



Search for a new light scalar in the $S \rightarrow \mu^+ \mu^-$ decay

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Outline

- Standard Model and beyond
- Models with a light scalar
- Scalar production
- Scalar decay
- Search for a light scalar in $K \rightarrow \pi S$, $S \rightarrow \mu\mu$
- Search for a light scalar in $B \rightarrow KS$, $S \rightarrow \mu\mu$
- Conclusions

Standard Model

Standard Model of Elementary Particles

		three generations of matter (elementary fermions)			three generations of antimatter (elementary antifermions)			interactions / force carriers (elementary bosons)	
		I	II	III	I	II	III		
mass		$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
	charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	0	0
spin		$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	0	0
		$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
LEPTONS		u up	c charm	t top	\bar{u} antiup	\bar{c} anticharm	\bar{t} antitop	g gluon	H higgs
	QUARKS	d down	s strange	b bottom	\bar{d} antidown	\bar{s} antistrange	\bar{b} antibottom	γ photon	
		e electron	μ muon	τ tau	e^+ positron	$\bar{\mu}$ antimuon	$\bar{\tau}$ antitau	Z Z ⁰ boson	GAUGE BOSONS VECTOR BOSONS
		ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	$\bar{\nu}_e$ electron antineutrino	$\bar{\nu}_\mu$ muon antineutrino	$\bar{\nu}_\tau$ tau antineutrino	W⁺ W ⁺ boson	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	$\bar{\nu}_e$ electron antineutrino	$\bar{\nu}_\mu$ muon antineutrino	$\bar{\nu}_\tau$ tau antineutrino	W⁻ W ⁻ boson		

Phenomena not explained by the Standard Model (SM)

- Neutrino masses/oscillations
- Baryon asymmetry of the Universe
- Dark Matter
- Dark Energy
- Gravity

Intrinsic difficulties

- Hierarchy problem
- Strong CP problem
- ...



SM is incomplete

Physics beyond the SM

Simplest model extension

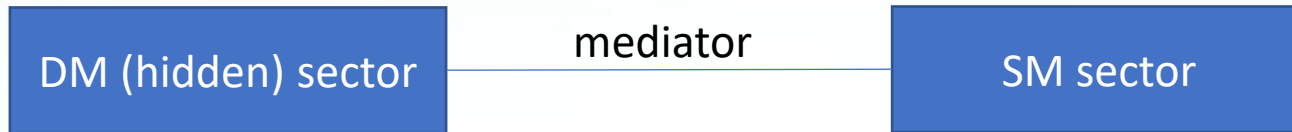
- Add one singlet scalar field
- New scalar couples to the Higgs boson (scalar portal, Higgs portal)
- Mass could be $O(\text{GeV})$


What could explain:

- Dark Matter
- Inflation
- Hierarchy problem

Dark Matter and a new scalar

Hidden sector concept

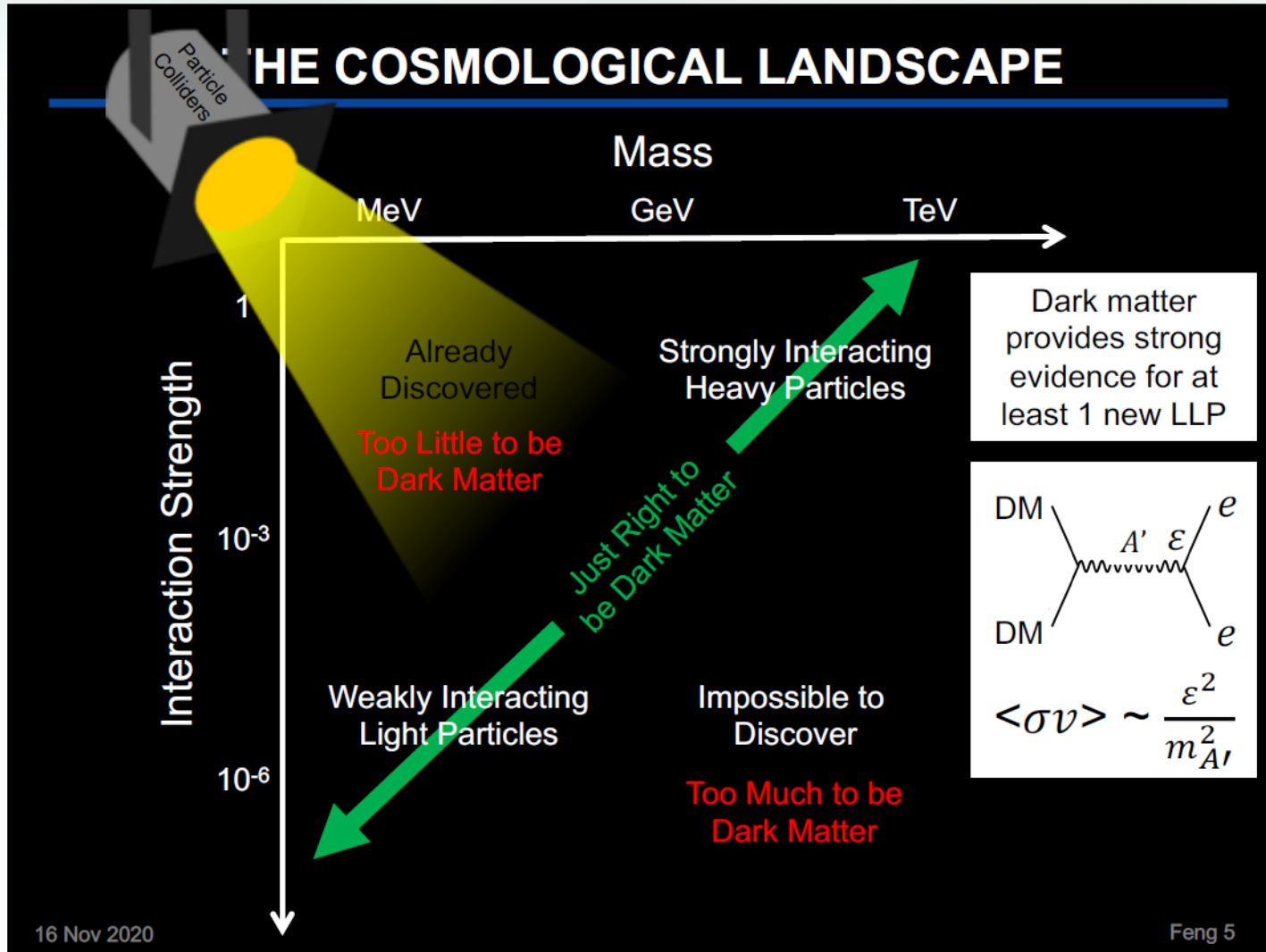


	Portal	Coupling
	Vector: Dark Photon, A'	$-\frac{\varepsilon}{2 \cos \theta_W} F'_{\mu\nu} B^{\mu\nu}$
	Scalar: Dark Higgs, S	$(\mu S + \lambda_{HS} S^2) H^\dagger H$
	Fermion: Heavy Neutral Lepton, N	$y_N L H N$
	Pseudo-scalar: Axion, a	$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, \frac{a}{f_a} G_{i,\mu\nu} \tilde{G}_i^{\mu\nu}, \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$



