

RENZO L. RICCA

Department of Mathematics & Applications, U. Milano-Bicocca renzo.ricca@unimib.it

Joint work with ALICE ROITBERG (U. Milano-Bicocca)



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• Gross-Pitaevskii (1961) equation: $\Psi = \Psi(\mathbf{x}, t)$

$$\mathrm{i}\,\hbar\,\partial_t\Psi = \left(-\frac{\hbar^2}{2m}\,\boldsymbol{\nabla}^2 + \mathfrak{g}|\Psi|^2 + V
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 $\mathfrak{g} = (4\pi\hbar^2 a_s)/m$, $V = V(\boldsymbol{x}, t)$.



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(Anderson et al. 1995)



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• Madelung (1926) transformation:

 $\Psi = \sqrt{\rho} \exp \left[\mathrm{i} (m/\hbar) \theta \right]$.



(Anderson et al. 1995)

Hydrodynamic derivation of the Gross–Pitaevskii equation in general Riemannian metric

Alice Roitberg¹ and Renzo L Ricca^{1,2,*}

 ¹ Department of Mathematics & Applications, University of Milano-Bicocca, Via Cozzi 55, 20125 Milano, Italy
 ² BDIC, Beijing University of Technology, 100 Pingleyuan, Beijing 100124, People's Republic of China

E-mail: renzo.ricca@unimib.it

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Abstract

Here we show that the Gross–Pitaevskii equation (GPE) for Bose–Einstein condensates (BECs) admits hydrodynamic interpretation in a general Riemannian metric, and show that in this metric the momentum equation has a new term that is associated with local curvature and density distribution profile. In particular conditions of steady state a new Einstein's field equation is determined in presence of negative curvature. Since GPE governs BECs defects that are useful, analogue models in cosmology, a relativistic form of GPE is also considered to show connection with models of analogue gravity, thus providing further grounds for future investigations of black hole dynamics in cosmology.

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• Right-hand side:

$$i\hbar\partial_t\Psi=\frac{i\hbar}{2\sqrt{\rho}}(\partial_t\rho)\,\mathrm{e}^{\mathrm{i}\left(\frac{m}{\hbar}\theta\right)}-m\sqrt{\rho}\,\mathrm{e}^{\mathrm{i}\left(\frac{m}{\hbar}\theta\right)}\partial_t\theta\,.$$

• Right-hand side:

$$i\hbar\partial_t\Psi=\frac{i\hbar}{2\sqrt{\rho}}(\partial_t\rho)e^{i(\frac{m}{\hbar}\theta)}-m\sqrt{\rho}e^{i(\frac{m}{\hbar}\theta)}\partial_t\theta.$$

• Left-hand side:

$$\boldsymbol{\nabla}^{2}\Psi = \frac{1}{\sqrt{|g|}}\partial_{i}\left[\sqrt{|g|}\,g^{ij}\partial_{j}\left(\sqrt{\rho}\,\mathrm{e}^{\mathrm{i}\left(\frac{m}{\hbar}\theta\right)}\right)\right]$$

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4

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$$= \left[\nabla^{2}\sqrt{\rho} + i\frac{m}{\hbar}\left(\frac{1}{\sqrt{\rho}}\nabla\rho\cdot\nabla\theta + \sqrt{\rho}\nabla\cdot(\nabla\theta)\right) - \frac{m^{2}}{\hbar^{2}}\sqrt{\rho}\left|\nabla\theta\right|^{2}\right]e^{i\left(\frac{m}{\hbar}\theta\right)}.$$

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• Equating the imaginary parts, we have:

$$\partial_t \rho = -\boldsymbol{\nabla} \cdot (\rho \boldsymbol{\nabla} \theta)$$
 . (C)

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• Equating the real parts:

$$\partial_t \theta + \frac{1}{2} |\nabla \theta|^2 = U + Q$$

(B)

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• Equating the imaginary parts, we have:

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• Equating the real parts: $\partial_t \theta + \frac{1}{2} |\nabla \theta|^2 = U + Q$ $\begin{cases}
U = -(\mathfrak{g}\rho + V)/m, \\
Q = (\hbar^2/2m^2)(\nabla^2\sqrt{\rho})/\sqrt{\rho}.
\end{cases}$ (B)

• Theorem (GPE in Euler form). GPE admits hydrodynamic description in the form of an Euler equation, given by

$$\partial_t \boldsymbol{u} + \boldsymbol{\nabla}_{\boldsymbol{u}} \boldsymbol{u} = \boldsymbol{\nabla} H$$
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where $\nabla_u u$ denotes the covariant derivative along itself and H = U + Q is the sum of classical and quantum potentials.

