


Search for neutrinos in coincidence with GWTC-2 events in the Super-Kamiokande detector

NeuTel conference

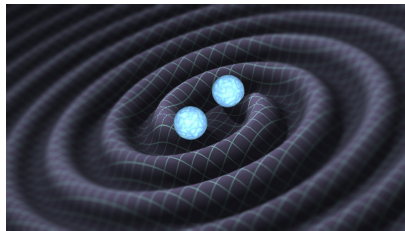
Mathieu Lamoureux 

(INFN Sezione di Padova, Italy)

for the Super-Kamiokande collaboration

Since 2015, the LIGO/Virgo Collaboration (LVC) is detecting and sending alerts for gravitational waves from the merger of binary objects.

- Binary Neutrino Star (BNS): may produce short Gamma-Ray Bursts (GRB) with neutrino production*
- Binary Black Hole (BBH): neutrino production in the accretion disks of the black holes†
- Neutron Star - Black Hole (NSBH)



Detecting coincident neutrinos from these objects would allow better understanding of the mechanisms behind them.

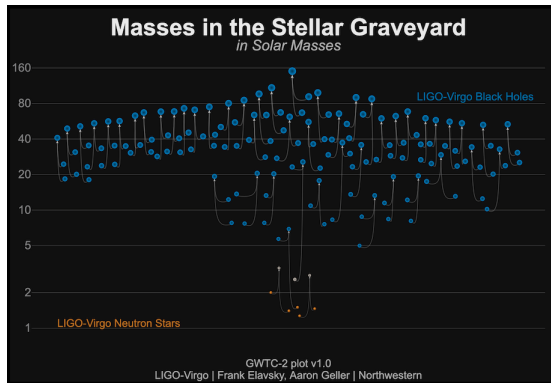
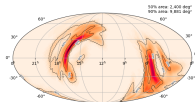
*Foucart, F., et al (2016). *Low mass binary neutron star mergers: Gravitational waves and neutrino emission*. Physical Review D, 93(4). [10.1103/PhysRevD.93.044019](https://arxiv.org/abs/10.1103/PhysRevD.93.044019)

†Caballero, O. L., et al (2016). *Black hole spin influence on accretion disk neutrino detection*. [10.1103/PhysRevD.93.123015](https://arxiv.org/abs/10.1103/PhysRevD.93.123015)

- LIGO-Virgo Third Observing Run (O3) covered April 2019 to March 2020
⇒ 56 alerts provided in realtime through GCN
- GWTC-2 covers the first half of O3 (April 2019 - September 2019)
⇒ 39 confirmed detections ⇐ **focus of this talk**

For each GW, we have:

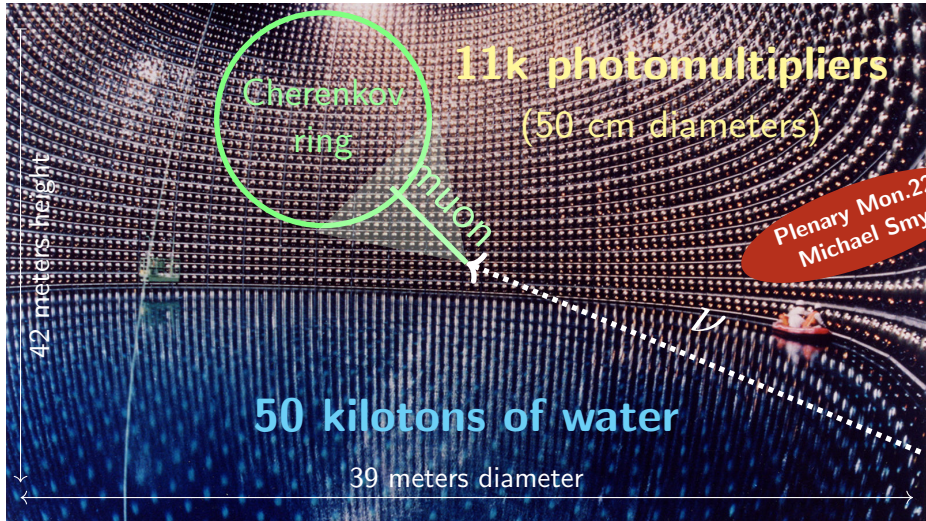
- time of the event
- sky localisation
- estimated distance
- estimated masses of the two objects
- can be roughly classified based on masses ($m < 3 M_{\odot}$ =NS, $m > 3 M_{\odot}$ =BH)



The Super-Kamiokande experiment

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Experiment running since 1998, located in the Mozumi mine in Japan.



