## The Xth Heavy Quarks and Leptons

## Recent Results from BESIII

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## Outline

- BESIII data

Light Hadron Spectroscopy:
$\rightarrow$ Measurement of the matrix element for the decay
$\eta^{\prime}(958) \rightarrow \eta \pi^{+} \pi^{-}$
$\rightarrow$ Study of $a_{0}(980)-f_{0}(980)$ mixing

- Confirmation of $X(1835)$ and observation of two new structures

Charmonium decay:

* Evidence for $\psi(2 S)$ decays into $\gamma \pi^{0}$ and $\gamma \eta$
* Two-photon transition from $\psi(2 S)$ to $J / \psi$
* Study of $\chi_{c J}$ radiative decays into a vector meson
* Obervation of $\chi_{c J} \rightarrow V V(V=\omega, \phi)$

All results are preliminary!

## World $J / \psi$ and $\psi(2 S)$ Samples $\left(\times 10^{6}\right)$




BESIII: $J / \psi$ 2009: ~226M, and $\psi(2 S)$ 2009: ~106M .

## Measurement of the Matrix Element for the Decay $\eta^{\prime}(958) \rightarrow \eta \pi^{+} \pi^{-}$

## Motivation:

$\rightarrow$ Important for deeper insight into the dynamics of the process and the structure of the particles.

- Important for studies devoted to chiral theory, the effect of the gluon component, and the possible nonet of light scalars.
- Important for the determination of a possible contribution from $f_{0}(600)$ (or $\sigma$ ) resonance (even though the $a_{0}(980)$ is also present).
- Precision measurements on $\eta$ and $\eta^{\prime}$ (958) provide useful information in understanding low energy QCD.


$$
\begin{aligned}
& \mathcal{B}\left(J / \psi \rightarrow \gamma \eta^{\prime}\right)=(4.84 \pm \\
& 0.03(\text { stat }) \pm 0.25(\text { sys })) \times 10^{-3}
\end{aligned}
$$

Measurement of the Matrix Element:

- $X=\frac{\sqrt{3}}{Q}\left(T_{\pi^{+}}-T_{\pi^{-}}\right), \quad Y=\frac{m_{\eta}+2 m_{\pi}}{m_{\pi}} \frac{T_{\eta}}{Q}-1$,
$T_{\pi, \eta}$ denote the kinetic energies of mesons in the $\eta^{\prime}(958)$ rest frame and $Q=T_{\eta}+T_{\pi^{+}}+T_{\pi^{-}}=m_{\eta^{\prime}(958)}-m_{\eta}-2 m_{\pi}$.
- general parametrization: $M^{2}=A\left(1+a Y+b Y^{2}+c X+d X^{2}\right)$
- linear parametrization: $M^{2}=A\left(|1+\alpha Y|^{2}+c X+d X^{2}\right)$ $\alpha$ is a complex parameter. A non-zero value of $\alpha$ may represent the contribution of a gluon component in the wave function of the $\eta^{\prime}(958)$ in the decay.



$$
\begin{aligned}
& \text { 1-d fit: } \operatorname{Re}(\alpha)=\underline{x} 0.034 \pm 0.005, \operatorname{Im}(\alpha)=0.00 \pm \mathrm{d}^{\mathrm{y}} .09, \\
& c=0.019 \pm 0.009, d=-0.058 \pm 0.012 .
\end{aligned}
$$

- $\chi^{2}(N, a, b, c, d)=\sum_{i}^{n_{b i n}} \frac{\left(D_{i}-N M_{i}\right)^{2}}{\sigma_{i}^{2}}$
$M_{i}$ and $D_{i}$ are the numbers of (weighted) entries in the $i$-th bin of the 2-d Dalitz plot for MC and data, respectively.
- general parametrization: $M_{i}=\sum_{j=1}^{N_{\text {ev }}}\left(1+a Y_{j}+b Y_{j}^{2}+c X_{j}+d X_{j}^{2}\right)$ $j$ is an index over the MC events, and $X_{j}$ and $Y_{j}$ are the true generated values of Dalitz variables. Similarly for the linear parametrization


Experimental distributions of the variable $Y$ in various intervals of $X$ with fitting function (histogram) for the general decomposition parametrization.