Searching for a mediator: portal towards the light DM

New scalar: mass and lifetime

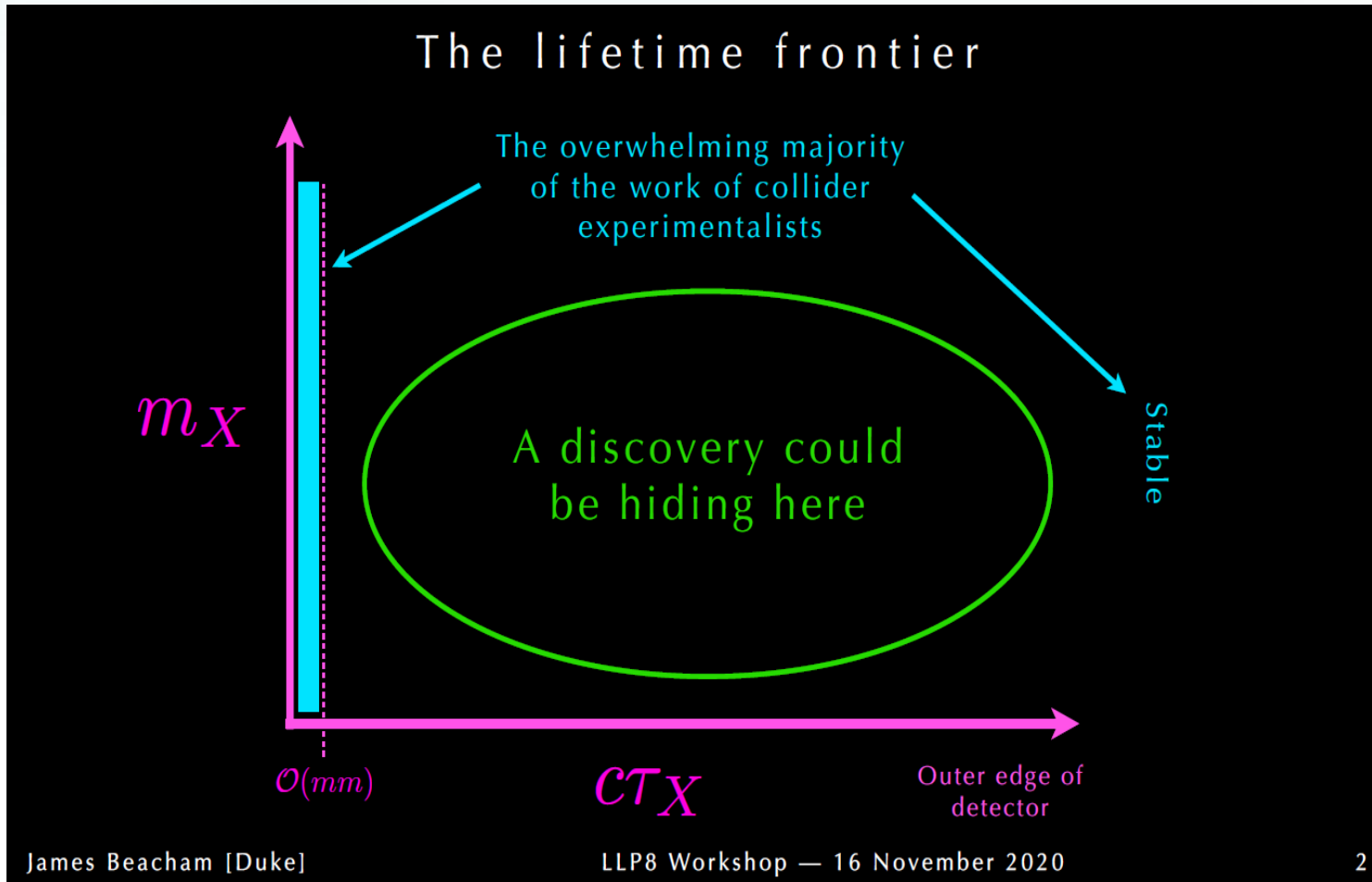


- $M \sim O(\text{GeV})$
- Interaction strength $\sim 1/\tau$



Long lived particle (LLP)

New scalar: lifetime



Sensitivity to large τ needed



- Experiments with a long decay volume
- High intensity beams

Hierarchy problem and a new scalar

Hierarchy problem for the Higgs mass:

- Measured mass = bare mass + quantum corrections
- Good model: small quantum corrections
- SM: very big quantum corrections



- Fine tuning of corrections to obtain the measured mass



Other formulation of the same problem:

- Why is the Higgs mass so small compared to the Planck mass?

Possible solution: relaxion mechanism

- Dynamic solution of the hierarchy problem
- Relaxion is a [scalar particle](#)
- Evolution of the relaxion field reduces the initially big Higgs mass to the measured value

Inflaton model

Models with a light inflaton

- Light scalar field
- Inflaton feebly coupling to the Higgs boson
- Can explain the inflation
- Can be incorporated into the vMSM model (explains neutrino masses, BAU, DM)
- 2 free parameters (mass and coupling constant θ)
- (m, θ) values limited from the cosmology
- Most sensitive processes: K and B decays



will be used here as the reference model

Model lagrangian

arXiv: 0912.0390

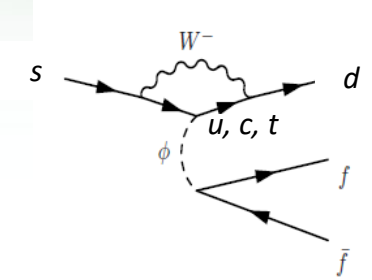
$$\mathcal{L}_{XSM} = \mathcal{L}_{SM} + \mathcal{L}_{XN} ,$$

$$\mathcal{L}_{XN} = \frac{1}{2} \partial_\mu X \partial^\mu X + \frac{1}{2} m_X^2 X^2 - \frac{\beta}{4} X^4 - \lambda \left(H^\dagger H - \frac{\alpha}{\lambda} X^2 \right)^2$$

Scalar production in K decays

Kaon decay $K \rightarrow \pi S$

- 3 terms in the amplitude
- Dominant contribution from the FCNC d-s transition
- Non-negligible interference with 4-fermion operator contribution



amplitude

$$A(K^+ \rightarrow \pi^+ \chi) \simeq \theta \frac{M_K^2}{v} \left\{ \underbrace{\gamma_1 \frac{1-\kappa}{2} \left(1 - \frac{m_\chi^2 - M_\pi^2}{M_K^2}\right)}_{\text{small}} - \underbrace{\gamma_2(1-\kappa)}_{\text{negligible}} + \frac{1}{2} \frac{3G_F\sqrt{2}}{16\pi^2} \sum_{i=c,t} V_{id}^* m_i^2 V_{is} \right\} \quad A(K_L \rightarrow \pi^0 \chi) = -A(K^+ \rightarrow \pi^+ \chi)$$

t-quark contribution dominant

Branching ratio

$$\text{Br}(K^+ \rightarrow \pi^+ \chi) = \frac{1}{\Gamma_{\text{total}}(K^+)} \frac{|A(K^+ \rightarrow \pi^+ \chi)|^2}{16\pi M_K} \frac{2|\mathbf{p}_\chi|}{M_K} \approx 1.3 \times 10^{-3} \cdot \left(\frac{2|\mathbf{p}_\chi|}{M_K}\right) \theta^2$$

arXiv: 0912.0390

BR($K \rightarrow \pi S$) proportional to θ^2

Scalar production in B decays

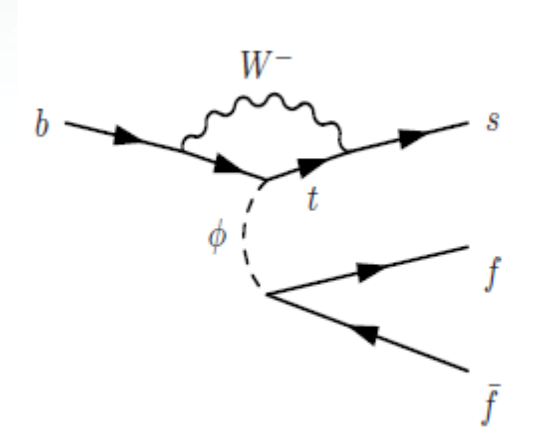
B decays $B^0 \rightarrow K^{0*} S$ and $B^+ \rightarrow K^+ S$:

- FCNC b-s transition

t-quark dominant contribution

$$\mathcal{L}_{\phi sb} = g_{\phi sb} \phi \bar{s}_L b_R + \text{h.c.},$$

$$g_{\phi sb} = \frac{s_\theta m_b}{v} \frac{3\sqrt{2} G_F m_t^2 V_{ts}^* V_{tb}}{16\pi^2}$$



Decay width

BR($B \rightarrow K S$) proportional to θ^2

$$\Gamma_{B \rightarrow K^{(*)} \phi} = |g_{\phi sb}|^2 \left| \langle K^{(*)} | \bar{s}_L b_R | B \rangle \right|^2 \frac{\lambda_{B, K^{(*)} \phi}^{1/2}}{16\pi m_B}$$

$$|\langle K^* | \bar{s}_L b_R | B \rangle|^2 = \frac{1}{4} \frac{m_B^4 \lambda_{B, K^{(*)} \phi}}{(m_b + m_s)^2} A_{K^*}^2$$

$$|\langle K | \bar{s}_L b_R | B \rangle|^2 = \frac{1}{4} \frac{(m_B^2 - m_K^2)^2}{(m_b - m_s)^2} f_K^2$$

$$A_{K^*} = \frac{1.36}{1 - q^2/27.9 \text{ GeV}^2} - \frac{0.99}{1 - q^2/36.8 \text{ GeV}^2}$$

$$f_K = \frac{0.33}{1 - q^2/37.5 \text{ GeV}^2}$$

arXiv: 1809.01876

Scalar decay

Decay width to leptons:

$$\Gamma(\phi \rightarrow \bar{\ell}\ell) \equiv \Gamma_{\bar{\ell}\ell} = \frac{s_\theta^2 G_F m_\phi}{4\sqrt{2}\pi} m_\ell^2 \beta_\ell^3$$

$$\beta_\ell = \sqrt{1 - 4m_\ell^2/m_\phi^2}$$

S → ee:

- Negligible for $m > 2 \cdot m_\mu$ wrt muon channel
- Bkg from photon conversion

