal{A}^{(1)} + \left(\frac{\alpha}{\pi}\right)^2 \mathcal{A}^{(2)} \right]$$

$$\mathcal{A}^{(0)} = \mathcal{A}_b^{(0)}$$

$$\mathcal{A}^{(1)} = \mathcal{A}_b^{(1)} + \left( \delta Z_\alpha^{(1)} + \delta Z_F^{(1)} \right) \mathcal{A}_b^{(0)}$$

$$\begin{aligned} \mathcal{A}^{(2)} = & \mathcal{A}_b^{(2)} + \left( 2\delta Z_\alpha^{(1)} + \delta Z_F^{(1)} \right) \mathcal{A}_b^{(1)} \\ & + \left( \delta Z_\alpha^{(2)} + \delta Z_F^{(2)} + \delta Z_f^{(2)} + \delta Z_F^{(1)} \delta Z_\alpha^{(1)} \right) \mathcal{A}_b^{(0)} \\ & + \delta Z_M^{(1)} \mathcal{A}_b^{(1, \text{mass CT})} \end{aligned}$$



$$\delta Z_f^{(2)} = n_h \left( \frac{L_\mu}{8} + \frac{1}{16\epsilon} - \frac{5}{96} \right)$$

[Czakon, Mitov, Moch (2007)]



# IR Factorization

$$\mathcal{M}^{(1)} \Big|_{\text{poles}} = \frac{1}{2} Z_1^{\text{IR}} \mathcal{M}^{(0)} \Big|_{\text{poles}}$$

$$\mathcal{M}^{(2)} \Big|_{\text{poles}} = \frac{1}{8} \left[ \left( Z_2^{\text{IR}} - (Z_1^{\text{IR}})^2 \right) \mathcal{M}^{(0)} + 2 Z_1^{\text{IR}} \mathcal{M}^{(1)} \right] \Big|_{\text{poles}}$$

[Becher, Neubert (2009)]

[Hill (2017)]

## IR Renormalization Factor

$$\ln Z_{\text{IR}} = \frac{\alpha}{4\pi} \left( \frac{\Gamma'_0}{4\epsilon^2} + \frac{\Gamma_0}{2\epsilon} \right) + \left( \frac{\alpha}{4\pi} \right)^2 \left( -\frac{3\beta_0\Gamma'_0}{16\epsilon^3} + \frac{\Gamma'_1 - 4\beta_0\Gamma_0}{16\epsilon^2} + \frac{\Gamma_1}{4\epsilon} \right) + \mathcal{O}(\alpha^3)$$

$$\Gamma'(\alpha) \equiv \frac{\partial}{\partial \ln \mu} \Gamma(\alpha) = \sum_{n=0}^{\infty} \Gamma'_i \left( \frac{\alpha}{4\pi} \right)^{n+1}$$

$$\Gamma = \gamma_{\text{cusp}}(\alpha) \ln \left( -\frac{s}{\mu^2} \right) + 2\gamma_{\text{cusp}}(\alpha) \ln \left( \frac{t - M^2}{u - M^2} \right) + \gamma_{\text{cusp},M}(\alpha, s) + 2\gamma_h(\alpha) + 2\gamma_\psi(\alpha)$$

Cusp Anomalous dimension

Anomalous dimension

# Numerical Result

$$\mathcal{M}^{(0)} = \frac{1}{s^2} [2(1 - \epsilon)s^2 + 4(t - M^2)^2 + 4st]$$

$$\mathcal{M}^{(1)} = A^{(1)} + n_l B_l^{(1)} + n_h C_h^{(1)}$$

$$\mathcal{M}^{(2)} = A^{(2)} + n_l B_l^{(2)} + n_h C_h^{(2)} + n_l^2 D_l^{(2)} + n_h n_l E_{hl}^{(2)} + n_h^2 F_h^{(2)}$$

- ◆ At 2 Loop there are 4063 GPLs up to weight 4

$$G(w_n, \dots, w_1; \tau) \equiv \int_0^\tau \frac{dt}{t - w_n} G(w_{n-1}, \dots, w_1; t)$$

- ◆ 18 Letters  $w_i = w_i(x, y, z)$

$$\begin{aligned} -t/M^2 &= x \\ -s/M^2 &= (1 - y)^2/y \\ -(u - M^2)/(t - M^2) &= z^2/y \end{aligned}$$

- ◆ The GPLs are evaluated by Ginac [PolyLogTools interface] and HandyG

[Vollinga, Weinzierl (2004)] [Duhr, Dulat (2019)]

[Naterop, Signer, Y. Ulrich (2019)]

☑ We have obtained complete agreement between the predicted IR poles and the 2-Loop UV renormalized amplitude

# Further checks

Kinematical point  $s/M^2 = 5, t/M^2 = -5/4, \mu = M$

	$\epsilon^{-4}$	$\epsilon^{-3}$	$\epsilon^{-2}$	$\epsilon^{-1}$	$\epsilon^0$	$\epsilon$
$\mathcal{M}^{(0)}$	-	-	-	-	$\frac{181}{100}$	-2
$A^{(1)}$	-	-	$-\frac{181}{100}$	1.99877525	22.0079572	-11.7311017
$B_l^{(1)}$	-	-	-	-	-0.069056030	4.94328573
$C_h^{(1)}$	-	-	-	-	-2.24934027	2.54943566
$A^{(2)}$	$\frac{181}{400}$	-0.499387626	-35.4922919	19.4997261	48.8842283	-
$B_l^{(2)}$	-	$-\frac{181}{400}$	0.785712779	-16.1576674	-3.75247701	-
$C_h^{(2)}$	-	-	1.12467013	-9.50785825	-25.8771503	-
$D_l^{(2)}$	-	-	-	-	-3.96845688	-
$E_{hl}^{(2)}$	-	-	-	-	-4.88512563	-
$F_h^{(2)}$	-	-	-	-	-0.158490810	-

