

THERMAL SELECTION
OF PHENOMENOLOGICALLY VIABLE
~~SUSY~~ VACUA

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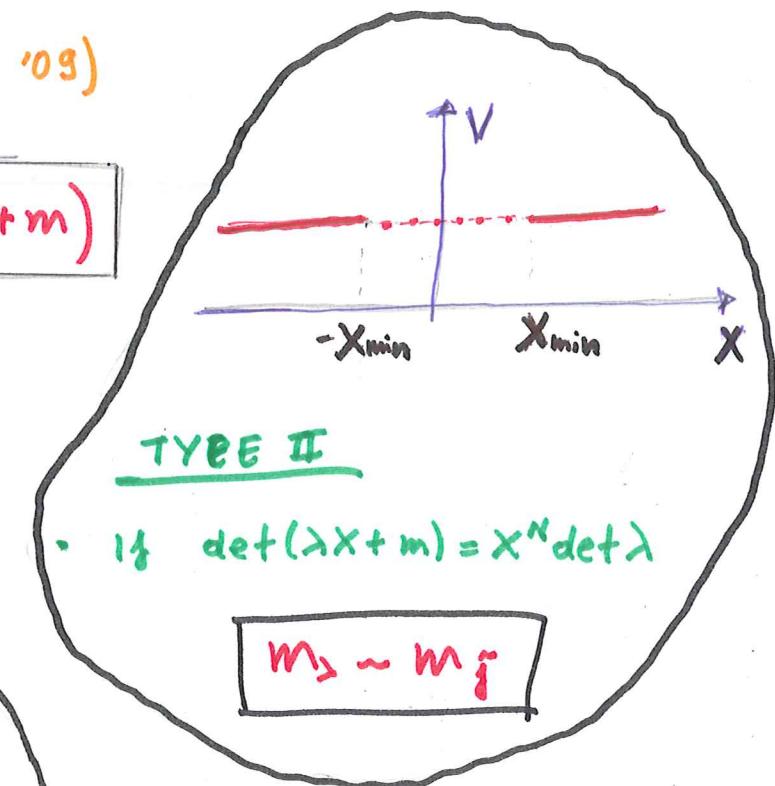
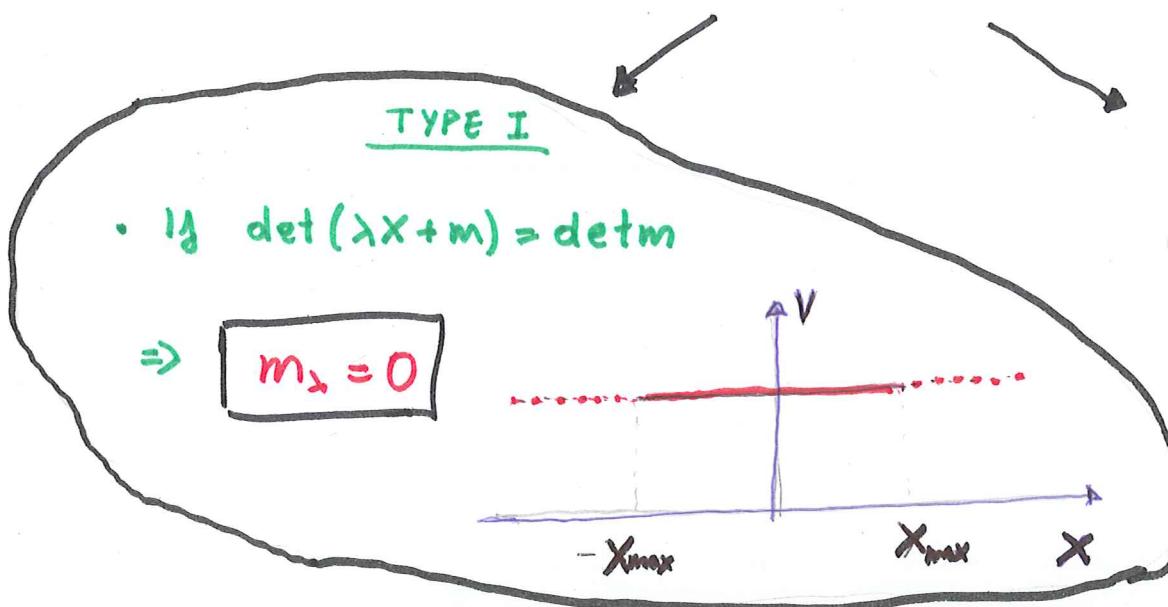
GAUGINO MASSES