Partially-Contained

 $E_\nu = 3.5 - 100 \text{ MeV}$

Fully-Contained

 $E_\nu = 0.1 - 10 \text{ GeV}$

Four samples covering the neutrino energy range from few MeV to $\mathcal{O}(\text{TeV})$:

- low-energy (LOWE)
- fully-contained events (FC)
- partially-contained events (PC)
- upgoing muons (UPMU)

LOWE is usually used for solar/supernova analyses.

Flash Fri.19
Alberto, DSνBFlash Wed.24
Alice, Spallation

The other samples are mainly used for atmospheric analysis.

Parallel Wed.24
Pablo, Atmospheric

$$E_\nu = 0.1 - 100 \text{ GeV}$$

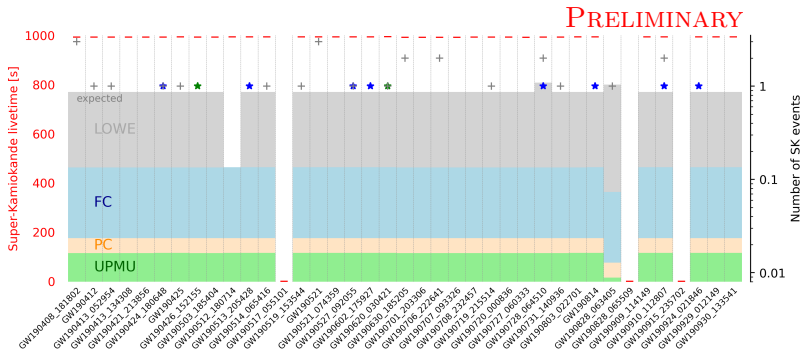
Upgoing muons

 $E_\nu > 1.6 \text{ GeV}$

- Define a ± 500 s centered on GW time
 - Search for events within this time window, in the four SK samples
 - Compare observation with expected background and extract neutrino flux upper limits
- and compute eventual signal significance by comparing neutrino directions and GW localisation (only for high-energy SK samples)

Low-energy sample	High-energy samples		
	FC	PC	UPMU
Standard solar/SRN selection + 7 MeV energy threshold to ensure stable bkg rate	Standard atmospheric selection		
expected background in 1000 seconds = 0.729	0.112	0.007	0.016

Performed the analysis for the 39 GW in GWTC-2. Three of them were associated to SK downtime (due to calibration) (one less for low-energy due to HV issues).



In total:

Sample	N_{obs}	N_{exp}
LOWE+	24	24.97
FC★	8	3.95
PC★	0	0.26
UPMU★	2	0.58

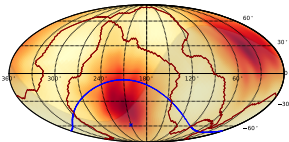
PRELIMINARY

No significant excess was observed in the follow-up analysis.

