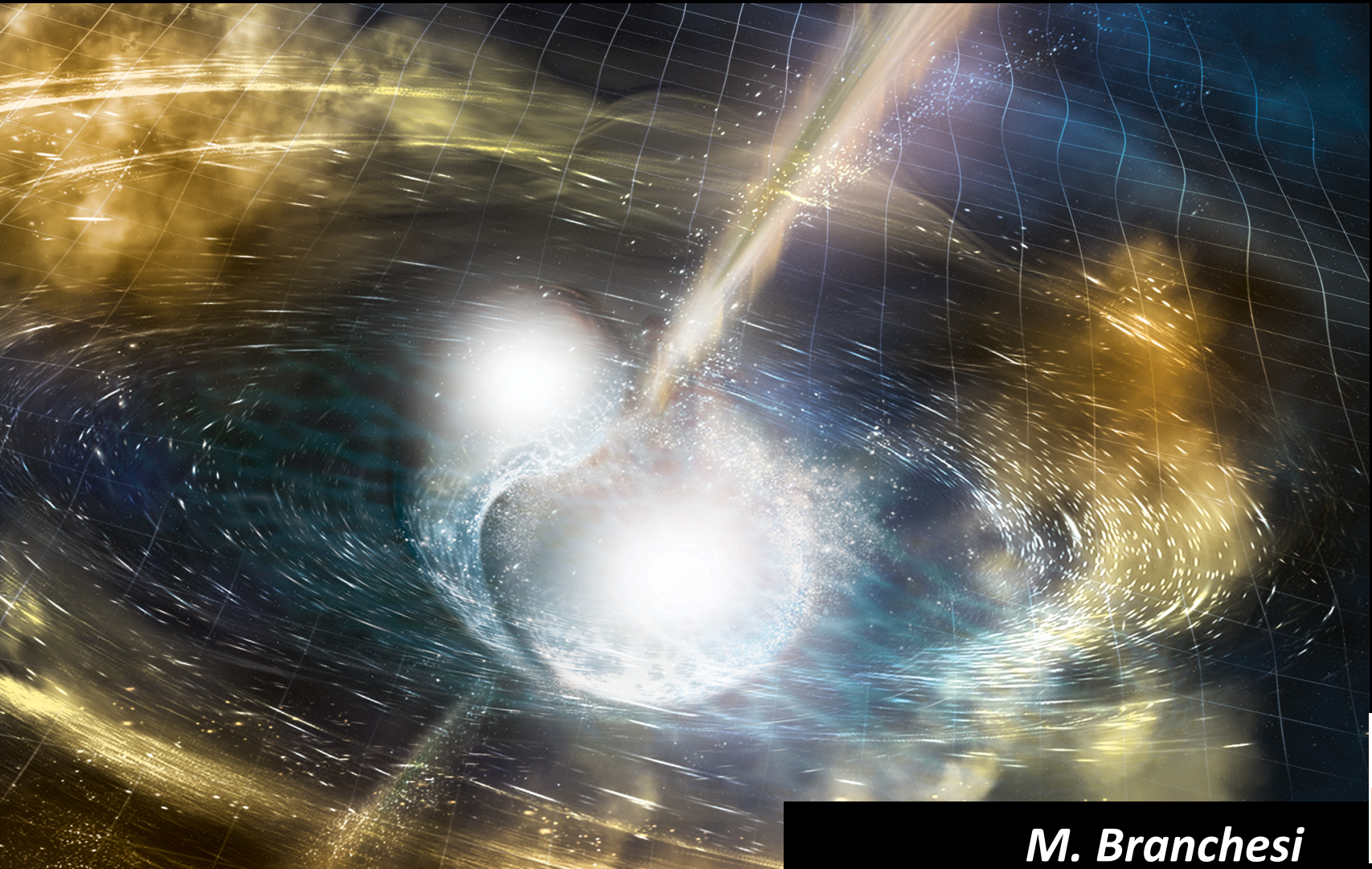


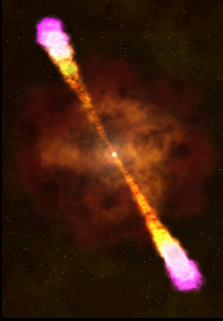
Multimessenger Astronomy with Gravitational Waves and Electromagnetic Signals



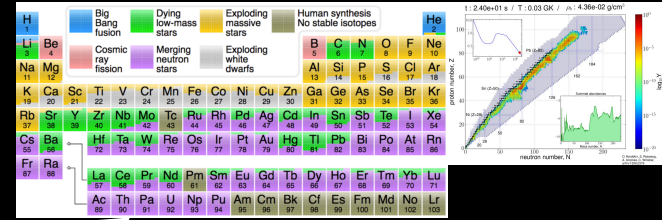
M. Branchesi
Gran Sasso Science Institute
INFN/LNGS and INAF

Radioactively powered transients

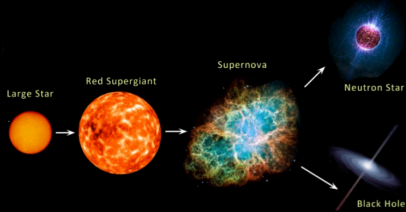
Relativistic astrophysics



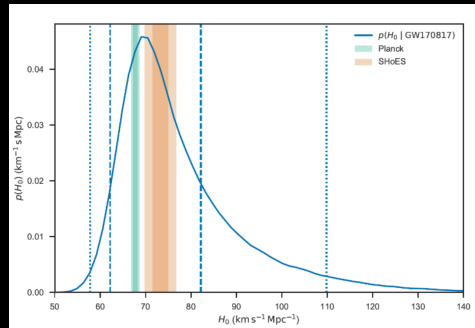
Nucleosynthesis and enrichment of the Universe



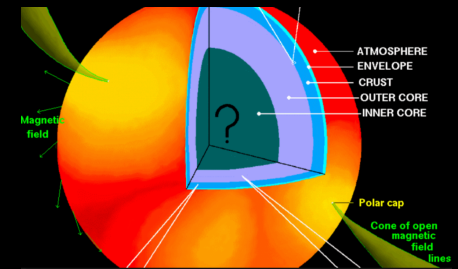
Compact object formation and evolution



Cosmology



Nuclear matter physics



A new window into the Universe

Earth



KAGRA, Japan



Credit: LIGO–Virgo



LIGO, Livingston, LA



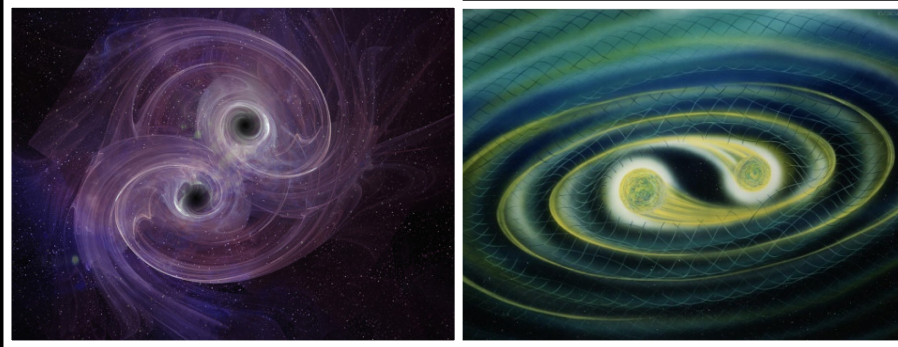
LIGO, Hanford, WA



Virgo, Cascina, Italy

ASTROPHYSICAL SOURCES emitting transient GW signals detectable by LIGO and Virgo (10-1000 Hz)

Coalescence of binary system of neutron stars and/or stellar-mass black-hole



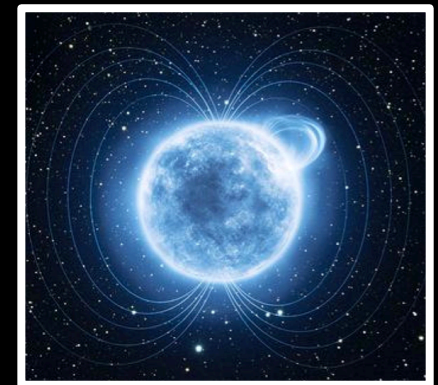
- Orbital evolution and GW signals accurately modeled by post-Newtonian approximation and numerical simulations
→ precise waveforms
- Energy emitted in GWs (BNS): $\sim 10^{-2} M_{\odot} c^2$

Isolated neutron-star instabilities

- Modeling of the GW shape and strength is complicated → uncertain waveforms
- Energy emitted in GWs:
 $\sim 10^{-9} M_{\odot} c^2$ for the core-collapse
 $\sim 10^{-16} - 10^{-6} M_{\odot} c^2$ for isolated NSs

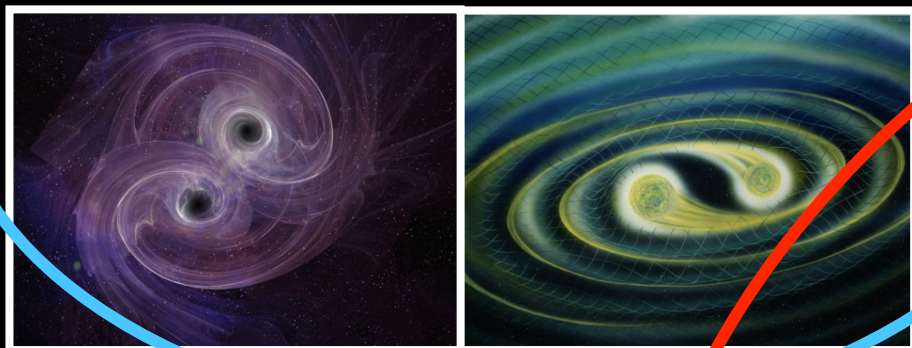


Core-collapse of massive stars



ASTROPHYSICAL SOURCES emitting transient GW signals detectable by LIGO and Virgo (10-1000 Hz)

Coalescence of binary system of neutron stars and/or stellar-mass black-hole

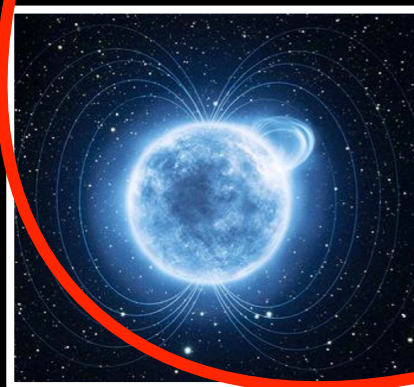


MATCHED-FILTER MODEL SEARCHES

Core-collapse of massive stars

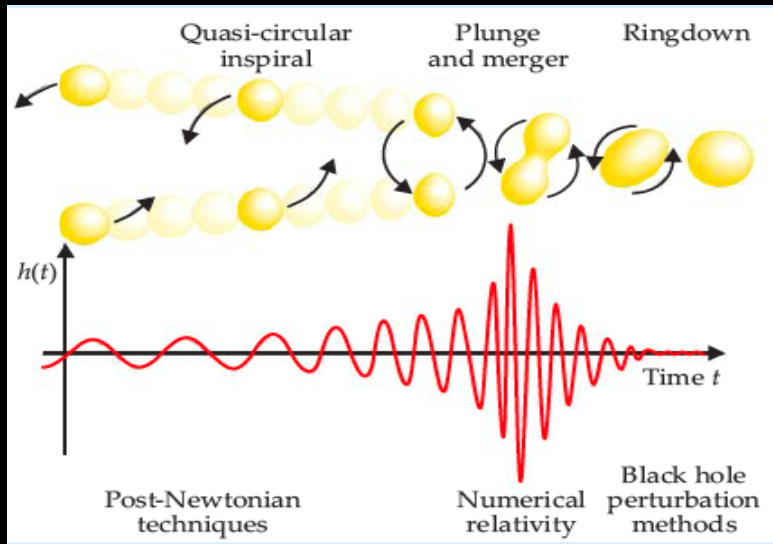


Isolated neutron-star instabilities



UNMODELED SEARCHES

Modelled compact binary coalescence searches



Waveforms depend on

- **intrinsic parameters: masses and spins** of the binary system (plus eccentricity, NS compactness, tidal deformability)
- **extrinsic parameters** that describe location, distance, merger time and system orientation with respect to an observer

Detection phase: known waveforms → **MATCHED FILTERING**

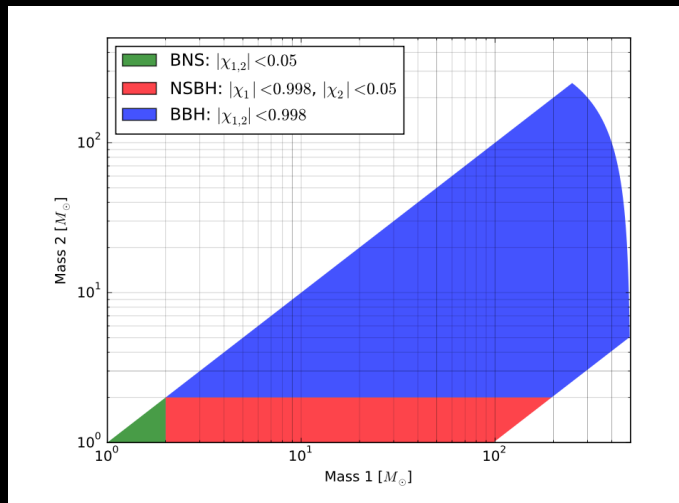
- Using waveform templates for a range of intrinsic parameters (masses and spin)
- “Extrinsic” parameters absorbed in overall amplitude

After detection → **Source PARAMETER RECONSTRUCTION:**

- Algorithms to explore the full-parameter space and find most likely values for sky location, masses, distance, orientation, spin...

Matched filtering searches

Template bank



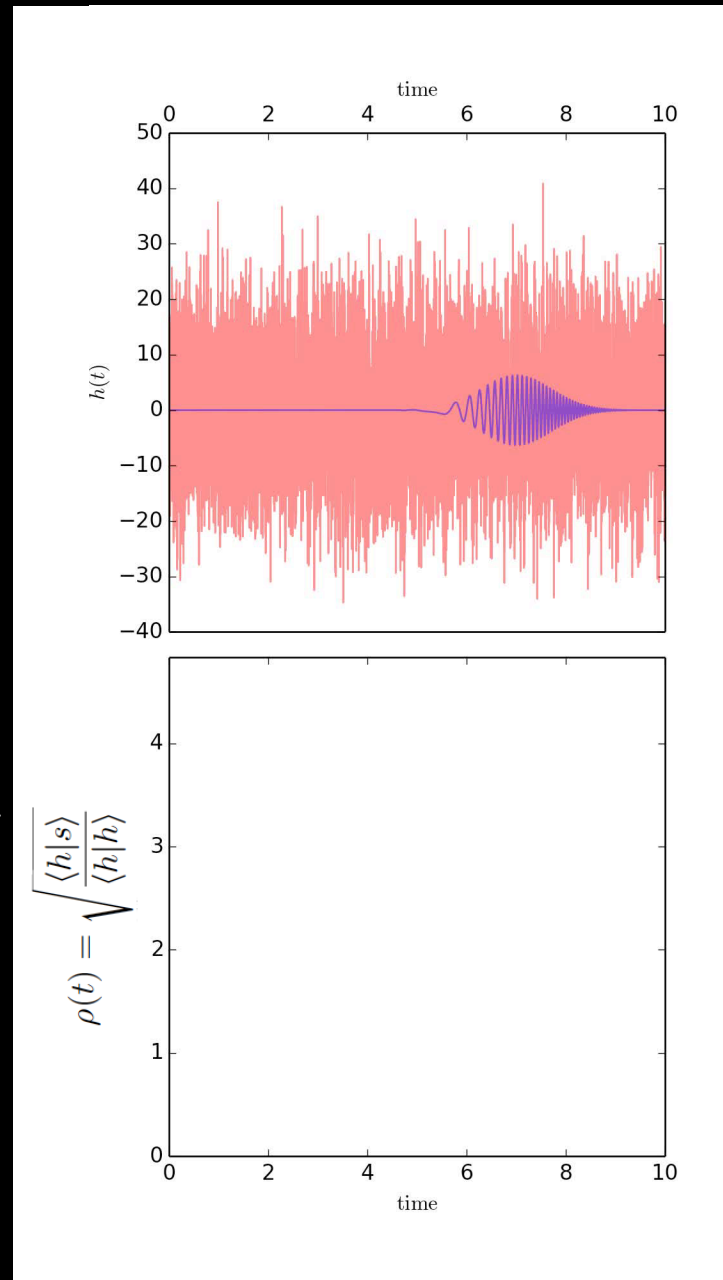
LVC Phys. Rev. X 6 (2016)

TEMPLATE

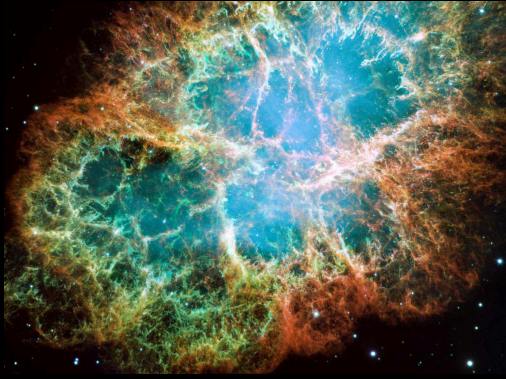
OBSERVATIONS

$$\langle h|s \rangle(t) = 4\Re \int_0^\infty \frac{\tilde{h}^*(f)\tilde{s}(f)e^{2\pi ift}}{S_n(f)} df$$

NOISE MODEL



Unmodeled GW transient searches



Transient sources:

- Core-collapse of massive stars
- Cosmic strings
- Neutron star instabilities
- Intermediate Massive BH
- ... the unknown