S → μμ

- dominant for $m < 2 \cdot m_\pi$
- Non-negligible BR for $m > 2 \cdot m_\pi$

S → γγ

- negligible

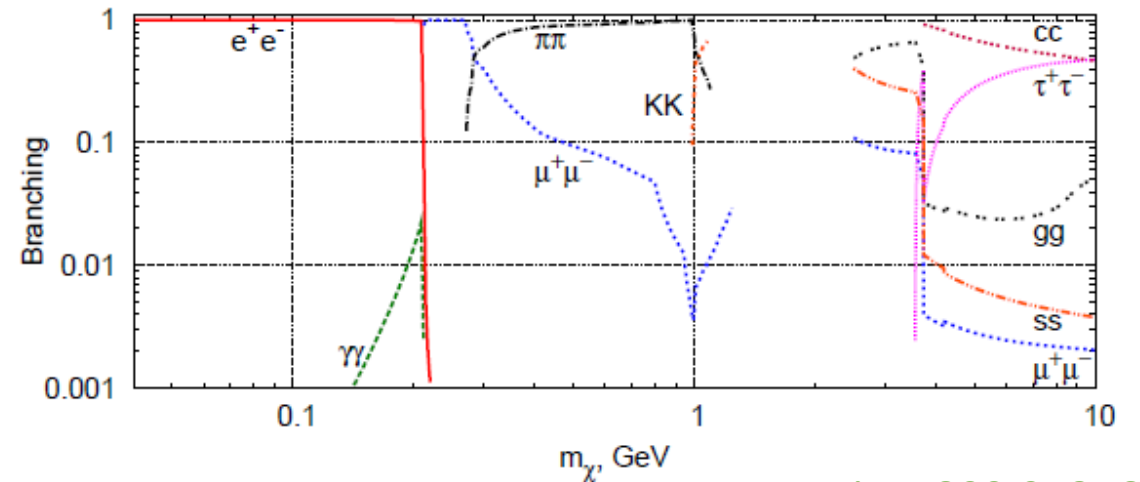
S → π π: non-trivial theoretical calculations

- ChPT for $m < 1.3$ GeV
- Dispersive analysis for $1.3 < m < 2$ GeV
- Perturbative spectator model for $m > 2$ GeV



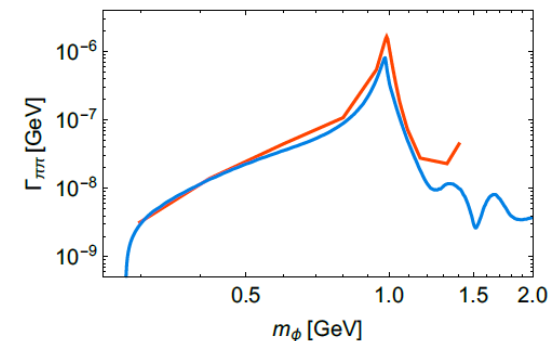
Muon channel seems to be the cleanest

→ $\Gamma(S \rightarrow \dots)$ proportional to θ^2



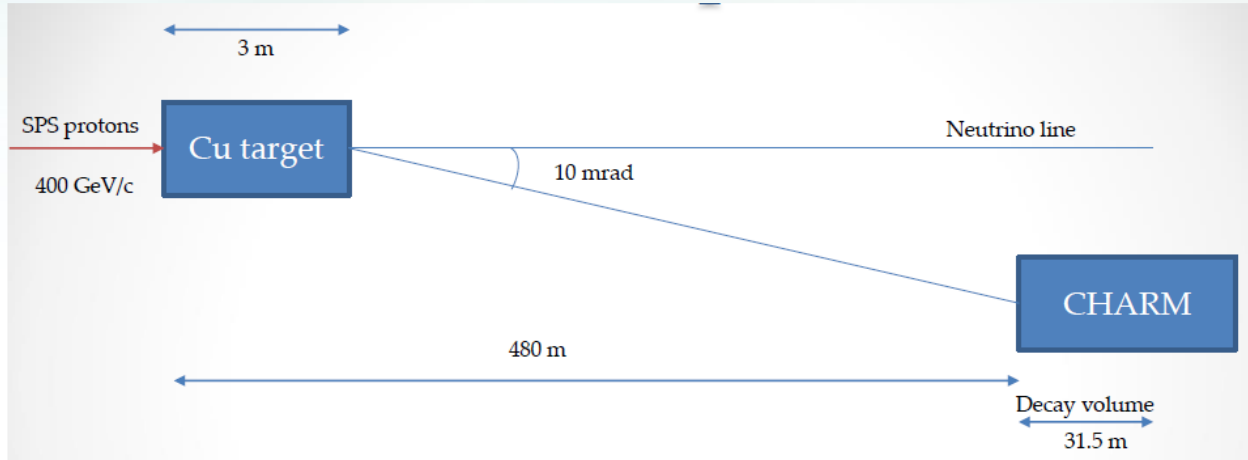
arXiv: 0912.0390

Donoghue this work arXiv: 1809.01876



Scalar search in CHARM

CHARM setup



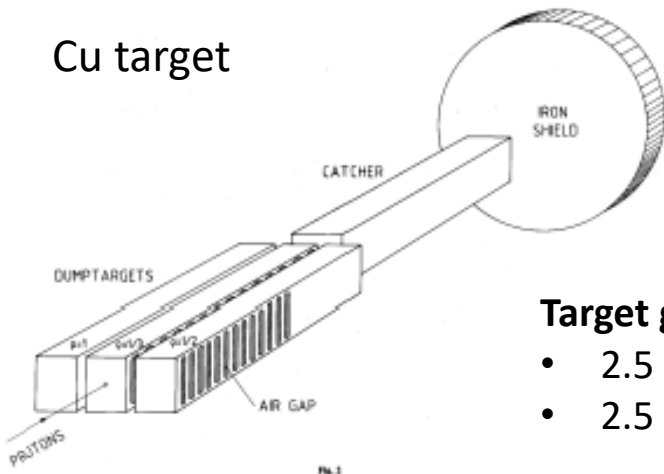
Scalar production @ CHARM:

- B decays
- K decays

Kaon absorption:

- $\lambda = 13$ cm in Cu
- Effective density smaller due to air gaps
- No kaons escape the target but some of them decay

Cu target



Target geometry

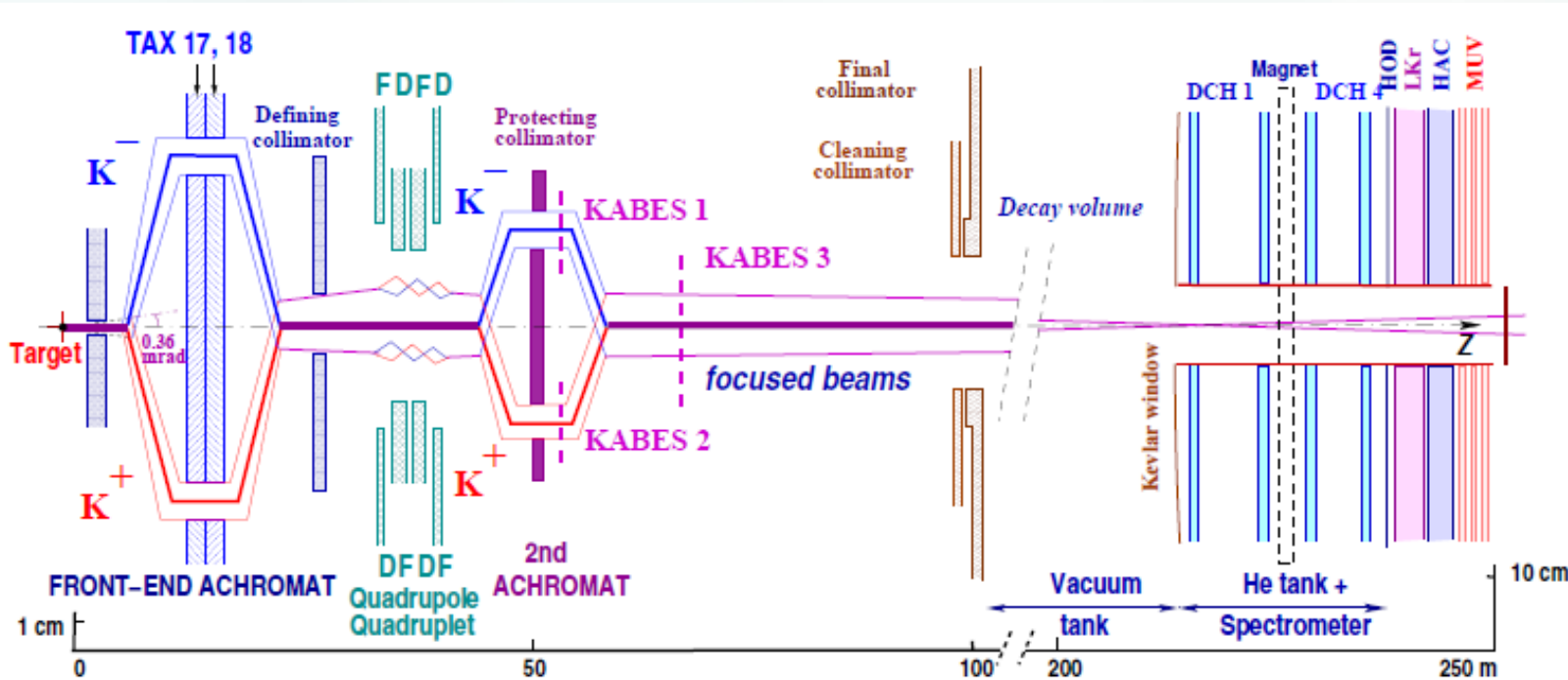
- 2.5 cm Cu + 2.5 cm air ($1.7 \cdot 10^{18}$ POT)
- 2.5 cm Cu + 5 cm air ($0.7 \cdot 10^{18}$ POT)

Data collected:

- $2.4 \cdot 10^{18}$ POT
- Signatures: $S \rightarrow \gamma\gamma$, $S \rightarrow \mu\mu$, $S \rightarrow ee$
- No signal candidates observed