[Czakon(2008)]

☑ We recovered the Abelian part of the QCD result of  $q\bar{q} \rightarrow t\bar{t}$  at 1- and 2-Loop [Bonciani, Ferroglia, Gehrmann, Maitre, Studerus (2008)]

[Bärnreuther, Czakon, Fiedler (2014)]

☑  $n_h$  contributions were checked independently [Fael, Passera (2019)]  
[Fael (2018)]

# Results: One Loop

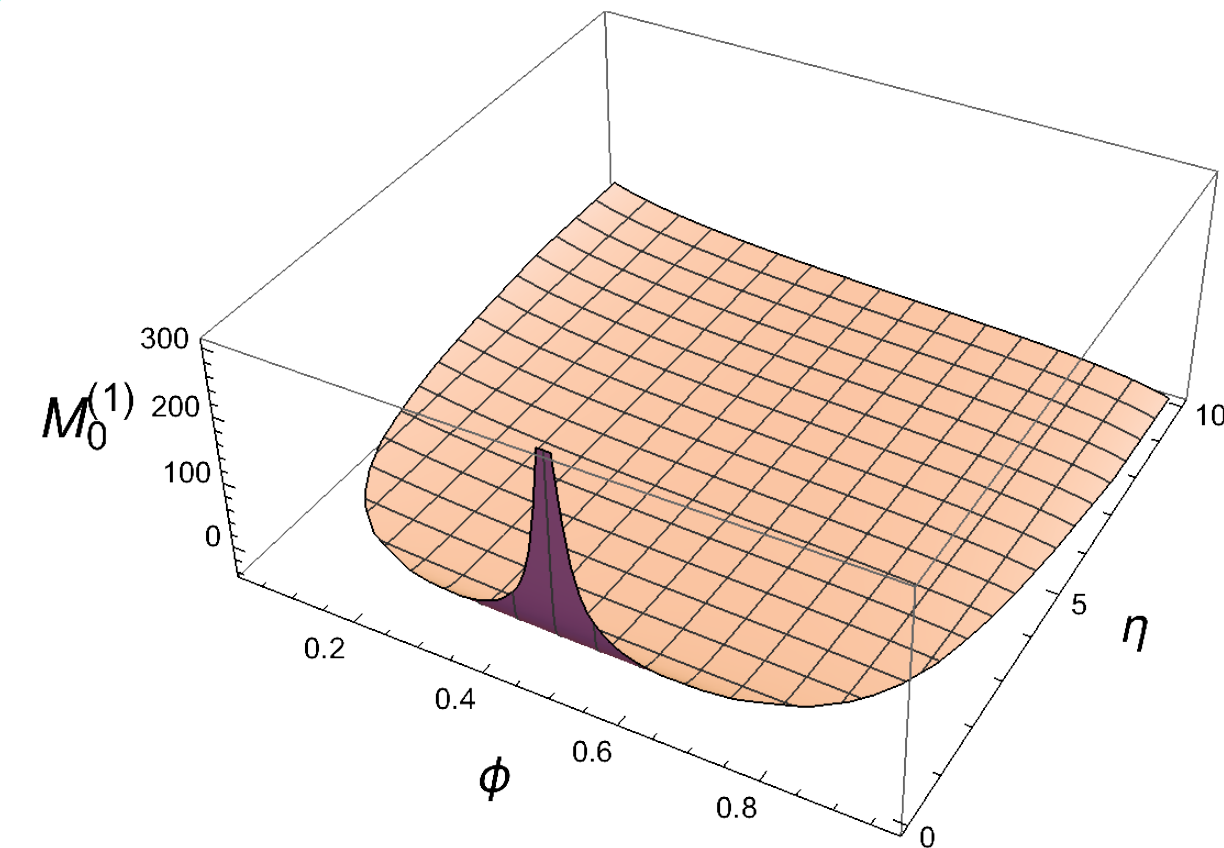
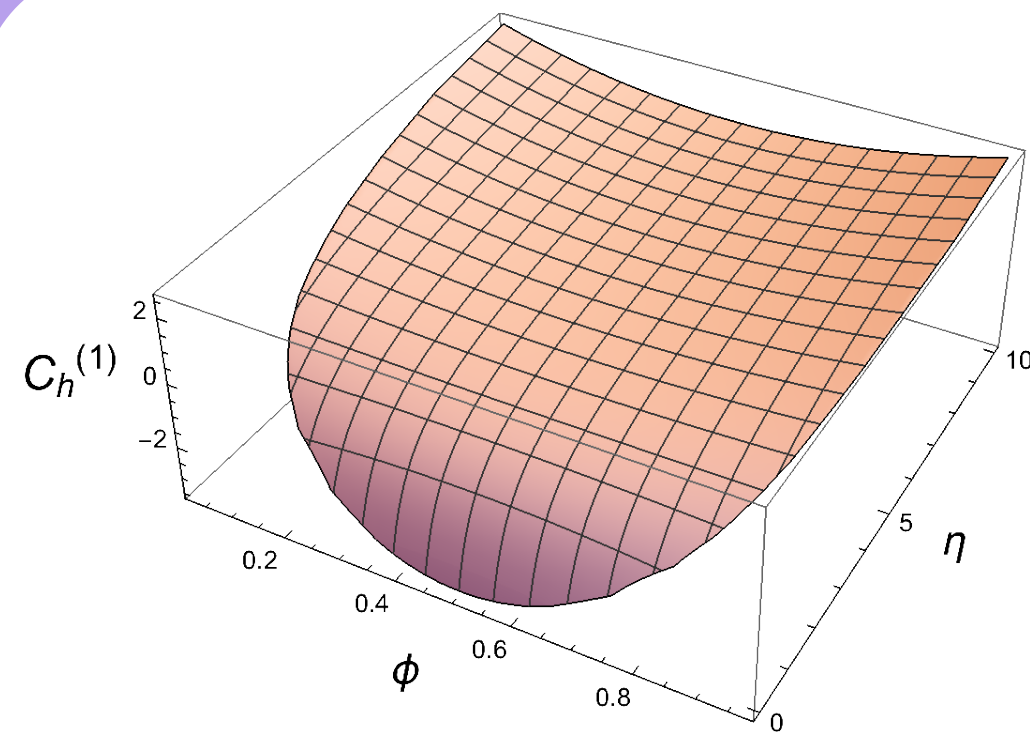
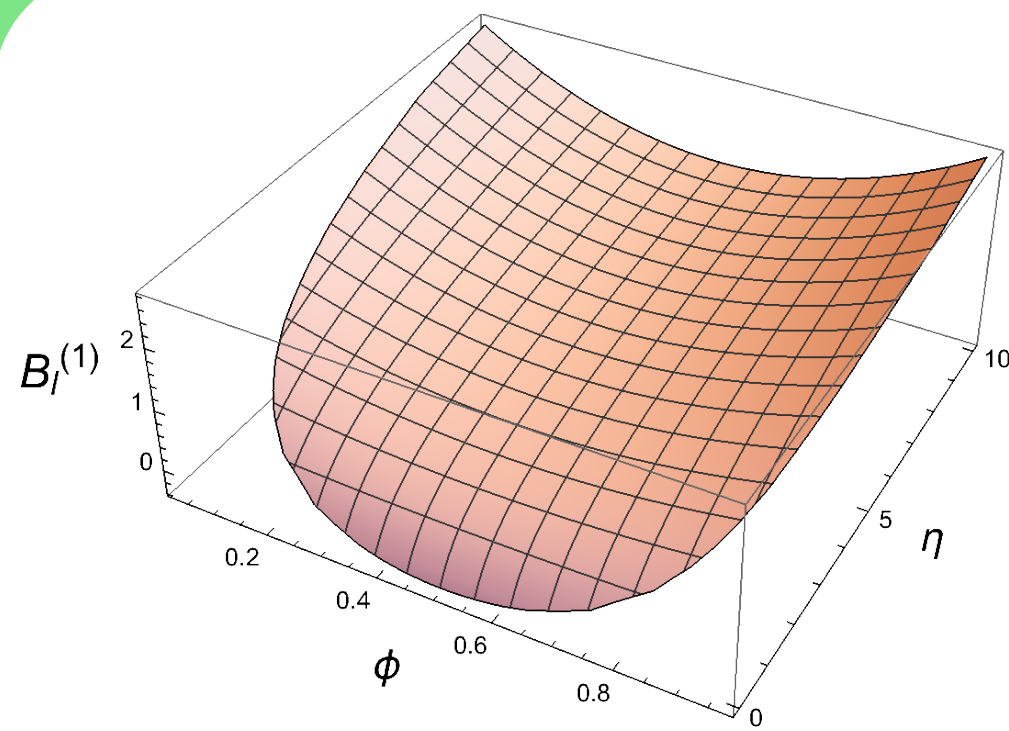
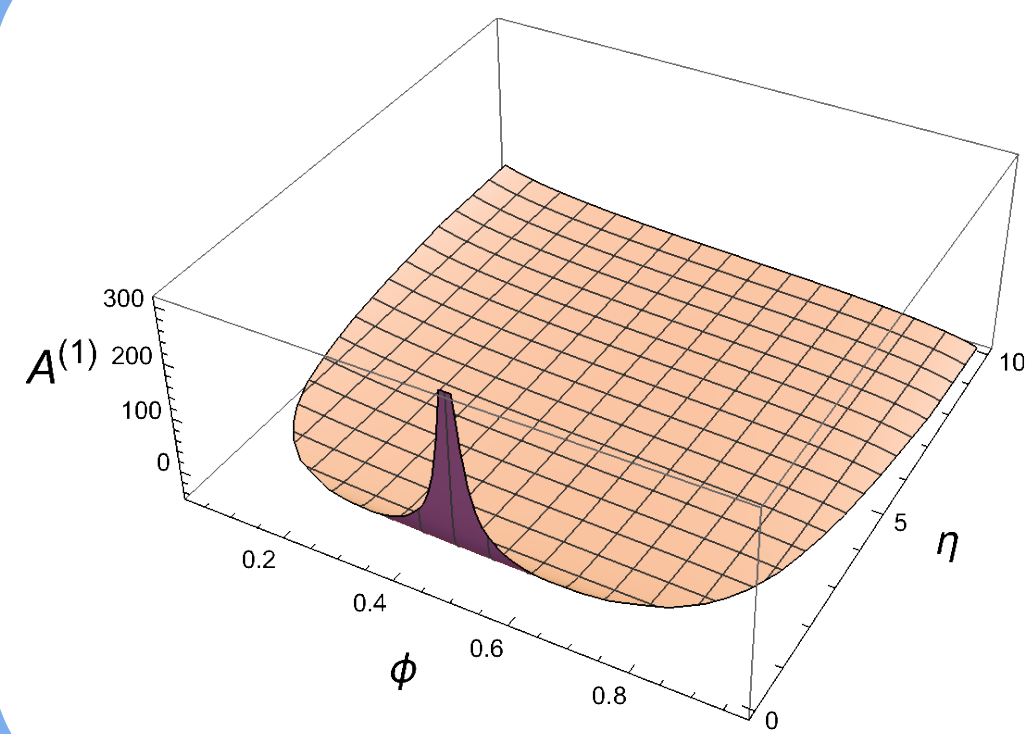
$$\mathcal{M}^{(1)} = A^{(1)} + n_l B_l^{(1)} + n_h C_h^{(1)}$$