Poorly modelled
→ Can't use matched filtering

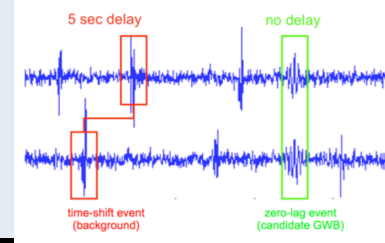
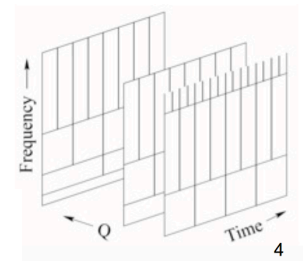
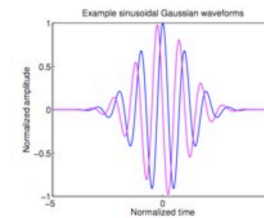
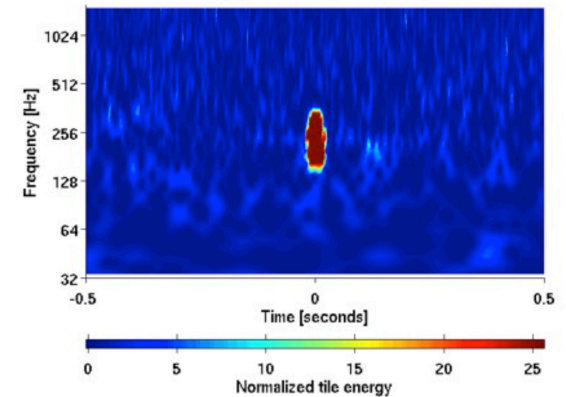
Detection without unknown waveform

→ **LOOK FOR "EXCESS POWER"**

All-sky, all-time search for transient as increase in power (hot pixels) in time-frequency map, minimal assumptions:

1. Duration: 1 ms to 1 s (characteristic time scale for stellar mass objects) → now also to a few hundreds of sec
2. Frequency: 10 to 5000 Hz (determined by detector's sensitivity)
3. **Signal appears coherently** in multiple detectors, consistent with antenna pattern → coincidence, coherent statistics, sky location

Noise fluctuations can be eliminated based on their non-correlation between detectors





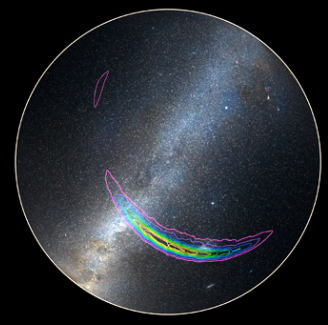
Low-latency GW data analysis pipelines to promptly identify GW candidates and send GW alerts



GW candidates

Sky Localization

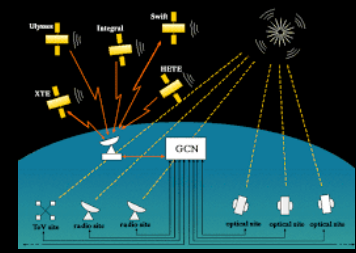
EM facilities



Low-latency search
to identify the GW-candidates

Software to

- select statistically significant triggers wrt background
- check detector sanity and data quality
- determine source localization

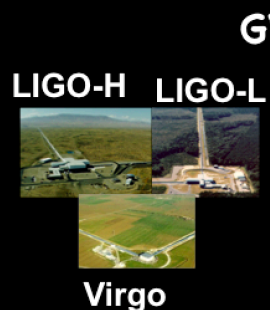


—————> **a few min** —————> **30 min**

O1, O2 run

LVC 2019, ApJ, 875, 161

O3 run started in April



GW candidates



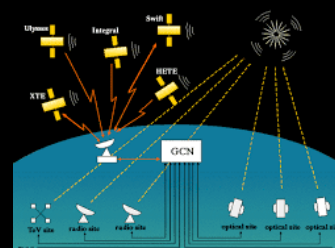
Sky Localization



EM facilities



Event validation
RETRACTION
or
CONFIRMATION
GCN



→ within min → 30 min- few hrs

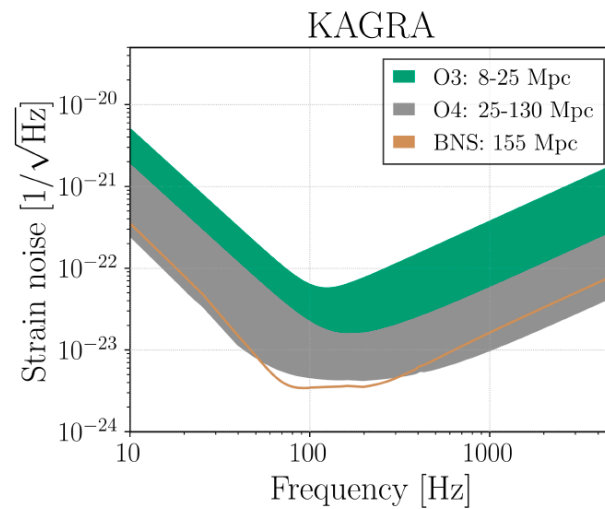
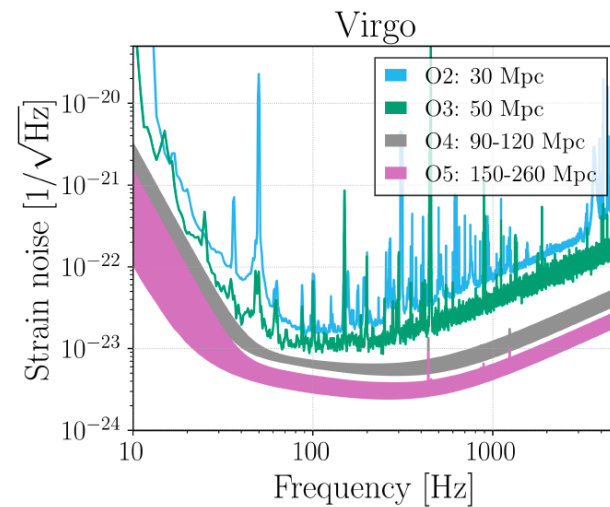
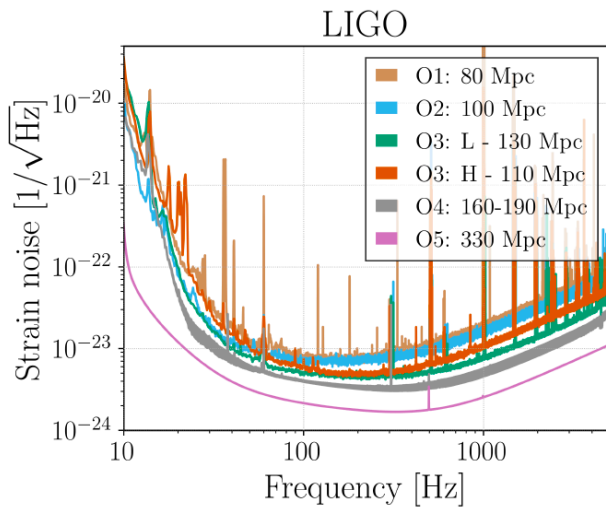
Parameter estimation codes

Hours, days

GW candidate updates

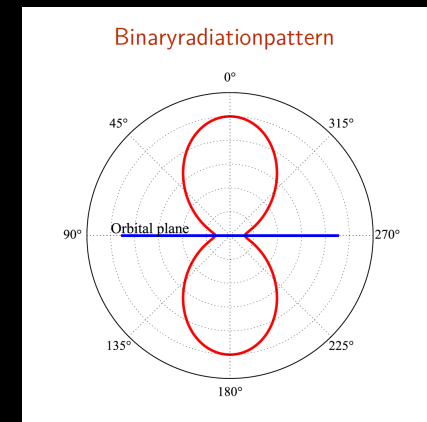
What is important to know to set-up the EM counterpart search?

Strain sensitivities as a function of frequency



SENSITIVITY IN TERMS OF RANGE/HORIZON DISTANCE

- **Range**: the volume- and orientation-averaged distance at which a compact binary coalescence gives a matched filter SNR of 8 in a single detector
- **Distance for face-on system**: distance at which an optimally oriented system (orbital plane perpendicular to the line of sight) would be observed with an SNR of 8: **range x 1.5**
- **Horizon**: distance at which an optimally oriented and located binary system would be observed with an SNR of 8: **range x 2.26**



RANGES corresponding to the orientation-averaged spacetime volumes surveyed per unit detector time

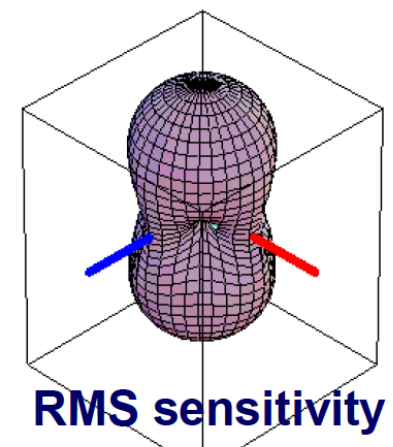
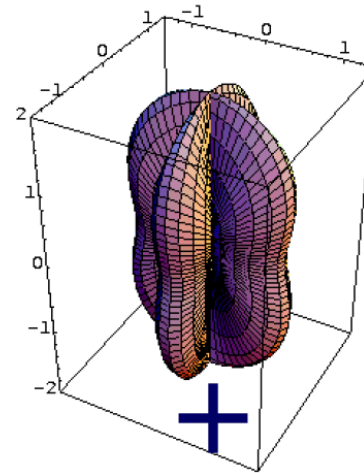
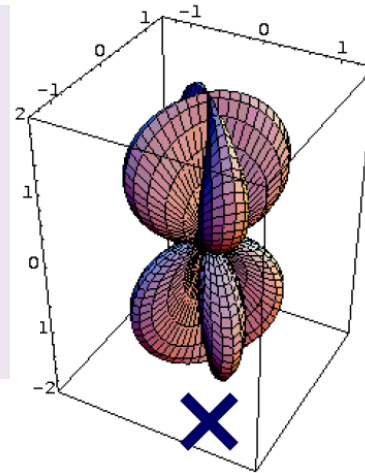
SNR = 8 in each detector

		O1	O2	O3	O4	O5	
1.4 Mo+1.4 Mo	BNS Range (Mpc)	aLIGO	80	100	110 – 130	160 – 190	330
		AdV	-	30	50	90 – 120	150 – 260
		KAGRA	-	-	8 – 25	25 – 130	130+
30 Mo+30 Mo	BBH Range (Mpc)	aLIGO	740	910	990 – 1200	1400 – 1600	2500
		AdV	-	270	500	860 – 1100	1300 – 2100
		KAGRA	-	-	80 – 260	260 – 1200	1200+
1.4 Mo+10 Mo	NSBH Range (Mpc)	aLIGO	140	180	190 – 240	300 – 330	590
		AdV	-	50	90	170 – 220	270 – 480
		KAGRA	-	-	15 – 45	45 – 290	290+
Burst Range (Mpc) [$E_{\text{GW}} = 10^{-2} M_{\odot} c^2$]		aLIGO	50	60	80 – 90	110 – 120	210
		AdV	-	25	35	65 – 80	100 – 155
		KAGRA	-	-	5 – 25	25 – 95	95+
Burst Range (kpc) [$E_{\text{GW}} = 10^{-9} M_{\odot} c^2$]		aLIGO	15	20	25 – 30	35 – 40	70
		AdV	-	10	10	20 – 25	35 – 50
		KAGRA	-	-	0 – 10	10 – 30	30+

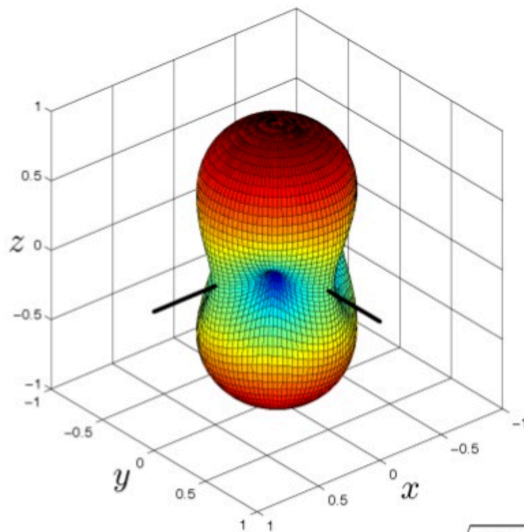
Sky location - single GW detector directional sensitivity

$$\frac{\Delta L}{L} = h_{\text{det}}(t) = F_+ h_+(t) + F_x h_x(t)$$

The **antenna pattern** depends on the polarization in a certain (x,+) basis



$$\sqrt{F_+(\theta, \phi)^2 + F_x(\theta, \phi)^2}$$

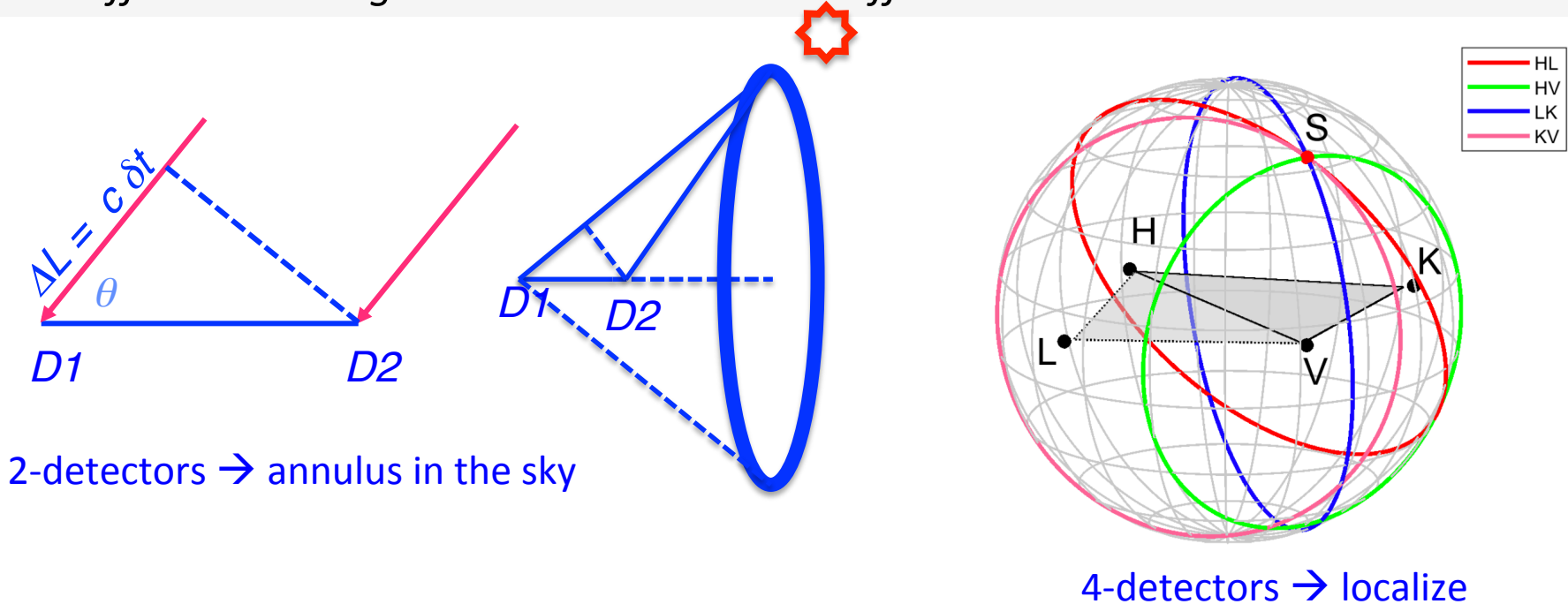


$$\sqrt{F_+^2(\theta, \phi, \psi = 0) + F_x^2(\theta, \phi, \psi = 0)}$$

- Single GW detector is a **good all-sky monitor**, nearly omni-directional (the transparency of Earth to GWs)
- But does not have good directional sensitivity, **not a pointing instrument!** It has a very poor angular resolution (about 100 deg)

The source localization requires a network of GW detectors

The **sky position** of a GW source is mainly **evaluated by triangulation**, measuring the differences in signal arrival times at the different network detector sites



The localization capability improves with signal SNR \rightarrow the sky localization area scales inversely with the square of the SNR

Compact binary Coalescence (CBC) Sky localization map

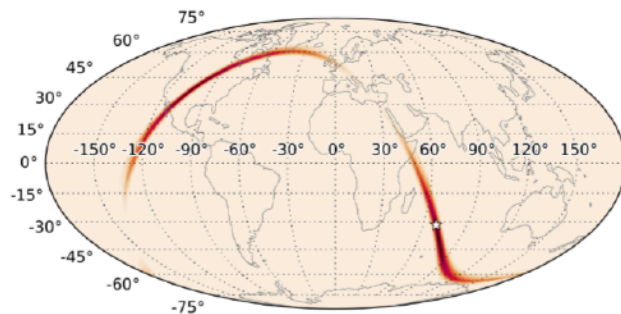
Arrival time
Amplitudes
Phase

→ sky location

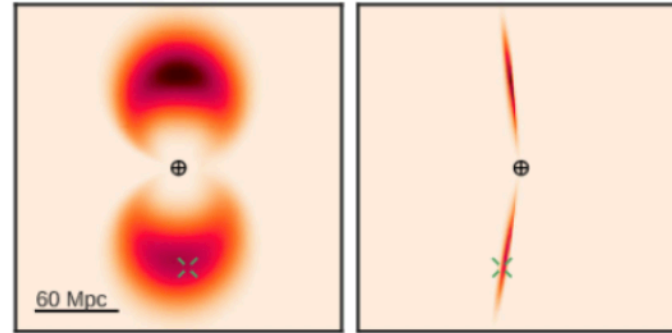
→ distance to the source

→ binary orientation

→ **Sky location also in 3 D**



Sky direction



Projections of 3d location

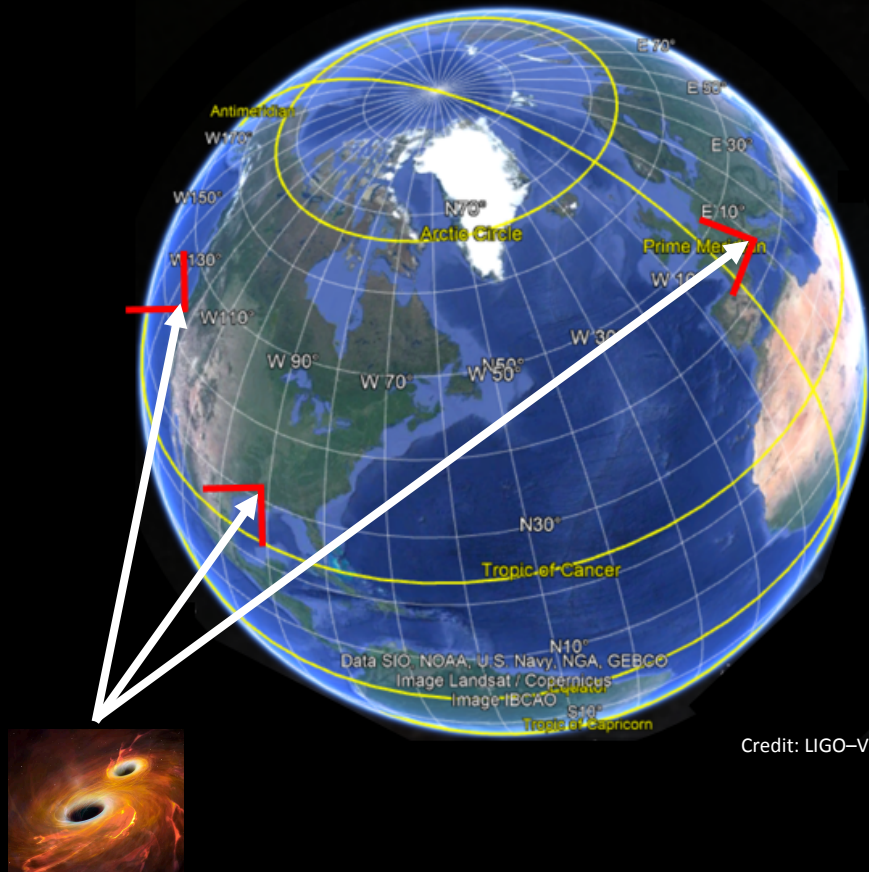
Online pipelines estimate → arrival time, phase, signal amplitude at each detector

