# EXTREMELY METAL-POOR Asymptotic Giant Branch Stars 

THE 13TH TORINO WORKSHOP ON AGB Stars \& the 3rd Perugia Workshop ON NUCLEAR ASTROPHYSICS
DSA3, University of Perugia, Perugia
June 19-24, 2022


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## 1

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## But...

## Population III stars

- Not directly observed yet, but many observational constraints on their existence:
a)detection of metals in Ly- $\alpha$ forest spectra of distant QSOs
b)detection of metals in damped Lyman systems
c) abundance ratios in extremely metal-poor (EMP) stars ([Fe/H] $<-3$ )


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It is crucial to study the evolution and nucleosynthesis of EMP stars!

## My work

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Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic

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- At the subsequent TDU, all the surface main CNO abundances are raised from one to three orders of magnitude
- But what happens inside the star

Hydrogen luminosity


Surface abundances

 during a PIE?

## AGB and PIE



## AGB and PIE

- The star has now a double shell structure, with the two shells advancing alternatively



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from Lattanzio (2003)


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$\mathrm{M}=2 \mathrm{M}_{\odot}, \mathrm{Z}=10^{-5}$, PIE: CNO
Before the PIE

PIE at maximum luminosity


Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic Giant Branch Stars, Universe 2022, 8, 44
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- c)CNO cycle activated in the second convective zone


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- c)CNO cycle activated in the second convective zone
- d)penetration of the convective envelope

PIE at maximum luminosity

TDU at maximum penetration


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Before the PIE

PIE at maximum luminosity

TDU at maximum penetration



## $\mathrm{M}=2 \mathrm{M}_{\odot}, \mathrm{Z}=10^{-5}$, PIE: ${ }^{7} \mathrm{Li}$

Before the PIE
a) ${ }^{11} \mathrm{~B}$ is produced via ${ }^{7} \mathrm{Li}(\alpha, \gamma){ }^{11} \mathrm{~B}$ and ${ }^{7} \mathrm{Be}(\alpha, \gamma){ }^{11} \mathrm{C}\left(\beta^{+}, v\right){ }^{11} \mathrm{~B}$

PIE at maximum luminosity

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b) ${ }^{7} \mathrm{Be}$ is produced via ${ }^{3} \mathrm{He}(\alpha, \gamma)^{7} \mathrm{Be}$

PIE at maximum luminosity

TDU at maximum penetration



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- b) ${ }^{7} \mathrm{Be}$ is produced via ${ }^{3} \mathrm{He}(\alpha, \gamma)^{7} \mathrm{Be}$ c) Production of ${ }^{7} \mathrm{Li}$ via ${ }^{7} \mathrm{Be}\left(\mathrm{e}^{-}, v\right)^{7} \mathrm{Li}$ ( $\mathrm{T} \lesssim 20 \mathrm{MK}$ ) luminosity

TDU at maximum penetration



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- b) ${ }^{7} \mathrm{Be}$ is produced via ${ }^{3} \mathrm{He}(\alpha, \gamma)^{7} \mathrm{Be}$
- c) Production of ${ }^{7} \mathrm{Li}$ via ${ }^{7} \mathrm{Be}\left(\mathrm{e}^{-}, v\right)^{7} \mathrm{Li}$ ( $\mathrm{T} \lesssim 20 \mathrm{MK}$ )


$\mathrm{M}=6 \mathrm{M}_{\odot}$, different Z
Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic ${ }_{50}$ Giant Branch Stars, Universe 2022, 8, 44

- Progressive reduction of core masses at each convective episode as Z decreases (more compact structures)
$\mathrm{M}=6 \mathrm{M}_{\odot}$, different Z
Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic ${ }_{51}$ Giant Branch Stars, Universe 2022, 8, 44

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- Two convective episodes (CNO cycle and He burning) in $\mathrm{Z}=10^{-4}$ and $\mathrm{Z}=10^{-6}$
$\mathrm{M}=6 \mathrm{M}_{\odot}$, different Z
Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic ${ }_{52}$ Giant Branch Stars, Universe 2022, 8, 44

