

# Status of the MUonE project

Giovanni Abbiendi  
(INFN Bologna)

# Muon anomalous magnetic moment

$$\vec{M}_l = g_l \frac{e}{2m_l} \vec{S}$$

Dirac eq :  $g_l = 2$

Quantum corrections  $\rightarrow$  the anomaly

$$a_l \equiv \frac{g_l - 2}{2}$$

This observable can be both precisely measured experimentally and predicted in the Standard Model, providing a stringent test of the SM.

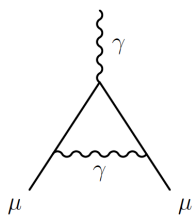
E821 experiment at BNL :  
[Phys.Rev.D73 (2006) 072003]

$$a_\mu^{\text{E821}} = 11659209.1(5.4)(3.3) \times 10^{-10}$$

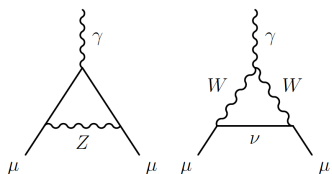
0.54 ppm

Dominated by statistics

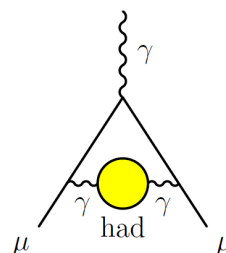
$$a_\mu^{\text{SM}} = a_\mu^{\text{QED}} + a_\mu^{\text{EWK}} + a_\mu^{\text{had}}$$



QED corrections known up to 5 loops, rel. precision  $\sim 7 \times 10^{-10}$   
LO term (Schwinger) =  $\alpha/2\pi \sim 0.00116$



EWK corrections  $\sim 10^{-9}$   
rel. uncertainty  $< 1\%$



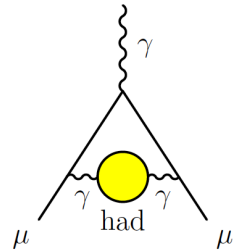
**Hadronic contribution  $\sim 7 \times 10^{-8}$**   
**-not calculable by pQCD-**

**Main contribution: LO Vacuum Polarization**  
estimated rel. uncertainty 0.6%

**➡ Dominant Theoretical uncertainty**

# $a_\mu^{\text{HLO}}$ : standard data-driven approach (time-like)

Dispersion relations, optical theorem:



$$a_\mu^{\text{HLO}} = \left( \frac{\alpha m_\mu}{3\pi} \right)^2 \int_{4m_\pi^2}^{\infty} ds \frac{\hat{K}(s) R_{\text{had}}(s)}{s^2}$$

$$R_{\text{had}}(s) = \sigma(e^+e^- \rightarrow \text{hadrons}) / \sigma(e^+e^- \rightarrow \mu^+\mu^-)$$

$K$  smooth function

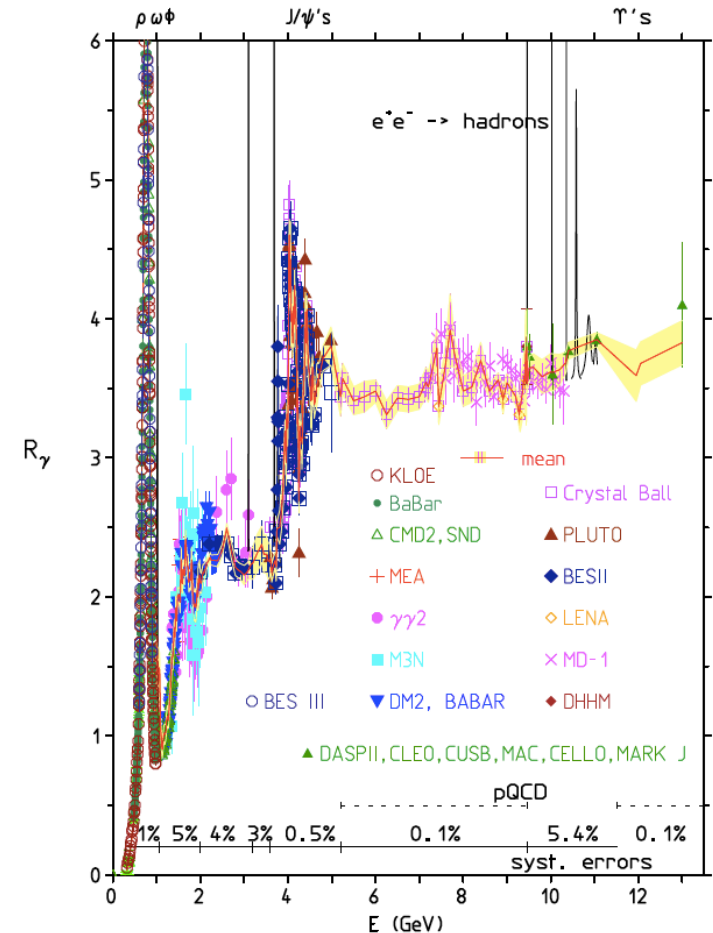
*Traditionally the integral is calculated by using the experimental measurements up to an energy cutoff, beyond which perturbative QCD can be applied.*

Main contribution: low-energy region ( $1/s^2$  enhancement), highly fluctuating due to hadron resonances and thresholds effects

**Alternative:** Lattice QCD calculations

continuously progressing, expected to become more and more competitive in the near future

F.Jegerlehner, EPJ Web Conf. 118 (2016) 01016



# $a_\mu$ measurement versus SM

Status report:

*T. Aoyama et al., Phys.Rept.887(2020)1*

$$a_\mu^{\text{HVP,LO}} = 693.1(4.0) \times 10^{-10}$$

**3.7  $\sigma$  discrepancy: new physics ?**

Recent lattice (still unpublished)

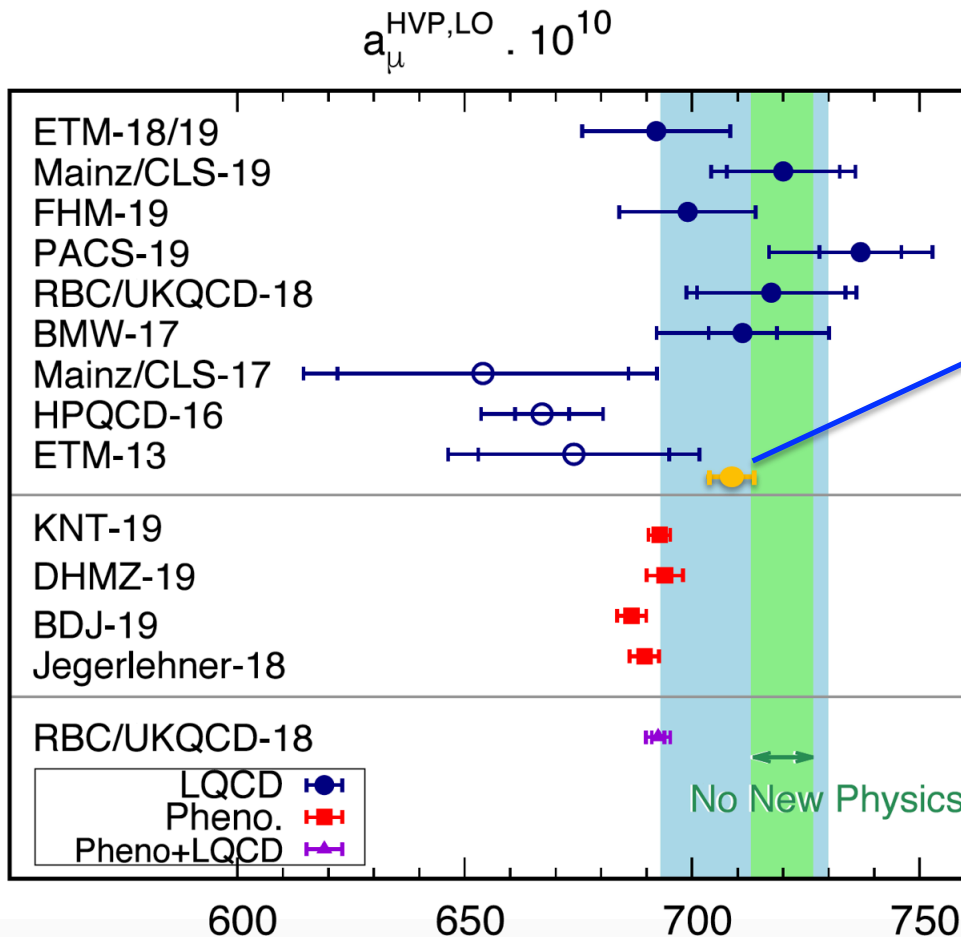
BMW20 arXiv:2002.12347

$$a_\mu^{\text{HVP,LO}} = 708.7(5.3) \times 10^{-10}$$

New g-2 experiment on-going at Fermilab aims at a **reduction of the experimental error by a factor of 4**

**Theory error** dominated by the **LO Hadronic contribution to the Vacuum Polarization: 0.6%**

*should be improved*



**MUonE experiment proposal: independent method, competitive precision**



# New measurement of $\alpha$ at 81ppt

Nature 588 (2020) 61

## Article

# Determination of the fine-structure constant with an accuracy of 81 parts per trillion

<https://doi.org/10.1038/s41586-020-2964-7>

Received: 7 May 2020

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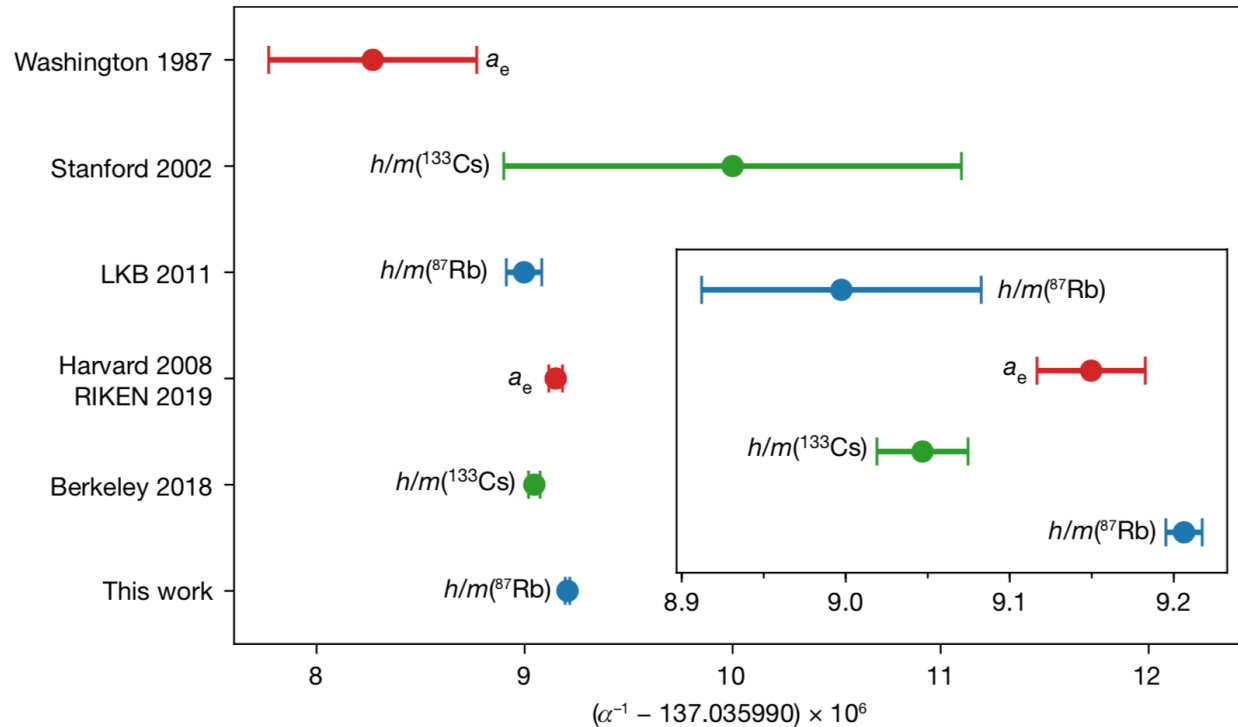


Check for updates

Léo Morel<sup>1</sup>, Zhibin Yao<sup>1</sup>, Pierre Cladé<sup>1</sup> & Saïda Guellati-Khélifa<sup>1,2✉</sup>

The standard model of particle physics is remarkably successful because it is consistent with (almost) all experimental results. However, it fails to explain dark matter, dark energy and the imbalance between matter and antimatter in the Universe. Because discrepancies between standard-model predictions and experimental observations may provide evidence of new physics, an accurate evaluation of these predictions requires highly precise values of the fundamental physical constants. Among them, the fine-structure constant  $\alpha$  is of particular importance because it sets the strength of the electromagnetic interaction between light and charged elementary particles, such as the electron and the muon. Here we use matter-wave interferometry to measure the recoil velocity of a rubidium atom that absorbs a photon, and determine the fine-structure constant  $\alpha^{-1} = 137.035999206(11)$  with a relative accuracy of 81 parts per trillion. The accuracy of eleven digits in  $\alpha$  leads to an electron  $g$  factor<sup>1,2</sup>—the most precise prediction of the standard model—that has a greatly reduced uncertainty. Our value of the fine-structure constant differs by more than 5 standard deviations from the best available result from caesium recoil measurements<sup>3</sup>. Our result modifies the constraints on possible candidate dark-matter particles proposed to explain the anomalous decays of excited states of  $^8\text{Be}$  nuclei<sup>4</sup> and paves the way for testing the discrepancy observed in the magnetic moment anomaly of the muon<sup>5</sup> in the electron sector<sup>6</sup>.

# $\alpha$ and $a_e$

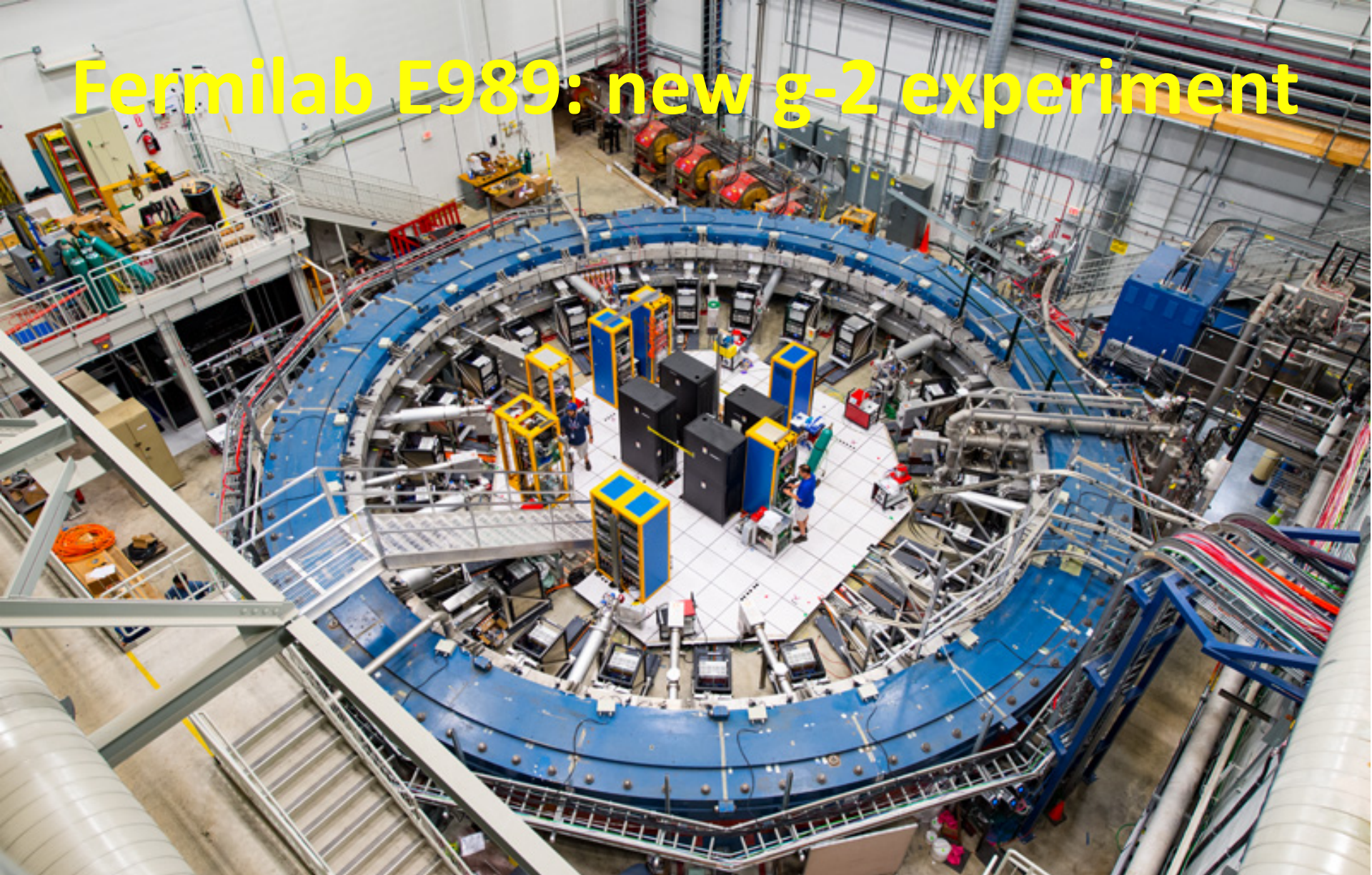


**Fig. 1 | Precision measurements of the fine-structure constant.** Comparison of most precise determinations of the fine-structure constant so far. The red points are from  $g_e - 2$  measurements and QED calculations, and the green and blue points are obtained from measurements of caesium and rubidium atomic

recoils, respectively. Errors bars correspond to  $\pm 1\sigma$  uncertainty. Previous data are from ref. <sup>34</sup> (Washington 1987), ref. <sup>10</sup> (Stanford 2002), ref. <sup>18</sup> (LKB 2011), ref. <sup>9</sup> (Harvard 2008), ref. <sup>2</sup> (RIKEN 2019) and ref. <sup>3</sup> (Berkeley 2018). Inset, magnification of the most accurate values of the fine-structure constant.



# Fermilab E989: new g-2 experiment



Fermilab and JPARC experiments aim at a reduction of the experimental error by a factor of 4



# Breaking News

## First results from the Muon g-2 experiment at Fermilab

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April 7, 2021, 10:00 am US/Central

The first results from the [Muon g-2](#) experiment at Fermilab will be unveiled and discussed in a special seminar to be held Wednesday, April 7, 2021, at 10:00 AM US Central Time.