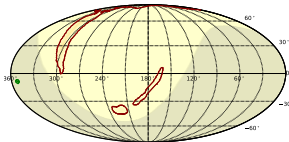
Ten SK high-energy events in time coincidence

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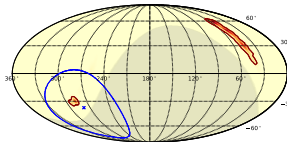
GW190424_180648



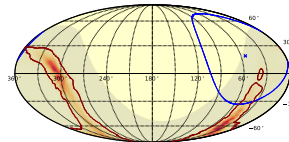
GW190426_152155



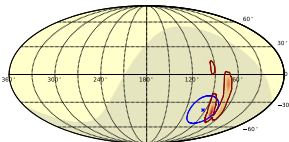
GW190513_205428



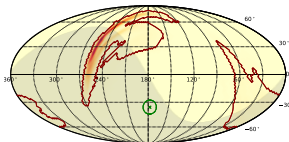
GW190527_092055



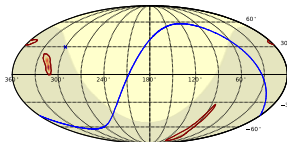
GW190602_175927



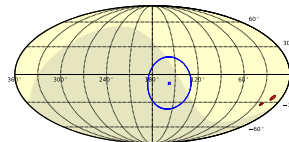
GW190620_030421



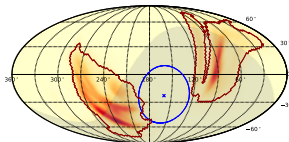
GW190728_064510



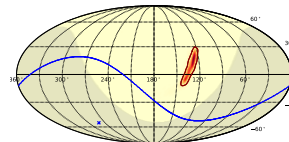
GW190814



GW190910_112807



GW190924_021846



All plots are **PRELIMINARY**

Skymaps in equatorial coordinates

Red: GW localisation and 90% contour

Blue: SK FC events with 1σ angular uncertainty

Green: SK UPMU events.

How likely the SK observation is associated to background, given time+space correlations?

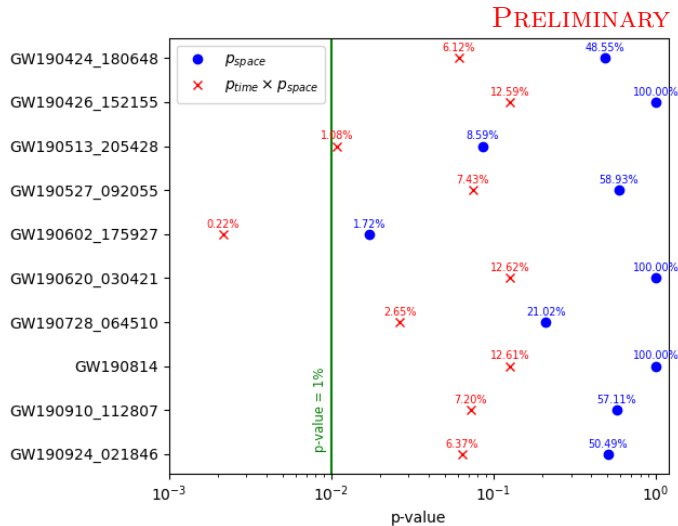
The p-value can be dissociated in $p = p_{\text{time}} \times p_{\text{space}}$, with:

- $p_{\text{time}} = \text{Prob}(N \geq 1) = 1 - e^{-n_B} \sim 12.6\%$ for $n_B = \text{total background}_{(\text{FC}+\text{PC}+\text{UPMU})} = 0.13$
- p_{space} is obtained by comparing neutrino direction and GW localisation*
 - For each sample ($k = \text{FC}, \text{PC}$ or UPMU), define the point-source likelihood $\mathcal{L}_\nu^{(k)}(n_S^{(k)}, \gamma; \Omega_S)$ that separates background from signal ($dn/dE \propto E^{-\gamma}$, direction Ω_S).
 - Compute the maximum log-likelihood ratio Λ (GW localisation \mathcal{P}_{GW} used as prior) and find the source direction Ω_S that maximises it:

$$\Lambda(\Omega_S) = 2 \sum_k \ln \left[\frac{\mathcal{L}_\nu(\widehat{n_S^{(k)}}, \widehat{\gamma^{(k)}}; \Omega_S)}{\mathcal{L}_\nu(n_S^{(k)} = 0; \Omega_S)} \right] + 2 \ln \mathcal{P}_{\text{GW}}(\Omega_S) \text{ and } \boxed{\text{TS} = \max_{\Omega} [\Lambda(\Omega)]}$$

- Compare TS_{data} with the expected background distribution (with $N \geq 1$) to obtain p_{space} .

