

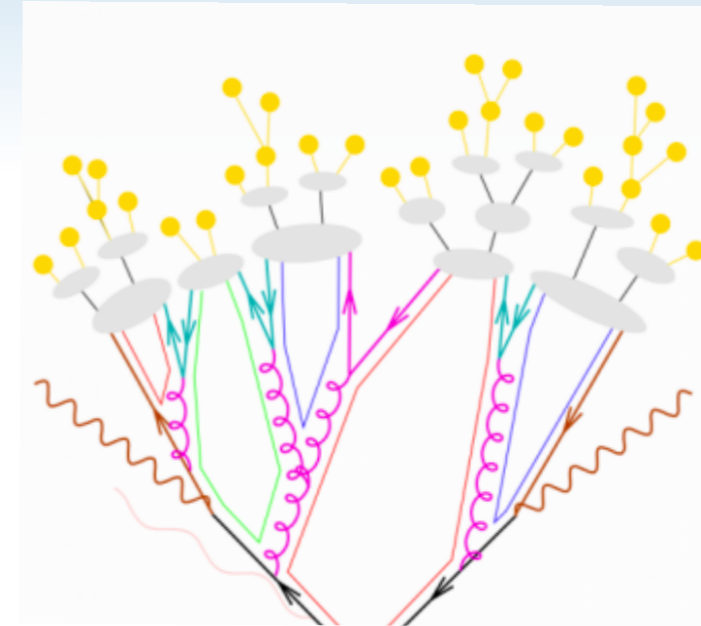
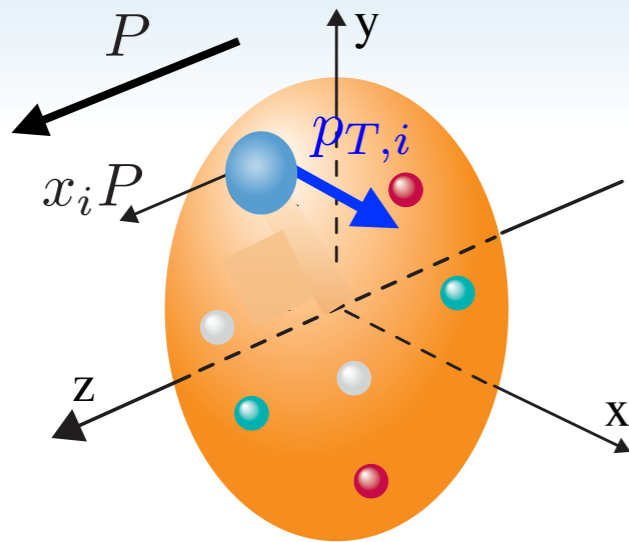
# Probing TMDs with jets

Kyle Lee  
LBNL

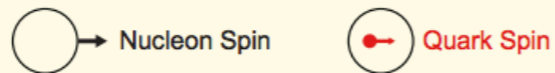
Sar Wors  
6 - 8 September 2021



# TMD structure



## Leading Twist TMDs



		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \circ \bullet$		$h_1^\perp = \circ \uparrow - \circ \downarrow$ Boer-Mulders
	L		$g_{1L} = \circ \rightarrow - \circ \leftarrow$ Helicity	$h_{1L}^\perp = \circ \nearrow - \circ \nwarrow$ Worm gear
	T	$f_{1T}^\perp = \circ \uparrow - \circ \downarrow$ Sivers	$g_{1T} = \circ \rightarrow - \circ \leftarrow$ Worm gear	$h_1 = \circ \uparrow - \circ \downarrow$ Transversity $h_{1T}^\perp = \circ \nearrow - \circ \nwarrow$

## Quark polarization

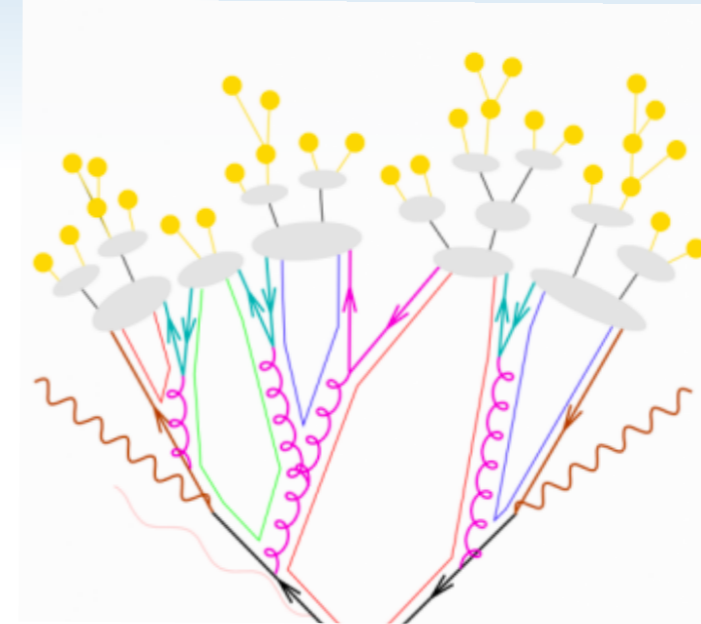
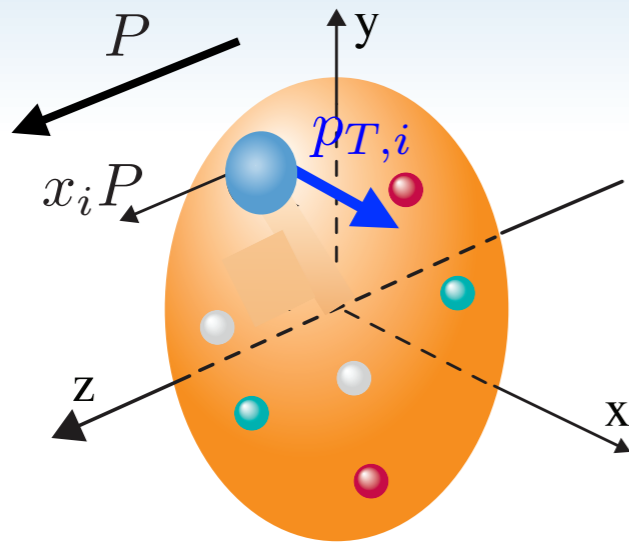
Hadron polarization

	U	L	T
U	$D^{h/q}$		$H^\perp h/q$
L		$G^{h/q}$	$H_L^\perp h/q$
T	$D_T^\perp h/q$	$G_T^{h/q}$	$H^{h/q} \quad H_T^\perp h/q$

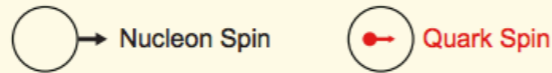
Quark **TMDPDF** inside spin- $\frac{1}{2}$  hadron

Quark **TMDFF** inside spin- $\frac{1}{2}$  hadron

# TMD structure



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	T	$f_{1T}^\perp = \bigcirc \leftarrow - \bigcirc \leftarrow$ Sivers	$g_{1T} = \bigcirc \leftarrow - \bigcirc \leftarrow$ Worm gear	$h_1 = \bigcirc \leftarrow - \bigcirc \leftarrow$ Transversity

Collinear analogs!

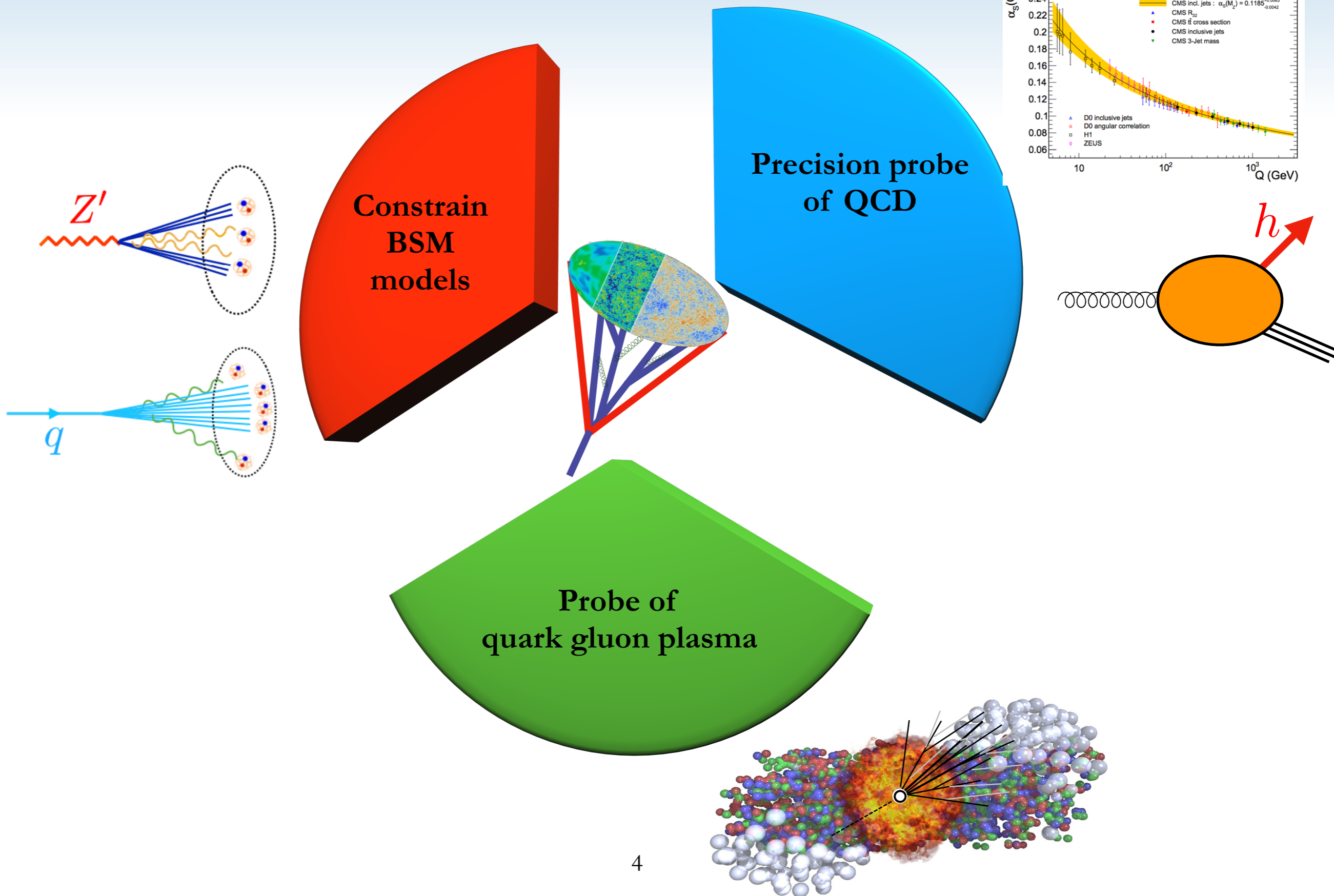
## Quark polarization

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		U	L	T
Hadron polarization	U	$D^{h/q}$		$H^\perp h/q$
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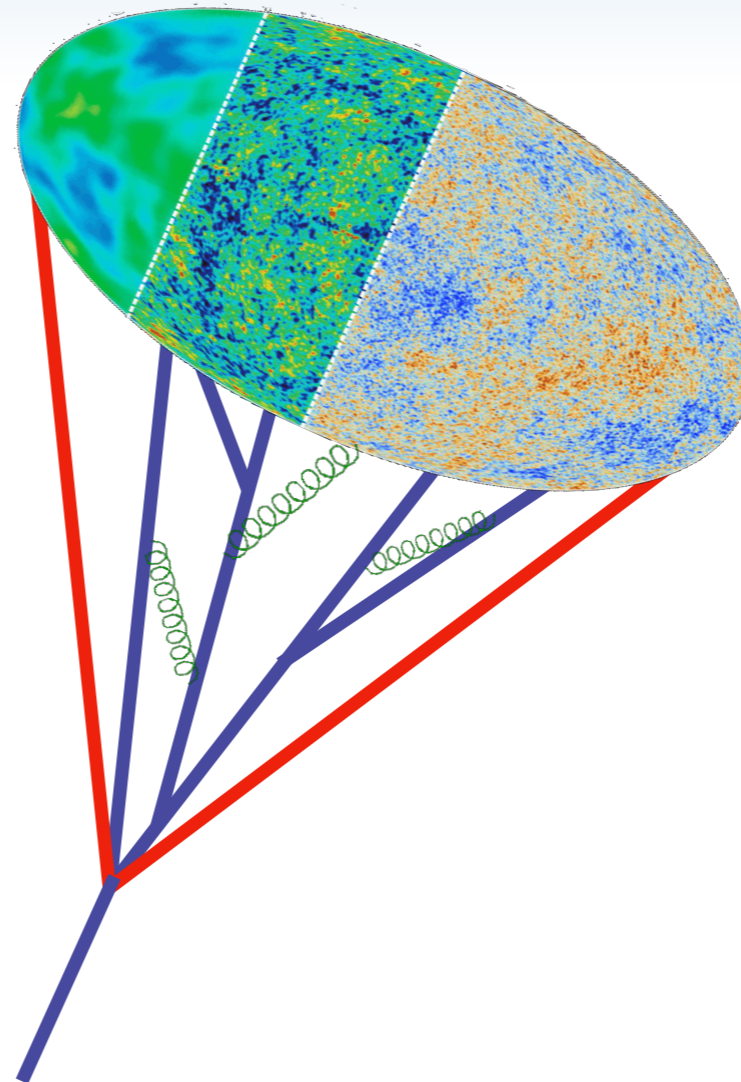
Quark **TMDPDF** inside spin- $\frac{1}{2}$  hadron

Quark **TMDFF** inside spin- $\frac{1}{2}$  hadron

# Applications of jets



# Study of hadron structures

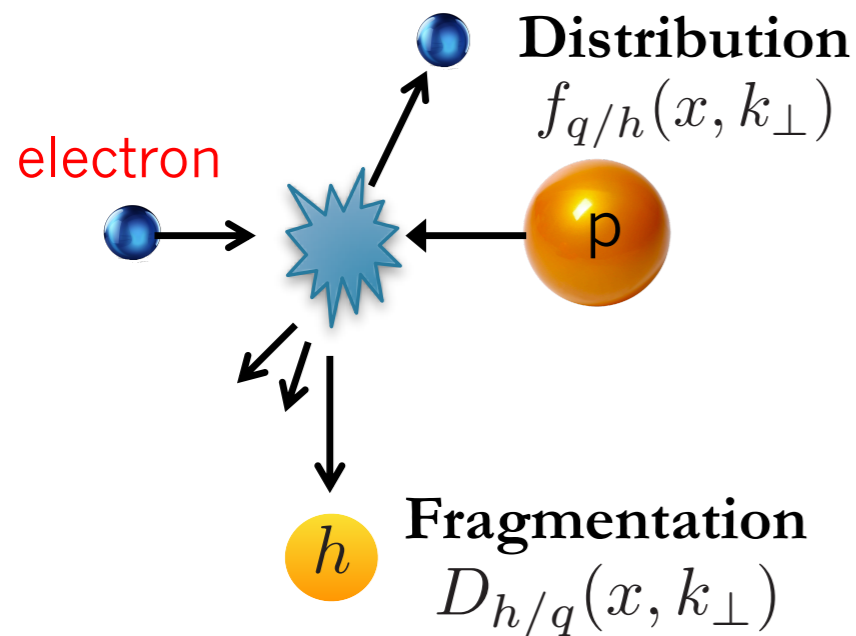


“Can we use **jets** to probe these TMD structure?”

# Standard processes to study TMD structure

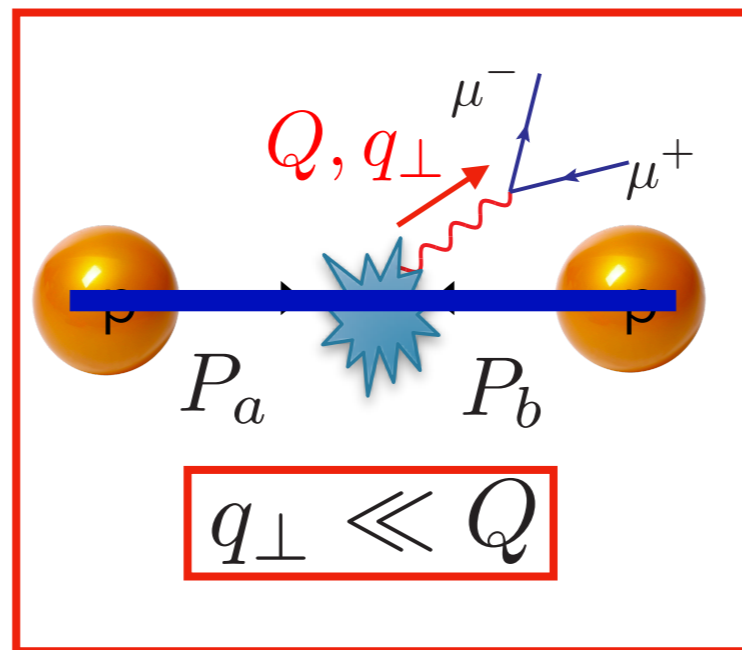
## Semi-Inclusive DIS (SIDIS)

$$\sigma \sim f_{q/P}(x, k_{\perp}) D_{h/q}(x, k_{\perp})$$



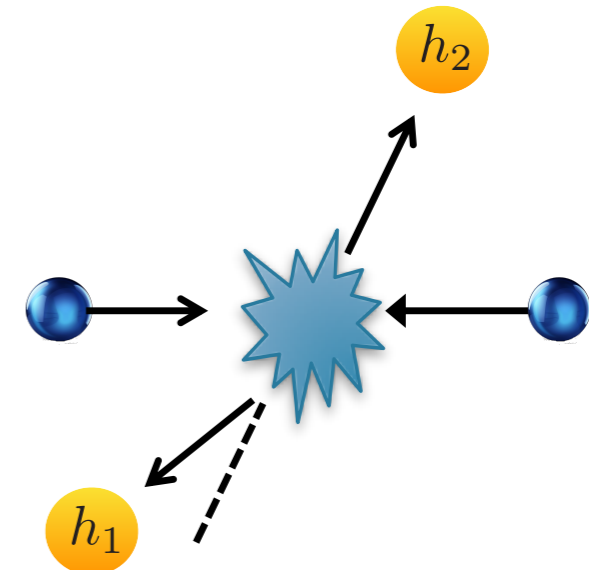
## Drell-Yan

$$\sigma \sim f_{q/P}(x, k_{\perp}) f_{\bar{q}/P}(x, k_{\perp})$$



## Dihadrons in $e^+e^-$

$$\sigma \sim D_{h_1/q}(x, k_{\perp}) D_{h_2/q}(x, k_{\perp})$$



- They have a well-established factorization formalism
- Small transverse momentum measured with respect to **an axis**

(CSS) Collin, Soper, Sterman '81-'85

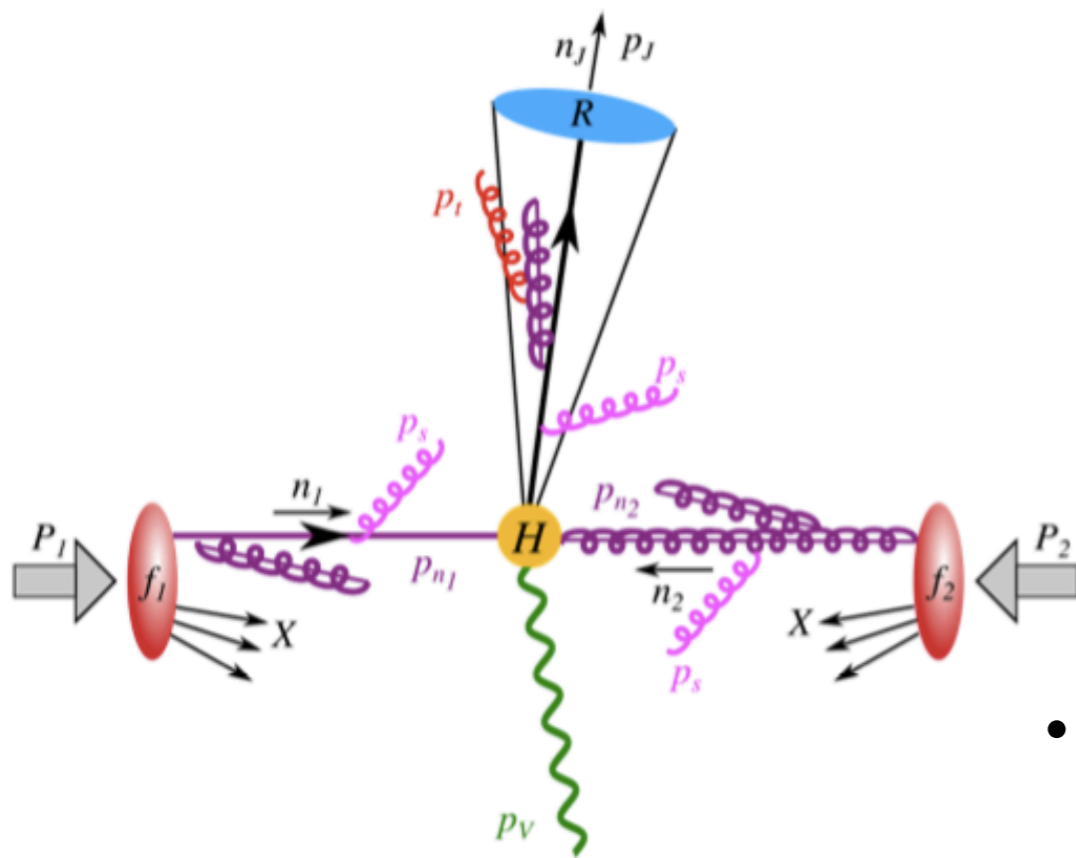
Ji, Ma, Yuan '04

Becher, Neubert, Wilhelm '11-'13

Echevarria, Idilbi, Scimemi '11-'14

# Beyond the standard processes

- Many other imaginable processes with sensitivity to the TMD structure



$$\begin{aligned}
 & PP \rightarrow J_1 + J_2 + X, \\
 & PP \rightarrow J + V + X, \\
 & PP \rightarrow J(h) + X, \dots
 \end{aligned}$$

LHC / RHIC

$$\begin{aligned}
 & eP \rightarrow e + J + X \\
 & eP \rightarrow Q + \bar{Q} + X, \\
 & eP \rightarrow J(h) + X, \dots
 \end{aligned}$$

EIC

- Many experiments sensitive to such processes
- Standard processes have low sensitivity to gluon TMDs.
- Standard processes sensitive to **two** TMDs simultaneously; many involving jets will only be sensitive to a single TMD.