Scalar search in NA48/2

NA48 setup

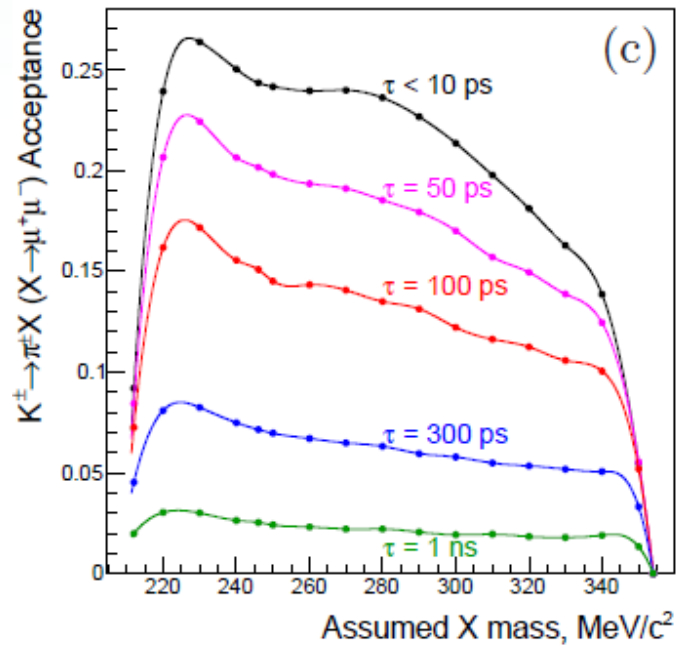


Scalar search in NA48/2

- $K^{\pm} \rightarrow \pi^{\pm} S$, $S \rightarrow \mu^+ \mu^-$
- 2003-2004 data
- Only prompt decays accessible due to the trigger configuration