Table: The left four columns are for $\mathrm{a}, \mathrm{b}, \mathrm{c}$ and d . The right for $\operatorname{Re}(\alpha), \operatorname{Im}(\alpha), \mathrm{c}$ and d .

| VES $^{1}$ | Theory | This work | CLEO | VES $^{2}$ | This work |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $-0.127 \pm 0.018$ | $-0.116 \pm 0.011$ | $-0.047 \pm 0.012$ | $-0.021 \pm 0.025$ | $-0.072 \pm 0.014$ | $-0.033 \pm 0.006$ |
| $-0.106 \pm 0.032$ | $-0.042 \pm 0.034$ | $-0.068 \pm 0.021$ | 0.000 (fixed) | $0.000 \pm 0.100$ | $0.000 \pm 0.050$ |
| $+0.015 \pm 0.018$ | - | $+0.020 \pm 0.012$ | 0.000 (fixed) | $+0.020 \pm 0.019$ | $+0.018 \pm 0.010$ |
| $-0.082 \pm 0.019$ | $+0.010 \pm 0.019$ | $-0.073 \pm 0.013$ | 0.000 (fixed) | $-0.066 \pm 0.034$ | $-0.058 \pm 0.013$ |

VES ${ }^{1}$ : Phys. Lett. B 651, 22 (2007) Theory: Eur. Phys. J A 26, 383 (2005)
CLEO: Phys. Rev. Lett. 84, 26 (2000) VES ${ }^{2}$ :Phys. Atom. Nucl. 68, 372 (2005).
Some comments:
The errors of our fitted parameter values are smaller than previous published results.
In the general parametrization, the values of $a$ and $b$ are consistent with the results from GAMS-4 (PLB177,115), however the values of $c$ and $d$ are consistent with the results from VES ${ }^{1}$.
A negative value of the coefficient $b$ indicates that two kinds of parametrization are not equivalent. This conclusion is consistent with that from GAMS-4 $\pi$. VES ${ }^{1}$ found the fit with linear parametrization yields unsatisfactory $\chi^{2} / N D F=170.5 / 114$ ratio.
(T) The quadratic term in $X$ is unambiguously different from zero. Similarly for the quadratic term in $Y$. The dynamical nature of this term needs clarification.
The value of the parameter $c$ testing C parity violation in strong interaction is consistent with zero within $2 \sigma$ in both parametrizations.

## Study of $a_{0}(980)-f_{0}(980)$ mixing

## Motivation:

- There has been much argument whether the $a_{0}^{0}(980)$ and the $f_{0}(980)$ are part of the ground-state quark-antiquark family or whether they are 4-quark states, hybrids or $K \bar{K}$ molecules.
$\rightarrow$ The mixing between $a_{0}^{0}(980)$ and $f_{0}(980)$ is expected to shed light on the nature of these two resonances.
* Two kinds of mixing intensities $\xi_{a f}$ and $\xi_{f a}$ for the $a_{0}^{0}(980) \rightarrow f_{0}(980)$ and $f_{0}(980) \rightarrow a_{0}^{0}(980)$ transitions are expressed as:

$$
\begin{gathered}
\xi_{f a}=\frac{\operatorname{Br}\left(J / \psi \rightarrow \phi f_{0}(980) \rightarrow \phi a_{0}^{0}(980) \rightarrow \phi \eta \pi^{0}\right)}{\operatorname{Br}\left(J / \psi \rightarrow \phi f_{0} \rightarrow \phi \pi \pi\right)}, \\
\xi_{a f}=\frac{\operatorname{Br}\left(\psi^{\prime} \rightarrow \gamma \chi_{c 1} \rightarrow \gamma \pi^{0} a_{0}^{0}(980) \rightarrow \gamma \pi^{0} f_{0}(980) \rightarrow \gamma \pi^{0} \pi^{+} \pi^{-}\right)}{\operatorname{Br}\left(\psi^{\prime} \rightarrow \gamma \chi_{c 1} \rightarrow \gamma \pi^{0} a_{0}^{0} \rightarrow \gamma \pi^{0} \pi^{0} \eta\right)} .
\end{gathered}
$$

