Beam Monitor for FLASH

Tools and the Role of Radioluminescence and Cherenkov Emission. Front. Phys. 8:328. doi: 10.3389/fphy.2020.00328

Conventional vs Flash Beams

Parameters in Flash Beams

Review di detectors and Ashraf MR et al . doi: 10.3389/fphy.2020.00328

For FLASH, instantaneous dose and the dose per pulse quantities can be orders of magnitude higher than in conventional therapy, leading to issues of saturation, and non-linear response of standard dosimeters at large doses.

- 1. Risoluzione spaziale \sim 1 mm
- 2. Risoluzione temporale compatibile con il fascio (1 Hz di fascio, pulse di 1 µm) ~ns
- 3. Risoluzione sulla dose.. 3%?
- 4. Indipendente dall'energia del fascio
- 5. Indipendente dalla dose-rate istantanea
- 6. Sensibile alla quantità di ossigeno punto per punto

- **• Dose Rate Linearity**
- **• Spatial Resolution**
- **• Time Resolution**

From Ashraf MR et al . doi: 10.3389/fphy.2020.00328

Ashraf MR et al . doi: 10.3389/fphy.2020.00328 **6**

From Ashraf MR et al . doi: 10.3389/fphy.2020.00328

7

From Ashraf MR et al . doi: 10.3389/fphy.2020.00328

Fluorescence in Air F_{II} in A_{II} in A_{II} and F_{III} is proportional to the number of F deposited energy of the fluorescence efficiency can also be formulated using the decay lifetimes. The decay li
contract the decay lifetimes of the decay lifetimes of the decay lifetimes. The decay lifetimes of the decay li Here, *n^e* is the number of electrons and *dE* **IFILIO CORCANCA.** IN AIR THE NUMBER OF END Thesis: The Fluorescond molecules is proportional to the number of the number o

The fluorescence yield Y is the **number of photons produced by one electron per 1 m of path-length**. The unit of the fluorescence wield, [photons/m], has practical reasons and was introduced to use the fluorescence yield for air shower measurements. number of photons produced by one electron per 1 m of path-length. The unit of the fluorescence *r* ence yield Y is the **number of photons produced by one electron per 1:**
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as practical reasons and was introduced to use the fluorescence yiel *n µ a r a a i a a a i a a i a* \overline{c} ence *p*! rield for air shower measure $\frac{1}{2}$
is/m], has ractical rea ber or priotoris produced by one electron per in the pati-lengt
sons and was introduced to use the fluorescence yield for air showe

order to get The fluorescence efficiency Φ has to be defined. The fluorescence $\Phi_{\lambda} = \frac{1}{\Delta}$ *Given wavelength (with number of emitted photons n_λ):*
dependence spectrum, because the fluorescence spectrum, because the fluorescence spectrum, because the fluorescence spectrum, because the fluorescence spectrum, be In order to get a hold on the fluorescence yield, another quantity, and radiated energy the fluorescence efficiency Φ has to be defined. The fluorescence the fluorescence efficiency Φ has to be defined. The fluorescence $\Phi_{\lambda} = \frac{\text{reduced error}}{\text{deno-sided end}}$ given wavelength (with number of emitted photons nλ): ^λ = τ0*/*τ*F L*, at wavelength λ without quenching, and can then be combined with rold on the

with a mitted of emitted approximated approximate interest and interest in the formulated as:
Produce is proportional to the deposited energy, the fluorescence efficiency can also be formulated as: If the number of excited molecules is proportional to the deposited energy, the fluorescence efficiency can also be formulated as:

Equation (2.38)

 $\overline{}$

length. The unit of the unit of the fluorescence yield, \tilde{p}_{h} reasons and was introduced reasons and was introduced by the first correction that the first collaboration that needs to use the fluorescence yield of Air excited by Electron to use the fluorescence yield of Air excited by Electron to use the fluorescence yield of Air sure in the fluorescence of Physi
atmospheric gases January 2006 Czechoslovak Journal of Physi **divided a pHD Thesis: The Fluorescence Yield of Air excited by Electron** to the Bethe-Bloch equation along a path \sim in a gaseous measuremedium of density ρ **VU II ANI GEAL NOW BOOKS AND REPLACE CAN NOW BOOKS ARELY: Measurement of the fluorescence yield in** atmospheric gases January 2006 Czechoslovak Journal of Physics 56:A361-A367 **measured with the AIRFLY Experiment** deposited energy of the fluorescence energy of the fluorescence efficiency can also be formulated using the decay lifetimes. The decay lifetimes of the decay lifetimes of the decay lifetimes of the decay lifetimes. The dec