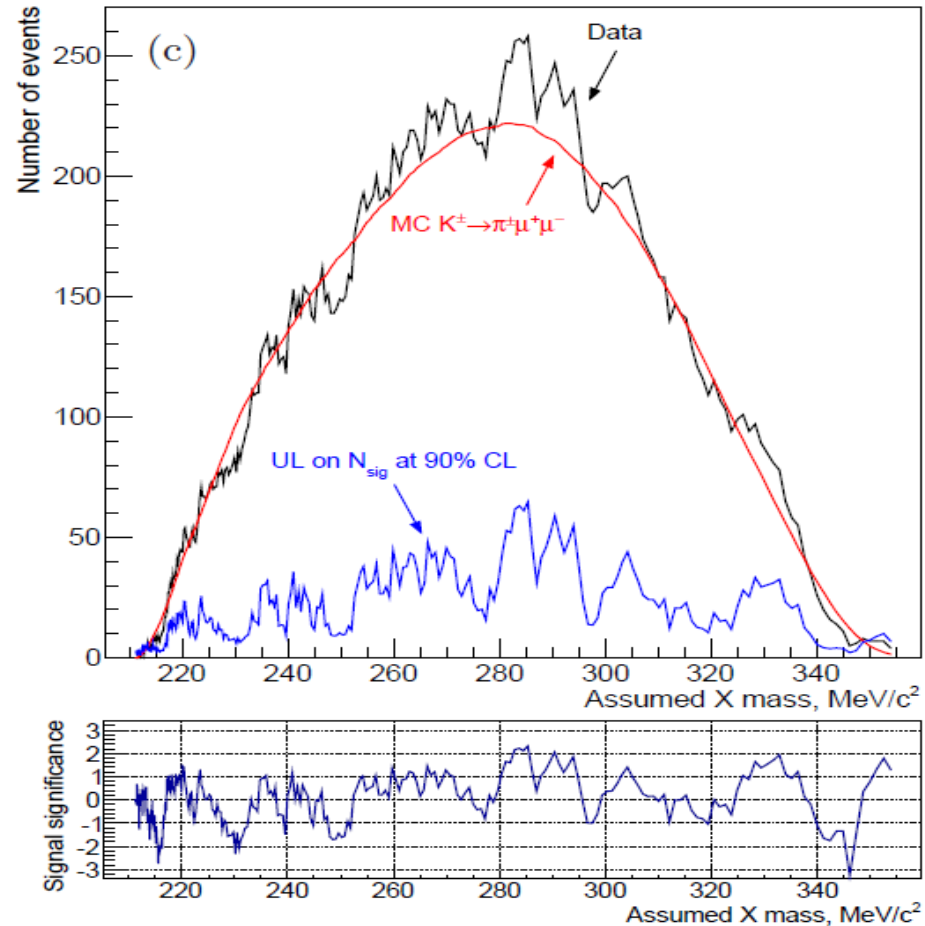
Limits from NA48/2

Signal acceptance



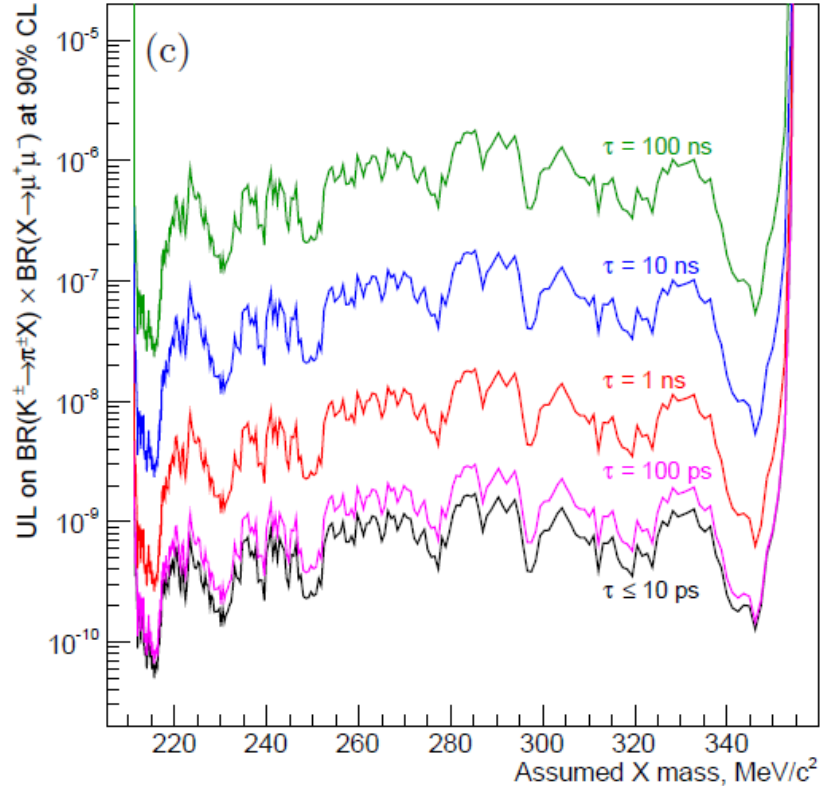
arXiv: 1612.04723

Mass scan

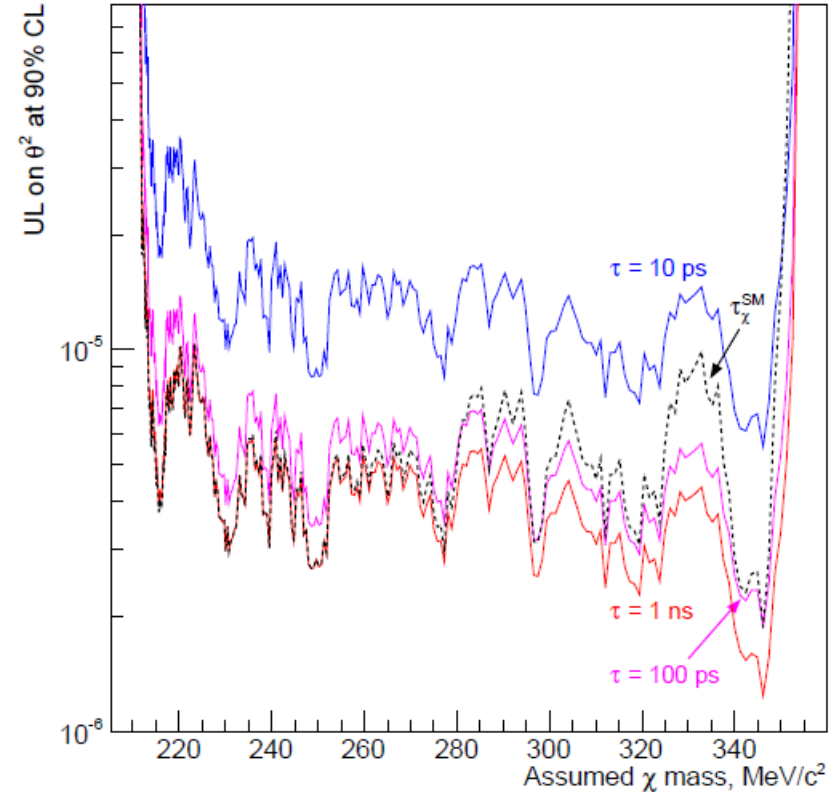


Limits from NA48/2

UL on BR



UL on θ^2



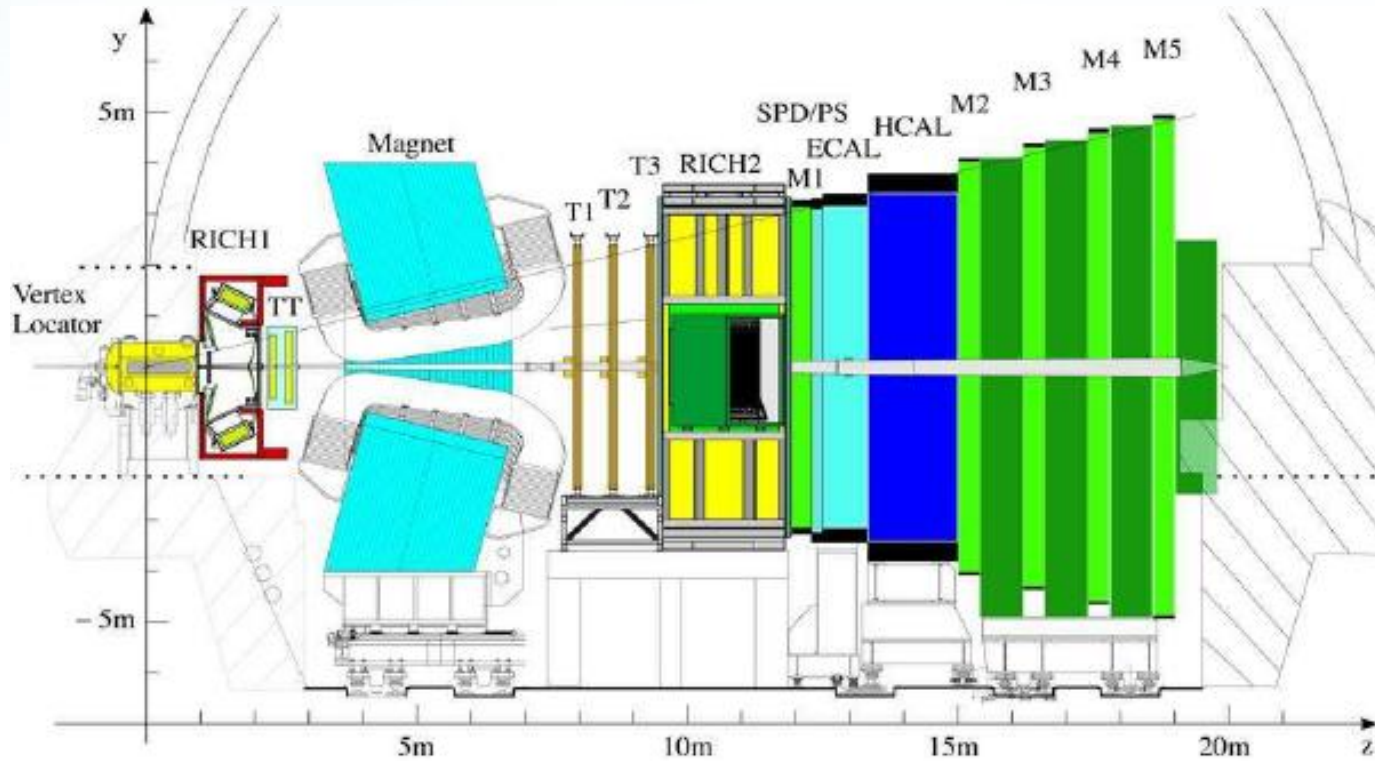
Prompt analysis: loose UL's for large τ

arXiv: 1612.04723

$$\theta^2 = \sqrt{\frac{8\pi\hbar v^2}{\alpha_\chi}} \sqrt{\frac{\mathcal{B}(K^\pm \rightarrow \pi^\pm \chi)\mathcal{B}(\chi \rightarrow \mu^+ \mu^-)}{\tau_\chi M_\chi^3 \lambda^{\frac{1}{2}}(1, r_\pi^2, r_\chi^2) \lambda^{\frac{1}{2}}(1, \tilde{\rho}_\mu^2, \tilde{\rho}_\mu^2) \tilde{\chi}_{\mu\mu}}}$$

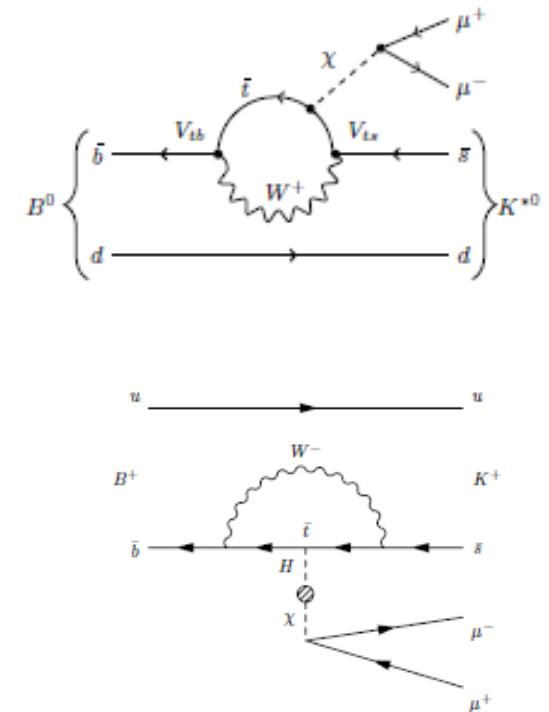
Scalar search in LHCb

LHCb setup



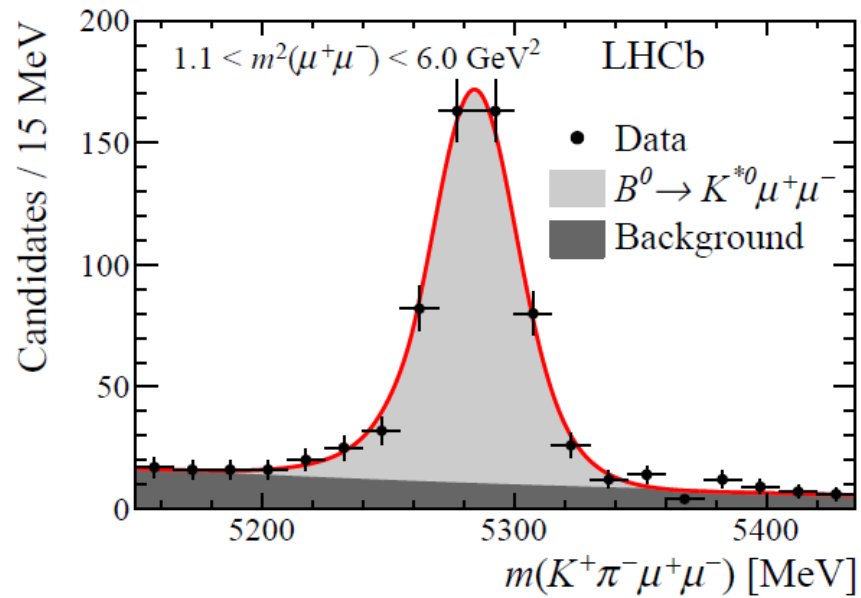
Scalar search in LHCb

- Run 1 data, 3 fb^{-1}
- $B^0 \rightarrow K^* S, S \rightarrow \mu^+ \mu^-$
- $B^+ \rightarrow K^+ S, S \rightarrow \mu^+ \mu^-$

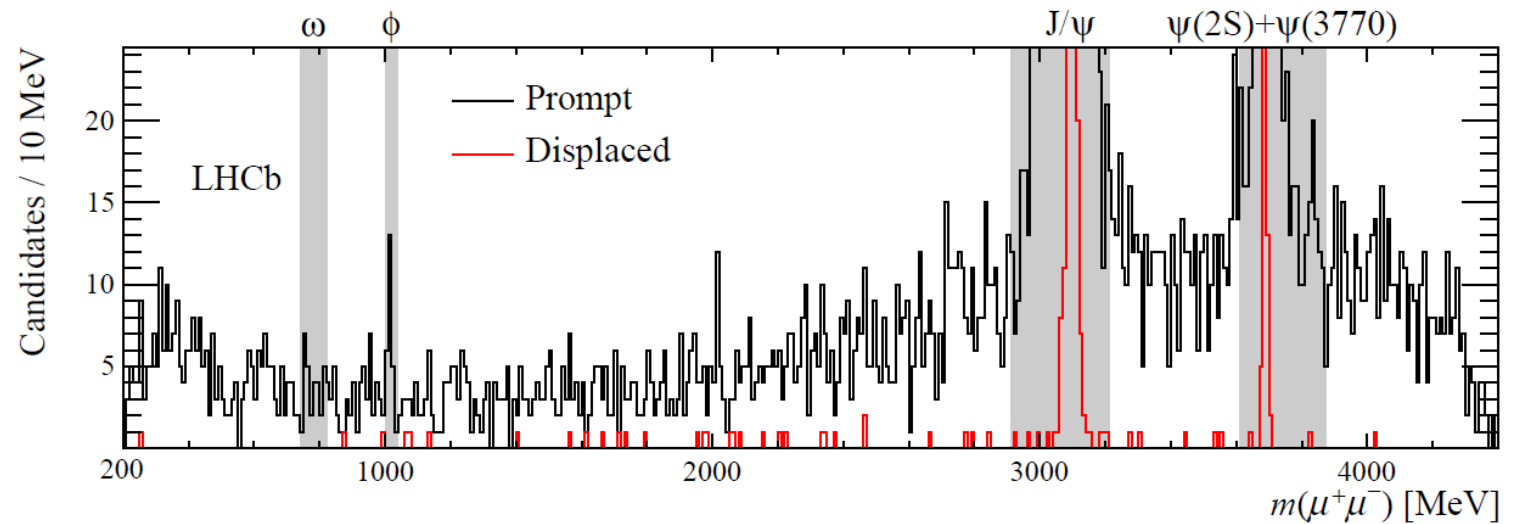


Scalar search in LHCb: B^0

invariant mass for the normalization $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