These estimates + template masses : constrain direction of GW arrival and distance to the source

→ **BAYESTAR** (Singer et al 2014, ApJ, 795, 2016 ApJL, 829): estimate 3D location in <1 minute

→ **LALInference, full PE Bayesian MCMC** (Veitch 2015; Berry et al. 2015), modeling the inspiral-merger-ring down phase and taking into account the calibration uncertainty

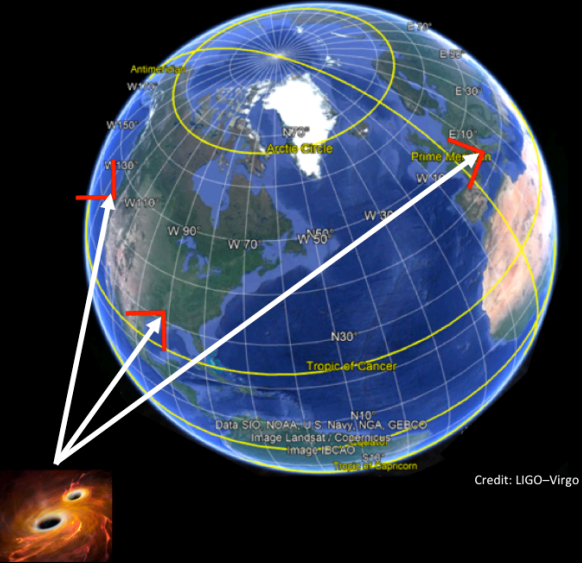
2017 August 14, 10:30:43 UT



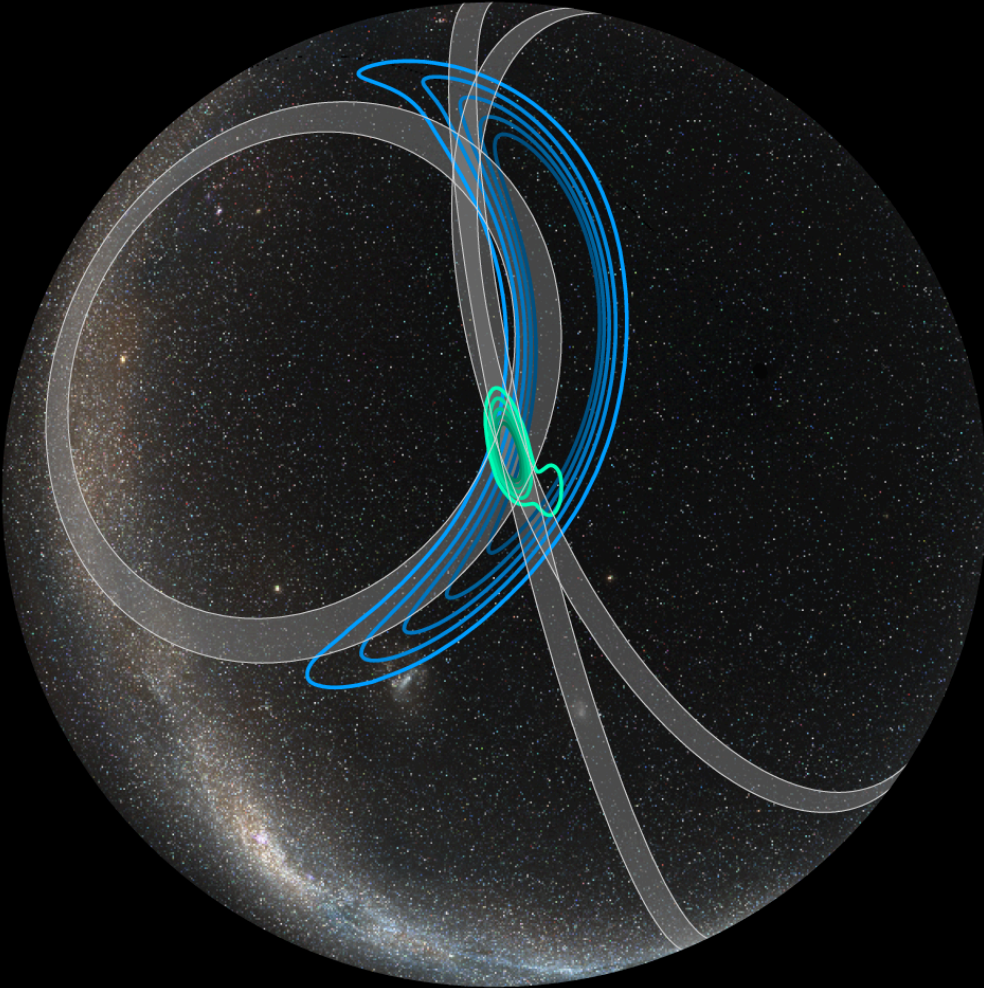
Credit: LIGO-Virgo

Virgo observed its first BBH
coalescence, GW170814

2017 August 14

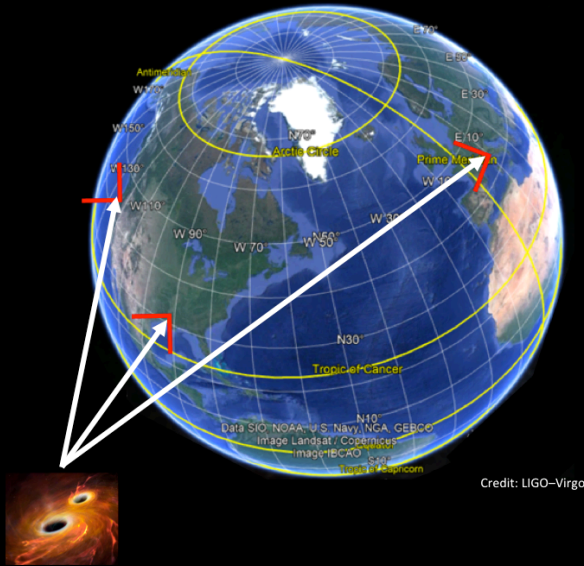


Credit: LIGO-Virgo

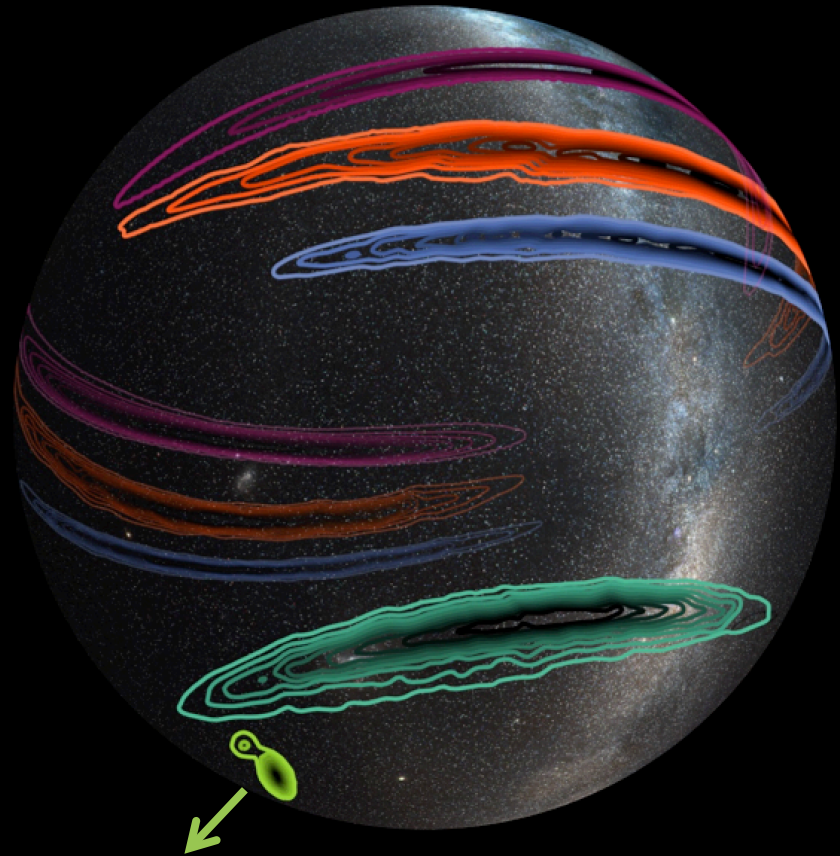


Credit: Leo Singer

2017 August 14



Credit: LIGO-Virgo



GW170814

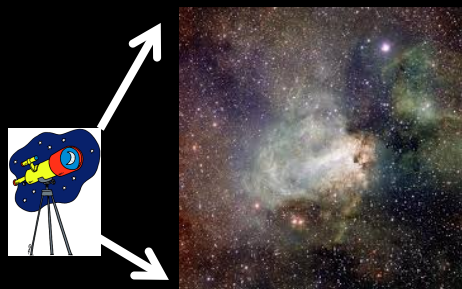
Credit: LIGO/Virgo/NASA/Leo Singer
(Milky Way image: Axel Mellinger)

LH 1160 square degrees

LHV 60 square degrees

LHV 60 square degrees

Hunt the elusive EM-counterpart!



Wide-field telescope
FOV >1 sq.degree

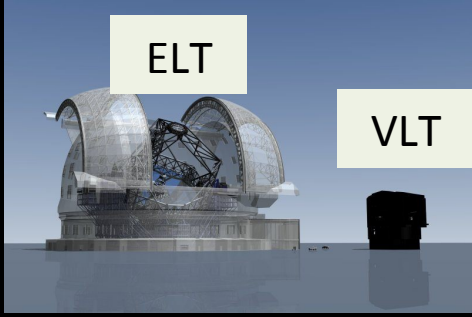


to cover hundreds/thousands
of square degrees



**“Fast” and “smart”
software** to select a
sample of candidate
counterparts

to remove transients
contaminants

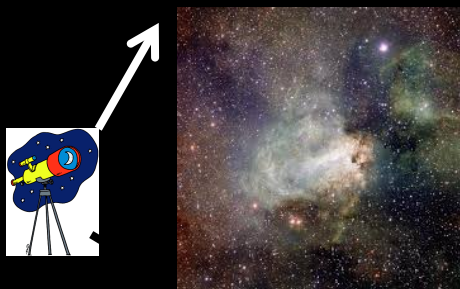


**Larger telescope to
characterize
the candidate nature**

To obtain observational time
for the characterization

**The EM
Counterpart!**

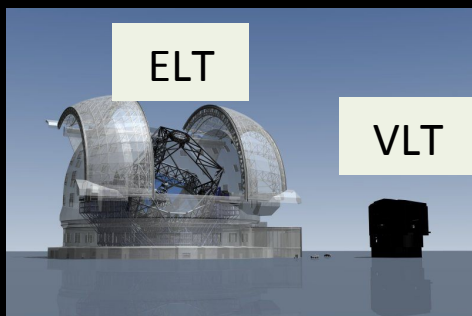
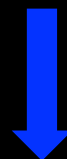
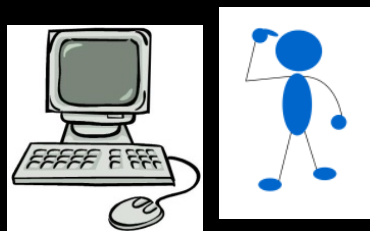
Hunt the elusive EM-counterpart!



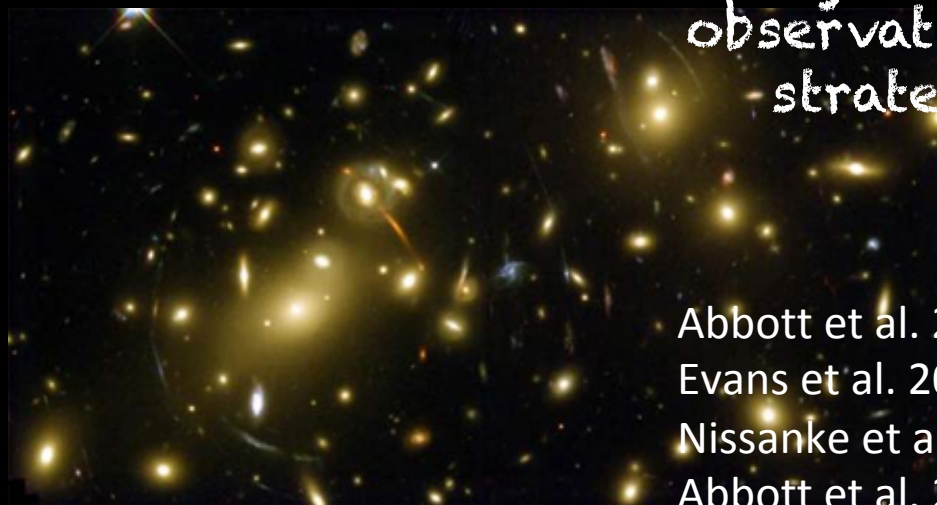
Wide-field telescope
FOV >1 sq.degree



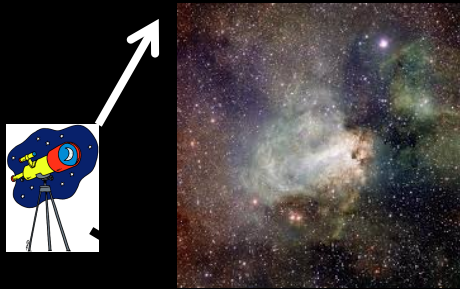
to cover hundreds/thousands
of square degrees



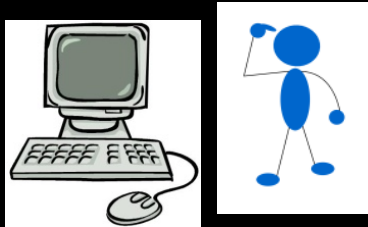
Galaxy-targeting
observational
strategy



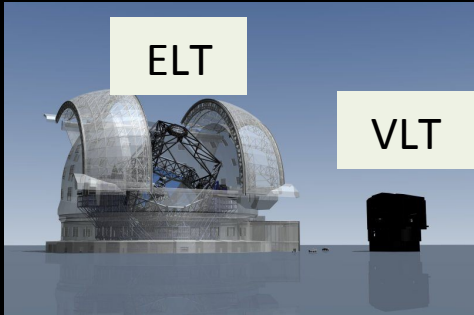
Abbott et al. 2012 A&A
Evans et al. 2012
Nissanke et al. 2013
Abbott et al. 2014 ApJS
Gehrels et al. 2016



Wide-field telescope
FOV >1 sq.degree



“Fast” and “smart” software to select a sample of candidate counterparts



Larger telescope to characterize the candidate nature

Optical/NIR band

10^4 - 10^5 variable objects over 100 sq. degrees



Artifacts and many astrophysical contaminants

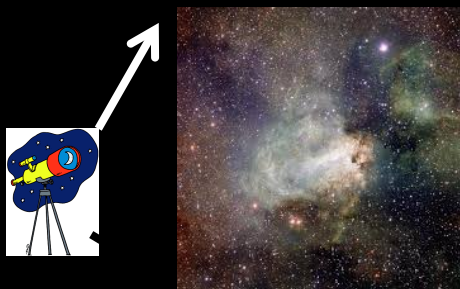
M-dwarf flares (min to hrs)
3 (0.3) deg⁻² up to red mag 24
at 20 (80) deg latitude
(Ridgway et al., 2014)

Supernovae (days to month)
7 deg⁻² up to red mag 24
(Graur et al., 2014; Dahlen et al., 2012; Cappellaro, 2014)



A few tens of candidate counterparts

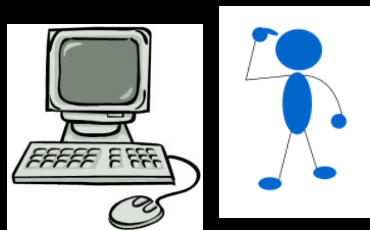
Hunt the elusive EM-counterpart!