$$\mathcal{M}^{(1)} = \frac{\mathcal{M}_{-2}^{(1)}}{\epsilon^2} + \frac{\mathcal{M}_{-1}^{(1)}}{\epsilon} + \mathcal{M}_0^{(1)}$$

$$\eta = s/(4M^2) - 1$$

$$\phi = -(t - M^2)/s$$

$$\frac{1}{2} \left( 1 - \sqrt{\frac{\eta}{1+\eta}} \right) \leq \phi \leq \frac{1}{2} \left( 1 + \sqrt{\frac{\eta}{1+\eta}} \right)$$

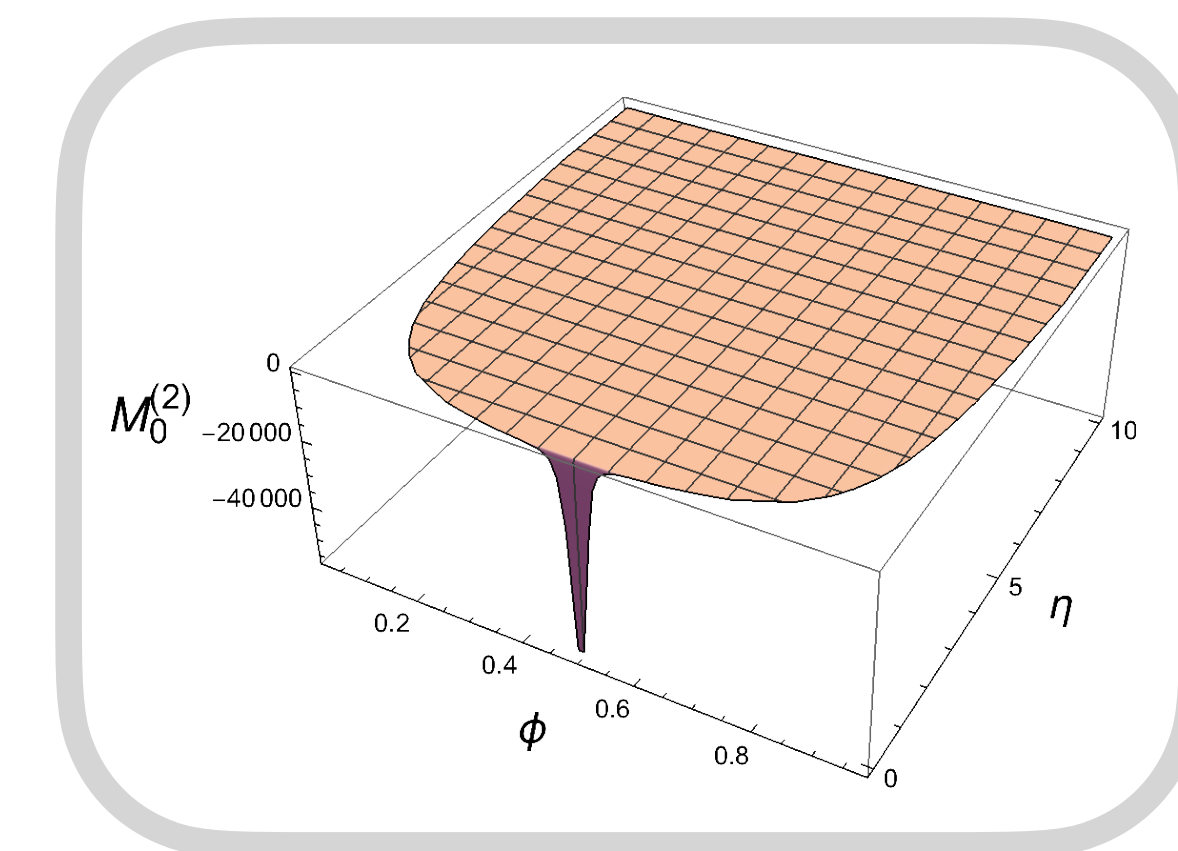
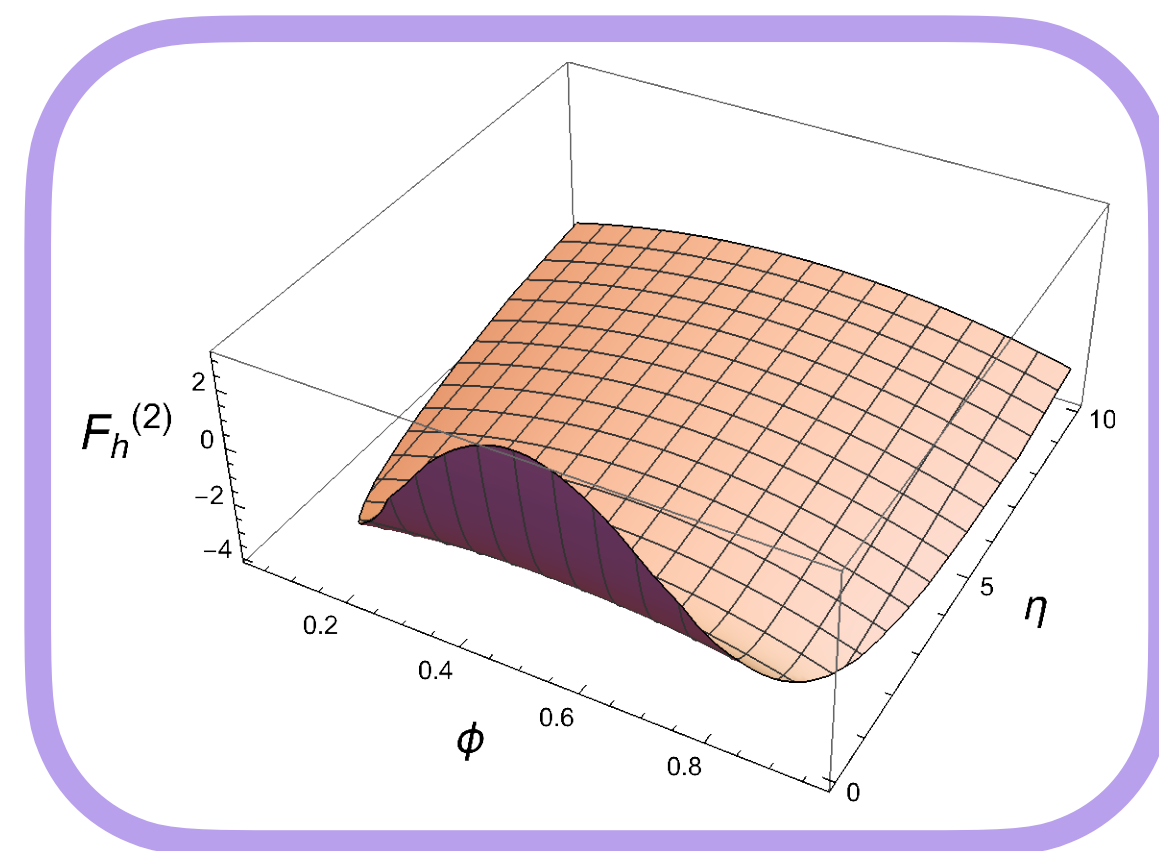
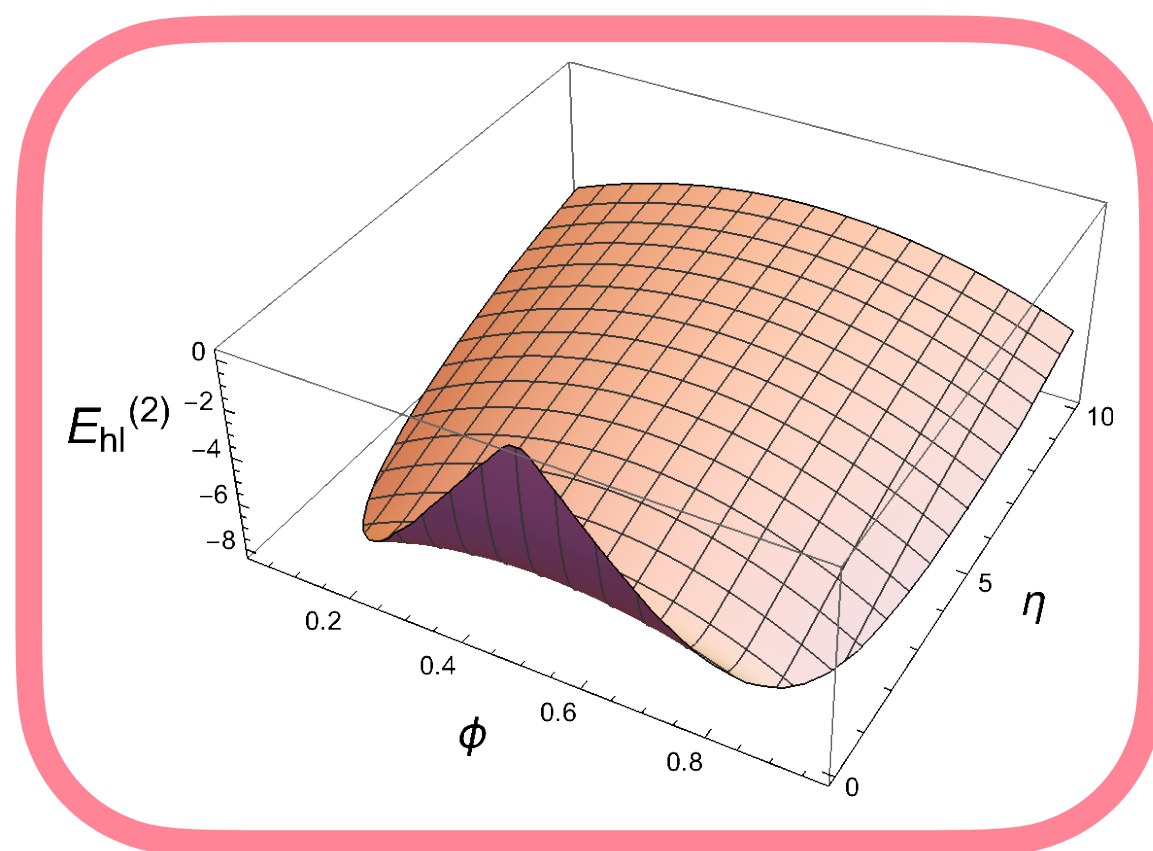