Measurement of $J / \psi \rightarrow \phi f_{0}(980) \rightarrow \phi a_{0}^{0}(980) \rightarrow \phi \eta \pi^{0}$ and $\psi^{\prime} \rightarrow \gamma \chi_{c 1} \rightarrow \gamma \pi^{0} a_{0}^{0}(980) \rightarrow \gamma \pi^{0} f_{0}(980) \rightarrow \gamma \pi^{0} \pi^{+} \pi^{-}$


(a). $\phi$ sianal reaion. (b). $\phi$ sidebands reaion. The fitted results:

(a). $\chi_{c 1}$ signal region. (b). $\chi_{c 1}$ sidebands region.

$$
\begin{aligned}
& \mathcal{B}\left(J / \psi \rightarrow \phi f_{0}(980) \rightarrow \phi a_{0}^{0}(980) \rightarrow \phi \eta \pi^{0}\right)=(3.2 \pm 1.1 \pm 0.8) \times 10^{-6} \\
& \left(<5.1 \times 10^{-6}\right) @ 90 \% \mathrm{C} . \mathrm{L} . \\
& \mathcal{B}\left(\psi^{\prime} \rightarrow \gamma \chi_{c 1} \rightarrow \gamma \pi^{0} a_{0}^{0}(980) \rightarrow \gamma \pi^{0} f_{0}(980) \rightarrow \gamma \pi^{0} \pi^{+} \pi^{-}\right)= \\
& (2.7 \pm 1.4 \pm 0.7) \times 10^{-7}\left(<5.9 \times 10^{-7} @ 90 \% \text { C.L. }\right) .
\end{aligned}
$$

## Discussion:

The mixing intensity $\xi_{f a}$ for the $f_{0}(980) \rightarrow a_{0}^{0}(980)$ transition is calculated to be:

$$
\begin{aligned}
& \xi_{f a}=\frac{\operatorname{Br}\left(J / \psi \rightarrow \phi f_{0}(980) \rightarrow \phi a_{0}^{0}(980) \rightarrow \phi \eta \pi^{0}\right)}{\operatorname{Br}\left(J / \psi \rightarrow \phi f_{0} \rightarrow \phi \pi \pi\right)} \\
&=0.6 \pm 0.2(\text { stat. }) \pm 0.2(\text { sys. }) \%
\end{aligned}
$$

The mixing intensity $\xi_{a f}$ for the $a_{0}^{0}(980) \rightarrow f_{0}(980)$ transition is calculated to be:

$$
\begin{aligned}
&\left.\xi_{a f}=\frac{\operatorname{Br}\left(\psi^{\prime} \rightarrow \gamma \chi_{c 1} \rightarrow \gamma \pi^{0} a_{0}^{0}(980) \rightarrow \gamma \pi^{0} f_{0}(980) \rightarrow \gamma \pi^{0} \pi^{+} \pi^{-}\right)}{\operatorname{Br}\left(\psi^{\prime} \rightarrow \gamma \chi_{c 1}\right.} \rightarrow \gamma \gamma \pi^{0} a_{0}^{0} \rightarrow \gamma \pi^{0} \pi^{0} \eta\right) \\
&=0.3 \pm 0.2(\text { stat. }) \pm 0.1(\text { sys. }) \%
\end{aligned}
$$



The mixing intensities and predictions with various theoretical predictions. The shaded region is our measurement with error bars and the red lines are our limits.

## Confirmation of X(1835) and observation of two new structures in $J / \psi \rightarrow \gamma \eta^{\prime}(958) \pi^{+} \pi^{-}$

Motivation:

* Confirmation of $X(1835)$ is necessary with high statistic data sample.

