# |V<sub>ud</sub>| & |V<sub>us</sub>|determinationfrom kaon decays

### Paolo Massarotti

INFN Naple -Naples University "Federico II", Flavour Physics and CP Violation 2010 Turin, May 26 2010



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negligeble (10<sup>-3</sup>) see talk by Robert Kowalewski

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negligeble (10<sup>-3</sup>) see talk by Robert Kowalewski

# Need a precise determination of |Vus|

# Vud determination

- O Best result: from superallowed 0<sup>+</sup>→0<sup>+</sup> nuclear transitions.
   (comprehensive review [Towner & Hardy Rep. Prog. Phys. 73 (2010) 046301])
- Master formula  $\mathcal{F} = \frac{K}{2G^2_F |V_{ud}|^2 (1 + \Delta_R)}$



3

Vud determination



From most recent neutron  $\beta$  decay result: 0.9758(13) From pion  $\beta$  decay (PDG08):0.9742(26)

# The FlaviaNet Kaon working group

# •The most precise measurement of |Vus| is obtained from the charged and neutral kaon channel

•FlaviaNet Kaon WG (www.lnf.infn.it/wg/vus/). Recent kaon physics results come from many experimental (BNL-E869, KLOE, KTeV, ISTRA+, NA48) and theoretical (Lattice,  $\chi_{PT}$ ,) improvements. The main purpose of this working group is to perform precision tests of the Standard Model and to determine with high accuracy fundamental couplings (such as  $V_{us}$ ) using only published data on kaon decays, taking correlations into account.

# Vus determination

Physics results:

•  $|\mathbf{V}_{\mathrm{us}}| \times f_{+}(0)$ 

$$\begin{split} \Gamma_{K_{\ell 3}} &= \frac{G_F^2 m_K^5}{192\pi^3} C_K^2 S_{\rm EW} \left( |V_{us}| f_+^{K^0 \pi^-}(0) \right)^2 I_{K\ell} \left( 1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi} \right)^2 \\ & |V_{us}| / |V_{ud}| \times f_K / f_{\pi}. \end{split}$$

$$\begin{split} \frac{\Gamma_{K_{\ell 2}}}{\Gamma_{\pi_{\ell 2}}} &= \frac{|V_{us}|^2}{|V_{ud}|^2} \frac{f_K^2}{f_{\pi}^2} \frac{m_K (1 - m_{\ell}^2 / m_K^2)^2}{m_{\pi} (1 - m_{\ell}^2 / m_{\pi}^2)^2} \left( 1 + \delta_{\rm EM} \right) \end{split}$$

Global fits and averages:

- $K_L$ ,  $K_S$ , and  $K^{\pm}$ , dominant BRs and lifetime.
- Parameterization of the  $K \rightarrow \pi$  interaction (form factor)

# Vus determination

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$$\begin{split} \frac{\Gamma_{K_{\ell 2}}}{I_{K_{\ell 2}}} &= \frac{|V_{us}|^2}{I_K^2} \frac{f_K^2}{m_K (1 - m_\ell^2 / m_K^2)^2} \left( 1 + \delta_{\rm EM} \right) \end{split}$$

$$\frac{1}{\Gamma_{\pi_{\ell^2}}} = \frac{1}{|V_{ud}|^2} \frac{r_{R}}{f_{\pi}^2} \frac{1}{m_{\pi}(1 - m_{\ell}^2/m_{\pi}^2)^2} (1 + \delta_{\rm EM})$$

Global fits and averages:

- $K_L$ ,  $K_S$ , and  $K^{\pm}$ , dominant BRs and lifetime.
- Parameterization of the  $K \rightarrow \pi$  interaction (form factor)