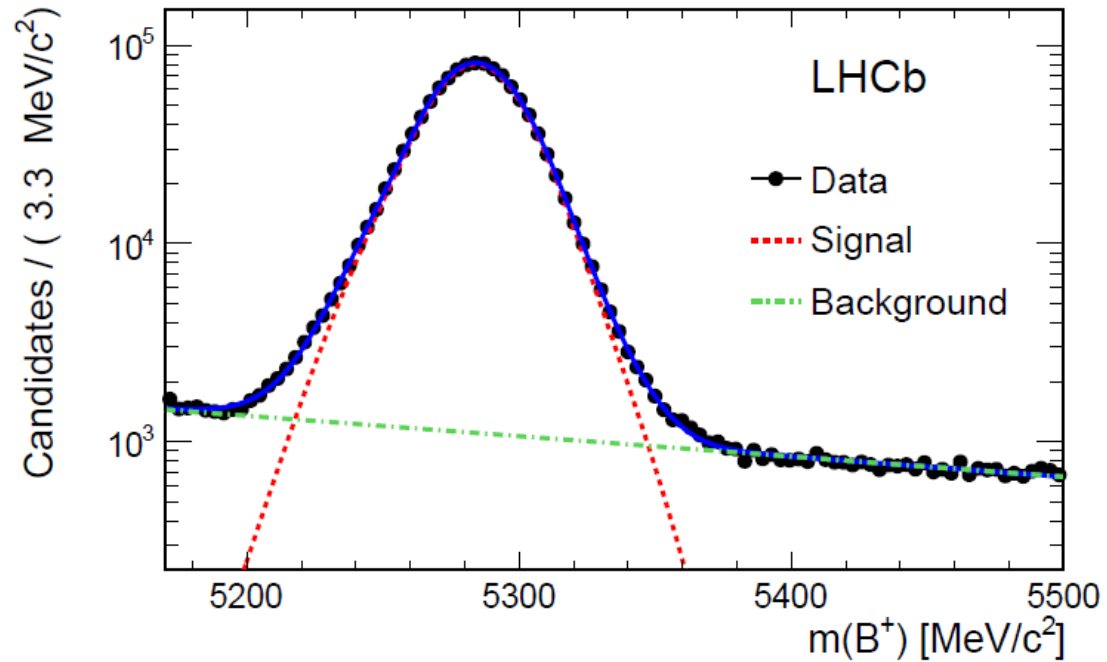
Mass scan



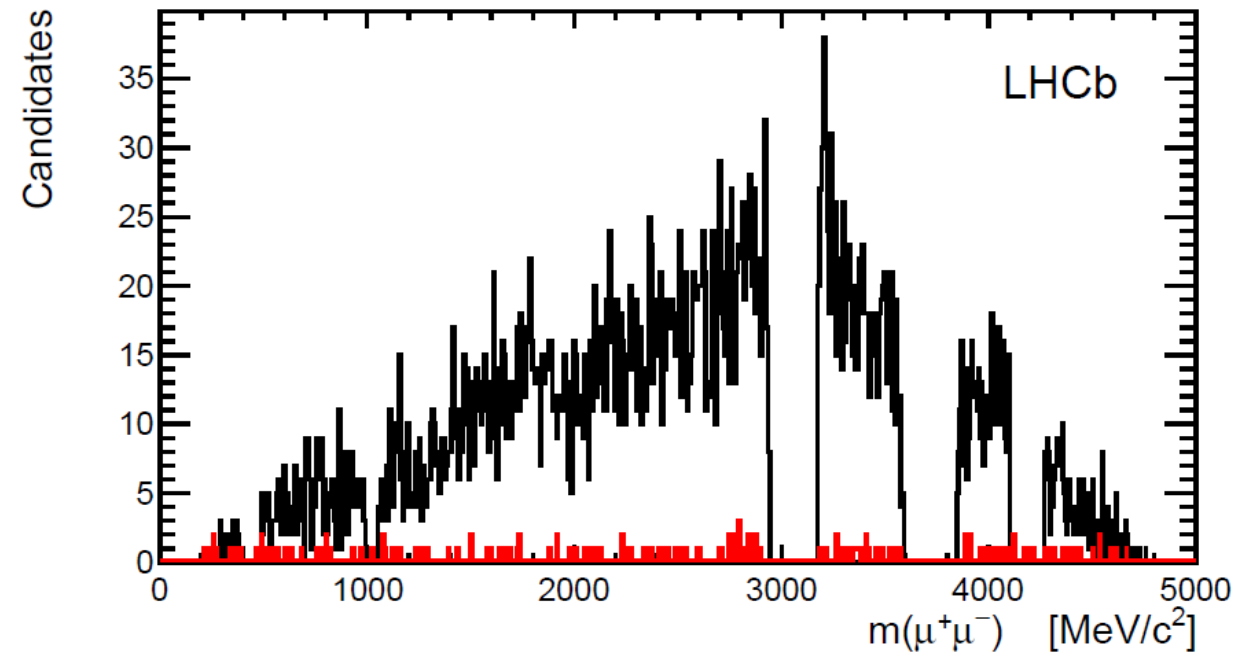
arXiv: 1508.04094

Scalar search in LHCb: B^+

Invariant mass for the normalization $B^+ \rightarrow K^+ J/\psi(\mu^+ \mu^-)$



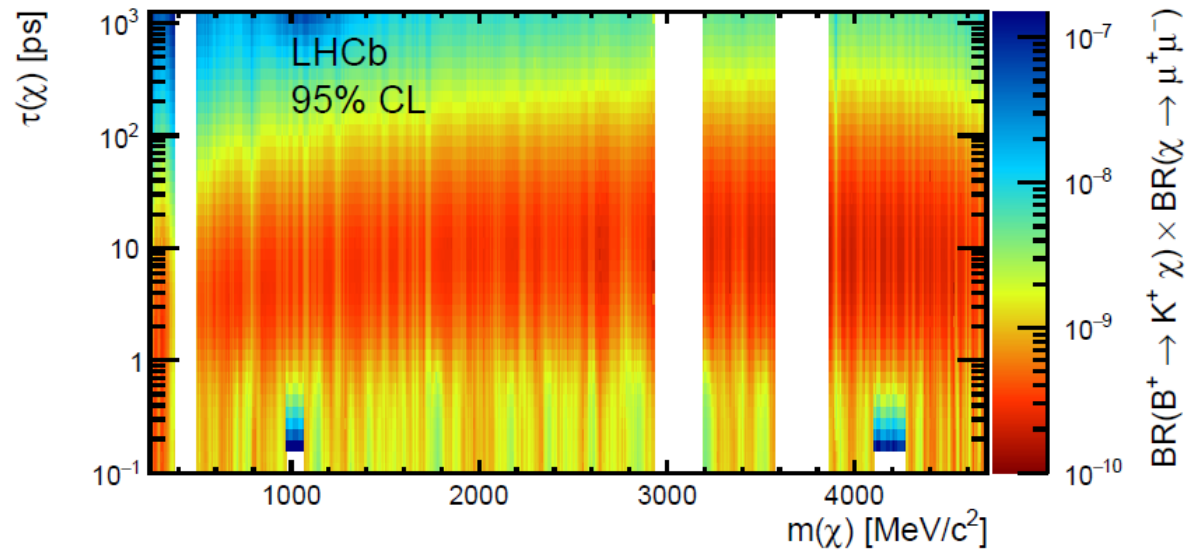
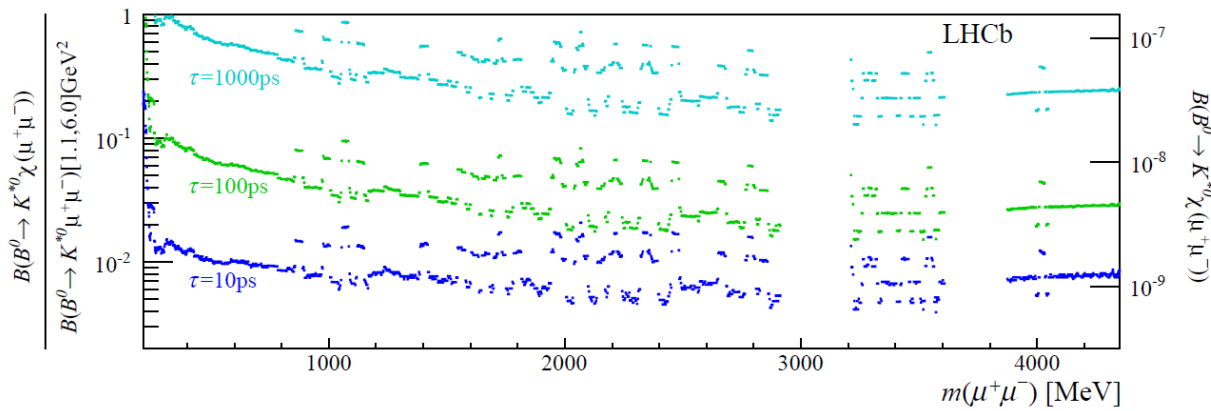
Mass scan



