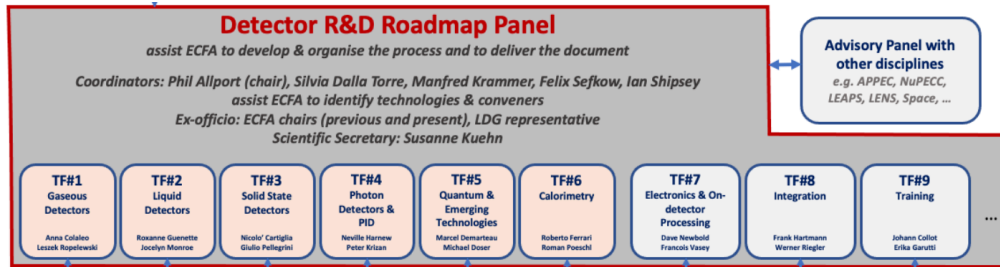


TF#5 - Quantum and Emerging Technologies



Convenors

Marcel Demarteau (ORNL)
Michael Doser (CERN)

Expert members

Caterina Braggio (Padova)
Andy Geraci (NWU)
Peter Graham (Stanford)
Anna Grasselino (Fermilab)
John March Russell (Oxford)
Stafford Withington (Cambridge)

Symposium date

Monday 12.4.2021

09 : 15 → 11 : 00 **Science targets – Overview and Landscape**

- Axions and other DM searches
- EDM searches and tests of fundamental symmetries
- Tests of Quantum Mechanics
- Multimessenger detection (atom interferometer networks and gravitational waves)

11 : 30 → 13 : 00 **Experimental methods and techniques - Overview and Landscape**

- Precision spectroscopy and clocks
- Novel ionic, atomic and molecular systems

13 : 30 → 16 : 00 **Experimental and technological challenges, New Developments**

- Superconducting platforms (detectors: TES, SNSPD, Haloscopes including single photon detection and squeezing)
- High sensitivity superconducting cryogenic electronics, low noise amplifiers including TWJPA
- Broadband axion detection
- Mechanical/optomechanical detectors
- Spin-based techniques

SEARCHES AT THE PRECISION FRONTIER OF PARTICLE PHYSICS

sensitive, lab-scale experiments
search for tiniest deviations
arising from **higher energy**
scales (/weaker couplings) than
can be reached at colliders

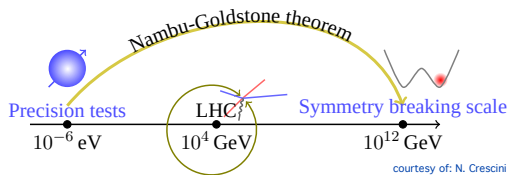
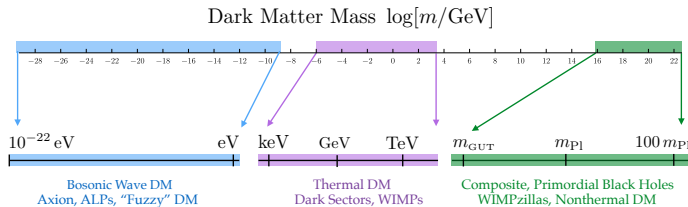


Figure credit: N. Crescini

quantum sensors
can reach
unprecedented
sensitivity and probe
ultralight DM

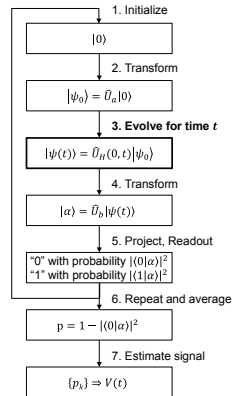


QUANTUM SENSING

"Quantum sensing" describes the use of a quantum system, quantum properties or quantum phenomena to perform a measurement of a physical quantity

Rev. Mod. Phys. 89, 035002 (2017)

1. Use of a **quantum object** to measure a physical quantity (classical or quantum). The quantum object is characterized by quantized energy levels, i.e. electronic, magnetic or vibrational states of superconducting or spin qubits, neutral atoms, or trapped ions.
2. Use of **quantum coherence** (i.e., wave-like spatial or temporal superposition states) to measure a physical quantity
3. Use of **quantum entanglement** to improve the sensitivity or precision of a measurement, beyond what is possible classically.



QUANTUM SENSING (IN THIS TALK)

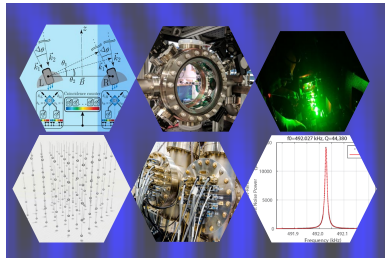
quantum sensing techniques and devices in quantum information science being applied to particle detection

Quantum Sensing for High Energy Physics

Report of the first workshop to identify approaches and techniques in the domain of quantum sensing that can be utilized by future High Energy Physics applications to further the scientific goals of High Energy Physics.

