## THEORY OF HADRONIC B DECAYS

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## CAVEAT

IN 20-25 MINUTES I WILL NOT BE ABLE TO COVER ALL OF THE LATEST DEVELOPMENTS IN THIS WIDESPREAD FIELD.

AS I AM GOING TO DISCUSS CONCEPTUAL AS WELL AS PHENOMENOLOGICAL ASPECTS, THE PHENOMENOLOGICAL PART OF MY TALK WILL NECESSARILY BE SELECTIVE AND SHORT.

I APOLOGIZE FOR NOT COVERING YOUR FAVOURITE TOPIC(S)!

## OUTLINE

## THEORY

REVIEW OF QCDF / PQCD / SCET
RECENT DEVELOPMENTS: CHARMING PENGUINS
GLAUBER GLUONS IN PQCD
PERTURBATIVE CORRECTIONS IN QCDF

## PHENOMENOLOGY

$$
\begin{array}{ll}
\text { TREE-DOMINATED DECAYS: } & B \rightarrow \pi \pi / \pi \rho / \rho \rho \\
\text { PENGUIN-DOMINATED DECAYS: } & B \rightarrow \pi K
\end{array}
$$

## THEORY

## Charmless hadronic $B$ decays

Particularly rich laboratory to probe the nature of flavour-changing

$$
b \rightarrow u \bar{u} d / u \bar{u} s \quad b \rightarrow d \bar{q} q / s \bar{q} q
$$

quark transitions and to test the CKM mechanism of CP violation

The challenge: quantitative control over $\left\langle M_{1} M_{2}\right| Q_{i}|B\rangle$


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Two complementary strategies:

- flavour symmetries (isospin, $\mathrm{SU}(3), \ldots$ )
extract hadronic matrix elements from data
use approximate symmetries of QCD $\Rightarrow$ relate strong decay amplitudes
- dynamical approaches
calculate hadronic matrix elements in QCD
exploit factorization in heavy quark limit $m_{b} \gg \Lambda_{Q C D}$
some additional insights from light-cone sum rules


## Factorization

Exploit $m_{b} \gg \Lambda_{Q C D}$ to disentangle

- short-distance effects $\sim m_{b} \quad \Rightarrow \quad$ perturbatively calculable
- long-distance effects $\sim \Lambda_{Q C D} \quad \Rightarrow \quad$ universal hadronic parameters

This is not a model, but describes QCD in the well-defined limit $m_{b} \rightarrow \infty$

Why are there three different incarnations of factorization?

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Why are there three different incarnations of factorization?

Factorization of hard exclusive processes pioneered more than 30 years ago for $\pi \gamma^{*} \rightarrow \pi$

$$
\langle\pi| J|\pi\rangle \simeq \int d u d v T(u, v) \phi_{\pi}(u) \phi_{\pi}(v)
$$

[Brodsky, Lepage 79]
For charmless hadronic $B$ decays expect

$$
\left\langle M_{1} M_{2}\right| Q_{i}|B\rangle \simeq \int d \omega d u d v T_{i}(\omega, u, v) \phi_{B}(\omega) \phi_{M_{1}}(v) \phi_{M_{2}}(u) \quad ?
$$

Problem: convolutions over $\omega$ and $v$ diverge $\Rightarrow$ not dominated by hard gluon exchange!?
but still: convolution over $u$ finite $\quad \Rightarrow$ "colour transparency"

## QCDF / pQCD / SCET

- QCD factorization

$$
\begin{aligned}
\left\langle M_{1} M_{2}\right| Q_{i}|B\rangle \simeq & F^{B M_{1}}(0) \int d u T_{i}^{\prime}(u) \phi_{M_{2}}(u) \\
& +\int d \omega d u d v T_{i}^{\prime \prime}(\omega, u, v) \phi_{B}(\omega) \phi_{M_{1}}(v) \phi_{M_{2}}(u)
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$$

