Enhanced violation of Leggett-Garg Inequality in three flavour neutrino oscillations via non-standard interactions

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Introduction

- Neutrino oscillate among themselves and these oscillations have their origin in the non-zero neutrino masses and mixing among the neutrino flavors.
- The standard paradigm of neutrino oscillations involves three flavours of neutrinos which are superpositions of the mass states carrying well-defined masses.
- Effective Hamiltonian for neutrino porpagation

$$\mathcal{H} = \mathcal{H}_{\rm vac} + \mathcal{H}_{\rm SI} + \mathcal{H}_{\rm NSI}$$

where $\mathcal{H}_{\rm vac}$ is the vacuum Hamiltonian and $\mathcal{H}_{\rm SI}, \mathcal{H}_{\rm NSI}$ are the effective Hamiltonians in presence of standard interaction (SI) and NSI respectively.

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In 1985, Leggett and Garg derived a class of inequalities which have the following assumptions:

Macroscopic realism (MR):

A macroscopic system with two or more macroscopically distinct states available to it will at all times be in one or the other of these states.

Non-Invasive measurability (NIM):

It is possible, in principle, to determine which of the states the system is in, without affecting the states itself or the system's subsequent dynamics.

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• Dichotomic observable: $Q = \pm 1$

$$\begin{array}{ccc} \underline{\mathcal{Q}}_1 & \underline{\mathcal{Q}}_2 & \underline{\mathcal{Q}}_3 \\ \hline t_1 & t_2 & t_3 \end{array} \quad \text{time}$$

- Two time correlation functions $C_{ij} = \frac{1}{N} \sum_{q=1}^{N} \langle Q_i^q Q_j^q \rangle$
- Macrorealism restricts the following combination of two time correlation functions:

$$\begin{split} \mathcal{K}_3 &= \mathcal{C}_{12} + \mathcal{C}_{23} - \mathcal{C}_{31} &= \langle \mathcal{Q}_1 \mathcal{Q}_2 \rangle + \langle \mathcal{Q}_2 \mathcal{Q}_3 \rangle - \langle \mathcal{Q}_1 \mathcal{Q}_3 \rangle \\ \mathcal{K}_3 &= \langle \mathcal{Q}_1 \mathcal{Q}_2 \rangle + \langle [\mathcal{Q}_2 - \mathcal{Q}_1] \mathcal{Q}_3 \rangle \\ \mathcal{K}_3 &= \begin{cases} 1 + 0 = 1 \\ -1 + (\pm 2) = 1 & \text{or} - 3 \end{cases} \end{split}$$

This gives the condition

$$-3 \leq K_3 \leq 1$$

In general we have

$$-n \leq K_n \leq (n-2)$$
 $3 \leq n, \text{odd};$
 $-(n-2) \leq K_n \leq (n-2)$ $4 \leq n, \text{even}$

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Existing literature on LGI in Neutrino Sector

- D. Gangopadhyay, D. Home, and A. Sinha Roy. Probing the Leggett-Garg Inequality for Oscillating Neutral Kaons and Neutrinos. Phys. Rev., A88(2):022115 2013
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- Debashis Gangopadhyay and Animesh Sinha Roy. Three-flavoured neutrino oscillations and the LeggettGarg inequality. Eur. Phys. J., C77(4):260, 2017.
- Javid Naikoo, Ashutosh Kumar Alok, Subhashish Banerjee, S. Uma Sankar, Giacomo Guarnieri, Christiane Schultze, and Beatrix C. Hiesmayr. A quantum information the- oretic quantity sensitive to the neutrino mass-hierarchy. Nucl. Phys. B, 951:114872, 2020.
- Javid Naikoo, Ashutosh Kumar Alok, Subhashish Banerjee, and S. Uma Sankar. Leggett-Garg inequality in the context of three flavour neutrino oscillation. 2019

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• Let the initial state of neutrino be prepared in a specific flavor, say muon neutrino $|\nu_{\mu}\rangle$. Then we have

$$egin{array}{rcl} Q &=& egin{cases} +1 & ext{for }
u_\mu \ -1 & ext{for }
u_e ext{ or }
u_ au \end{array}$$

 The correlation function C₁₂ can be evaluated by using all the 9 joint probabilities as

$$C_{ij} = P_{\nu_e,\nu_e}(L_i, L_j) - P_{\nu_e,\nu_\mu}(L_i, L_j) - P_{\nu_e,\nu_\tau}(L_i, L_j) - P_{\nu_\mu,\nu_e}(L_i, L_j) + P_{\nu_\mu,\nu_\mu}(L_i, L_j) + P_{\nu_\mu,\nu_\tau}(L_i, L_j) - P_{\nu_\tau,\nu_e}(L_i, L_j) + P_{\nu_\tau,\nu_\mu}(L_i, L_j) + P_{\nu_\tau,\nu_\tau}(L_i, L_j)$$

where $P_{\nu_{\alpha}\nu_{\beta}}(L_i, L_j) = P_{\nu_{\mu} \rightarrow \nu_{\alpha}}(L_i)P_{\nu_{\alpha} \rightarrow \nu_{\beta}}(L_j)$

• For maximizing the LGI parameters we take $L_1 = 140.15$ km, and $(L_2 - L_1) = (L_3 - L_2) = (L_4 - L_3) = \Delta L$. Energy is taken to be 1 GeV and $\delta = 3\pi/2$. NH is assumed unless otherwise stated.

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Results



Figure: SI-K₃

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Figure: SI-K4

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Figure: No dependence on δ

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Figure: Mass hierarchy dependence

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$$\mathcal{H} = \frac{1}{2E} \mathcal{U} \begin{pmatrix} 0 & \\ & \delta m_{21}^2 & \\ & & \delta m_{31}^2 \end{pmatrix} \mathcal{U}^{\dagger}$$

$$+ \frac{A(x)}{2E} \begin{pmatrix} 1 + \varepsilon_{ee} & \varepsilon_{e\mu} & \varepsilon_{e\tau} \\ \varepsilon_{e\mu}^* & \varepsilon_{\mu\mu} & \varepsilon_{\mu\tau} \\ \varepsilon_{e\tau}^* & \varepsilon_{\mu\tau}^* & \varepsilon_{\tau\tau} \end{pmatrix}$$

where $A(x) = 2E\sqrt{2}G_F n_e(x)$ is the standard charged current potential and \mathcal{U} is the mixing matrix.



Figure: Non-standard senario

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Figure: SI and NSI

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Summary

- Exploring Leggett-Garg Inequalities in neutrino sector has been a topic of interest for the past decade.
- Neutrino oscillations, being a quantum mechanical phenomenon, violate LGI.
- LGI parameter does not depend on the CP-violating factor δ .
- It does depend on the mass hierarchy and at large values of ΔL, the curves for the two hierarchies are out of phase with each other.
- In the NSI senario, for $\epsilon_{e\mu} \neq 0$, we achieve an enhancement in the LGI parameter and for $\epsilon_{e\tau} \neq 0$, the LGI parameter is supressed.
- Varying the other NSI parameters, we maximise the value of the LGI parameter for NSI case and note a significant enhancement in the violation of LGI in the NSI senario as compared to the SI senario.
- For more details, please refer to arXiv:2009.12328 [hep-ph]

Thank You



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