The Muon g-2 experiment searches for telltale signs of new particles and forces by examining the muon's interaction with a surrounding magnetic field. By precisely determining the magnetic moment of the muon and comparing with similarly exact theoretical predictions, the experiment is sensitive to new physics lurking in the subatomic quantum fluctuations surrounding the muon. A previous experiment performed two decades ago at Brookhaven National Laboratory revealed an intriguing hint of such physics. The highly anticipated result from Fermilab pushes the precision of the experiment into uncharted territory in the quest to confirm or refute that finding.

The experimental result will be presented by Chris Polly, Fermilab physicist and co-spokesperson for the Muon g-2 scientific collaboration, following a summary of the current theoretical status given by Aida El-Khadra, a UIUC theoretical physicist and co-chair of the Muon g-2 Theory Initiative.

Seminar agenda:

10:00 – 10:05 Introduction

10:05 – 10:20 Theory overview — Aida El-Khadra, UIUC theoretical physicist

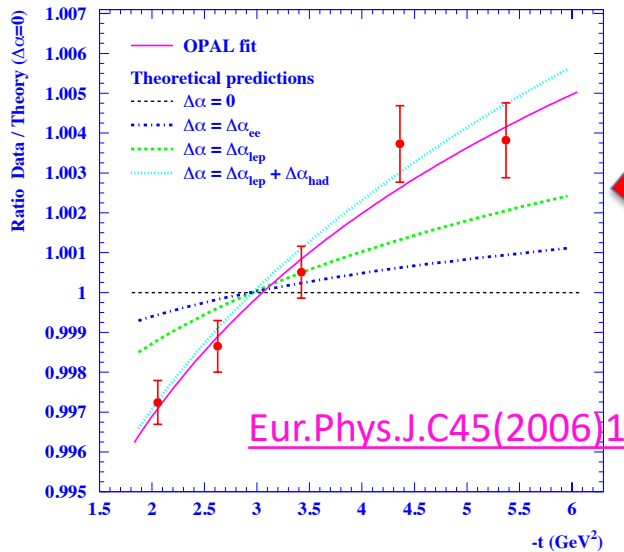
10:20 – 11:00 Muon g-2 results — Chris Polly, Fermilab experimental physicist

11:00 – 11:20 Question & Answer

# MUonE idea

## EXPERIMENT: OPAL @LEP

$$e^+e^- \rightarrow e^+e^- \quad \sqrt{s} \approx 91.2 \text{ GeV}$$



## THEORY

C.M. Carloni Calame, M. Passera, L. Trentadue, G. Venanzoni,  
[Phys.Lett.B746\(2015\)325](#)

-Initially proposed for use with Bhabha scattering data from  
 flavour factories-

$$a_{\mu}^{HLO} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{had}[t(x)]$$



*Idea tutta italiana: Gruppo misto di teorici e  
 sperimentali, inizio collaborazione 2016,  
 partecipazione al workshop del CERN Physics  
 Beyond Colliders*

Measuring the leading hadronic contribution to the muon  $g-2$  via  $\mu e$  scattering

[G. Abbiendi](#)<sup>1</sup>, C. M. Carloni Calame<sup>2</sup>, [U. Marconi](#)<sup>1</sup>, C. Matteuzzi<sup>3</sup>, G. Montagna<sup>4,2</sup>,  
 O. Nicosini<sup>2</sup>, M. Passera<sup>5</sup>, F. Piccinini<sup>2</sup>, [R. Tenchini](#)<sup>6</sup>, L. Trentadue<sup>7,3</sup>, and [G. Venanzoni](#)<sup>8</sup>

<sup>1</sup>INFN, Sezione di Bologna, Bologna, Italy

<sup>2</sup>INFN, Sezione di Pavia, Pavia, Italy

<sup>3</sup>INFN, Sezione di Milano Bicocca, Milano, Italy

<sup>4</sup>Dipartimento di Fisica, Università di Pavia, Pavia, Italy

<sup>5</sup>INFN, Sezione di Padova, Padova, Italy

<sup>6</sup>INFN, Sezione di Pisa, Pisa, Italy

<sup>7</sup>Dipartimento di Fisica e Scienze della Terra "M. Melloni",  
 Università di Parma, Parma, Italy

<sup>8</sup>INFN, Laboratori Nazionali di Frascati, Frascati, Italy

MUonE idea:

[Eur.Phys.J.C77\(2017\)139](#)

# MUonE experiment proposal

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH



June 5, 2019

## Letter of Intent: The MUonE Project

The MUonE Collaboration

**5 Giugno 2019:** Lol presentata al SPSC  
[Letter-Of-Intent SPSC-I-252](#)

**22 Gennaio 2020:** decisione del SPSC:  
riconosciuto l'interesse fondamentale  
della proposta; approvato il Test Run  
nel 2021

**20 Maggio 2020:** CSN1 INFN apre la  
sigla MUonE

*INFN Bologna*  
G. Abbiendi, L. Capriotti, G. Galli, U. Marconi, C. Patrignani, A. Principe,  
*INFN Firenze*  
G. Sguazzoni,  
*Imperial College, London*  
J. Borg, K. Uchida, G. Hall, A. Howard, M. Pesaresi,  
*Krakow IFJ Pan*  
M. Baszczyk, P. Dorosz, M. Kucharczyk, M. Witek, M. Zdybal,  
*INFN Milano Bicocca*  
A. Broggio, C. Matteuzzi, M. Paganoni, M. Soldani, L. Trentadue,  
*Budker Institute, Novosibirsk*  
S. Eidelman, I. Logashenko, F. Ignatov,  
*INFN Padova*  
A. Bragagnolo, E. Conti, S. Di Vita, M. Fael, S. Laporta, P. Mastrolia, G. Ossola,  
P. Paradisi, M. Passera, M. Presilla, A. Primo, P. Ronchese, U. Schubert, G. Simi,  
F. Simonetto, R. Stroili, W.J. Torres-Bobadilla,  
*INFN Pavia*  
C. Carloni Calame, M. Chiesa, G. Montagna, O. Nicrosini, F. Piccinini,  
*INFN Pisa*  
G. Bagliesi, C. Ferrari\*, M. Incagli, F. Ligabue, F. Palla, R.N. Pilato, P. Spagnolo, R. Tenchini,  
G. Venanzoni, P. G. Verdini,  
*INFN Trieste*  
G. Cantatore, M. Karuza,  
*Shanghai Jiao Tong University*  
L. Li,  
*Paul Scherrer Institute*  
A. Signer, Y. Ulrich  
*University of Dublin*  
M.K. Marinkovic,  
*University of Liverpool*  
T. Bowcock, K. Rinnert, T. Teubner,  
*University of Virginia*  
D. Pocanic  
*CERN*<sup>†</sup>  
D. Abbaneo, J. Bernhard, D. Banerjee<sup>‡</sup>, S. Mersi

\*Istituto Nazionale di Ottica del C.N.R. UOS Pisa

<sup>†</sup>Having contributed to several studies, the members of CERN personnel do not take position nor responsibility towards the required approval processes as established by the organization.

<sup>‡</sup>also at University of Illinois at Urbana-Champaign

still growing up



CERN

*Exp*



INFN +Univ. (Bologna,  
Milano-Bicocca, Padova,  
Pavia, Perugia, Pisa, Trieste)

*Exp-Th*



Imperial College (London),  
Liverpool U.

*Exp-Th*

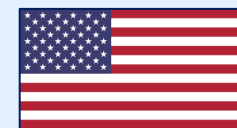


Krakow IFJ Pan

*Exp*

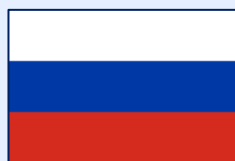


# The MUonE Collaboration



Northwestern U.,  
Virginia U.

*Exp*



Budker Inst.  
(Novosibirsk)

*Exp*



Demokritos INPP  
(Athens)

*Exp-Th*



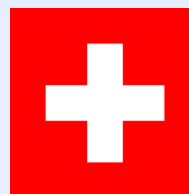
Shanghai  
Jiao Tong U.

*Exp*



LMU München

*Th*



PSI (Villigen),  
U.Zürich

*Th*

+ other involved theorists from: LAPH/Annecy (F), U.Valencia (E), KIT/Karlsruhe (D), New York City Tech (USA)

# $a_\mu^{HLO}$ : the MUonE approach (space-like data)

C.M. Carloni Calame, M. Passera, L. Trentadue, G. Venanzoni, [Phys.Lett.B746\(2015\)325](#)

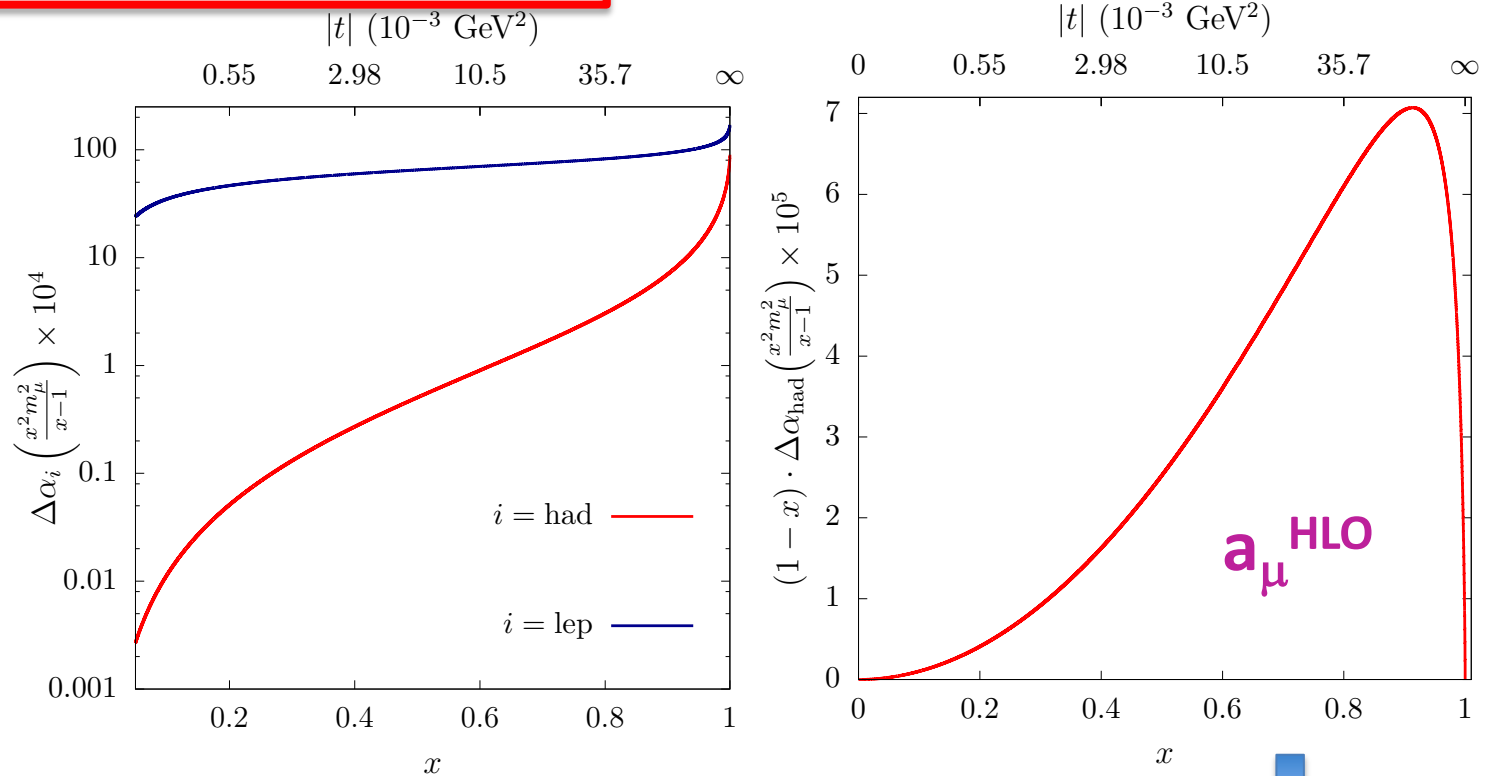
$$a_\mu^{HLO} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{had}[t(x)]$$

$$t(x) = -\frac{x^2 m_\mu^2}{1-x} \quad \begin{array}{l} 0 \leq -t < \infty \\ 0 \leq x < 1 \end{array}$$

$\Delta\alpha_{had}$  is the hadronic contribution to the running of  $\alpha$  in the space-like region ( $t < 0$ )

$$\alpha(t) = \frac{\alpha}{1 - \Delta\alpha(t)}$$

$$\Delta\alpha = \Delta\alpha_{lep} + \Delta\alpha_{had}$$



Integrand function smooth: no resonances

Low-energy enhancement:

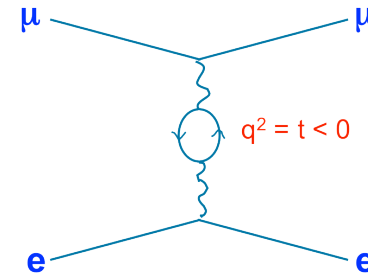
peak of the integrand at  $x \cong 0.9 \rightarrow t = -0.11 \text{ GeV}^2 \rightarrow \Delta\alpha_{had} \sim 10^{-3}$



# MUonE experiment idea

[Eur.Phys.J.C77\(2017\)139](#)

Very precise measurement of the running of  $\alpha_{\text{QED}}$  from the shape of the differential cross section of elastic scattering of  $\mu(150\text{-}160\text{GeV})$  on atomic electrons of a fixed target with low Z (Be or C) at the CERN SPS

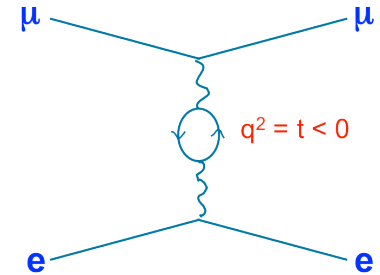


$$\frac{d\sigma}{dt} \approx \frac{d\sigma_0}{dt} \left| \frac{\alpha(t)}{\alpha(0)} \right|^2 \approx \frac{d\sigma_0}{dt} \left| \frac{1}{1 - \Delta\alpha(t)} \right|^2 \quad \xrightarrow{\text{running of } \alpha} \quad \Delta\alpha(t) = \underbrace{\Delta\alpha_{\text{lep}}(t)}_{\text{known from QED}} + \underbrace{\Delta\alpha_{\text{had}}(t)}_{\text{to be measured}}$$

From  $\Delta\alpha_{\text{had}}(t)$  determine  $a_\mu^{\text{HLO}}$  by the space-like approach: [Phys.Lett.B746\(2015\)325](#)

$$a_\mu^{\text{HLO}} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{\text{had}}[t(x)]$$

# $\mu$ -e elastic scattering



At LO

$$\frac{d\sigma}{dt} = \frac{4\pi\alpha^2}{\lambda(s, m_e^2, m_\mu^2)} \left[ \frac{(s - m_e^2 - m_\mu^2)^2}{t^2} + \frac{s}{t} + \frac{1}{2} \right]$$

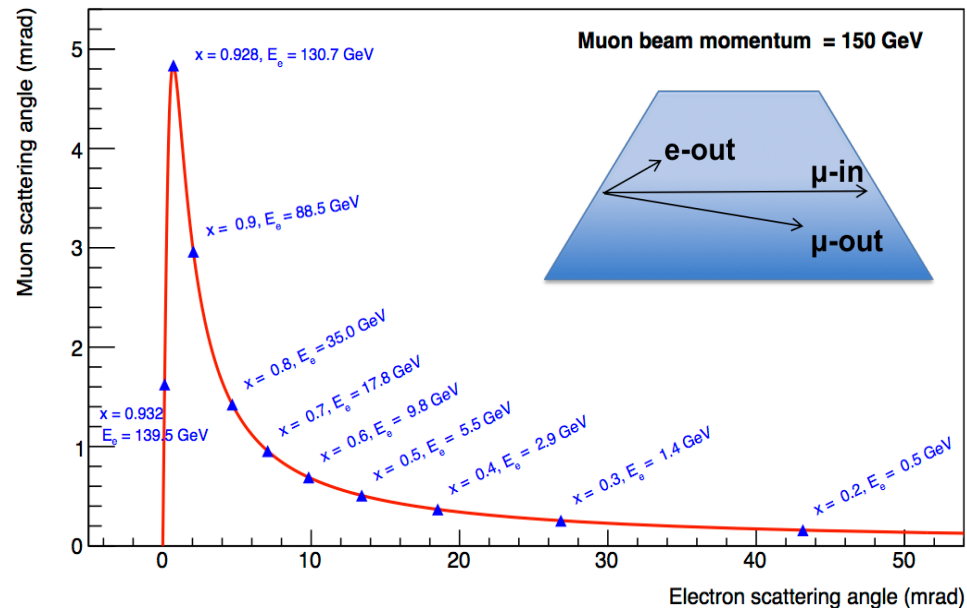
$$\frac{d\sigma}{dt} = \frac{d\sigma_0}{dt} \left| \frac{\alpha(t)}{\alpha(0)} \right|^2 \quad \alpha(t) = \frac{\alpha(0)}{1 - \Delta\alpha(t)} \quad \Delta\alpha(t) = \Delta\alpha_{\text{lep}}(t) + \Delta\alpha_{\text{had}}(t)$$

## ➤ Elastic scattering: simple kinematics

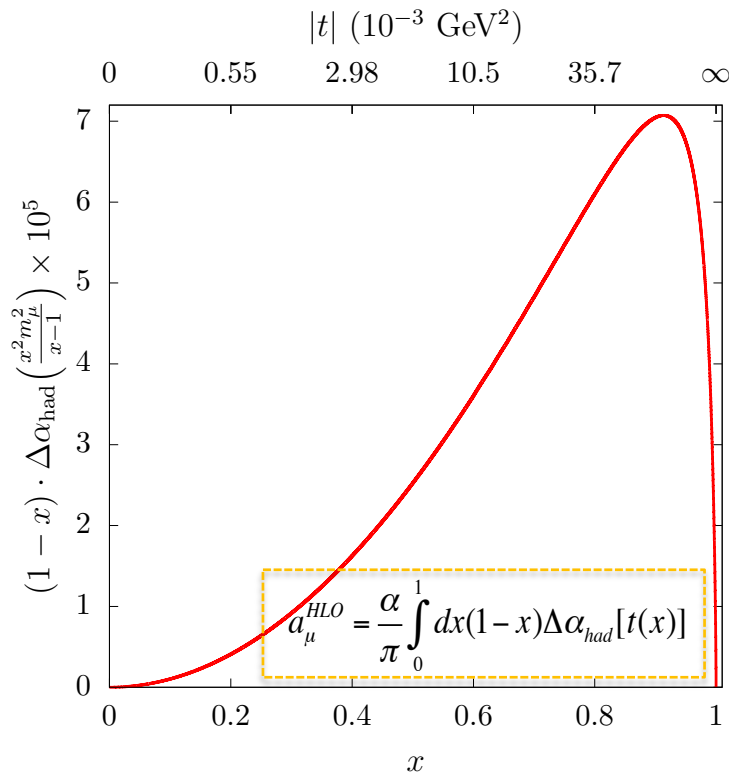
❖  $t \approx -2 m_e E_e$   $E_e$  can be determined from the scattering angle  $\theta_e$  and the beam energy