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- perturbative QCD
[Keum, Li, Sanda 00]
$\left\langle M_{1} M_{2}\right| Q_{i}|B\rangle \simeq \int d \omega d u d v d k_{i \perp} T_{i}\left(\omega, u, v, k_{i \perp}\right) \phi_{B}\left(\omega, k_{1 \perp}\right) \phi_{M_{1}}\left(v, k_{2 \perp}\right) \phi_{M_{2}}\left(u, k_{3 \perp}\right)$
+ Sudakov resummation $\Rightarrow$ endpoint divergence smeared out by transverse momenta


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+ Sudakov resummation $\Rightarrow$ endpoint divergence smeared out by transverse momenta
- Soft-collinear effective theory

SCET = QCDF: simply EFT vs diagrammatical formulation
BPRS $\neq \mathrm{BBNS}$ : phenomenological implementation for $B \rightarrow M_{1} M_{2}$ quite different issues: $F^{B M_{1}}$ traded for $\xi^{B M_{1}}, \alpha_{s}\left(\sqrt{\lambda m_{b}}\right)$ non-perturbative, $\ldots$ (minor)
long-distance charm loops, zero-bin subtractions (major)

## Long-distance charm loops

Old charming penguin story - two different questions:

- power-suppressed but numerically important
$\Rightarrow$ not supported by light-cone sum rule estimate
[Colangelo et al. 89; Ciuchini et al. 97+]
[Khodjamirian, Mannel, Melic 03]
- leading power spoiling factorization does the threshold region with a non-relativistic $c \bar{c}$ pair require a special treatment?

Recent work addresses second question
(a) $e^{+} e^{-} \rightarrow$ hadrons

$\int d q^{2} \ldots \operatorname{lm} \Pi\left(q^{2}\right)$
(b) $B \rightarrow X_{s} \ell^{+} \ell^{-}$

$\int d q^{2} \ldots\left|\Pi\left(q^{2}\right)\right|^{2}$
(c) charming penguins

$\int d q^{2} \ldots \Pi\left(q^{2}\right)$
$\Rightarrow$ global quark-hadron duality holds in (a) and (c), but breaks down in (b)
$\Rightarrow$ no special treatment required in (c), long-distance charm loops are power-suppressed

## Glauber gluons in PQCD

It has been realized recently that $k_{T}$-factorization breaks down in $p p \rightarrow h_{1} h_{2} X$ at high $p_{T}$

- problem related to a peculiar mode: Glauber gluons
- effect is a non-universal long-distance contribution $\Rightarrow$ ruins $k_{T}$-factorization
- problem not present in collinear factorization

Important for pQCD approach to non-leptonic $B$ decays

- confirmed that problem exists $\Rightarrow$ modification of pQCD approach
- claimed that it leads to an universal soft factor $e^{i S} \Rightarrow k_{T}$-factorization still holds
- $e^{i S}$ from fit to $\operatorname{Br}\left(\pi^{0} \pi^{0}\right) \Rightarrow$ large complex $C \quad \Rightarrow$ "solves" $\pi \pi / \pi K$ puzzles

Issues: operator definition of soft factor?

- if universal why associated to $\pi$ but not to $\rho$ ? $\Rightarrow$ would worsen $\operatorname{Br}\left(\rho^{0} \rho^{0}\right)$
- at present I consider this rather as a fit than as a dynamical explanation


## Perturbative corrections in QCDF

Ongoing effort to compute NNLO corrections in QCDF

$$
\left\langle M_{1} M_{2}\right| Q_{i}|B\rangle \simeq F^{B M_{1}} T_{i}^{\prime} \otimes \phi_{M_{2}}+T_{i}^{\prime \prime} \otimes \phi_{B} \otimes \phi_{M_{1}} \otimes \phi_{M_{2}}
$$

strong phases $\sim \mathcal{O}\left(\alpha_{s}\right) \quad \Rightarrow \quad$ NNLO is first correction for direct CP asymmetries!