Fig. from Chien, Shao, Wu '19

1) Inclusive jet production

**TMDFFs**  $PP / eP \rightarrow J(h) + X$

2) Lepton + Jet imbalance

**TMDPDFs**  $eP \rightarrow e + J + X$

3) Lepton + Jet imbalance  
with hadron in jet

$$eP \rightarrow e + J(h) + X$$

**TMDFFs / TMDPDFs**

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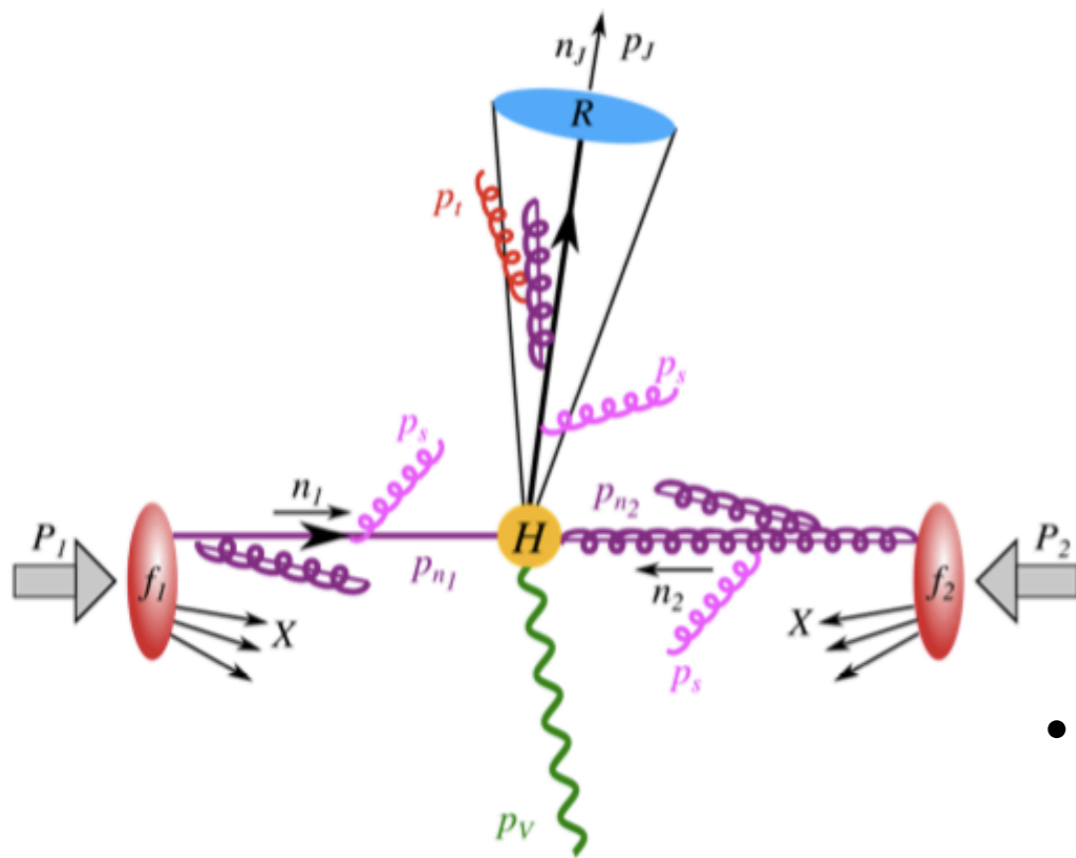


Fig. from Chien, Shao, Wu '19

collinear structure

1) Inclusive jet production

**TMDFFs**  $PP / eP \rightarrow J(h) + X$

2) Lepton + Jet imbalance

**TMDPDFs**  $eP \rightarrow e + J + X$

3) Lepton + Jet imbalance  
with hadron in jet

$$eP \rightarrow e + J(h) + X$$

**TMDFFs / TMDPDFs**



# Jet Fragmentation Functions (JFFs)

Unpolarized case:

(replace  $pp$  with  $ep$  for EIC)

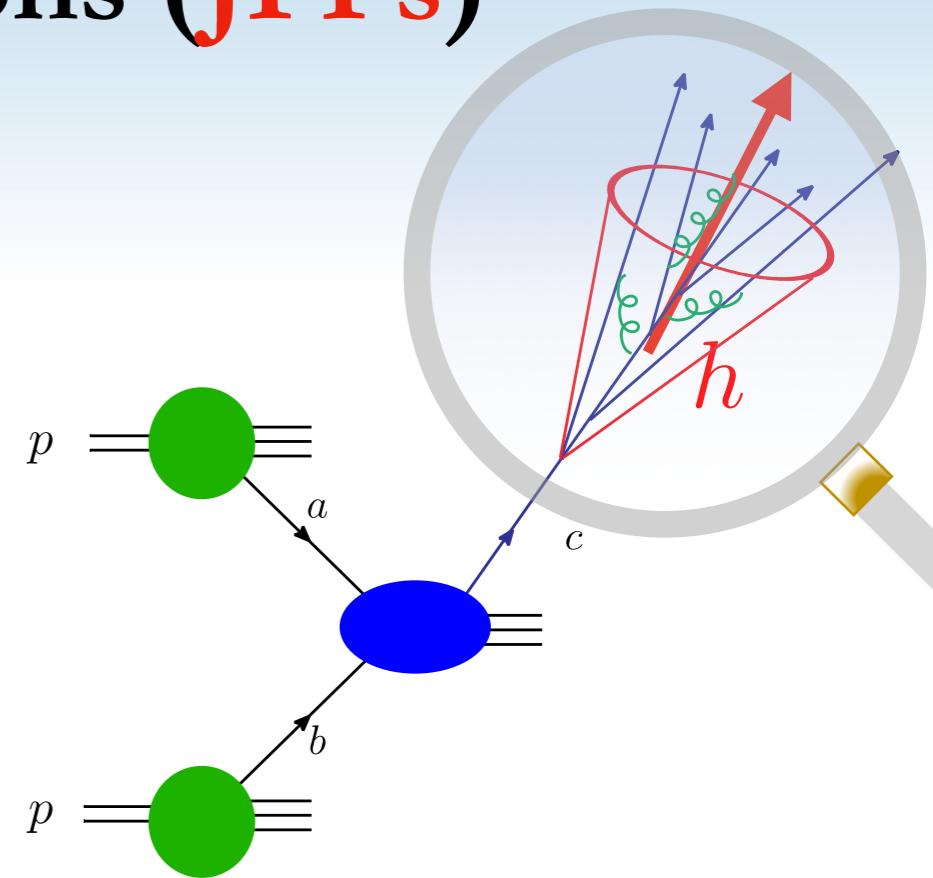
$$\frac{d\sigma^{pp \rightarrow \text{jet}(h)X}}{dp_T d\eta dz_h} = \sum_{a,b,c} f_{a/A} \otimes f_{b/B} \otimes H_{ab}^c \otimes \mathcal{G}_c^h(z_h)$$

$\Lambda_{\text{QCD}}$                        $p_T$                        $p_T R$   
 $\Lambda_{\text{QCD}}$

where

$$z = p_T^J / p_T^c$$

$$z_h = p_T^h / p_T^J$$



1) Inclusive jet production  
**TMDFFs**  
 collinear PDFs/FFs

$$\frac{d\sigma^{pp \rightarrow hX}}{dp_T d\eta} = \sum_{a,b,c} f_{a/A} \otimes f_{b/B} \otimes H_{ab}^c \otimes D_c^h$$

$\Lambda_{\text{QCD}}$                        $p_T$                        $\Lambda_{\text{QCD}}$

where  $z = p_T^h / p_T^c$

Procura, Stewart `10  
 Arleo, Fontannaz, Guillet, Nguyen `14  
 Kaufmann, Mukherjee, Vogelsang `15  
 Kang, Ringer, Vitev `16  
 Dai, Kim, Leibovich `16

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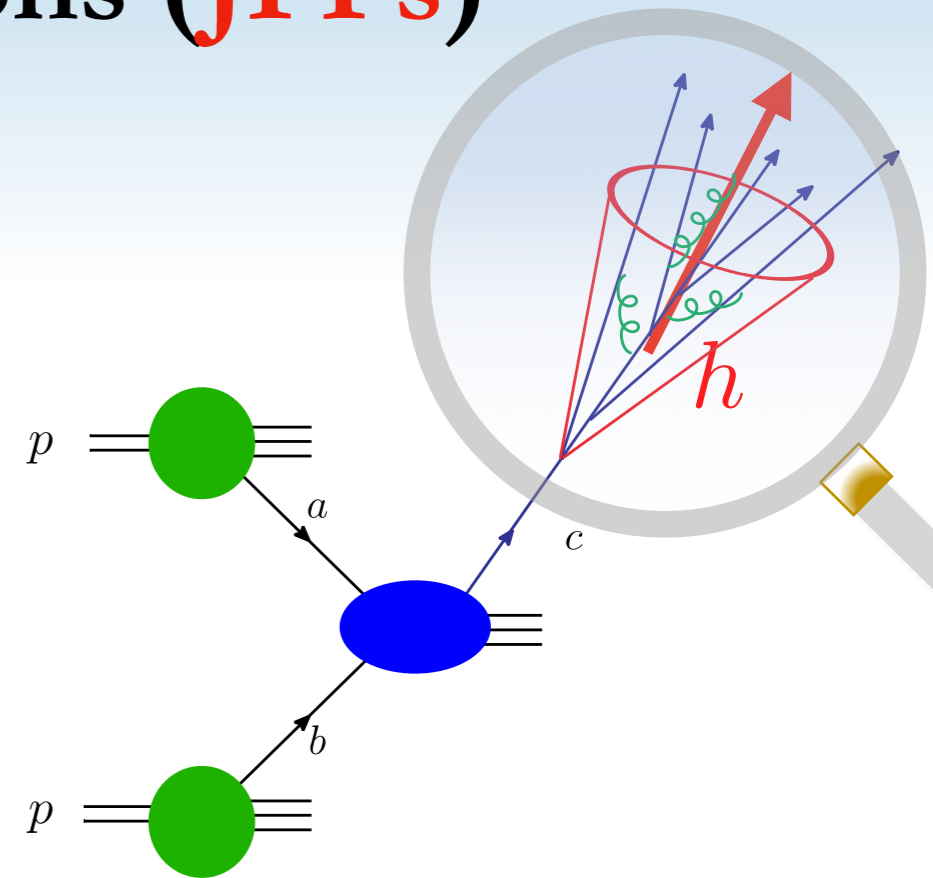
$\Lambda_{\text{QCD}}$        $p_T$        $p_T R$        $\Lambda_{\text{QCD}}$

collinear PDFs

where

$$z = p_T^J / p_T^c$$

$$z_h = p_T^h / p_T^J$$



1) Inclusive jet production

TMDFFs

collinear structure

$$\frac{d\sigma^{pp \rightarrow hX}}{dp_T d\eta} = \sum_{a,b,c} f_{a/A} \otimes f_{b/B} \otimes H_{ab}^c \otimes D_c^h$$

$\Lambda_{\text{QCD}}$        $p_T$        $\Lambda_{\text{QCD}}$

where  $z = p_T^h / p_T^c$

collinear FFs

# Jet Fragmentation Functions (JFFs)

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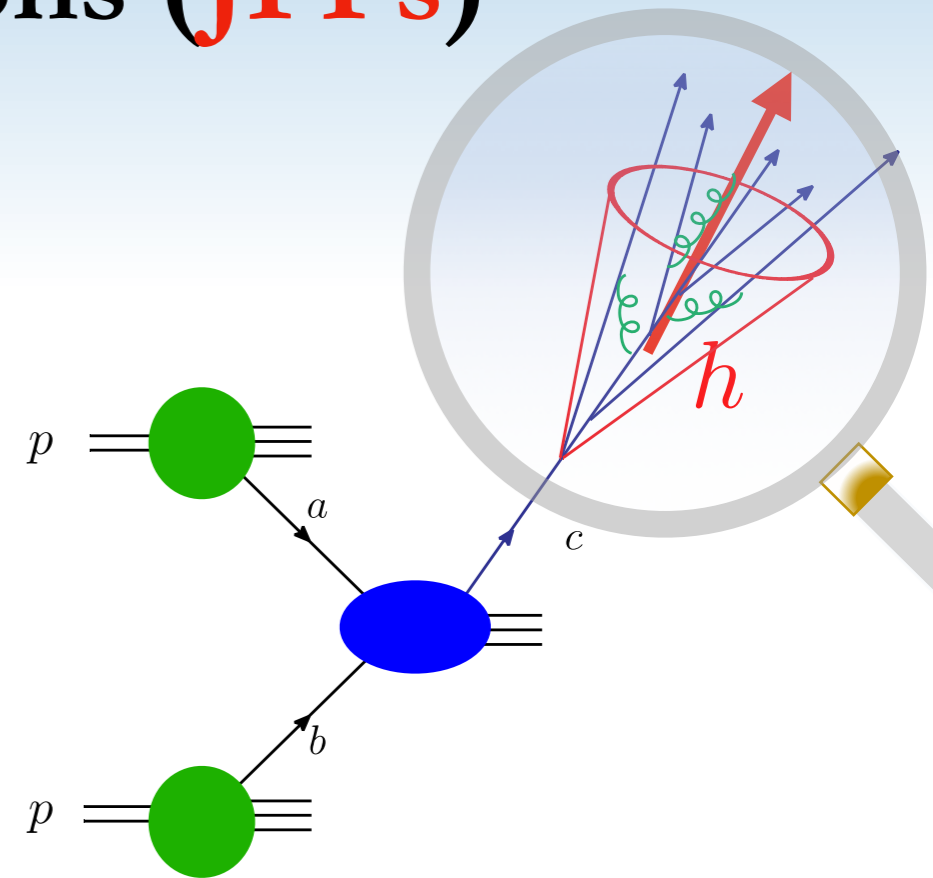
$\Lambda_{\text{QCD}}$        $p_T$        $p_T R$        $\Lambda_{\text{QCD}}$

collinear PDFs

where

$$z = p_T^J / p_T^c$$

$$z_h = p_T^h / p_T^J$$



IR sensitivity and require matching:

$$\mathcal{G}_c^h(z, z_h, p_T R, \mu) = \sum_j \mathcal{J}_{ij}(z, z_h, p_T R, \mu) \otimes D_j^h(z_h, \mu)$$

matching coefficients

$p_T R$

$\Lambda_{\text{QCD}}$

collinear FFs

1) Inclusive jet production

TMDFFs

collinear structure

- Collinear JFFs can be related to collinear FFs

Procura, Stewart `10

Arleo, Fontannaz, Guillet, Nguyen `14

Kaufmann, Mukherjee, Vogelsang `15

Kang, Ringer, Vitev `16

Dai, Kim, Leibovich `16

# Polarized Jet Fragmentation Functions (JFFs)

Polarized case:

$$\frac{d\sigma_{LL}^{\bar{p}p \rightarrow \text{jet}(\bar{h})X}}{dp_T d\eta dz_h} = \sum_{a,b,c} \Lambda_A g_{a/A} \otimes f_{b/B} \otimes \Delta_{LL} H_{\bar{a}b}^{\bar{c}} \otimes \Delta_L \mathcal{G}_c^h(z_h) \Lambda_h$$

$$\frac{d\sigma_{TT}^{\bar{p}p \rightarrow \text{jet}(\bar{h})X}}{dp_T d\eta dz_h} = \sum_{a,b,c} S_{A\perp}^i h_{a/A} \otimes f_{b/B} \otimes \Delta_{TT} (H_{ij})_{\bar{a}b}^{\bar{c}} \otimes \Delta_T \mathcal{G}_c^h(z_h) S_{h\perp}^j$$

Helicity PDF

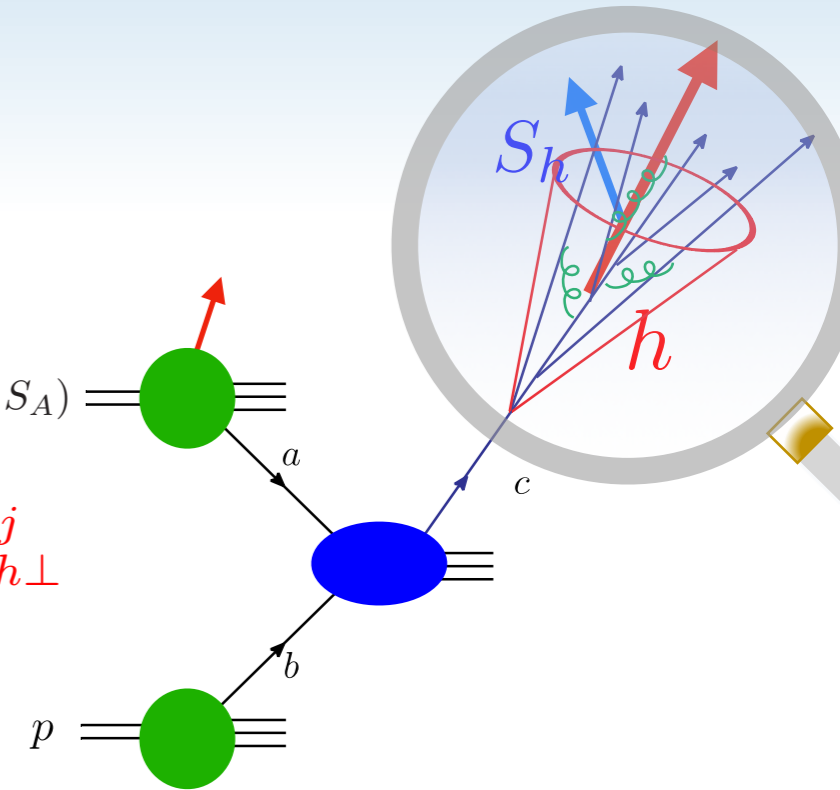
Helicity JFF

Transversity PDF

Transversity JFF

where  $z = p_T^J / p_T^c$

$z_h = p_T^h / p_T^J$



$$d\sigma_{LL} \equiv \frac{d\sigma(\Lambda_A, \Lambda_h) - d\sigma(\Lambda_A, -\Lambda_h)}{2}$$

$$d\sigma_{TT} \equiv \frac{d\sigma(\vec{S}_{A\perp}, \vec{S}_{h\perp}) - d\sigma(\vec{S}_{A\perp}, -\vec{S}_{h\perp})}{2}$$