- We consider general O'R models

$$W = F X + (\lambda_{ij} X + m_{ij}) q_i \bar{q}_j$$

- Gaugino masses (Komargodski & Shih '09)

$$m_\lambda \sim F^+ \frac{\partial}{\partial X} \log \det(\lambda X + m)$$



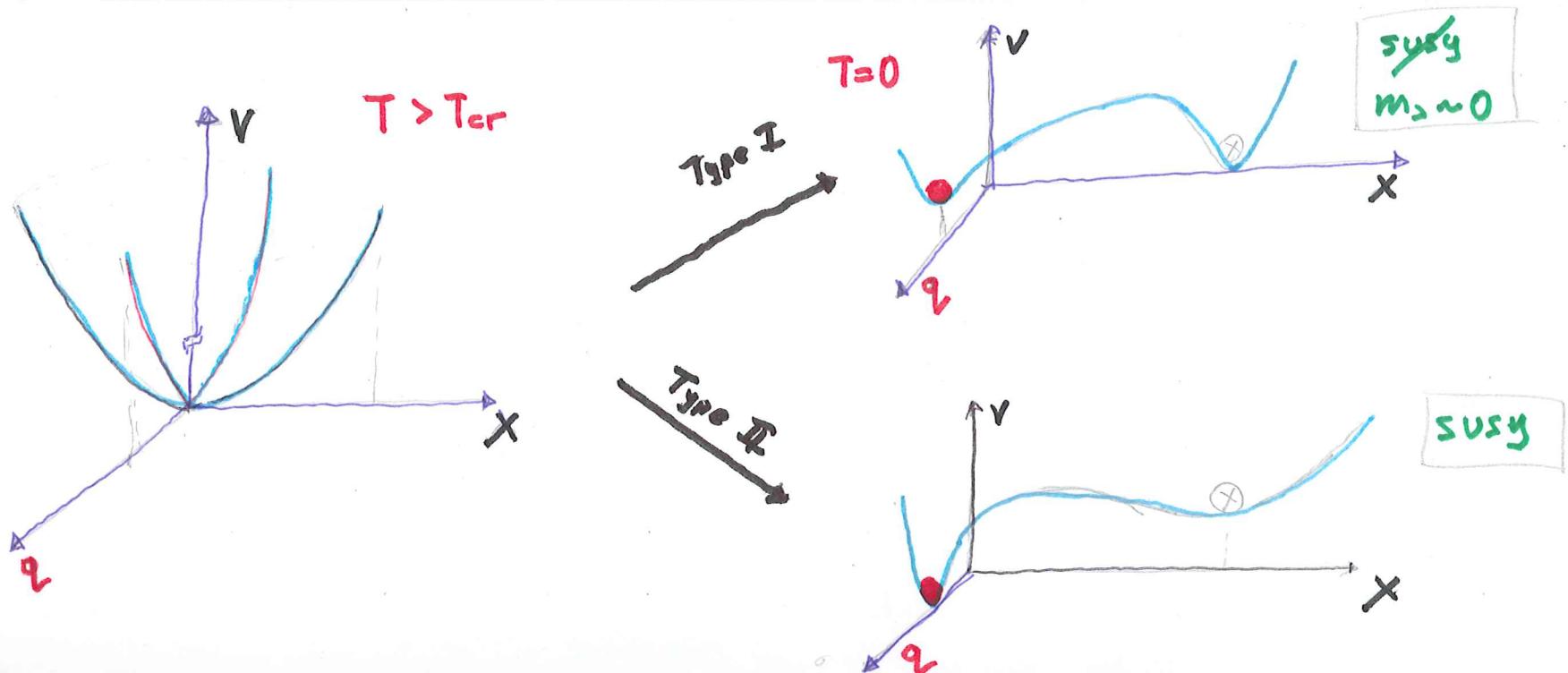
(Cheung et al '07)