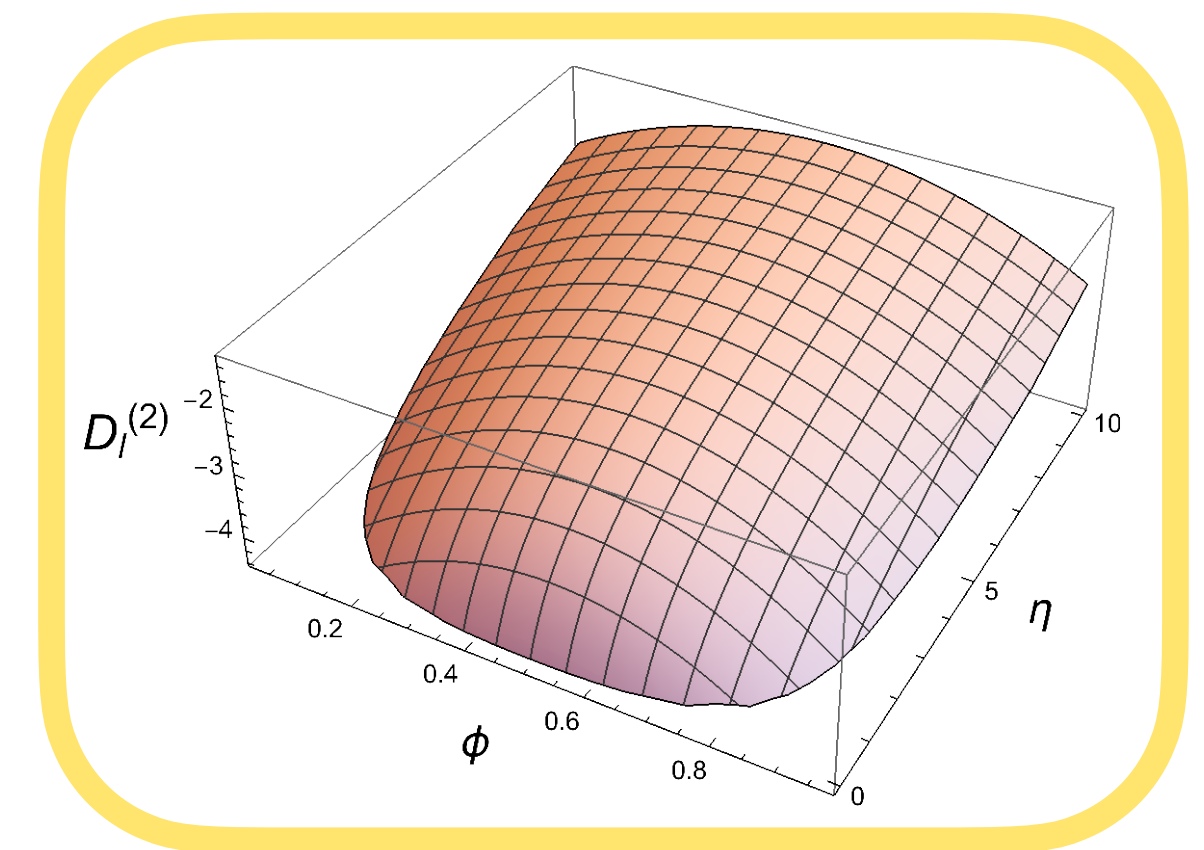
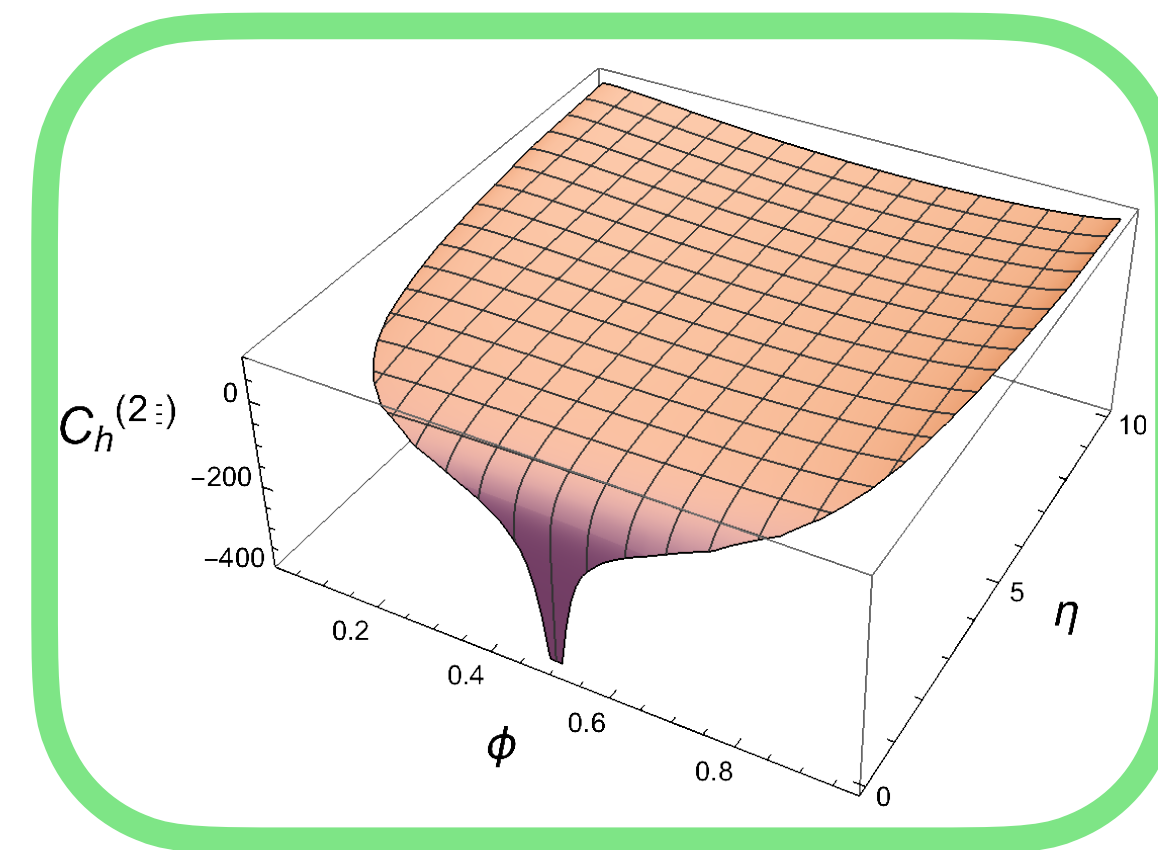
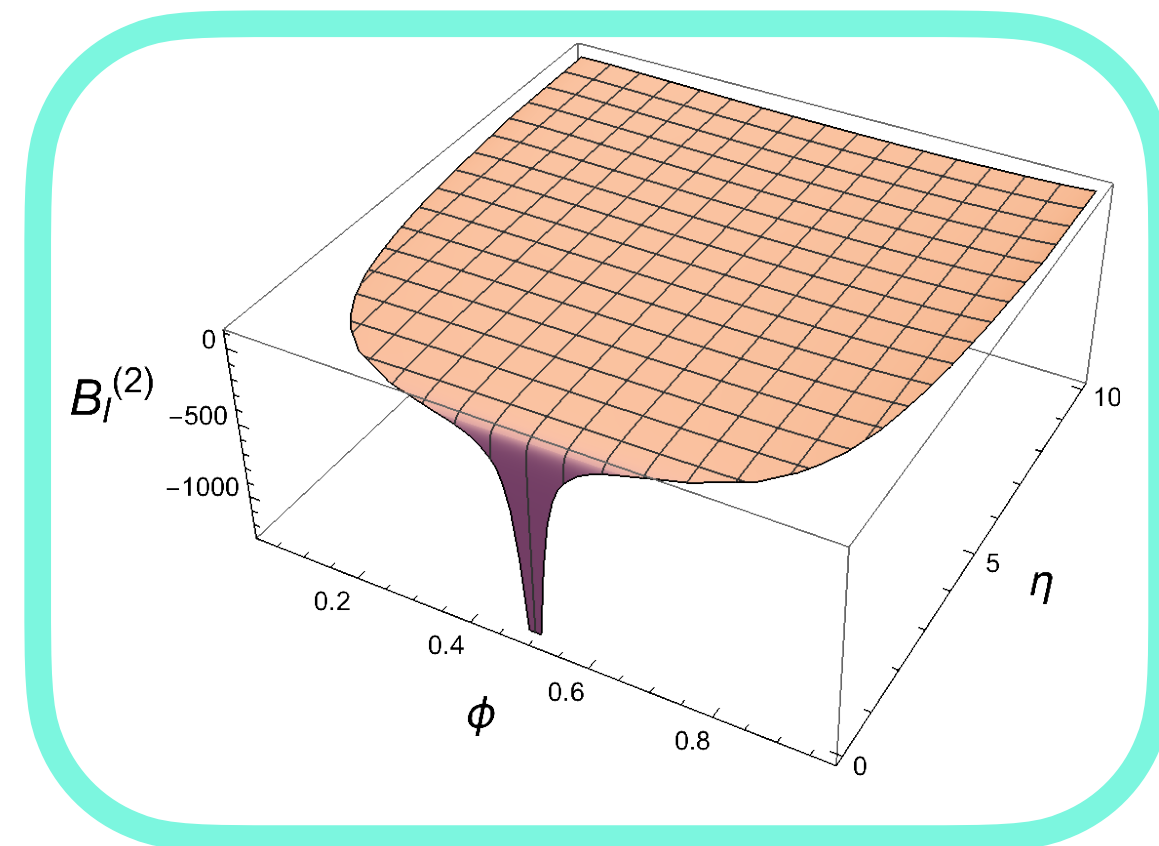
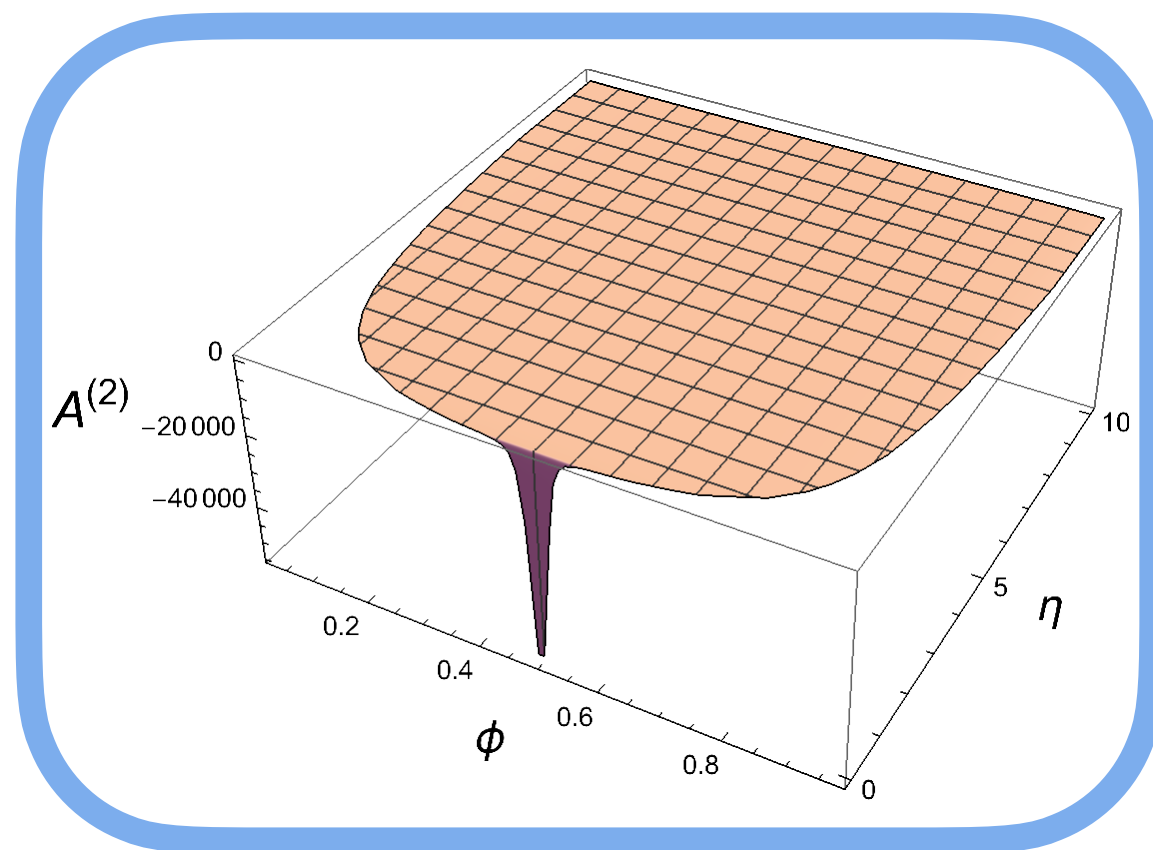




