

# Physics opportunities with an “energy upgrade” of DaΦne

[flexible c.o.m energy to run both below and above the  $\Phi$  peak]

Gino Isidori

[ *INFN - Frascati* ]

- ▶ Introduction
- ▶ The hadronic contribution to  $(g-2)_\mu$
- ▶ The hadronic contribution to  $\alpha_{\text{em}}$
- ▶ Spectroscopy and  $\gamma\gamma$  physics
- ▶ Exotics
- ▶ Conclusions

## ► Introduction

The *high-intensity frontier* in the LHC era:

- Study of rare & forbidden processes

Full complementarity with the LHC for the identification of nature of physics beyond the SM

- Precision measurements of fundamental SM couplings

Several SM parameters, which are likely to play a fundamental role in the identification of the underlying theory, can only be measured at low energies.

- Deeper studies of QCD

There are still interesting aspects of non-perturbative QCD which are not fully understood and need to be investigated further.

## ► Introduction

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- Study of rare & forbidden processes [  $(g-2)_\mu, \dots$  ]

Full complementarity with the LHC for the identification of nature of physics beyond the SM

- Precision measurements of fundamental SM couplings [  $\alpha_{\text{em}}, \dots$  ]

Several SM parameters, which are likely to play a fundamental role in the identification of the underlying theory, can only be measured at low energies.

- Deeper studies of QCD [  $\gamma\gamma \rightarrow \pi, \eta, \sigma, \dots$  ]

There are still interesting aspects of non-perturbative QCD which are not fully understood and need to be investigated further.

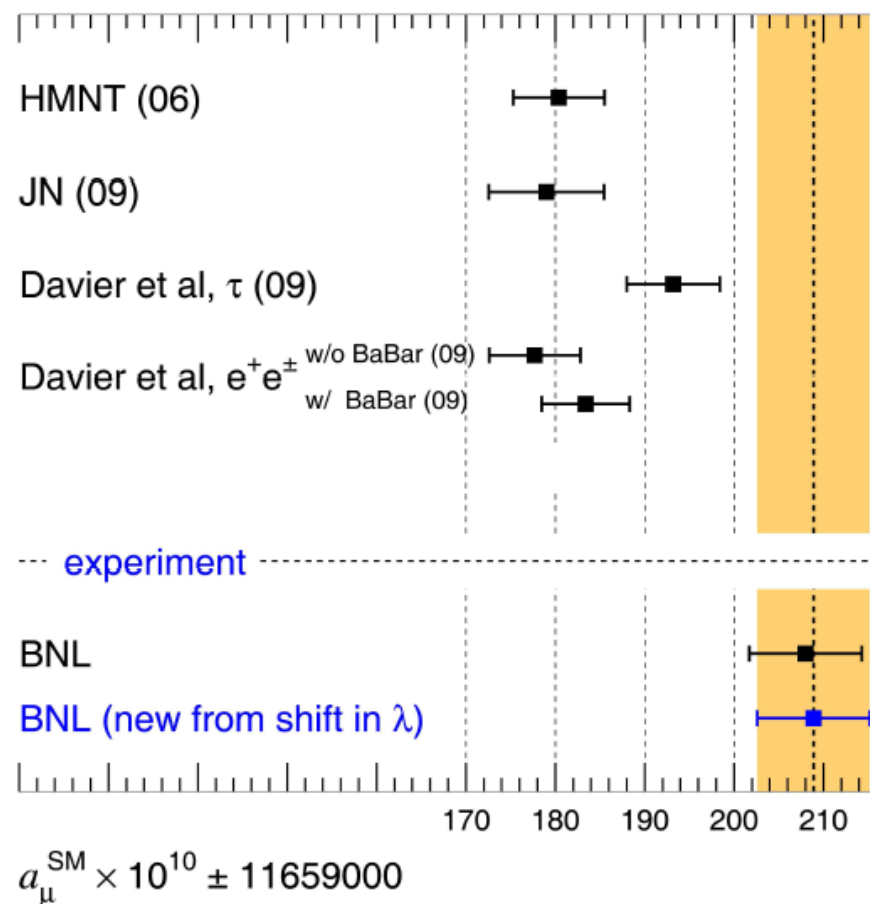
A DAΦNE upgrade with flexible c.o.m. energy [  $\sim (0.5-2.5)$  GeV ]  
could allow to perform significant measurements in all these 3 categories

## ► The hadronic contribution to $(g-2)_\mu$

The anomalous magnetic moment of the muon is one of the most (if not *the most*) significant low-energy constraint on a wide class of new-physics models

$$\Delta a_\mu^{\text{exp}} = a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (25 \pm 8)10^{-10}$$

- In the last few years the result for  $a_\mu^{\text{SM}}$  has become more reliable and the size of the discrepancy has (slightly) increased
- The discrepancy is large compared to  $a_\mu^{\text{ew-SM}}$  ( $a_\mu^{\text{ew-SM}} \sim 15 \times 10^{-10}$ )



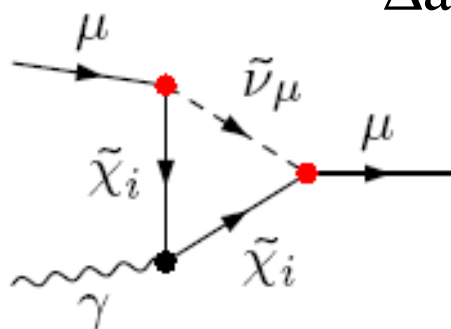
► The hadronic contribution to  $(g-2)_\mu$

The anomalous magnetic moment of the muon is one of the most (if not *the most*) significant low-energy constraint on a wide class of new-physics models. Its present  $\sim 3\sigma$  discrepancy from the SM expectation could naturally be accommodated in well motivated new-physics models, such as the MSSM

Unique example of helicity suppressed observable sensitive to electroweak physics