THERMAL EVOLUTION (1)

- Type I susy vacua are thermally preferred
Type II are thermally disfavoured
(Abel et al, Craig et al, Fischler et al '06
Katz '09, Dalianis & Lalak '10)

- To escape the problem: non-thermal evolution was considered

THERMAL EVOLUTION

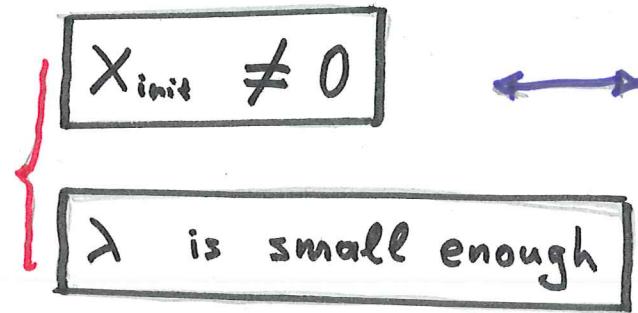


A WAY OUT

(Dalianis & Lalak 10)

- It is possible: SELECTION OF THE PHENO-VIABLE SUSY VACUUM
DURING A THERMAL PHASE WITH HIGH T_{RH}

• If:



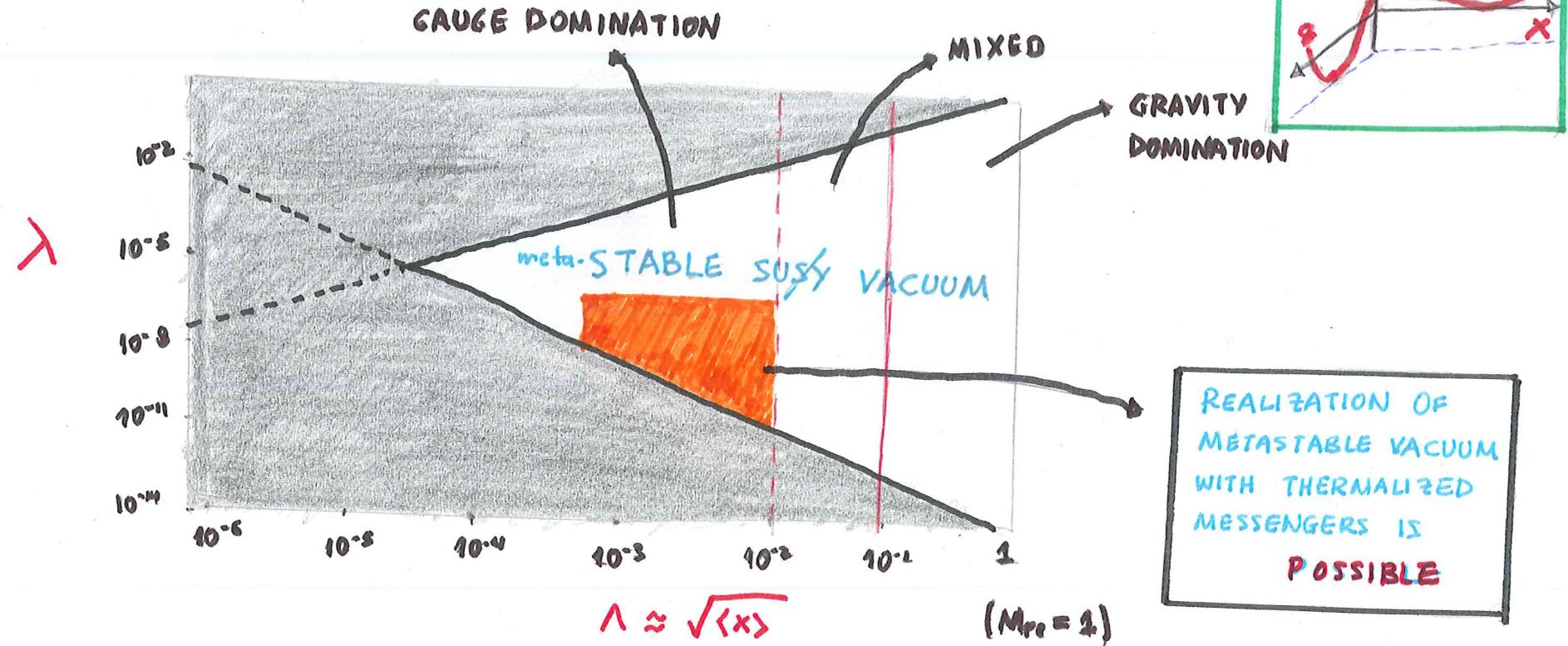
INFLATION { $w = w(I, x)$
 $\delta K = - \frac{g^* g I^* I}{M_{Pl}^2}$