# $K_L$ leading branching ratios and $\tau_L$

21 input measurements:	Parameter	Value	S
	$BR(K_{e3})$	0.4056(9)	1.3
5+3 Klev ratios	$BR(K_{\mu 3})$	0.2704(10)	1.5
<b>NA48</b> $K_{e3}$ /2tr and $\Gamma(3\pi^0)$	$BR(3\pi^0)$	0.1952(9)	1.2
<b>4 KLOE</b> BRs	$BR(\pi^+\pi^-\pi^0)$	0.1254(6)	1.3
<b>KLOE</b> , NA48 $\pi^+\pi^-/K_{l3}$	$BR(\pi^+\pi^-)$	$1.967(7) \times 10^{-3}$	1.1
<b>KLOE</b> , NA48 $\gamma\gamma/3\pi^0$	${ m BR}(\pi^+\pi^-\gamma)$	$4.15(9) \times 10^{-5}$	1.6
<b>PDG</b> ETAFIT for $\pi^+\pi^-/\pi^0\pi^0$	${ m BR}(\pi^+\pi^-\gamma_{ m DE})$	$2.84(8) \times 10^{-5}$	1.3
<b>KLOE</b> $\tau_{\rm L}$ from $3\pi^0$	${ m BR}(2\pi^0)$	$8.65(4) \times 10^{-4}$	1.4
$Vosburgh '72 \tau$	${ m BR}(\gamma\gamma)$	$5.47(4) \times 10^{-4}$	1.1
vusburgn 72 t <sub>L</sub>	$ au_{K_L}$	51.16(21)  ns	1.1

10 free parameters, 1 constraint: ΣBR=1

All  $\pi^+\pi^-/K_{l3}$  measurements are fully inclusive of inner breaksstrahlung KLOE measurement is fully inclusive of DE, negligible in KTeV one

PDG '04 –•

**BR(K**<sub>113</sub>) [%

# Evolution of the average BR values

This fit χ²/ndf = 19.8/12 (7.1%)
Minor differences wrt PDG04:
elimination of numerous old measurements

BR's shifted by  $6\sigma$ ,  $-6\sigma$ ,  $-5\sigma$ 



# $K_S$ leading branching ratios and $\tau_S$

6 input measurements:	Parameter	Value
<b>KLOE BR(Ke3)/BR(<math>\pi^+\pi^-</math>)</b>	$BR(\pi^+\pi^-)$	0.6920(5)
<b>KLOE</b> BR( $\pi^+\pi^-$ )/BR( $\pi^0\pi^0$ )	$BR(\pi^0\pi^0)$	0.3069(5)
Universal lepton coupling	$BR(K_{e3})$	$7.05(8) \times 10^{-4}$
NA48 BR( $K_{s}e3$ )/BR( $K_{L}e3$ )	$BR(K_{\mu3})$	$4.69(6) \times 10^{-4}$
$\tau_{\rm S}$ : non CPT-constrained fit value,	$ au_{K_S}$	89.59(6)  ps
2002 NA48 and 2003 KTeV measurements	$\mathcal{D}$	

5 free parameters:  $K_S \pi \pi$ ,  $K_S \pi^0 \pi^0$ ,  $K_S e^3$ ,  $K_S \mu^3$ ,  $\tau_S$ , 1 constraint:  $\Sigma BR=1$ 

**KLOE meas. completely determine the leading BR values**.

This fit  $\chi^2/ndf = 0.015/1$  (90%) S $\approx 1$  for any of the output values.

# K<sup>±</sup> leading branching ratios and $\tau^{\pm}$

D
$\Gamma$
B
В
В
В
В
В
$ au_{ m c}$

Parameter	Value	S
$BR(K_{\mu 2})$	63.47(18)%	1.3
$BR(\pi\pi^0)$	20.61(8)%	1.1
$BR(\pi\pi\pi)$	5.73(16)%	1.2
$BR(K_{e3})$	5.078(31)%	1.3
$BR(K_{\mu 3})$	3.359(32)%	1.9
$BR(\pi\pi^0\pi^0)$	1.757(24)%	1.0
$ au_{K^{\pm}}$	12.384(15)  ns	1.2

7 free parameters, 1 constraint: ΣBR=1 Don't use the result from Lobkowicz  $(\tau)$ , don't use the BRE from Chiang:

• 6 BRs constrained to sum to unit. **PDG**•, the correlation matrix not available

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# Evolution of the average BR values

- This fit  $\chi^2/ndf = 25.8/11$  (0.69%); PDG09 fit:  $\chi^2/ndf = 52/25$  (0.13%)
- some conflict among newer meas. involving BR(Ke3): the pulls are +0.6 and -2.1 for NA48 and KLOE respectively
- •some conflict among newer meas. involving BR(Kµ3):

the pulls are +1.0 and -3.2 for NA48 and KLOE respectively



# Vus determination

Physics results:

•  $|\mathbf{V}_{\mathrm{us}}| \times f_{+}(0)$ 

$$\begin{split} \Gamma_{K_{\ell 3}} &= \frac{G_F^2 m_K^5}{192\pi^3} C_K^2 S_{\rm EW} \left( |V_{us}| f_+^{K^0 \pi^-}(0) \right)^2 I_{K\ell} \left( 1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi} \right)^2 \\ &|V_{us}| / |V_{ud}| \times f_K / f_{\pi}. \end{split}$$

$$\begin{split} \frac{\Gamma_{K_{\ell 2}}}{I_{K_{\ell 2}}} &= \frac{|V_{us}|^2}{I_K^2} \frac{f_K^2}{I_K^2} \frac{m_K (1 - m_\ell^2 / m_K^2)^2}{I_K^2} \left( 1 + \delta_{\rm EM} \right) \end{split}$$

$$\frac{1}{\Gamma_{\pi_{\ell 2}}} = \frac{1}{|V_{ud}|^2} \frac{m_{\pi}}{f_{\pi}^2} \frac{1}{m_{\pi}(1 - m_{\ell}^2/m_{\pi}^2)^2} (1 + \delta_{\rm EM})$$

Global fits and averages:

- $K_L$ ,  $K_S$ , and  $K^{\pm}$ , dominant BRs and lifetime.
- Parameterization of the  $K \rightarrow \pi$  interaction (form factor)

# Parameterization of $K_{\ell 3}$ form factors

• Hadronic K $\rightarrow \pi$  matrix element is described by two form factors  $f_+(t)$  and  $f_0(t)$ defined by:  $\langle \pi(p_\pi) | \bar{s} \gamma_\mu u | K(p_K) \rangle = (p_\pi + p_K)_\mu f_+^{K\pi}(t) + (p_K - p_\pi)_\mu f_-^{K\pi}(t)$ 

$$f_0(t) = f_+(t) + \frac{t}{m_K^2 - m_\pi^2} f_-(t)$$

- Experimental or theoretical inputs to define *t*-dependence of *f*<sub>+,0</sub>(t). *f*<sub>-</sub>(t) term negligible for K<sub>e3</sub>.
  ➤ Taylor expansion: *f*<sub>+,0</sub><sup>Taylor</sup>(t) = 1 + λ'<sub>+,0</sub> t/m<sup>2</sup><sub>-+</sub> + 1/2 λ''<sub>+,0</sub> (t/m<sup>2</sup><sub>-+</sub>)<sup>2</sup>
- $\lambda'$  and  $\lambda''$  are strongly correlated: -95% for  $f_+(t)$ , and -99.96% for  $f_0(t)$ .

### One parameter parameterizations: ➤ Pole parameterization

$$\tilde{f}_{+,0}(t) = \frac{M_{V,S}^2}{M_{V,S}^2 - t}$$

> Dispersive approach plus  $K\pi$  scattering data for both  $f_+(t)$  and  $f_0(t)$ 

$$\bar{f}_{+}^{\text{disp}}(t) = \exp\left[\frac{t}{m_{\pi}^2}\left(\Lambda_{+} + H(t)\right)\right] \qquad \bar{f}_{0}^{\text{disp}}(t) = \exp\left[\frac{t}{\Delta_{K\pi}}(\ln C - G(t))\right]$$

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# Vector form factor from K<sub>e3</sub>