→ direct probe of **Yukawa interactions**

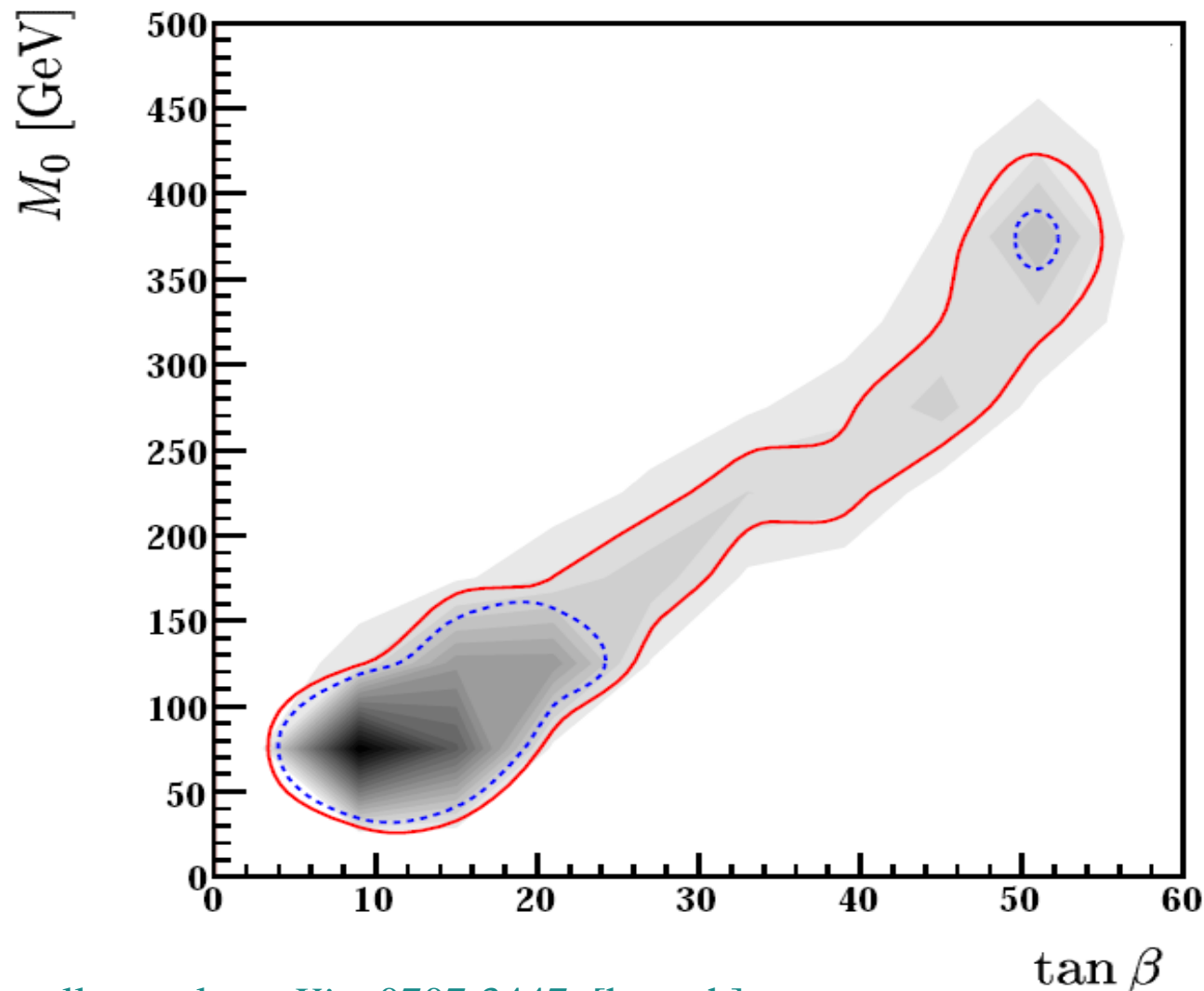
$$\mathcal{L}_{\text{eff}} = \frac{c_{\text{eff}}}{\Lambda^2} \bar{\mu}_R \sigma^{\mu\nu} \mu_L \phi F_{\mu\nu}$$

$$\Delta a_\mu^{\text{SUSY}} \sim \tan\beta \times (m_W/M_{\text{SUSY}})^2 \times (a_\mu^{\text{ew-SM}}) \times \text{sgn}(\mu)$$


← Couplings proportional to the muon Yukawa and not to its mass !!

► The hadronic contribution to  $(g-2)_\mu$

E.g.: Role of indirect constraints in a global fit of the CMSSM:

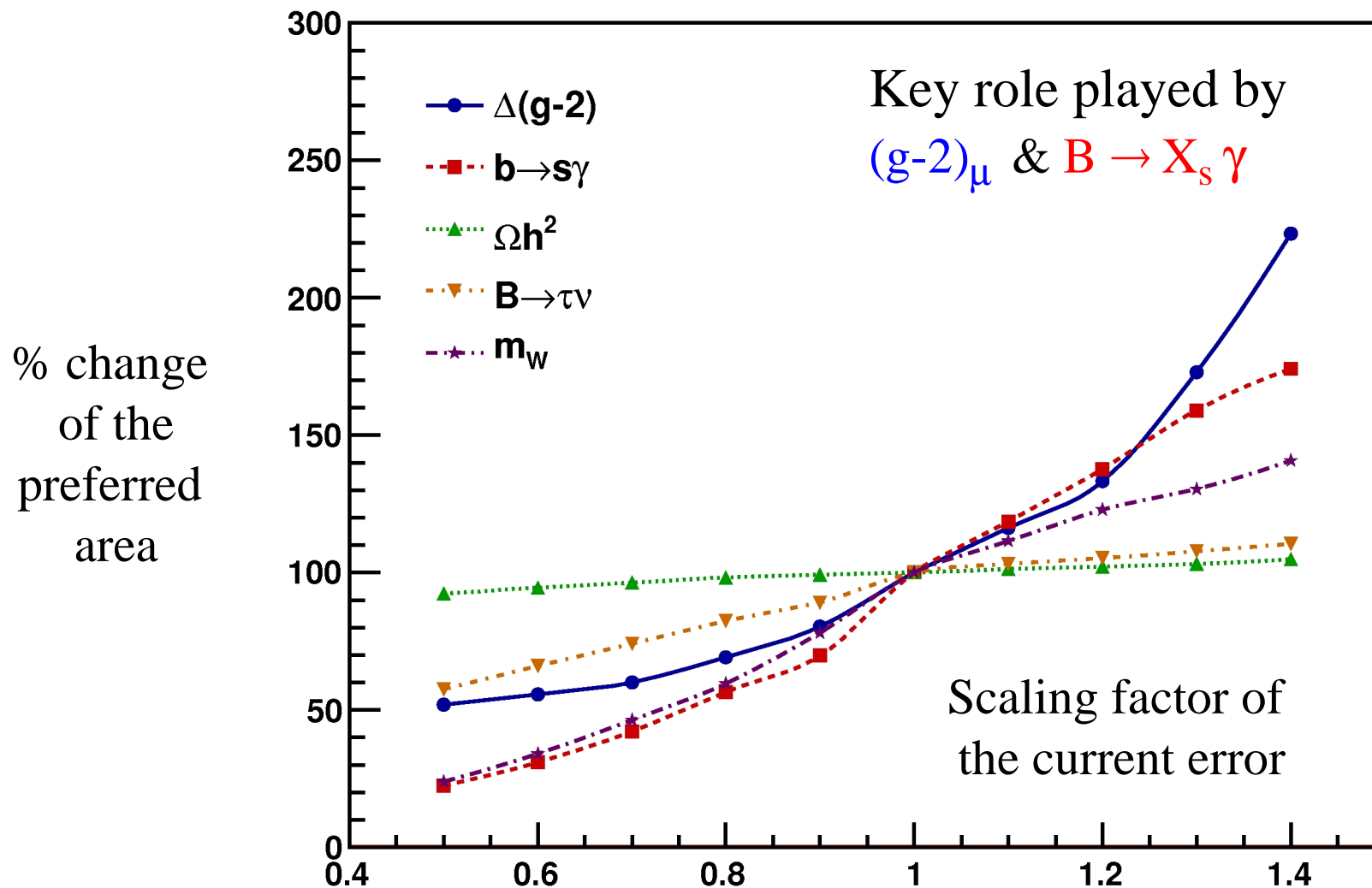


Key role played by  
 $(g-2)_\mu$  &  $B \rightarrow X_s \gamma$

**N.B.:** at LHC we hope  
to measure well  $M_0$   
(and other mass param.)  
but we don't expect a  
good sensitivity to  $\tan\beta$

► The hadronic contribution to  $(g-2)_\mu$

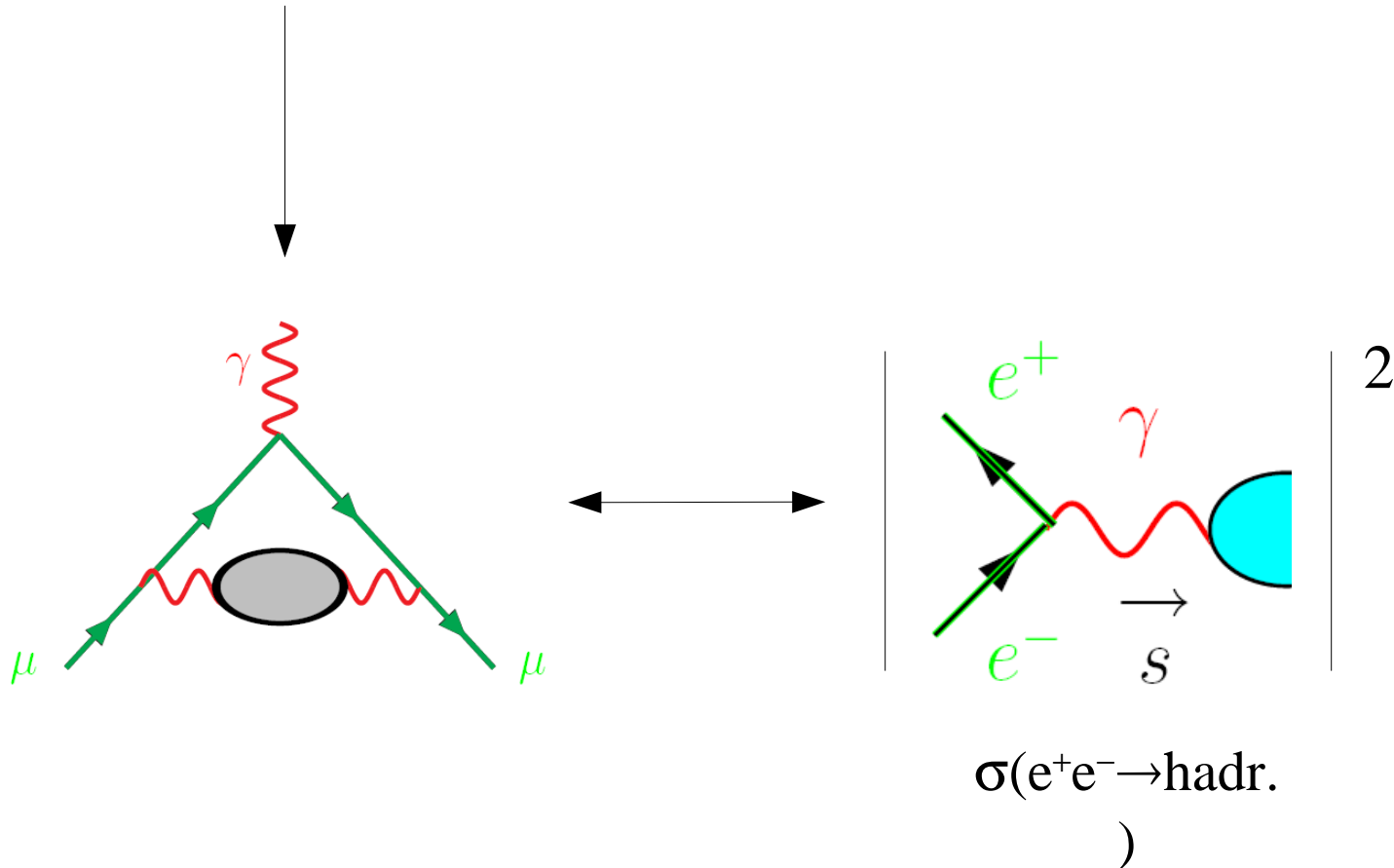
E.g.: Role of indirect constraints in a global fit of the CMSSM:



► The hadronic contribution to  $(g-2)_\mu$

$$\Delta a_\mu^{\text{exp}} = a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = \left[ \begin{array}{ll} (27.7 \pm 8.4) \cdot 10^{-10} & \text{Eidelman TAU08} \\ (24.6 \pm 8.0) \cdot 10^{-10} & \text{Davier et al. arXiv: 0908.4128} \end{array} \right]$$

$$8.4 \approx 5_{\text{HLO}} \oplus 3_{\text{LBL}} \oplus 6_{\text{BNL}}$$





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$$4 \approx 2.5 \oplus 2.5_{\text{LBL}} \oplus 1.6$$

new  
 $e^+e^- \rightarrow \text{hadr.}$

new  $g-2$

A substantial reduction of the present error is not a dream !