The Metastable Vacuum
"emerges" at high
Temperatures

$T_{RH} > M_{messengers}$
WANTED!
DOES NOT RULE OUT SUSY

X-THERMAL
MASS VERY
SMALL

(Dalianis & Lalak '10)



$(\langle x \rangle, \lambda)$	T_{RH} (GeV)
$(10^{14} \text{ GeV}, 10^{-7})$	$10^{8.5} < T_{RH} < 10^{10}$
$(10^4, 10^{-8})$	$10^8 < T_{RH}$
$(10^{12}, 10^{-7})$	$10^3 < T_{RH} < 10^5$

- COSMOLOGY
- ✓ Leptogenesis
 - ✓ Thermal Dark Matter
- (Olechowski et al '09)

$$\cancel{x}_{\text{initial}} \neq 0$$

→ DEPENDENCE ON
INITIAL CONDITIONS

$$\lambda \ll 1$$

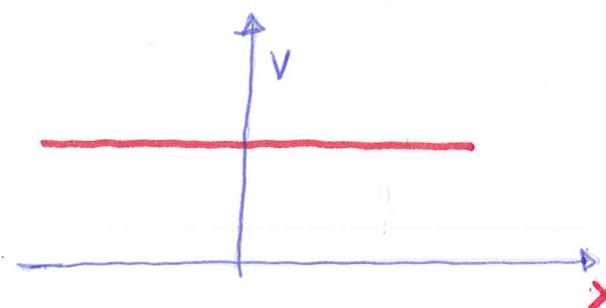
→ ANY OTHER
CONSEQUENCES?