### **Quadratic expansion**:

- Measurements from ISTRA+, KLOE, KTeV, NA48 with K<sub>L</sub>e3 and K-e3 decays.
- Good fit quality:  $\chi^2/ndf=5.3/6(51\%)$  for all data;  $\chi^2/ndf=4.7/4(32\%)$  for K<sub>L</sub> only
- The significance of the quadratic term is 4.2 $\sigma$  from all data and 3.5 $\sigma$  from  $K_L$  only.
- Using all data or  $K_L$  only changes the space phase integrals  $I_{e3}^0$  and  $I_{e3}^{\pm}$  by 0.06%.
- Errors on  $I_{e3}$  are significantly smaller when K- data are included.

A pole parameterization is in good agreement with present data:

 $\tilde{f}_{+}(t) = M_V^2/(M_V^2 - t)$ , with  $M_V \sim 892$  MeV  $\lambda' = (m_{\pi^+}/M_V)^2$ ;  $\lambda'' = 2\lambda'^2$ 

- All four experiments quote value for  $M_V$  for pole fit to Ke3 data. The average value is  $Mv = 871 \pm 5 \text{ MeV} (\chi^2/\text{ndf}=3.8/3)$
- The values for  $\lambda_{+}'$  and  $\lambda_{+}''$  from pole expansion are in agreement with quadratic fit results.
- Using quadratic averages or pole fit results changes  $I_{e3}^0$  by 0.11%.

Improvements: dispersive parameterization for  $f_+(t)$ , with good analytical and unitarity properties and a correct threshold behavior,

(e.g. Bernard, Oertel, Passemar, Stern Phys. Rev. Lett. D80 (2009) 034034)

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# Vector and scalar form factor from K<sub>µ3</sub>

- $\lambda_{+}'$ ,  $\lambda_{+}''$  and  $\lambda_{0}$  measured for Kµ3 from ISTRA+, KLOE, KTeV, and NA48.
- NA48 results are difficult to accomodate in the  $[\lambda_{+}', \lambda_{+}'', \lambda_{0}]$  space.
- Fit probability varies from  $3 \times 10^{-7}$  (with NA48) to 14.5% (without NA48).



• Because of correlation, is not possible measure  $\lambda_0''$  at any plausible level of stat.

• Neglecting a quadratic term in the param. of scalar FF implies:  $\lambda_0' \rightarrow \lambda_0' + 3.5 \lambda_0''$ 

# Vector and scalar form factor from $K_{\mu3}$

The dispersive form factor parameterization clearly illustrate the contrast between  $K_{\mu3}$  result from NA48 and those from the other experiment



|V<sub>ud</sub>| & |V<sub>us</sub>| determination from kaon decays P. Massarotti FP&CPV torino May 26, 2010 16

# Vector and scalar form factor from $K_{\ell 3}$

•Comparison of phase-space integrals evaluated from our averages of results of quadratic-linear and dispersive fits

Integral	$\lambda'_+,\lambda''_+,\lambda_0$	$\Lambda_+, \ln C$	Rel. diff.
$I(K_{e3}^0)$	0.15457(20)	0.15476(18)	+0.12%
$I(K_{e3}^{\pm})$	0.15894(21)	0.15922(18)	+0.18%
$I(K^0_{\mu 3})$	0.10266(20)	0.10253(16)	-0.13%
$I(K_{\mu3}^{\pm})$	0.10564(20)	0.10559(17)	-0.05%
$ ho(\dot{K_{e3}},K_{\mu3})$	+0.56	+0.38	

The intergals, when evaluated from the dispersive fit results, tend to be slightly greater (no more than 0.2%) than from the the quadratic fit results.