Φ^λ = radiated energy deposited energy = *n*^λ *· E*^λ *Edep ,* (2.38) *T*. The influence of the gas composition is also hidden in the reference pressure, since it depends on the total collisional cross section σ. The wavelength

explicitly mentioned, defined in units of photons/
m but some authors express it in photons/MeV: **d** io *Y* The fluorescence yield is always, unless explicitly mentioned, defined in units of photons/

s always, unless
in units of photons/
$$
Y\left[\frac{\text{photons}}{\text{m}}\right] = Y\left[\frac{\text{photons}}{\text{MeV}}\right] \cdot \frac{dE}{dX} \rho
$$
.

dE

 α deposite the number of electrons and dEoAx the deposited energy of one electron -. according to the Bethe-Bloch equation along a postument of the Bethe-Bloch equation along a postu density ρ. The fluorescence yield can be determined for the number of emitted photons Here, ne is the number of electrons and dEp∆x the deposited energy of one electron according to the Bethe-Bloch equation along a path of length ∆x in a gaseous medium of **nλ per electron and path length.** m $r_{\rm{max}}$ is the normal of cleatures and d Γ on Λ the depend in present in pressure. λ Indict, in this this indimition of disolations and depend in unposition of λ *Here*
- *acco dE density ρ*
nλ per el *a* to the Bethe-Bloch equation along a path of length *i*

| λ | System (ν_i, ν_f) | $A_{\nu_i, \, \nu_f}$ | λ | System (ν_i, ν_f) | |
|---------------------|-------------------------|-----------------------|---------------------|-------------------------|--|
| (nm) | | $(10^6 s^{-1})$ | (nm) | | |
| 310.40 | 2P(4,3) | 3.02 | 358.21 | 1N(1,0) | |
| 1311.67 | 2P(3,2) | 5.94 | 364.17 | 2P(4,6) | |
| $\frac{1}{3}$ 13.60 | 2P(2,1) | 10.1 | 367.19 | 2P(3,5) | |
| $\frac{1}{3}$ 15.93 | 2P(1,0) | 11.9 | $\frac{1}{3}$ 71.05 | 2P(2,4) | |
| 326.81 | 2P(4,4) | 3.71 | \dagger 375.54 | 2P(1,3) | |
| $\frac{1}{3}$ 28.53 | 2P(3,3) | 2.85 | $\frac{1}{380.49}$ | 2P(0,2) | |
| 329.34 | 1N(4,2) | 3.19 | 385.79 | 2P(4,7) | |
| 329.84 | 1N(3,1) | 2.08 | 385.79 | 1N(2,2) | |
| 330.80 | 1N(2,0) | 0.90 | 388.43 | 1N(1,1) | |
| $\frac{1}{3}30.90$ | 2P(2,2) | 0.80 | 389.46 | 2P(3,6) | |
| $\frac{1}{3}33.90$ | 2P(1,1) | 0.59 | $\frac{1}{3}$ 91.44 | 1N(0,0) | |
| 1337.13 | 2P(0,0) | 13.1 | \uparrow 394.30 | 2P(2,5) | |
| 344.60 | 2P(4,5) | 0.12 | \ddagger 399.84 | 2P(1,4) | |
| $\frac{1}{3}346.90$ | 2P(3,4) | 0.12 | $\frac{405.94}{ }$ | 2P(0,3) | |
| $\frac{1}{3}$ 50.05 | 2P(2,3) | 1.71 | 409.48 | 2P(4,8) | |
| 353.26 | 1N(5,4) | 6.63 | 414.18 | 2P(3,7) | |
| $\frac{1353.67}{ }$ | 2P(1,2) | 5.54 | 416.68 | 1N(3,4) | |
| 353.83 | 1N(4,3) | 7.46 | 419.91 | 1N(2,3) | |
| 354.89 | 1N(3,2) | 8.09 | 420.05 | 2P(2,6) | |
| 356.39 | 1N(2,1) | 7.88 | 423.65 | 1N(1,2) | |
| $\frac{1357.69}{ }$ | 2P(0,1) | 8.84 | 427.81 | 1N(0,1) | |