COUPLING λ

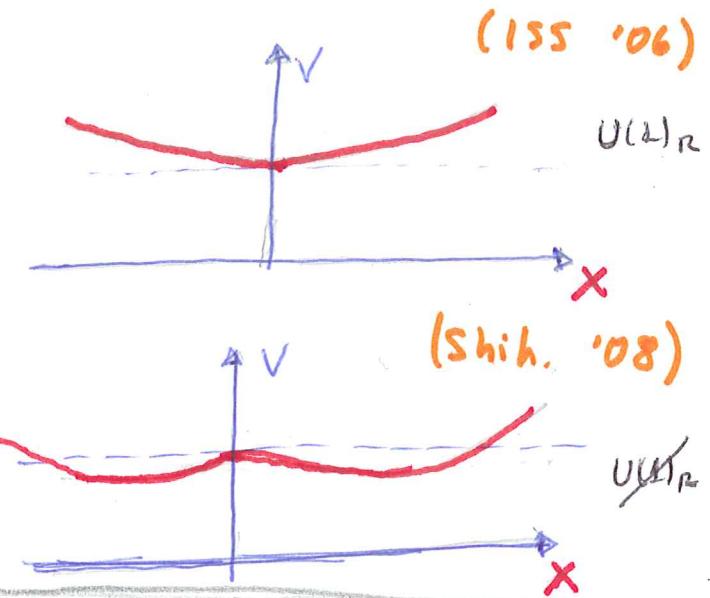
+ QUANTUM CORRECTIONS

$$\delta W = \lambda_{ij} \times q_i \bar{q}_j$$

- 1) It lifts the flat X -direction

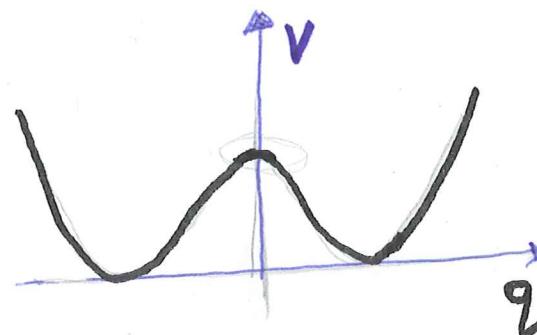


1-loop



- 2) In the q -direction?

$$W = Fx - \lambda \times q\bar{q}$$



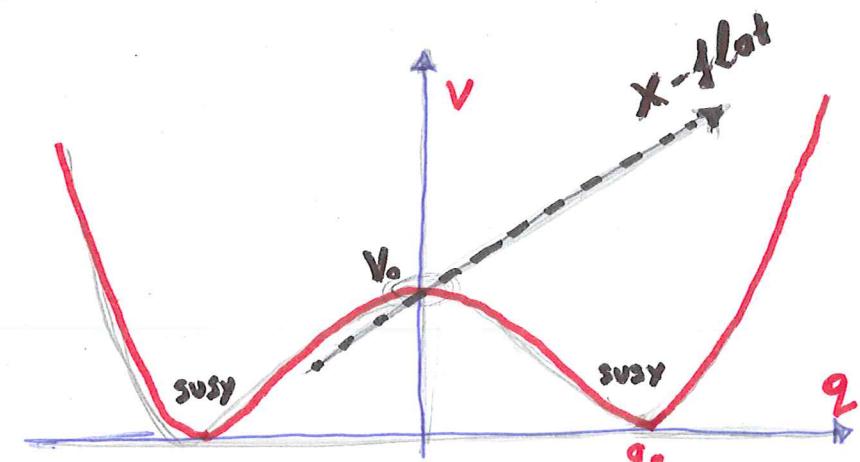
6.

WHAT IF q IS
CHARGED UNDER
AN ABELIAN GAUGE
GROUP $U(L)_a$ WITH
CHARGE e ?

more on QUANTUM CORRECTIONS

- $V^{1\text{-loop}} = \frac{1}{64\pi^2} \left(\text{Tr } m_E^4 \log \frac{m_E^2}{\Lambda^2} - \text{Tr } m_F^4 \log \frac{m_F^2}{\Lambda^2} \right)$

- WE ASSUME q CHARGED
UNDER $U(L)_q$ (GAUGED)

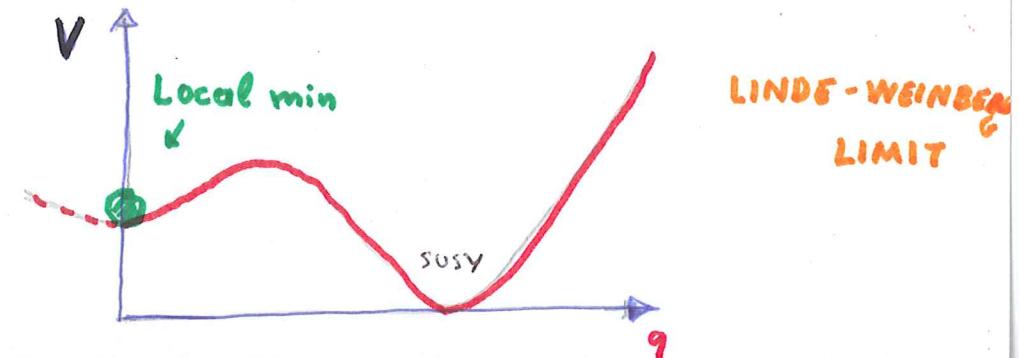


$$\Rightarrow V_{eff}^q = V_{0,R} + m_R^2 |\tilde{q}|^2 + \lambda_R^2 |\tilde{q}|^4 + (\alpha \lambda^4 + \beta e^4) |\tilde{q}|^4 \log \frac{|\tilde{q}|^2}{q_0^2}$$

- $\lambda \ll 1 \Rightarrow \lambda/e \ll 1$

and $m_R^2 > 0$

7.