*IceCube collaboration. IceCube Search for Neutrinos Coincident with Compact Binary Mergers from LIGO-Virgo's First Gravitational-wave Transient Catalog. *Astrophys.J.Lett.* 898 (2020) 1, L10



The most significant GW+ ν is for GW190602_175927:

$$p_{\text{space}} = 1.72\%, p = 0.22\%$$

Considering the number of trials ($N = 10$ GW+ ν time coincidences), we get **post-trial** p-value:

$$P_{\text{space}} = 15.9\%$$

For $p = p_{\text{time}} \times p_{\text{space}}$, the number of trials is the total number of GW follow-ups ($N = 36$) and we get:

$$P = 7.8\%$$

The neutrino flux is assumed as $\frac{dn}{dE_\nu} = \phi_0 E_\nu^{-2}$ and
 $N_{\text{expected signal}} = \int_{E_{\min}}^{E_{\max}} dE_\nu A_{\text{eff}}^{s,f}(E_\nu, \theta) \times \frac{dn}{dE_\nu}.$

Sample-by-sample flux limits

For each sample and flavour ($\nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu$), we define the flux likelihood:

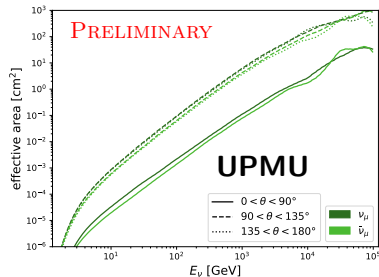
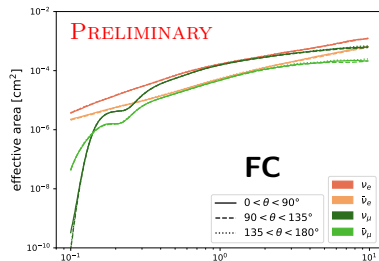
$$\mathcal{L}(\phi_0; n_B, N) = \int \frac{(c(\Omega)\phi_0 + n_B)^N}{N!} e^{-(c(\Omega)\phi_0 + n_B)} \mathcal{P}_{\text{GW}}(\Omega) d\Omega$$

with $c(\Omega) = \int_{E_{\min}}^{E_{\max}} dE_\nu A_{\text{eff}}(E_\nu, \theta) E_\nu^{-2}$ and the 90% U.L on the flux ϕ^{up} is obtained by solving $\int_0^{\phi^{\text{up}}} \mathcal{L}(\phi) d\phi = 0.9$

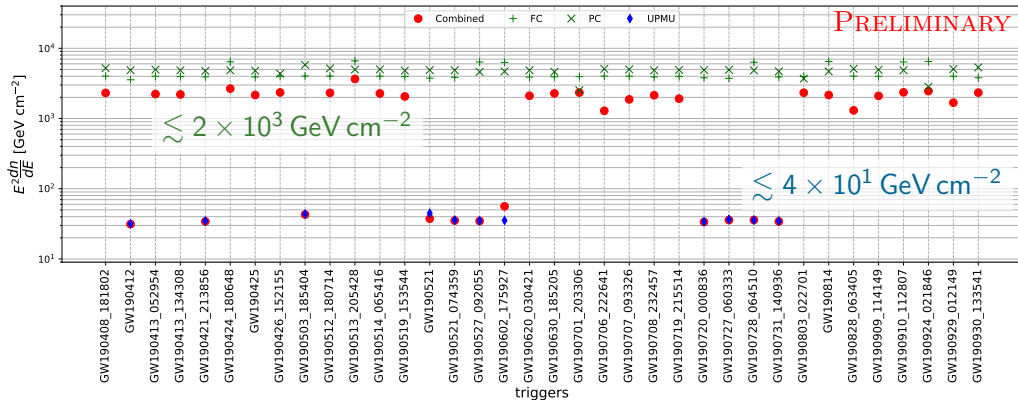
Combined flux limits

Limits combining FC, PC and UPMU are obtained by using the combined TS defined before (details in backup).