Similar definitions for

$$\Delta_{LL} H_{\bar{a}b}^{\bar{c}} \text{ and } \Delta_{TT} (H_{ij})$$

# Polarized Jet Fragmentation Functions (JFFs)

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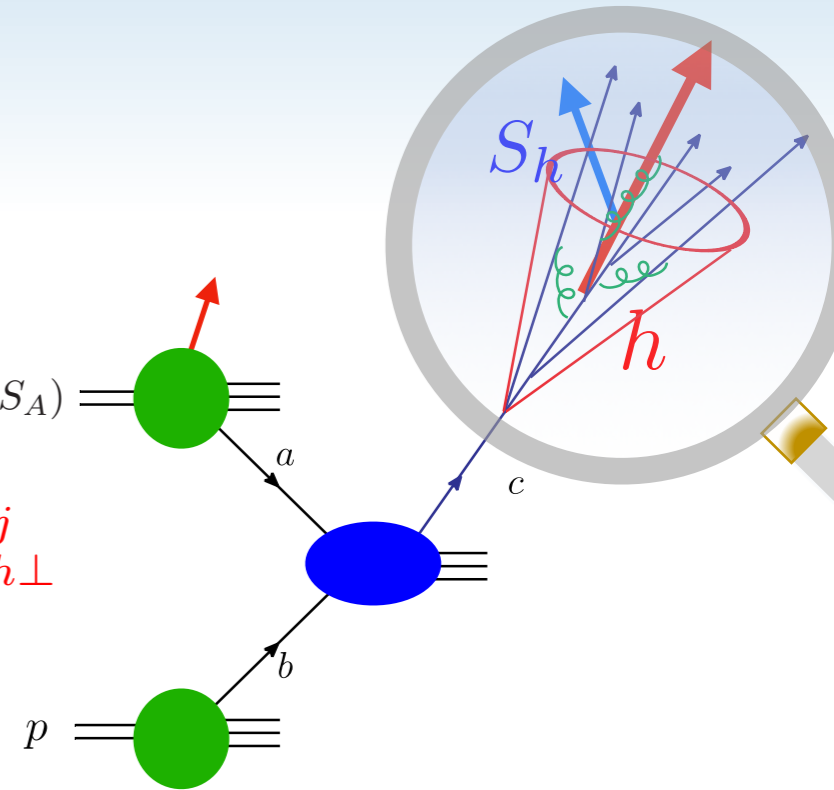
**Helicity PDF**
**Helicity JFF**

$$\frac{d\sigma_{TT}^{\bar{p}p \rightarrow \text{jet}(\bar{h})X}}{dp_T d\eta dz_h} = \sum_{a,b,c} S_{A\perp}^i h_{a/A} \otimes f_{b/B} \otimes \Delta_{TT} (H_{ij})_{\bar{a}b}^{\bar{c}} \otimes \Delta_T \mathcal{G}_c^h(z_h) S_{h\perp}^j$$

**Transversity PDF**
**Transversity JFF**

where  $z = p_T^J / p_T^c$

$z_h = p_T^h / p_T^J$



$$\frac{d\sigma_{LL}^{\bar{p}p \rightarrow \bar{h}X}}{dp_T d\eta} = \sum_{a,b,c} \Lambda_A g_{a/A} \otimes f_{b/B} \otimes \Delta_{LL} H_{\bar{a}b}^{\bar{c}} \otimes G_c^h$$

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**Quark polarization**

	U	L	T
U	$D^{h/q}$		
L		$G^{h/q}$	
T			$H^{h/q}$

**Hadron polarization**

# Polarized Jet Fragmentation Functions (JFFs)

Polarized case:

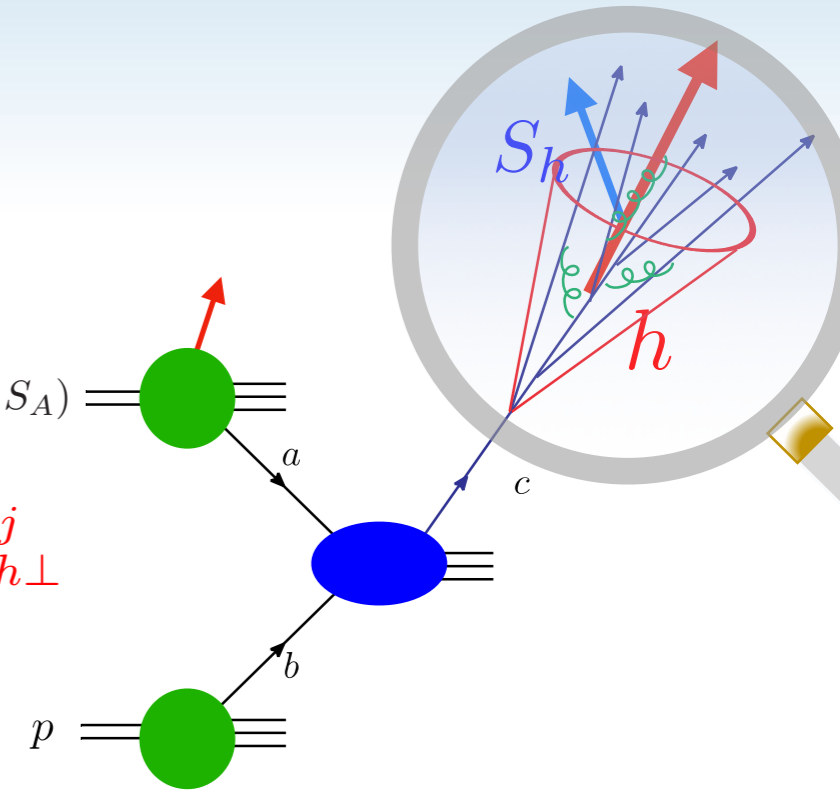
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Helicity PDF                      Helicity JFF  
Transversity PDF                      Transversity JFF

$$\frac{d\sigma_{TT}^{\bar{p}p \rightarrow \text{jet}(\bar{h})X}}{dp_T d\eta dz_h} = \sum_{a,b,c} S_{A\perp}^i h_{a/A} \otimes f_{b/B} \otimes \Delta_{TT} (H_{ij})_{\bar{a}b}^{\bar{c}} \otimes \Delta_T \mathcal{G}_c^h(z_h) S_{h\perp}^j$$

where  $z = p_T^J / p_T^c$

$z_h = p_T^h / p_T^J$



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$p_T R$                        $\Lambda_{\text{QCD}}$

$$\Delta_L \mathcal{G}_c^h(z, z_h, p_T R, \mu) = \sum_j \Delta_L \mathcal{J}_{ij}(z, z_h, p_T R, \mu) \otimes G_j^h(z_h, \mu)$$

$$\Delta_T \mathcal{G}_c^h(z, z_h, p_T R, \mu) = \sum_j \Delta_T \mathcal{J}_{ij}(z, z_h, p_T R, \mu) \otimes H_j^h(z_h, \mu)$$

		Quark polarization		
		U	L	T
Hadron polarization	U	$D^{h/q}$		
	L		$G^{h/q}$	
	T			$H^{h/q}$

# Jet Fragmentation Functions (JFFs)

- **Light charged hadrons**

Arleo, Fontannaz, Guillet, Nguyen `14

Kaufmann, Mukherjee, Vogelsang `15

Kang, Ringer, Vitev `16

Neill, Scimemi, Waalewijn `16

- **Photons**

Kaufmann, Mukherjee, Vogelsang `16

- **Heavy flavor mesons**

Chien, Kang, Ringer, Vitev, Xing `15

Bain, Dai, Hornig, Leibovich, Makris, Mehen `16

Anderle, Kaufmann, Stratmann, Ringer, Vitev `17

- **Quarkonia**

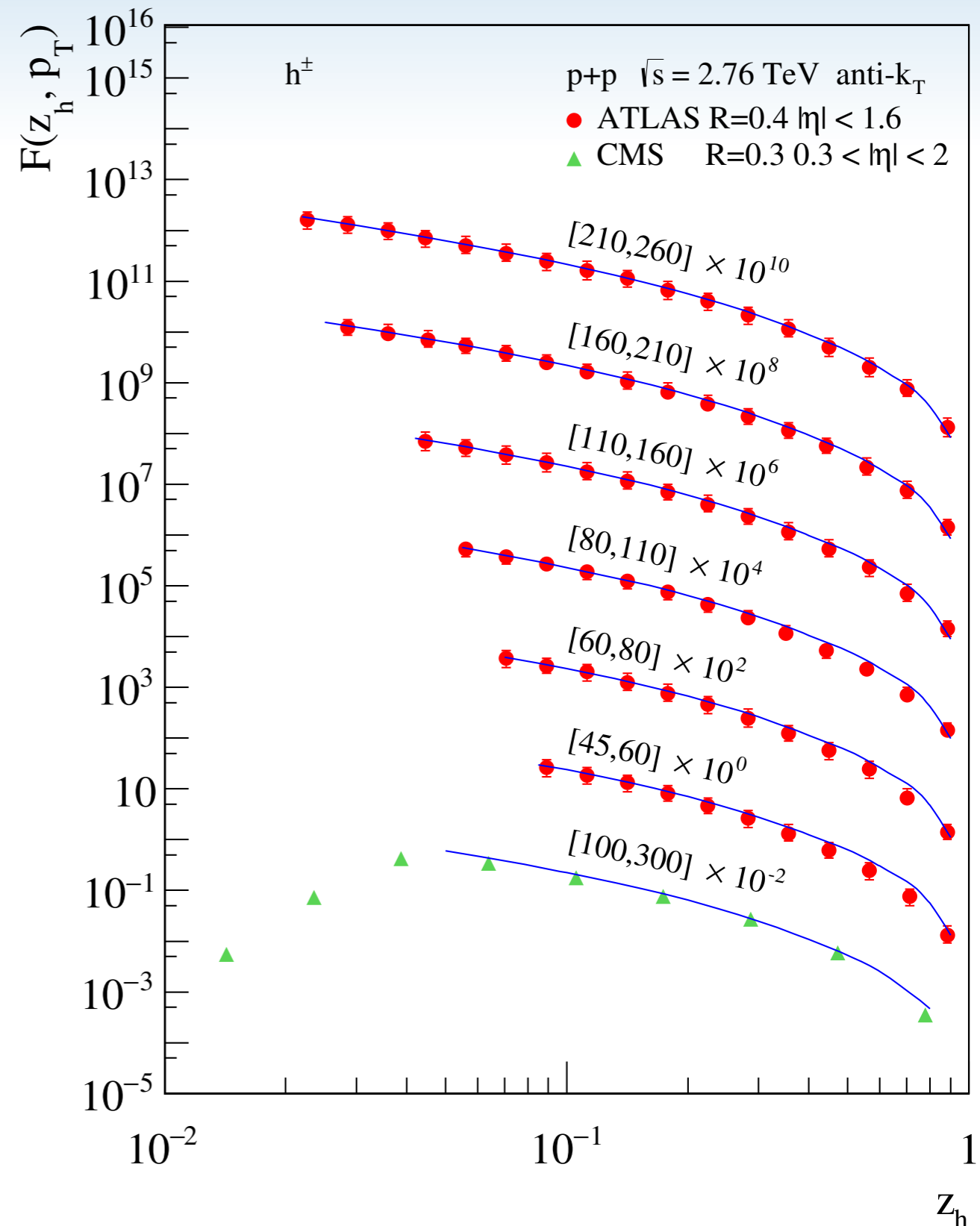
Baumgart, Leibovich, Mehen, Rothstein `14

Bain, Dai, Hornig, Leibovich, Makris, Mehen `16

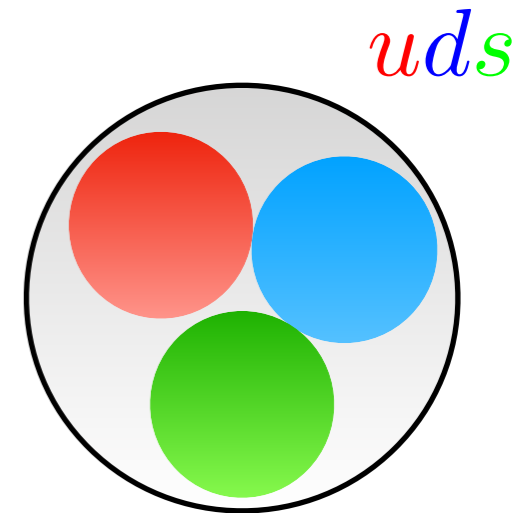
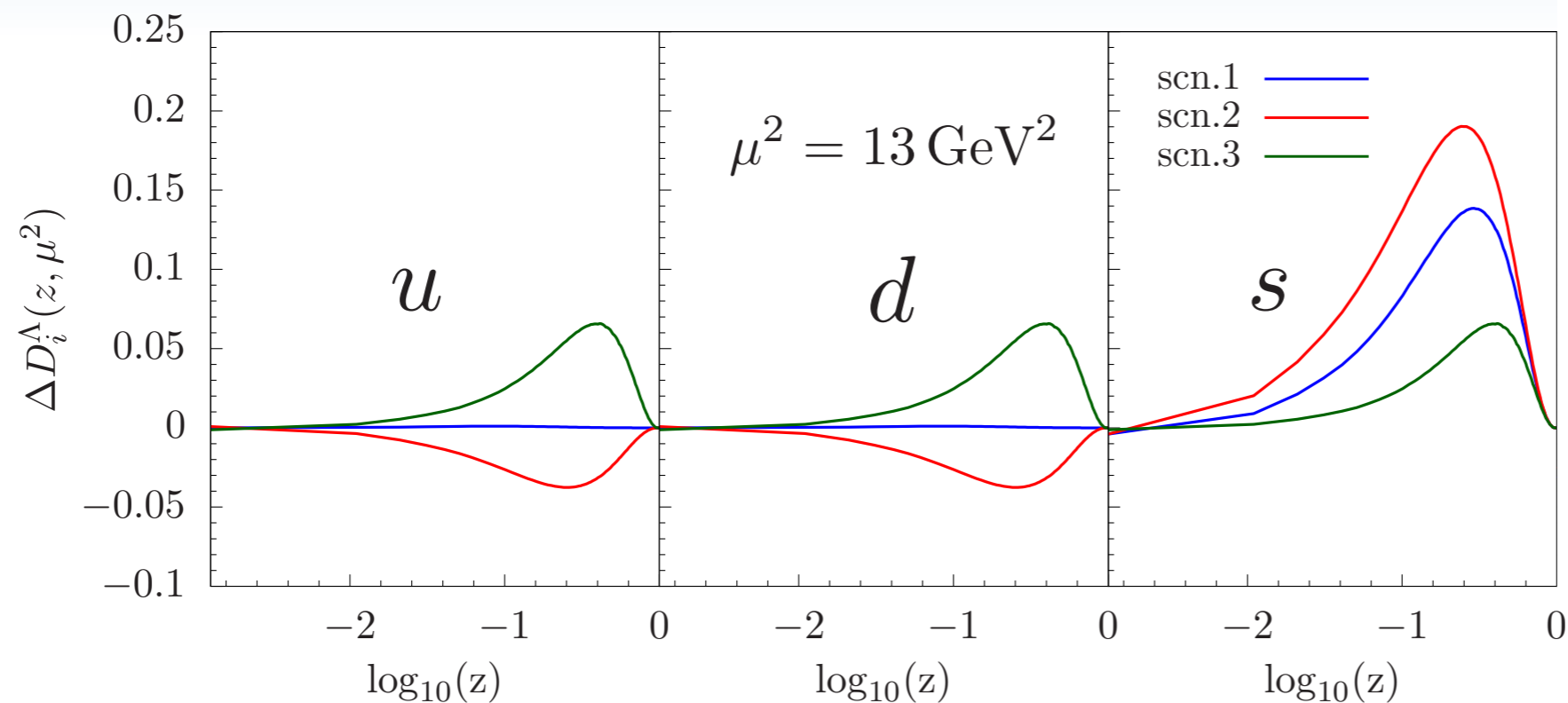
Kang, Qiu, Ringer, Xing, Zhang `17

Bain, Dai, Leibovich, Makris, Mehen `17

$$F(z_h, p_T) = \frac{d\sigma^{pp \rightarrow (\text{jet}h)X}}{dp_T d\eta dz_h} \bigg/ \frac{d\sigma^{pp \rightarrow \text{jet}X}}{dp_T d\eta}$$



# Longitudinally polarized $\Lambda$

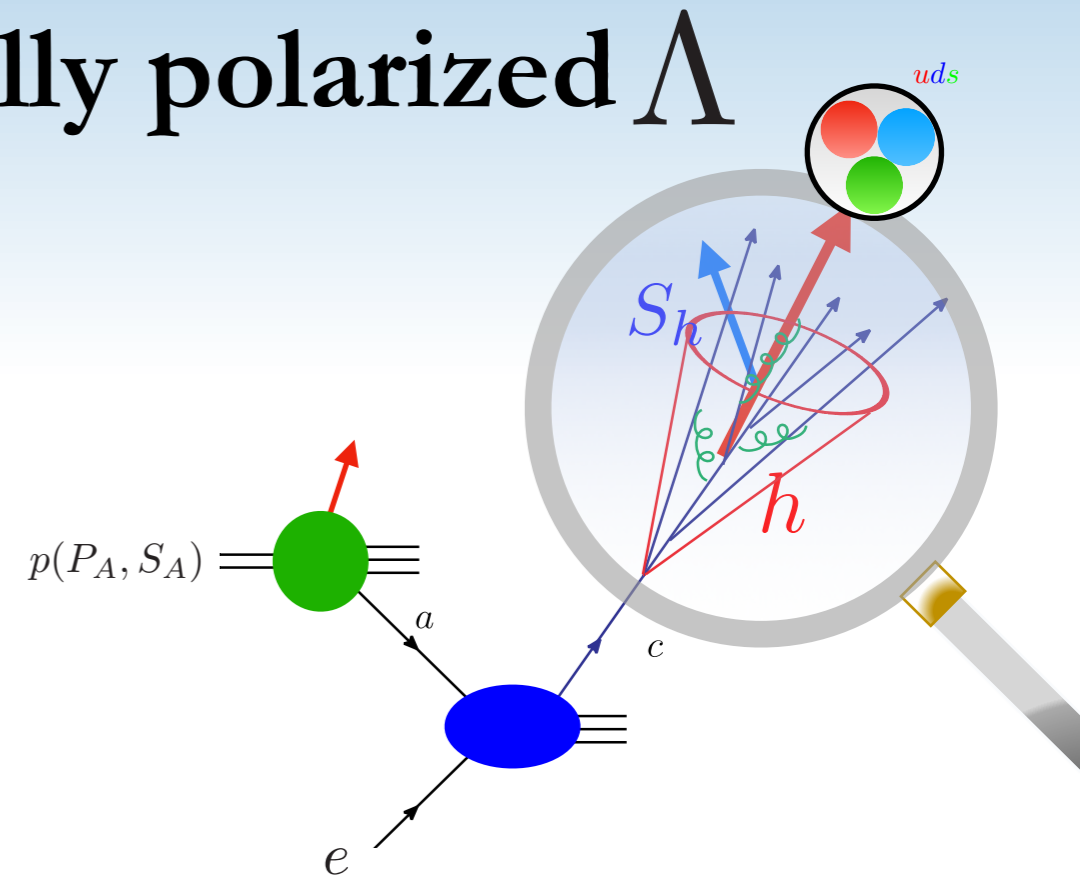
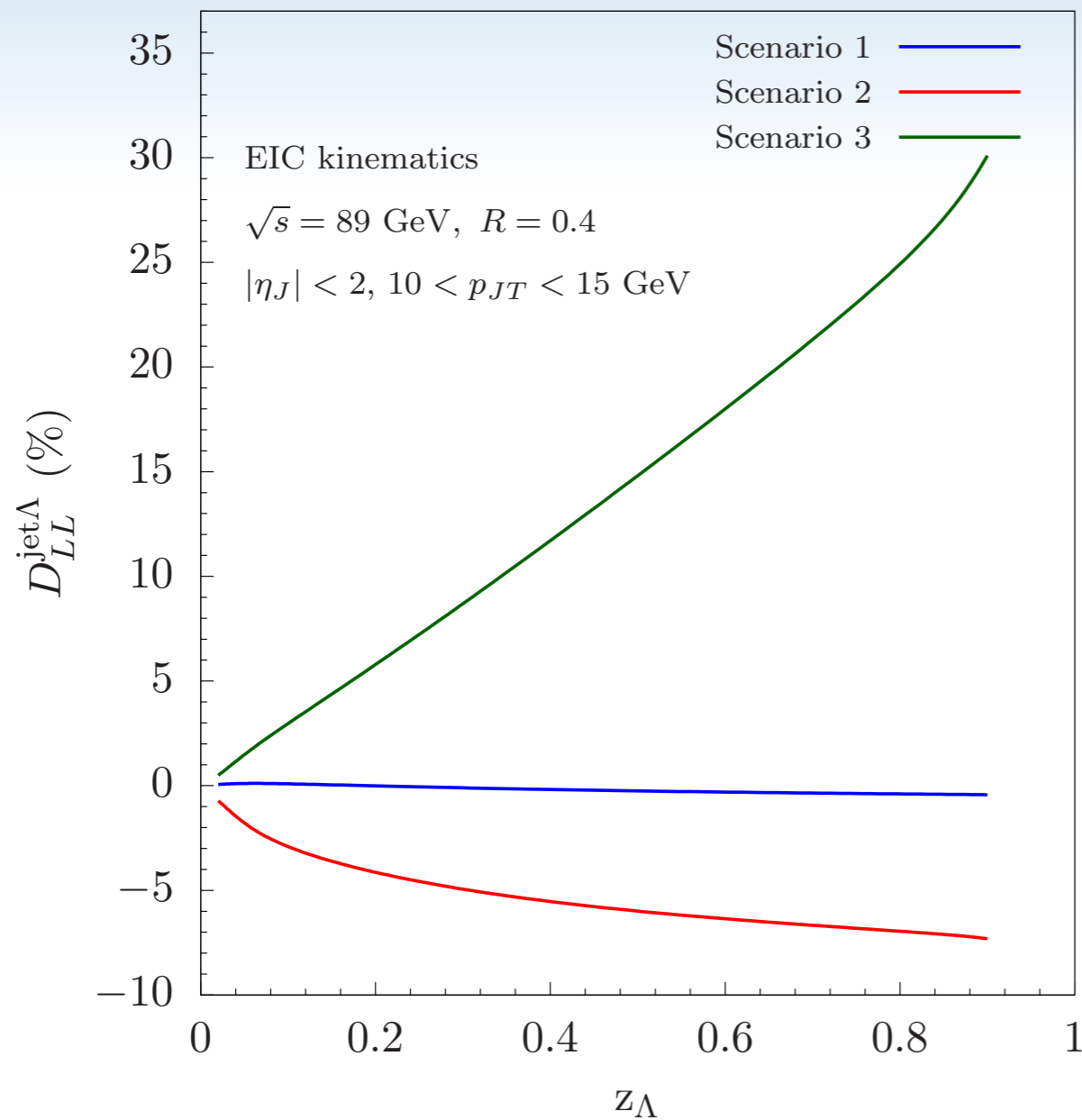


Three different scenarios can explain LEP data equally well

- **Scenario 1:** Only strange quarks have contribution to the fragmentation.
- **Scenario 2:** Negative distributions of up and down quarks are assumed.
- **Scenario 3:** Same fragmentation for up, down, and strange quarks.