LQCD predicts the $0^{-+}$glueball mass is $2.3 \mathrm{GeV} / \mathrm{c}^{2}$.

- A $0^{-+}$glueball may have similar property as $\eta_{c}$ (the main $\eta_{c}$ decay mode is $\left.\eta^{\prime}(958) \pi^{+} \pi^{-}\right)$.


BESII results:
signal significance is $7.7 \sigma$
$M=1833.7 \pm 6.1$ (stat) $\pm 2.7$ (sys)
$\mathrm{MeV} / c^{2}$
$\Gamma=67.7 \pm 20.3($ stat $) \pm 7.7$ (sys)
$\mathrm{MeV} / \mathrm{c}^{2}$
Phys. Rev. Lett. 95, 262001
(2005)

## Mass spectrum of $\pi^{+} \pi^{-} \eta^{\prime}$


$>X(1835)$ and $\eta_{\mathrm{c}}$ is observed.
> Two additional structures at $\mathrm{M} \sim 2.1 \mathrm{GeV}$ and 2.3 GeV
> There maybe some $\mathrm{f}_{1}(1510)$.


## Fitting the mass spectrum:

$>$ Three background components:
(1) Contribution from non- $\eta$ ' events estimated by $\eta$ ' mass sideband
(2) Contribution from $J / \psi \rightarrow \pi^{0} \pi^{+} \pi^{-} \eta^{\prime}\left(\eta^{\prime} \rightarrow \gamma \rho\right)$ with re-weighting method
(3) Contribution from "PS background"

$$
f_{\text {bkg }}(x)=\left(x-m_{0}\right)^{1 / 2}+a_{0}\left(x-m_{0}\right)^{3 / 2}+a_{1}\left(x-m_{0}\right)^{5 / 2}, m_{0}=2 m_{\pi}+m_{\eta^{\prime}}
$$

 Red line: estimated contribution of (1)+ (2) Black line: total background

| resonance | $M\left(\mathrm{MeV} / c^{2}\right)$ | $\Gamma\left(\mathrm{MeV} / c^{2}\right)$ | Stat. sig. |
| :---: | :---: | :---: | :---: |
| $\mathrm{X}(1835)$ | $1838.1 \pm 2.8$ | $179.5 \pm 9.1$ | $>25 \sigma$ |
| $\mathrm{X}(2120)$ | $2124.8 \pm 5.6$ | $101 \pm 14$ | $>7.2 \sigma$ |
| $\mathrm{X}(2370)$ | $2371.0 \pm 6.4$ | $108 \pm 15$ | $>6.7 \sigma$ |

Stat. sig. is conservatively estimated:
fit range, background shape, contribution of extra resonances
$\Leftrightarrow X(1835)$ resonance is confirmed at BESIII, but the width is significantly larger than that measured at BESII with one resonance in the fit.
$\Leftrightarrow$ Two new resonances, $X(2120)$ and $X(2370)$, are observed.
$\Leftrightarrow$ PWA is needed

## Evidence for $\psi(2 S)$ decays into $\gamma \pi^{0}$ and $\gamma \eta$

## Motivation:

$\rightarrow$ Important tests for various phenomenological mechanisms, such as vector meson dominance model, two-gluon couplings to $q \bar{q}$ states, mixing of $\eta_{c}-\eta^{(\prime)}$, and final-state radiation by light quarks.
$\rightarrow$ The ratio of $R_{J / \psi}=\mathcal{B}(J / \psi \rightarrow \gamma \eta) / \mathcal{B}\left(J / \psi \rightarrow \gamma \eta^{\prime}\right)$ can be predicted by the first order of perturbation theory, and $R_{J / \psi}=R_{\psi(2 S)}$ is expected (CLEO: $R_{\psi(2 S)}<1.8 \%$ at $90 \%$ C.L. and $R_{J / \psi}=(21.1 \pm 09) \%$. PRD79,111101 (2009))
$\rightarrow$ The decay $\psi(2 S) \rightarrow \gamma \pi^{0}$ is suppressed because the photon can only be from final state radiation off one of the quarks.
) $\mathcal{B}\left(\psi(2 S) \rightarrow \gamma \pi^{0}\right)=2.19 \times 10^{-7}$ : calculated in PRD79,097301. CLEO: $<5.0 \times 10^{-6}$ at $90 \%$ C.L. (PRD79,111101).
${ }^{+} e^{-} \rightarrow \psi(2 S) / \gamma^{\star} \rightarrow \gamma \pi^{0}$ will be very useful in testing the form factor for timelike photons $Q^{2}=-q^{2}<0$ (PRD79,097301).