Effective area A_{eff}



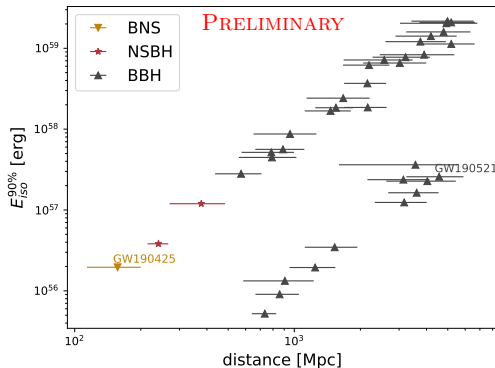
Example of limits for ν_μ flavour:



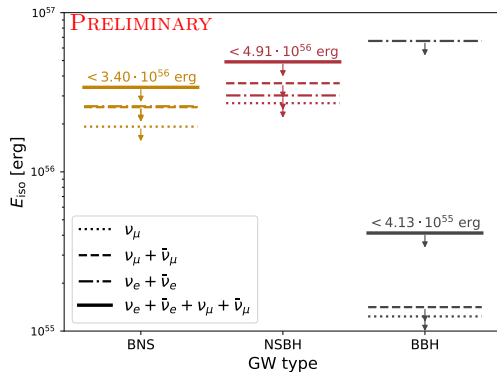
Better limits with the UPMU sample when the GW is below the local horizon. Combined limits is close to the best individual one.

- The total energy in ν from the source (assuming isotropic) is $E_{\text{iso}} = 4\pi d^2 \int \frac{dn}{dE} \times E dE$
 $\Rightarrow E_{\text{iso}}$ limits obtained by using the 3D localisation skymap from the LVC data release.
- We can stack events by nature, assuming same emission (or $E_{\text{iso}} \propto M_{\text{source}}$ in backup).

Individual limits on $E_{\text{iso}}^{\nu_\mu}$



Stacked limits on $E_{\text{iso}}^{\text{all-flavours}*}$



*This is done assuming the flux at Earth is equally distributed between the flavours ($\nu_e : \nu_\mu : \nu_\tau = 1 : 1 : 1$)

- For low-energy analysis, the case is simpler as effective area does not depend on direction.
- Upper limits on fluence are obtained assuming Fermi-Dirac ($\langle E \rangle = 20 \text{ MeV}$) or flat spectrum:

$$\Phi_{90} = \frac{N_{90}}{N_{\text{Target}} \int \lambda(E_\nu) \sigma(E_\nu) R(E_e, E_{\text{vis}}) \epsilon(E_{\text{vis}}) dE_\nu} \text{ with } \lambda(E_\nu) = \text{F.-D. or Const.}$$

- Typical fluence limits: $\begin{cases} \Phi(\nu_e) \lesssim 5 \times 10^9 \text{ cm}^{-2}, & \Phi(\bar{\nu}_e) \lesssim 1 \times 10^8 \text{ cm}^{-2} \\ \Phi(\nu_x) \lesssim 3 \times 10^{10} \text{ cm}^{-2}, & \Phi(\bar{\nu}_x) \lesssim 4 \times 10^{10} \text{ cm}^{-2} \end{cases} (\nu_x = \nu_{\mu, \tau})$
- E_{iso} limits are obtained as in the high-energy case, using the LVC distance estimate:
 $E_{\text{iso}}^{\bar{\nu}_e} < 9.59 \times 10^{57} \text{ erg for GW190425 } (d \sim 160 \text{ Mpc})$