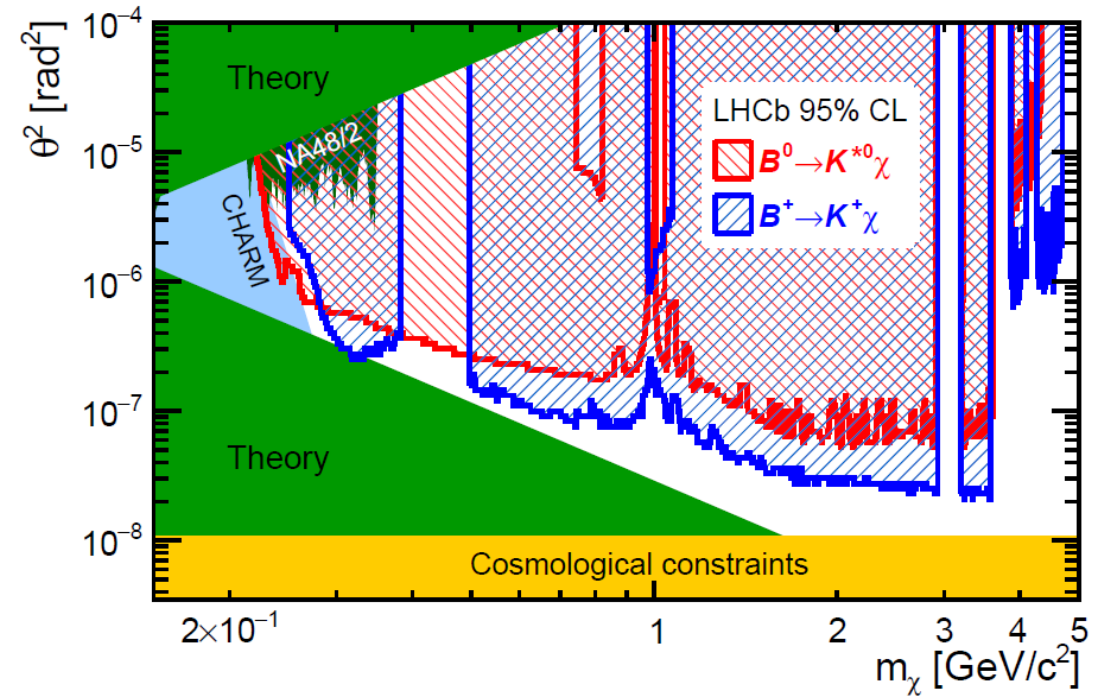
[arXiv: 1612.07818](https://arxiv.org/abs/1612.07818)