| Status | 2-loop vertex corrections ( $T_{i}^{\prime}$ ) | 1-loop spectator scattering ( $T_{i}^{\prime \prime}$ ) |
| :---: | :---: | :---: |
| Trees | [GB 07, 09] [Beneke, Huber, Li 09 | [Beneke, Jäger 05] <br> [Kivel 06] <br> [Pilipp 07] |
| Penguins | in progress | [Beneke, Jäger 06] [Jain, Rothstein, Stewart 07] |

- factorization found to hold (as expected) at highly non-trivial order
- direct CP asymmetries not yet available at NNLO
- first NNLO results for CP-averaged branching ratios of tree-dominated decays
[GB, Pilipp 09; Beneke, Huber, Li 09]


## Summary of theory part

|  | BBNS (QCDF) | pQCD | BPRS (SCET) |
| :--- | :---: | :---: | :---: |
| $\alpha_{s}\left(\sqrt{\Lambda m_{b}}\right.$ ) | perturbative | perturbative | non-perturbative |
| charm loops | perturbative <br> (small phase) | perturbative <br> (small phase) | non-perturbative <br> (large phase from fit to data) |
| weak annihilation <br> (power correction) | non-perturbative <br> (crude model, arbitrary phase) | perturbative <br> (large phase) | perturbative <br> (with zero bins, small phase) |
| strong phases | generically small <br> $\left(\sim \alpha_{s}, 1 / m_{b}\right)$ | can be sizeable <br> (annihilation, Glaubers) | charm be sizeable |
| perturbative <br> calculation | partially NNLO |  |  |
| hadronic input | prom lattice + <br> QCD sum rules | NLO <br> from QCD sum rules <br> + data, model $\phi_{B}(x, b)$ | from QCD sum rules <br> + data, model $\xi_{J}^{B M}(z)$ |

- theory predictions for direct CP asymmetries can unfortunately differ a lot!
- measurements (even bounds) of pure annihilation decays highly appreciated:
$B_{d} \rightarrow K^{-} K^{+}, B_{s} \rightarrow \pi \pi / \pi \rho / \rho \rho$


## PHENOMENOLOGY

## $B \rightarrow \pi \pi / \pi \rho / \rho \rho$

Large $\operatorname{Br}\left(\pi^{0} \pi^{0}\right)=(1.55 \pm 0.19) \cdot 10^{-6}$ challenging for dynamical approaches

- seems to indicate large $C / T \quad \Rightarrow \quad$ interesting in view of $\pi K$ puzzle
- Babar: $(1.83 \pm 0.21 \pm 0.13) \cdot 10^{-6} \quad$ Belle: $(1.1 \pm 0.3 \pm 0.1) \cdot 10^{-6}$
$\Rightarrow$ reconsider in the light of NNLO calculation in QCDF and $\pi \rho / \rho \rho$ data

Tree amplitudes now completely determined to NNLO

$$
\begin{aligned}
T \sim \alpha_{1}(\pi \pi)= & {[1.008] v_{0}+[0.022+0.009 i] v_{1}+[0.024+0.026 i] v_{2} } \\
& -[0.014]]_{S_{1}}-[0.016+0.012 i]_{S_{2}}-[0.008]_{1 / m_{b}} \\
= & 1.015_{-0.029}^{0.020}+\left(0.023_{-0.015}^{+0.015}\right) i \\
C \sim \alpha_{2}(\pi \pi)= & {[0.224] v_{0}-[0.174+0.075 i] v_{1}-[0.029+0.046 i] v_{2} } \\
& +[0.084]_{s_{1}}+[0.037+0.022 i] s_{2}+[0.052]_{1 / m_{b}} \\
= & 0.194_{-0.095}^{+0.130}-\left(0.099_{-0.056}^{+0.057}\right) i
\end{aligned}
$$