# JFFs to study longitudinally polarized $\Lambda$



Gives shape expected from the scenarios

$$D_{LL}^{\text{jet}\Lambda} = \frac{d\Delta_{LL}\sigma}{d\sigma}$$

- **Scenario 1:** Only strange quarks have contribution to the fragmentation.
- **Scenario 2:** Negative distributions of up and down quarks are assumed.
- **Scenario 3:** Same fragmentation for up, down, and strange quarks.

# TMD hadron in jet (**TMDJFFs**)

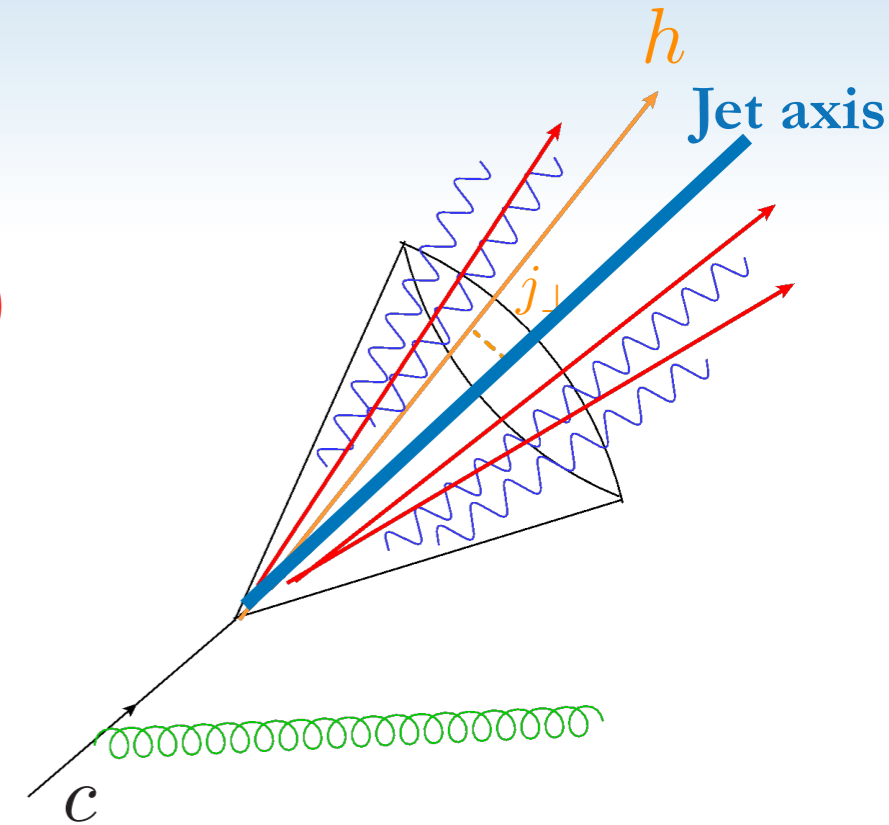
Unpolarized case:

(replace  $pp$  with  $ep$  for EIC)

$$\frac{d\sigma^{pp \rightarrow \text{jet}(h)X}}{dp_T d\eta dz_h d^2 j_\perp} = \sum_{a,b,c} f_{a/A} \otimes f_{b/B} \otimes H_{ab}^c \otimes \mathcal{G}_c^h(z_h, j_\perp)$$

$\Lambda_{\text{QCD}}$                        $p_T$                        $p_T R$   
 $\Lambda_{\text{QCD}}$

where  $z = p_T^J / p_T^c$   
 $z_h = p_T^h / p_T^J$

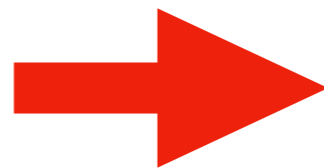


(including polarized jet fragmentation functions)

TMD Fragmentation Functions (**TMDFFs**)

TMD Jet Fragmentation Functions (**TMDJFFs**)

		Quark polarization		
		U	L	T
Hadron polarization	U	$D^{h/q}$		$H^{\perp h/q}$
	L		$G^{h/q}$	$H_L^{\perp h/q}$
	T	$D_T^{\perp h/q}$	$G_T^{h/q}$	$H^{h/q} \quad H_T^{\perp h/q}$



		Quark polarization		
		U	L	T
Hadron polarization	U	$\mathcal{D}^{h/q}$		$\mathcal{H}^{\perp h/q}$
	L		$\mathcal{G}^{h/q}$	$\mathcal{H}_L^{\perp h/q}$
	T	$\mathcal{D}_T^{\perp h/q}$	$\mathcal{G}_T^{h/q}$	$\mathcal{H}^{h/q} \quad \mathcal{H}_T^{\perp h/q}$

# TMD hadron in jet (TMDJFFs)

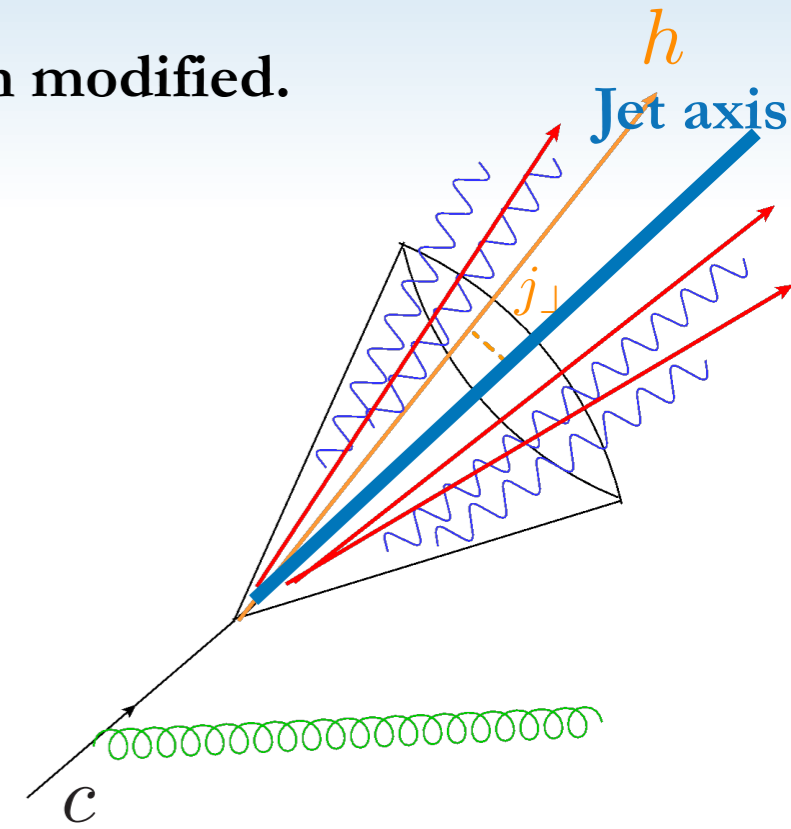
- Still hard-collinear factorization structure other than jet function modified.

$$\frac{d\sigma^{pp \rightarrow \text{jet}(h)X}}{dp_T d\eta dz_h d^2 j_\perp} = \sum_{a,b,c} f_{a/A} \otimes f_{b/B} \otimes H_{ab}^c \otimes \mathcal{G}_c^h(z_h, \mathbf{j}_\perp)$$

When  $\Lambda_{\text{QCD}} \lesssim j_\perp \ll p_T R$ ,  $\lambda \sim j_\perp/p_T$

collinear  $k_c \sim p_T(\lambda^2, 1, \lambda)$

soft  $k_s \sim p_T(\lambda R, \lambda/R, \lambda)$



## Unpolarized TMDJFF

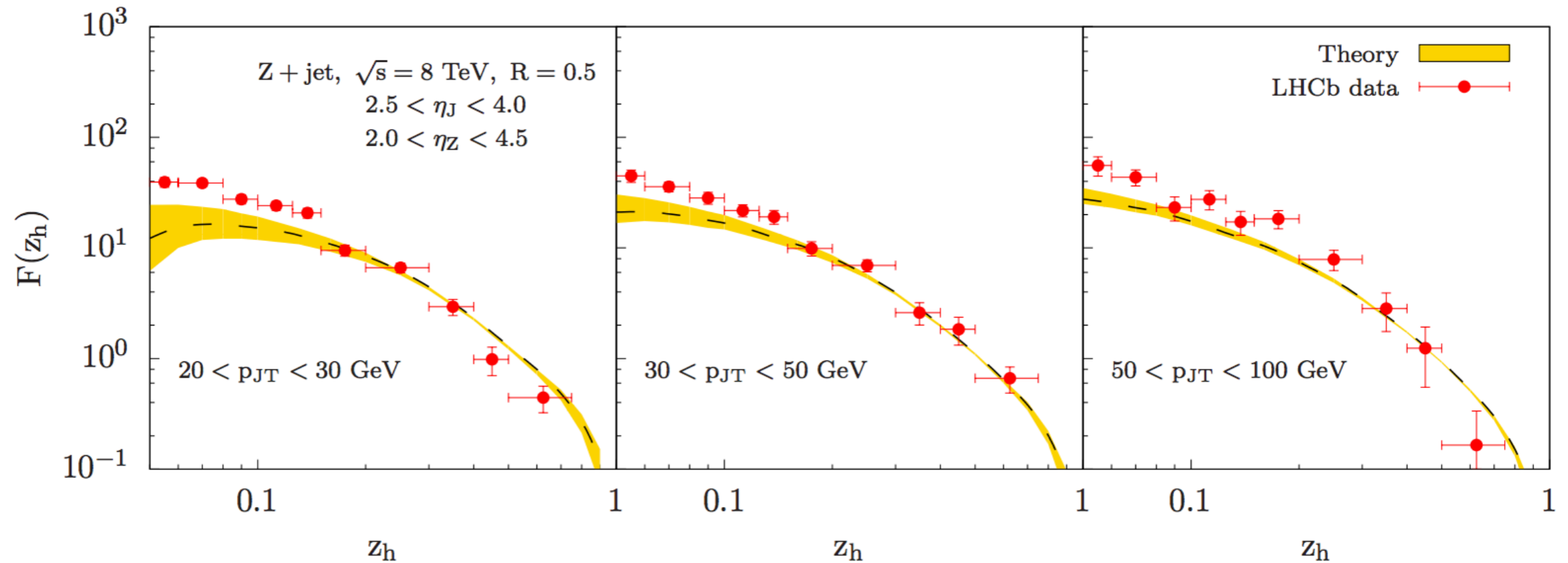
$$\begin{aligned} \mathcal{D}_1^{h/q}(z, z_h, j_\perp^2, \mu, \zeta_J) &= \mathcal{H}_{c \rightarrow i}(z, p_T R, \mu) \int_{\mathbf{k}_\perp, \lambda_\perp} D_1^{h/q, \text{unsub}}(z_h, \mathbf{k}_\perp^2, \mu, \zeta'/\nu^2) S_q(\lambda_\perp^2, \mu, \nu R) \\ &= \mathcal{H}_{c \rightarrow i}(z, p_T R, \mu) \int \frac{b db}{2\pi} J_0\left(\frac{j_\perp b}{z_h}\right) \tilde{D}_1^{h/q, \text{unsub}}(z_h, b^2, \mu, \zeta'/\nu^2) S_q(b^2, \mu, \nu R) \\ &= \mathcal{H}_{c \rightarrow i}(z, p_T R, \mu) \underbrace{D_1^{h/q}(z_h, j_\perp^2, \mu, \zeta_J)} \end{aligned}$$

Standard subtracted TMDFFs, say in SIDIS

Relation also holds for other TMDJFFs.

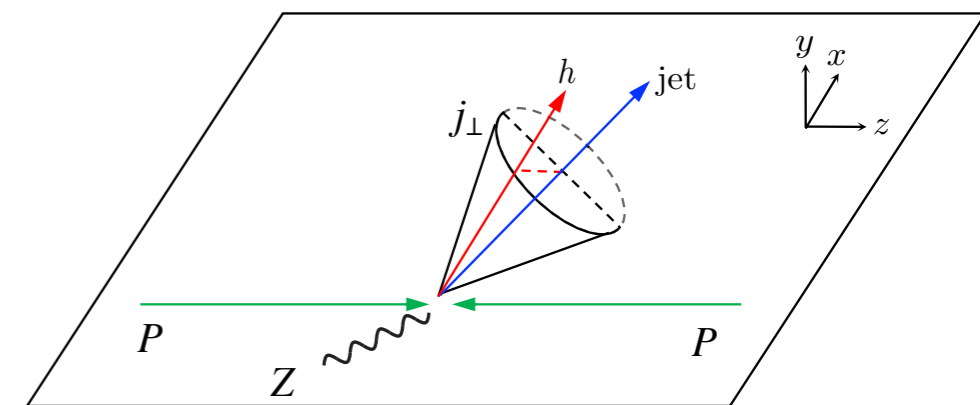
- TMDJFFs can be related to TMDFFs

# Z-tagged jet

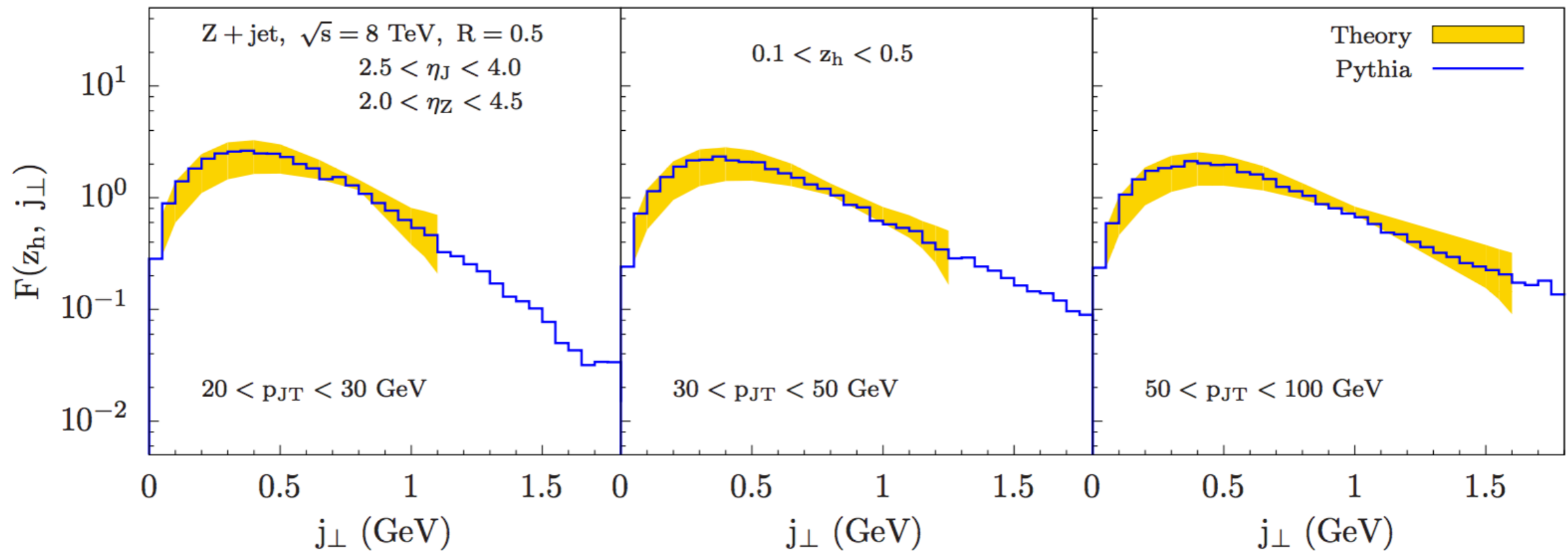


- LHCb collaboration measured collinear FFs and  $j_{\perp}$

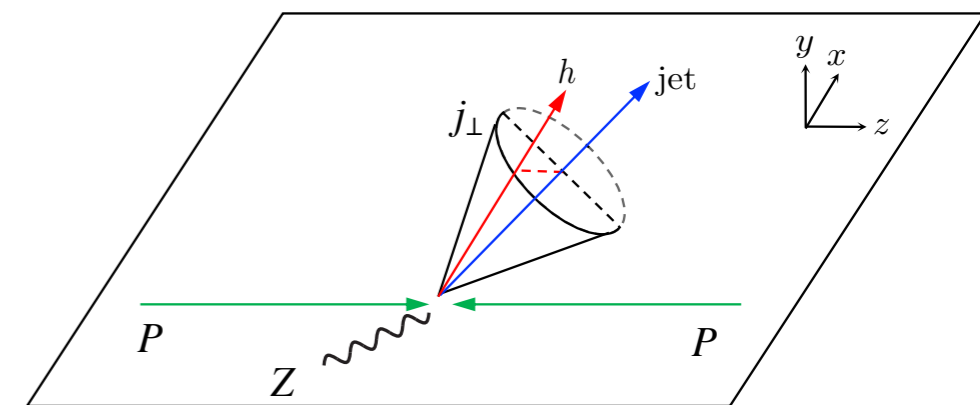
Agrees well in the  $0.1 < z_h < 0.5$  region