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$\mathrm{M}=6 \mathrm{M}_{\odot}$, different Z
Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic 53 Giant Branch Stars, Universe 2022, 8, 44

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- Two convective episodes (CNO cycle and He burning) in $\mathrm{Z}=10^{-4}$ and $\mathrm{Z}=10^{-6}$
- Three convective episodes (CNO cycle, CNO cycle $+3 \alpha$ reactions and He burning) in $\mathrm{Z}=10^{-10}$
- The star contracts until it reaches $\mathrm{T} \approx 10^{8} \mathrm{~K}$ and $3 \alpha$ reactions take place producing ${ }^{12} \mathrm{C}$
$\mathrm{M}=6 \mathrm{M}_{\odot}$, different Z
Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic ${ }_{54}$ Giant Branch Stars, Universe 2022, 8, 44

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Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic ${ }_{56}$

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- When CNO central abundance reaches $\approx 2 \times 10^{-10}$, CNO cycle is activated and a second convective episode appears $\Rightarrow \mathrm{CNO}+3 \alpha$
- This leaves a footprint in the HR diagram
$\mathrm{M}=6 \mathrm{M}_{\odot}$, different Z
Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic ${ }_{57}$ Giant Branch Stars, Universe 2022, 8, 44

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- This leaves a footprint in the HR diagram
- Metal-poor stars are hotter than metal-rich stars $\Rightarrow$ He burning in the blue part of the HR diagram


## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}:$ SDU

# $\mathbf{M}=6 \mathbf{M}_{\odot}$, different $7: S D T$ 

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Straniero, O. Extremely Metal-
Poor Asymptotic Giant Branch Stars, Universe 2022, 8, 44

|  | $Z=\mathbf{1 0}^{-4}$ |  | $Z=\mathbf{1 0}^{-\mathbf{6}}$ |  | $Z=\mathbf{1 0}^{-\mathbf{1 0}}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Before | After | Before | After | Before | After |
| H | $7.50 \times 10^{-1}$ | $6.51 \times 10^{-1}$ | $7.50 \times 10^{-1}$ | $6.50 \times 10^{-1}$ | $7.50 \times 10^{-1}$ | $6.14 \times 10^{-1}$ |
| ${ }^{4} \mathrm{He}$ | $2.50 \times 10^{-1}$ | $3.49 \times 10^{-1}$ | $2.50 \times 10^{-1}$ | $3.50 \times 10^{-1}$ | $2.50 \times 10^{-1}$ | $3.86 \times 10^{-1}$ |
| ${ }^{12} \mathrm{C}$ | $1.74 \times 10^{-5}$ | $8.21 \times 10^{-6}$ | $1.74 \times 10^{-7}$ | $5.35 \times 10^{-7}$ | $1.74 \times 10^{-11}$ | $8.78 \times 10^{-8}$ |
| ${ }^{13} \mathrm{C}$ | $1.97 \times 10^{-7}$ | $3.66 \times 10^{-7}$ | $1.97 \times 10^{-9}$ | $2.58 \times 10^{-9}$ | $1.97 \times 10^{-13}$ | $1.29 \times 10^{-12}$ |
| ${ }^{14} \mathrm{~N}$ | $4.90 \times 10^{-6}$ | $2.69 \times 10^{-5}$ | $4.90 \times 10^{-8}$ | $3.87 \times 10^{-7}$ | $4.90 \times 10^{-12}$ | $1.45 \times 10^{-9}$ |
| ${ }^{16} \mathrm{O}$ | $4.25 \times 10^{-5}$ | $3.19 \times 10^{-5}$ | $4.25 \times 10^{-7}$ | $2.19 \times 10^{-7}$ | $4.25 \times 10^{-11}$ | $1.81 \times 10^{-10}$ |
| CNO | $4.46 \times 10^{-6}$ | $4.60 \times 10^{-6}$ | $4.46 \times 10^{-8}$ | $8.59 \times 10^{-8}$ | $4.46 \times 10^{-12}$ | $7.43 \times 10^{-9}$ |