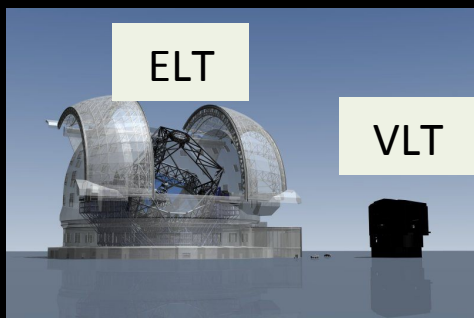
Wide-field telescope
FOV >1 sq.degree



“Fast” and “smart” software to select a sample of candidate counterparts



Larger telescope to characterize
the candidate nature



X-rays

✓ less contaminants

Transient rate $2.5 \times 10^{-3} \text{ deg}^{-2}$
 $\text{flux}_{0.2-2\text{KeV}} > 3 \times 10^{-12} \text{ ergs}^{-1} \text{ cm}^{-2}$
(Kanner et al. 2013)

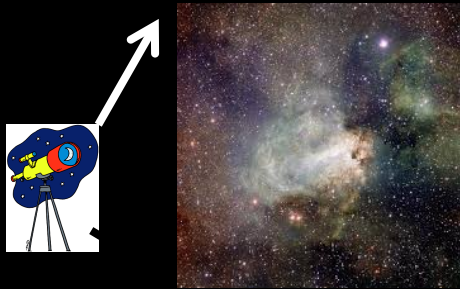
✗ no wide-field telescope

Gamma-rays

✓ less contaminants

✓ all-sky monitors

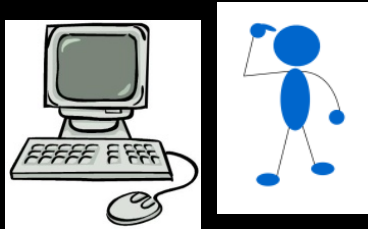
✗ beamed emission



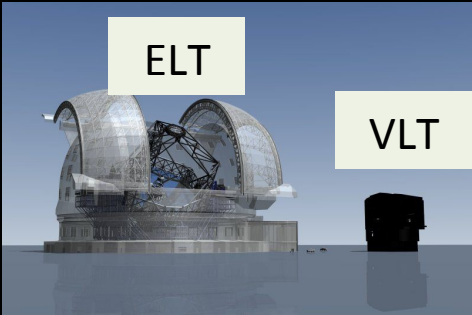
Wide-field telescope
FOV >1 sq.degree



“Fast” and “smart” software to select a sample of candidate counterparts



Larger telescope to characterize the candidate nature



The EM Counterpart!

Radio

✓ less contaminants

Transient rate < 0.37 deg⁻²
peak-flux_1.4 GHz > 0.21 mJy
timescales 1 d – 3 m
(Mooley et al., 2013)

✓ wide-field array at low frequencies (MHz)

✗ faint sources

✗ long delay GW → radio emission

What can help to define an
observational strategy?

ALERT CONTENTS to support observing strategy

- Estimate of **FALSE ALARM RATE** (FAR) of the event candidate
FAR=Rate of noise events louder than the candidate event
- Event **TIME** and **LOCALIZATION** given as a posterior probability distribution of the source's sky position(HEALPix FITS file)

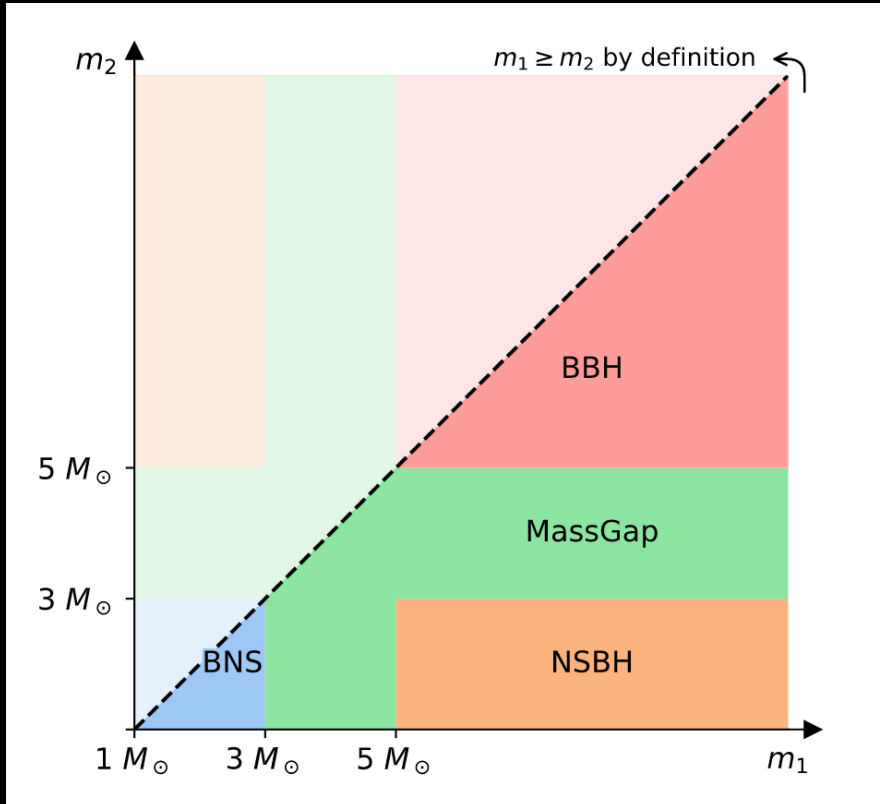
For CBC candidates:

- **3-D skymap** (Singer et al. 2016, ApJL 829, L15), with direction dependent luminosity distance
- **Luminosity distance** marginalized over whole sky
(mean+/-standard deviation)



ALERT CONTENTS

- For CBC candidates, **CLASSIFICATION** and **PROPERTIES**



Credit: User Guide

Categories in terms of
component masses



CLASSIFICATION:

→ P_{astro} probability that the signal is astrophysical

This probability evaluates whether the source belongs to one of five categories: BNS, mass gap, NSBH, BBH, Terrestrial

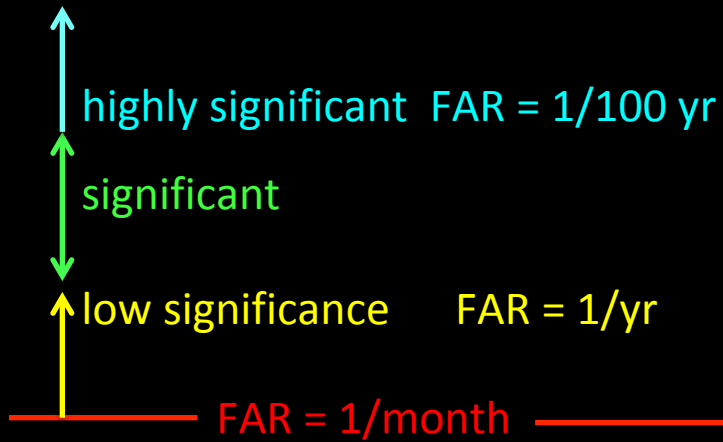
Based on our knowledge of trigger distribution, assumptions about signal distribution (such as that sources are uniformly distributed in volume), and knowledge and assumptions about merger rate per unit volume for each class of sources
(See Kapadia et al 2019)

PROPERTIES

- **HasNS** → probability that the mass of one or more of the binary's two companion compact objects is consistent with a neutron star.
- **HasRemnant** → probability that a non-zero amount of neutron star material remained outside the final remnant compact object (a necessary but not sufficient condition to produce certain kinds of electromagnetic emission such as a short GR or a kilonova)

(Foucart 2012, 2018, PhRvD, Pannarale & Ohme, 2014, ApJ)

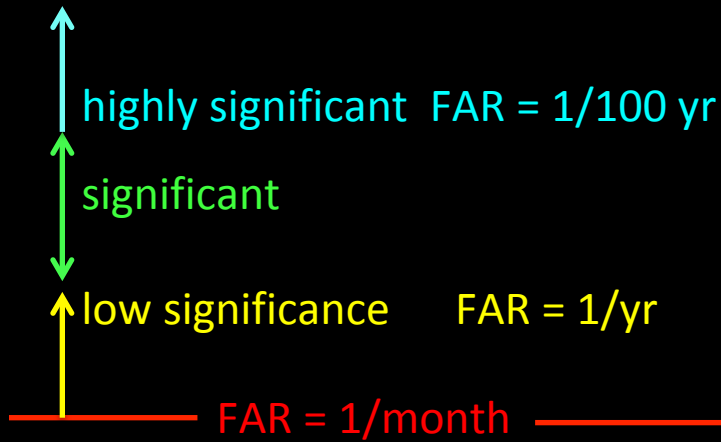




→ **FAR** = Rate of noise events louder than the candidate event



Candidates to be observed selected based on the observer's choice of FAR threshold



➔ **FAR** = Rate of noise events louder than the candidate event

↓

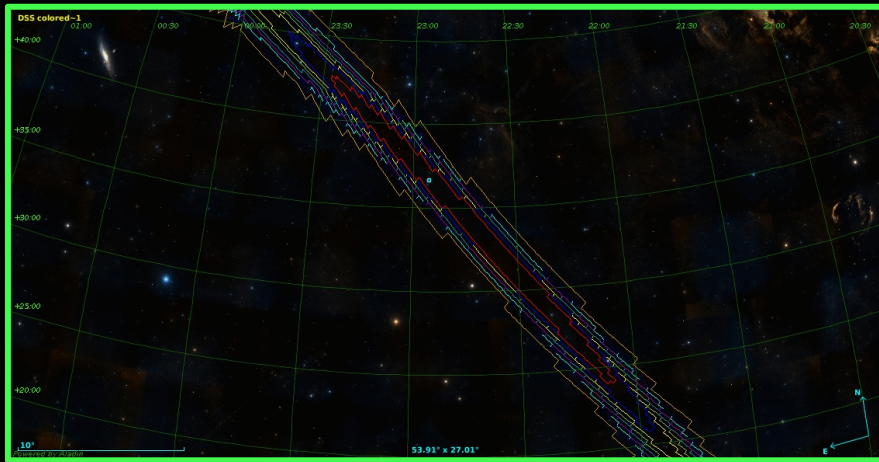
Candidates to be observed selected based on the observer's choice of FAR threshold

➔ Sky map + basic source classification

Credit: G. Greco, GWsky <https://github.com/ggreco77/GWsky>

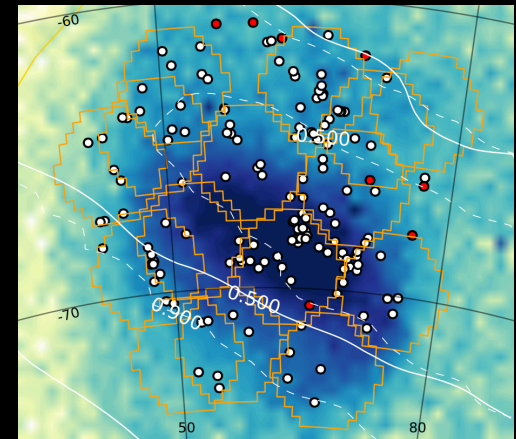


To decide the search type



Tiling the sky map to maximize the enclosed localization probability

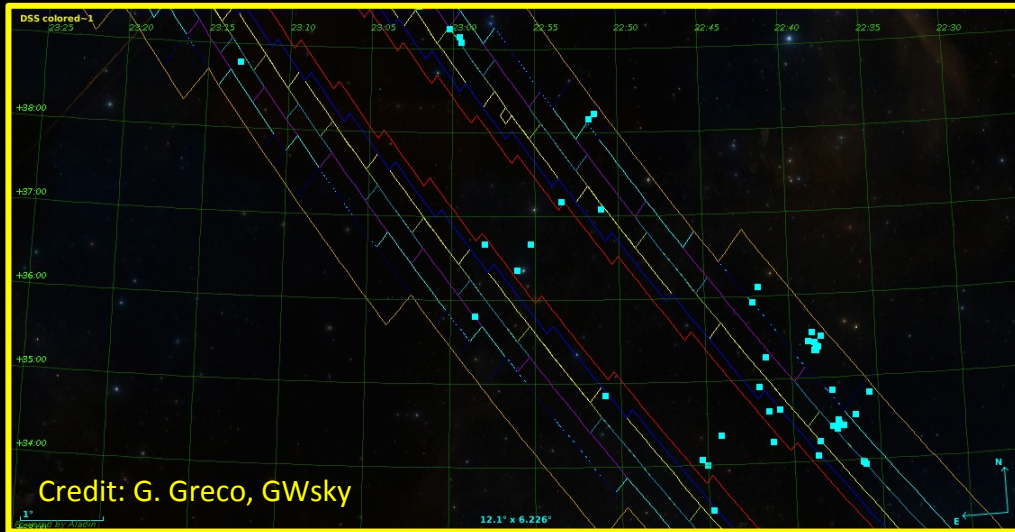
DES, Annis et al. 2016, ApJL



Burst → failed-SNe

Search for missing Supergiants in the LMC

→ Sky map + source classification + distance



Targeting ranked galaxies
(Small FoV instruments)

← Targeting ranked FoV pointings
(Instruments FoV > 1 deg²)

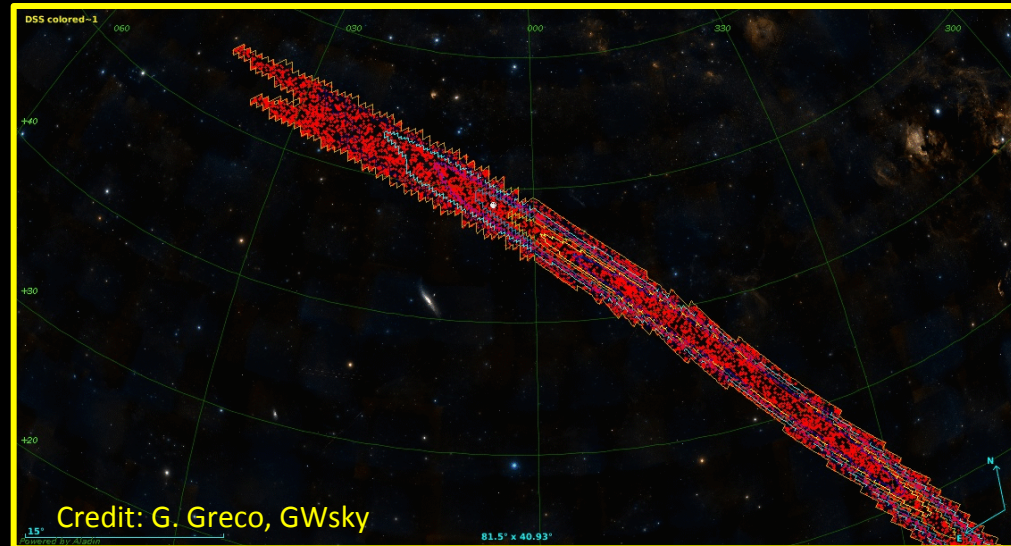
Sky map weighted by galaxy luminosity

For each FoV →

$$P = \sum \frac{L_i}{L_{tot}} P_{GW}$$

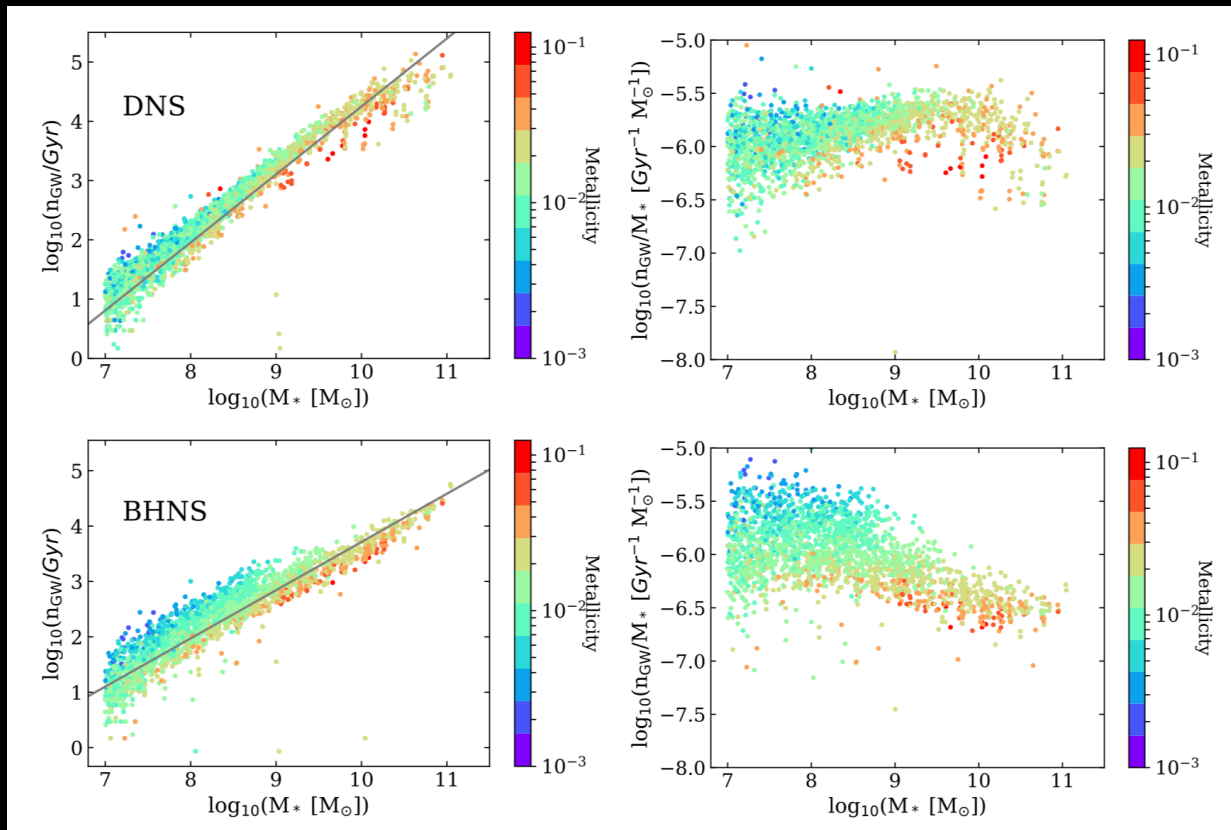
P_{GW} = probability that GW candidate lies within the FoV

See e.g Evans et al. 2016, MNRAS



HOW TO RANK THE GALAXIES?