- Follow-up analysis of GWTC-2 events have been done using the four SK samples (low-energy, FC, PC, UPMU).
- No excess has been observed with respect to expected background.
- Most significant observation is for GW190602_175927 (FC event ~ 300 s before the GW, with compatible direction) \Rightarrow **post-trial p-value is 7.8% (1.4σ)**
- Flux limits have been computed:
 - **High-Energy:** $E^2 \frac{dn}{dE} \Big|_{\nu_\mu} \lesssim 4 \times 10^1 \text{ GeV cm}^{-2}$ if GW is below the horizon (2×10^3 otherwise)
 - **Low-energy:** $\Phi(\bar{\nu}_e) \lesssim 1 \times 10^8 \text{ cm}^{-2}$
- E_{iso} were also extracted, independently event-by-event or stacking events of the same nature (assuming $E_{\text{iso}} = \text{Const.}$ or $E_{\text{iso}} \propto \mathcal{M}_{\text{tot}}$)
- **Publication coming soon** with detailed results for O3a events.
- **Future:** possible realtime follow-up (within few days) from O4

This presentation was made on behalf of the Super-Kamiokande collaboration



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Backups

PRELIMINARY

Trigger	Sample	Δt [s]	E [GeV]	RA [deg]	Dec [deg]	δ [deg]	p-value [%]
GW190424_180648	FC	104.03	0.57	210.82	-58.74	52.08	48.55
GW190426_152155	UPMU	278.99	9.52	352.37	-8.46	2.15	100.00
GW190513_205428	FC	-183.27	0.68	279.34	-37.27	41.19	8.59
GW190527_092055	FC	248.41	0.48	54.09	18.80	52.08	58.93
GW190602_175927	FC	-286.52	2.75	93.67	-38.90	16.22	1.72
GW190620_030421	UPMU	-327.70	2.33	177.69	-35.59	8.04	100.00
GW190728_064510	FC	102.99	0.19	300.45	29.71	92.51	21.02
GW190814	FC	250.36	1.21	157.59	-9.47	28.26	100.00
GW190910_112807	FC	301.42	1.08	160.13	-22.70	32.09	57.11
GW190924_021846	FC	411.87	0.30	281.38	-54.52	73.58	50.49

- **Flux:** We define the following likelihood by using the TS defined before:

$$\mathcal{L}(\phi_0; TS_{\text{data}}, \mathcal{P}_{GW}) = \int \sum_{k=0}^{\infty} \left[\frac{(c(\Omega)\phi_0)^k}{k!} e^{-c(\Omega)\phi_0} \times \mathcal{P}_k(TS_{\text{data}}) \right] \times \mathcal{P}_{GW}(\Omega) d\Omega$$

where $\mathcal{P}_i(TS)$ is the distribution of the test statistic assuming the signal consists in i events, assuming E^{-2} spectrum ($dn/dE = \phi_0 E^{-2}$). The 90% upper limit is obtained as above ($\int_0^{\phi_0^{\text{up}}} \mathcal{L}(\phi_0) d\phi_0 = 0.90$).

- **Total energy:** Same for E_{iso} limits:

$$\mathcal{L}(E_{\text{iso}}; TS_{\text{data}}^{(i)}, \mathcal{V}_{GW}^{(i)}) = \int \sum_{k=0}^{\infty} \left[\frac{(c'(r, \Omega)E_{\text{iso}})^k}{k!} e^{-c'(r, \Omega)E_{\text{iso}}} \times \mathcal{P}_k^{(i)}(TS_{\text{data}}^{(i)}) \right] \times \mathcal{V}_{GW}^{(i)}(r, \Omega) d\Omega$$

For each sample k , we define the likelihood:

$$\mathcal{L}_\nu^{(k)}(n_S^{(k)}, \gamma; \Omega_S) = \frac{e^{-(n_S^{(k)} + n_B^{(k)})} (n_S^{(k)} + n_B^{(k)})^{N^{(k)}}}{N^{(k)}!} \prod_{i=1}^{N^{(k)}} \frac{n_S^{(k)} \mathcal{S}^{(k)}(\vec{x}_i, E_i; \Omega_S, \gamma) + n_B^{(k)} \mathcal{B}^{(k)}(\vec{x}_i, E_i)}{n_S^{(k)} + n_B^{(k)}}$$

where $\mathcal{S}^{(k)}$ and $\mathcal{B}^{(k)}$ are the signal/background p.d.f. (characterizing detector response).