# $M=6 M_{\odot}$, different $Z: S D U$ 

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Straniero, O. Extremely MetalPoor Asymptotic Giant Branch Stars, Universe 2022, 8, 44

|  | $\boldsymbol{Z}=\mathbf{1 0}^{-4}$ |  | $\boldsymbol{Z}=\mathbf{1 0}^{-\mathbf{6}}$ |  | $\boldsymbol{Z}=\mathbf{1 0}^{-\mathbf{1 0}}$ |  |
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- In the $\mathrm{Z}=10^{-4}$ model, the surface abundances of ${ }^{12} \mathrm{C}$ and ${ }^{16} \mathrm{O}$ decrease, while those of ${ }^{13} \mathrm{C}$ and ${ }^{14} \mathrm{~N}$ increase, but the total number of $\mathrm{C}+\mathrm{N}+\mathrm{O}$ nuclei is almost conserved


# $\mathbf{M}=6 \mathbf{M}_{\odot}$, different $7: S D T$ 

Cirillo, M.; Piersanti, L.;
Straniero, O. Extremely Metal-
Poor Asymptotic Giant Branch Stars, Universe 2022, 8, 44

|  | $\boldsymbol{Z}=\mathbf{1 0}^{-\mathbf{4}}$ |  |  | $\boldsymbol{Z = \mathbf { 1 0 } ^ { - \mathbf { 6 } }}$ |  | $\boldsymbol{Z}=\mathbf{1 0}^{\mathbf{- 1 0}}$ |  |
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|  | Before | After | Before | After | Before | After |  |
| H | $7.50 \times 10^{-1}$ | $6.51 \times 10^{-1}$ | $7.50 \times 10^{-1}$ | $6.50 \times 10^{-1}$ | $7.50 \times 10^{-1}$ | $6.14 \times 10^{-1}$ |  |
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- In the $\mathrm{Z}=10^{-6}$ model, the surface abundance of ${ }^{12} \mathrm{C}$ increases after the SDU


# $\mathbf{M I}^{\circ}=6 \mathbf{M}_{\odot}$ different $7: S D T$ 

Cirillo, M.; Piersanti, L.;
Straniero, O. Extremely Metal-
Poor Asymptotic Giant Branch
Stars, Universe 2022, 8, 44

|  | $\boldsymbol{Z}=\mathbf{1 0}^{-\mathbf{4}}$ |  |  | $\boldsymbol{Z = \mathbf { 1 0 } ^ { - \mathbf { 6 } }}$ |  | $\boldsymbol{Z}=\mathbf{1 0}^{\mathbf{- 1 0}}$ |  |
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- In the $\mathrm{Z}=10^{-6}$ model, the surface abundance of ${ }^{12} \mathrm{C}$ increases after the SDU $\Rightarrow$ new phenomenon!


# $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}: \mathrm{SDU}$ 

Cirillo, M.; Piersanti, L.;
Straniero, O. Extremely MetalPoor Asymptotic Giant Branch Stars, Universe 2022, 8, 44

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- Enhancement of $\mathrm{C}+\mathrm{N}+\mathrm{O}$ in the envelope


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Cirillo, M.; Piersanti, L.;
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- This is a consequence of the coexistence of core-H and He burning
- Enhancement of $\mathrm{C}+\mathrm{N}+\mathrm{O}$ in the envelope $\longrightarrow$ efficiency of the shell-H burning increases


## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}:$ TP-AGB



Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-

## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}: \mathrm{TP}-\mathrm{AGB}$

## - H-flash or PIE



Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-

## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}:$ TP-AGB

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- The occurrence of the first PIE leads to stronger TPs (helium luminosity)


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- ${ }^{12} \mathrm{C}$ and ${ }^{16} \mathrm{O} \rightarrow{ }^{14} \mathrm{~N}$


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Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-

## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}: \mathrm{TP}-\mathrm{AGB}$

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- Weaker TPs $\square$ PIEs and TDUs cease $\longrightarrow$ new phenomenon!
- Shell-H burning tends to become stationary $\longrightarrow$ new phenomenon!
- Moderate HBB and CNO equilibrium


Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-

## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}: \mathrm{TP}-\mathrm{AGB}$

$$
\mathrm{Z}=10^{-10}
$$



## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}:$ TP-AGB

## - Weaker TP



Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-

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## - Weaker TP $\square$ no PIEs!