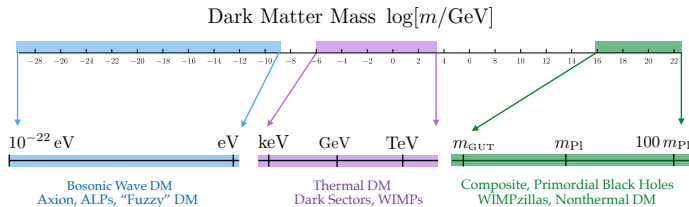
Organized by the Coordinating Panel for Advanced Detectors of the Division of Particles and Fields of the American Physical Society

March 27, 2018



arXiv.org > physics > arXiv:2102.10996
"Opportunities for DOE National Laboratory-led
QuantISED Experiments"
DOE program in High Energy Physics

ULTRALIGHT BOSONIC DARK MATTER

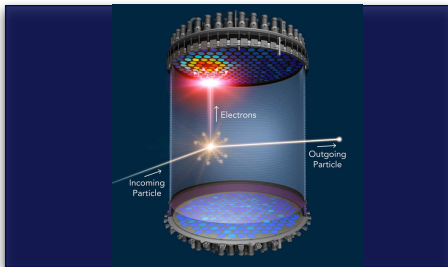


- BSM theories predict the existence of **ultralight bosons** which could be the dark matter
- $m_\chi \lesssim 1 \text{ eV} \iff$ **wavelike** dark matter
- wavelike dark matter can interact with Standard Model particles: photons, fermions, gluons, ...
- experiments target different portals, mass (/frequency) ranges
- blooming phase

AXION vs WIMP DETECTION

"The axion would be something of a spiritual cousin to the photon, but with just a hint of mass"

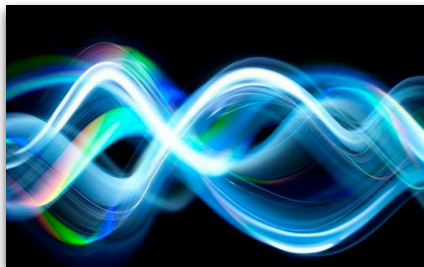
P. Sikivie



WIMP [4-1000] GeV

- number density is small
- tiny wavelength
- no detector-scale coherence

⇒ observable: **scattering of individual particles**

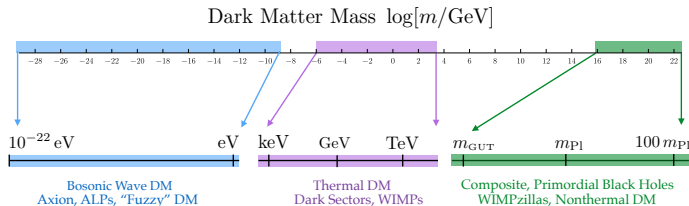


AXION [$m_A \lesssim \text{eV}$]

- number density is large (bosons)
- long wavelength
- coherence within detector

⇒ observable: **classical, oscillating, background field**

ULTRALIGHT BOSONIC DARK MATTER

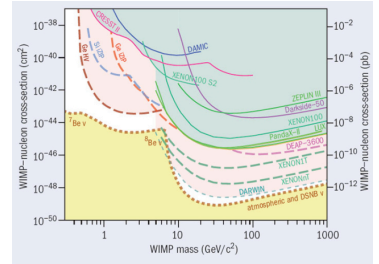


- BSM theories predict the existence of **ultralight bosons** which could be the dark matter
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- wavelike dark matter can interact with Standard Model particles: photons, fermions, gluons, ...
- experiments target different portals, mass (/frequency) ranges
- blooming phase!

IS DM MADE OF AXIONS?

⇒ a well motivated scenario:

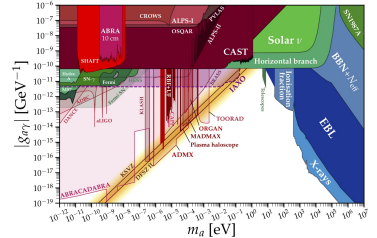
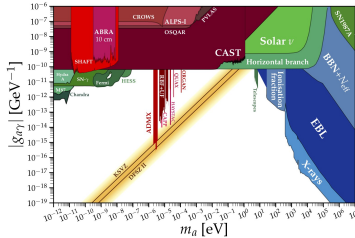
- “three birds with one particle”
 1. a CP problem solution
 2. Dark Matter candidate
 3. baryon asymmetry [PRL 124, 111602 (2020)]
- SUSY is failing tests at LHC
- WIMPs searches with next generation of experiments