Then, we compute the log-likelihood ratio:

$$\Lambda(\Omega_S) = 2 \sum_k \ln \left[\frac{\mathcal{L}_\nu(\widehat{n_S^{(k)}}, \widehat{\gamma^{(k)}}; \Omega_S)}{\mathcal{L}_\nu(n_S^{(k)} = 0; \Omega_S)} \right] + 2 \ln \mathcal{P}_{GW}(\Omega_S)$$

The final test statistic and p-value are:

$$TS = \max_{\Omega} [\Lambda(\Omega)] \text{ and } p_{\text{space}} = \int_{TS_{\text{data}}}^{\infty} \mathcal{P}_{\text{bkg}}(TS) dTS$$

where $\mathcal{P}_{\text{bkg}}(TS)$ is the expected background distribution.

Trigger name	Sample		ν_e	$\bar{\nu}_e$	ν_μ (ν_x)	$\bar{\nu}_\mu$ ($\bar{\nu}_x$)
GW190425	HE $E^2 \frac{dn}{dE}$	FC	$2.22 \cdot 10^3$	$4.32 \cdot 10^3$	$3.91 \cdot 10^3$	$9.42 \cdot 10^3$
		PC	$3.32 \cdot 10^4$	$1.12 \cdot 10^5$	$4.81 \cdot 10^3$	$8.74 \cdot 10^3$
		UPMU	—	—	—	—
		Combined	$2.09 \cdot 10^3$	$4.28 \cdot 10^3$	$2.16 \cdot 10^3$	$4.20 \cdot 10^3$
	HE E_{iso}	Per-flavour	$1.98 \cdot 10^{56}$	$3.85 \cdot 10^{56}$	$1.96 \cdot 10^{56}$	$3.69 \cdot 10^{56}$
		$\nu + \bar{\nu}$	$2.62 \cdot 10^{56}$		$2.52 \cdot 10^{56}$	
		All		$3.47 \cdot 10^{56}$		
	LE Φ	Flat	$1.49 \cdot 10^9$	$1.83 \cdot 10^7$	$9.35 \cdot 10^9$	$1.11 \cdot 10^{10}$
		Fermi-Dirac	$3.92 \cdot 10^9$	$9.57 \cdot 10^7$	$2.43 \cdot 10^{10}$	$2.87 \cdot 10^{10}$
	LE E_{iso}	Per-flavour	$3.92 \cdot 10^{59}$	$9.59 \cdot 10^{57}$	$2.43 \cdot 10^{60}$	$2.87 \cdot 10^{60}$
		All		$5.54 \cdot 10^{58}$		

PRELIMINARY $E^2 \frac{dn}{dE}$ [in GeV cm^{-2}], Φ [in cm^{-2}], E_{iso} [in erg]

Trigger name	Sample		ν_e	$\bar{\nu}_e$	ν_μ (ν_x)	$\bar{\nu}_\mu$ ($\bar{\nu}_x$)
GW190521	HE $E^2 \frac{dn}{dE}$	FC	$2.27 \cdot 10^3$	$4.71 \cdot 10^3$	$3.76 \cdot 10^3$	$9.60 \cdot 10^3$
		PC	$3.66 \cdot 10^4$	$3.68 \cdot 10^4$	$4.89 \cdot 10^3$	$8.35 \cdot 10^3$
		UPMU	—	—	$4.48 \cdot 10^1$	$5.04 \cdot 10^1$
		Combined	$2.21 \cdot 10^3$	$4.60 \cdot 10^3$	$3.75 \cdot 10^1$	$4.82 \cdot 10^1$
	HE E_{iso}	Per-flavour	$1.69 \cdot 10^{59}$	$3.46 \cdot 10^{59}$	$2.58 \cdot 10^{57}$	$3.72 \cdot 10^{57}$
		$\nu + \bar{\nu}$	$2.26 \cdot 10^{59}$		$3.00 \cdot 10^{57}$	
		All		$8.94 \cdot 10^{57}$		
	LE Φ	Flat	$2.63 \cdot 10^9$	$3.22 \cdot 10^7$	$1.65 \cdot 10^{10}$	$1.95 \cdot 10^{10}$
		Fermi-Dirac	$6.89 \cdot 10^9$	$1.68 \cdot 10^8$	$4.27 \cdot 10^{10}$	$5.04 \cdot 10^{10}$
	LE E_{iso}	Per-flavour	$5.85 \cdot 10^{62}$	$1.43 \cdot 10^{61}$	$3.63 \cdot 10^{63}$	$4.28 \cdot 10^{63}$
		All		$8.26 \cdot 10^{61}$		