Venanzoni, MCW '09

$$\delta a_\mu^{\text{HLO}} = 5.29 = 3.0 (\sqrt{s} < 1 \text{ GeV}) \oplus 3.9 (1 < \sqrt{s} < 2 \text{ GeV})$$

$$\delta a_\mu^{\text{HLO}} \rightarrow 3 = 2.5 (\sqrt{s} < 1 \text{ GeV}) \oplus 1.5 (1 < \sqrt{s} < 2 \text{ GeV})$$

This means:

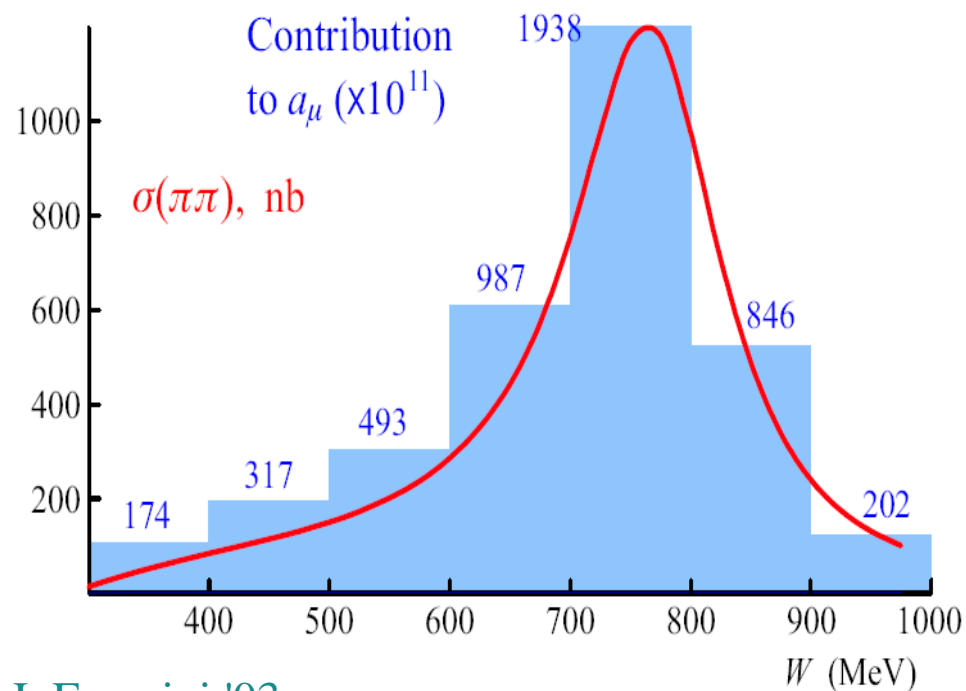
$$\delta \sigma_{\text{HAD}} \sim 0.4\% \quad \sqrt{s} < 1 \text{ GeV} \quad (\text{instead of } 0.7\% \text{ as now})$$

$$\delta \sigma_{\text{HAD}} \sim 2\% \quad 1 < \sqrt{s} < 2 \text{ GeV} \quad (\text{instead of } 6\% \text{ as now})$$

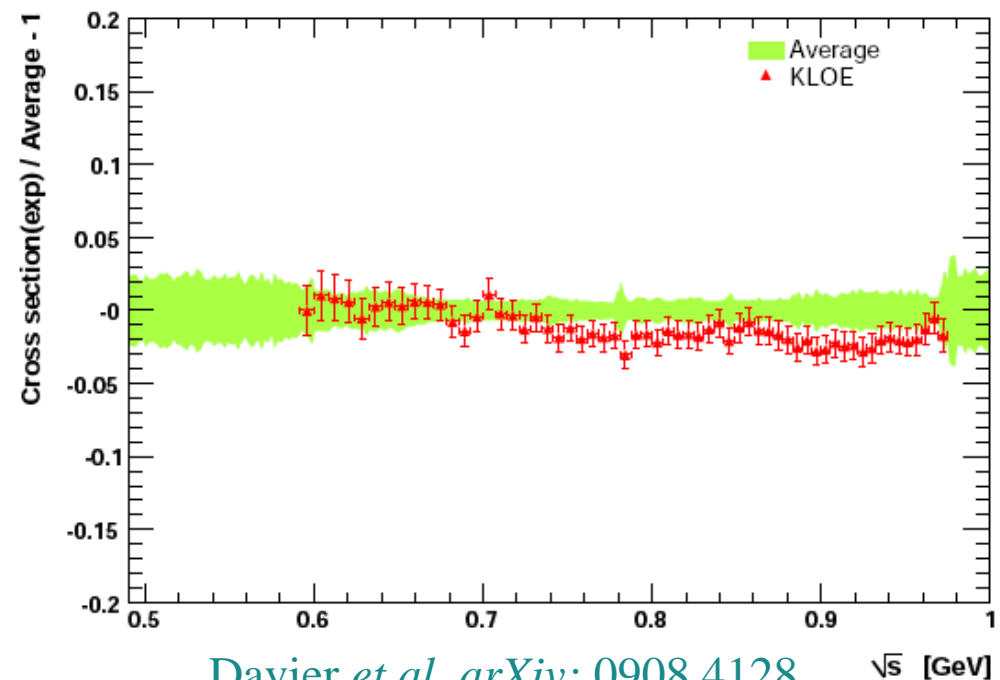
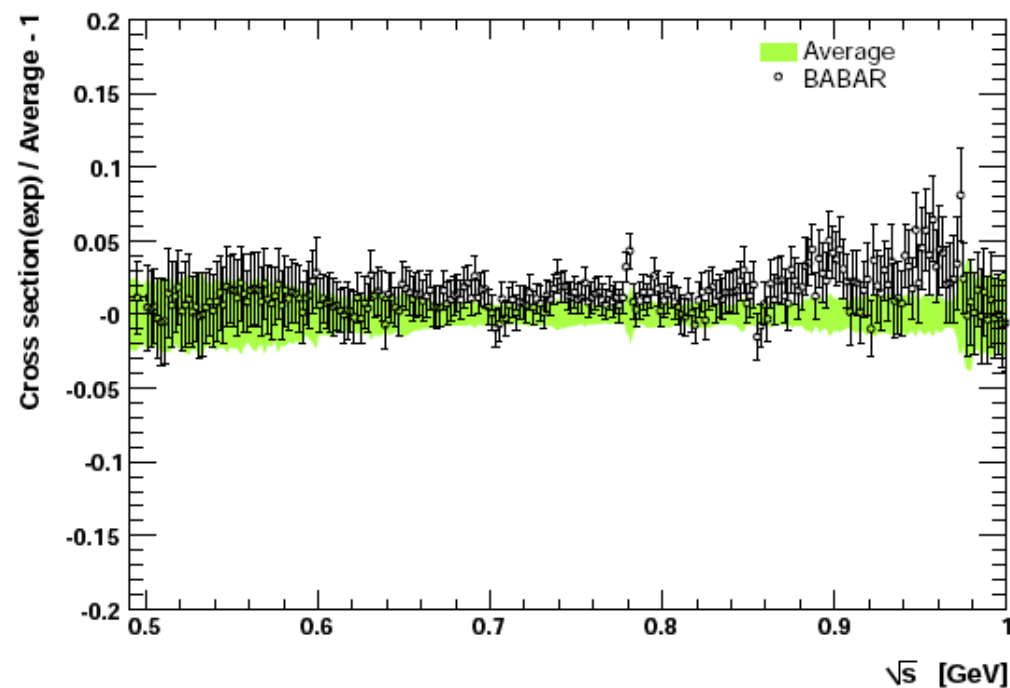
But we need better data on  $\sigma(e^+e^- \rightarrow \text{hadr.})$  both below and above 1 GeV