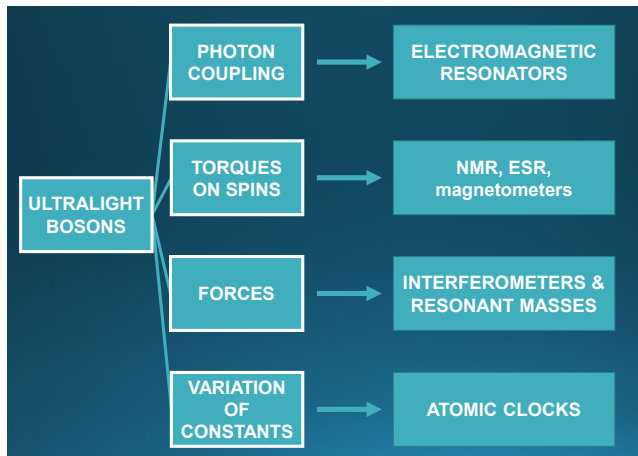
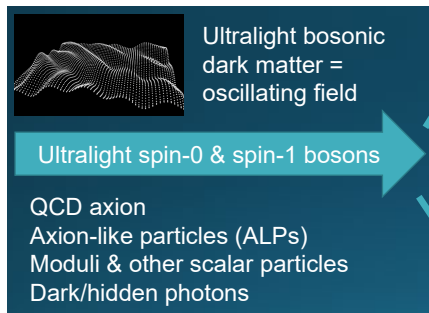
⇒ blooming phase:

- new detection concepts
- commissioning of small-scale setups
- upgrade to *large scale* experiments for established techniques

<https://cajohare.github.io/AxionLimits/>

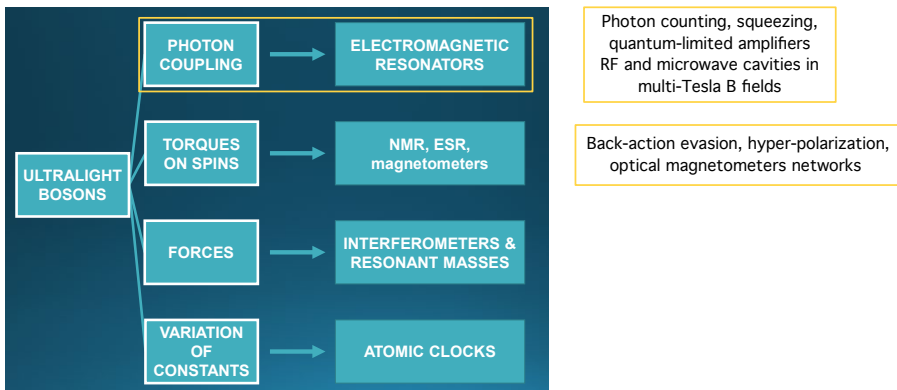


EXPERIMENTAL TECHNIQUES -TF5



from Derek F. Jackson Kimball “Quantum sensors for wavelike Dark Matter Searches” Snowmass 2021

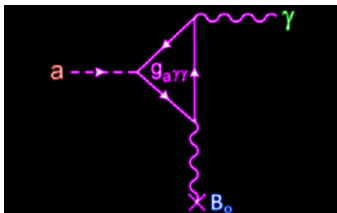
EXPERIMENTAL TECHNIQUES -TF5



from Derek F. Jackson Kimball “*Quantum sensors for wavelike Dark Matter Searches*” Snowmass 2021

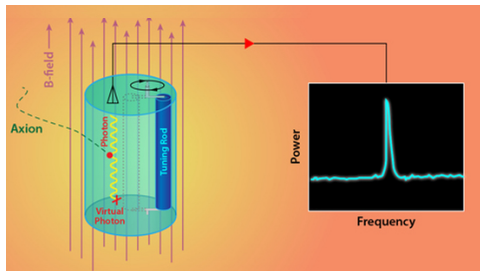
HALOSCOPE - resonant search for axion DM in the Galactic halo

- original proposal by P. Sikivie (1983)
- search for axions as cold dark matter constituent: SHM from Λ_{CDM} , local DM density ρ
 - signal is a **line** with 10^{-6} relative width in the energy(→ frequency) spectrum
 - + sharp (10^{-11}) components due to non-thermalized
- an **axion** may interact with a **strong \vec{B} field** to produce a **photon** of a specific frequency (→ m_a)



HALOSCOPE - resonant search for axion DM in the Galactic halo

– if axions are *almost monochromatic* then their conversion to detectable particles (photons) can be accomplished using *high-Q* microwave cavities.



$$\omega_{\text{TM}0nl} = \sqrt{\left(\frac{\epsilon_n}{r}\right)^2 + \left(\frac{l\pi}{h}\right)^2}$$

TM_{0nl} are the cavity modes that couple with the axion

- resonant amplification in $[m_a \pm m_a/Q]$
- data in thin slices of parameter space; typically $Q < Q_a \sim 1/\sigma_v^2 \sim 10^6$
- signal power $P_{a \rightarrow \gamma}$ is model-dependent

$$P_{a \rightarrow \gamma} \propto (B^2 V Q) (g_{a\gamma}^2 \frac{\rho}{m_a})$$

exceedingly tiny ($\sim 10^{-23}$ W)

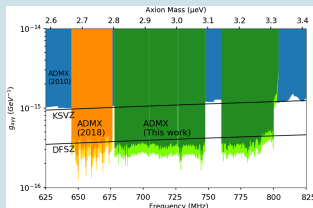
“The last signal ever received from the 7.5 W transmitter aboard Pioneer 10 in 2002, then 12.1 billion kilometers from Earth, was a prodigious 2.5×10^{-21} W. And unlike with the axion, physicists knew its frequency!”