Global fits and averages:

- $K_L$ ,  $K_S$ , and  $K^{\pm}$ , dominant BRs and lifetime.
- Parameterization of the  $K \rightarrow \pi$  interaction (form factor)

Physics results:

- $|\mathbf{V}_{\mathrm{us}}| \times f_{+}(0)$
- $|\mathbf{V}_{us}|/|\mathbf{V}_{ud}| \times f_K/f_{\pi}$ .
- Theoretical estimations of  $f_+(0)$  and  $f_K/f_{\pi}$ .
- $\bullet$  V<sub>us</sub> and V<sub>ud</sub> determinations.
- Bounds on helicity suppressed amplitudes.

# Determination of $|V_{us}| \times f_{+}(0)$

$$\Gamma(K_{l3(\gamma)}) = \frac{C_{K}^{2} G_{F}^{2} M_{K}^{5}}{192\pi^{3}} S_{EW} |V_{us}|^{2} |f_{+}^{K^{0}\pi^{-}}(0)|^{2} I_{K\ell}(\lambda_{+,0}) (1 + \delta_{SU(2)}^{K} + \delta_{em}^{K\ell})^{2}$$

with  $K = K^+$ ,  $K^0$ ;  $\ell = e, \mu$  and  $C_K^2 = 1/2$  for  $K^+$ , 1 for  $K^0$ 

### **Inputs from theory:**

- $S_{EW}$  Universal short distance EW correction (1.0232)
- $\frac{\delta^{K}_{SU(2)}}{\text{strong SU}(2) \text{ breaking}}$
- $\frac{\delta^{K\ell}}{e^{m}}$  Long distance EM effects

 $f_{+}^{K^{0}\pi^{-}}(0)$  Form factor at zero momentum transfer (t=0)

Callan-Treiman

 $I_{K\ell}(\lambda)$ 

**Inputs from experiment:** 

 $\Gamma(K_{l3(\gamma)})$  Branching ratios properly inclusive of radiative effects; lifetimes

Phase space integral:  $\lambda$ 's parameterize form factor dependence on *t* :  $K_{e3}$ : only  $\lambda_+$ 

 $K_{\mu\beta}$ : need  $\lambda_+$  and  $\lambda_0$ 

 $|V_{ud}| \& |V_{us}|$  determination from kaon decays

FP&CPV torino May 26, 2010 19

# SU(2) and *em* corrections



(values used to extract  $|V_{us}|f_{+}(0)$ )

- $\delta_{em}$  for full phase space: all measurements assumed fully inclusive.
- Different estimates of  $\delta_{em}$  agree within the quoted errors.
- Available correlation matrix between different corrections for  $\delta_{em}$ .

Determination of  $|V_{us}| \times f_{+}(0)$ 

$\Gamma(K_{l3(\gamma)}) = \frac{C_{K}^{2} G_{F}^{2} M_{K}^{5}}{192\pi^{3}} S_{EW}  V_{uS} ^{2}  f_{+}^{K^{0}\pi^{-}}(0) ^{2} I_{K\ell}(\lambda_{+,0}) (1 + \delta^{K}_{SU(2)} + \delta^{K\ell}_{em})^{2}$								
0.213 0.214 0.215	0.216 0.217	-		% err ×10⁻²	<b>BR</b> ×10 <sup>-2</sup>	τ×10-2	∆×10-2	I <sub>Kt</sub> ×10-2
K <sub>L</sub> e3	<b></b> -	$K_L e3$	0.2163(6)	26	9	20	11	6
<b>Κ<sub>L</sub> μ3</b>		<i>К<sub>L</sub></i> µ3	0.2166(6)	29	15	18	11	8
K <sub>S</sub> e3		K <sub>s</sub> e3	0.2155(13)	61	60	3	11	6
K <sup>±</sup> e3	•	K±e3	0.2160(11)	52	31	8	40	6
<b>Κ</b> <sup>±</sup> μ <b>3</b>		<i>К</i> ±µ3	0.2158(14)	63	47	8	39	8
Average: $ V_{us}  f_+(0) = 0.2163(5) \qquad \chi^2/\text{ndf} = 0.77/4 (94\%)$								
<b>V</b> <sub>ud</sub> & <b>V</b> <sub>us</sub> determination from kaon decays <b>P. Massarotti FP&amp;CPV torino May 26, 2010 21</b>								

Theoretical estimate of  $f_{\perp}(0)$ 

Leutwyler & Roos estimate still widely used:  $f_{+}(0) = 0.961(8).$ 

Lattice evaluations generally agree well with this value; use RBC-UKQCD10 value:  $f_{+}(0) = 0.959(5) (0.5\%)$ accuracy, total err.).