➤ Scattering angles  $\theta_e$  and  $\theta_\mu$  correlated (helps selection: rejection of radiative/inelastic events)

➤ Elastic events are planar



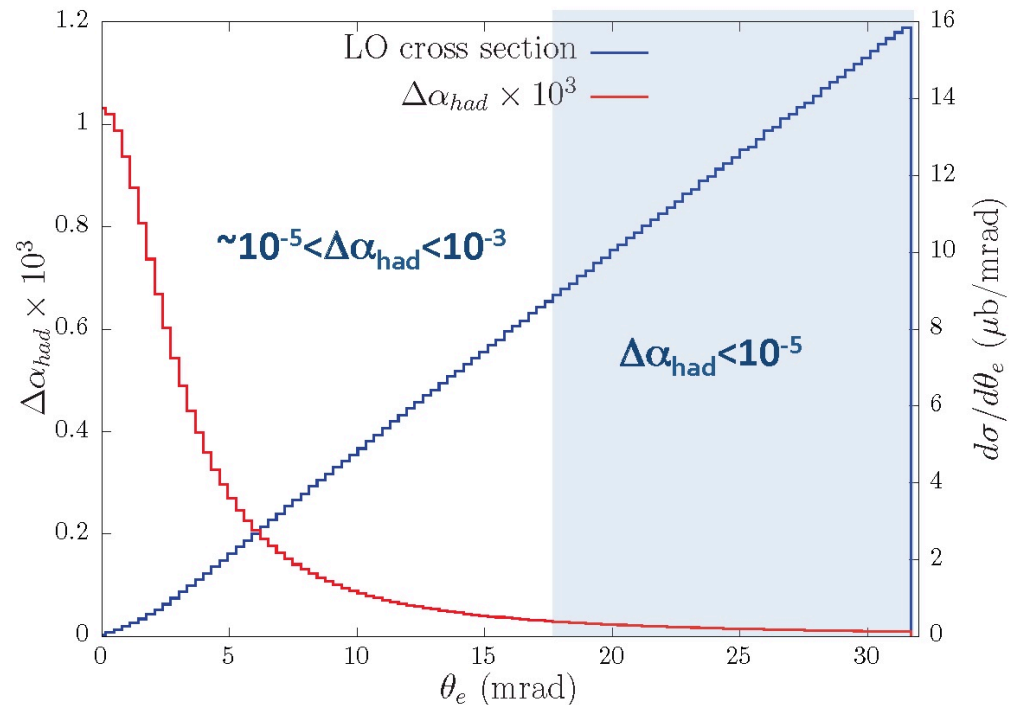
# $\mu$ -e elastic scattering (2)



**THIS IS CHALLENGING:**  
 Observable effect  $\sim 10^{-3}$   
 wanted accuracy  $\sim 10^{-2}$   
 $\rightarrow$  Required precision  $\sim 10^{-5}$   
 on the shape of  $d\sigma/dt$

For  $E(\text{beam})=150 \text{ GeV}$  the phase space covers 87% of the  $a_\mu^{\text{HLO}}$  integral.

Smooth extrapolation to the full integral with a proper fit model

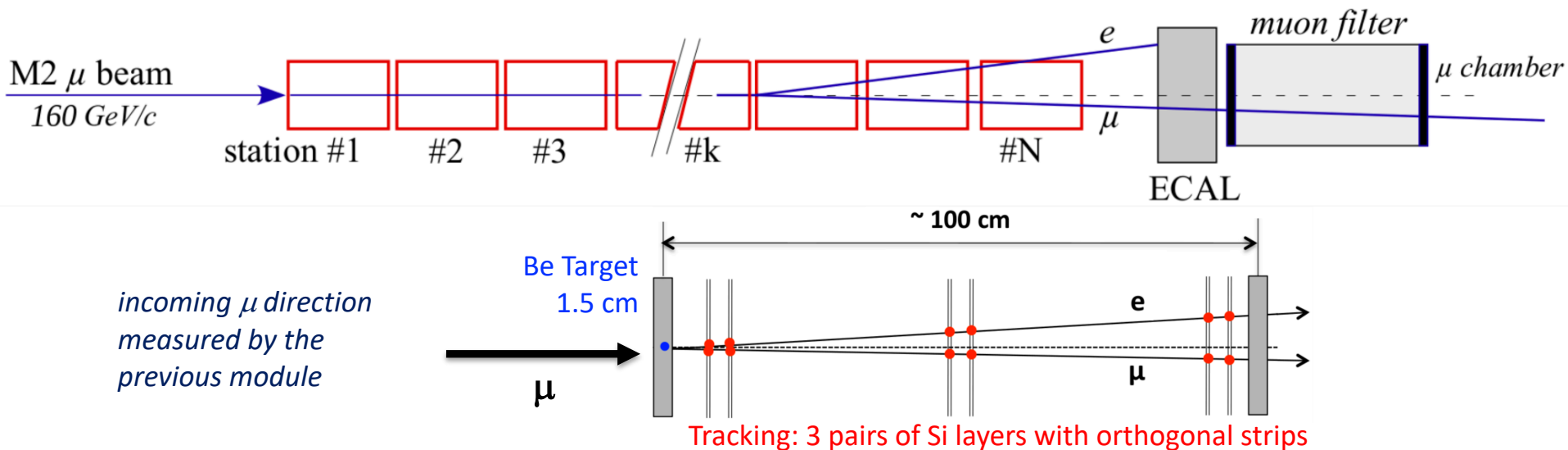


# MUonE Detector Layout

Letter-Of-Intent SPSC-I-252

*The detector concept is simple, the challenge is to keep the systematics at the same level as the statistical error .*

- Large statistics to reach the necessary sensitivity
- Minimal distortions of the outgoing  $e/\mu$  trajectories within the target material and small rate of radiative events
- Modular structure of 40 independent and precise tracking stations, with split light targets equivalent to 60cm Be
- ECAL and Muon filter after the last station, to help the ID and background rejection



Boosted kinematics:  $\theta_e < 32 \text{ mrad}$  (for  $E_e > 1 \text{ GeV}$ ),  $\theta_\mu < 5 \text{ mrad}$ : the whole acceptance can be covered with a  $10 \times 10 \text{ cm}^2$  silicon sensor at 1m distance from the target, reducing many systematic errors

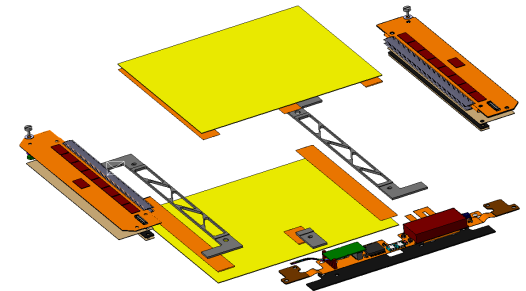
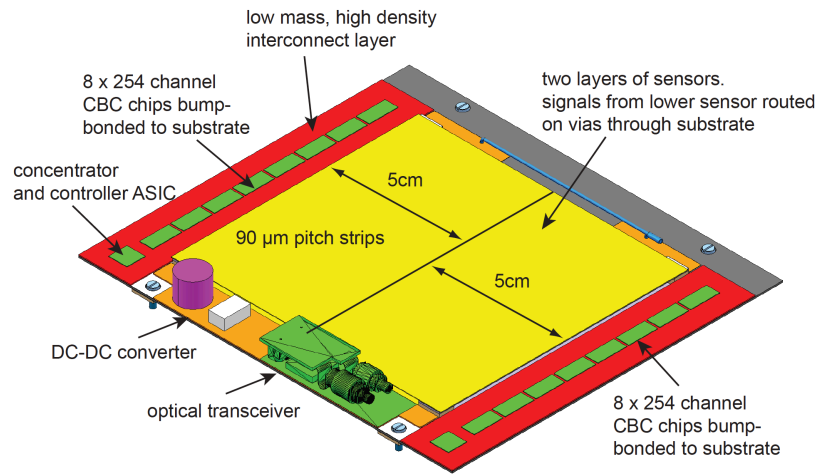
# Detector choice: CMS-upgrade Outer Tracker 2S

[MUonE Letter-Of-Intent SPSC-I-252](#)

Details: see [CMS Tracker Upgrade TDR](#)

Two close-by planes of strips reading the same coordinate, providing track elements (**stubs**)

suppression of background from single-layer hits or large-angle tracks



- Large active area  $10 \times 10 \text{ cm}^2$ 
  - > complete/uniform angular coverage with a single sensor
- Good position resolution  $\sim 20 \mu\text{m}$ 
  - > further improvable with a  $15^\circ$ - $20^\circ$  tilt around the strip axis and/or with effective staggering of the planes (with a microrotation)

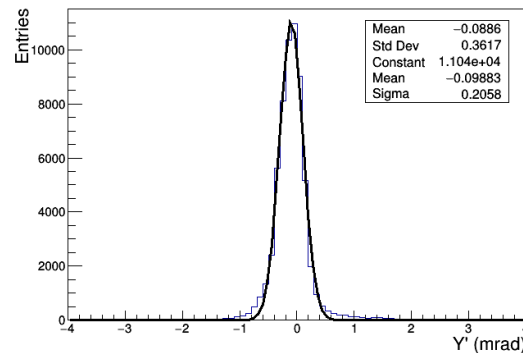
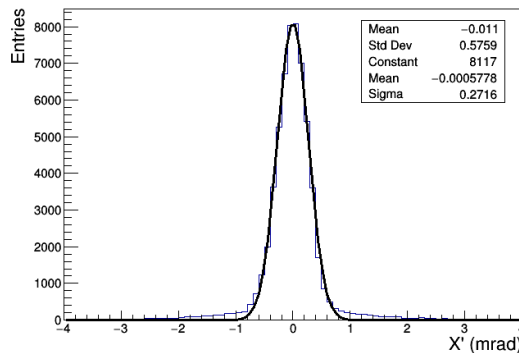
**MAIN Difference w.r.t. LHC operation:** signal is asynchronous while sampling has fixed clock at 40MHz -> can be overcome with a specific configuration of the FE

# Location @ CERN & M2 beam parameters

MUonE Letter-Of-Intent SPSC-I-252

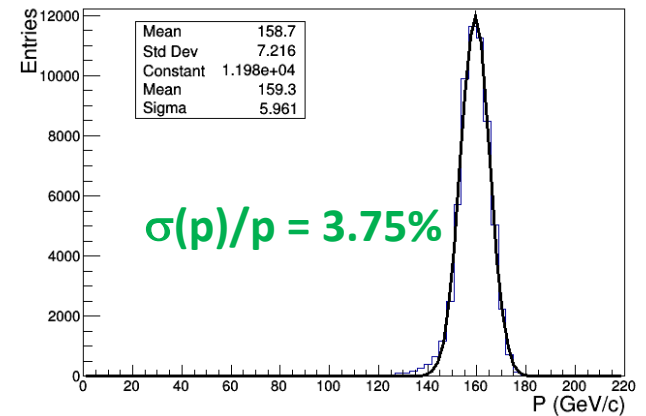


Very small divergence  $\sim 0.2\text{-}0.3$  mrad

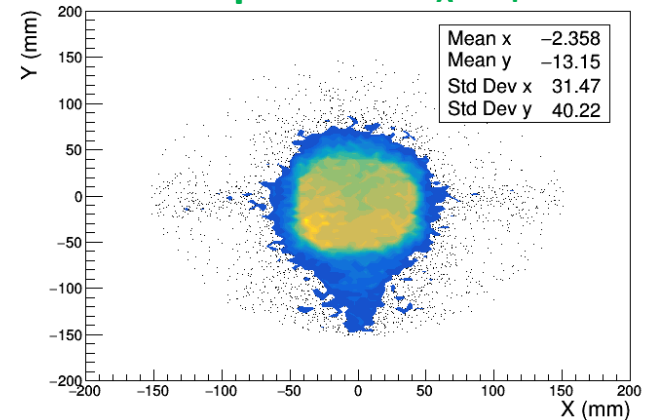


Upstream of the COMPASS detector,  
after its Beam Momentum Station  
(BMS), on the M2 beam line :  
available  $\sim 40$  m

Beam Momentum



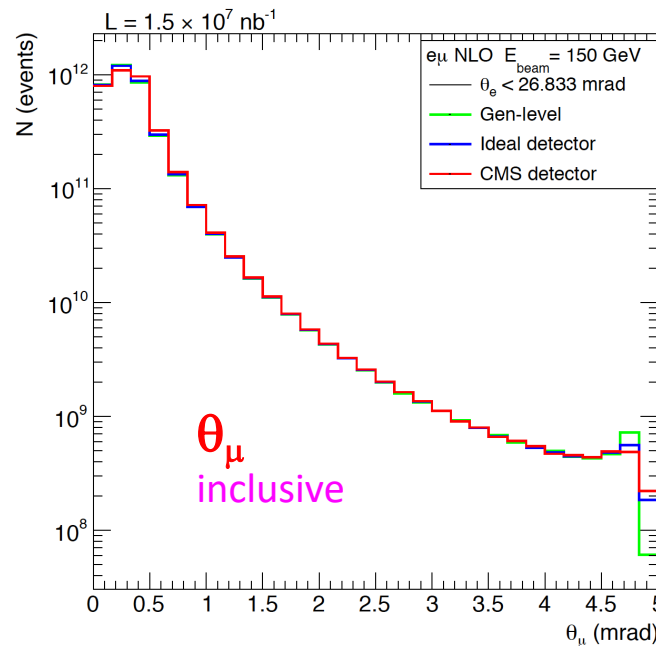
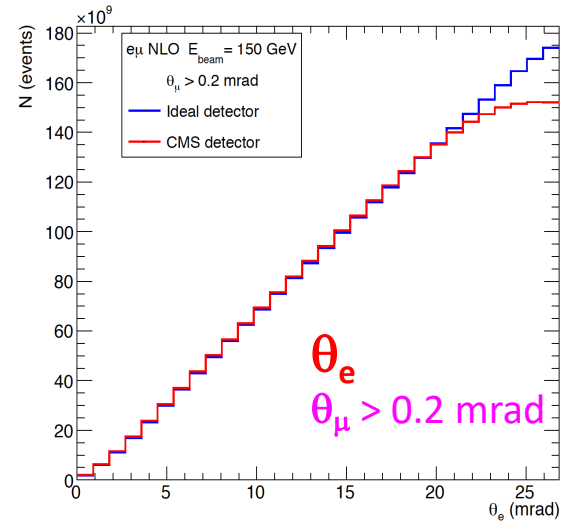
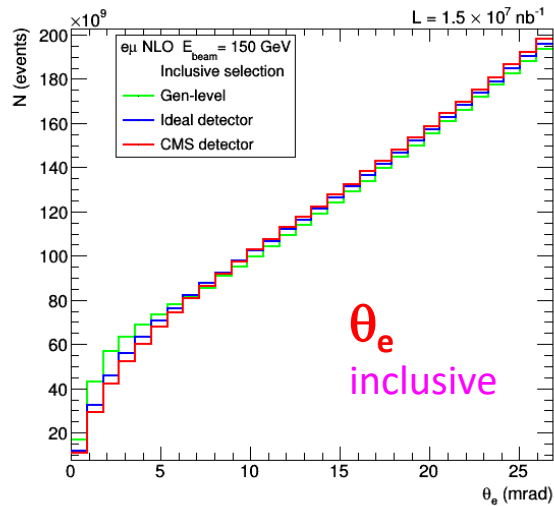
Beam spot size  $\sigma_x \sim \sigma_y \sim 3\text{cm}$



M2 beam typical max intensity:  $5 \times 10^7 \mu\text{s}$

SPS Fixed Target cycle  $\sim 15\text{-}20$  s / Spill duration  $\sim 5$  s

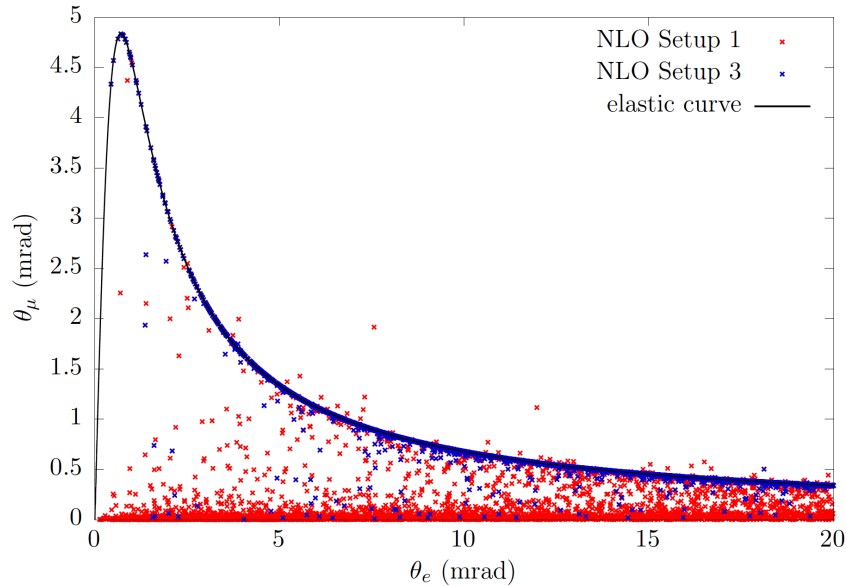
# NLO Angular distributions



Event yield corresponding to  
the nominal integrated  
Luminosity  $L = 1.5 \times 10^7 \text{ nb}^{-1}$