K. V. Bibber and L. Rosenberg, Physics Today 59, 8, 30 (2006)

HALOSCOPES: 2020 RESULTS

ADMX

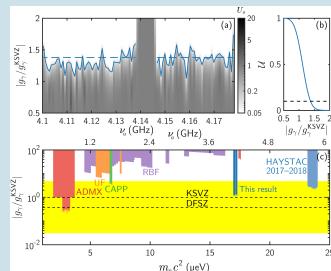
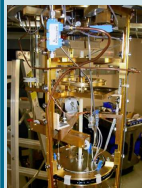
PRL **124**, 101303 (2020)



- sensitivity to QCD axion
- 2.82 – 3.31 μeV mass range
- cylindrical cavity at $T = 100$ mK,
 $Q_0 \sim 30000$, $V = 1361$
- 7.6 T field

HAYSTAC

Nature **590**, 238 (2021)



- 17.14 – 17.28 μeV axion mass range
- **squeezing** doubled the search rate
- cylindrical cavity at $T = 60$ mK
 $Q_0 \sim 50000$, $V = 1.51$
- 8T magnet

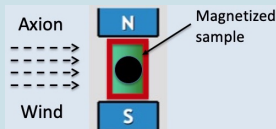
“Roberto and I spent a few months cooking up this theory, and now the experimentalists have spent 40 years looking for it”
H. Quinn

QUAX - QUAERERE AXIONS

Detection of cosmological axions through their **coupling to electrons** or **photons**

ELECTRON COUPLING – QUAX

the FMR haloscope

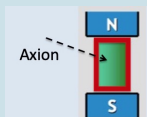


the axion DM cloud acts as an **effective RF magnetic field** on the electron spin exciting **magnetic transitions** in a magnetized sample (YIG) → **RF photons**

$$P_{\text{out}} = \frac{P_{\text{in}}}{2} = 8 \times 10^{-26} \left(\frac{m_a}{2 \cdot 10^{-4} \text{ eV}} \right)^3 \left(\frac{V_s}{1 \text{ liter}} \right) \left(\frac{n_S}{10^{28}/\text{m}^3} \right) \left(\frac{\tau_{\text{min}}}{10^{-6} \text{ s}} \right) \text{ W},$$

PHOTON COUPLING – QUAX $a\gamma$

high frequency range (8.5 – 11) GHz



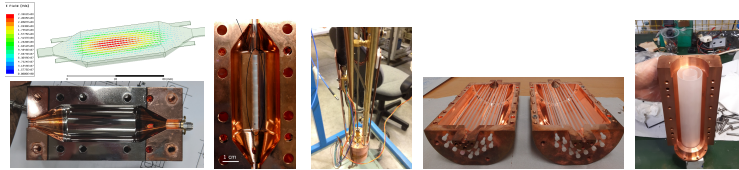
DM axions are converted into **RF photons** inside a **resonant cavity** immersed in a **strong magnetic field**

$$P_{\text{ax}} = 3.3 \cdot 10^{-24} \text{ W} \left(\frac{V_{\text{eff}}^{\text{Sa}}}{2.3 \cdot 10^{-5} \text{ m}^3} \right) \left(\frac{B}{8 \text{ T}} \right)^2 \times \left(\frac{g_\gamma}{-0.97} \right)^2 \left(\frac{\rho_a}{0.45 \text{ GeV cm}^{-3}} \right) \left(\frac{f}{13.5 \text{ GHz}} \right) \left(\frac{Q_L}{145\,000} \right)$$

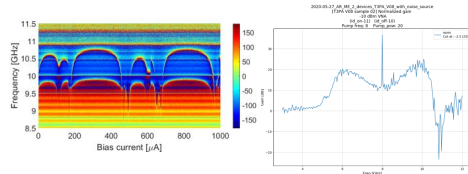
QUANTUM SENSING

To get to **cosmologically relevant sensitivity**:

- develop **high Q** ($\gtrsim 10^6$) microwave cavities, compatible with **strong magnetic fields** ($B \gtrsim 8$ T)

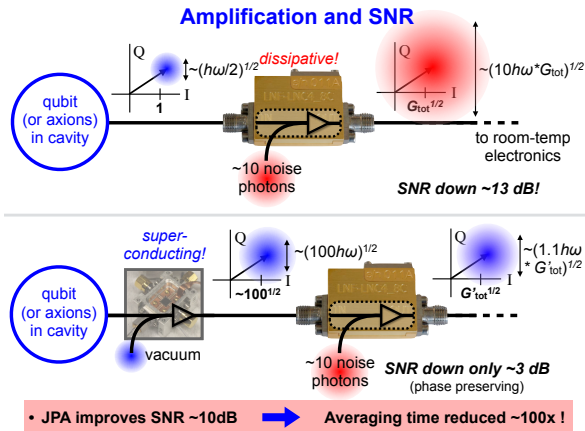
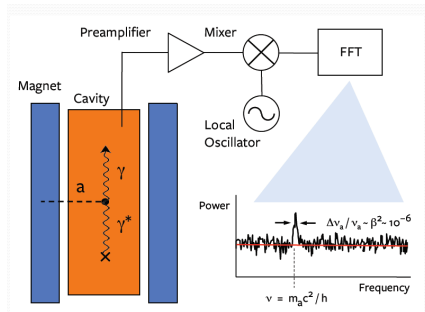