## Results:



## Two-photon transition from $\psi(2 S)$ to $J / \psi$

## Motivation:

On experimental side:

- not seen previously in $\psi(2 S)$ decays
* analogous process to positronium and hydrogen two-photon transition
© CLEO reported $\Upsilon(3 S) \rightarrow \gamma \gamma \Upsilon(2 S)$ (Phys. Rev. D 49, 40 (1994))
On theoretical side:
* order $\alpha^{2}$ QED transition between two hadrons
- Similar process has been studied in heavy-light quark system
© improve understanding of heavy quarkonium such as spectrum, decay et al, and the strong interaction
* possibility of testing the hadron-loop effect


## Signal Estimation

## unbinned maximum likelihood fit with composition of three PDFs:

- signal (red): shape from phase-space-like MC simulation
- $\boldsymbol{\psi ( 2 S )}$ bkg.(blue): shape and magnitude from exclusive MC simulation
- other bkg.(green): $1^{\text {st }}$-order polynominal



Combined with $e^{+} e^{-}$and $\mu^{+} \mu^{-}$modes, the branching fraction is measured to be $\mathcal{B}(\psi(2 S) \rightarrow \gamma \gamma J / \psi)=\left(1.02 \pm 0.05(\text { stat })_{-0.20}^{+0.18}(\right.$ sys $\left.)\right) \times 10^{-3}$.

## Study of $\chi_{a /}$ radiative decays into a vector meson

## Motivation:

$\rightarrow \psi \rightarrow \gamma X \rightarrow \gamma \gamma V\left(\rho^{0}, \omega, \phi\right)$ provide a favorable place to extract information on the flavor content of the $C$-even resonance $X$ to study gluon hadronization dynamics.
$\rightarrow$ By including hadronic loop contributions, a recent pQCD calculation (arXiv:1005.0066) obtains results in agreement with the experimental measurements of $\mathcal{B}\left(\chi_{c 1} \rightarrow \gamma V\right)$.
Table: Theoretical predictions(in units of $10^{-6}$ ) and results from the CLEO.

| Mode | CLEO $^{1}$ | pQCD $^{2}$ | QCD $^{3}$ | QCD+QED $^{3}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\chi_{c 0} \rightarrow \gamma \rho^{0}$ | $<9.6$ | 1.2 | 3.2 | 2.0 |
| $\chi_{c 1} \rightarrow \gamma \rho^{0}$ | $243 \pm 19 \pm 22$ | 14 | 41 | 42 |
| $\chi_{c 2} \rightarrow \gamma \rho^{0}$ | $<50$ | 4.4 | 13 | 38 |
| $\chi_{c 0} \rightarrow \gamma \omega$ | $<8.8$ | 0.13 | 0.35 | 0.22 |
| $\chi_{c 1} \rightarrow \gamma \omega$ | $83 \pm 15 \pm 12$ | 1.6 | 4.6 | 4.7 |
| $\chi_{c 2} \rightarrow \gamma \omega$ | $<7.0$ | 0.5 | 1.5 | 4.2 |
| $\chi_{c 0} \rightarrow \gamma \phi$ | $<6.4$ | 0.46 | 1.3 | 0.03 |
| $\chi_{c 1} \rightarrow \gamma \phi$ | $<26$ | 3.6 | 11 | 11 |
| $\chi_{c 2} \rightarrow \gamma \phi$ | $<13$ | 1.1 | 3.3 | 6.5 |

1. PRL 101,151801 (2008). 2. Chin. Phys. Lett. 23, 2376 (2006). 3. hep-ph/0701009



Invariant mass distributions of (a) $\gamma \phi$, (b) $\gamma \rho^{0}$, and (c) $\gamma \omega$. Dots with error bars are data; histograms are the best fit; dashed lines are signal shapes; and the shaded histograms are vector meson sideband background plus a 2nd order polynomial background.