# NLO MC and elastic selection

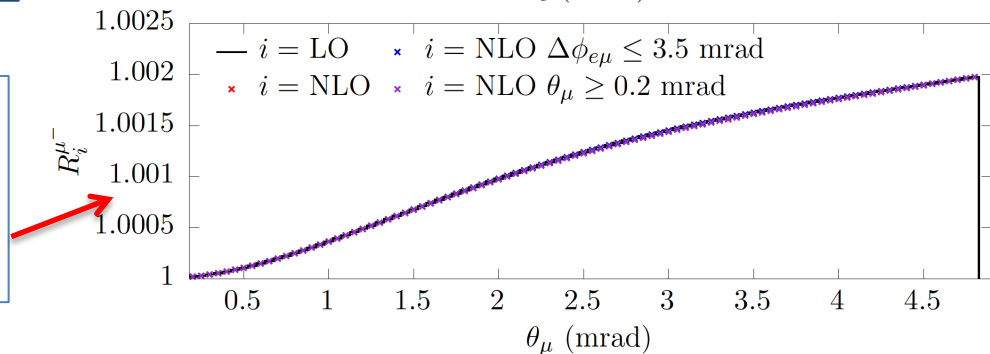
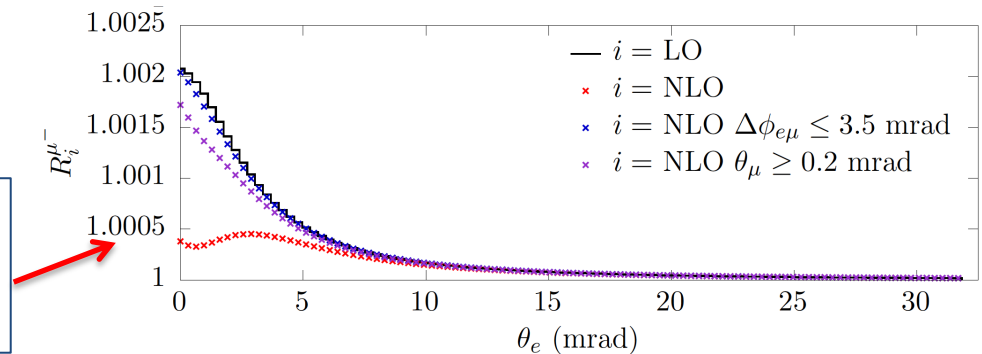
[M.Alacevich et al, JHEP02\(2019\)155](#)



Without any selection the signal sensitivity of the electron angle is destroyed -> necessary to implement an “elastic” selection

Instead the muon angle is a robust observable, stable w.r.t. radiative corrections -> it can be used with an inclusive selection (theoretically advantageous)

NLO Setup 1 is the inclusive selection (no cuts)  
Setup 3 has an acoplanarity cut  $|\pi - (\phi_e - \phi_\mu)| < 3.5$  mrad





# $\Delta\alpha_{had}$ parameterization

Physics-inspired from the calculable contribution of lepton-pairs and top quarks at  $t < 0$

$$\Delta\alpha_{had}(t) = k \left\{ -\frac{5}{9} - \frac{4M}{3t} + \left( \frac{4M^2}{3t^2} + \frac{M}{3t} - \frac{1}{6} \right) \frac{2}{\sqrt{1 - \frac{4M}{t}}} \log \left| \frac{1 - \sqrt{1 - \frac{4M}{t}}}{1 + \sqrt{1 - \frac{4M}{t}}} \right| \right\}$$

$M$  with dimension of mass squared, related to the mass of the fermion in the vacuum polarization loop

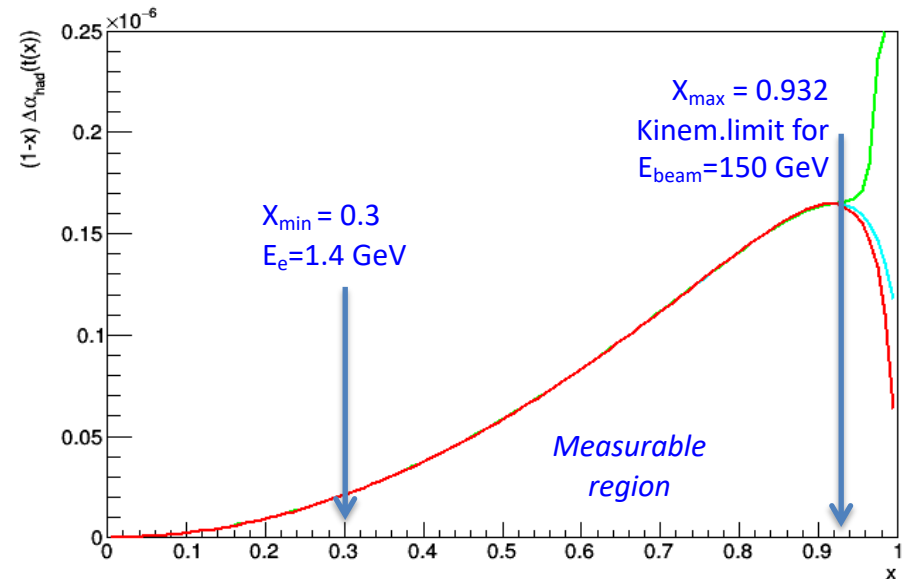
$k$  depending on the coupling  $\alpha(0)$ , the electric charge and the colour charge of the fermion

Low- $|t|$  behavior dominant in the MUonE kinematical range:

$$\Delta\alpha_{had}(t) \simeq -\frac{1}{15} \frac{k}{M} t$$

$a_\mu^{HLO}$  calculable from the master integral in the FULL phase space with this parameterization.

Instead simple polinomials diverge for  $x \rightarrow 1$  (green is a cubic polinomial in  $t$ )

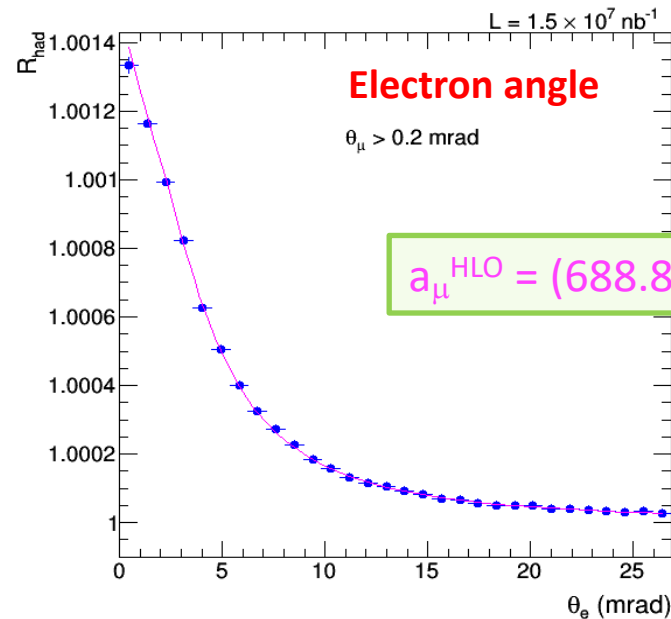
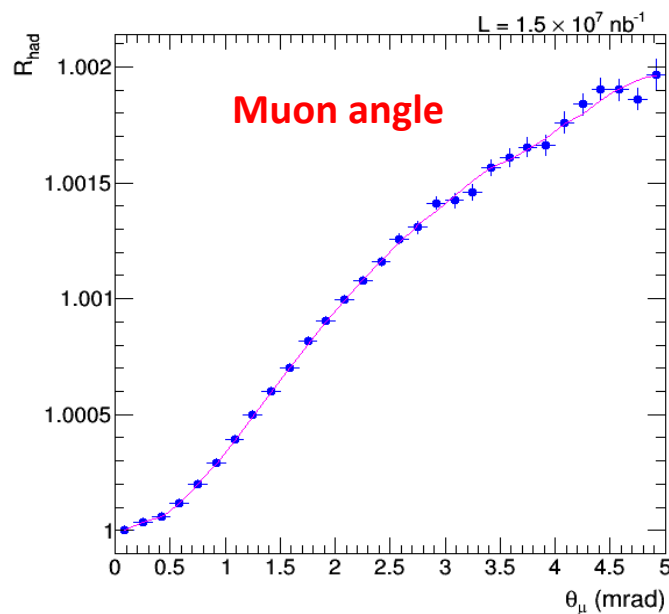


# Extraction of the hadronic running of $\alpha$

Most easily displayed by taking **ratios** of the observed angular distributions and the theory predictions evaluated for  $\alpha(t)$  corresponding to only the leptonic running.

Observable effect  $\sim 10^{-3}$  / wanted precision  $\sim 10^{-2} \rightarrow$  required precision  $\sim 10^{-5}$

Example toy experiment shown with statistics corresponding to the nominal integrated Luminosity  $L = 1.5 \times 10^7 \text{ nb}^{-1}$  (corresponding to 3-year run)



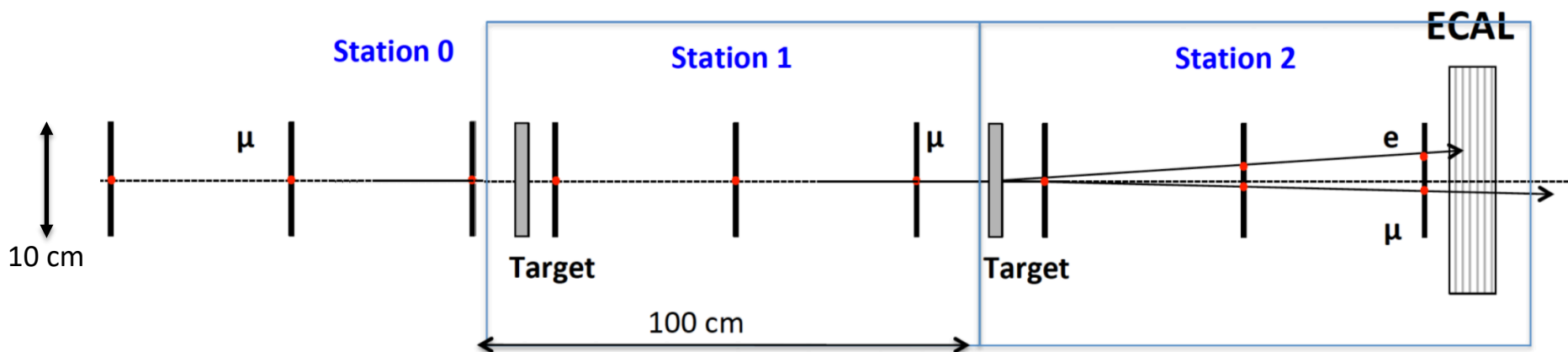
*Template fit to the 2D angular distribution from NLO MC generator with parameterised detector resolution.*

*$\Delta\alpha_{\text{had}}(t)$  parameterised according to the “Lepton-Like” form. Shape-only  $\chi^2$  fit.*

# Test Run setup

*To be held at CERN in Fall 2021: 3 weeks allocated with full intensity  $\mu$  beam*

*Location: M2 beam line, upstream of the COMPASS detector, after its BMS (available  $\sim 40$  m)*

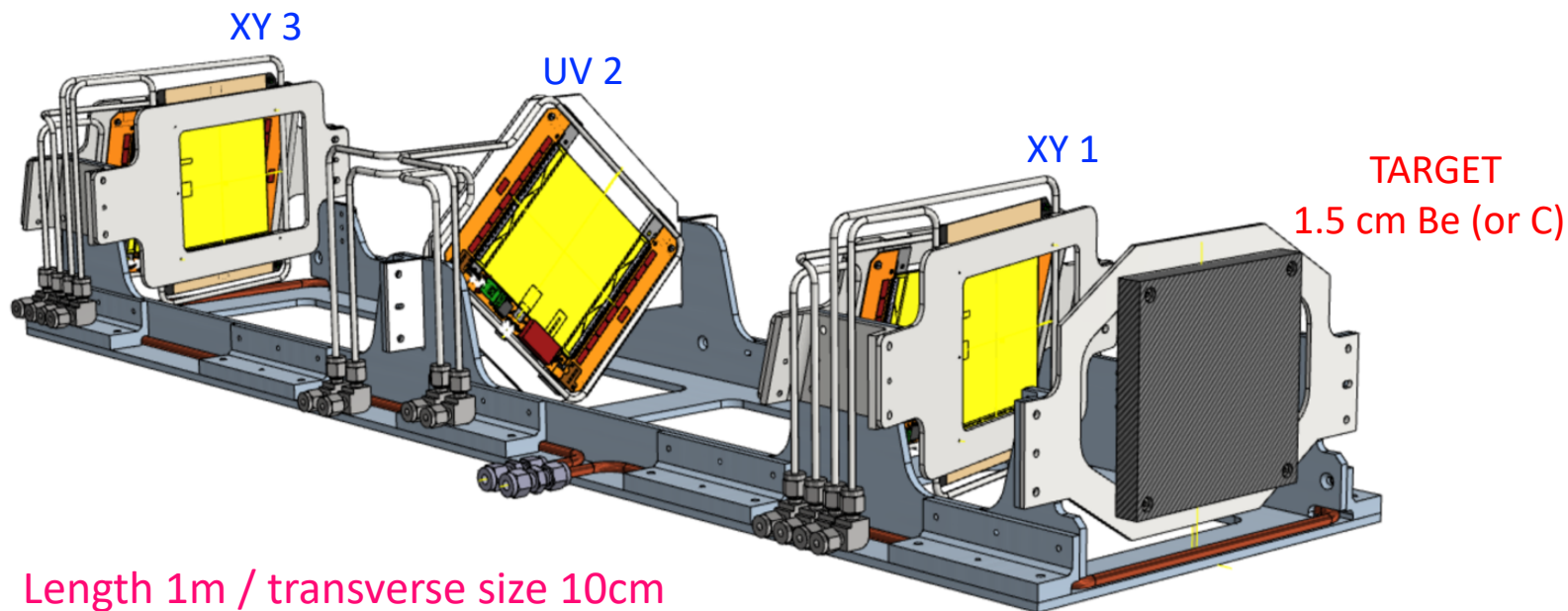


## Main objectives:

- Confirm the system engineering
- Check mechanical and thermal stability.
- Test the alignment procedure
- Assess the detector counting rate capability.
- Check the DAQ system.
- Validate the trigger strategy (FPGA real-time processing to identify and reconstruct  $\mu$ -e events).
- Assess the systematic errors
- After commissioning, take data to measure the leptonic contribution to the running of  $\alpha(q^2)$ .

**If the results are satisfactory proceed to full-scale experiment to be deployed during LHC Run3**

# MUonE tracking station



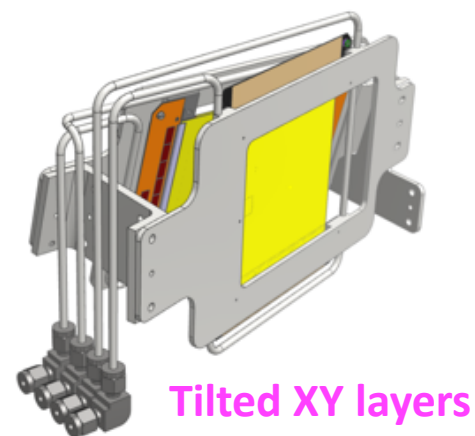
Length 1m / transverse size 10cm

Target followed by 3 tracking layers:  
each one is a pair of close-by 2S modules  
with orthogonal strips, tilted by 233mrad

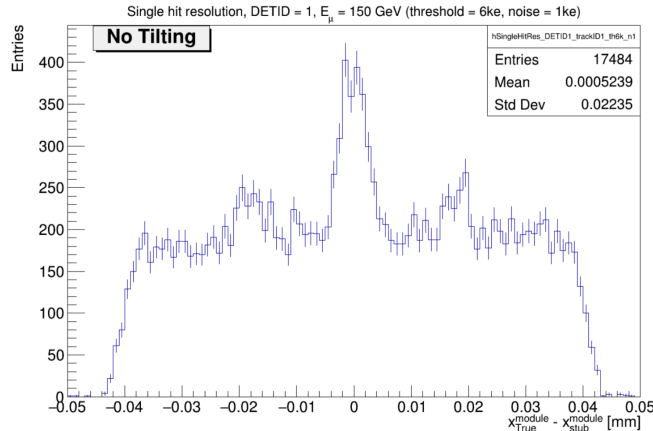
**Stringent request: relative positions within the station stable to better than 10 $\mu$ m**

Low CTE support structure: INVAR (alloy of 65%Fe, 35%Ni)

Cooling system, tracker enclosure, Room temperature  
stabilized within 1-2  $^{\circ}$ C



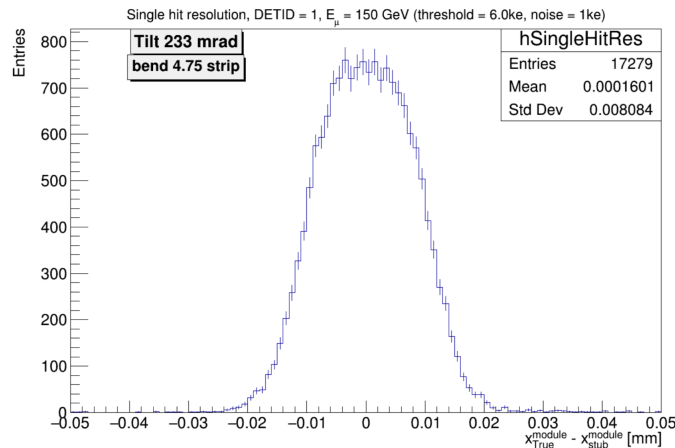
# Simulation: Intrinsic Resolution – Tilted geometry



Strip digital readout: with  $90\mu\text{m}$  pitch the expected resolution is  $90/\sqrt{12} \cong 26\mu\text{m}$  on a single sensor layer for single-strip clusters

Tilting a sensor around an axis parallel to the strips →  
Charge sharing between adjacent strips, improving the resolution

The best is obtained when  $\langle \text{cluster width} \rangle \sim 1.5$  (same number of clusters made of 1 or 2 strips) for a tilt angle  $\sim 15$  degrees

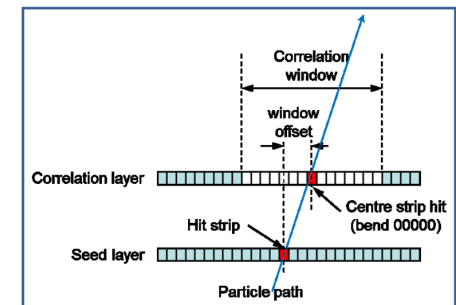


Further improvement: a small tilt of 25mrad is equivalent to an half-strip staggering of the two sensor layers of a 2S module

**Final resolution:**  
 **$22\mu\text{m} \rightarrow 8\text{--}11\mu\text{m}$**

measured coordinate (x) determined by hit position on one layer and direction of the track stub

Tilt angle [mrad]	$\langle \text{bend} \rangle$ [strips]	threshold $[\sigma]$	resolution $[\mu\text{m}]$	$\langle \text{cluster width} \rangle$ [strips]
210	4.25	5	7.8	1.51
221	4.5	5.5	11.5	1.51
233	4.75	6	8.0	1.50
245	5	6.5	11.2	1.51
257	5.25	7	8.7	1.50
268	5.5	7.5	11.0	1.49