# Z-tagged jet



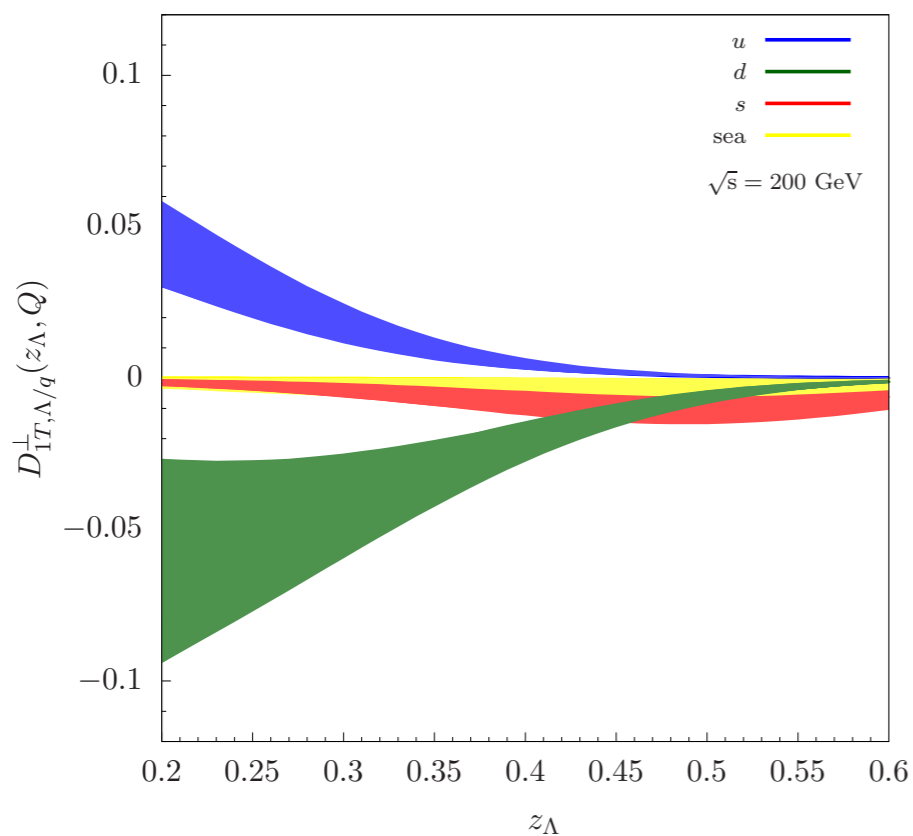
- LHCb collaboration measured collinear FFs and  $j_\perp$   
 Agrees well in the  $0.1 < z_h < 0.5$  region
- LHCb collaboration measured  $j_\perp$  recently.  
 but with entire  $0 < z_h < 1$



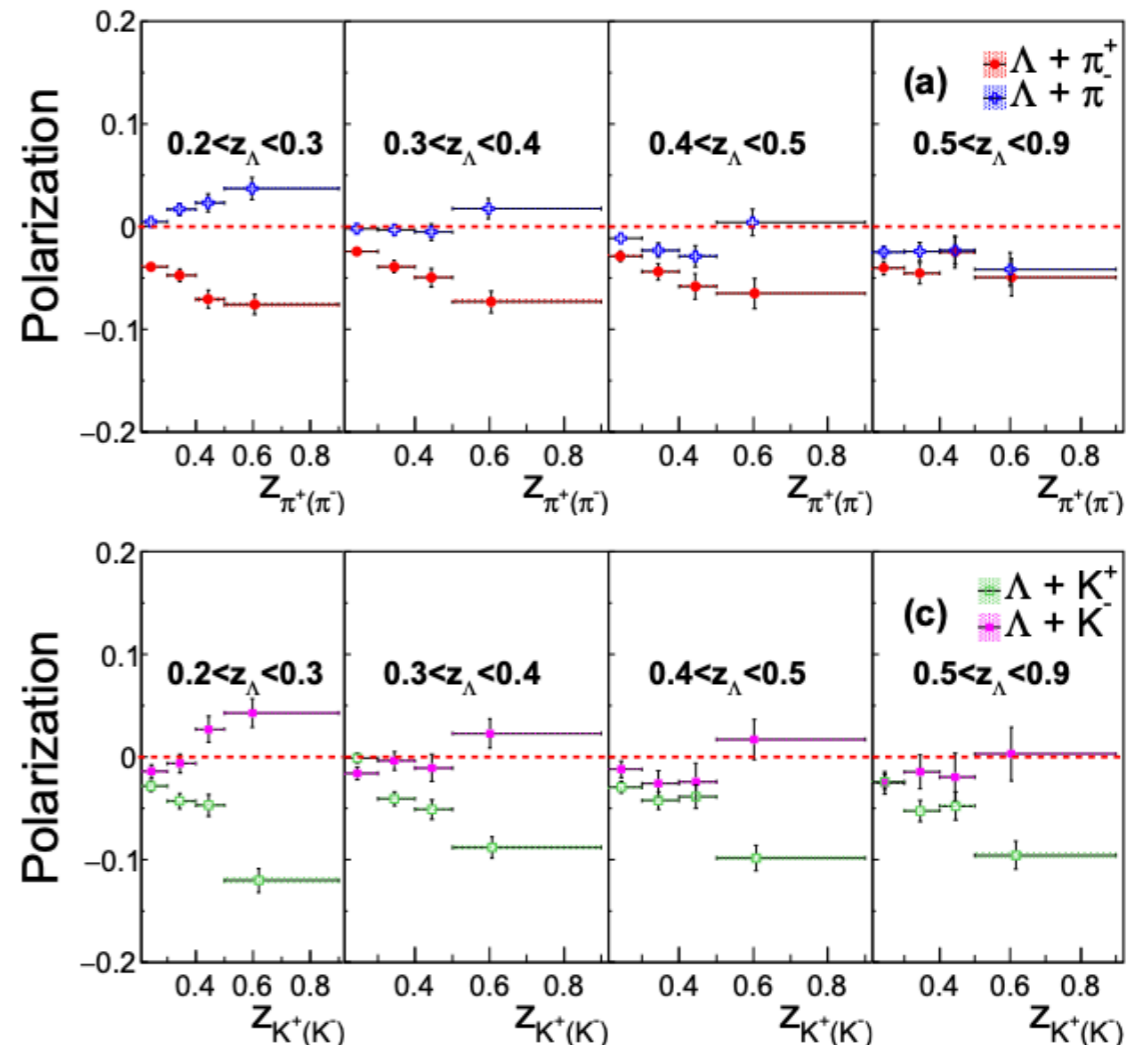
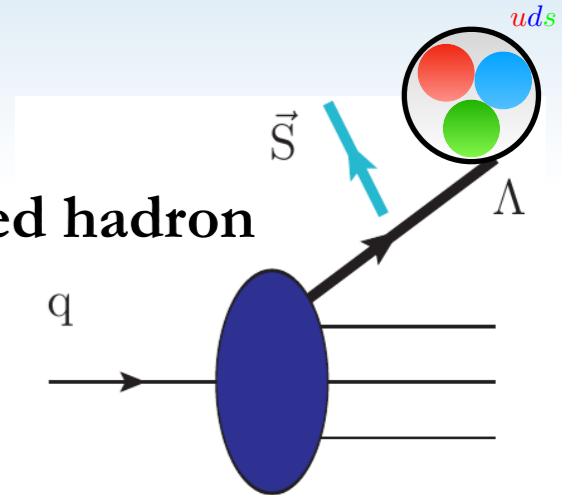
# Polarizing FF

## TMD Fragmentation Functions

		Quark polarization		
		U	L	T
Hadron polarization	U	$D^{h/q}$		$H^{\perp h/q}$
	L		$G^{h/q}$	$H_L^{\perp h/q}$
	T	$D_T^{\perp h/q}$	$G_T^{h/q}$	$H^{h/q}$ $H_T^{\perp h/q}$

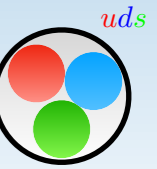


- Describes transversely polarized hadron inside unpolarized parton.



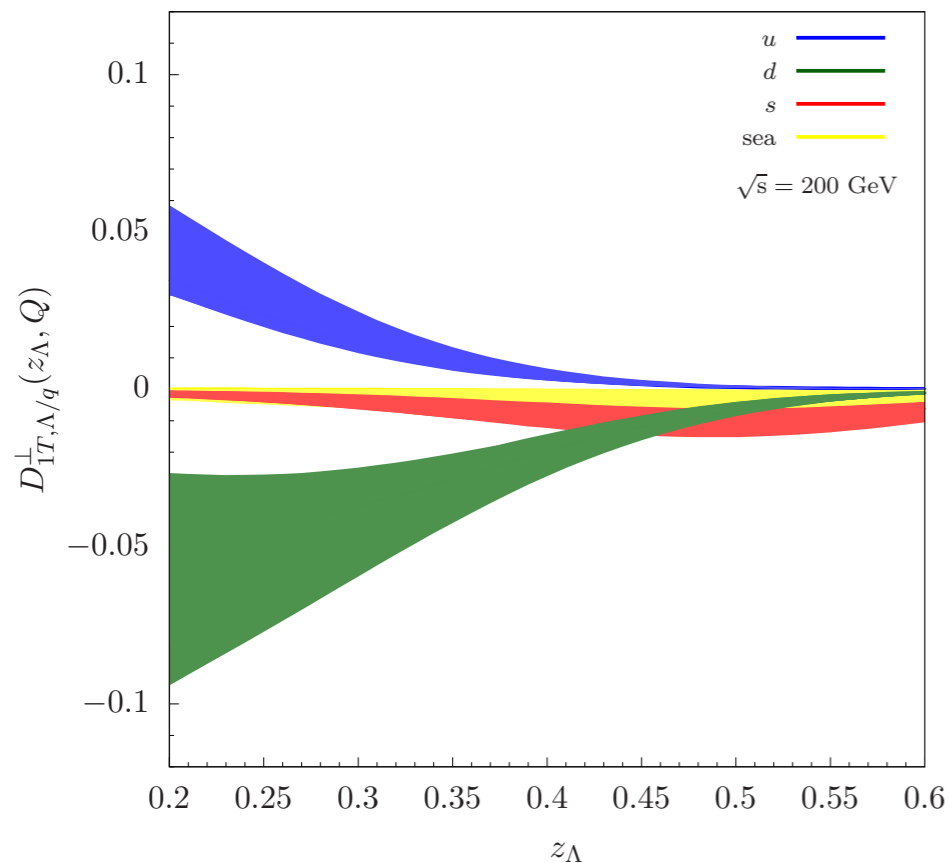
- Used PFF fits from Belle data  
Callos, Kang, Terry, '20

# Polarizing JFF



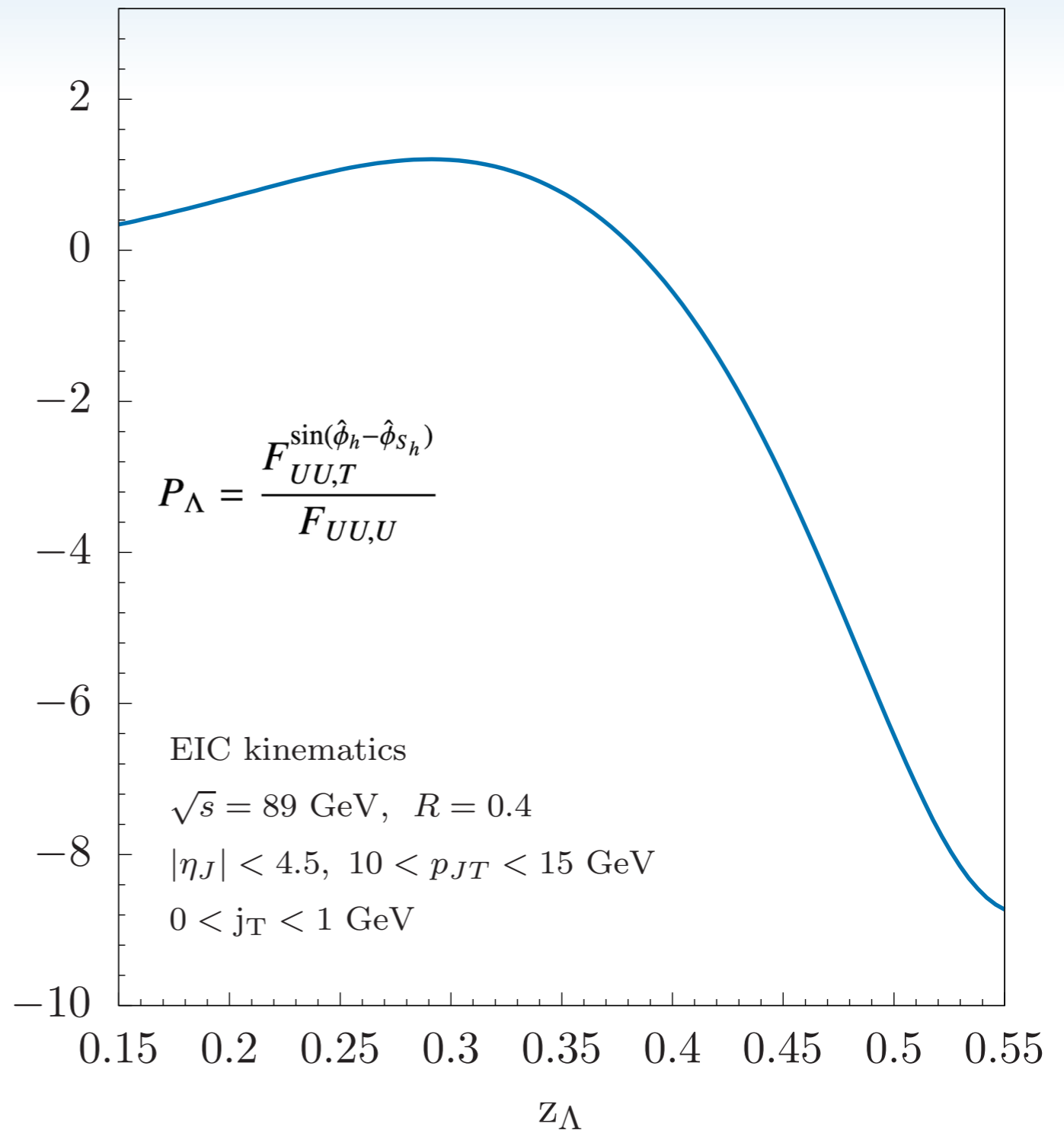
- Used PFF fits from Belle data

Callos, Kang, Terry, '20



- Predictions at the LHC kinematics
- Positive from up quark PFF at small  $z_\Lambda$
- Negative from down quark PFF at  $z_\Lambda \gtrsim 0.3$

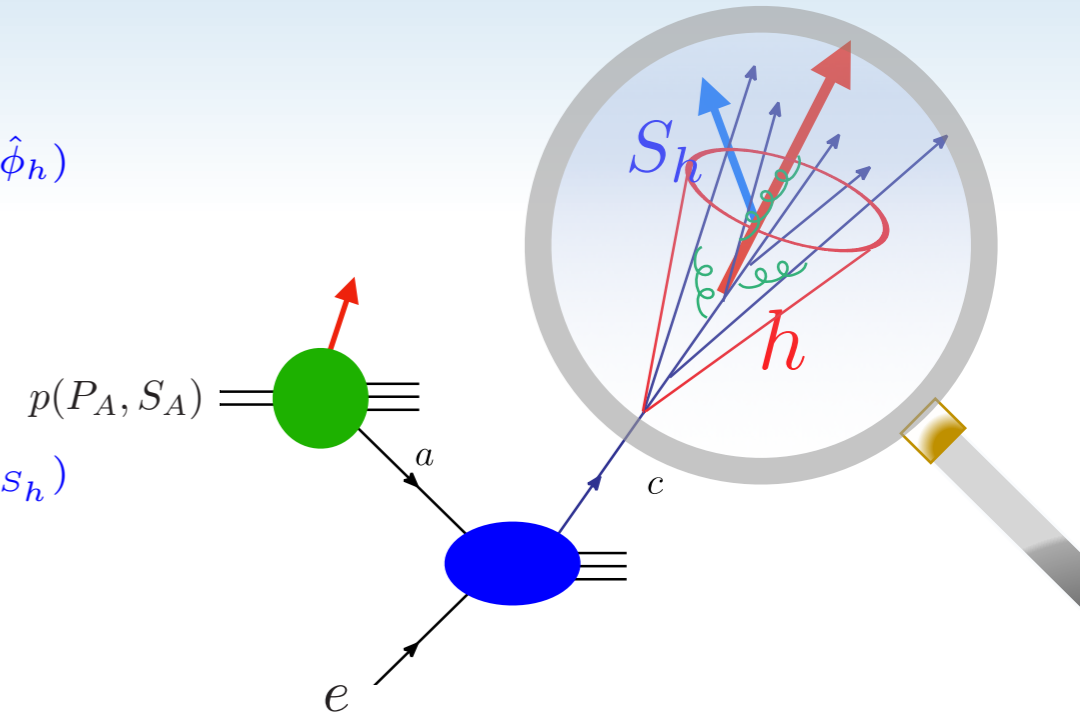
$P_\Lambda$  (%)



# Azimuthal angular dependence

$$\begin{aligned}
 \frac{d\sigma^{p(S_A)+p/e\rightarrow(\text{jet } h(S_h))X}}{dp_{JT}d\eta_J dz_h d^2\mathbf{j}_\perp} &= F_{UU,U} + |\mathbf{S}_T| \sin(\phi_{S_A} - \hat{\phi}_h) F_{TU,U}^{\sin(\phi_{S_A} - \hat{\phi}_h)} \\
 &+ \Lambda_h \left[ \lambda F_{LU,L} + |\mathbf{S}_T| \cos(\phi_{S_A} - \hat{\phi}_h) F_{TU,L}^{\cos(\phi_{S_A} - \hat{\phi}_h)} \right] \\
 &+ |\mathbf{S}_{h\perp}| \left\{ \sin(\hat{\phi}_h - \hat{\phi}_{S_h}) F_{UU,T}^{\sin(\hat{\phi}_h - \hat{\phi}_{S_h})} + \lambda \cos(\hat{\phi}_h - \hat{\phi}_{S_h}) F_{LU,T}^{\cos(\hat{\phi}_h - \hat{\phi}_{S_h})} \right. \\
 &\quad + |\mathbf{S}_T| \left( \cos(\phi_{S_A} - \hat{\phi}_{S_h}) F_{TU,T}^{\cos(\phi_{S_A} - \hat{\phi}_{S_h})} \right. \\
 &\quad \left. \left. + \cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_{S_A}) F_{TU,T}^{\cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_{S_A})} \right) \right\},
 \end{aligned}$$

$F_{S_A S_B, S_h}$   
 $\uparrow$   
**Polarization of**  $A, B, h$

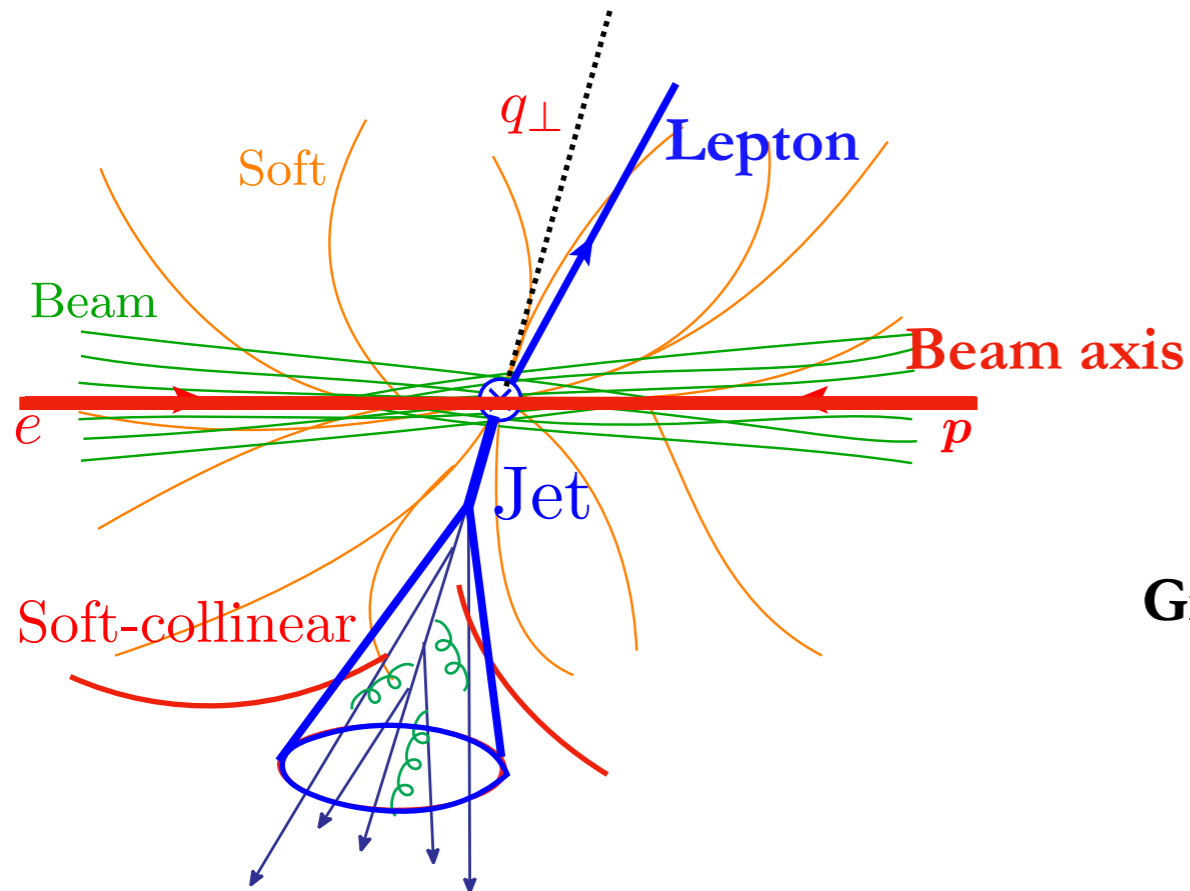


- Different structures come with different characteristic angular dependence.