Table: Results of $\chi_{c J} \rightarrow \gamma V$. The upper limits are at the $90 \%$ C.L.

| Decay <br> mode | Number of <br> events | Efficiency <br> $(\%)$ | Systematic <br> error $(\%)$ | Branching <br> fraction $\left(\times 10^{-6}\right)$ | Statistical <br> significance |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\chi_{c 0} \rightarrow \gamma \phi$ | $15.0 \pm 6.6$ | 32.4 | 8.1 | $9.5 \pm 4.2 \pm 0.8$ | $2.9 \sigma$ |
| $\chi_{c 1} \rightarrow \gamma \phi$ | $42.6 \pm 8.6$ | 34.6 | 7.8 | $25.8 \pm 5.2 \pm 2.0$ | $6.4 \sigma$ |
| $\chi_{c 2} \rightarrow \gamma \phi$ | $4.6 \pm 4.9$ | 32.6 | 8.8 | $<8.0$ |  |
| $\chi_{c 0} \rightarrow \gamma \rho^{0}$ | $6 \pm 12$ | 22.6 | 7.4 | $<10.2$ |  |
| $\chi_{c 1} \rightarrow \gamma \rho^{0}$ | $432 \pm 25$ | 19.4 | 7.2 | $228 \pm 13 \pm 16$ | $\gg 10 \sigma$ |
| $\chi_{c 2} \rightarrow \gamma \rho^{0}$ | $13 \pm 11$ | 15.7 | 7.9 | $<20.4$ |  |
| $\chi_{c 0} \rightarrow \gamma \omega$ | $5 \pm 11$ | 18.6 | 8.3 | $<12.7$ |  |
| $\chi_{c 1} \rightarrow \gamma \omega$ | $136 \pm 14$ | 22.7 | 8.0 | $69.7 \pm 7.2 \pm 5.6$ | $\gg 10 \sigma$ |
| $\chi_{c 2} \rightarrow \gamma \omega$ | $1 \pm 6$ | 19.2 | 8.9 | $<6.0$ |  |

The longitudinal (transverse) polarization exhibits a $\cos ^{2} \Theta\left(\sin ^{2} \Theta\right)$ dependence, and the angular distribution is expressed as:

$$
\frac{d N}{d \cos \theta} \propto\left|A_{L}\right|^{2} \cos ^{2} \theta+\frac{1}{2}\left|A_{T}\right|^{2} \sin ^{2} \Theta
$$

where $A_{L}$ and $A_{T}$ are the longitudinal and transverse polarization amplitudes, and $\Theta$ is defined as the angle between the vector meson flight direction in the $\chi_{\sigma \omega}$ rest frame and either the $\pi^{+} / K^{+}$direction in the $\rho^{0} / \phi$ rest frame or the normal to the $\omega$ decay plane in the $\omega$ rest frame.




Results: The transverse component fraction:
$f_{T}=\left|A_{T}\right|^{2} /\left(\left|A_{T}\right|^{2}+\left|A_{L}\right|^{2}\right)=N_{T} /\left(N_{T}+R * N_{L}\right)$, where $R=\varepsilon_{T} / \varepsilon_{L}$
$f_{T}$ are $0.29_{-0.12-0.09}^{+0.13+0.10}$ for $\chi_{c 1} \rightarrow \gamma \phi, 0.158 \pm 0.034_{-0.014}^{+0.015}$ for $\chi_{c 1} \rightarrow \gamma \rho^{0}$, and $0.247_{-0.087-0.026}^{+0.090+0.044}$ for $\chi_{c 1} \rightarrow \gamma \omega$.