- individual NNLO corrections quite significant, but substantial cancellations
- $\operatorname{Re}\left(\alpha_{2}\right)$ still uncertain, mainly due to poor knowledge of $\lambda_{B}^{-1}=\int \frac{d \omega}{\omega} \phi_{B}(\omega)$


## CP-averaged branching ratios

| Mode | QCDF | B | Experiment |
| :---: | :---: | :---: | :---: |
| $\pi^{-} \pi^{0}$ | $6.22_{-2.01}^{+2.37}$ | 5.46 | $5.59_{-0.40}^{+0.41}$ |
| $\rho_{L}^{-} \rho_{L}^{0}$ | $21.0_{-7.3}^{+8.5}$ | 21.3 | $22.5_{-1.9}^{+1.9}$ |
| $\pi^{-} \rho^{0}$ | $9.34_{-3.23}^{+4.00}$ | 10.4 | $8.3_{-1.3}^{+1.2}$ |
| $\pi^{0} \rho^{-}$ | $15.1_{-5.0}^{+5.7}$ | 11.9 | $10.9_{-1.5}^{+1.4}$ |
| $\pi^{+} \pi^{-}$ | $8.96_{-3.32}^{+3.78}$ | 5.21 | $5.16_{-0.22}^{+0.22}$ |
| $\pi^{0} \pi^{0}$ | $0.35_{-0.21}^{+0.37}$ | 0.63 | $1.55_{-0.19}^{+0.19}$ |
| $\pi^{+} \rho^{-}$ | $22.8_{-8.0}^{+9.1}$ | 13.2 | $15.7_{-1.8}^{+1.8}$ |
| $\pi^{-} \rho^{+}$ | $11.5_{-4.3}^{+5.1}$ | 8.41 | $7.3_{-1.2}^{+1.2}$ |
| $\pi^{ \pm} \rho^{\mp}$ | $34.3_{-10.0}^{+11.5}$ | 21.6 | $23.0_{-2.3}^{+2.3}$ |
| $\pi^{0} \rho^{0}$ | $0.52_{-0.42}^{+0.76}$ | 1.64 | $2.0_{-0.5}^{+0.5}$ |
| $\rho_{L}^{+} \rho_{L}^{-}$ | $30.3_{-11.2}^{+12.9}$ | 22.3 | $23.6_{-3.2}^{+3.2}$ |
| $\rho_{L}^{0} \rho_{L}^{0}$ | $0.44_{-0.37}^{+0.66}$ | 1.33 | $0.69_{-0.30}^{+0.30}$ |

- theo. uncertainties highly correlated ( $\left.F^{B M_{1}},\left|V_{u b}\right|\right)$
- colour-suppressed modes $\pi^{0} \pi^{0} / \pi^{0} \rho^{0} / \rho^{0} \rho^{0}$ rather uncertain ( $\lambda_{B}$ and $1 / m_{b}$ )
- $\rho^{0} \rho^{0}$ and $\pi^{0} \rho^{0}$ (with smaller penguins) fit better than $\pi^{0} \pi^{0}$
- preference for enhanced colour-suppressed amplitude

B: enhanced colour-suppressed amplitude with

$$
\lambda_{B}=0.2, F_{+}^{B \pi}=0.21 \text { and } A_{0}^{B \rho}=0.27
$$

## Testing factorization

Eliminate dependence on $F^{B M_{1}}$ and $\left|V_{u b}\right|$ via

$$
\mathcal{R}_{M_{3}}\left(M_{1} M_{2}\right)=\frac{\Gamma\left(B \rightarrow M_{1} M_{2}\right)}{d \Gamma\left(B \rightarrow M_{3} \ell \nu\right) /\left.d q^{2}\right|_{q^{2}=0}}
$$

$\Rightarrow$ requires measurement of semileptonic decay spectrum and extrapolation to $q^{2}=0$

$$
B \rightarrow \pi \ell \nu
$$


[Babar 07]
$\Rightarrow \quad\left|V_{u b}\right| F_{+}^{B \pi}(0)=(9.1 \pm 0.7) \cdot 10^{-4}$
$B \rightarrow \rho \ell \nu$