- use **quantum-limited amplifiers**
JPA (Josephson Parametric Amplifier),
TW (Travelling Wave) JPA (collaboration with N. Roche, Grenoble) → DARTWARS (call gr5)



FROM HEMT TO JPA

the Dicke's receiver noise (\rightarrow sensitivity to ultralight DM) is determined by the preamp



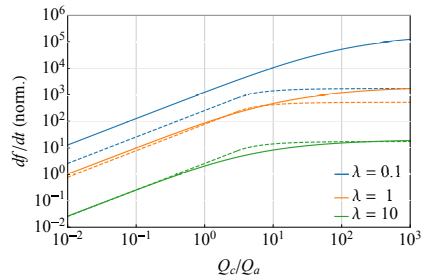
from A. Eddins "Josephson Parametric Amplifiers: Theory and Application"

SCAN RATE

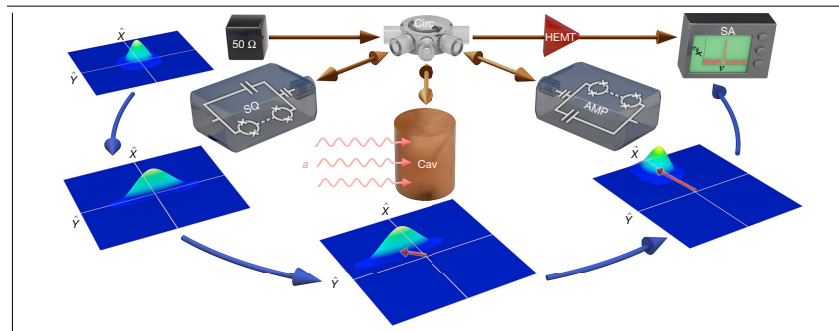
Revisiting the detection rate for axion haloscopes D. Kim et al JCAP03, 066 (2020)

Abstract. The cavity haloscope has been employed to detect microwave photons resonantly converted from invisible cosmic axions under a strong magnetic field. In this scheme, the axion-photon conversion power has been formulated to be valid for certain conditions, either $Q_{\text{cavity}} \ll Q_{\text{axion}}$ or $Q_{\text{cavity}} \gg Q_{\text{axion}}$. This remedy, however, fails when these two quantities are comparable to each other. Furthermore, the noise power flow has been treated independently of the impedance mismatch of the system, which could give rise to misleading estimates of the experimental sensitivity. We revisit the analytical approaches to derive a general description of the signal and noise power. We also optimize the coupling strength of a receiver to yield the maximal sensitivity for axion search experiments.

$$\frac{df}{dt} = \frac{1}{\text{SNR}^2} \left(\frac{P_0}{k_B T_{\text{eff}}} \right)^2 \left(\frac{\frac{\beta}{(1+\beta)}}{\frac{4\beta}{(1+\beta)^2} + \lambda} \right)^2 \frac{Q_l Q_a^2}{Q_l + Q_a}$$

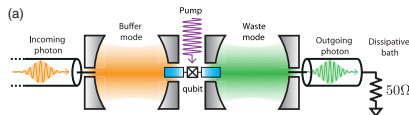


(BLUE LINE): SQUEEZING



QUANTUM SENSING

- (next step in sensitivity*) we need to be able to **count photons** \iff develop the **microwave SPD**
NOTE: photon detection at optical frequencies relies on ionization, microwave photons are 5 orders of magnitude lower!



PHYSICAL REVIEW X **10**, 021038 (2020)

Quantronics Group, FR (E. Flurin)

qubit frequency \iff number of photons in the cavity

The electric field of even a single photon will excite the non-linear response of the qubit oscillator and shift its frequency

JPA readout vs counter

Example:

linear amplifier σ^{LA} , and quantum counter σ^{QC} power sensitivity

$$T_N^{LA} = 1 \text{ K},$$

$$QC \text{ noise rate } R_n = 100 \text{ Hz},$$

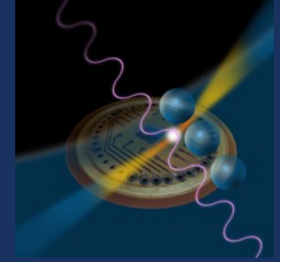
$$f = 10 \text{ GHz},$$

$$Q_a = 10^6$$

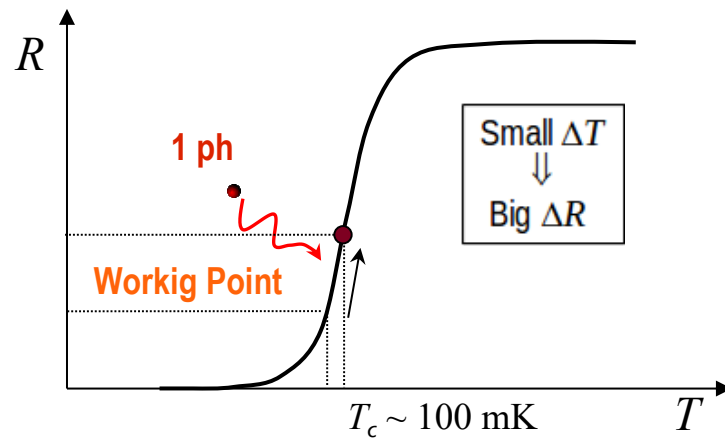
\rightarrow sensitivity improved by a factor 14 \iff 200 in measurement time!