 $|V_{ud}| \& |V_{us}|$  determination from kaon decays

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FP&CPV torino May 26, 2010 22  $V_{us}/V_{ud}$  determination from BR(K<sub>µ2</sub>)



K12:  $|V_{us}|/|V_{ud}| f_K / f_{\pi} = 0.2758(5)$  and  $f_K / f_{\pi} = 1.193(6)$ , obtain  $|V_{us}|/|V_{ud}| = 0.2312(13)$ 

 $|V_{ud}| \& |V_{us}|$  determination from kaon decays

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**FP&CPV torino May 26, 2010 23** 



 $K_{\mu 2}$ : sensitivity to NP

Comparison of  $V_{us}$  from  $K_{\ell 2}$  (helicity suppressed) and from  $K_{\ell 3}$  (helicity allowed) To reduce theoretical uncertainties study the quantity:

$$R_{l23} = \left| \frac{V_{us}(K_{\ell 2})}{V_{us}(K_{\ell 3})} \times \frac{V_{ud}(0^+ \to 0^+)}{V_{ud}(\pi_{\ell 2})} \right|$$

Within SM  $R_{\ell 23} = 1$ ; NP effects can show as scalar currents due to a charged Higgs:

$$R_{\mu 23} \approx \left| 1 - \frac{m_{K^+}^2}{m_{H^+}^2} \frac{\tan^2 \beta}{1 + \epsilon_0 \tan \beta} \right|$$

 $K_{\mu 2}$ : sensitivity to NP!

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 $R_{\ell 23}$  is accessible via  $BR(K_{\mu 2})/BR(\pi_{\mu 2})$ ,  $V_{us}f_{+}(0)$ , and  $V_{ud}$ , and  $f_{K}/f_{\pi}/f_{+}(0)$  determinations.

- Using K<sup>±</sup> fit results, assuming unitarity for Vus(K<sub> $\ell 3$ </sub>) and using f<sub>K</sub>/f<sub> $\pi$ </sub>/f<sub>+</sub>(0) from lattice: R<sub> $\ell 23$ </sub>= 0.999(7)
- Uncertainty dominated by  $f_{\rm K}/f_{\pi}/f_{+}(0)$ .
- 95% CL excluded region (with  $\varepsilon_0 \sim 0.01$ ).
- In tan $\beta$ -M<sub>H±</sub> plane, R<sub>µ23</sub> fully cover the region uncovered by BR(B $\rightarrow \tau \nu$ ).



# Conclusions

- Dominant  $K_S$ ,  $K_L$ , and  $K^{\pm}$  BRs, and lifetime known with very good accuracy.
- Dispersive approach for form factors.
- Constant improvements from lattice calculations of  $f_+(0)$  and  $f_K/f_{\pi}$ :
- $|V_{us}| f_{+}(0)$  at 0.2% level.
- $|V_{us}|$  measured with 0.6% accuracy (with  $f_{+}(0)=0.959(5)$ ) Dominant contribution to uncertainty on  $|V_{us}|$  still from  $f_{+}(0)$ . CKM unitarity test satisfied at 0.17 $\sigma$  level
- Comparing  $|V_{us}|$  values from Kµ2 and K13, exclude large region in the  $(m_{H^+}, \tan\beta)$  plane, complementary to results from  $B \rightarrow \tau \nu$  decays.