[Babar 05; Belle 07; CLEO 07; figure from Flynn et al 08]
currently insufficient to extract $\left|V_{u b}\right| A_{0}^{B \rho}(0)$

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$$

| Mode | QCDF | B | Experiment |
| :---: | :---: | :---: | :---: |
| $\mathcal{R}_{\pi}\left(\pi^{-} \pi^{0}\right)$ | $0.70_{-0.08}^{+0.12}$ | 0.95 | $0.81_{-0.14}^{+0.14}$ |
| $\mathcal{R}_{\rho}\left(\rho_{L}^{-} \rho_{L}^{0}\right)$ | $1.91_{-0.23}^{+0.32}$ | 2.38 | n.a. |
| $\mathcal{R}_{\rho}\left(\pi^{-} \rho^{0}\right)$ | $0.85_{-0.14}^{+0.22}$ | 1.16 | n.a. |
| $\mathcal{R}_{\pi}\left(\pi^{0} \rho^{-}\right)$ | $1.71_{-0.24}^{+0.27}$ | 2.07 | $1.57_{-0.32}^{+0.32}$ |
| $\mathcal{R}_{\pi}\left(\pi^{+} \pi^{-}\right)$ | $1.09_{-0.20}^{+0.22}$ | 0.97 | $0.80_{-0.13}^{+0.13}$ |
| $\mathcal{R}_{\pi}\left(\pi^{+} \rho^{-}\right)$ | $2.77_{-0.31}^{+0.32}$ | 2.46 | $2.43_{-0.47}^{+0.47}$ |
| $\mathcal{R}_{\rho}\left(\pi^{-} \rho^{+}\right)$ | $1.12_{-0.14}^{+0.20}$ | 1.01 | n.a. |
| $\mathcal{R}_{\rho}\left(\rho_{L}^{+} \rho_{L}^{-}\right)$ | $2.95_{-0.35}^{+0.37}$ | 2.68 | n.a. |
| $R\left(\rho_{L}^{-} \rho_{L}^{0} / \rho_{L}^{+} \rho_{L}^{-}\right)$ | $0.65_{-0.11}^{+0.16}$ | 0.89 | $0.89_{-0.14}^{+0.14}$ |
| $R\left(\pi^{-} \pi^{0} / \pi^{+} \pi^{-}\right)$ | $0.65_{-0.14}^{+0.19}$ | 0.98 | $1.01_{-0.09}^{+0.09}$ |

- $\pi^{-} \pi^{0} / \rho^{-} \rho^{0}$ provide clean access to $|T+C|$
- good overall description (in particular for small $\lambda_{B}$ )
- only exception $\operatorname{Br}\left(\pi^{0} \pi^{0}\right)$ which has a substantial penguin contribution and large theo. uncertainties
$B$ : enhanced colour-suppressed amplitude with

$$
\lambda_{B}=0.2, F_{+}^{B \pi}=0.21 \text { and } A_{0}^{B \rho}=0.27
$$

## The $B \rightarrow \pi K$ puzzle

Since 2006 branching ratios do no longer look puzzling

|  | QCDF | Experiment |
| :---: | :---: | :---: |
| $R_{C}$ | $1.15_{-0.18}^{+0.21}$ | $1.12 \pm 0.07$ |
| $R_{n}$ | $1.16_{-0.20}^{+0.24}$ | $1.02 \pm 0.07$ |

$$
R_{c}=2 \frac{\Gamma\left(B^{-} \rightarrow \pi^{0} K^{-}\right)}{\Gamma\left(B^{-} \rightarrow \pi^{-} \bar{K}^{0}\right)} \quad R_{n}=\frac{1}{2} \frac{\Gamma\left(\bar{B}_{d}^{0} \rightarrow \pi^{+} K^{-}\right)}{\Gamma\left(\bar{B}_{d}^{0} \rightarrow \pi^{0} \bar{K}^{0}\right)}
$$