From theoretical simulations: identify the most probable host by combining the results of population-synthesis models together with galaxy catalogs from galaxy cosmological simulation



Artale et al. 2019 MNRAS

- strong correlation between host galaxy mass and merger rate
- low mass galaxies have a more efficient merger rate per galaxy of NSBH systems

REAL OBSERVATIONS → GALAXY CATALOG + 3D SKY LOCALIZATION MAP

The overall probability of the merger occurring in a galaxy is given by

1) Localization probability

$$P_{loc} = P_{(RA,DEC)} P_{(D_l)}$$

$$\Rightarrow P = P_{loc} P_{lum}$$

2) Mass/SFR probability

$$P_{lum} = \frac{L_{K,B}}{L_{tot}} = \frac{L_{K,B}}{\sum L_{K,B}}$$

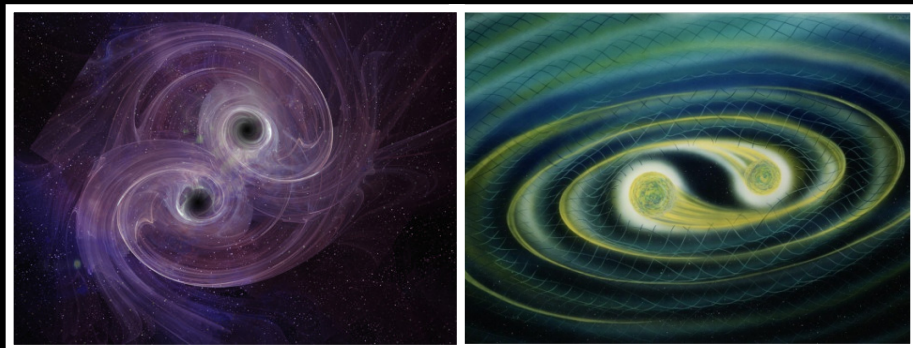
After the observations the same formalism is used to evaluate the probability covered by the galaxy targeted search (including catalog incompleteness)

See e.g. Gehrels et al. 2016,
Arcavi et al. 2017, Salmon 2019

*Electromagnetic emissions from
gravitational wave sources detectable by
ground-based detectors (10-1000 Hz)*

ASTROPHYSICAL SOURCES emitting transient GW signals detectable by LIGO and Virgo (10-1000 Hz)

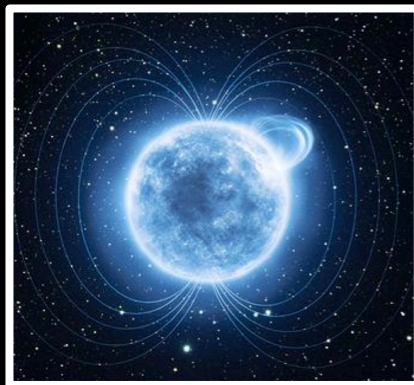
Coalescence of binary system of neutron stars and/or stellar-mass black-hole



Core-collapse of massive stars



Isolated NSs instabilities



EM emissions

NS-NS and NS-BH mergers

Short Gamma Ray Burst (sGRB)

Ultra-relativistic outflow

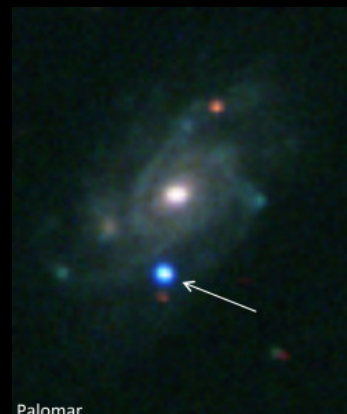
Beamed emission

Sub-relativistic dynamical ejecta

Isotropic emission
kilonova

disk wind outflow
spin-down luminosity

Core-collapse of massive stars



Palomar

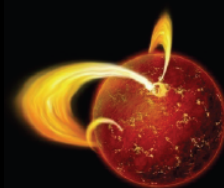
SBO X-ray/UV

Optical

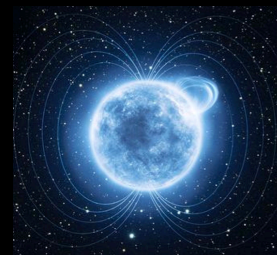
Radio

+ Long GRB

Isolated NS instabilities

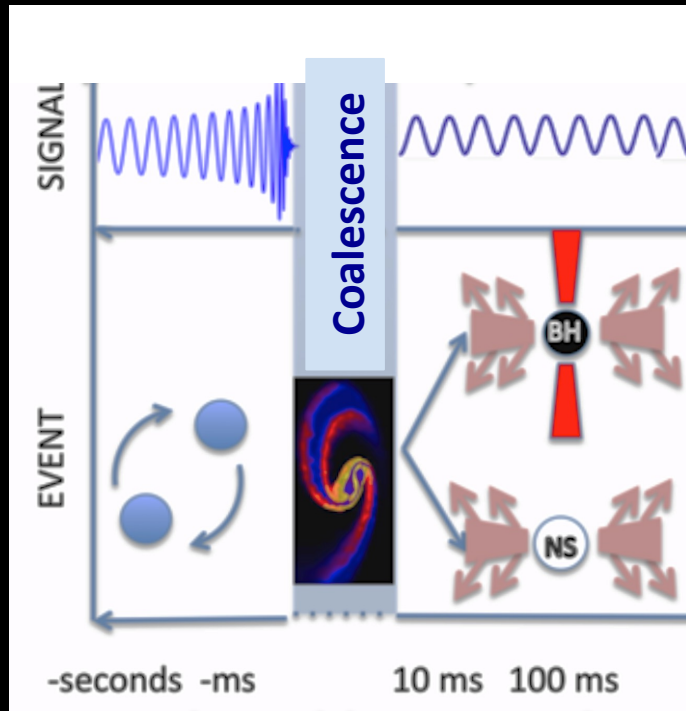


Soft Gamma Ray Repeaters and Anomalous X-ray Pulsars



Radio/gamma-ray Pulsar glitches

NS-NS and NS-BH inspiral and merger



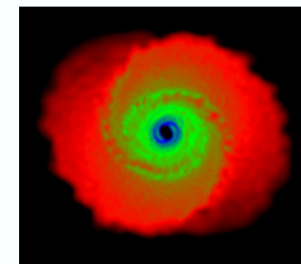
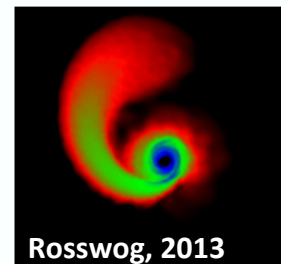
The merger gives rise to:

- **dynamically ejected unbound mass**
- **ejected mass gravitationally bound** to the central remnant either falls back or circularizes into an **accretion disk**

NS-NS binary → unbound mass of 10^{-4} - 10^{-2} M_{\odot} ejected at 0.1 - $0.3c$, which depends on **total mass, mass ratio, EOS NS and binary eccentricity**

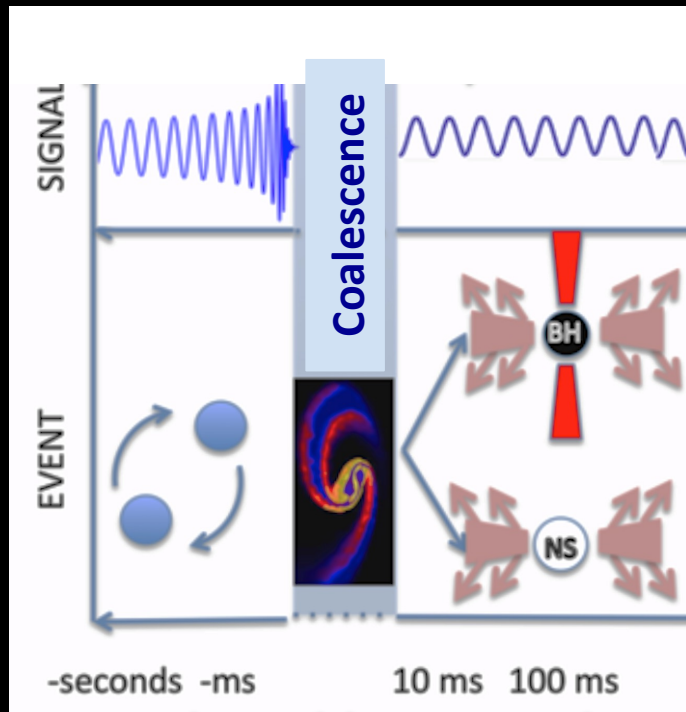
Dynamical Phase

Accretion phase



Fernandez & Metzger 2016, ARNPS, 66

NS-NS and NS-BH inspiral and merger



The merger gives rise to:

- **dynamically ejected unbound mass**
- **ejected mass gravitationally bound** to the central remnant either falls back or circularizes into an **accretion disk**

NS-BH binary → unbound mass up to 0.1 Mo depends on **ratio of the tidal disruption radius to the innermost stable circular orbit**

If < 1 → NS swallowed by the BH no mass ejection

If > 1 NS → tidally disrupted, long spiral arms

which depends on **the mass ratio, the BH spin and the NS compactness**

Dynamical Phase

Accretion phase

NSBH COALESCENCE



- Before the merger, the BH is described by its mass M_{BH} and spin χ_{BH} which determine the radius of the ISCO, R_{ISCO}
- Once the NS approaches the BH, the tidal forces increase. The objects' internal structure become important as the orbital separation approaches the size of the bodies

• IF

$$d_{\text{tidal}} < R_{\text{ISCO}}$$



NS direct plunge into BH

→ little or no mass left outside of the BH

→ **NO EM COUNTERPART**

• IF

$$d_{\text{tidal}} > R_{\text{ISCO}}$$



NS effectively disrupted

→ BH remnant surrounded by baryon matter

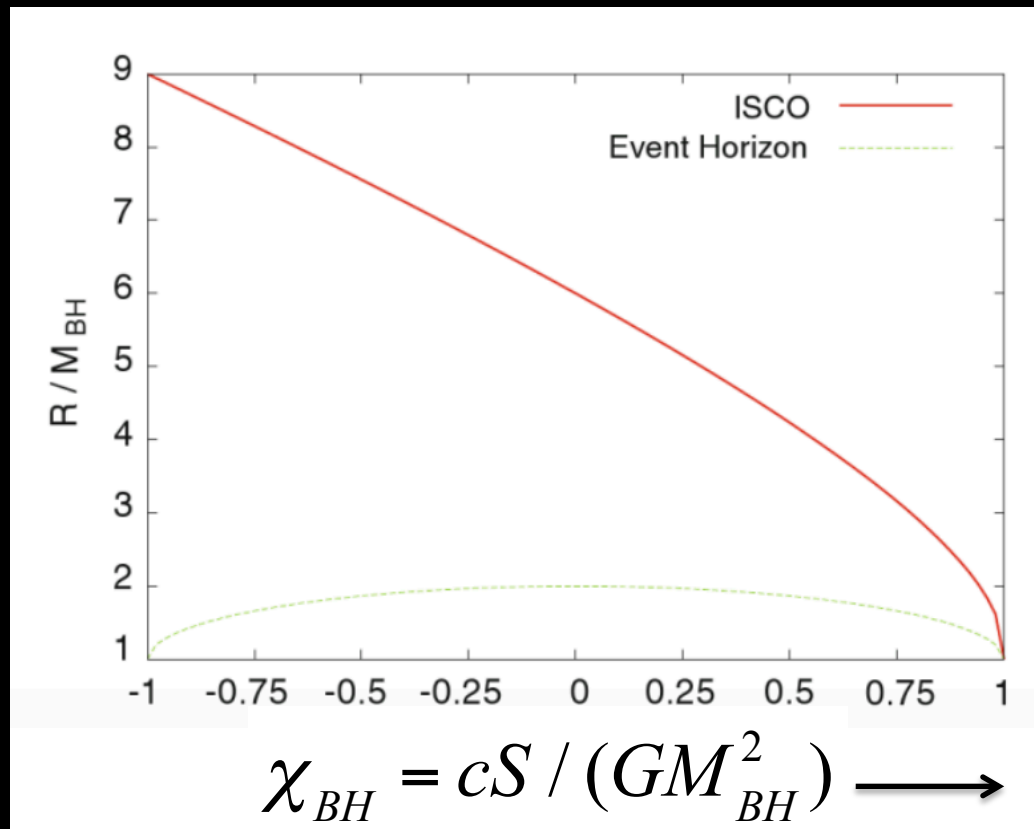
→ **EM COUNTERPARTS**

ISCO = innermost stable circular orbit of the BH, inside which no material have a stable circular orbit around the BH

For a non rotating Schwarzschild BH

$$R_{ISCO} = 6GM_{BH} / c^2 = 3R_S$$

For a rotating BH the equatorial ISCO also depends on the spin angular momentum



Kyutoku 2013

Non dimensional spin parameter

$$\chi = cS / (GM_{BH}^2)$$

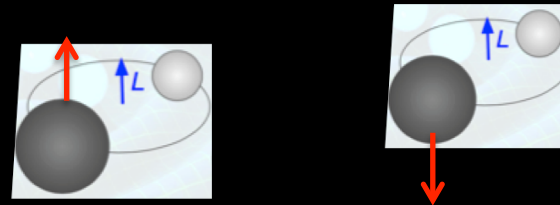
dimensionless spin parameter

- NS spin negligible \rightarrow typically assumed

$$\chi_{NS} \sim 0$$

NSs are expected to be born rapidly rotating but before NSBH coalescence (which requires long time from their birth) they have time to spin down by dipole-emission (the lack of matter accreting onto the NS prevents spin-up by recycling)

- Assumed non-precessing binaries \rightarrow BH spin vector aligned or anti-aligned with the orbital angular momentum



- Anti-aligned configurations \rightarrow larger ISCO, favour direct plunge of the NS into the BH \rightarrow no baryonic mass left outside the final BH to power an EM counterpart

Tidal disruption radius
occurs

d_{tidal}

radius at which tidal disruption

The tidal disruption occurs when the tidal force of the BH is stronger than the self-gravity of the NS

$$\frac{M_{BH}}{r^3} R_{NS} \geq C^2 \frac{M_{NS}}{R_{NS}^2}$$

C=non dimensional coefficient
r=orbital separation

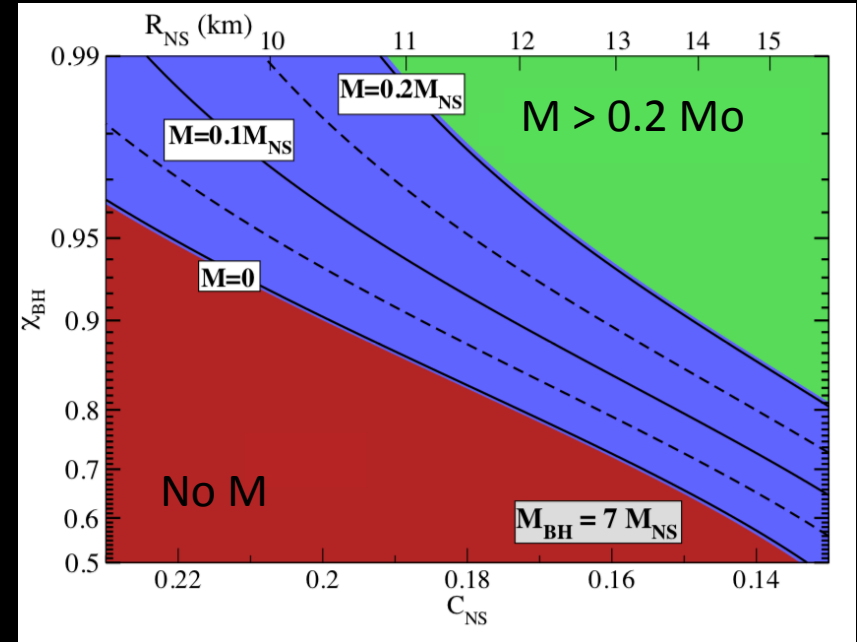
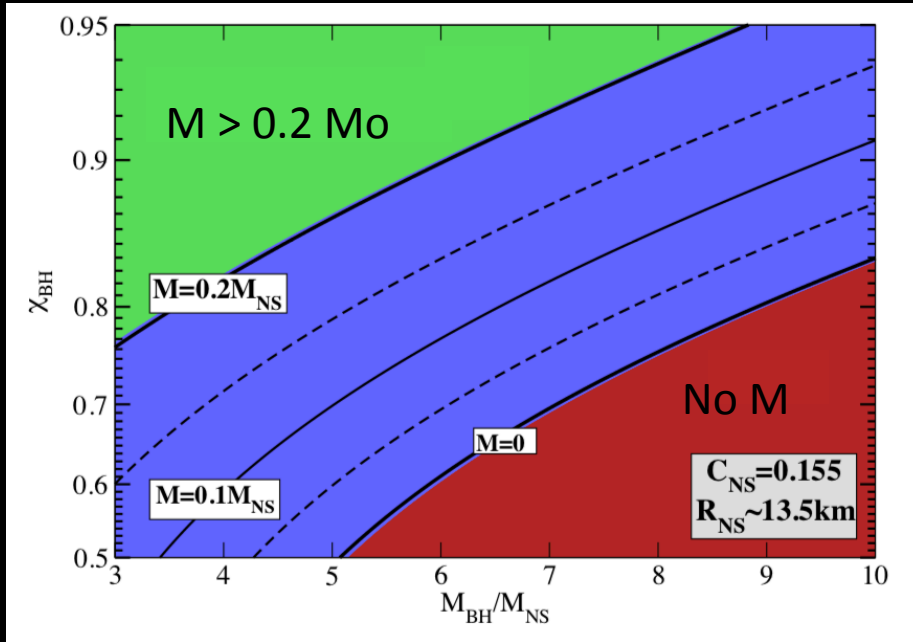
Newtonian theory