$$\frac{\sigma^{QC}}{\sigma^{LA}} = \frac{h}{k_B T_n^{LA}} \sqrt{f Q_a R_n}$$

SIMP: SINGLE MICROWAVE PHOTON DETECTION

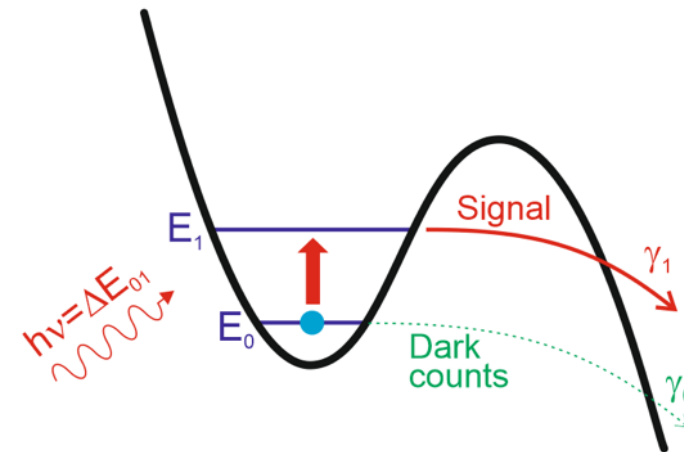


- 3 years project funded by CSNV of INFN
- Goal: development of single microwave photon detectors with nanowire-TES and Current Biased Josephson Junction (CBJJ)
- In collaboration with CNR-IFN, CNR-Nano and INRiM



Nanowire TES

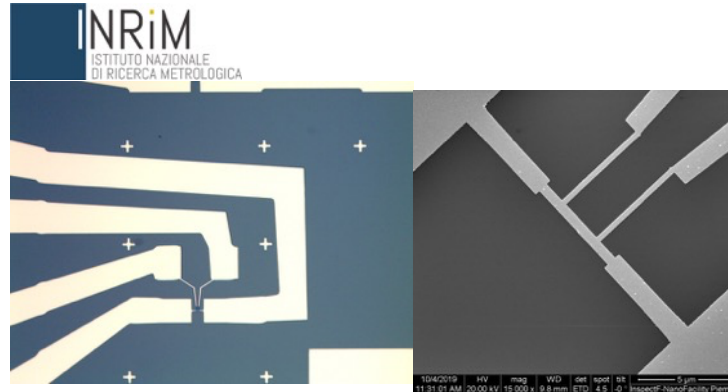
Small volume and low critical temperature allow steep variation of resistance when a single photon of 50-100 GHz is absorbed.



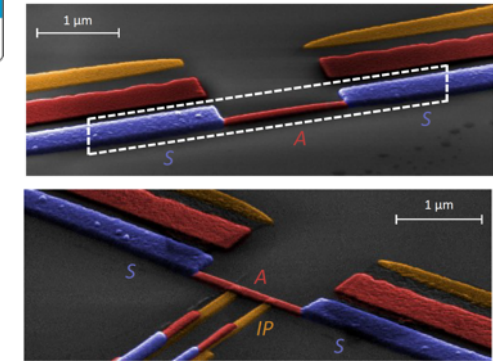
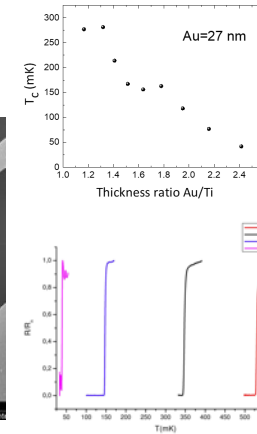
CBJJ

The absorption of a single 10 GHz photon causes the transition of the JJ from the superconducting to the resistive state.

1. Nanowire TES fabricated at INRiM (Ti/Au) and CNR-Nano (Al/Cu) with typical size 20 nm x 100 nm x 1000 nm and $T_c \sim 100$ mK.