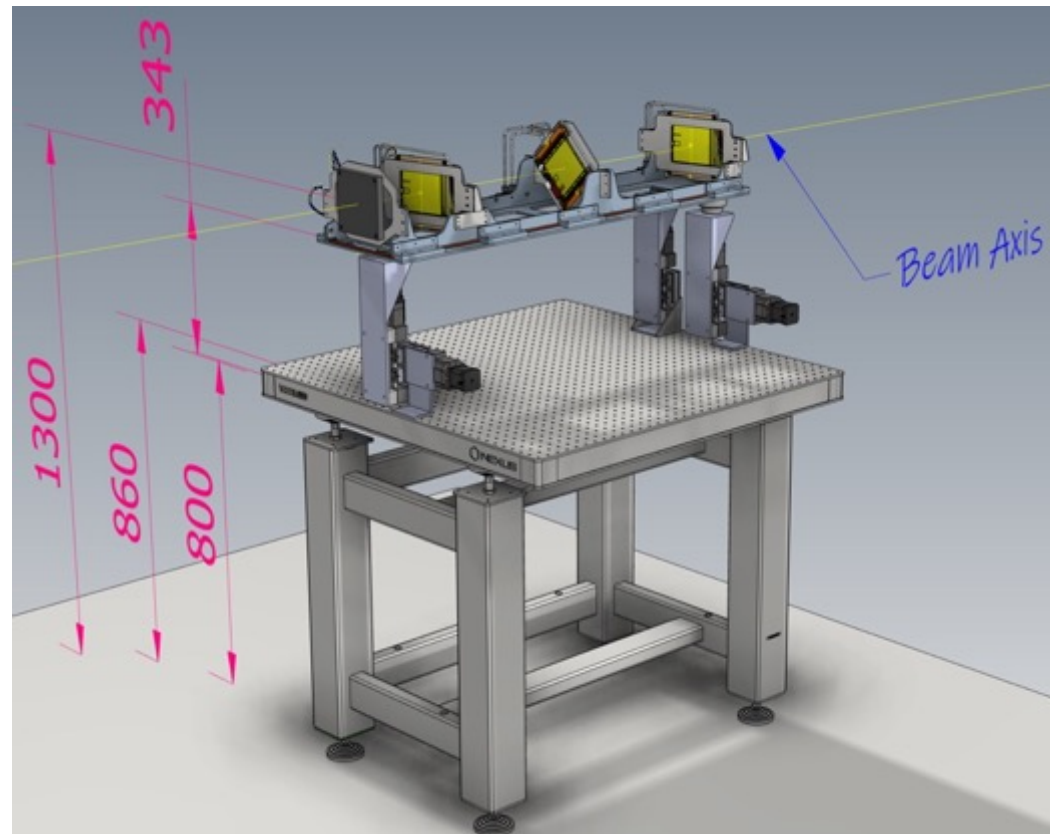


# Tracker mechanics

Two aluminium mockups have been built:  
test mounting of dummy stations, planarity,  
alignment, cooling system, precision  
movement system and holographic system

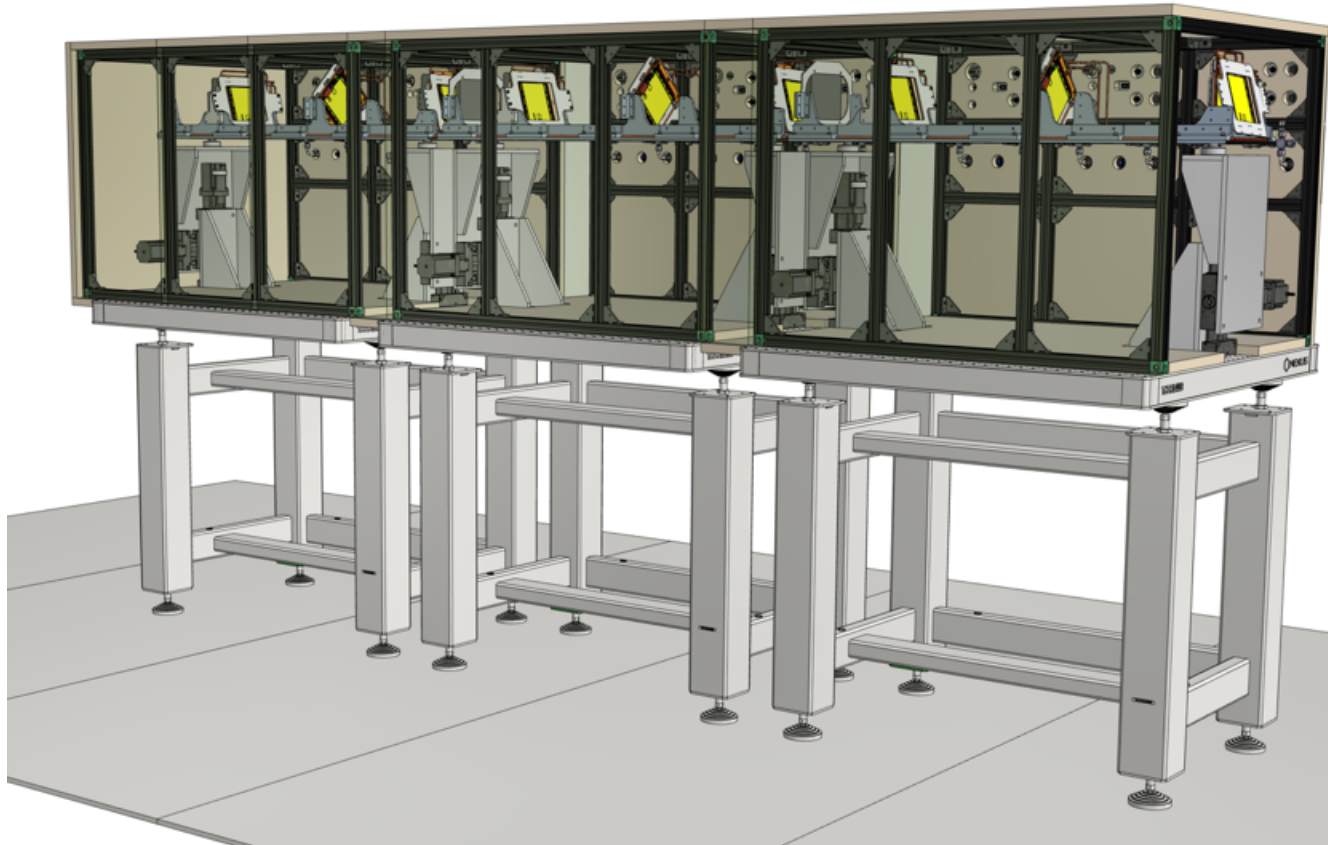


Each station's position/orientation will be precisely adjustable with 3 motorized linear stages allowing to shift on X, Y axes by up to 3cm in steps of  $5\text{ }\mu\text{m}$  (by kinematic coupling)





# Tracker mechanics (2)



- Tracker enclosure shielding from light and to stabilise thermally
- Electrical, optical and hydraulic connections on the top, removable side panels
- Further complemented by a surrounding tent containing also the calorimeter, with chiller stabilising the room temperature

# Calorimeter

## PbWO<sub>4</sub> crystals used by the CMS ECAL

Small 5x5 array, size 14x14 cm<sup>2</sup>, length 22cm (24.7 X<sub>0</sub>)

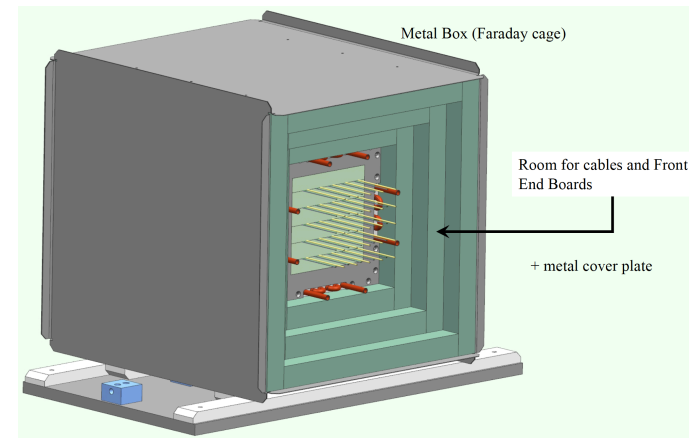
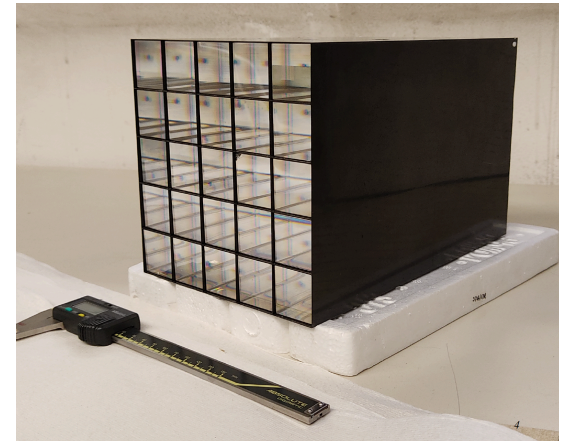
**Mechanics:** Carbon fiber alveolar structure with

- cooling system
- thermal insulation by polyurethane rigid foam panels and temperature control ( $\Delta T < \sim 0.1$  °C)
  - Both crystal light yield and APD gain depend on temperature:  
( $\approx -2\%/^{\circ}\text{C}$  for the crystals, and  $\approx -4\%/^{\circ}\text{C}$  for the APDs)
- all cables and fibers on the back face
- movable with mm precision in the two axes perpendicular to beam

Hamamatsu APD sensor (1 cm<sup>2</sup>)

FE electronics linking with Serenity board for DAQ

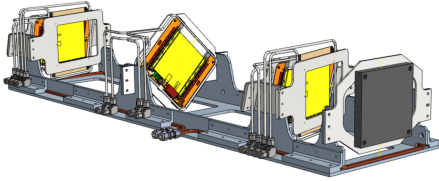
Laser calibration /monitoring system for APD and FEE gain





# DAQ

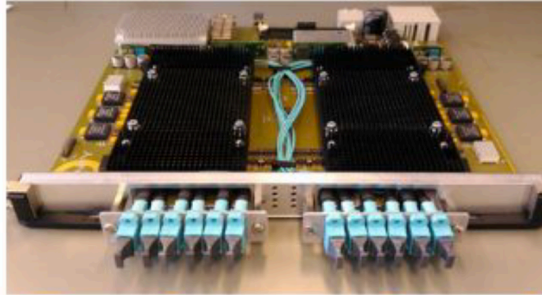
Tracker: event size ~1kb



35 Gbps total

*in-spill*

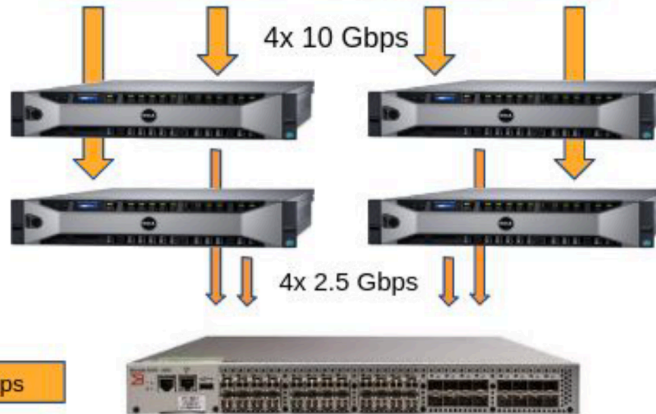
DAQ heart:  
SERENITY board



ECAL (safe factor 2)



Test Run:  
~0.5 PB of data



*SPS Duty cycle  
~0.25*

**Plan for the Test Run: NO online selection, read out all data (3 stations)**

FPGA algos will be run online just to tag events and replayed offline for detailed studies

Data taking for ~two weeks, SPS efficiency ~2/3 → ~0.5 PB of data

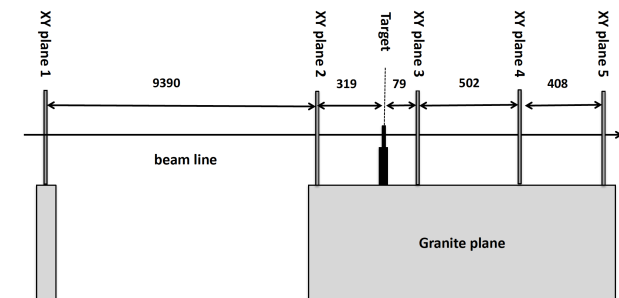
**The Test Run will be a proof of concept for the MUonE DAQ**

# 2017 Beam Test: Multiple Coulomb scattering

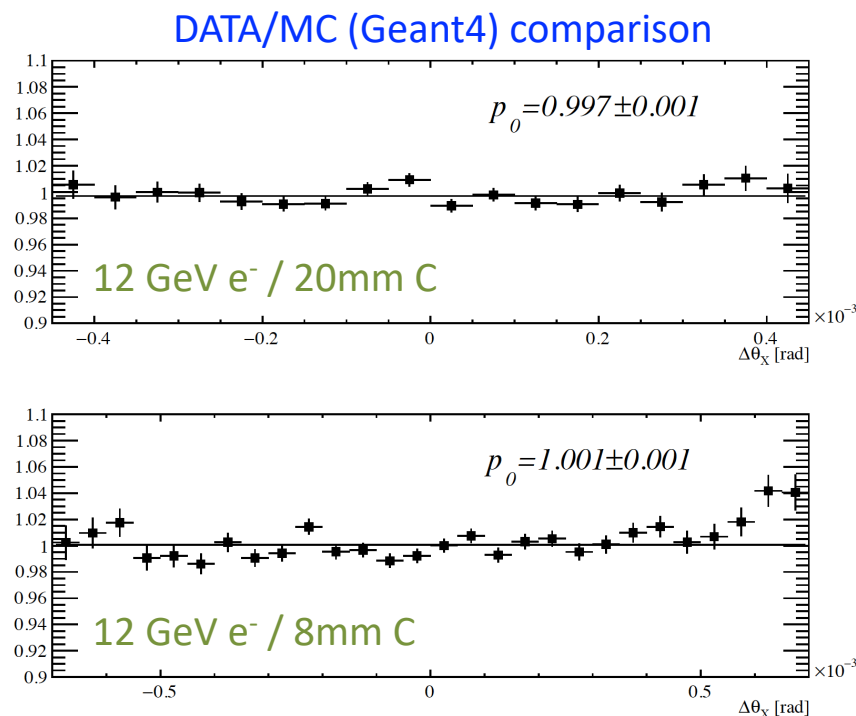
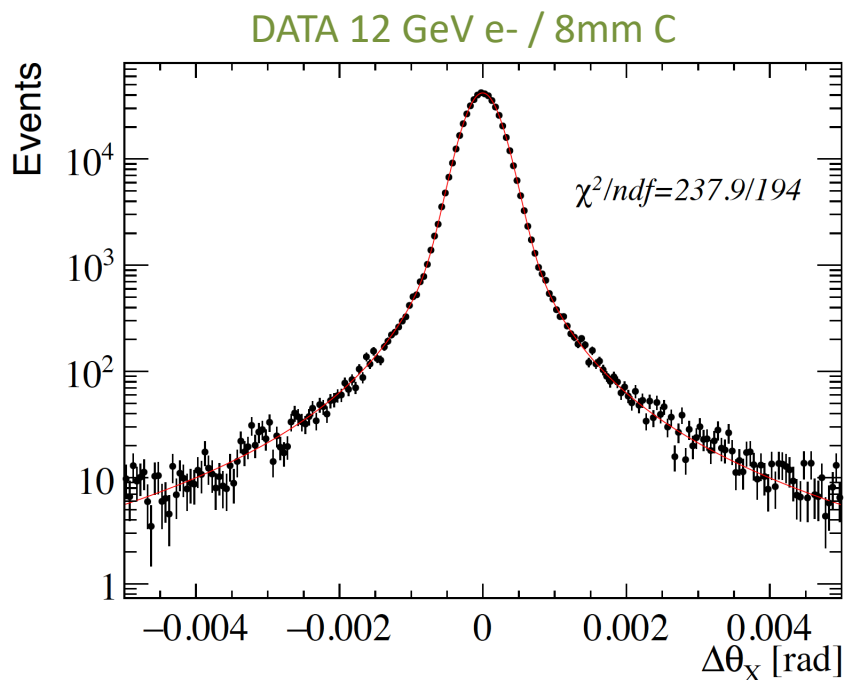
Studied in a Beam Test in 2017:

[JINST 15 \(2020\) P01017](#)

12–20 GeV electrons  
on 8-20 mm C targets



Adapted UA9  
detector at CERN  
H8 Beam Line



- Good description of data with a fit.
- Distribution core within 1-few % from GEANT.