2.2. The Spectrum of Molecular Nitrogen 21 Details on band heads of N2 and N+2 between 300nm and 430nm

Eigure 2.0: Finorescence spectrum of an between 200 nm and 450 nm recorded by AIRFLY with transi-
Figure 2.0: The cas was evolved by 2 MoV electrons at a massure of 200 bDs. In the right upper sermer to Figure 2.8: Fluorescence spectrum of air between 280 nm and 430 nm recorded by AIRFLY with transition labels. The gas was excited by 3 MeV electrons at a pressure of 800 hPa. In the right upper corner, the spectrum reported by Bunner (1967) is shown [15].

Pick−Up coil

AIRFLY Exp. SetUP $\overline{}$ and $\overline{}$ are AIRFLY Experimental experimental

sensitivity at 420 nm. The bialkali catode has a quantum efficiency at 337 nm of about 8% and a very low dark current. - Hamamatsu H7195P model and was chosen for low background. It has a bialkali photocatode with a diameter of 46 mm and a peak

The interference filter, used to delimit the wavelength region of the light reaching the PMT, has a peak transmission of about 50% a nm and a width of 10 nm (FWHM). With this filter, the light reaching the PMT originates to 98.3% from the 2P(0,0)-transition at 337 nm. $\boxed{10}$ - The interference filter, used to delimit the wavelength region of the light reaching the PMT, has a peak transmission of about 50% at 340

Energy Dependence

https://www.researchgate.net/publication/215566660_AIRFLY_Measurement_of_the_fluorescence_yield_in_atmospheric_gases 11 F.Arciprete et al. AIRFLY: Measurement of the fluorescence yield in atmospheric gases January 2006 Czechoslovak Journal of Physics 56:A361-A367

Expected photons

Per elettroni da "13" MeV ci aspettiamo:

- 4 ph. al m di fluorescenza (@4 π)
- Cherenkov sotto soglia

- 4 ph. al m di fluorescenza (@4 π)
- 6 ph. al m di Cherenkov $({\sim} 0.1^{\circ})$

- 5 ph. al m di fluorescenza (@4 π)
- 70 ph. al m di Cherenkov $($ 1.4 \circ)

Per elettroni da 20 MeV ci aspettiamo :

Per elettroni da 130 MeV ci aspettiamo :

Expected photons

According to the detector we decide to implement.. SIT LIAC machine

(1)

Periscopio@SIT: 30.7.2020

Per elettroni da "13" MeV ci aspettiamo:

- 4 ph. al m di fluorescenza $(\omega 4\pi)$
- Cherenkov sotto soglia

According to the detector we decide to implement.. SIT LIAC machine

(1) :

- 0.04 ph x 10 cm x 0.1 x 10^10 (elet./pulse) \sim 4 10^8 ph
- $\bullet \sim 8$ 10^7 ph.el

PERISCOPIO

Charge

Preliminary analysis

Preliminary Conclusion with PERISCOPIO

- **We do see the beam IN and OUT difference in the detector.**
- **We observe some 'dependency' in the detector response as a function of the beam, probably we are mainly sensitive to the secondary products.**
- **It can be everything :)**
- **We are too close with the PMT.**
- **The "applicatore" produces background**

Expected photons

According to the detector we decide to implement.. SIT LIAC machine

- 0.04 ph x 10 cm x 0.1 x 10^10 (elet./pulse) \sim 4 10^8 ph
- ~ 8 10^7 ph.el
- ➡ **PMT troppo dentro l'alone del fascio**

➡ **L'applicatore produce troppo fondo che arriva ai PMT!** ➡ **Senza applicatore siamo in grado di vedere l'IN/OUT beam**

* riflessione on/off

• 0.04 ph/cm x 5 cm x 0.005 (sottostima/sovrastima geometria*) x 10^{10} (elet./pulse) ~ 1 10^7 ph

• \sim 2 10^6 ph.el

MC FLUKA SIMULATION

By Gaia, Antonio (and Battistoni for support)

Energy of particles reaching the PMT

MC FLUKA SIMULATION

- **• # Electrons: 107**
- **• Energy: 13 MeV**

MC FLUKA SIMULATION: Optic photons card activation

Released energy in the Air target

Optical Photons Fluence - FLUKA generated [# ph./(cm2 primary)]

The wallets are reflective at 90% (Teflon sheet from Thorelabs)

MC FLUKA SIMULATION: Optic photons card activation

The wallets are reflective at 90% (Teflon sheet from Thorelabs). Black configuration decreases the number of photons of a factor 100.