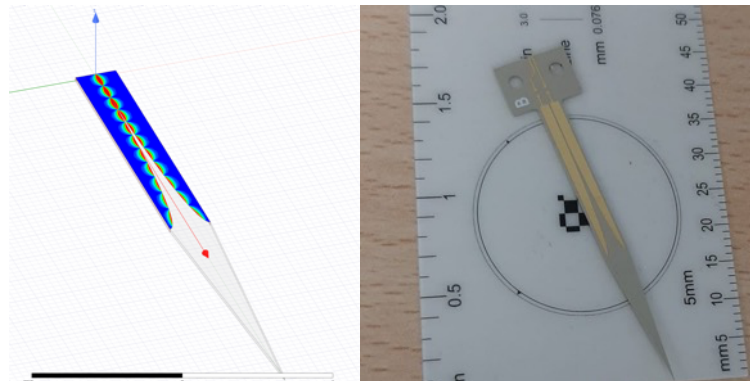


Transition measured with nano-TES coupled to DC-SQUID at INRiM

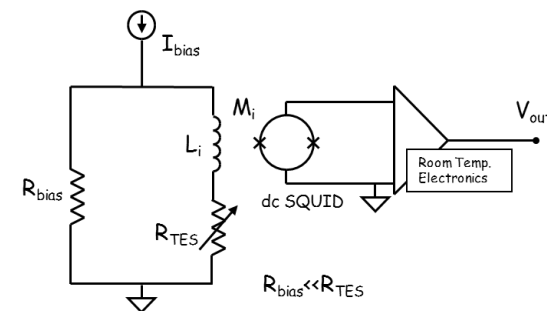


Full DC characterization in F.Paolucci et al arXiv:2007.08320

2. Waveguide photon collected with a finline terminated with the nano-TES.



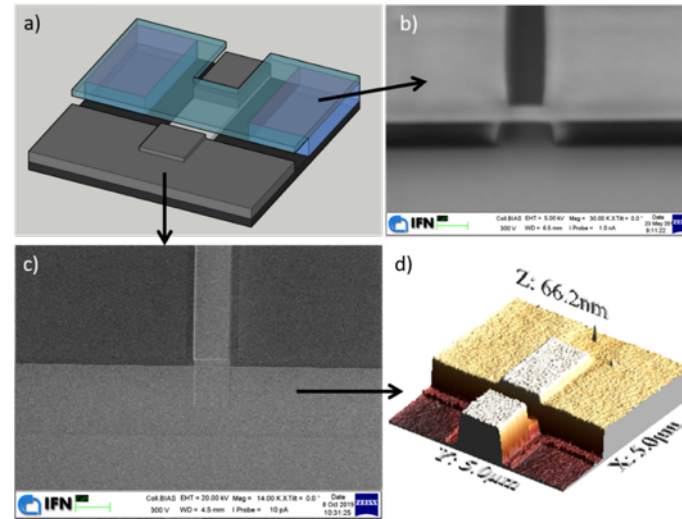
3. DC-SQUID readout and RF tests in preparation in the TIFPA (Trento) dilution refrigerator.