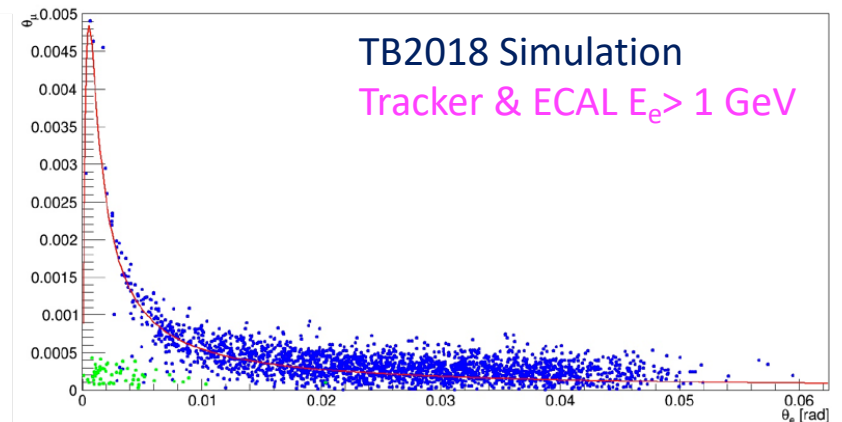
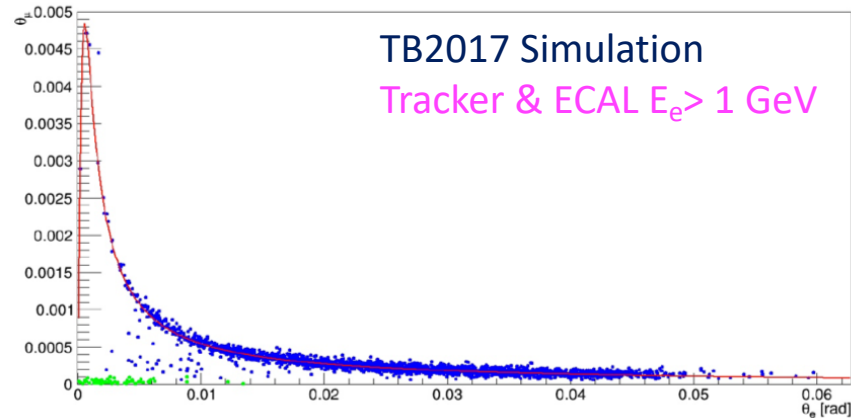
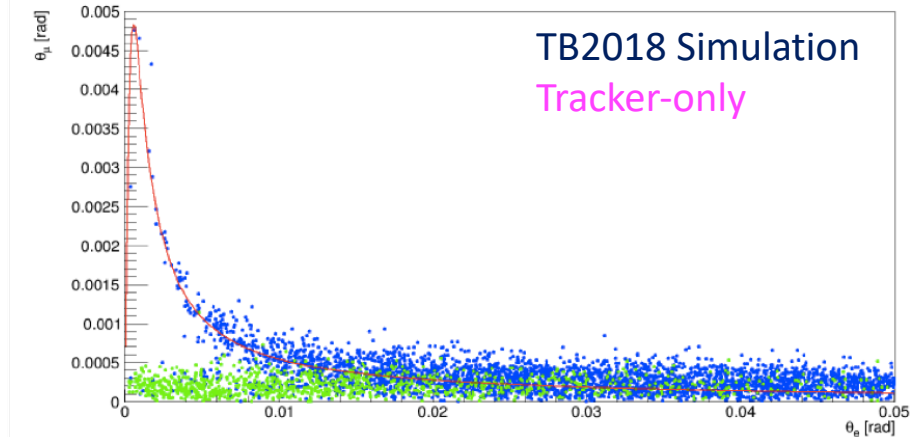
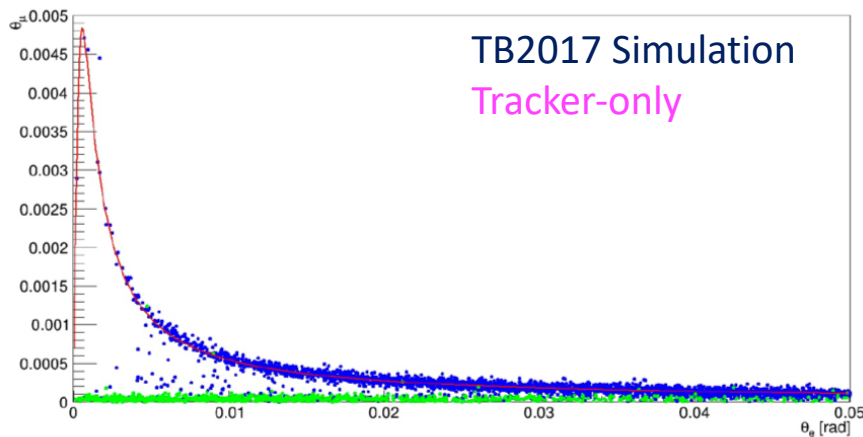
# GEANT4 simulations

Effect of the tracker position resolution on  $\theta_\mu$  vs  $\theta_e$  distribution:

(Left) TB2017: UA9 resolution  $7\mu\text{m}$  ; (Right) TB2018: resolution  $\sim 35\text{-}40\mu\text{m}$

Signal: elastic  $\mu e$

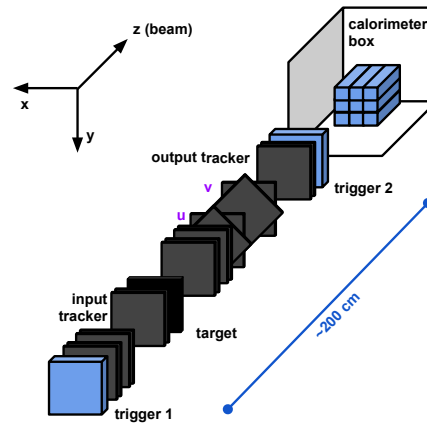
Background:  $e^+e^-$  pair production



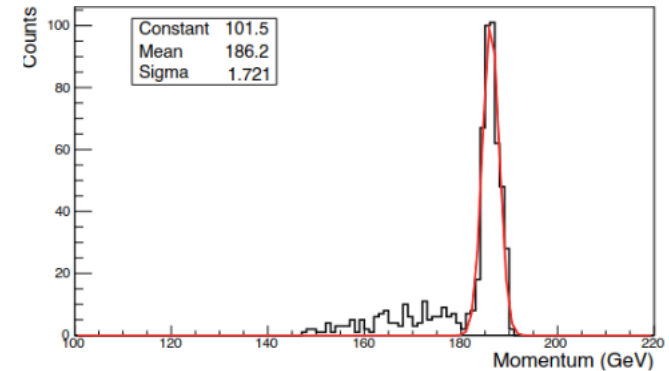
# 2018 Beam Test: $\mu e$ elastic scattering

[arXiv:2102.11111](https://arxiv.org/abs/2102.11111)

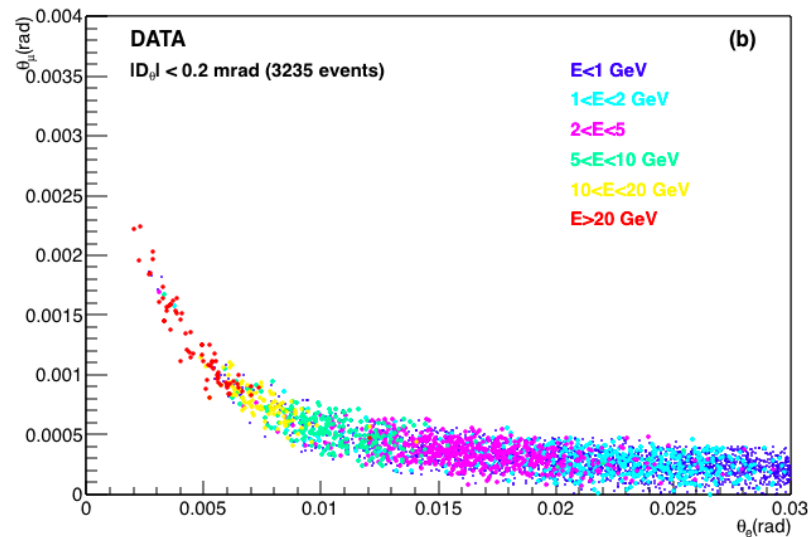
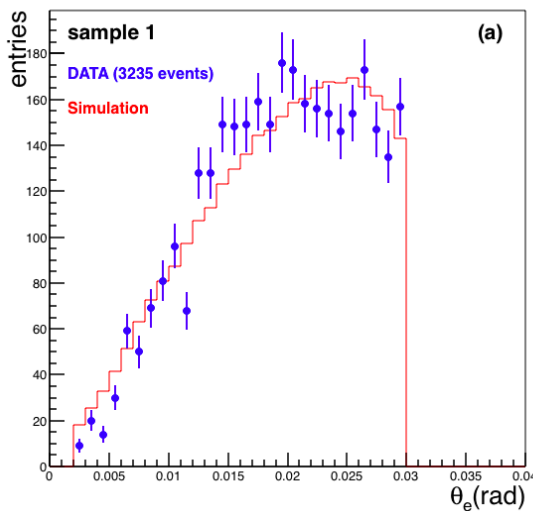
CERN North Area,  
downstream COMPASS  
8mm C target  
Si strip tracking (sensors  
from AGILE, with worse  
resolution than MUonE)  
Small BGO ECAL



$\mu$  spectrum peaked at 187 GeV  
From decays of 190 GeV beam  $\pi$   
1m W dump absorbing all surviving  $\pi$



Setup with lower performance than MUonE ( $\sigma_x \sim 35 \mu\text{m}$ )  
Selection of a clean sample of elastic events



Important:  
Simulation of  
Background  
processes in part.  
 $e^+e^-$  pair  
production

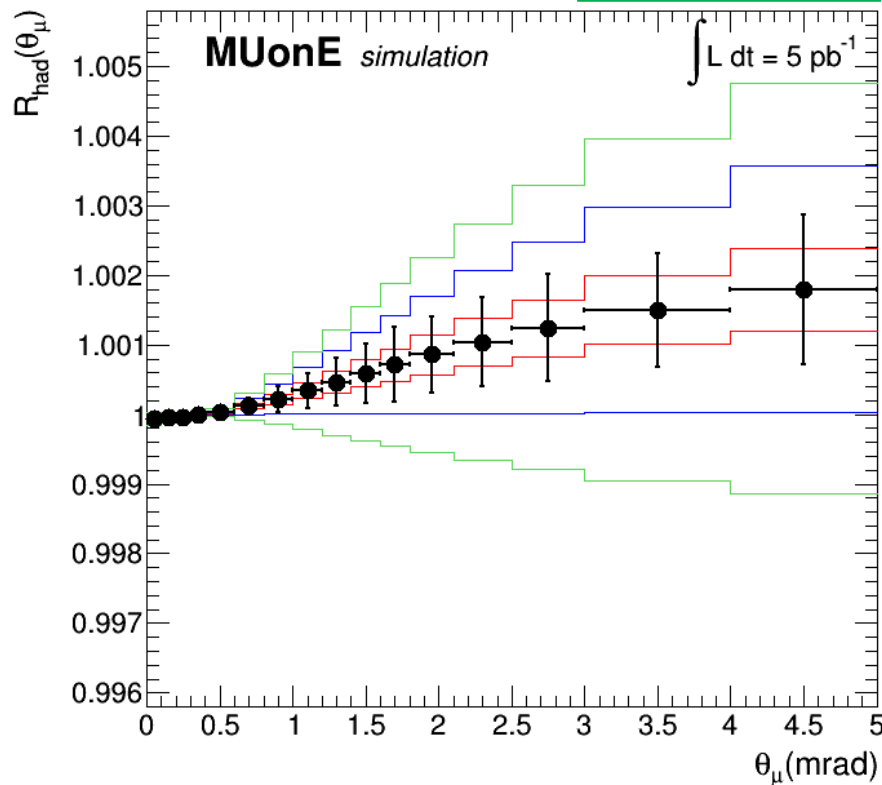
New GEANT4  
version 10.7  
(validation ongoing)

# Expected sensitivity of a First Physics Run

Expected integrated Luminosity with the Test Run setup with full beam intensity & detector efficiency  $\sim 1\text{pb}^{-1}/\text{day}$

In one week  $\sim 5\text{pb}^{-1} \rightarrow \sim 10^9$   $\mu\text{e}$  scattering events with  $E_e > 1$  GeV  
( $\theta_e < 30$  mrad)

[arXiv:2012.07016](https://arxiv.org/abs/2012.07016)



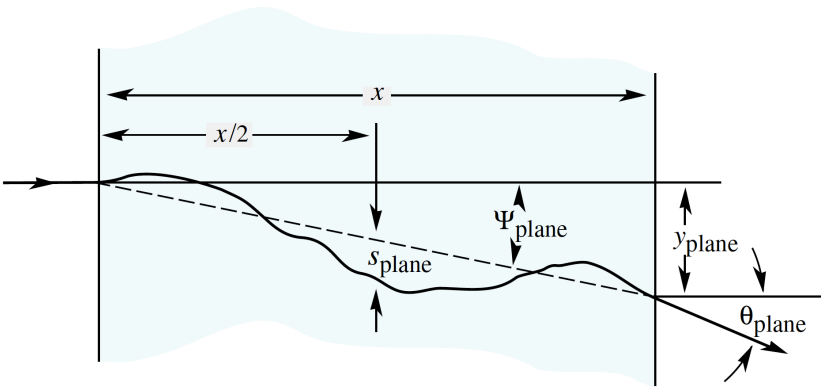
Initial sensitivity to the hadronic running of  $\alpha$ .

Pure statistical level:  $5.2\sigma$   
2D ( $\theta_\mu, \theta_e$ )  $K = 0.136 \pm 0.026$

Definitely we will have sensitivity to the leptonic running (ten times larger)

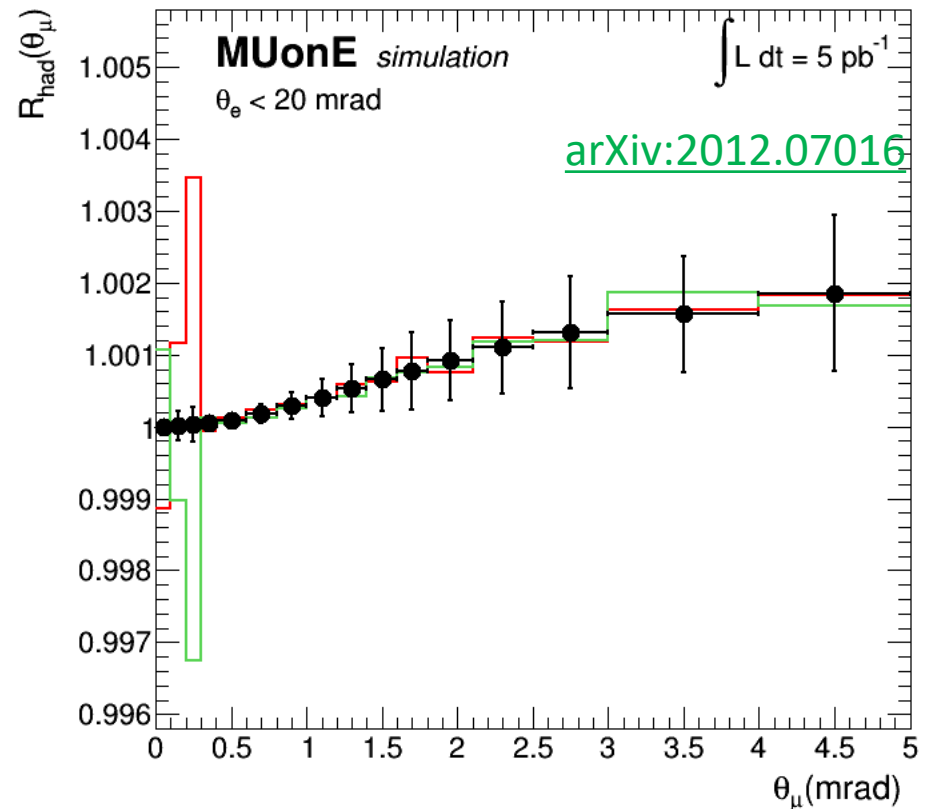
Template fit with just one fit parameter  $K = k/M$  in the  $\Delta\alpha_{\text{had}}$  parameterization.  
The other parameter fixed at its expected value:  $M = 0.0525 \text{ GeV}^2$

# Systematic Effects: Multiple Coulomb Scattering



*Particle trajectories disturbed:  
especially low-energy electrons*

Effects of a flat error of  $\pm 1\%$  on the  
core width of multiple scattering



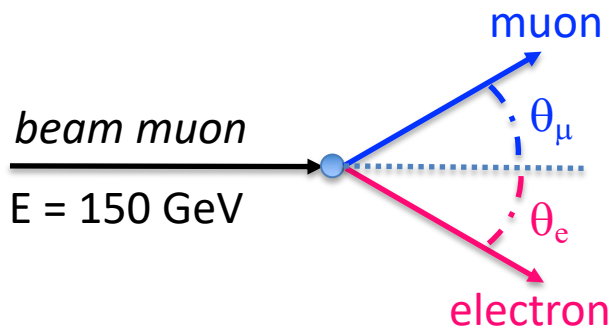
Multiple scattering previously studied in a Beam Test in 2017: [JINST 15 \(2020\) P01017](#)  
with 12–20 GeV electrons on 8-20 mm C targets

# Systematic Effects: Beam Energy scale

Time dependency of the beam energy profile has to be continuously monitored during the run:

- SPS monitor
  - COMPASS BMS
- } needed external infos

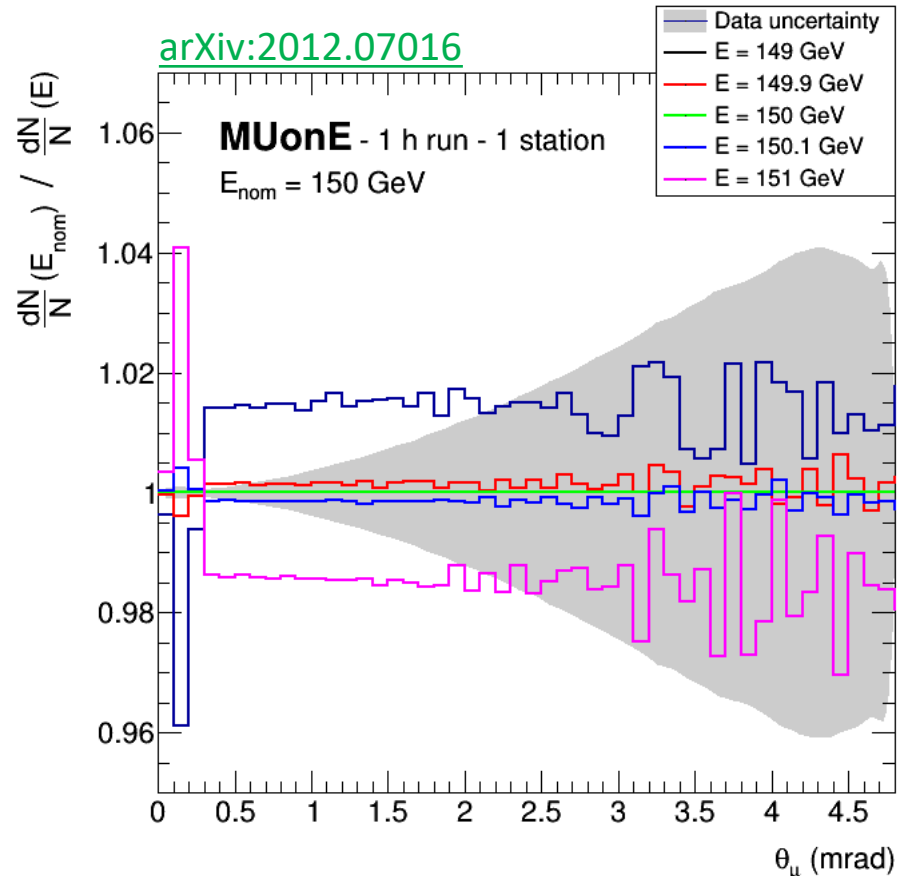
However, the absolute beam energy scale has to be calibrated by a physics process:  
kinematical method on elastic  $\mu e$  events