GAUGED q -DIRECTION

- It can generate MASS HIERARCHIES IN SUSY THEORIES
- HIGGED {
 - $\langle x \rangle$ exponentially large (WITTEN '81)
 - Or NO-HIERARCHICALLY LARGE $\langle x \rangle$ (ISS '07)
 - If GAUGE SYMMETRY NOT HIGGED $\rightarrow \langle x \rangle = 0$
 $\nwarrow q=0$
- R-SYMMETRY can be broken in various ways:

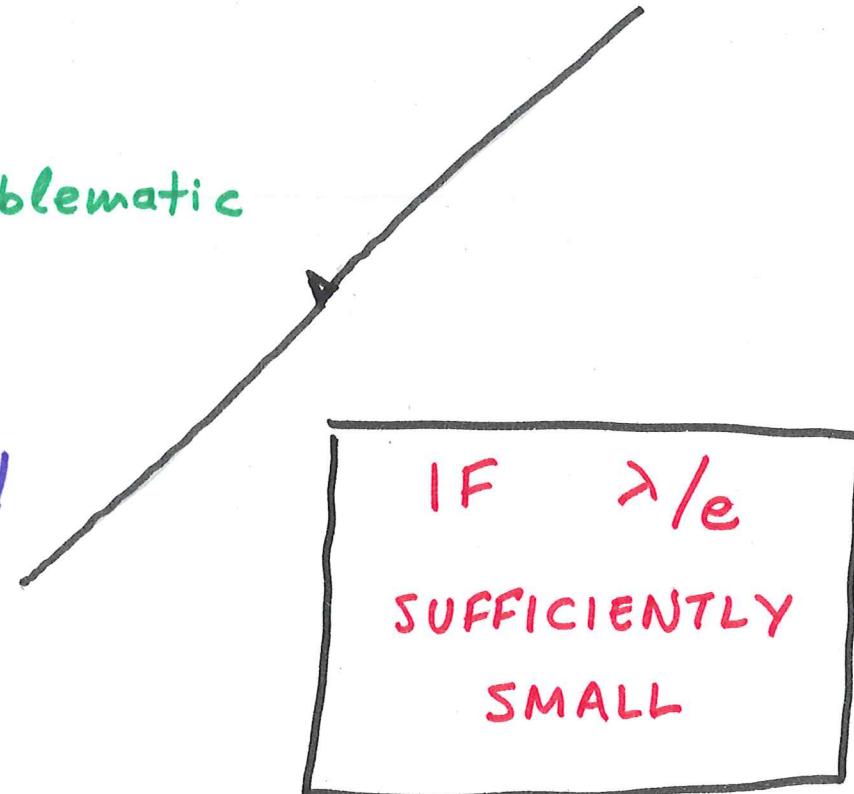
NON-CANONICAL KÄHLER, GRAVITY, EXOTIC R-CHARGES

THERMAL EVOLUTION (2)

- THERE IS NO 2nd ORDER PHASE TRANSITION
TO THE SUSY VACUUM $(\delta q_T < T/\sqrt{2g})$

- $X_{init} = 0$ is NOT problematic

- T_{RH} NOT constrained



CONCLUSIONS

- ★ There is NO THERMAL TRANSITION to SUSY Vacuum via tachyonic directions if $\gamma/e \ll 1$ and in a WINDOW of values
- ★ GAUGINO MASSES ARE NOT VANISHING ($m_{\tilde{g}} \approx m_{\tilde{\tau}}$)
- ★ But, λ CANNOT BE ARBITRARILY SMALL