$$\frac{3M_{BH}}{d_{tidal}^3} R_{NS} \sim \frac{M_{NS}}{R_{NS}^2}$$



$$d_{tidal} \sim R_{NS} \left(\frac{3M_{BH}}{M_{NS}} \right)^{1/3}$$

Foucart 2012



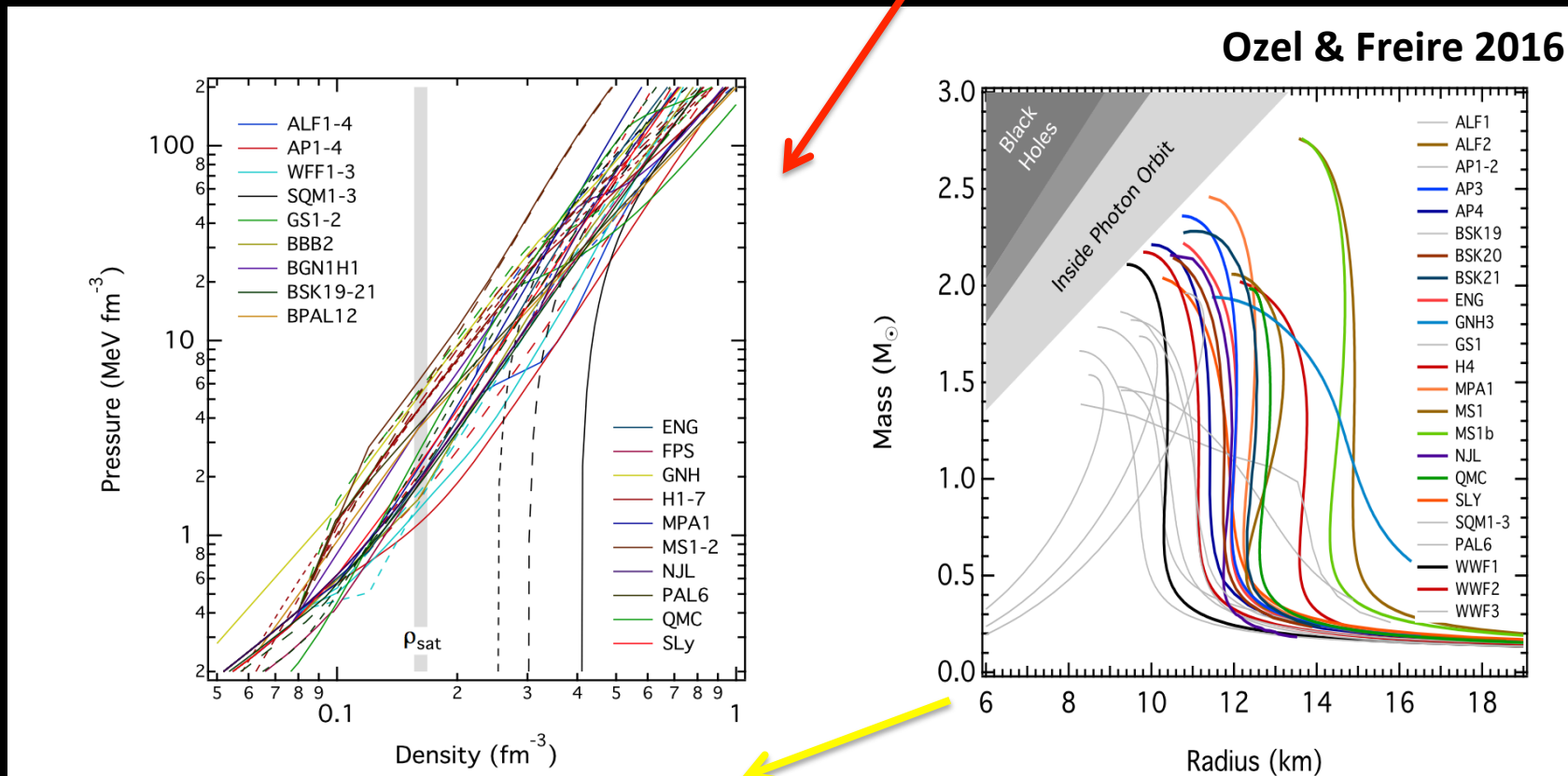
Large baryon mass left outside the merger remnant:

- Mass ratio BH/NS small \rightarrow small BH mass
- Large BH spin angular momentum
- Small NS compactness

In the degenerate interiors of neutron stars **EOS: $P \propto \rho^\alpha$**

Small $\alpha \rightarrow$ soft EOS (easier to compress)

High $\alpha \rightarrow$ stiff EOS (harder to compress)



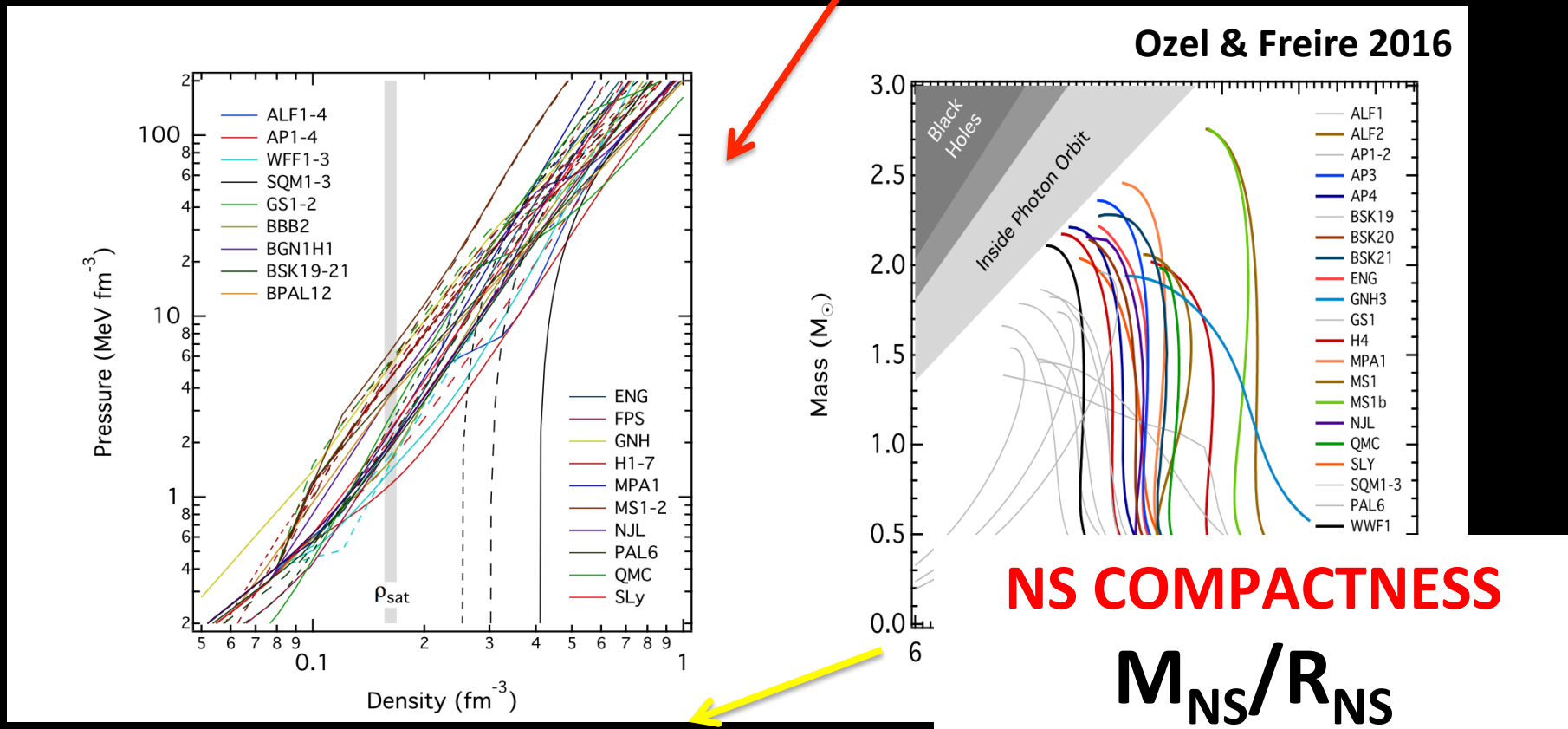
Mass-Radius relation is “unique” to the underlying EoS

- **Soft EoS:** low maximum M and smaller R for the same M (**more compact**)
- **Stiff EoS:** high maximum M and larger R for the same M (**less compact**)

In the degenerate interiors of neutron stars **EOS: $P \propto \rho^\alpha$**

Small $\alpha \rightarrow$ soft EOS (easier to compress)

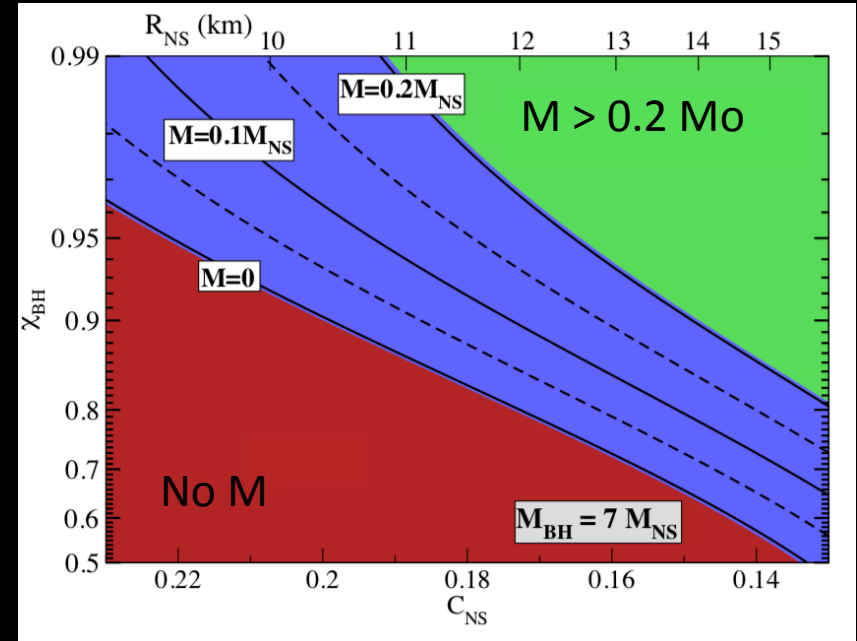
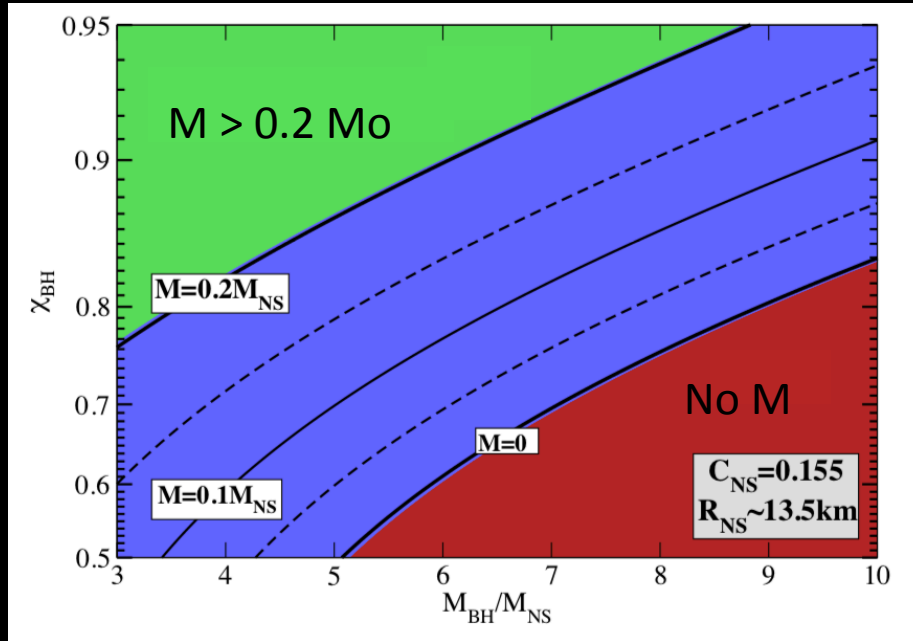
High $\alpha \rightarrow$ stiff EOS (harder to compress)



Mass-Radius relation is “unique” to the underlying EoS

- **Soft EoS:** low maximum M and smaller R for the same M (**more compact**)
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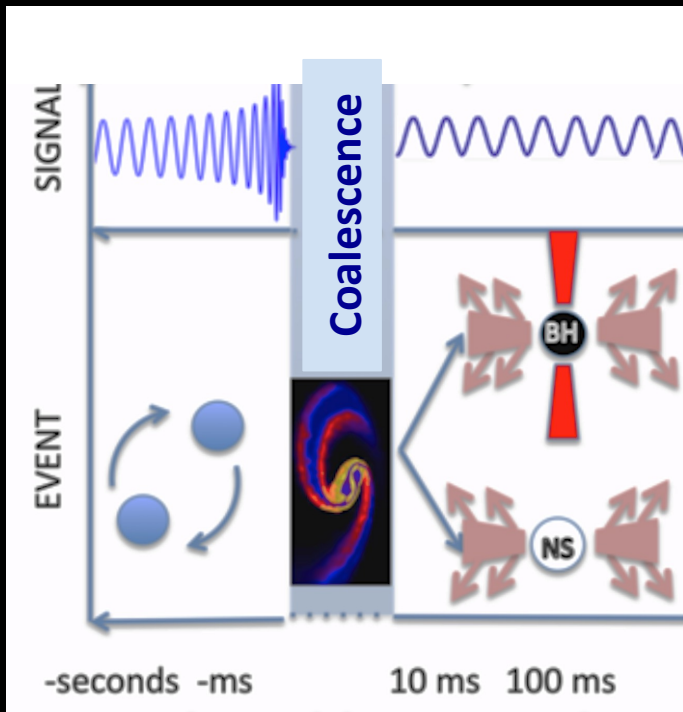
Foucart 2012



Large baryon mass left outside the merger remnant:

- Mass ratio BH/NS small \rightarrow small BH mass
- Large BH spin angular momentum
- Small NS compactness \rightarrow **same M large NS radius, stiff EOS**
(harder to compress, easier to be disrupted)

NS-NS and NS-BH inspiral and merger



- Ejected material gravitationally bound from the central remnant can fall back or circularizes into an accretion disk

Disk mass up to $\sim 0.3M_{\odot}$

Disk mass depends on the **mass ratio of the binary, the spins of the binary components, the EOS, and the total mass of the binary**

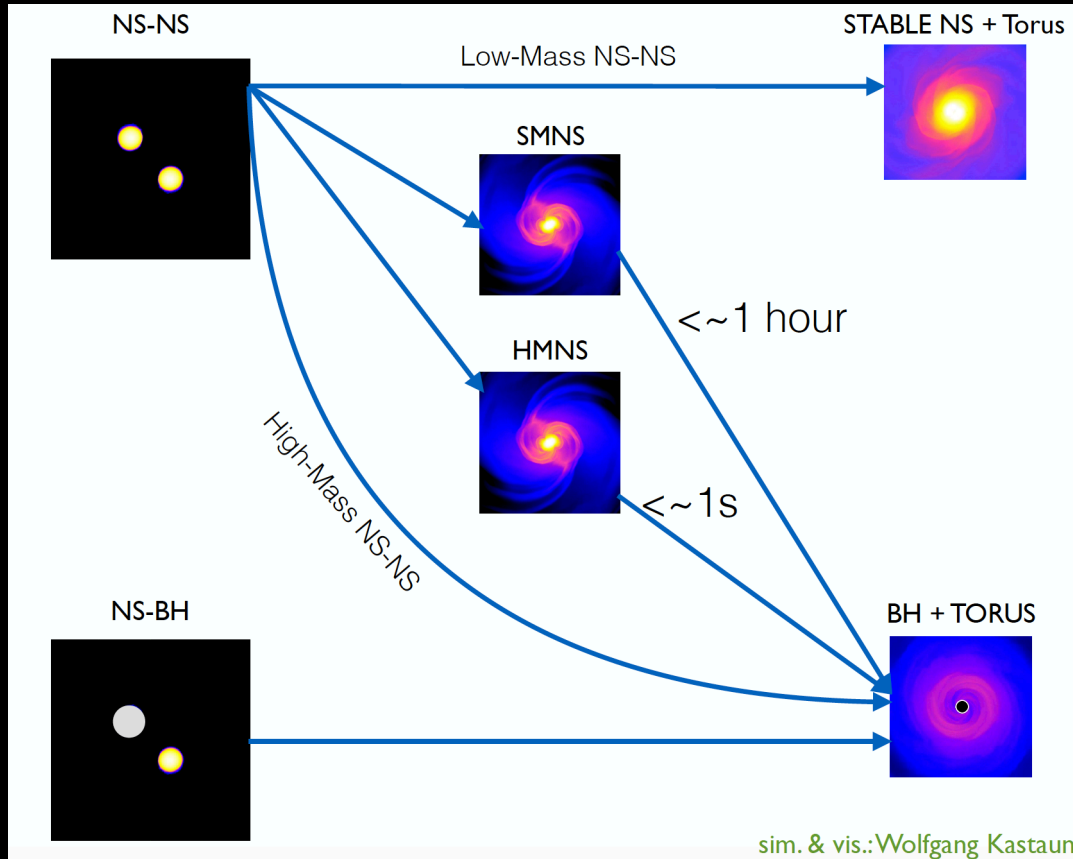
For NS-BH see e.g. Foucart 2012, PhRvD, 86;
Maselli & Ferrari, PhRvD, 89;
Pannarale & Ohme, ApJL, 791

Dynamical Phase

Accretion phase

Outflow mass and geometry influence the EM emission

Central remnant of NS-NS or NS-BH merger



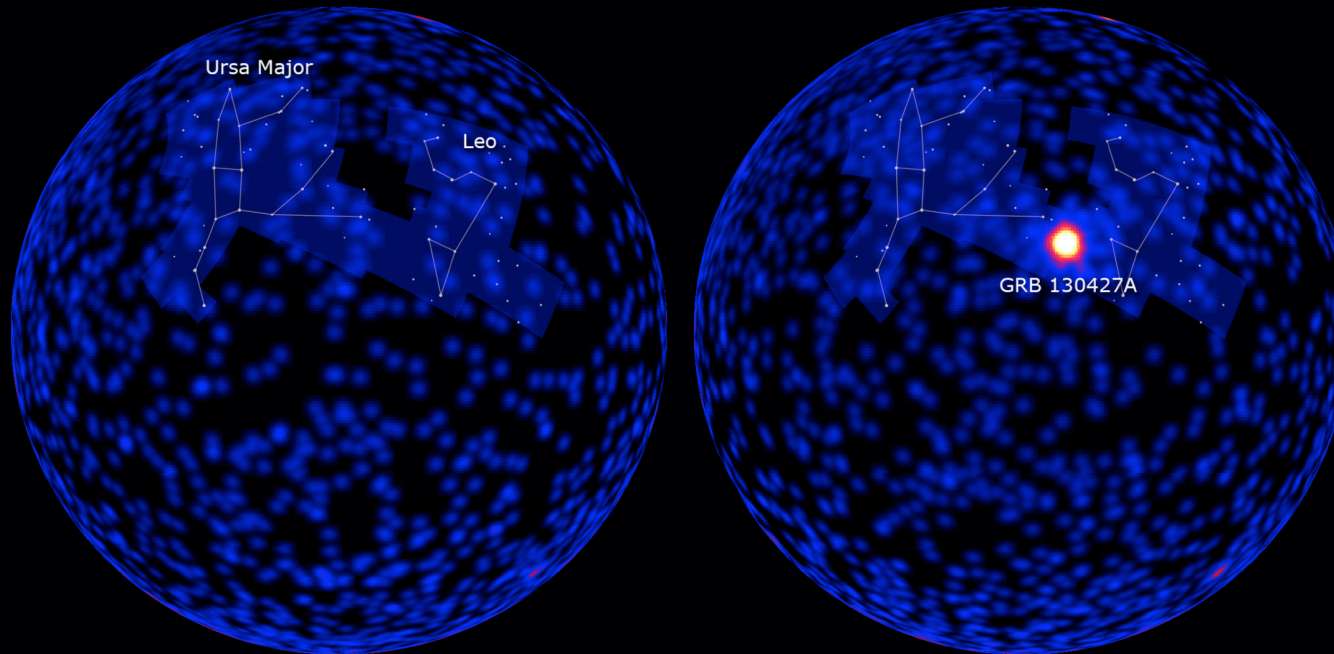
The central remnant influences GW and EM emission

What is central remnant?