# Lepton + Jet imbalance

- One of the simplest process  $e + P \rightarrow e + \text{Jet} + X$



$$q_{\perp} \equiv |\vec{p}_{e\perp} + \vec{p}_{J\perp}|, \quad p_{\perp} \equiv |\vec{p}_{e\perp} - \vec{p}_{J\perp}|/2$$

$q_{\perp} \ll p_{\perp}$ , **sensitive** to the large logs of  $\ln(q_{\perp}/p_{\perp})$  and **TMD structures** of the hadrons.

$$q_{\perp} = p_{X,\perp} = |\vec{k}_{c,\perp} + \vec{k}_{gs,\perp} + \vec{k}_{sc,\perp}|$$

Giving relevant modes :  $(+, -, \perp)$   $\lambda = q_{\perp}/p_{\perp}$

$$n\text{-collinear} \quad k_n \sim p_{\perp} (\lambda^2, 1, \lambda)_{n\bar{n}}$$

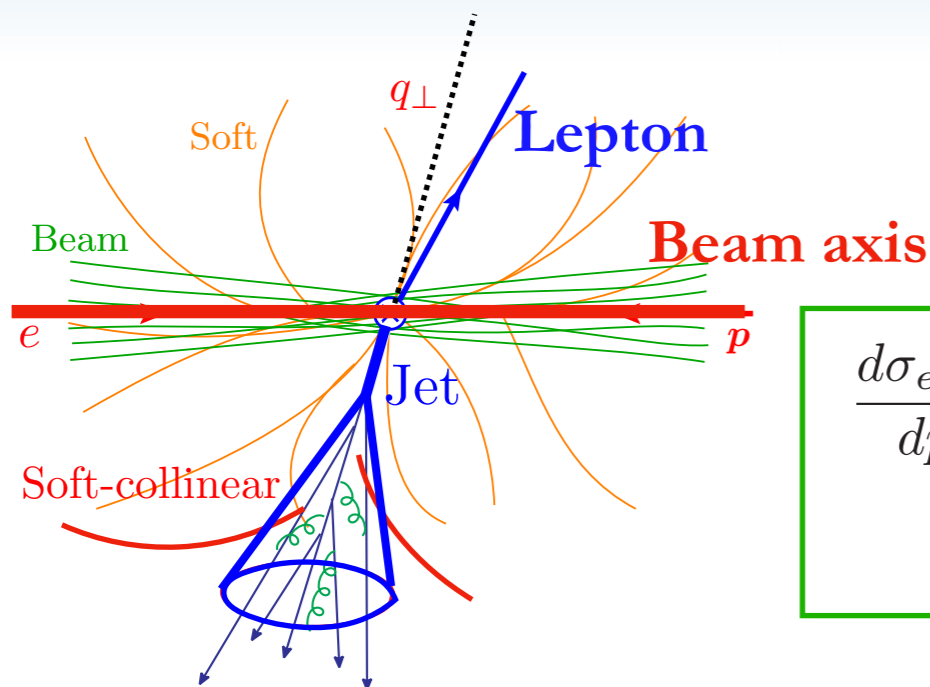
$$\text{global soft} \quad k_{gs} \sim p_{\perp} (\lambda, \lambda, \lambda)$$

$$\text{soft-collinear} \quad k_{sc} \sim p_{\perp} R(\lambda R, \lambda/R, \lambda)_{n_J, \bar{n}_J}$$

$$n_J\text{-collinear} \quad k_J \sim p_{\perp} (R^2, 1, R)_{n_J, \bar{n}_J}$$

2) Lepton + Jet imbalance  
TMDPDFs

# Lepton + Jet imbalance



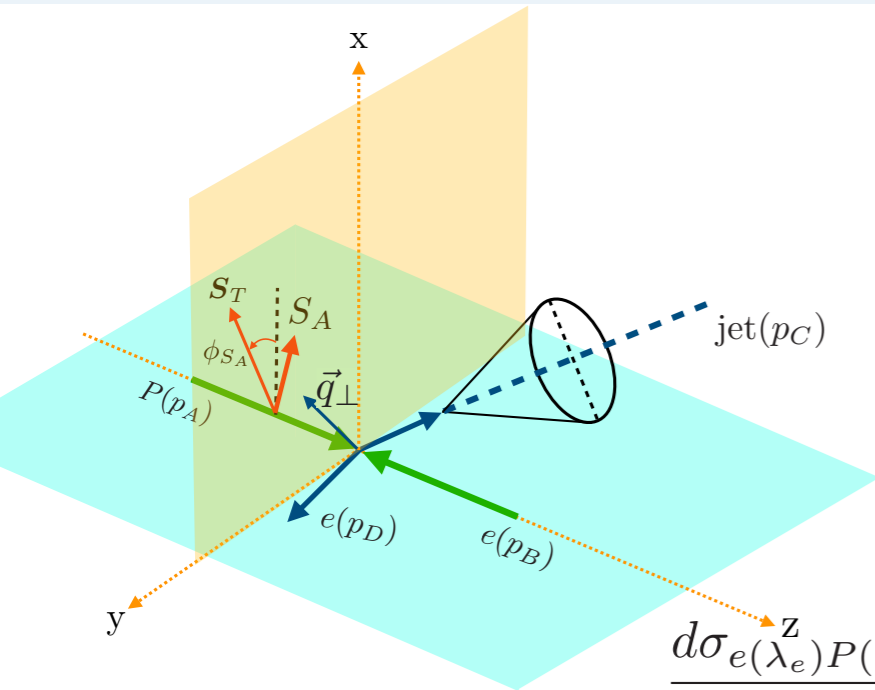
$$\frac{d\sigma_{eP \rightarrow e+\text{jet}}}{dp_{\perp} dq_{\perp}} = \int \prod_i^3 d^2 k_{i\perp} H(Q) \delta^{(2)}(\vec{k}_{1\perp} + \vec{k}_{2\perp} + \vec{k}_{3\perp} - q_{\perp}) \times f_a(x, \vec{k}_{1\perp}) S^{\text{global}}(\vec{k}_{2\perp}) S_{J_c}(\vec{k}_{3\perp}) J_c(p_{\perp} R)$$

We arrive at factorization using SCET

$n$ -collinear	$k_n \sim p_{\perp} (\lambda^2, 1, \lambda)_{n\bar{n}}$	} TMDPDFs
global soft	$k_{gs} \sim p_{\perp} (\lambda, \lambda, \lambda)$	
soft-collinear	$k_{sc} \sim p_{\perp} R (\lambda R, \lambda/R, \lambda)_{n_J, \bar{n}_J}$	} Soft functions
$n_J$ -collinear	$k_J \sim p_{\perp} (R^2, 1, R)_{n_J, \bar{n}_J}$	
		} Jet function

Liu, Ringer, Vogelsang, Yuan '18, '20  
Arratia, Kang, Prokudin, Ringer '20

# Lepton + Jet imbalance



$$\frac{d\sigma_{eP \rightarrow e+jet}}{dp_{\perp} dq_{\perp}} = \int \prod_i^3 d^2 k_{i\perp} H(Q) \delta^{(2)}(\vec{k}_{1\perp} + \vec{k}_{2\perp} + \vec{k}_{3\perp} - q_{\perp}) \times f_a(x, \vec{k}_{1\perp}) S^{\text{global}}(\vec{k}_{2\perp}) S_{J_c}(\vec{k}_{3\perp}) J_c(p_{\perp} R)$$

$$\frac{d\sigma_{e(\lambda_e)P(S) \rightarrow e+jet}^Z}{dp_{\perp} dq_{\perp}} \sim f_1 \sim g_{1L} = \underbrace{F_{UU}}_{\text{circled}} + \lambda_p \lambda_e F_{LL} + S_T \left\{ \sin(\phi_{S_A} - \phi_q) F_{TU}^{\sin(\phi_{S_A} - \phi_q)} + \lambda_e \cos(\phi_{S_A} - \phi_q) F_{TL}^{\cos(\phi_{S_A} - \phi_q)} \right\},$$

$$\sim f_{1T}^{\perp} \sim g_{1T}$$

**Leading Twist TMDs**      Nucleon Spin      Quark Spin

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \text{circle with red dot}$		$h_1^{\perp} = \text{circle with red dot} - \text{circle with red dot}$ Boer-Mulders
	L		$g_{1L} = \text{circle with red arrow} \rightarrow - \text{circle with red arrow} \rightarrow$ Helicity	$h_{1L}^{\perp} = \text{circle with red arrow} \rightarrow - \text{circle with red arrow} \rightarrow$ Worm gear
	T	$f_{1T}^{\perp} = \text{circle with red dot} \uparrow - \text{circle with red dot} \downarrow$ Sivers	$g_{1T} = \text{circle with red arrow} \uparrow - \text{circle with red arrow} \downarrow$ Worm gear	$h_1 = \text{circle with red dot} \uparrow - \text{circle with red dot} \downarrow$ Transversity $h_{1T}^{\perp} = \text{circle with red arrow} \uparrow - \text{circle with red arrow} \downarrow$

- With jet, only sensitive to single TMDs (compared to standard processes)
- We do not get sensitivity to all TMDPDFs (only to chiral-even TMDPDFs)

Liu, Ringer, Vogelsang, Yuan '18, '20  
Arratia, Kang, Prokudin, Ringer '20  
Kang, KL, Shao, Zhao '21

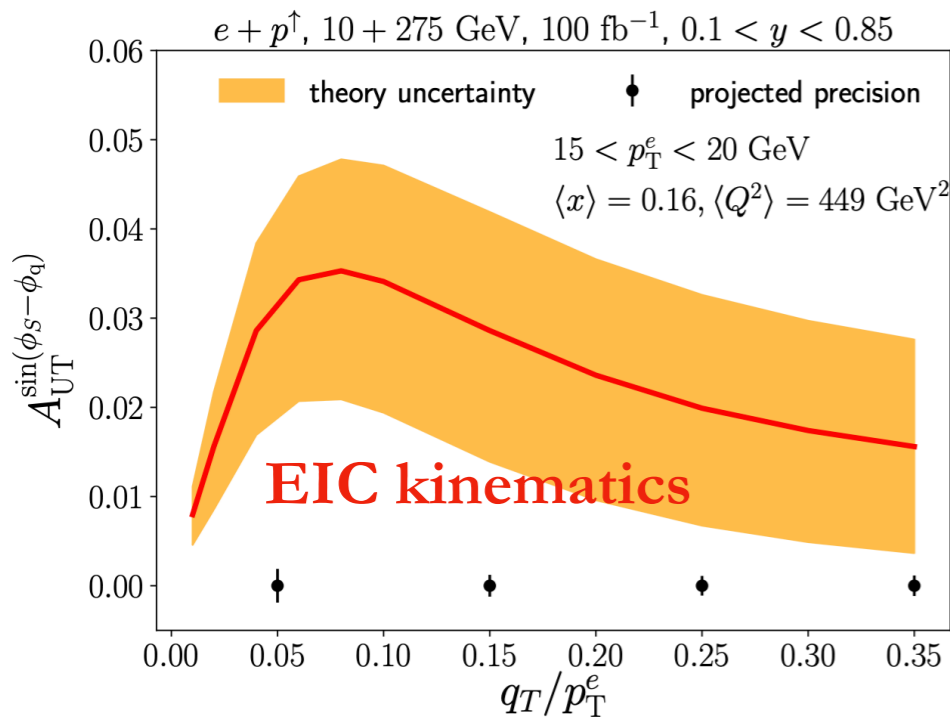
# Sivers asymmetry

$$\frac{d\sigma_{e(\lambda_e)P(S)\rightarrow e+\text{jet}}}{dp_{\perp}dq_{\perp}} \sim f_1 \quad \sim g_{1L}$$

$$= F_{UU} + \lambda_p \lambda_e F_{LL}$$

$$+ S_T \left\{ \sin(\phi_{S_A} - \phi_q) F_{TU}^{\sin(\phi_{S_A} - \phi_q)} + \lambda_e \cos(\phi_{S_A} - \phi_q) F_{TL}^{\cos(\phi_{S_A} - \phi_q)} \right\},$$

$\sim f_{1T}^{\perp}$   $\sim g_{1T}$



$$f_a(x_a, k_{\perp}) \rightarrow \frac{\epsilon_{\perp}^{\rho\sigma} S_{\perp\rho} k_{\perp\sigma}}{M} f_{aT}^{\perp}(x_a, k_{\perp})$$

$$\frac{d\Delta\sigma_{eP\rightarrow e+\text{jet}}}{dp_{\perp}dq_{\perp}} = \frac{\epsilon_{\perp}^{\rho\sigma} S_{\perp\rho}}{M} \int \prod_i^3 d^2k_{i\perp} H(Q) \delta^{(2)}(\vec{k}_{1\perp} + \vec{k}_{2\perp} + \vec{k}_{3\perp} - q_{\perp})$$

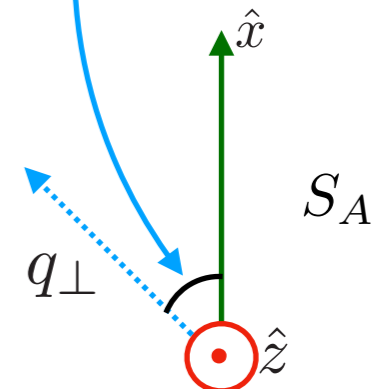
$$\times k_{1\perp\sigma} f_{aT}^{\perp}(x_a, k_{1\perp}) S^{\text{global}}(\vec{k}_{2\perp}) S_{J_c}(\vec{k}_{3\perp}) J_c(p_{\perp}R)$$

$$\propto \sin(\phi_{S_A} - \phi_q)$$

$$A_{UT}^{\sin(\phi_{S_A} - \phi_q)} = \frac{F_{UT}^{\sin(\phi_{S_A} - \phi_q)}}{F_{UU}}$$

- **Positive  $\Delta\sigma \implies$  a preference of imbalance to be on left**

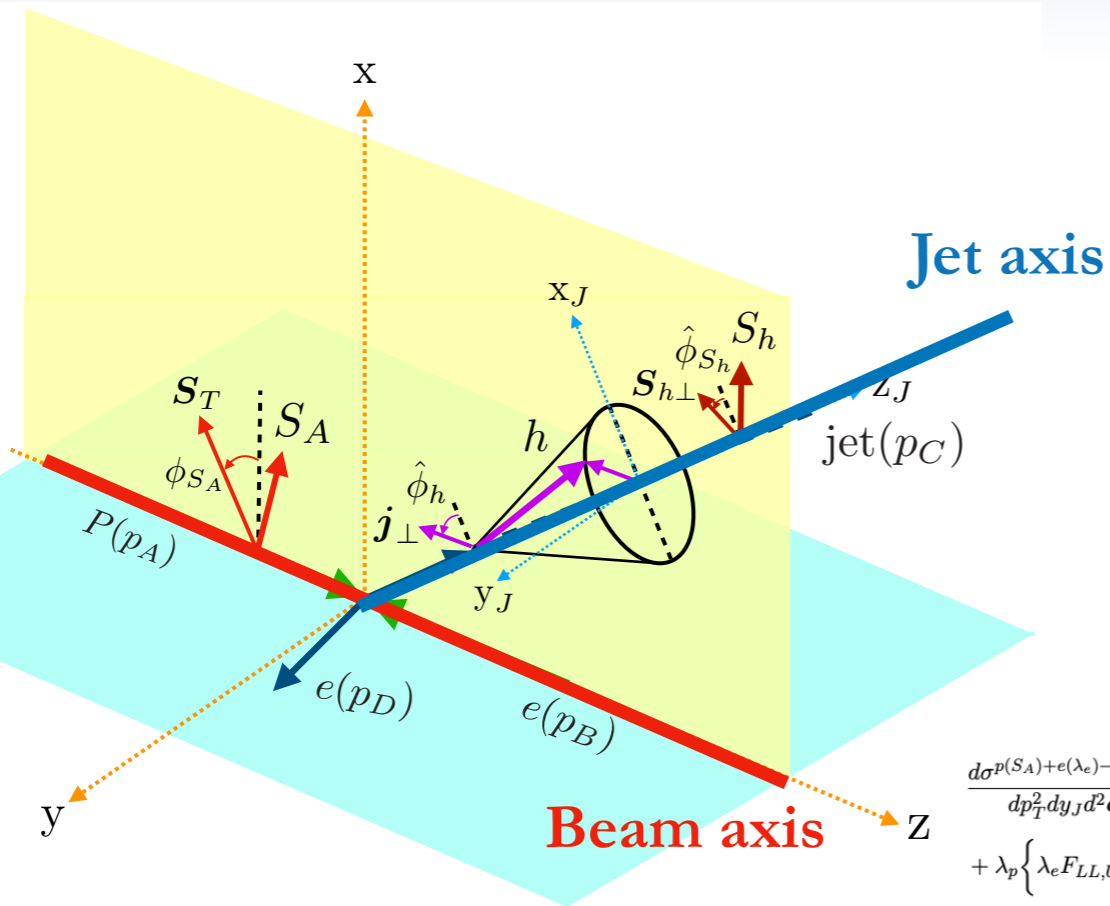
(When polarized proton moving towards us and transverse spin pointing up.)