But direct CP asymmetries look somewhat odd

| $A_{C P}[\%]$ | QCDF | Experiment |
| :---: | :---: | :---: |
| $\pi^{0} K^{-}$ | $9.4_{-14.6}^{+12.1}$ | $5.0_{-2.5}^{+2.5}$ |
| $\pi^{+} K^{-}$ | $5.6_{-15.1}^{+12.2}$ | $-9.8_{-1.1}^{+1.2}$ |
| $\Delta A_{C P}$ | $3.8_{-2.7}^{+2.9}$ | $14.8_{-2.8}^{+2.7}$ |

$\Delta A_{C P}$ can be predicted quite precisely
$\Rightarrow$ could be NP but also large + complex
$C / T$ beyond factorization

More robust SM test given by isospin sum rule

| $A_{C P}[\%]$ | Sum rule | Experiment |
| :---: | :---: | :---: |
| $\pi^{0} \bar{K}^{0}$ | $-0.15 \pm 0.04$ | $-0.01 \pm 0.10$ |

$$
A_{C P}\left(\pi^{+} K^{-}\right)+A_{C P}\left(\pi^{-} \bar{K}^{0}\right) \simeq A_{C P}\left(\pi^{0} K^{-}\right)+A_{C P}\left(\pi^{0} \bar{K}^{0}\right)
$$

$\Rightarrow$ predicts large $A_{C P}\left(\pi^{0} \bar{K}^{0}\right)$, but discrimination of NP requires much more data

## Mixing-induced CP asymmetry

Naive expectation neglecting doubly Cabibbo-suppressed terms

$$
S_{\pi^{0} K_{S}} \simeq \sin (2 \beta)_{J / \psi K_{S}}=0.681 \pm 0.025 \quad S_{\pi^{0} K_{S}} \stackrel{\exp }{=} 0.57 \pm 0.17
$$

"tree pollution" estimated to increase $S_{\pi^{0} K_{S}}$ by 0.04-0.12
[Beneke 05; Williamson, Zupan 06]

Isospin relations reveal some tension with $\pi K$ data

uses two $\mathrm{SU}(3)$ relations $\Rightarrow R_{T+C}$ and $R_{q}$ SU(3) breaking estimated with QCDF
$\Rightarrow S_{\pi^{0} K_{S}}=\left.\left.\left.\left.0.99_{-0.08}^{+0.01}\right|_{\exp -0.001} ^{+0.000}\right|_{R_{T+C}-0.11}\right|_{R_{q}-0.07} ^{+0.00}\right|_{\gamma}$
uncertainty reducible by lattice determination of $F^{B K} / F^{B \pi}$

Additional information on $\pi K$ puzzle can be obtained from $\pi K^{*} / \rho K$ sector
$\Rightarrow$ at present data still insufficient, but expect sizeable direct CP asymmetries
[Chiang, London 09; Gronau, Pirjol, Zupan 10]

## Conclusion

Factorization is based on a twofold expansion in $\alpha_{s}\left(m_{b}\right)$ and $1 / m_{b}$

- perturbative calculations are reaching NNLO precision
- unfortunately no similar progress (yet) on power corrections

Satisfactory overall description of $\pi \pi / \pi \rho / \rho \rho$ branching ratios

- large $\operatorname{Br}\left(\pi^{0} \pi^{0}\right)$ still puzzling, but not necessarily a failure of factorization
- $A_{C P}\left(\pi^{+} \pi^{-}\right)$from Belle also somewhat large, but $\sim 2 \sigma$ above Babar value

The $B \rightarrow \pi K$ puzzle continues to be exciting

- depends crucially on experimental progress on $\pi^{0} \bar{K}^{0}$ observables

There are many interesting topics which I could not discuss

- $B_{s}$ decays, $B \rightarrow V V$ polarization puzzle, $B$ decays with scalars, baryons, 3-body, $\ldots$