- It depends on the total mass of the binary
- The mass threshold above which a BH forms directly depends on EOS

EM non-thermal emission

Gamma-Ray Bursts

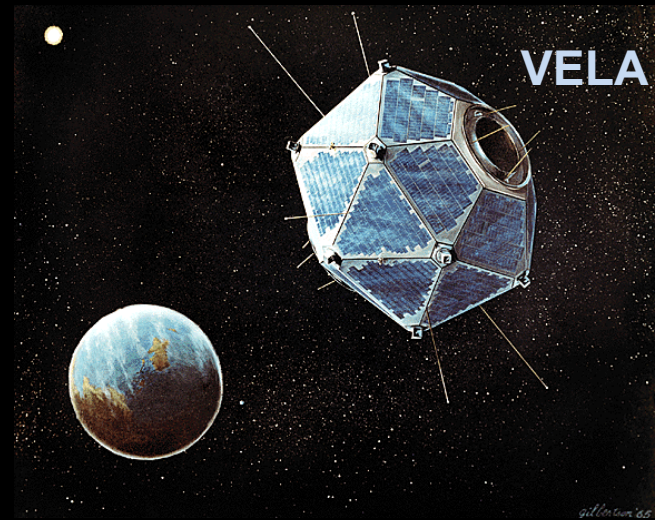
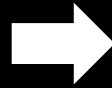


Before and after Fermi LAT observation of GRB 130427A

Brief, sudden, intense flashes of gamma ray radiation which release energy up to $\sim 10^{53}$ erg (isotropic-equivalent)

Duration: **from few ms to hundreds of s**
Observational band: **10 keV – 1 MeV**
Flux: **10^{-8} - 10^{-4} erg cm⁻² s⁻¹**

GRBs were discovered serendipitously in the late 1960s by U.S. military satellites looking out for Soviet nuclear testing



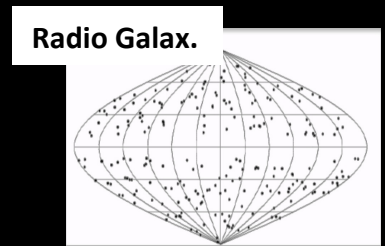
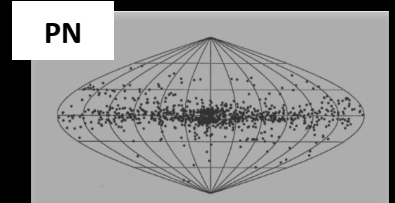
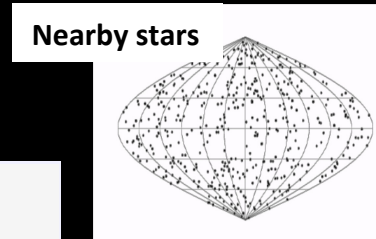
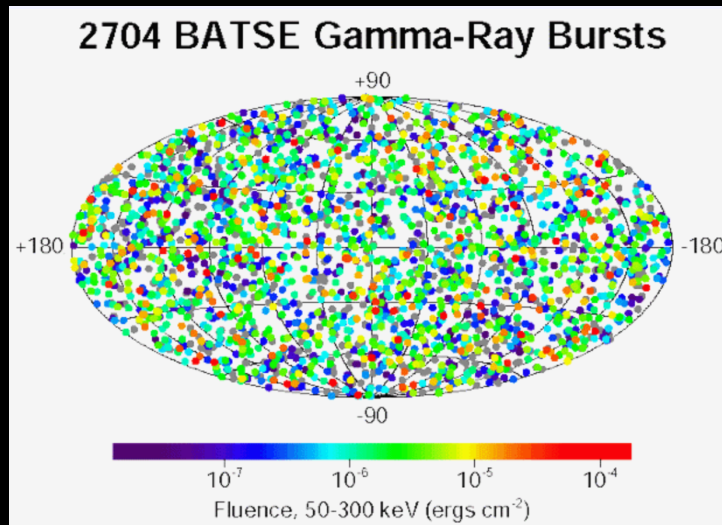
Galactic or cosmological?



BATSE
20 keV-MeV
(1991-2000)

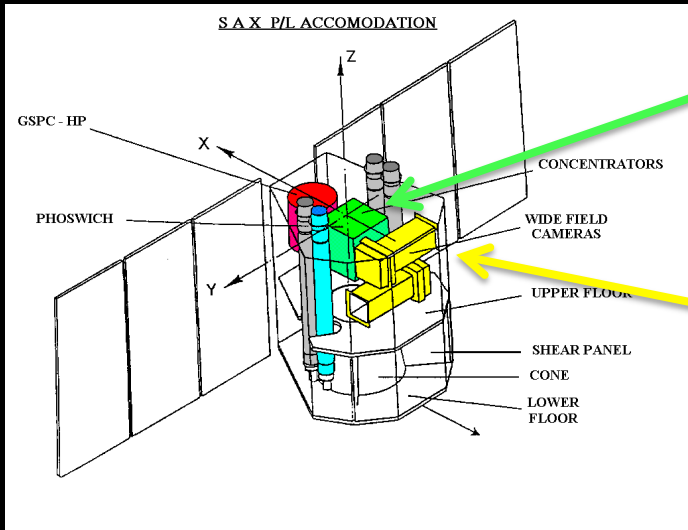


Isotropic angular distribution



BeppoSAX

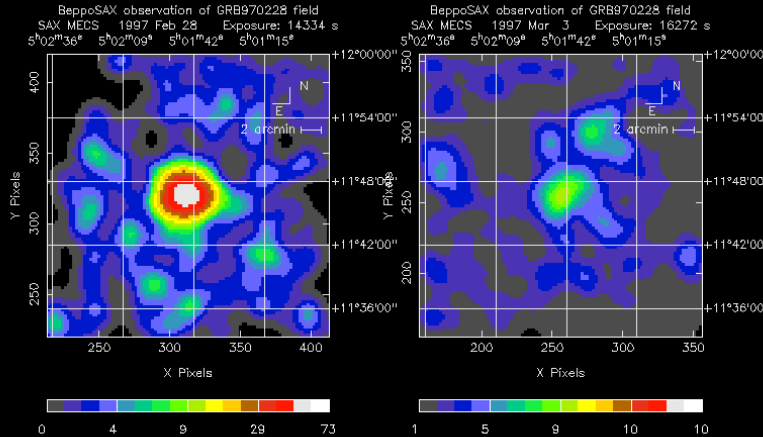
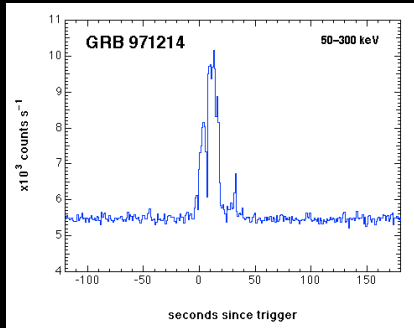
(1996-2002) *Italian-Dutch satellite for X-ray astronomy
resolved the origin of gamma-ray bursts*



*Scintillator for gamma-rays
60-600 keV, poor angular resolution*

*Wide Field Camera (WFC)
2-30 keV; 20x20 degree FoV
5 arcmin angular resolution*

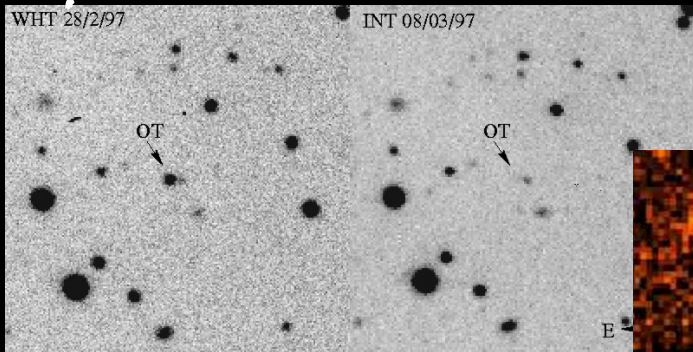
GRB 970228 in the FOV of the WFC



Well localized fading X-ray afterglow!

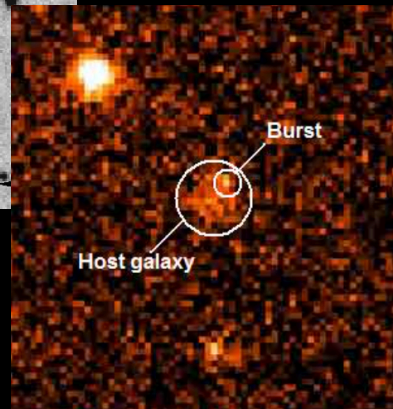
Costa et al., 1997

Optical afterglow/host galaxy



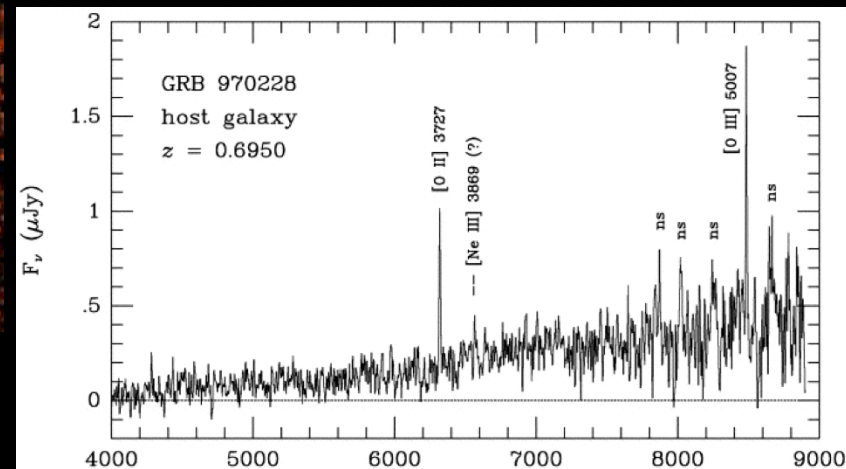
$z=0.695$, $D_L=3.6$ Gpc

Cosmological redshift



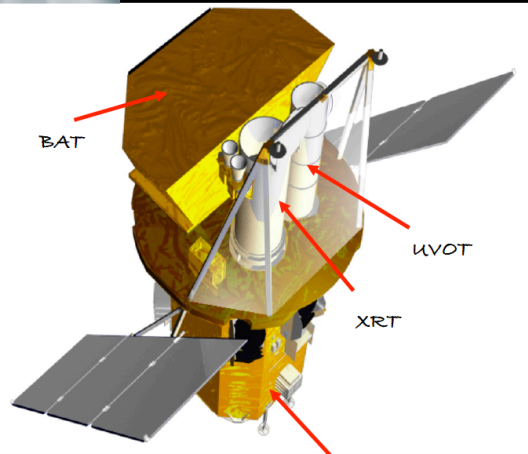
Groot, Galama, van Paradijs, et al IAUC 6584, March 12, 1997

van Paradijs et al., 1997



Swift: "everything in space"

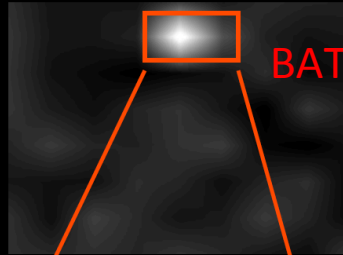
Since Nov. 2004



Spacecraft

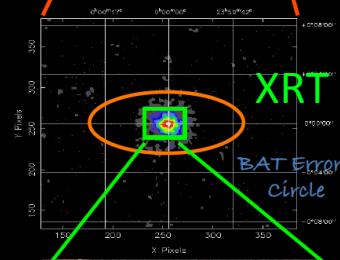
Satellite slews (1 min) and repoints its X ray (XRT) and UV telescopes to observe the error region of the GRB.

$T < 10 \text{ sec}$
 $\theta < 4'$
 $E > 15 \text{ keV}$



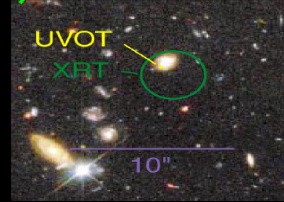
BAT

$T < 100 \text{ sec}$
 $\theta < 5''$
 $E < 10 \text{ keV}$



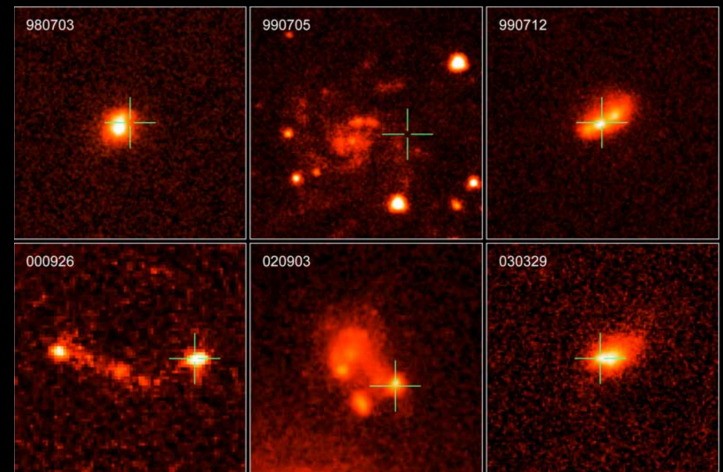
XRT

$T < 300 \text{ sec}$
Optical/NIR



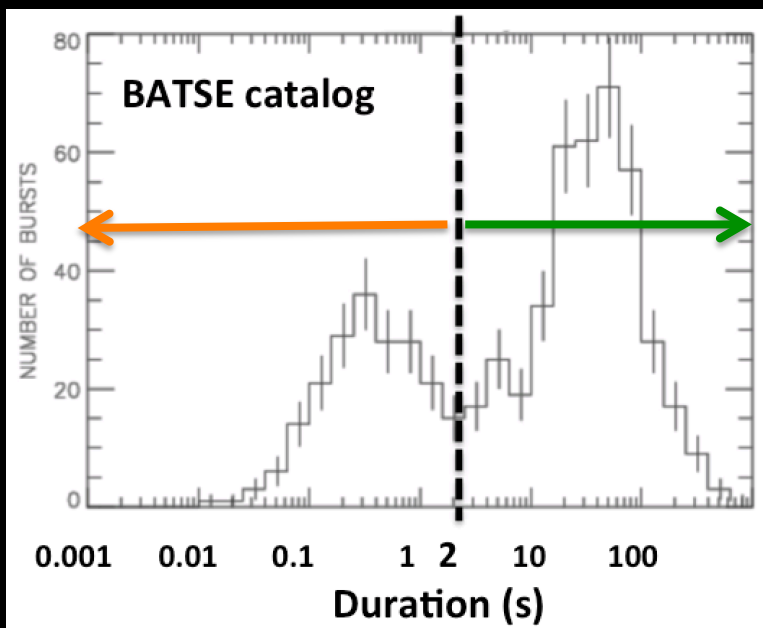
GRBs are extragalactic, cosmological, and occur in galaxies

GRBs host galaxies observed by HST



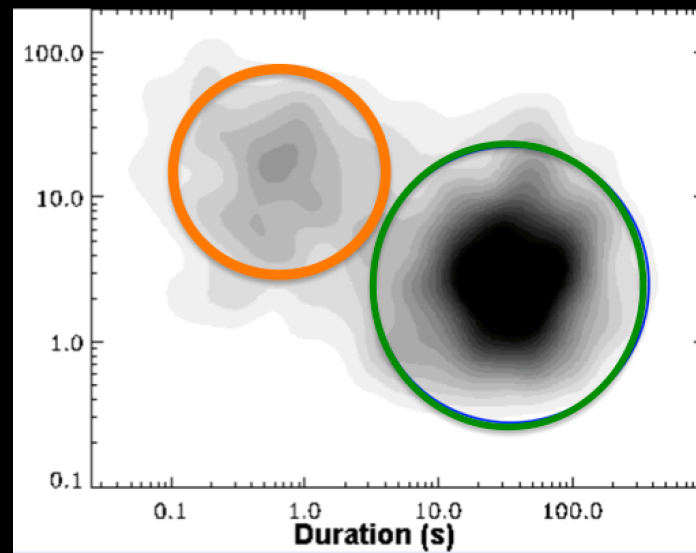
GRB progenitors

Bimodal duration distribution



Kouveliotou et al. 1993

Hardness ratio



Different Progenitor

Short Hard GRB

- lack of observed SN
- association with older stellar population
- larger distance from the host galaxy center ($\sim 5-10$ kpc)
- accretion timescale of disk in binary merger model is short ($t \sim 1$ s)