MC FLUKA SIMULATION: Optic photons card activation

- ๏ The LIAC HWL shoots 1010 electrons with energy 6 MeV in a pulse 4µs long at a frequency of 10Hz.
- ๏ The first run was performed with the beam shooting straight at the center of the dosemeters (ref. up, abs. down).

- ๏ **1st question**: can we observe in-beam/off-beam difference?
	-

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- ๏ **2rd question**: What if one of the PMTs is **farther** from the aperture?
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- ๏ **3rd question**: What difference do **reflective** and **absorber** make?
	-

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-
- ➤ **Note (2):** In real word only two faces of

- ➤ **We can see the IN beam from the OFF beam**
- ➤ **The response of the detector is sensitive to position and geometry**
- **probably a not-reflective detector is a better choise .**
- ➤ **Filters study => to do..**

➤ **The study reflectivity/absorption with MC and data is coherent with expectancy..**

Preliminary Conclusion with TUBONE

geometry is a better choice . \Box Filters study => to do..

M We can see the IN beam from the OFF beam

M The response of the detector is sensitive to position and

M The study reflectivity/absorption with MC and data is coherent with expectancy.. probably a not-reflective detector

Expected photons

According to the detector we decide to implement.. SIT LIAC machine

- 0.04 ph x 10 cm x 0.1 x 10^10 (elet./pulse) \sim 4 10^8 ph
- ~ 8 10^7 ph.el
- ➡ **PMT troppo dentro l'alone del fascio**

• 0.04 ph/cm x 5 cm x 0.005 (sottostima/sovrastima geometria*) x 10^{10} (elet./pulse) ~ 1 10^7 ph

• \sim 2 10^6 ph.el

(1):

• 0.04 ph/cm x 2 cm x 0.003 x 10^{12*} (elet./pulse) ~ 2 10^8 ph $\bullet \sim 510^{17}$ ph.el ➡**Senza applicatore PMT: H6524** *FLASH machine in 100 times more intense

➡ **L'applicatore produce troppo fondo che arriva ai PMT** ➡ **Senza applicatore siamo stati in grado di vedere l'IN/OUT beam**

* riflessione on/off

PMT: H10580

Tubino@SIT: 05.07.2021 (dal futuro con < 3)

- ๏ We plan to take the FLASH beam..
- ๏ The inside is not reflective. Less saturation effects, less geometry dependence.

Vediamo se questa mano vi piace quanto l'altra..

Tubino@SIT: 05.07.2021 (dal futuro con < 3)

- ๏ We plan to take the FLASH beam..
- ๏ The inside is not reflective. Less saturation effects, less geometry dependence.
- ๏ For the moment we just equalised the PMT response..
- ๏ "New".. less noice and more stable
- ๏ Test with UV Filters..

SUPPLY VOLTAGE (V)

GAIN

Tubino@SIT: 05.07.2021 (dal futuro con < 3)

- ๏ We plan to take the FLASH beam..
- ๏ The inside is not reflective. Less saturation effects, less geometry dependence.
- ๏ For the moment we just equalised the PMT response..
- ๏ "New".. less noice and more stable
- ๏ Test with UV Filters..
- ๏ We plan to change the current of the beam and see if we can follow the

Matteo Pacitti:

A parità di lunghezza dell'impulso possiamo variare la corrente di fascio da 73 mA fino a 130 mA. Considerando un impulso rettangolare di durata 4 us, quindi approssimando leggermente, andiamo da 1.8 x 1012 elettroni a 2.9 x 1012 elettroni.

Future prospective..

- ๏ I definitely find very 'romantic' that an empty detector can produce light..
- ๏ It is certainly true that this light can be correlated with the electron beam (on/off, shift, …)
- ๏ We have to make the step further and go from a very 'qualitative measure' a a raw 'quantitative measure'.. and than decide what to do..
- ๏ The new geometry should help.. but it is not the best choice ever..
	- ๏ => charge average study
	- \bullet \Rightarrow wavelength study
	- ๏ => geometry optimisation

We still need to prove that we can correlate with the number of electrons.. !!