# **KLOE-2** experiment



28

# **KLOE-2** experiment



 $e^+e^- \rightarrow e^+e^-\gamma^*\gamma^* \rightarrow e^+e^-X$ 

# **KLOE-2** experiment

### Minimal detector upgrade

 Tagger for γγ physics: to detect off-momentum e<sup>±</sup> from

$$e^+e^- \rightarrow e^+e^-\gamma^*\gamma^* \rightarrow e^+e^-X$$

- Low Energy Tagger (130-230 MeV) calorimeters, LYSO + SiPM
- High Energy Tagger (E > 400 MeV) position sensitive detectors (strong energy-position correlation ⇒ use the DAΦNE magnets as e<sup>±</sup> spectrometer)
- No QCAL on quadrupoles (Pb shields)
- Luminosity goal: 5 fb<sup>-1</sup> @ √s ≈M<sub>φ</sub>
   Roll-in (Dec 2009) and alignment (Jan 2010)
   Commisioning Mid June 2010





# Determination of $|V_{us}| \times f_+(0)$ : improvements

		% err ×10⁻²	<b>BR</b> ×10 <sup>-2</sup>	τ×10-2	Δ×10-2	I <sub>Kl</sub> ×10-2
$K_L e3$	0.2163(6)	26	9	20	11	6
<i>К<sub>L</sub></i> µ3	0.2166(6)	29	15	18	11	8
K <sub>S</sub> e3	0.2155(13)	61	60	3	11	6
K±e3	0.2160(11)	52	31	9	40	6
<i>K</i> ±µ3	0.2158(14)	63	47	8	39	8

**V**<sub>ud</sub> & **V**<sub>us</sub> determination from kaon decays **P. Massarotti FP&CPV torino May 26, 2010 30** 

# Determination of $|V_{us}| \times f_+(0)$ : improvements

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$K_L e3$	0.2163(6)	26	9	13	11	6
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K±e3	0.2160(11)	52	25	5	25	6
<i>K</i> ±µ3	0.2158(14)	63	23	8	25	8

arXiv: 1003.3868v2 [hep-ex] 29 Mar 2010: Approved by EPJ

# Additional information

|Vud| & |Vus| determination from kaon decaysP. MassarottiFP&CPV torino May 26, 201031

# Vud: error budget

### **OPPORTUNITIES FOR IMPROVEMENT**

- Goal remains to tighten the window for new physics by reducing the uncertainty on V<sub>ud</sub>.
- Uncertainty on calculated radiative correction ∆<sub>R</sub> is the dominant contribution to the error budget.
- Nuclear-structure-dependent corrections, δ<sub>c</sub> and δ<sub>NS</sub>, can be tested by experiment; this has already led to improvements, but more are still possible.

Data on "well known" transitions can be made more precise, and new cases can be measured.



# Vud : data

### WORLD DATA FOR $0^+ \rightarrow 0^+$ DECAY, 2008



 $|V_{ud}| \& |V_{us}|$  determination from kaon decays

P. Massarotti

FP&CPV torino May 26, 2010 33

Dispersive parameterization: a test of lattice calculations

Scalar form factor  $f_0(t) = \widetilde{f_0}(t) f_+(0)$  extrapolation at **Callan-Treiman** point:

$$\bar{f}_0(\Delta_{K\pi}) = \frac{f_K}{f_\pi} \frac{1}{f_+(0)} + \Delta_{CT} \qquad \Delta_{CT} = (-3.5 \pm 8) \times 10^{-3}$$

links  $f_{+}(0)$  and  $f_{\rm K}/f_{\pi}$  with  $\lambda_0$  measured in Kµ3 decays.

 $\widetilde{f}_0(\Delta_{K\pi})$  is evaluated fitting K<sub>L</sub>µ3 with a dispersive parameterization

$$\tilde{f}_0(t) = \exp\left(\frac{t}{\Delta_{K\pi}}\log(C - G(t))\right)$$

G(t) from  $K\pi$  scattering data. To fit we use a 3<sup>rd</sup> order expansion



From CT, using  $f_K/f_{\pi}=1.193(6)$  obtain:  $f_+(0)=0.974(12)$  in agreement with RBC/UKQCD10 value:  $f_+(0) = 0.959(5)$ .