For equal angles:

$$\theta_\mu = \theta_e \equiv \theta \quad \theta \simeq \sqrt{\frac{2m_e}{E}}$$

Can reach <3 MeV uncertainty in a single station in less than one week  
From SPS E scale  $\sim 1\%$  : 1.5 GeV



Effect of a syst shift of the average beam energy on the  $\theta_\mu$  distribution: 1h run / 1 station

# Theory

Impressive progress

## STATUS: report of the **MUonE theory initiative**

*“Theory for muon-electron scattering @ 10ppm”, [P.Banerjee et al, Eur.Phys.J.C80\(2020\)591](#)*

**NLO exact calculation** including masses ( $m_\mu$ ,  $m_e$ ) and EWK corrections in a fully differential MC code  
[M.Alacevich et al, JHEP02\(2019\)155](#) cross-checked with independent calculation by Fael & Passera

## Full NNLO not yet available

- Two-loop master integrals ( $m_e=0$ ,  $m_\mu \neq 0$ )  
planar [P.Mastrolia et al, JHEP11\(2017\)198](#) and non-planar [S.Di Vita et al, JHEP09\(2018\)016](#)
- NNLO hadronic corrections: [M.Fael, M.Passera, Phys.Rev.Lett.122\(2019\)192001](#); [M.Fael, JHEP02\(2019\)027](#)
- Framework to recover leading  $m_e$  terms at NNLO from amplitudes calculated with massless electrons:  
T.Engel et al., [JHEP02\(2019\)118](#), [JHEP01\(2020\)085](#)
- Two independent fully exclusive NNLO MC codes, featuring the exact NNLO photonic corrections on the leptonic legs, including all mass terms: [C.Carloni Calame et al., arXiv:2007.01586](#); [P.Banerjee et al, arXiv:2007.01654](#)

➡ **VERY GOOD AGREEMENT between the two codes**

**Resummations (Parton shower and YFS) matched to (N)NLO fixed order under way**

## Study of possible contaminations from NEW physics on MUonE:

[A.Masiero, P.Paradisi and M.Passera, arXiv:2002.05418](#)

[P.S.Bhupal Dev et al., JHEP05\(2020\)53](#)

➔ **MUonE is NOT vulnerable !**



# Summary



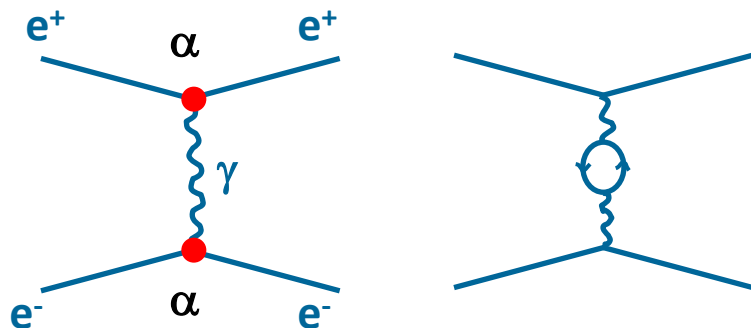
- **Long-standing puzzle of muon  $g-2$ :**
  - Experiment-Theory(SM) discrepancy  $3-4\sigma$
  - sensitive to BSM physics
  - Ongoing/future experiments will reduce the exp.error by a factor of 4
  - Theory error dominated by the Leading Hadronic contribution  $a_\mu^{\text{HLO}}$
- **MUonE experiment proposal:** measuring the running of  $\alpha_{\text{QED}}$  from the shape of the differential cross section for elastic scattering of  $\mu(150\text{GeV})$  on atomic electrons at the CERN SPS [Eur.Phys.J.C77\(2017\)139](#)
  - Getting  $a_\mu^{\text{HLO}}$  with a novel method integrating over the space-like region
  - Independent and complementary to the standard method integrating over the time-like region and to lattice QCD calculations
  - Competitive precision  $\sim 0.35-0.5\%$  on  $a_\mu^{\text{HLO}}$  allowing to better constrain the theory prediction , will help to solve the puzzle
- [Letter-Of-Intent SPSC-I-252](#) submitted to CERN in June 2019
- **CERN has recognized the fundamental interest and approved a Test Run to be carried out at the end of 2021, which should verify the detector design and assess the potential to achieve a competitive measurement, as a condition to move on towards the full-scale experiment.**
  - Main challenge: control of systematic effects at the level of the statistical precision
- Full-scale experiment foreseen during LHC Run3 (2022-2024) if results of the Test Run are satisfactory
- Delays in the Test Run preparation related to the Covid-19 pandemic, need to follow up the evolving situation

# BACKUP

# Measurement of $\Delta\alpha_{\text{had}}(t)$ spacelike at LEP

OPAL measurement: Bhabha scattering  
at small angle, with  $1.8 < -t < 6.1 \text{ GeV}^2$

about  $10^7$  events  
precision at the per mille level



$$\frac{d\sigma}{dt} = \frac{d\sigma^{(0)}}{dt} \left[ \frac{\alpha(t)}{\alpha_0} \right]^2 (1 + \varepsilon)(1 + \delta_\gamma) + \delta_Z$$

Born term for t-channel single  $\gamma$  exchange

$$\left( \frac{1}{1 - \Delta\alpha(t)} \right)^2$$

Effective coupling

factorized

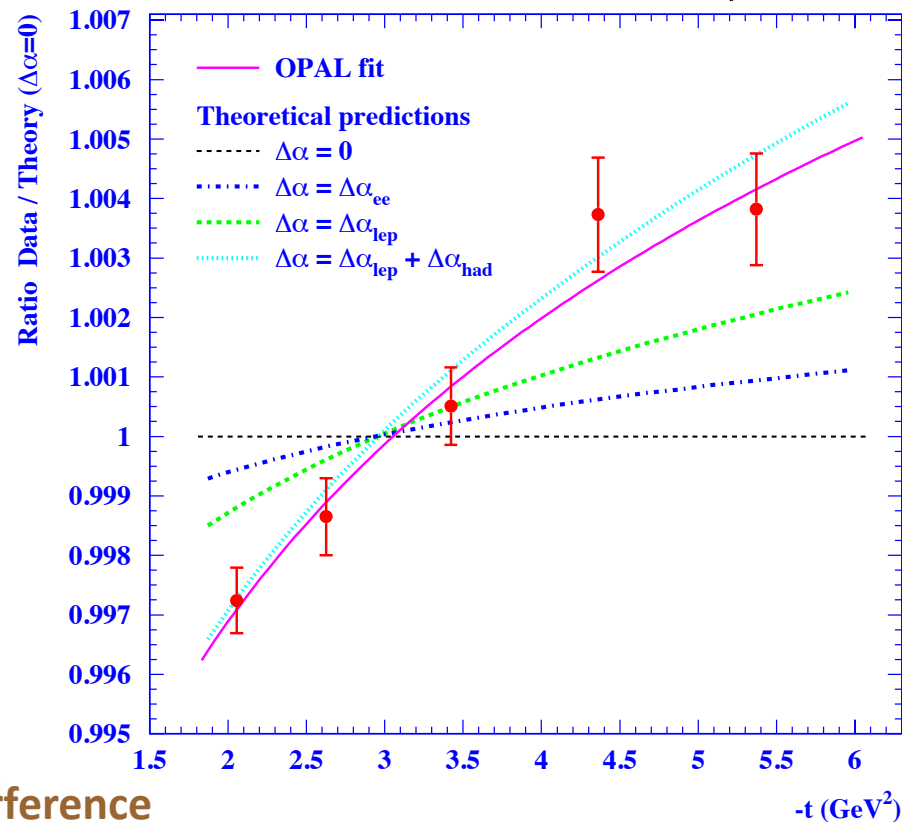
Photonic  
radiative  
corrections

Z interference  
correction

s-channel  $\gamma$  exchange  
correction

Eur.Phys.J.C45(2006)1

OPAL  $e^+e^- \rightarrow e^+e^-$   $\sqrt{s} \approx 91.2 \text{ GeV}$



Other measurements in the  
space-like region by L3, VENUS

# MUonE: alcune date

- **7 Settembre 2016**: prima presentazione al Workshop «*Physics Beyond Colliders*» al CERN
- **Paper** sottomesso a rivista, pubblicato in **Eur.Phys.J. C(2017)77**  
doi:10.1140/epjc/s10052-017-4633-z
- **12 Maggio 2017**: prima presentazione all'INFN CSN1
- **2017 Test Beam al CERN** (setup UA9)
  - Paper pubblicato in JINST. 15 (2020) P01017
- **2018 Test Beam al CERN (North Area)** – paper to be submitted
- **Maggio 2019**: INFN referees: T.Dorigo, T.Lari, P.Meridiani, S.My, F.Petrucchi
- **5 Giugno 2019**: **Letter of Intent**, presentata al SPSC, CERN-SPSC-2019-026 ,  
SPSC-I-252, <http://cds.cern.ch/record/2677471>
- **1 Ottobre 2019**: INFN CSN1 ha approvato un finanziamento di 105 KEu per il 2020, per la preparazione di un Pilot Run nel 2021
- **Ottobre 2019**: referees del CERN: U.Wiedemann, A.Ferrari, M.Bona
- **21 Gennaio 2020**: prima **presentazione MUonE nella sessione aperta del SPSC** (C.Matteuzzi)
- **22 Gennaio 2020**: decisione del SPSC: riconosciuto l'interesse fondamentale della proposta; approvato il Test Run nel 2021
- **20 Maggio 2020**: INFN CSN1 apre la sigla MUonE

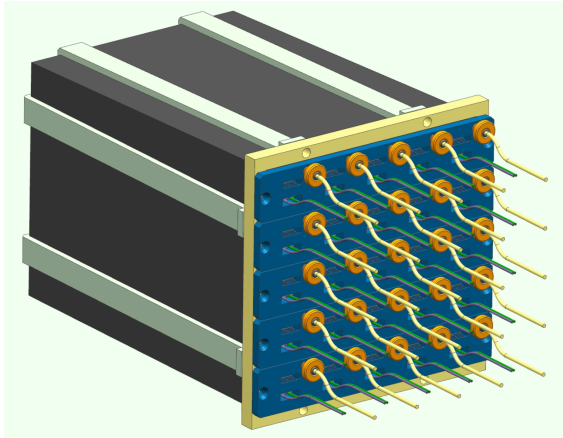
# Status / plans for the Test Run

- **Tracker:** delays (few months) in the procurement of the 2S modules (bottleneck: hybrids' production) due to Covid-19
  - Unlikely to have more than one MUonE tracking station fully integrated and ready for beam test in Fall 2021
  - Situation still subject to unpredictable changes
- **Calorimeter:** tight schedule but original plan still feasible
- **DAQ:** good progress, but partly related to the availability of tracker modules

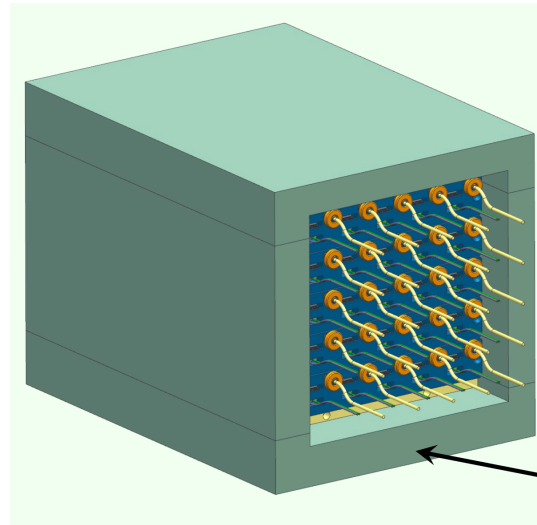
## MUonE plans to have the Test Run at the end of 2021 even with a partial setup

- *In this case with reduced objectives, mainly detector commissioning in real conditions of beam and environment*
- *If so we will consider the request of some time in early 2022, according to which conditions will have been realized in 2021.*

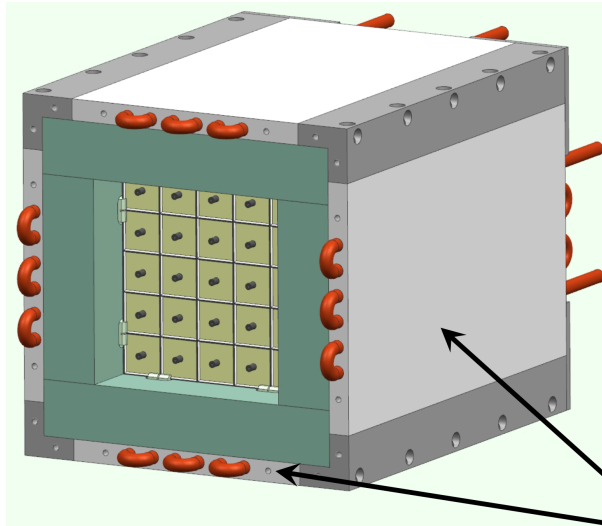
# Calorimeter (2)



**BACK Side**



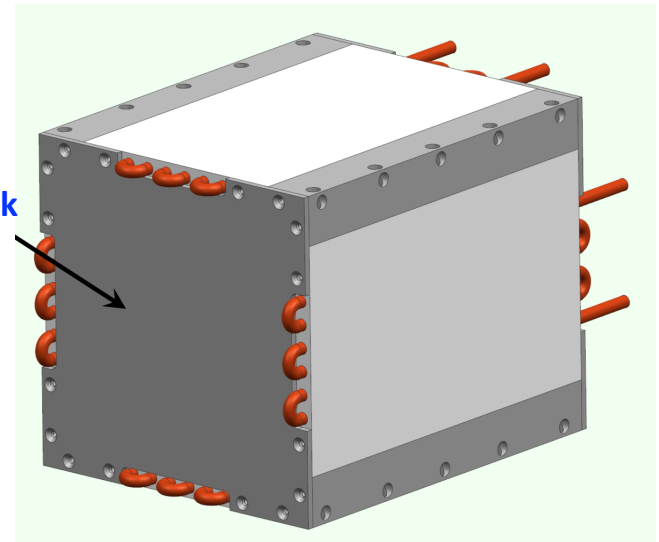
**Thermal insulator**



**FRONT Side**

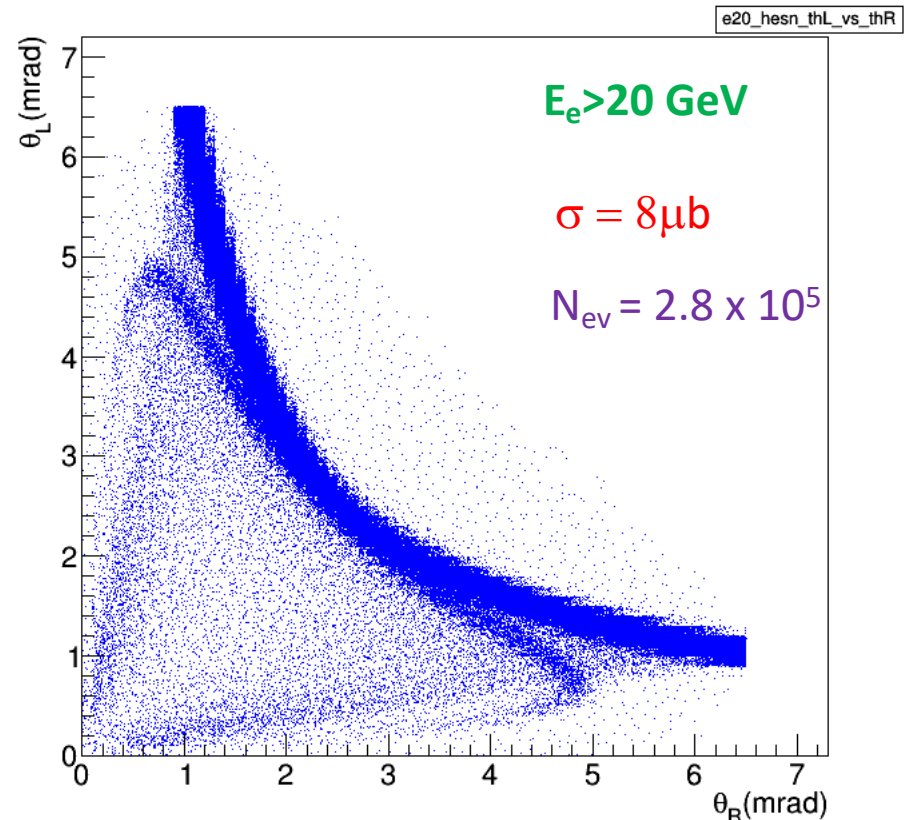
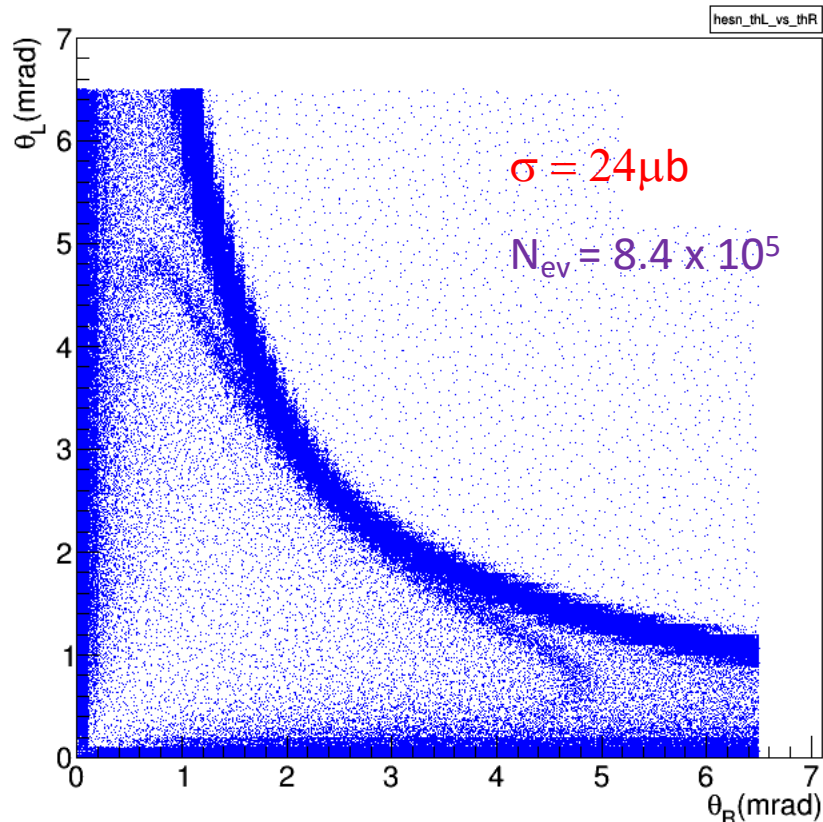
**Aluminum plates  
with embedded  
cooling pipes**