## Observation of $\chi_{a J} \rightarrow \omega \omega, \phi \phi$ and $\omega \phi$

## Motivation:

Important laboratory to test QCD:

- Previous measurements from BESII.

Only $\chi_{\mathrm{c} 0}$ and $\chi_{\mathrm{c} 2}$ decays into $\phi \phi$ and $\omega \omega$ are observed.

| $\operatorname{BR}\left(10^{-3}\right)$ | $\chi_{\mathrm{c} 0}$ | $\chi_{\mathrm{c} 2}$ |
| :--- | :---: | :---: |
|  |  |  |
|  | $0.94 \pm 0.21 \pm 0.13$ | $1.70 \pm 0.30 \pm 0.25$ |
| BESII, PLB 642, 197 (2006) |  |  |
| $\rightarrow \omega \omega$ | $2.29 \pm 0.58 \pm 0.41$ | $1.77 \pm 0.47 \pm 0.36$ | BESII, PLB 630, 7 (2005)

$\chi_{\mathrm{c} 1} \rightarrow \mathrm{VV}$ is suppressed due to helicity selection rule in pQCD having different polarization, so it is suppressed.

$$
\text { So } \lambda_{1}+\lambda_{2} \neq 0
$$

$\square \chi_{c\lrcorner}->\omega \phi$ is doubly OZI suppressed.


First observation of $\omega \phi$ which is a doubly OZI suppressed decay, long distance contribution may be important in charmonium decays.

## Other recent results from BESIII that have not been included here.

§ Analysis of $J / \psi \rightarrow \omega \eta \pi^{+} \pi^{-}$: A structure denoted as $\mathrm{X}(1870)$ is seen in $\eta \pi^{+} \pi^{-}$mass spectrum. For details, see Yanping HUANG's report at ICHEP10:
http://indico.cern.ch/contributionDisplay.py?contribld=1210\&sessionld=46\&confld=73513

B Analysis of $\chi_{C J} \rightarrow 4 \pi^{0}$ : it is the first measurement. For details, see Ronggang PING's report at ICHEP10.
http://indico.cern.ch/contributionDisplay.py?contribld=1233\&sessionld=50\&confld=73513

## Summary

Some preliminary results from BESIII have been shown here.
$\leftrightarrow_{0}$ The Dalitz plot of $\eta^{\prime}(958) \rightarrow \eta \pi^{+} \pi^{-}$decay is studied in a generalized and a linear representation.
\& We perform direct measurements of $a_{0}^{0}(980)-f_{0}(980)$ mixing.
\& $X(1835) \rightarrow \eta^{\prime}(958) \pi^{+} \pi^{-}$is confirmed and two new resonances, $X(2120)$ and $X(2370)$, are observed.
$\psi(2 S) \rightarrow \gamma \pi^{0}$ and $\psi(2 S) \rightarrow \gamma \eta$ are observed for the first time with signal significance of $4.1 \sigma$ and $3.2 \sigma$, respectively.
A significant enhancement of two-photon transition of $\psi(2 S)$ to $\mathrm{J} / \psi$ was observed for the first time
$\Delta_{0}$ The decays $\chi_{c J} \rightarrow \gamma V\left(V=\phi, \rho^{0}, \omega\right)$ are studied. The fractions of the longitudinal polarization component of V in $\chi_{c 1} \rightarrow \gamma \mathrm{~V}$ are measured.
$\chi_{c \omega}$ signals are observed in the decays $\chi_{\omega J} \rightarrow \omega \omega, \phi \phi$ and $\omega \phi$.

## Thanks! 谢谢!

## Backup



The corresponding projections on variables $X$ and $Y$ in (b) and (c), respectively, where the dashed histograms are from MC signal sample with $\eta^{\prime}(958) \rightarrow \eta \pi^{+} \pi^{-}$events produced with phase space and the blank histograms are the fitted results described in the text.