**Sivers** from SIDIS extraction  
Echevarria, Idilbi, Kang, Vitev, '14

Arratia, Kang, Prokudin, Ringer '20  
Kang, KL, Shao, Zhao '21

# Polarized Jet Fragmentation Functions and lepton + jet imbalance



- Observation of polarized hadron inside jet gives sensitivity to **all** TMDPDFs and TMDFFs. (analogous correlations to standard SIDIS)
- Sensitivity to two TMDs, but sensitive to  $\vec{q}_\perp$  and  $\vec{j}_\perp$  separately (**advantage of two axes**)

## Many characteristic correlations

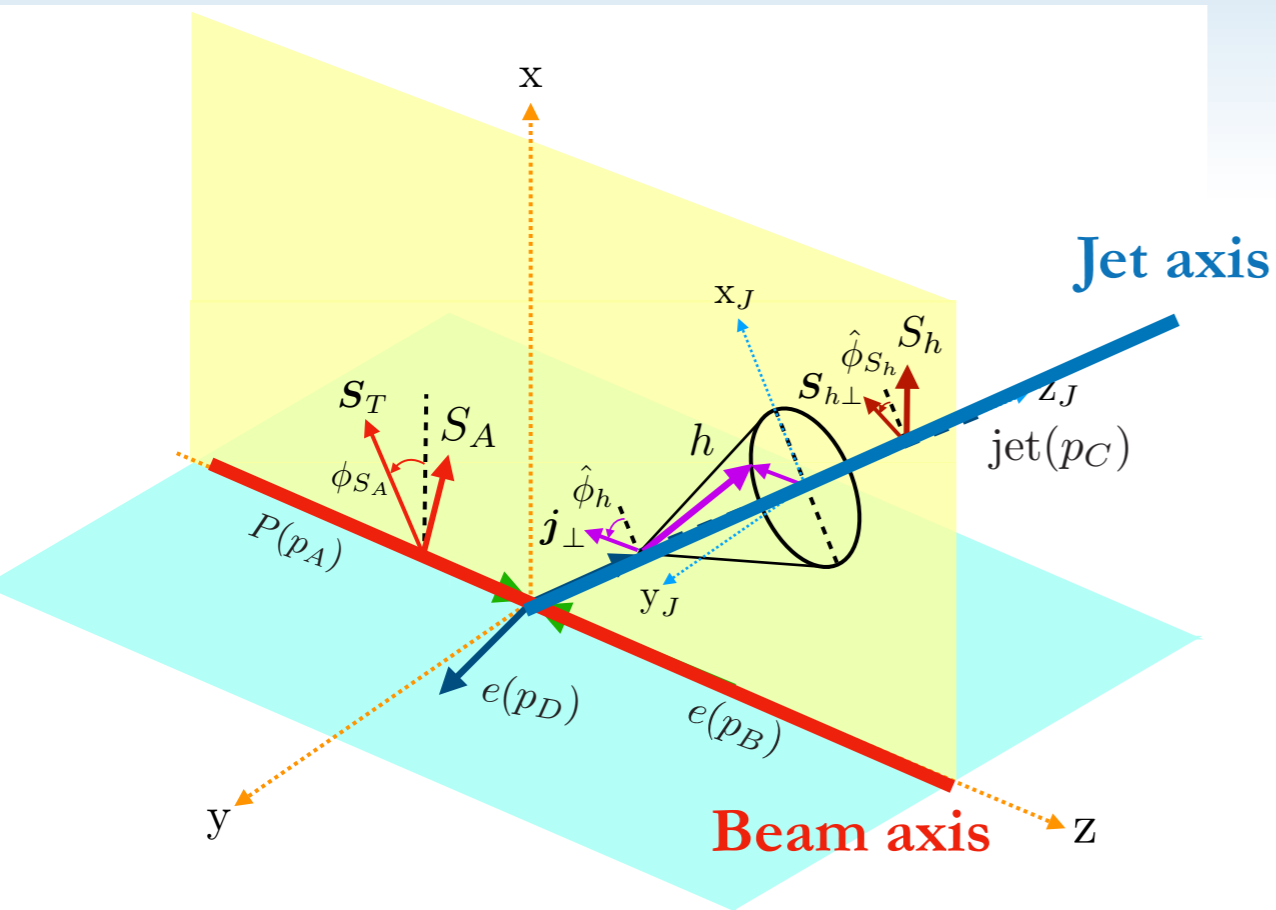
$$\begin{aligned}
 \frac{d\sigma^{p(S_A)+e(\lambda_e)\rightarrow e+(\text{jet } h(S_h))+X}}{dp_T^2 dy_J d^2 q_T dz_h d^2 j_\perp} &= F_{UU,U} + \cos(\phi_q - \hat{\phi}_h) F_{UU,U}^{\cos(\phi_q - \hat{\phi}_h)} \\
 &+ \lambda_p \left\{ \lambda_e F_{LL,U} + \sin(\phi_q - \hat{\phi}_h) F_{LU,U}^{\sin(\phi_q - \hat{\phi}_h)} \right\} \\
 &+ S_T \left\{ \sin(\phi_q - \phi_{S_A}) F_{TU,U}^{\sin(\phi_q - \phi_{S_A})} + \lambda_e \cos(\phi_q - \phi_{S_A}) F_{TL,U}^{\cos(\phi_q - \phi_{S_A})} \right. \\
 &\quad \left. + \sin(\phi_{S_A} - \hat{\phi}_h) F_{TU,U}^{\sin(\phi_{S_A} - \hat{\phi}_h)} + \sin(2\phi_q - \hat{\phi}_h - \phi_{S_A}) F_{TU,U}^{\sin(2\phi_q - \hat{\phi}_h - \phi_{S_A})} \right\} \\
 &+ \lambda_h \left\{ \lambda_e F_{UL,L} + \sin(\hat{\phi}_h - \phi_q) F_{UU,L}^{\sin(\hat{\phi}_h - \phi_q)} + \lambda_p \left[ F_{LU,L} + \cos(\hat{\phi}_h - \phi_q) F_{LU,L}^{\cos(\hat{\phi}_h - \phi_q)} \right] \right. \\
 &\quad \left. + S_T \left[ \cos(\phi_q - \phi_{S_A}) F_{TU,L}^{\cos(\phi_q - \phi_{S_A})} + \lambda_e \sin(\phi_q - \phi_{S_A}) F_{TL,L}^{\sin(\phi_q - \phi_{S_A})} \right. \right. \\
 &\quad \left. \left. + \cos(\phi_{S_A} - \hat{\phi}_h) F_{TU,L}^{\cos(\phi_{S_A} - \hat{\phi}_h)} + \cos(2\phi_q - \phi_{S_A} - \hat{\phi}_h) F_{TU,L}^{\cos(2\phi_q - \phi_{S_A} - \hat{\phi}_h)} \right] \right\} \\
 &+ S_{h\perp} \left\{ \sin(\hat{\phi}_h - \hat{\phi}_{S_h}) F_{UU,T}^{\sin(\hat{\phi}_h - \hat{\phi}_{S_h})} + \lambda_e \cos(\hat{\phi}_h - \hat{\phi}_{S_h}) F_{UL,T}^{\cos(\hat{\phi}_h - \hat{\phi}_{S_h})} \right. \\
 &\quad \left. + \sin(\hat{\phi}_{S_h} - \phi_q) F_{UU,T}^{\sin(\hat{\phi}_{S_h} - \phi_q)} + \sin(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_q) F_{UU,T}^{\sin(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_q)} \right. \\
 &\quad \left. + \lambda_p \left[ \cos(\hat{\phi}_h - \hat{\phi}_{S_h}) F_{LU,T}^{\cos(\hat{\phi}_h - \hat{\phi}_{S_h})} + \cos(\phi_q - \hat{\phi}_{S_h}) F_{LU,T}^{\cos(\phi_q - \hat{\phi}_{S_h})} \right. \right. \\
 &\quad \left. \left. + \cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_q) F_{LU,T}^{\cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_q)} + \lambda_e \sin(\hat{\phi}_h - \hat{\phi}_{S_h}) F_{LL,T}^{\sin(\hat{\phi}_h - \hat{\phi}_{S_h})} \right] \right. \\
 &\quad \left. + S_T \left[ \cos(\phi_{S_A} - \hat{\phi}_{S_h}) F_{TU,T}^{\cos(\phi_{S_A} - \hat{\phi}_{S_h})} + \cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_{S_A}) F_{TU,T}^{\cos(2\hat{\phi}_h - \hat{\phi}_{S_h} - \phi_{S_A})} \right. \right. \\
 &\quad \left. \left. + \sin(\hat{\phi}_h - \hat{\phi}_{S_h}) \sin(\phi_q - \phi_{S_A}) F_{TU,T}^{\sin(\hat{\phi}_h - \hat{\phi}_{S_h}) \sin(\phi_q - \phi_{S_A})} \right. \right. \\
 &\quad \left. \left. + \cos(\hat{\phi}_h - \hat{\phi}_{S_h}) \cos(\phi_q - \phi_{S_A}) F_{TU,T}^{\cos(\hat{\phi}_h - \hat{\phi}_{S_h}) \cos(\phi_q - \phi_{S_A})} \right. \right. \\
 &\quad \left. \left. + \cos(2\phi_q - \phi_{S_A} - \hat{\phi}_{S_h}) F_{TU,T}^{\cos(2\phi_q - \phi_{S_A} - \hat{\phi}_{S_h})} \right. \right. \\
 &\quad \left. \left. + \cos(2\hat{\phi}_h - \hat{\phi}_{S_h} + 2\phi_q - \phi_{S_A}) F_{TU,T}^{\cos(2\hat{\phi}_h - \hat{\phi}_{S_h} + 2\phi_q - \phi_{S_A})} \right. \right. \\
 &\quad \left. \left. + \lambda_e \cos(\hat{\phi}_h - \hat{\phi}_{S_h}) \sin(\phi_{S_A} - \phi_q) F_{TL,T}^{\cos(\hat{\phi}_h - \hat{\phi}_{S_h}) \sin(\phi_{S_A} - \phi_q)} \right. \right. \\
 &\quad \left. \left. + \lambda_e \sin(\hat{\phi}_h - \hat{\phi}_{S_h}) \cos(\phi_{S_A} - \phi_q) F_{TL,T}^{\sin(\hat{\phi}_h - \hat{\phi}_{S_h}) \cos(\phi_{S_A} - \phi_q)} \right] \right\},
 \end{aligned}$$

3) Lepton + Jet imbalance  
with hadron in jet  
TMDFFs / TMDPDFs

# Phenomenology : $A_{UU,U}^{\cos(\phi_q - \hat{\phi}_h)}$

$$A_{UU,U}^{\cos(\phi_q - \hat{\phi}_h)} \equiv \frac{F_{UU,U}^{\cos(\phi_q - \hat{\phi}_h)}(q_{\perp}, j_{\perp})}{F_{UU,U}(q_{\perp}, j_{\perp})} \sim \frac{h_1^{\perp}(q_{\perp})H_1^{\perp}(j_{\perp})}{f_1(q_{\perp})D_1(j_{\perp})}$$

- Boer-Mulders and Collins functions sensitive to transverse momentum measured with respect to different axes.
- “**Separation**” of the incoming and outgoing dynamics.



Leading Twist TMDs

○ → Nucleon Spin    ⊙ → Quark Spin

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \odot$		$h_1^{\perp} = \odot - \ominus$ Boer-Mulders
	L		$g_{1L} = \odot \rightarrow - \ominus \rightarrow$ Helicity	$h_{1L}^{\perp} = \odot \rightarrow - \ominus \rightarrow$
	T	$f_{1T}^{\perp} = \odot \uparrow - \ominus \downarrow$ Sivers	$g_{1T} = \odot \uparrow - \ominus \uparrow$	$h_1 = \odot \uparrow - \ominus \uparrow$ Transversity $h_{1T}^{\perp} = \odot \uparrow - \ominus \uparrow$

Quark polarization

Hadron polarization	Quark polarization			
	U	L	T	
	U	$D^{h/q}$		$H^{\perp h/q}$ Collins
	L		$G^{h/q}$	$H_L^{\perp h/q}$
T	$D_T^{\perp h/q}$	$G_T^{h/q}$	$H^{h/q}$ $H_T^{\perp h/q}$	

# Unpolarized $\pi$ in jet (Boer-Mulders, Collins)

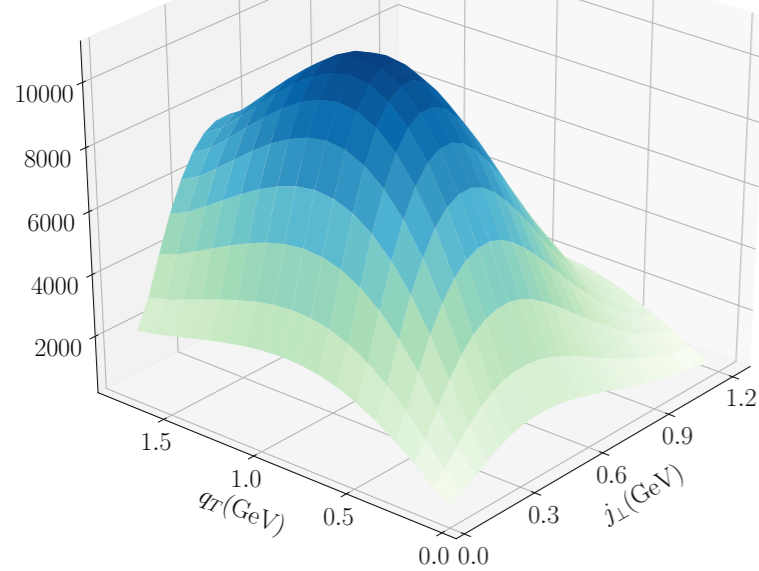
$\pi^-$   $q_T [0, 1.8], j_T [0, 1.2]$

Denominator

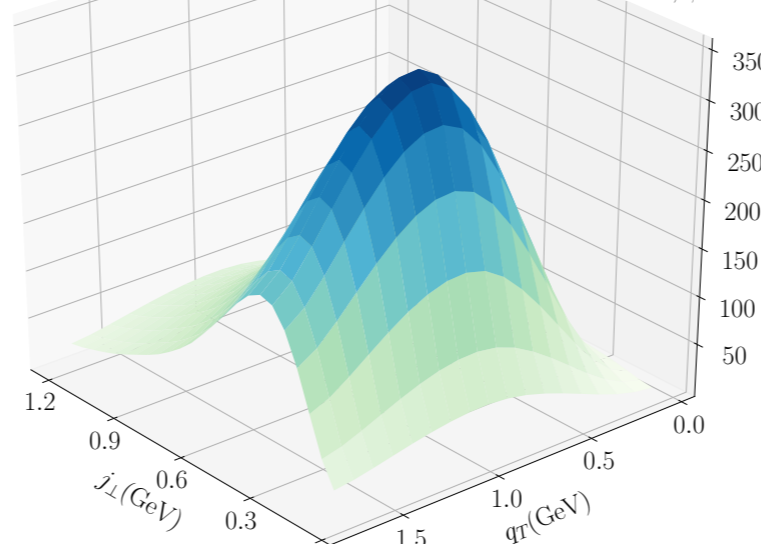
Numerator

$$A^{\cos(\phi_q - \hat{\phi}_h)} \equiv \frac{F_{UU,U}^{\cos(\phi_q - \hat{\phi}_h)}(q_\perp, j_\perp)}{F_{UU,U}(q_\perp, j_\perp)} \sim \frac{h_1^\perp(q_\perp)H_1^\perp(j_\perp)}{f_1(q_\perp)D_1(j_\perp)}$$

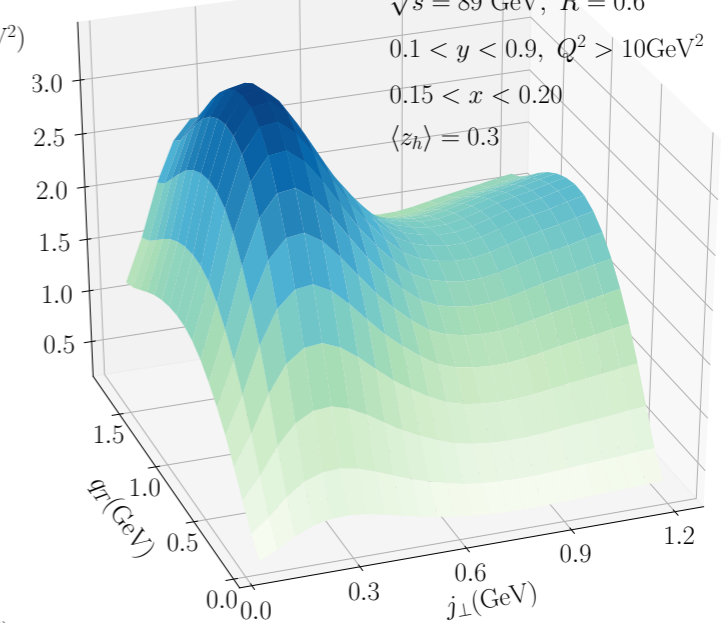
$F_{UU,U;\pi^-}$  (pb/GeV<sup>2</sup>)



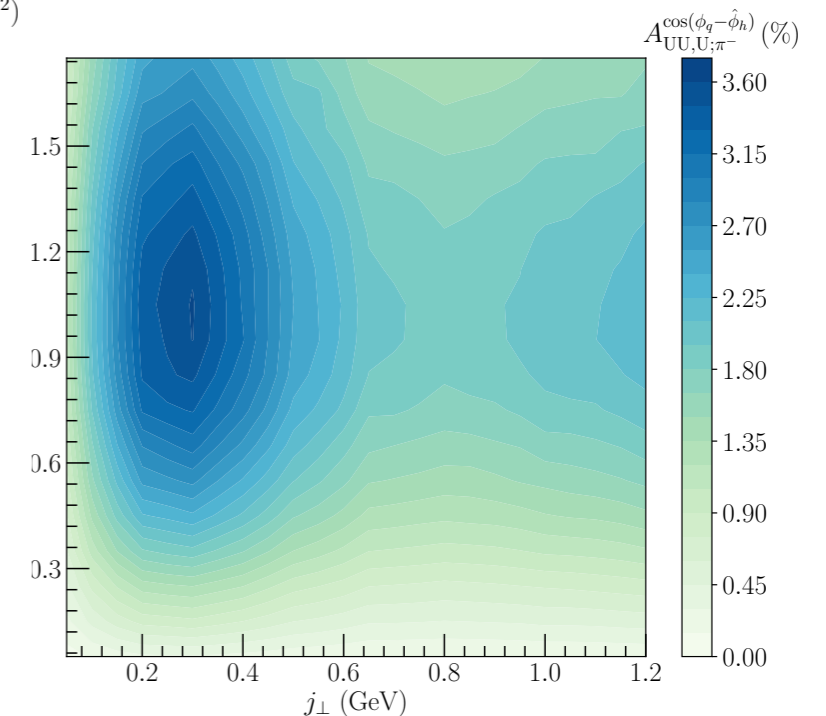
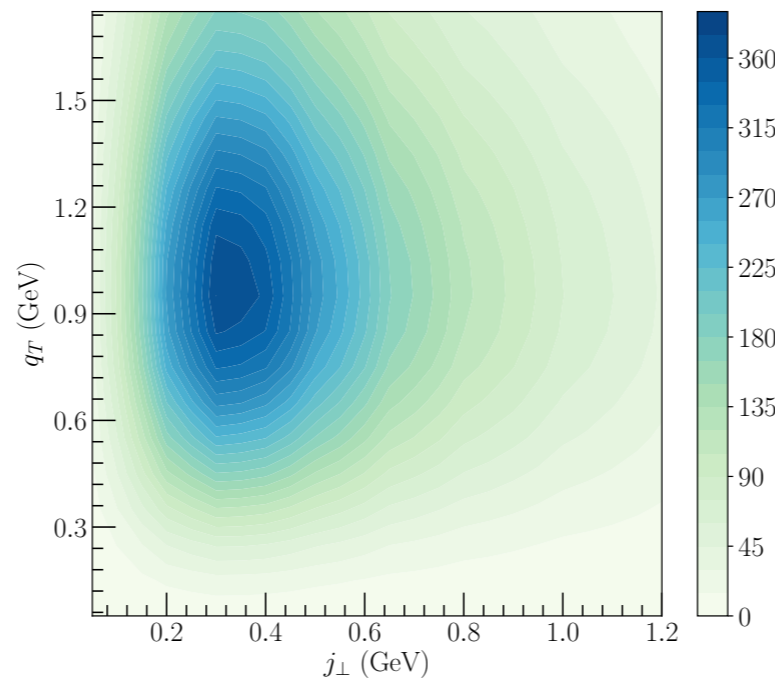
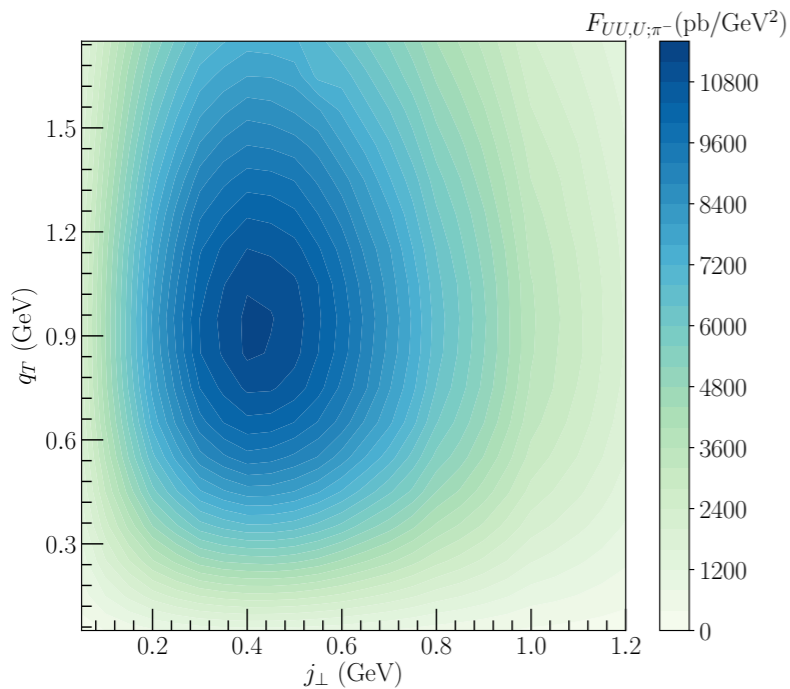
$F_{UU,U;\pi^-}^{\cos(\phi_q - \hat{\phi}_h)}$  (pb/GeV<sup>2</sup>)



$A_{UU,U;\pi^-}^{\cos(\phi_q - \hat{\phi}_h)}$  (%)