**5mm thick  
Al foil**



# 2D angular selection

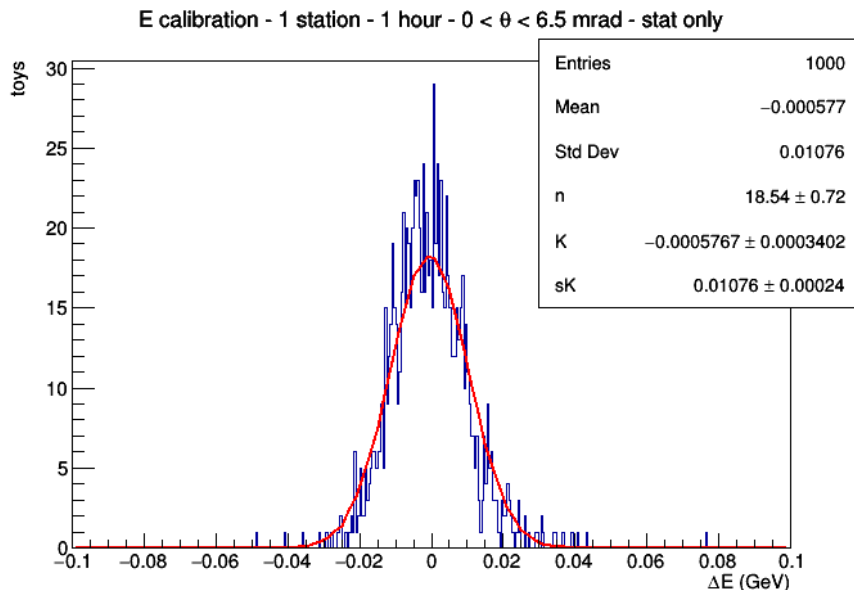
- $\theta_L, \theta_R < 6.5$  mrad
  - Additionally a possible calorimeter cut  $E_e > 20$  GeV





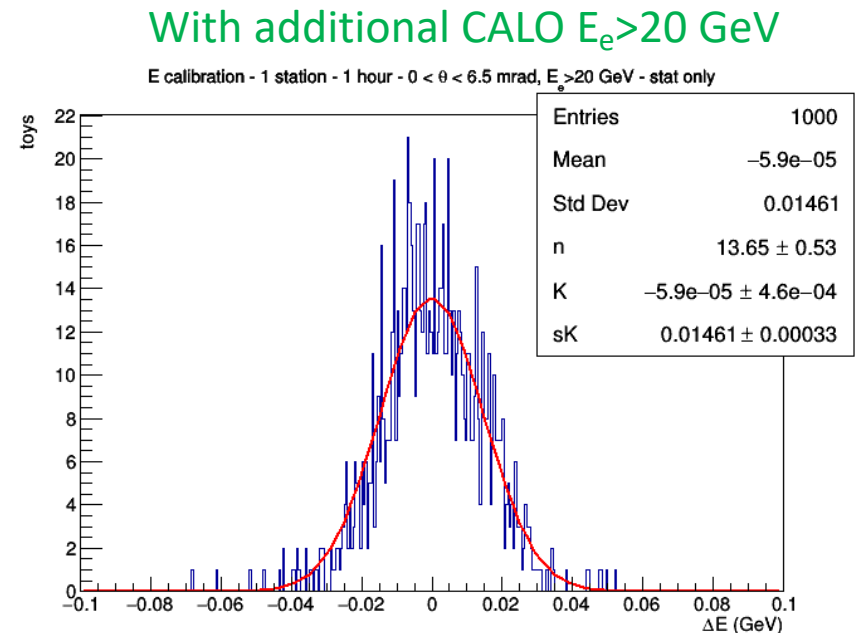
# Beam Energy scale: statistical accuracy

- Template fit of  $(\theta_L, \theta_R)$  with Beam energy as fit parameter in the range  $150 \text{ GeV} \pm 100 \text{ MeV}$ .
- Considering 1 hour run time in one station (1.5cm Be)  $\rightarrow \sim 35 \text{ nb}^{-1}$
- Angular selection:  $0 < (\theta_L, \theta_R) < 6.5 \text{ mrad} \rightarrow 24 \mu\text{b}$ 
  - With additional CALO  $E_e > 20 \text{ GeV} \rightarrow 8 \mu\text{b}$
- 1000 toys (each one with  $8.4 \times 10^5 / 2.8 \times 10^5$  events )



$\sigma \simeq 10.8 \text{ MeV}$

Bias  $\sim 0$



$\sigma \simeq 14.6 \text{ MeV}$

# Beam energy spread

- The beam energy profile has to be known. The M2 energy spread is  $\sim 3.75\%$ . Assuming to have the integrated profile corresponding to a given run from the BMS spectrometer. A (very) conservative assumption  $\sigma_{\text{BMS}} = (1.0 \pm 0.5)\%$  gives 19 MeV uncertainty on a 1-hour run on one station and a bias +16 / -5 MeV.
- The assumed BMS uncertainty is pessimistic, so this can certainly be better.
- If the event-by-event energy measurement would be available the impact of the energy spread could be effectively reduced by cutting the tails of the energy distribution, selecting only the central core.

# Beam energy calibration (1hour-1station)

Summary of statistical and systematic uncertainties and biases from different sources (educated guess), estimated for a 1-hour run on one station. Increasing the running time all the uncertainties scale with statistics. Biases stay constant.

		ANGULAR SELECTION $\theta_L, \theta_R < 6.5$ mrad		ANGULAR SELECTION and CALO $E_e > 20$ GeV	
		Bias (MeV)	Uncertainty (MeV)	Bias (MeV)	Uncertainty (MeV)
STATISTICAL	1hour / 1station	-	<b>11</b>	-	<b>15</b>
Longitudinal size	+10 $\mu$ m	<b>+3</b>	-		
	-10 $\mu$ m	<b>-3</b>	-		
Intrinsic Angular Resolution	+5%	<b>-3</b>	<b>8</b>		
	-5%	<b>+2</b>	<b>6</b>		
Beam Energy spread	BMS resolution +50%	<b>+16</b>	<b>14</b>		
	BMS resolution -50%	<b>-5</b>	<b>19</b>		
Multiple Coulomb Scattering	Target thickness +2%	<b>+7</b>	<b>15</b>	<b>+1</b>	<b>7</b>
	Target thickness -2%	<b>-3</b>	<b>13</b>	<b>0</b>	<b>5</b>
<b>TOTAL</b>		<b>+28 / -14</b>	<b>28</b>	<b>+22/-11</b>	<b>26</b>

# Beam energy calibration (4days-1station)

Summary of statistical and systematic uncertainties and biases from different sources (educated guess), estimated for 4-days (100 h) run on one station.

		ANGULAR SELECTION $\theta_L, \theta_R < 6.5$ mrad		ANGULAR SELECTION and CALO $E_e > 20$ GeV	
		Bias (MeV)	Uncertainty (MeV)	Bias (MeV)	Uncertainty (MeV)
STATISTICAL	1hour / 1station	-	<b>1.1</b>	-	<b>1.5</b>
Longitudinal size	+10 $\mu$ m	<b>+3</b>	-		
	-10 $\mu$ m	<b>-3</b>	-		
Intrinsic Angular Resolution	+5%	<b>-3</b>	<b>0.8</b>		
	-5%	<b>+2</b>	<b>0.6</b>		
Beam Energy spread	BMS resolution +50%	<b>+16</b>	<b>1.4</b>		
	BMS resolution -50%	<b>-5</b>	<b>1.9</b>		
Multiple Coulomb Scattering	Target thickness +2%	<b>+7</b>	<b>1.5</b>	<b>+1</b>	<b>0.7</b>
	Target thickness -2%	<b>-3</b>	<b>1.3</b>	<b>0</b>	<b>0.5</b>
<b>TOTAL</b>		<b>+28/-14</b>	<b>2.8</b>	<b>+22/-11</b>	<b>2.6</b>

# Acoplanarity

There are several possible quantities related to the deviation of a given event from perfect coplanarity. They have different properties and numerically are very different.

Let  $\mathbf{i}, \mathbf{m}, \mathbf{e}$  be unit vectors respectively along the directions of the incoming muon, the outgoing muon and the outgoing electron.

- 1) Triple product  $T = \mathbf{i} \cdot \mathbf{m} \times \mathbf{e}$

(used by NA7; geometrically the volume of the parallelepiped defined by the three vectors)

- 2) Angle between the incoming muon and the plane of the outgoing particles ( $\mathbf{m}, \mathbf{e}$ )

$$A = \frac{\pi}{2} - \cos^{-1} \left( \frac{\mathbf{i} \cdot \mathbf{m} \times \mathbf{e}}{|\mathbf{m} \times \mathbf{e}|} \right)$$

- 3) Angle between the scattering planes formed by the outgoing particles with the incoming muon

$$A_{\Phi} = \pm \left[ \pi - \cos^{-1} \left( \frac{(\mathbf{i} \times \mathbf{m}) \cdot (\mathbf{i} \times \mathbf{e})}{|\mathbf{i} \times \mathbf{m}| |\mathbf{i} \times \mathbf{e}|} \right) \right] \text{ for } \begin{cases} T > 0 \\ T < 0 \end{cases}$$

Notice:  $A_{\Phi}$  tests also that the outgoing electron and muon are directed on opposite sides in the transverse plane, while  $T$  and  $A$  do not depend on this.

- this can provide significantly different power in suppressing the backgrounds, in particular also the pair production

# Acoplanarity

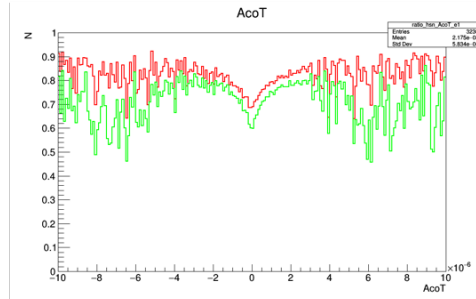
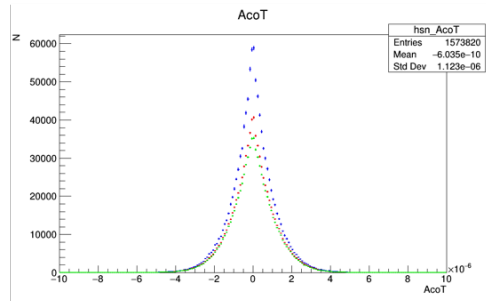
Blue: All Events

Red: Electron  $E > 1$  GeV

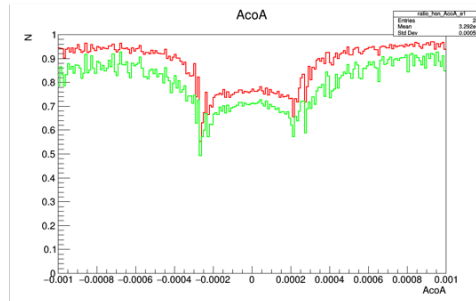
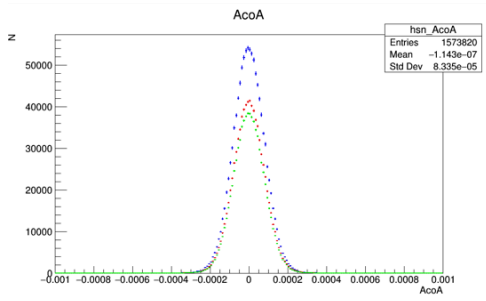
Green: Electron  $E > 2$  GeV

Ratios:  $E > X / \text{All}$

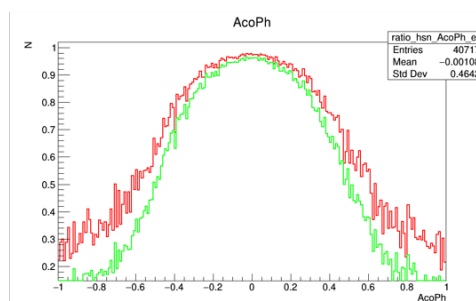
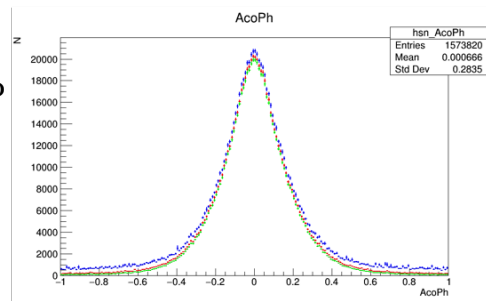
T



A



$A_\Phi$



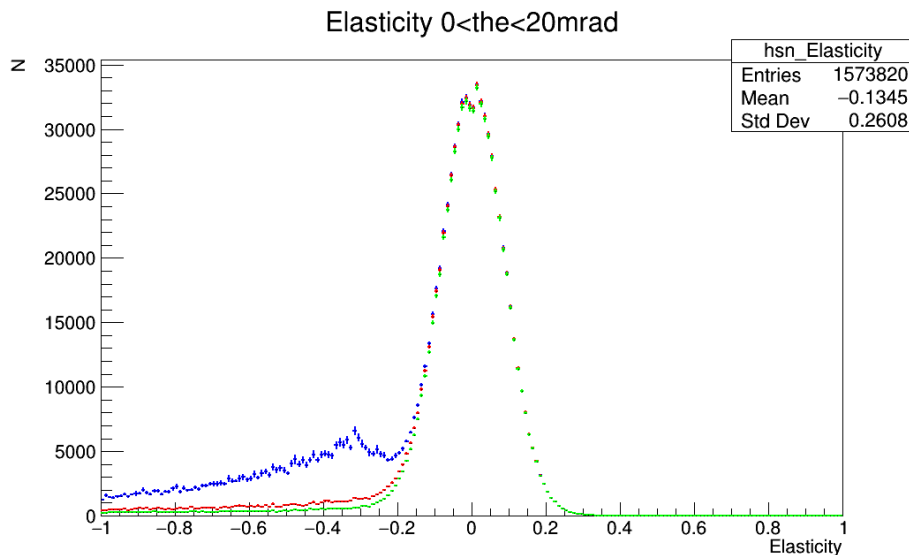
Energy cuts give unwanted pathological effects on T and A

Good behaviour

# Elasticity

The angular distance of a given event defined by the two angles  $P=(\theta_e, \theta_\mu)$  from the nearest point C on the curve corresponding to the elastic scattering (at a given c.m.s. energy) can be taken as a measurement of the elasticity of the event

- ideally for perfect elasticity  $D=0$



Blue: All Events

Red: Electron  $E > 1 \text{ GeV}$

Green: Electron  $E > 2 \text{ GeV}$

The left tail from radiative and detector smearing effects is effectively removed by energy cuts

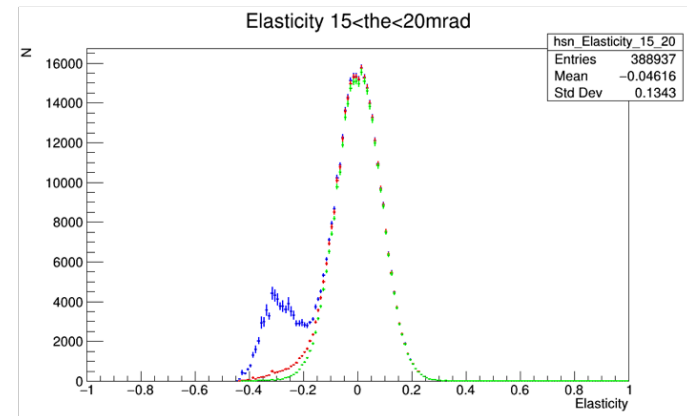
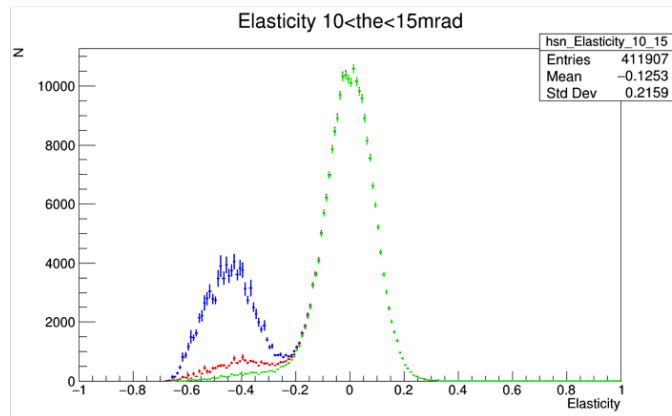
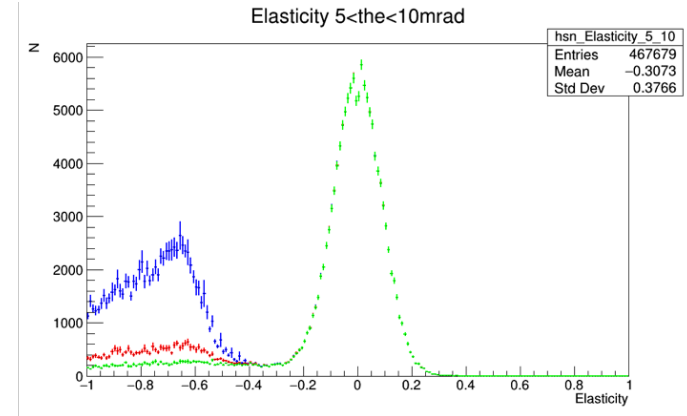
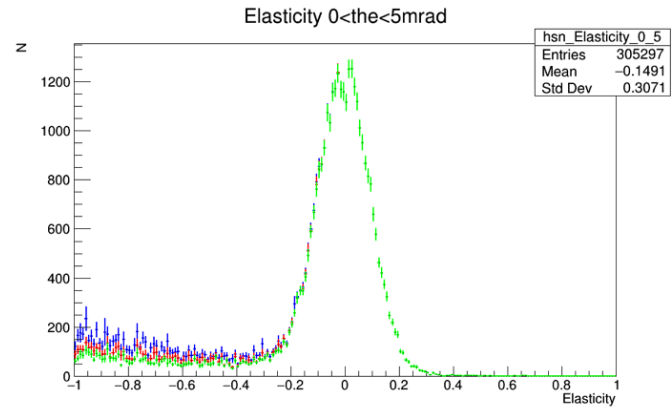


# Elasticity - angular regions

Blue: All Events

Red: Electron  $E > 1$  GeV

Green: Electron  $E > 2$  GeV



The left tail appears differently in different angular regions but is always suppressed effectively by energy cuts