The improvement in the 1-2 GeV region clearly requires an energy upgrade [useful also for  $\alpha_{\text{em}}$ ]

But an energy scan in the 600-900 MeV region is also the only way to get a *convincing* error ( $\delta\sigma < 0.5\%$ ) in dominant low-energy region



J. Franzini '03

Davier *et al.* arXiv: 0908.4128

► The hadronic contribution to  $\alpha_{\text{em}}$

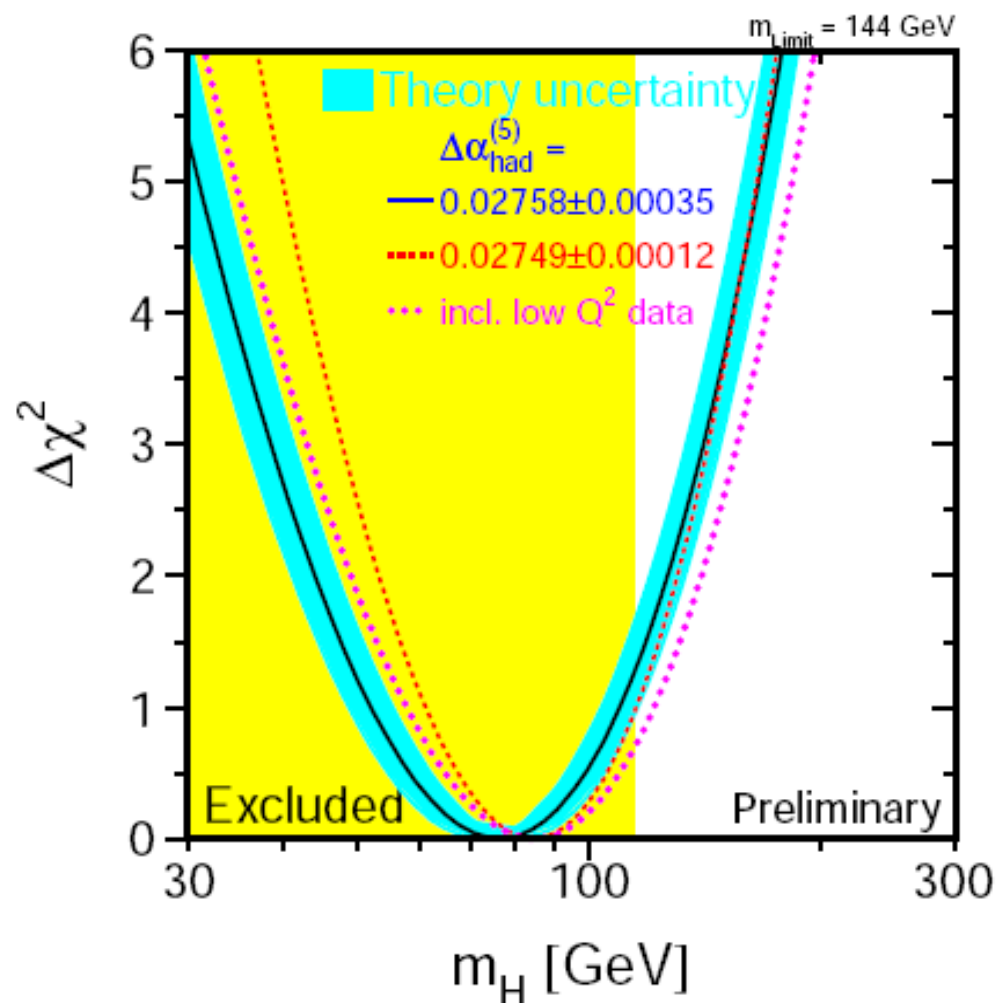
While  $\alpha_{\text{em}}(m_e)$  is known with an incredible precision [  $\sim 3 \times 10^{-9}$  !], the error on  $\alpha_{\text{em}}(M_Z)$  -the effective coupling relevant at the electroweak scale- is much larger because of hadronic uncertainties:

$$\frac{\Delta\alpha_{\text{em}}(M_Z)}{\alpha_{\text{em}}(M_Z)} \sim (1-4) \times 10^{-4}$$