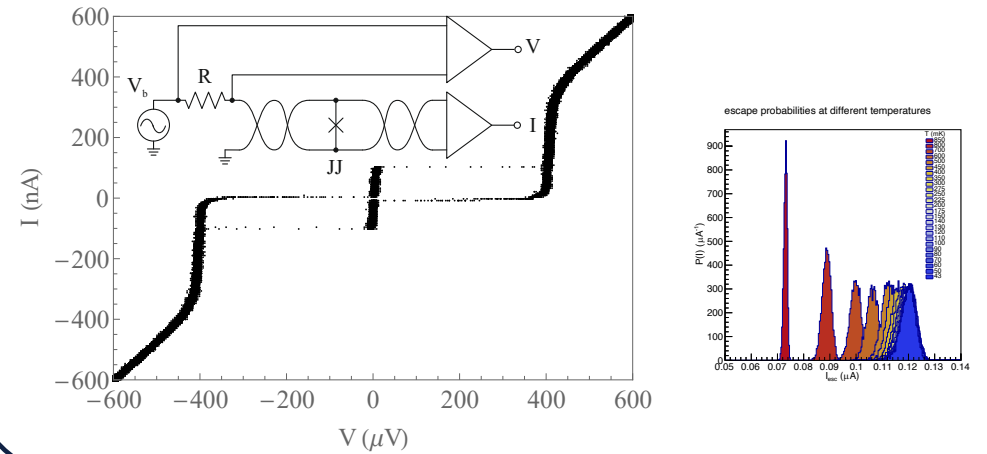
TIFPA Dilution Refrigerator



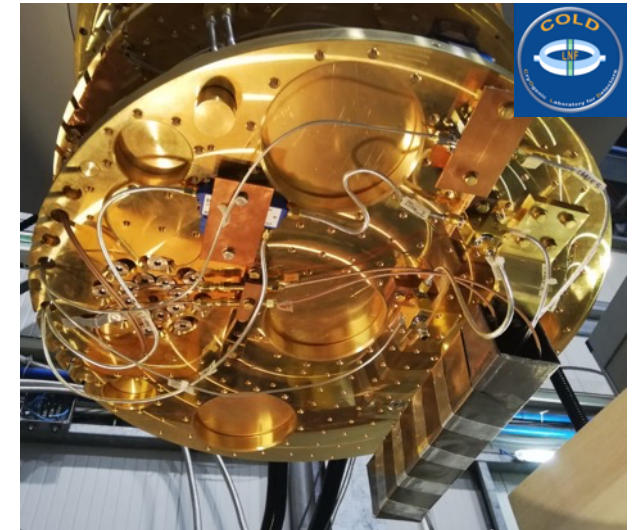
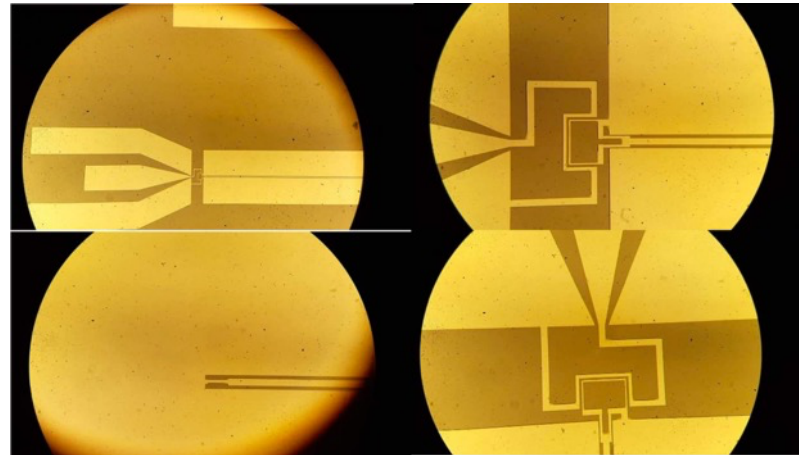
1. JJ fabricated with EBL and Al shadow evaporation technique at CNR-IFN



2. JJ characterization with IV measurements in a dilution refrigerators at 40 mK

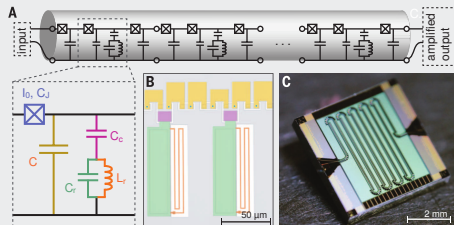


3. RF tests ongoing at LNF on chip with transmission lines terminated on CBJJ or DC-SQUID. Devices tested also as JPA. $T=10$ mK.



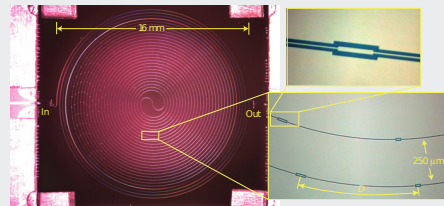
- P Falferi et al. «Status of the SIMP Project: Toward the Single Microwave Photon Detection» J Low Temp Phys (LTD proceedings) DOI 10.1007/s10909-020-02381-x
- S. Pagano et al. «Development of a Josephson junction based single photon microwave detector for axion detection experiments», J. Phys.: Conf. Ser. 1559 012020 (EUCAS proceedings)
- G. Filatrella et al. «Analysis of thermal and quantum escape times of Josephson junctions for signal detection» Chaotic Modeling and Simulation CMSIM An International Journal of Nonlinear Science.
- C. Barone et al, Analysis of Josephson junctions switching time distributions for the optimal detection of single microwave photons, Chaos, Solitons and Fractals 142 (2021) 110496.
- F. Paolucci et al. «Highly sensitive nano-TEs for gigahertz astronomy and dark matter search» J.App. Phys. DOI:10.1063/5.0021996 arXiv:2007.08320
- F. Paolucci et al Phys Rev Appl 14, 034055 (2020) Hypersensitive Tunable Josephson Escape sensor for Gigahertz astronomy

Traveling Wave Josephson Parametric Amplifiers



- TWJPA implemented as a nonlinear lumped-element transmission line;
- One unit cell consists of a Josephson junction and a capacitive shunt to ground;
- **Already demonstrated quantum-limited noise read out;**
- **High bandwidth: over a 4 GHz centred at 5 GHz;**
- **Limited gain: < 20 dB**
- **Small dynamic range: < -90 dBm**

Dispersion-engineered Traveling Wave Kinetic Inductance



- DTWKI exploits the Nonlinear kinetic inductance of TiN and NbTiN superconducting film;
- Phase matching by Coplanar Waveguide (CPW) or lumped-element artificial transmission line;
- **Noise near to quantum-limited**
- **High bandwidth: over a 4 GHz centred at 5 GHz;**
- **Limited gain: < 20 dB, and gain profile with large ripple**
- **High dynamic range: from -50 to -45 dBm**

The principal objectives of DART WARS are

1. The practical development of high performing parametric amplifiers following two different promising approaches (DTWKI, TWJPA) and exploring new design solutions, new materials and advanced fabrication processes;
2. The read out demonstration of various detectors/components (TESs, MKIDs, microwave cavities and qubits) with improved performances due to a parametric amplification with a noise at the quantum level;

The technical goal is to achieve

- a gain value around 20 dB, comparable to HEMT;
- a high saturation power (around -50 dBm);
- quantum limited or nearly quantum limited noise (noise temperature < 600 mK)
- reduction of the gain ripple and yield improvement.

These features will lead to reading out large arrays of detectors or qubits with no noise degradation by reaching unprecedented energy resolution, noise equivalent power, energy threshold, opening future prospective toward next generation of particle, astroparticle and astrophysics experiments.