Limits from LHCb

Limits on BR



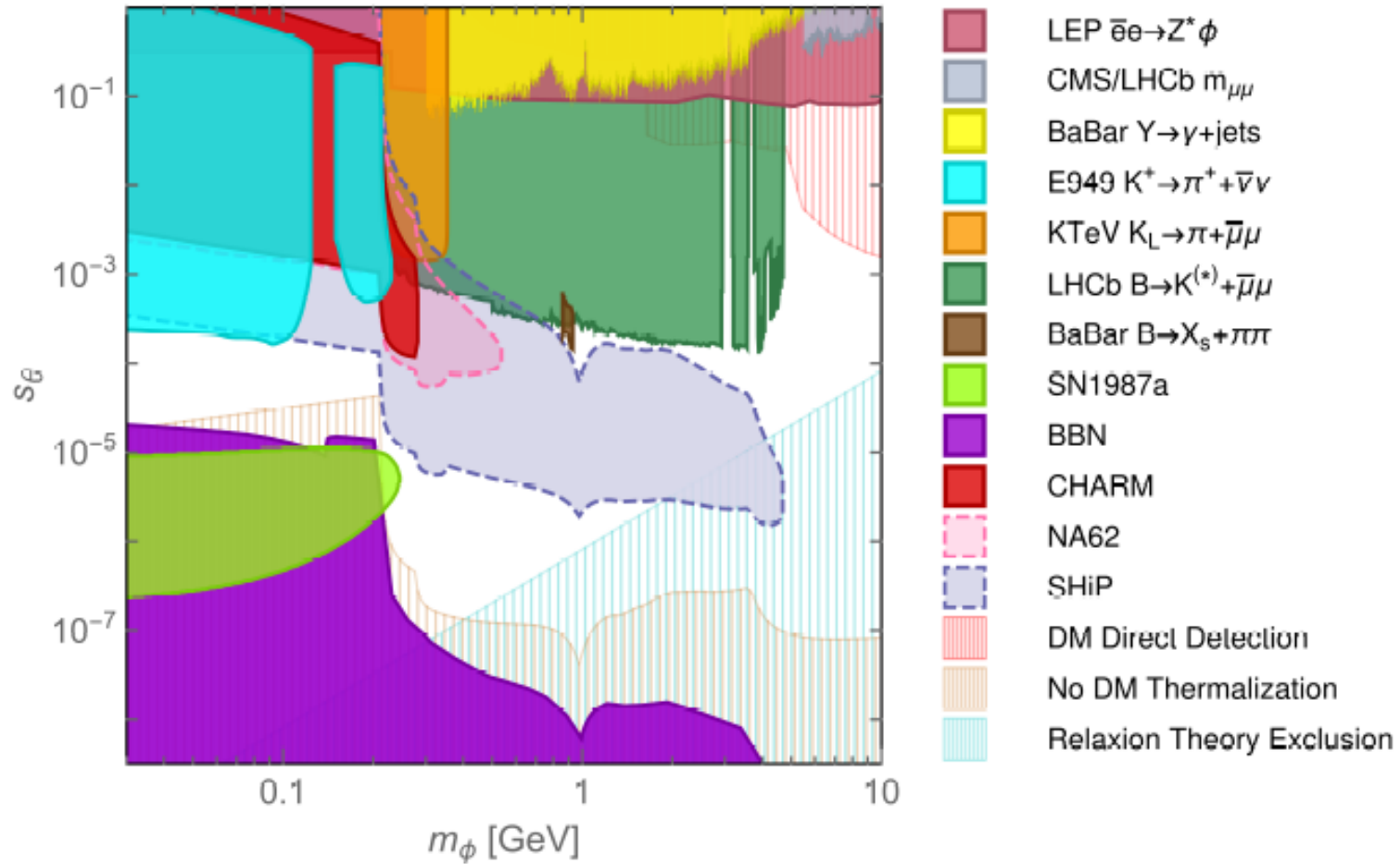
Limits on θ^2



arXiv: 1612.07818

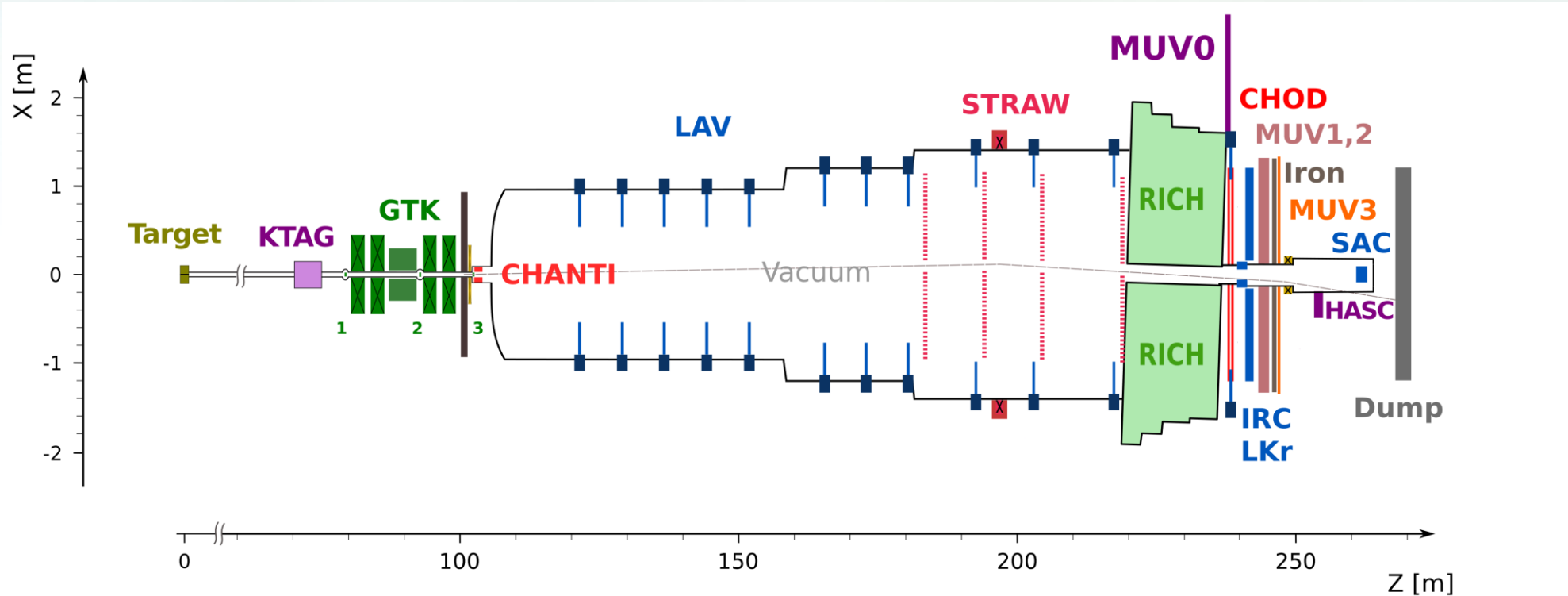
Combined limits

arXiv: 1809.01876



Scalar search in NA62

NA62 setup



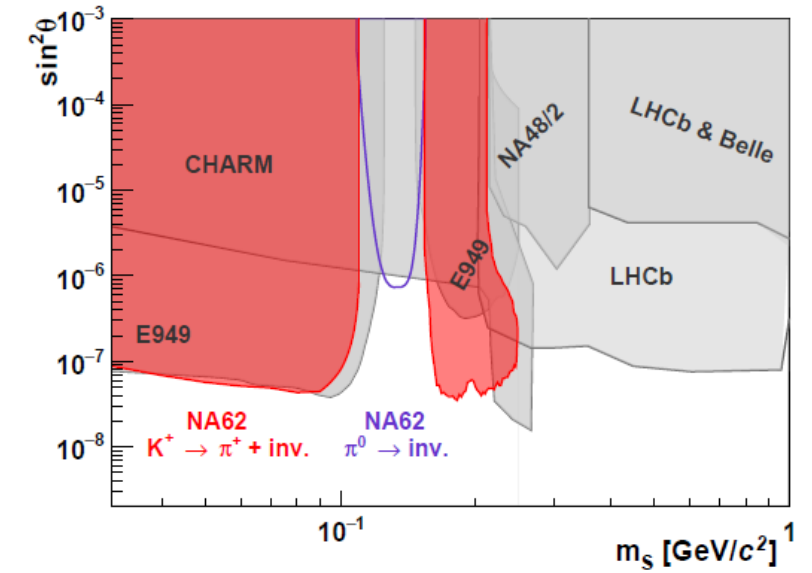
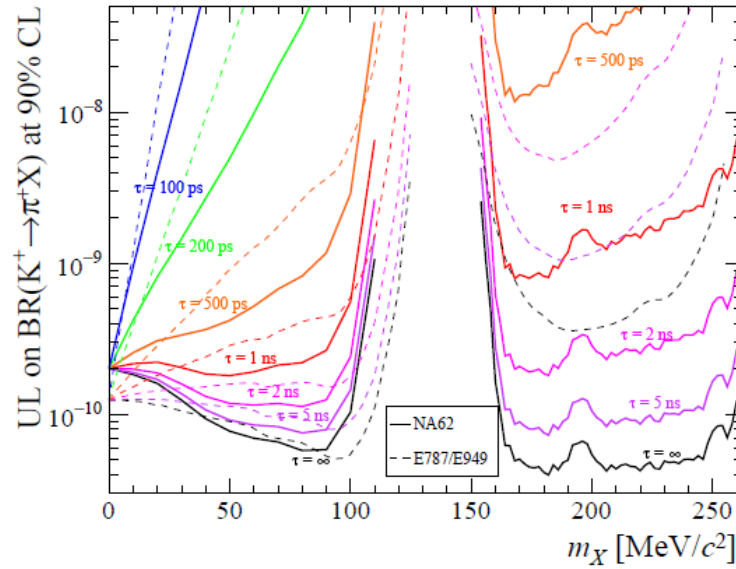
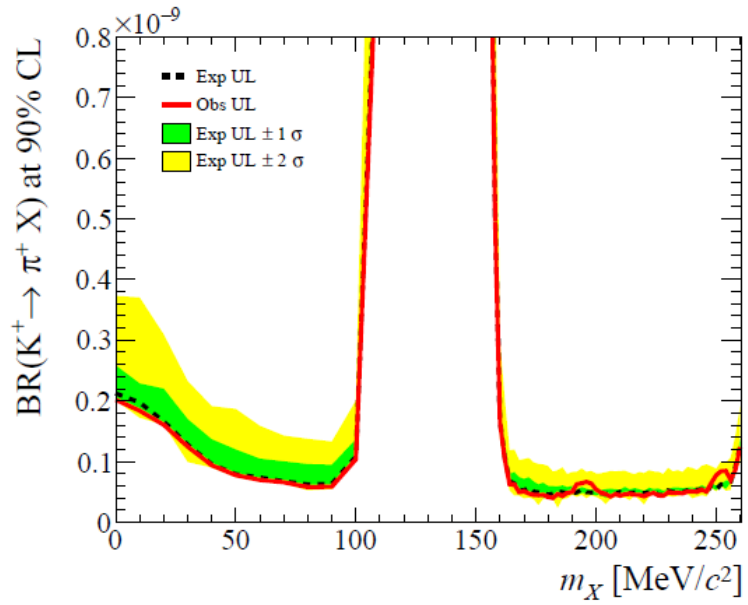
Scalar search in NA62

- Displaced analysis: $K^+ \rightarrow \pi^+ S$, $S \rightarrow \mu^+ \mu^-$ → Ongoing analysis
- Production: $K^+ \rightarrow \pi^+ S$ (byproduct of the PNN analysis) → Published
- Decay: $S \rightarrow \mu^+ \mu^-$ (beam dump mode) → Planned for the next run

Limits from NA62 on $BR(K \rightarrow \pi S)$

Limits on BR

Limits on θ^2



arXiv: 2011.11329

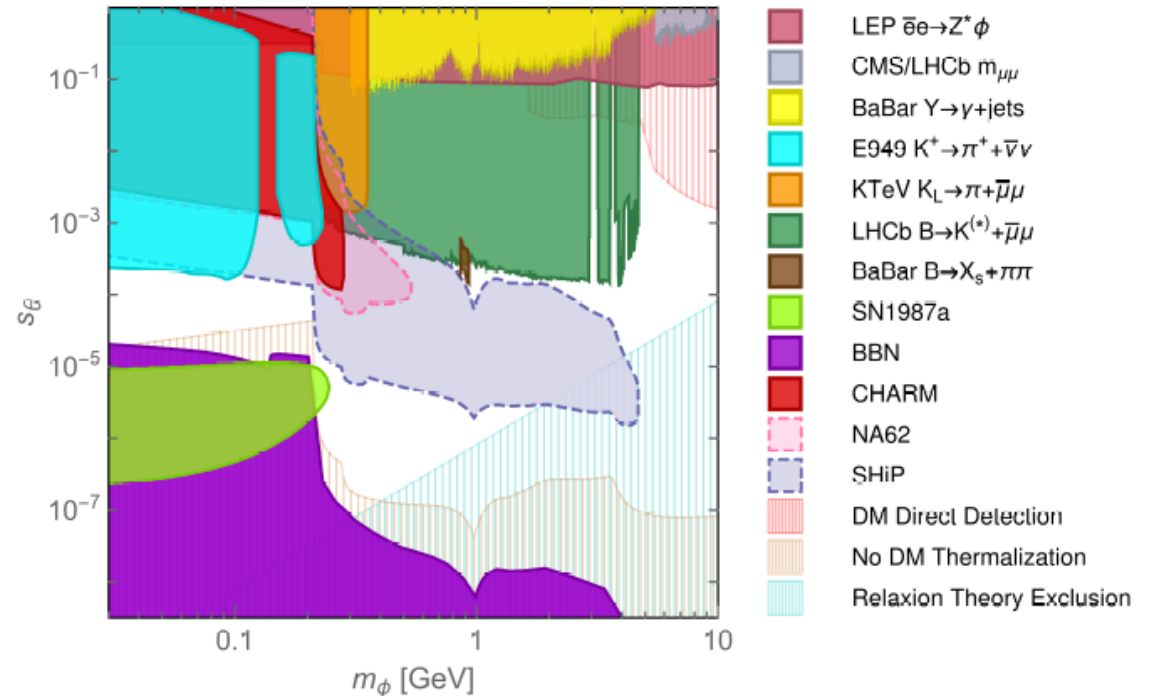
Conclusions

Models with a light scalar

- ❑ possible extension of SM
- ❑ Could explain Dark Matter, hierarchy problem, inflation
- ❑ Inflaton model: 2 parameters m , θ (coupling with Higgs)
- ❑ Best production processes: K and B decays
- ❑ Best decay signature: $S \rightarrow \mu^+ \mu^-$

Experimental searches for the light scalar

- Beam dump (CHARM, SHIP)
- K decays (NA48/2, NA62)
- B decays (LHCb)



Thank you !



Spare