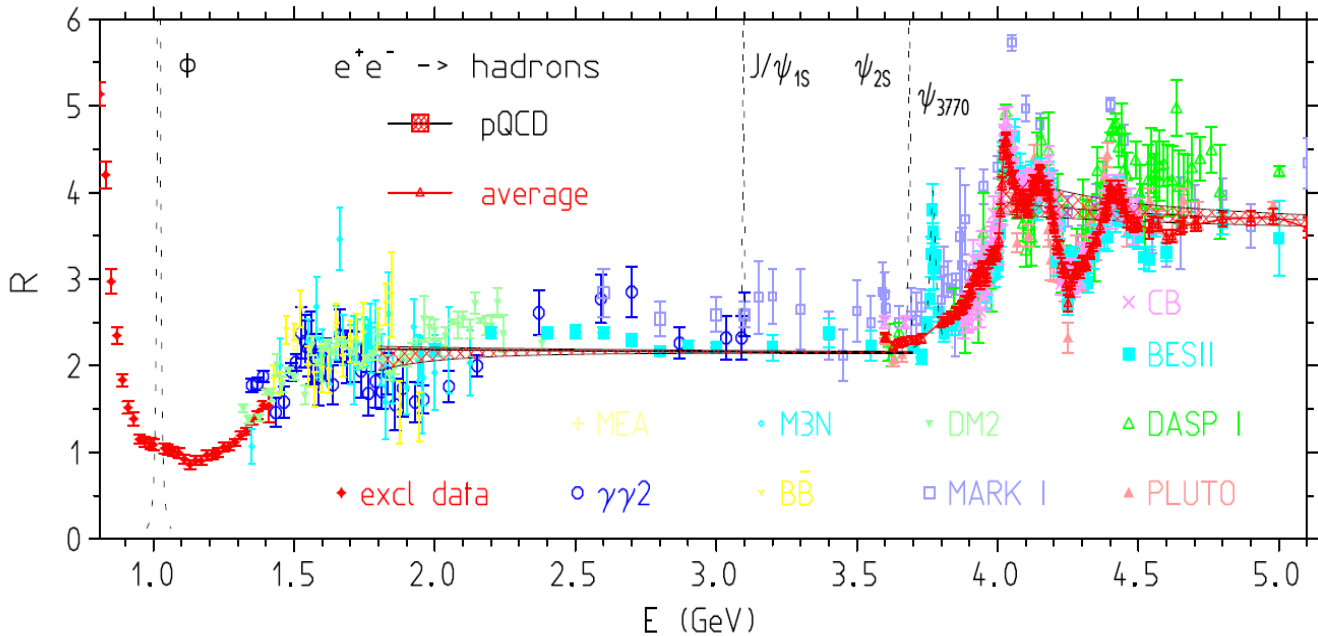
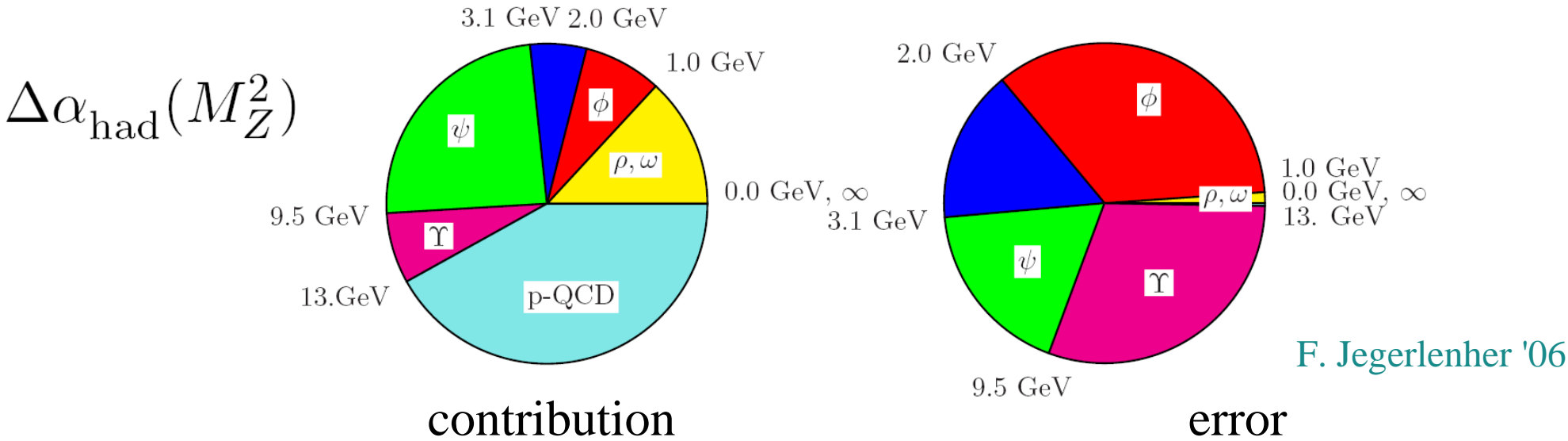
Significant source of uncertainty  
(given the precision reached on  
other e.w. fundamental couplings)

$$\frac{\delta G_\mu}{G_\mu} \sim 8.6 \times 10^{-6}$$

$$\frac{\delta M_Z}{M_Z} \sim 2.4 \times 10^{-5}$$



► The hadronic contribution to  $\alpha_{em}$

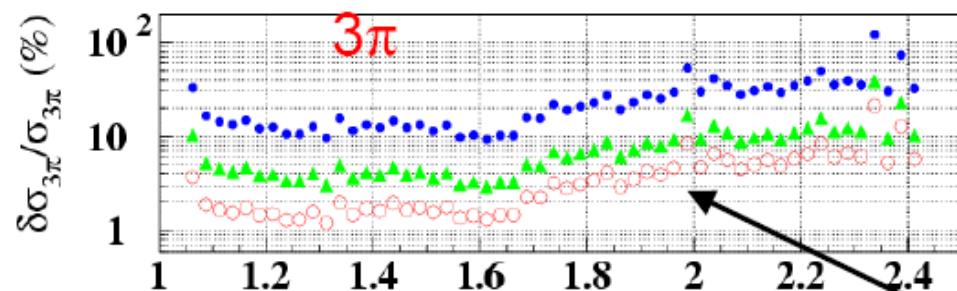


The high-energy range is not an issue: we can reduce the error using theory

The key problem is the  $E < 2.5$  GeV region

## ► The hadronic contribution to $\alpha_{\text{em}}$

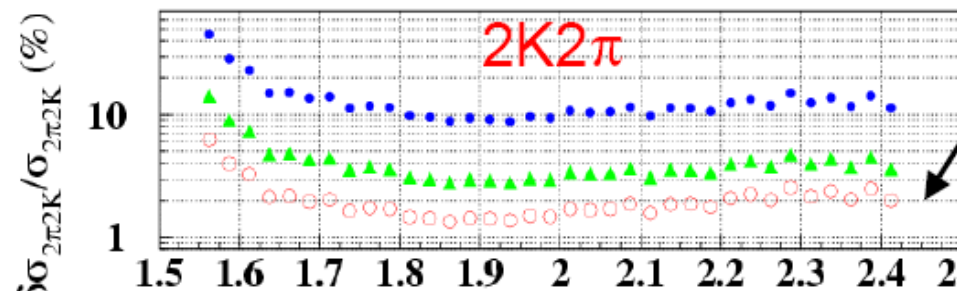
Improving  $\sigma(e^+e^- \rightarrow \text{had.})$  in the 1-2 GeV region is not an easy task (several channels), but an energy scan with KLOE (possibly with inner tracker) seems to be the best tool to achieve this important goal:



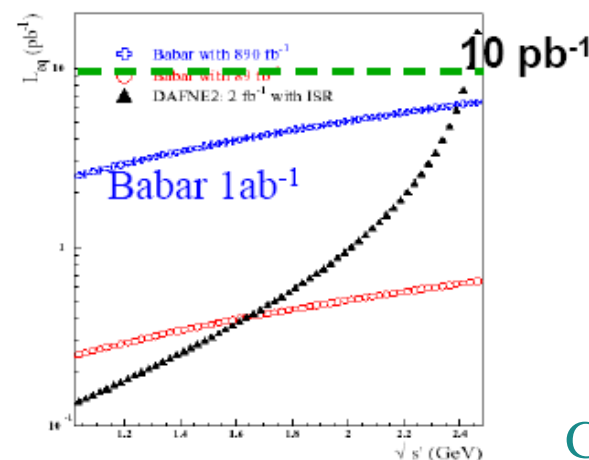
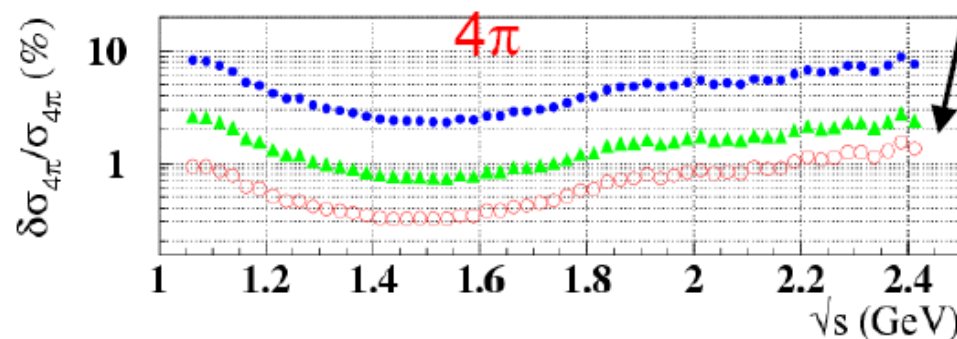
BaBar, with the published  $L_{\text{int}}$  per point ( $90 \text{ fb}^{-1}$ )

BaBar, with  $10 \times$  (the present  $L_{\text{int}}$ )

DAFNE-2, with  $20 \text{ pb}^{-1}$  per point [20 MeV bins]



- DAFNE-2 is **statistically** better than  $O(1 \text{ ab}^{-1})$  B-factories
- Improvement on systematics come as well

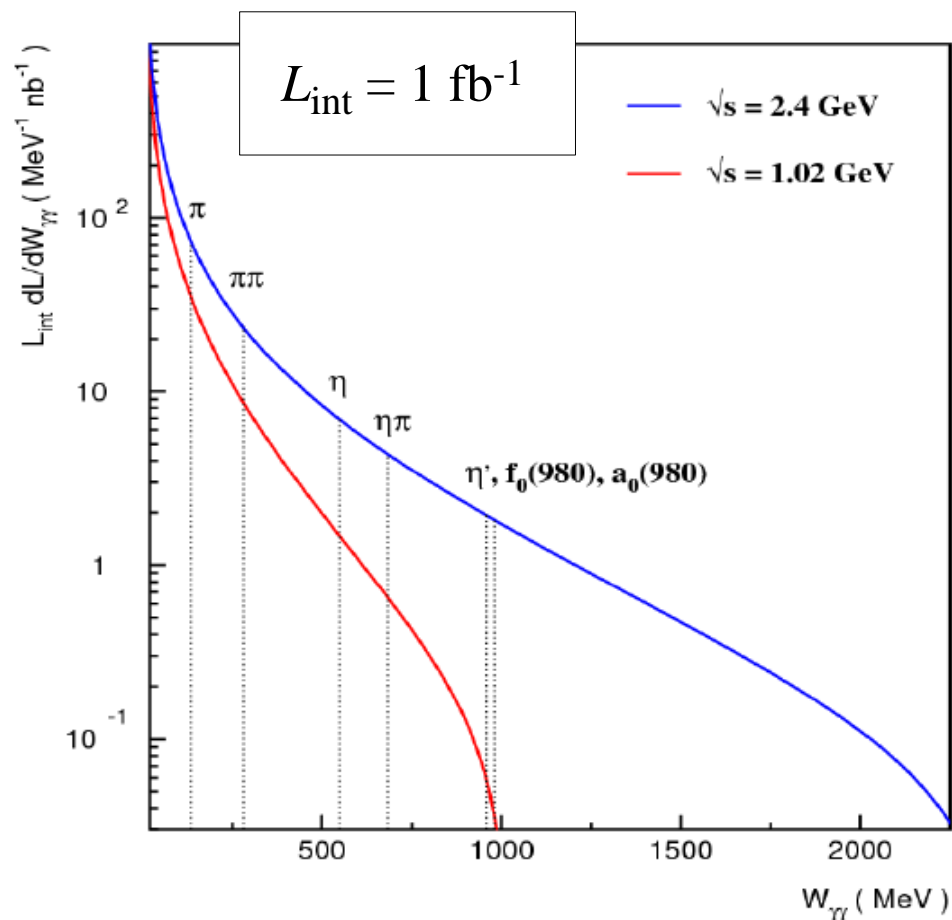


## ► Spectroscopy and $\gamma\gamma$ physics

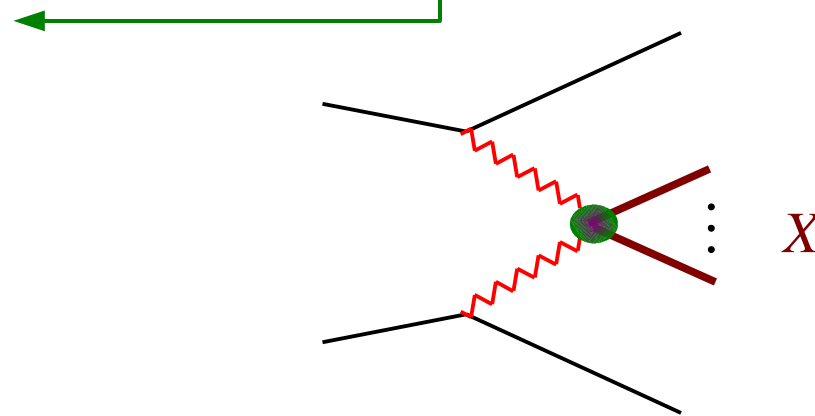
There are still several interesting open issues in low-energy spectroscopy (4q states, glueballs, etc...) which would strongly benefit from high-precision measurements of  $e^+e^- \rightarrow X$  and  $e^+e^- \rightarrow X + e^+e^-$  at  $\sqrt{s}$  above 1 GeV