# Results: Two Loop

$$\mathcal{M}^{(2)} = A^{(2)} + n_l B_l^{(2)} + n_h C_h^{(2)} + n_l^2 D_l^{(2)} + n_h n_l E_{hl}^{(2)} + n_h^2 F_h^{(2)}$$

$$\mathcal{M}^{(2)} = \frac{\mathcal{M}_{-4}^{(2)}}{\epsilon^4} + \dots + \frac{\mathcal{M}_{-1}^{(2)}}{\epsilon} + \mathcal{M}_0^{(2)}$$





# Conclusion and Outlook

- ☑ First complete analytic 2-Loop amplitude for four fermion scattering in QED with a pair of massive leptons
- ☑ The computation of the 2-Loop amplitude is done within the framework of AIDA [Automated]
- ☑ Complete agreement with the universal IR poles predicted by SCET
- ☑ We have created a grid of 10500 points for the 2-Loop amplitude
- ☑ One crucial input for the theory initiative at MUonE

📌 Inclusion of the effects of the mass of the Muon

Mitov, Moch (2006) Becher, Melnikov (2007) Engel, Gnendiger, Signer, Ulrich (2019)  
Heller (2021)

📌 Possible synergy with the MC efforts to include this matrix element

Pavia and PSI Group

📌 Roundtable discussion: To discuss / understand how the grids of the  $e\mu$  amplitude should be prepared to facilitate the inclusion to MC generators

Pavia and PSI Group

📌 All the ingredients are available for the full analytic evaluation of the 2-Loop amplitude of  $q\bar{q} \rightarrow t\bar{t}$