- Lower T at the base of the convective envelope



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- Lower T at the base of the convective envelope $\longrightarrow$ marginal activation of HBB



## $\mathrm{M}=6 \mathrm{M}_{\odot}$, different $\mathrm{Z}: \mathrm{TP}-\mathrm{AGB}$

- Weaker TP $\Rightarrow$ no PIEs!
- Lower T at the base of the convective envelope $\longrightarrow$ marginal activation of HBB
- Shell-H burning tends to become stationary as in the previous case


Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-

## Conclusions and ongoing activity

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- One model with $\mathrm{M}=2 \mathrm{M}_{\odot}$ and $\mathrm{Z}=10^{-5}$
- Three models with $\mathrm{M}=6 \mathrm{M}$ and $\mathrm{Z}=10^{-4}, 10^{-6}$ and $10^{-10}$
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- PIE
- Huge production of neutrons ( $\mathrm{n}_{\mathrm{n}} \gtrsim 10^{14}$ $\mathrm{cm}^{-3}$ )
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$\mathrm{M}=6 \mathrm{M}_{\odot}$ and $\mathrm{Z}=10^{-4}, 10^{-6}$ and $10^{-10}$
- Anomalous core-H and He burning as Z decreases
- Core-He burning at higher T as Z decreases
- Anomalous SDU as Z decreases
- Anomalous behaviour during the TPAGB phase


## Conclusions and ongoing activity

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- Anomalous behaviour during the TPAGB phase
- Ongoing activity: full grid of models of EMP stars with $1 \mathrm{M}_{\odot} \leq \mathrm{M} \leq 20 \mathrm{M}_{\odot}$ and $3.1 \times 10^{-4} \leq \mathrm{Z} \leq 10^{-10}$ in order to investigate their main properties and peculiar phenomena like Proton Ingestion Episode


# THANKS FOR YOUR ATTENTION! 

Cirillo, M.; Piersanti, L.; Straniero, O. Extremely Metal-Poor Asymptotic Giant Branch Stars, Universe 2022, 8, 44

MARIO CIRILLO<br>Università degli studi di Roma "Tor Vergata"<br>Istituto Nazionale di Astrofisica -<br>OSSERVATORIO ASTRONOMICO D’ABRUZZO

Equation of state

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- Important because it influences the extension of the convective zones in AGB (SDU, TDU and HBB)


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- In metal-poor stars convection is able to mix regions with different chemical compositions, leading to thermal runaways and thus reducing the characteristic timescales with respect to more metal-rich stars
- The reduction of the burning timescales influences the main physical quantities of the star $\longrightarrow$ coupling physics, burning and mixing affects the evolution of the whole structure


## Equation of state

- Two different equations of state adopted


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- Two different equations of state adopted $\log \mathrm{T}=6.5$

| EOS | Main properties | Temperature |
| :---: | :---: | :---: |
| range |  |  |

Straniero 2)Deviations from perfect gas (electron degeneracy, pair
and production, relativistic effects and Coulomb interactions) taken into
Prada account
$6<\log \mathrm{T}<10$
Moroni 3)Ideal for advanced burning phases and high temperatures and densities
1)Partially ionized matter
2)More accurate than Saha equation of state because of the
Opal treatment of all excited states, taking into account many-body $3.3<\log \mathrm{T}<8.3$ effects and Coulomb interactions
3)Ideal for atmospheric layers and low temperatures and densities