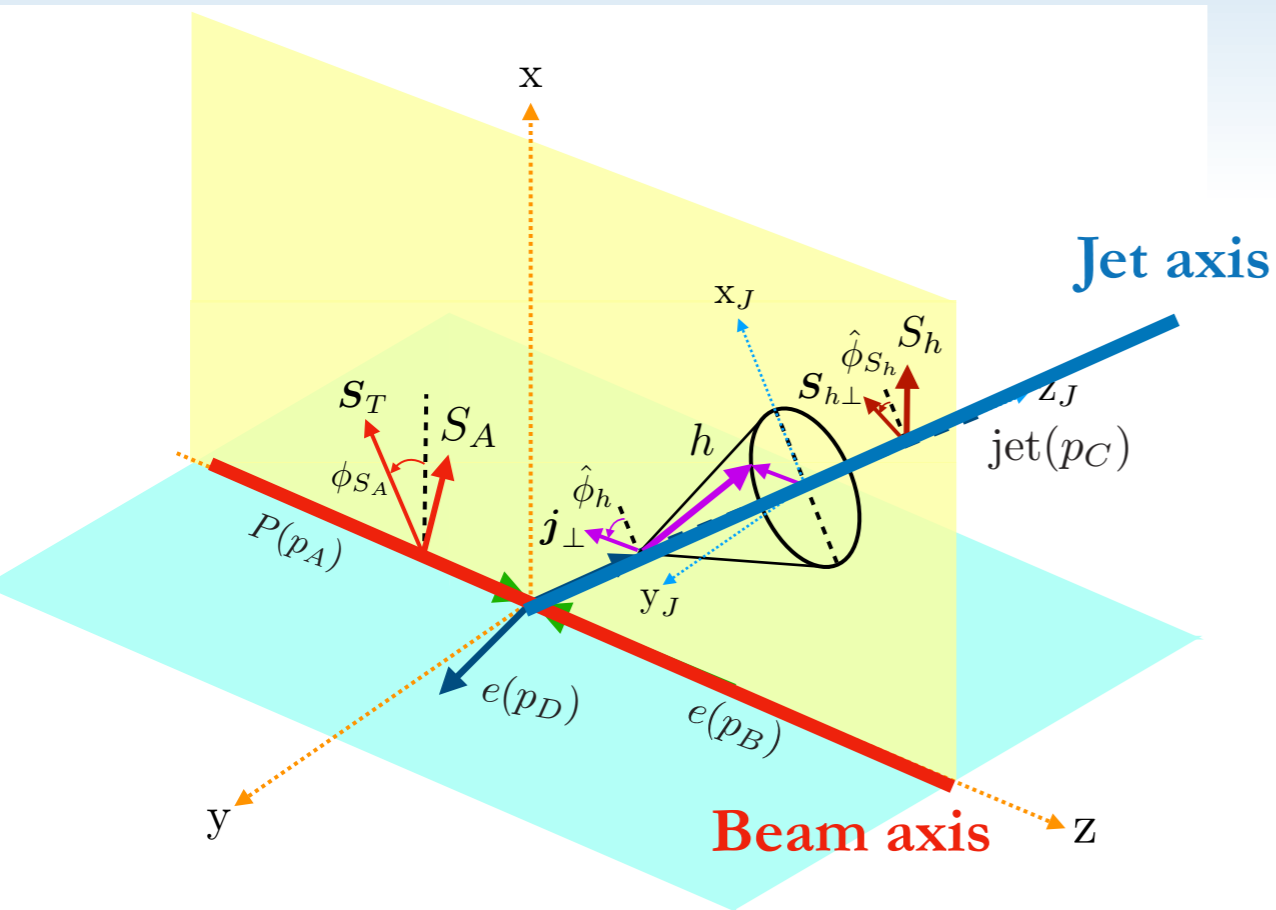
$F_{UU,U;\pi^-}^{\cos(\phi_q - \hat{\phi}_h)}$  (pb/GeV<sup>2</sup>)



Parametrization from *Barone, Melis, Prokudin '10* (Boer-Mulders)  
*Kang, Prokudin, Sun, Yuan '15* (Collins)

*Kang, KL, Shao, Zhao '21*

# Phenomenology : $A_{UU,T}^{\sin(\hat{\phi}_\Lambda - \hat{\phi}_{S_\Lambda})}$



$$A_{UU,T}^{\sin(\hat{\phi}_\Lambda - \hat{\phi}_{S_\Lambda})} = \frac{F_{UU,T}^{\sin(\hat{\phi}_\Lambda - \hat{\phi}_{S_\Lambda})}}{F_{UU,U}} \sim \frac{f_1(q_\perp) D_{1T}^\perp(j_\perp)}{f_1(q_\perp) D_1(j_\perp)}$$

- **“Separation”** of the incoming and outgoing dynamics cancel the  $q_T$  dependence for this case.

Leading Twist TMDs

○ → Nucleon Spin    ⊙ → Quark Spin

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \odot$		$h_1^\perp = \odot - \ominus$ Boer-Mulders
	L		$g_{1L} = \odot \rightarrow - \ominus \rightarrow$ Helicity	$h_{1L}^\perp = \odot \rightarrow - \ominus \rightarrow$
	T	$f_{1T}^\perp = \odot \uparrow - \ominus \downarrow$ Sivers	$g_{1T} = \odot \rightarrow - \ominus \rightarrow$	$h_1 = \odot \uparrow - \ominus \downarrow$ Transversity $h_{1T}^\perp = \odot \rightarrow - \ominus \rightarrow$

Quark polarization

		U	L	T
Hadron polarization	U	$D^{h/q}$		$H^\perp h/q$
	L		$G^{h/q}$	$H_L^\perp h/q$
	T	$D_T^\perp h/q$	$G_T^{h/q}$	$H^{h/q} H_T^\perp h/q$

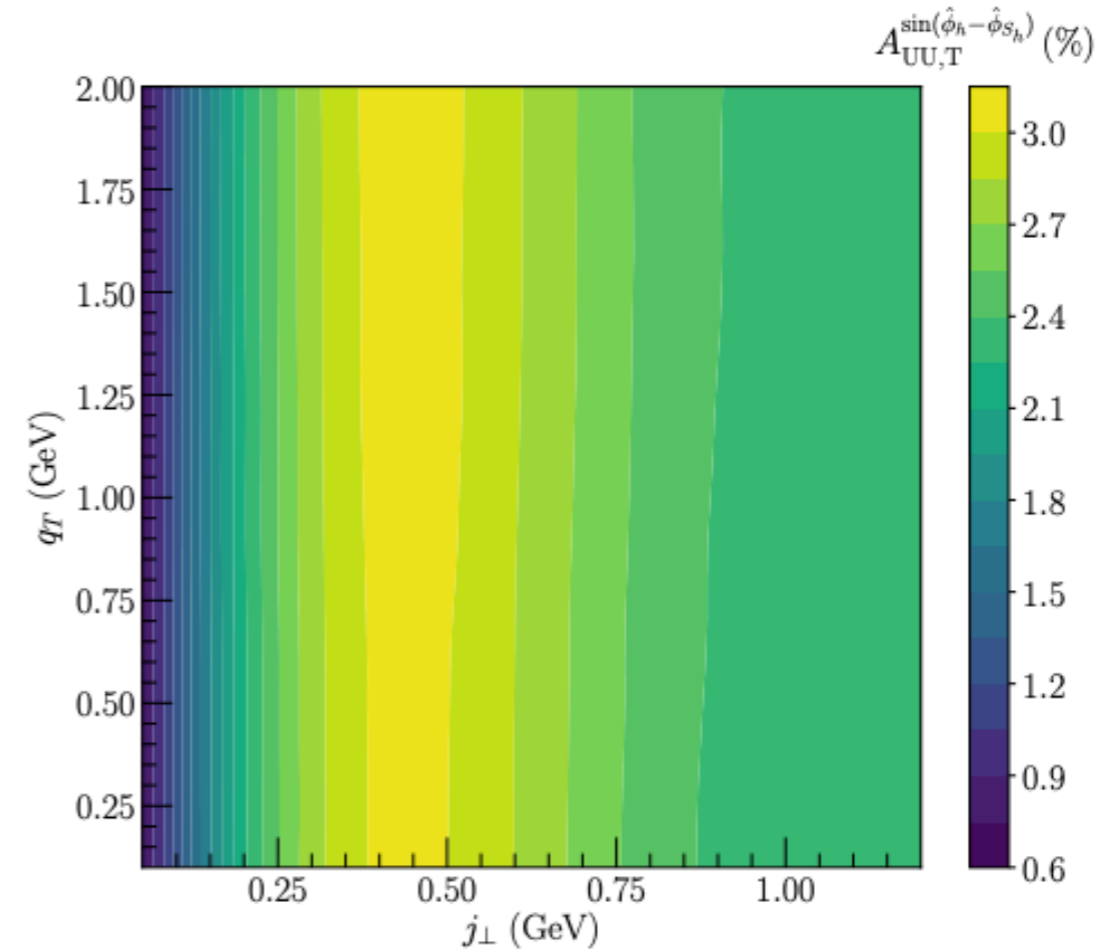
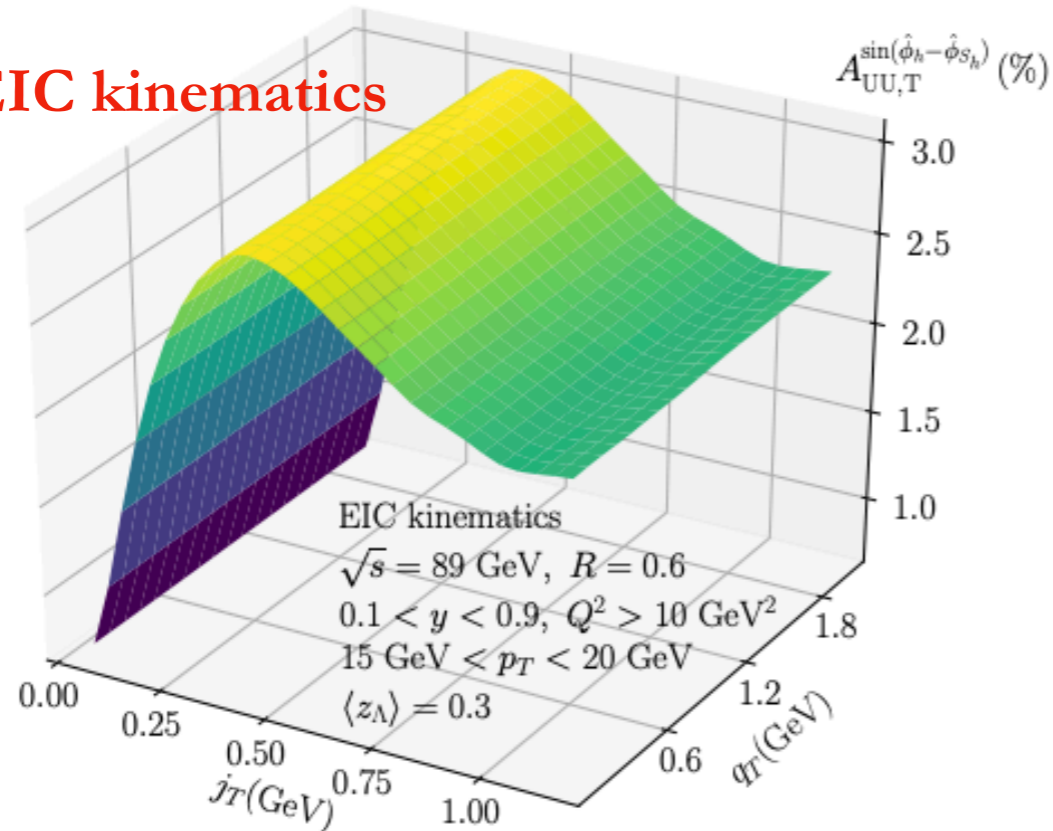
Polarizing FF



# Phenomenology : $A_{UU,T}^{\sin(\hat{\phi}_\Lambda - \hat{\phi}_{S_\Lambda})}$

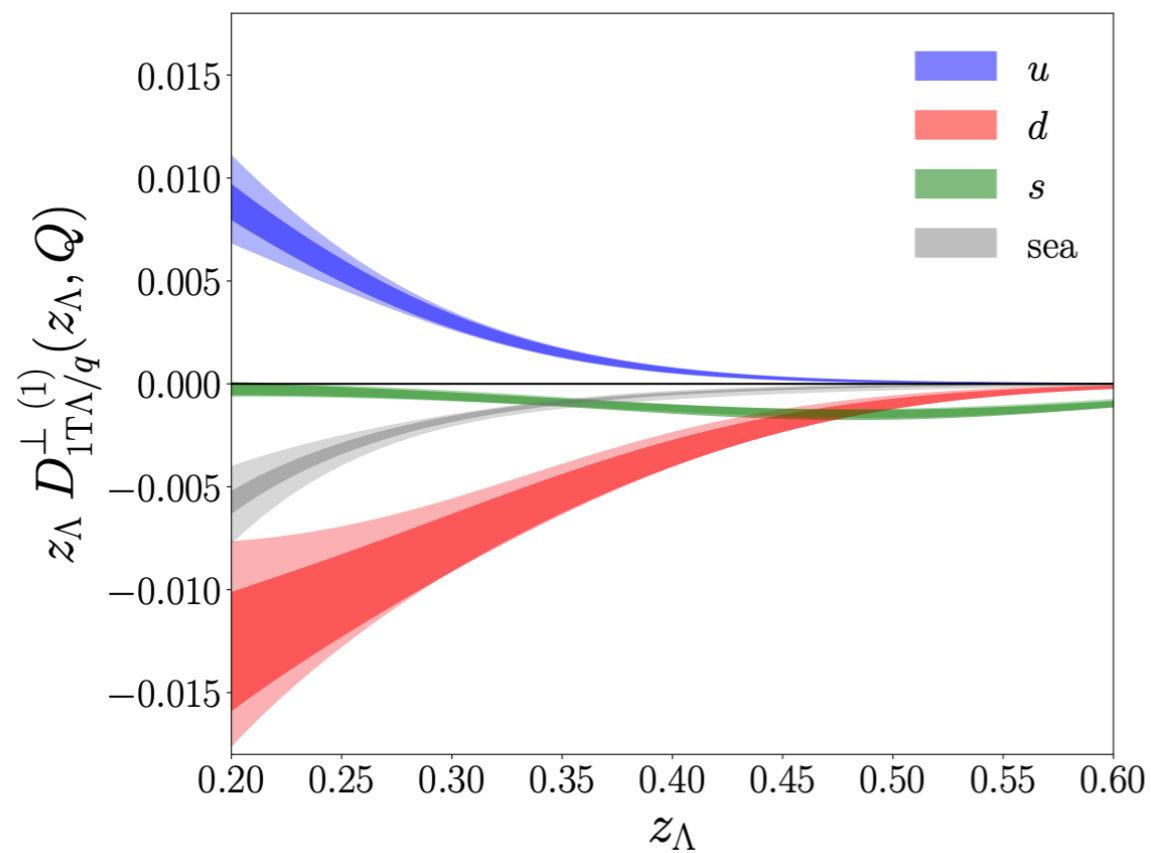
$$A_{UU,T}^{\sin(\hat{\phi}_\Lambda - \hat{\phi}_{S_\Lambda})} = \frac{F_{UU,T}^{\sin(\hat{\phi}_\Lambda - \hat{\phi}_{S_\Lambda})}}{F_{UU,U}} \sim \frac{f_1(q_\perp) D_{1T}^\perp(j_\perp)}{f_1(q_\perp) D_1(j_\perp)}$$

**EIC kinematics**



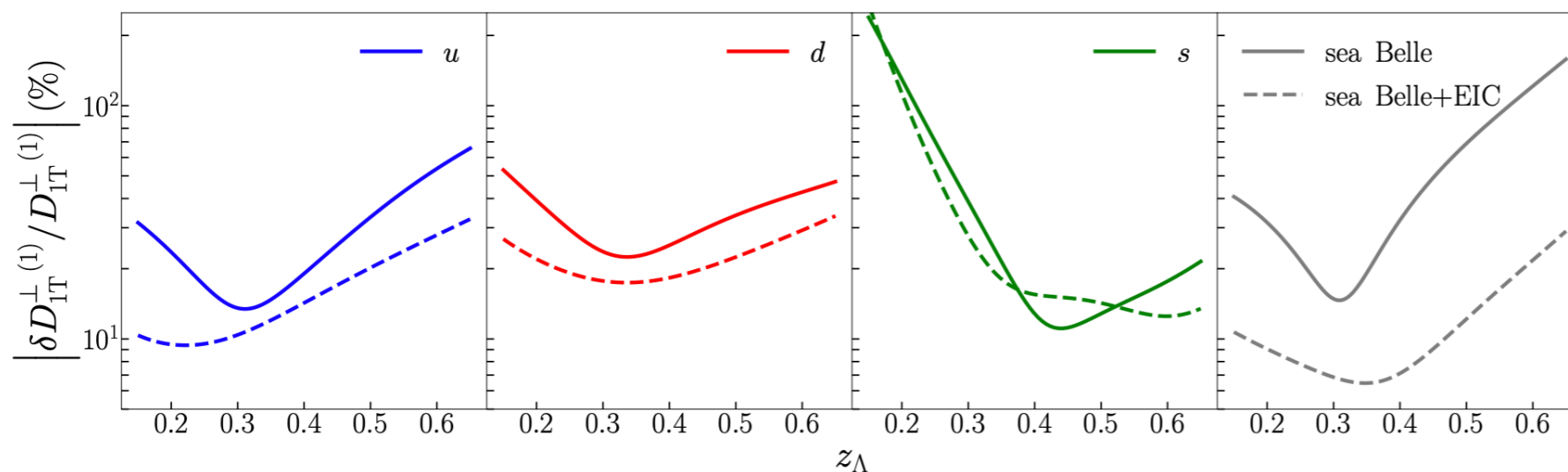
- $q_T$  dependence indeed cancels and is only sensitive to TMDFFs.

# Phenomenology : $A_{UU,T}^{\sin(\hat{\phi}_\Lambda - \hat{\phi}_{S_\Lambda})}$



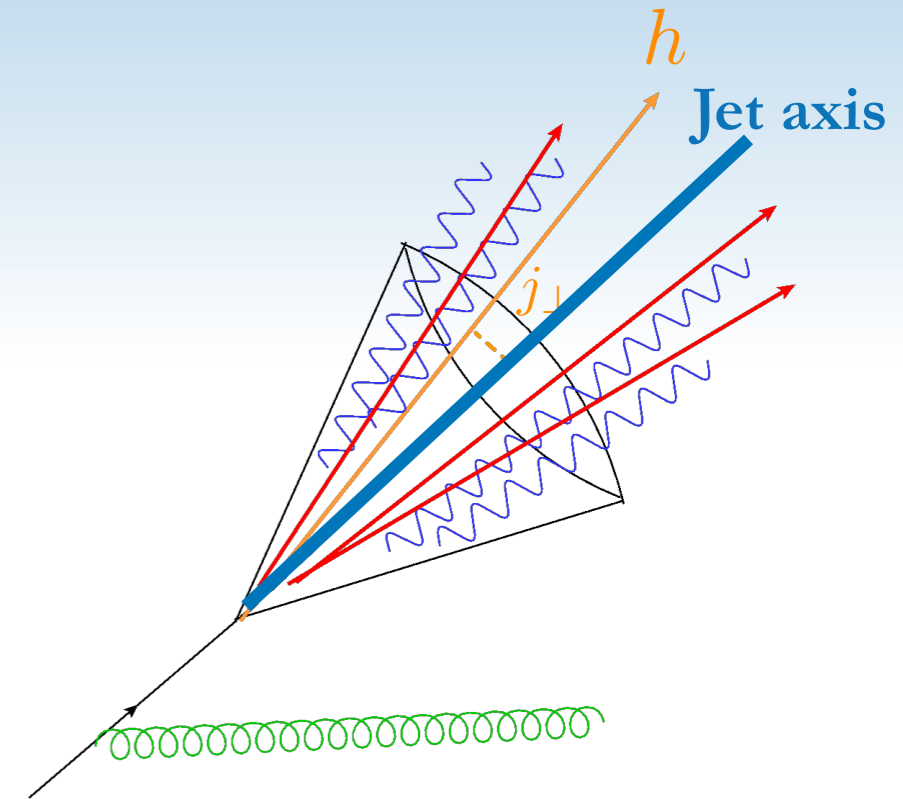
- **EIC pseudo-data significantly decreases the uncertainties in the determination of TMDPDFs.**

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# Conclusion



- New processes involving jets to non-perturbative structure

$$PP / eP \rightarrow J(h) + X, \quad eP \rightarrow e + J + X, \quad eP \rightarrow e + J(h) + X, \quad \dots$$

- Jet substructure techniques can be used to access information about **FFs**
  - Information differential in  $z_h$  allow more direct access to the FFs
  - Jet axis provides us a mean to access TMDFF structure
- Jet processes at the EIC can deconvolve the dependence between the **TMDPDF** and **TMDFF**.
  - Its high luminosity, wide energy range, and polarized beams will illuminate our understanding of the hadron structure and process of hadronization.
- Jets are great way to **'isolate'** and obtain **'differential information'** of the non-perturbative structure of interest!

**Thank you!**