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$$\frac{dN_X}{dW_{\gamma\gamma}} = L_{int} \frac{dL}{dW_{\gamma\gamma}} \sigma(\gamma\gamma \rightarrow X)$$

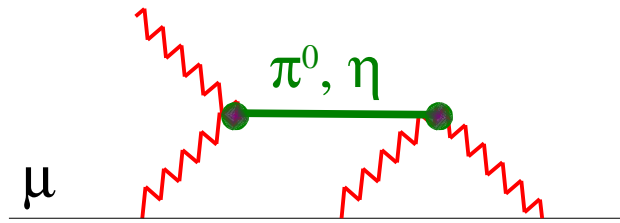


The high-energy requirement is particularly relevant for  $\gamma\gamma$  physics

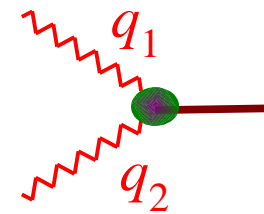
## ► Spectroscopy and $\gamma\gamma$ physics

One specific aspect worth to be recalled: the possibility to improve the error on the **LBL** contribution to  $(g-2)_\mu$

Largely dominant LBL contrib.:



Uncertainty due to our ignorance about  $\pi(\eta)\gamma^*\gamma^*$  form factors  $F(q_1^2, q_2^2)$

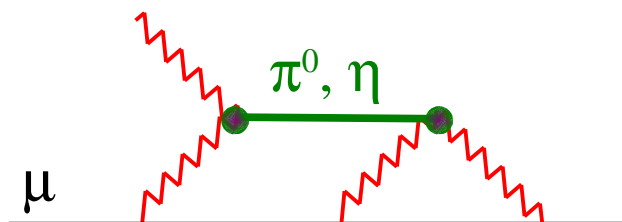




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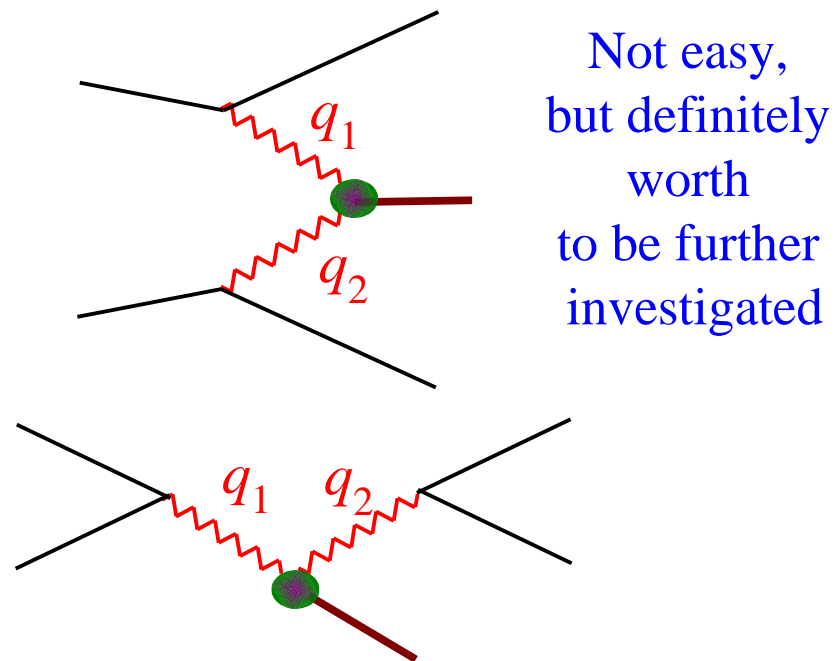


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Some infos on these f.f. are obtained from  
 -radiative Dalitz decays [  $dF(0, q^2)/dq^2$  ]  
 -leptonic decays [ **weighted integral of the f.f.** ]

What we really miss is  $d^2 F(q_1^2, q_2^2)/dq_1^2 dq_2^2$   
 which could possibly extracted from

$\gamma^*\gamma^* \rightarrow \mathbf{P}$  at large angle and/or  $\gamma^* \rightarrow \mathbf{P} \gamma^*$ ,  
 both at large energy [Bijnens, Persson, '01]



Not easy,  
 but definitely  
 worth  
 to be further  
 investigated

## ► Exotics

Recent astrophysical puzzles [**Pamela**, **Fermi**, ...] and the long-standing **Dark Matter** problem, have motivated the construction of various exotic extensions of the SM model characterized by

- new light states [**vectors and scalars with masses in the 0.1-1 GeV range**]
- weakly coupled to photons, or directly coupled to muons and/or electrons



Precision differential measurements of  $e^+e^- \rightarrow e^+e^- + \gamma$  and/or  $e^+e^- \rightarrow \mu^+\mu^- + \gamma$  and/or  $e^+e^- \rightarrow E_{\text{miss}} + \gamma$ , ... at low energies are the best way to constrain (or find evidences...) of such models [ **large parameter space is still unexplored !** ]  
F. Bossi

Probably a lot can be done already at  $\sqrt{s} = 1$  GeV; however, going to lower energies could possibly help [**further investigation needed**]

## ► Conclusions

There is a rich physics program both below and above 1 GeV:

- The low-energy program (0.6–1.4 GeV) is really unique (no competition from other facilities), quite interesting, and apparently not too expensive  $\Rightarrow$  worth to consider it even in parallel to SuperB (shorter time scale, already part of the KLOE2-LOI), provided there is no competition for the site
- The high-energy option (1.4–2.5 GeV) is more expensive and would not make sense if the SuperB project will go on. However, it is certainly an interesting mid-term low-cost alternative for the Lab, if the SuperB project will not be approved (if we want to have it ready as a relativistic “plan-B”, it need to be further investigated now: the project is particularly interesting if it can be realised soon)

## ► Conclusions

Final (partially-related) comment:

- If the SuperB will be approved, it would be wise to consider the possibility of a machine with “flexible c.o.m. energy” (*is it feasible without compromising peak luminosity, and with a “fast” energy-tune?*) able to run down to lower energies (e.g. from the  $\Upsilon$  to the  $\psi$ ):
  - On the physics side, this could allow to cover part of the program I just discussed ( + much more ...)
  - On a more strategical point of view, this could allow to avoid a “duplicate” of the Japanese machine