PRELIMINARY $E^2 \frac{dn}{dE}$ [in GeV cm^{-2}], Φ [in cm^{-2}], E_{iso} [in erg]

We combine the likelihoods within a given population*:

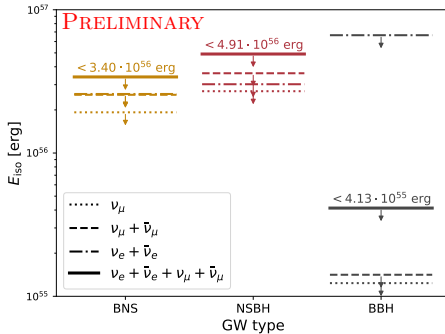
- Assuming same expected E_{iso} for all events:

$$\mathcal{L}^{\text{Pop}}(E_{\text{iso}}; \{TS_{\text{data}}\}^{(i)}, \{\nu_{\text{GW}}^{(i)}\}) = \prod_{i=1}^N \mathcal{L}(E_{\text{iso}}; TS_{\text{data}})^{(i)}, \nu_{\text{GW}}^{(i)})$$

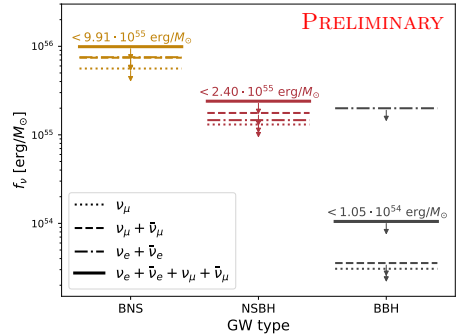
- Assuming neutrino emissionsscales with object total mass \mathcal{M}_{tot} :

$$\mathcal{L}^{\text{Pop}}(f_{\nu}; \{TS_{\text{data}}\}^{(i)}, \{\nu_{\text{GW}}^{(i)}\}, \{\mathcal{M}_{\text{tot}}^{(i)}\}) = \prod_{i=1}^N \int \mathcal{M}_{\text{tot}}^{(i)} \mathcal{L}(f_{\nu} \mathcal{M}_{\text{tot}}^{(i)}; TS_{\text{data}})^{(i)}, \nu_{\text{GW}}^{(i)}) p_{\text{GW}}(\mathcal{M}_{\text{tot}}^{(i)}) d\mathcal{M}_{\text{tot}}^{(i)}$$

assuming same E_{iso}



assuming $E_{\text{iso}} \propto \mathcal{M}_{\text{tot}}$



Experiment	Super-Kamiokande	ANTARES	IceCube
Energy range	0.1-10 ⁵ GeV	TeV-PeV	10-10 ^{9.5} GeV
$E^2 dn/dE$ limits (min)	4×10^1 GeV cm ⁻²	1 GeV cm ⁻²	0.03 GeV cm ⁻²
$E^2 dn/dE$ limits (max)	2×10^3 GeV cm ⁻²	9 GeV cm ⁻²	0.6 GeV cm ⁻²
Reference	this work	Poster @CRνMM	PoS-ICRC2019-918

This is assuming E^{-2} . The situation will be in favour of SK for $\gamma > 2$ (e.g. E^{-3}).