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• re-write covariant derivative using Koszul's formula, so that

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• Corollary. The relation between GPE and hydrodynamics is one-to-one on any manifold with generic metric g.

• Theorem (GPE in Navier-Stokes form). GPE admits hydrodynamic description in the form of a Navier-Stokes equation

$$\rho\left(\partial_t \boldsymbol{u} + \boldsymbol{\nabla}_{\boldsymbol{u}} \boldsymbol{u}\right) = -\boldsymbol{\nabla}\wp + \boldsymbol{\nabla} \cdot \boldsymbol{\tau} + \boldsymbol{\mathcal{E}}, \qquad (NS)$$

where V=0 , $\wp=(\mathfrak{g}/2m)\rho^2$, $\boldsymbol{\tau}=(\hbar^2/4m^2)
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and ${\cal E}$ is a density curvature tensor, ${\cal E}^j = -(\hbar^2/4m^2)R^{jk}\partial_k
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The second term on the r.h.s. can be written as

$$\frac{\hbar^2}{m^2}\rho \,\mathrm{d}\left(\frac{\boldsymbol{\nabla}^2\sqrt{\rho}}{2\sqrt{\rho}}\right) = \boldsymbol{\nabla}\cdot\left(\rho \operatorname{Hess}(\ln \rho)\right) + \tilde{\boldsymbol{\mathcal{E}}}^{\flat},$$

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where $\tilde{\boldsymbol{\mathcal{E}}}^{\flat} = (4m^2/\hbar^2)\mathcal{E}_j \,\mathrm{d}x^j$, $\tilde{\mathcal{E}}_j = g^{ik}R^r_{ikj}\partial_r\rho$ and $\boldsymbol{\mathcal{E}} = \frac{\hbar^2}{4m^2}\tilde{\boldsymbol{\mathcal{E}}}$.

 Theorem (Momentum conservation law). The momentum ρu associated with the hydrodynamic form of GPE satisfies the following conservation law

$$\partial_t(\rho \boldsymbol{u}) = -\boldsymbol{\nabla}\cdot\boldsymbol{\mathcal{M}} + \boldsymbol{\chi}$$
, (M)

where $\mathcal{M} = \mathcal{M}_{ii} dx^i \wedge dx^j$, $\mathcal{M}_{ij} = \mathcal{D}_{ij} + \prod_{ij} - \tau_{ij} - \mathcal{G}_{ij}$ with $\mathcal{D}_{ij} = \rho u_i u_j$, $u^{\flat} = u_j dx^j$, $\prod_{ij} = \wp g_{ij}$, $\mathcal{G} = \hbar^2 \rho / (4m^2) \mathcal{G}$ (G Einstein tensor) and χ is explicit function of the geometry of the manifold through the Ricci scalar R.

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Proof. • Use (C) and (NS) equations;

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Proof. • Use (C) and (NS) equations;

• compute divergence of terms;

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Proof. • Use (C) and (NS) equations;

- compute divergence of terms;
- use contracted form of Bianchi's identity

 $\nabla_k R_j^k \,\mathrm{d} x^j = (1/2) \partial_j R \,\mathrm{d} x^j$.

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 - compute divergence of terms;
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$$\nabla_k R^k_j \,\mathrm{d} x^j = (1/2) \partial_j R \,\mathrm{d} x^j$$
.

The new term $\chi = -\frac{\hbar^2}{8m^2} R d\rho$ breaks the standard form of a conservation law.

Thank you!