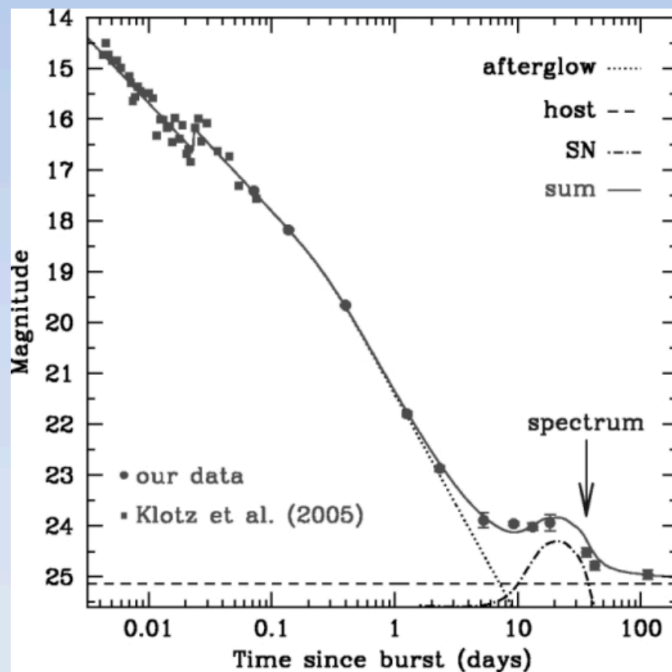
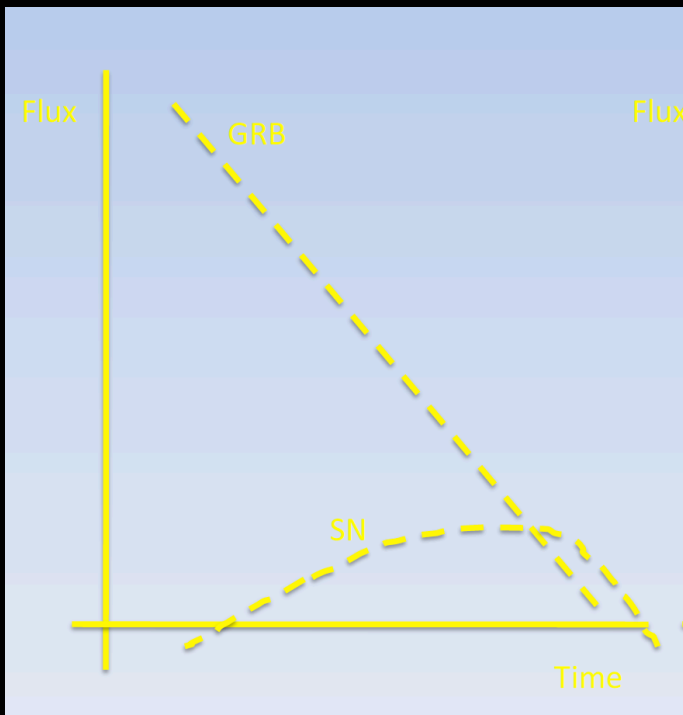
➔ **NS-NS NS-BH mergers**

Long Soft GRB

- observed Type Ic SN spectrum
- accretion disk is fed by fallback of SN material onto disk, timescale $t \sim 10-100$ s

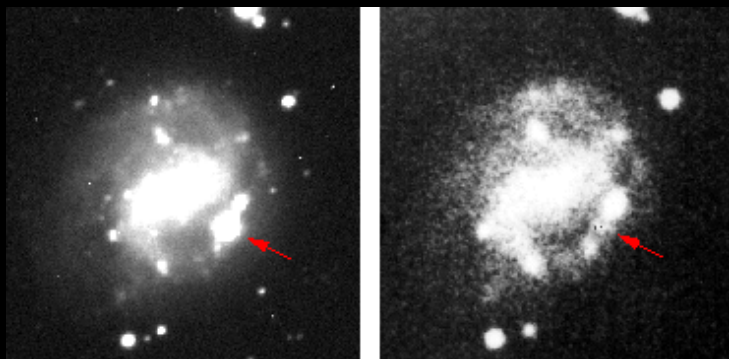
➔ **Core-collapse of massive stars**

Long GRB and Supernovae



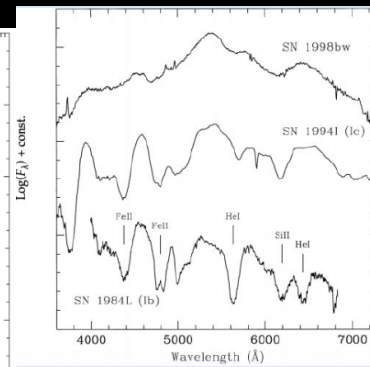
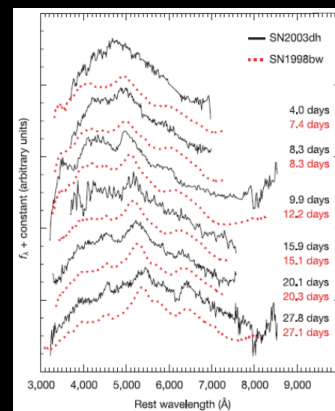
SN 1998bw/GRB 980425
Type Ic supernova

Galama et al. 1998; Stanek et al. 2003; Hjorth et al. 2003; Della Valle et al. 2003; Malesani et al. 2004; Soderberg et al. 2005; Pian et al. 2006; Campana et al. 2006; Della Valle et al. 2006, Bufano et al. 2012, Melandri et al. 2012, Schulze et al. 2014, Melnadri et al. 2014 and others...



Iwamoto et al 1998; Woosley et al. 1999

Spectra



GRB EMISSION

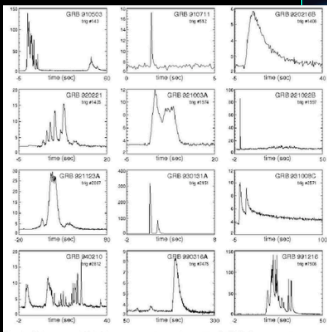
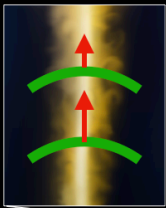


Image Credit: ESA/Hubble, M. Kornmesse

Prompt emission phase:

Energy range: keV-MeV

Variability time-scales: ms-s

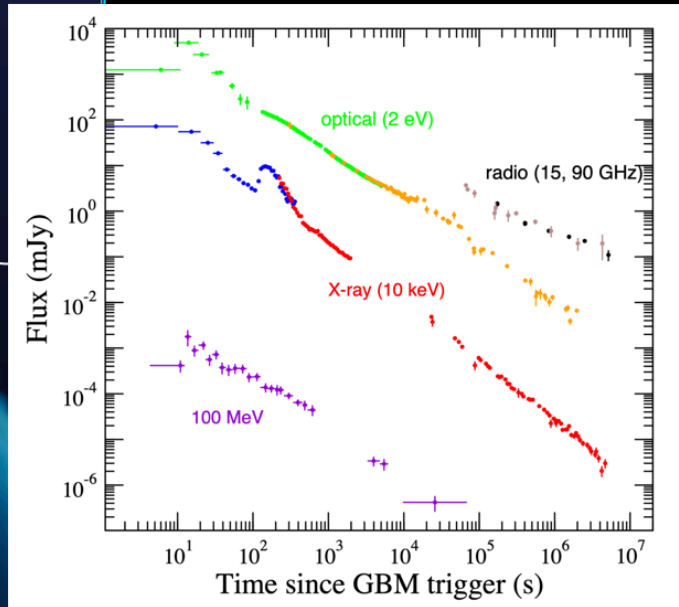


Shemi & Piran (1990)

Rees & Meszaros (1994)

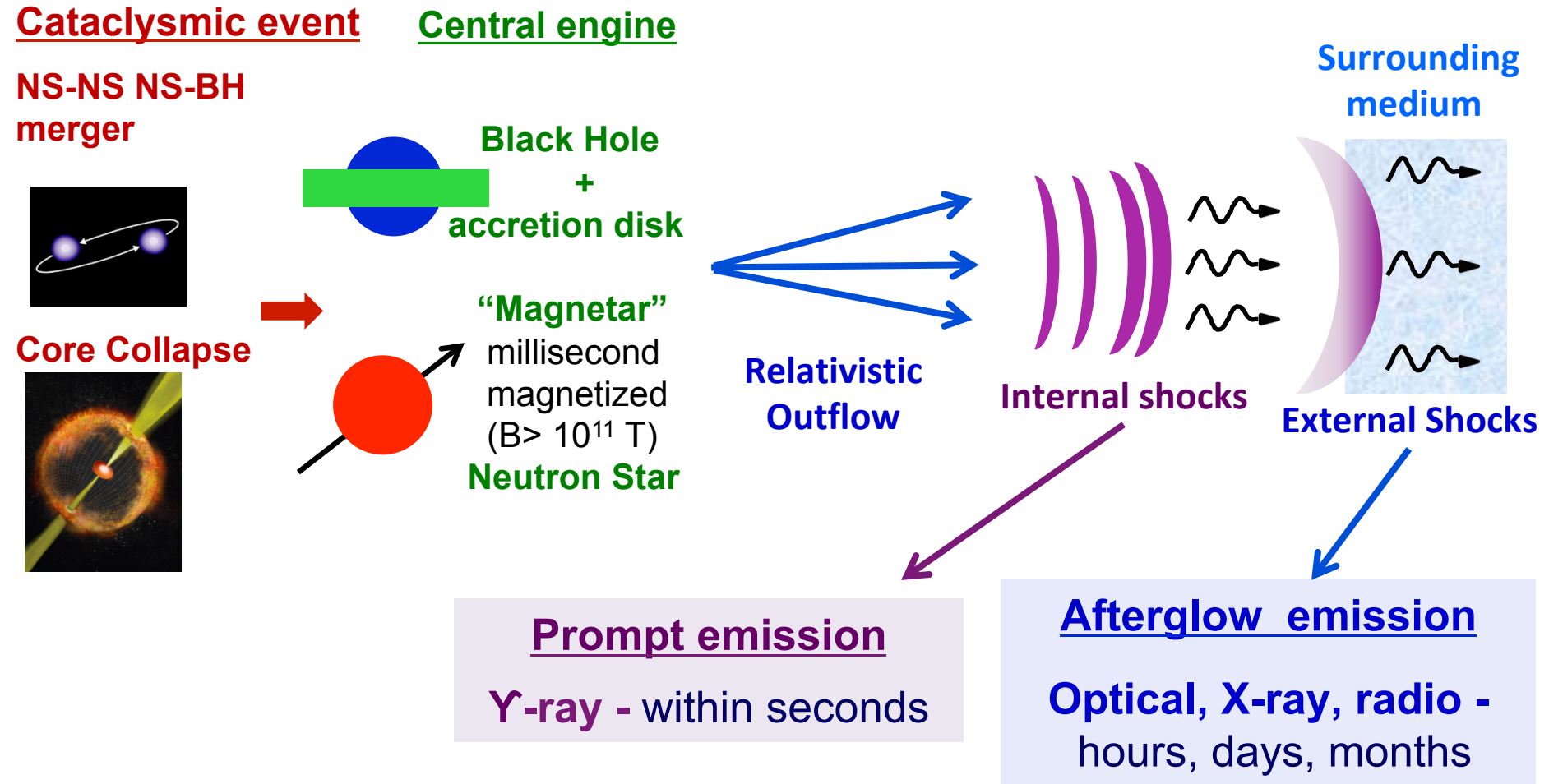
Credit: Ronchini

Afterglow phase



From Panaitescu et al (2013)

GRBs emission - Fireball Model

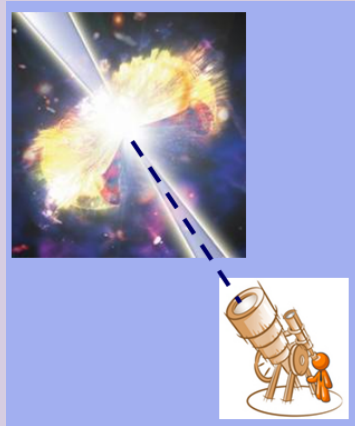


Kinetic energy of the relativistic jet converted into radiation

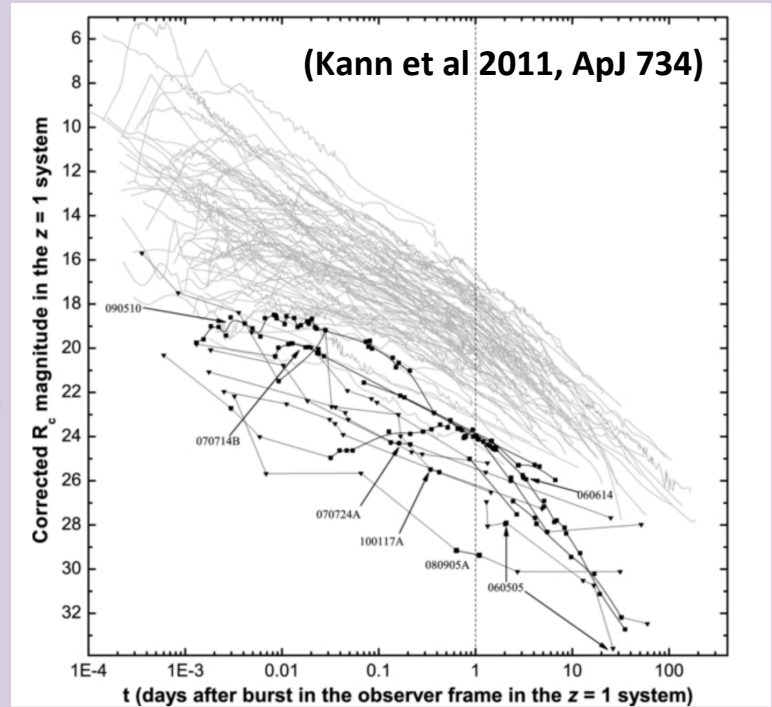
$$M_{\text{jet}} = 10^{-7} - 10^{-5} M_{\odot}, \Gamma \geq 100, E = 10^{48} - 10^{51} \text{ erg}$$

Optical afterglows of on-axis GRBs

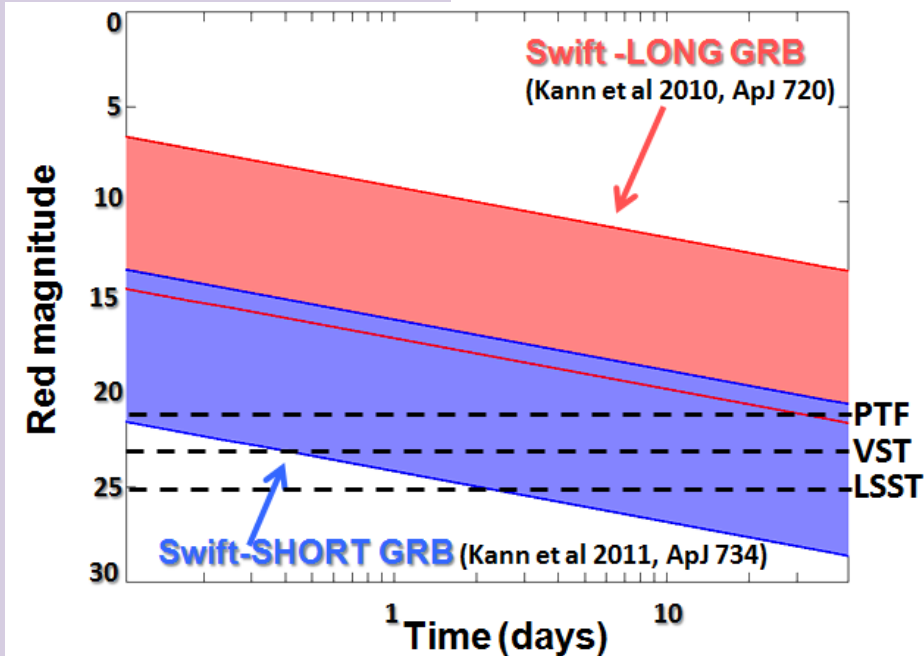
On-axis GRBs



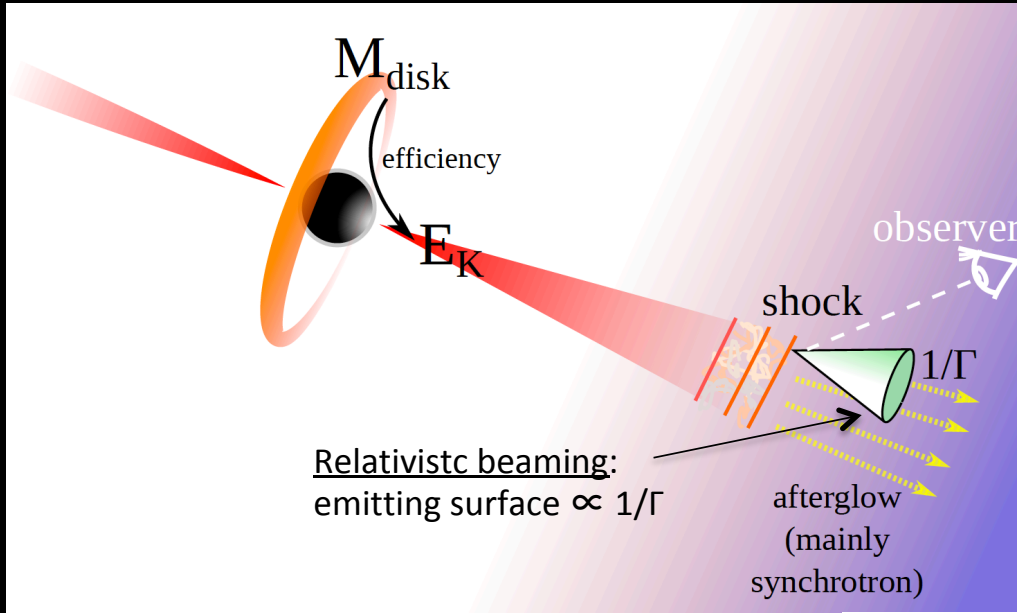
Observed GRB optical afterglows



Source at 200 Mpc



On average the optical afterglow decays as a power law $\text{time}^{-\alpha}$ with α in the range 1 to 1.5



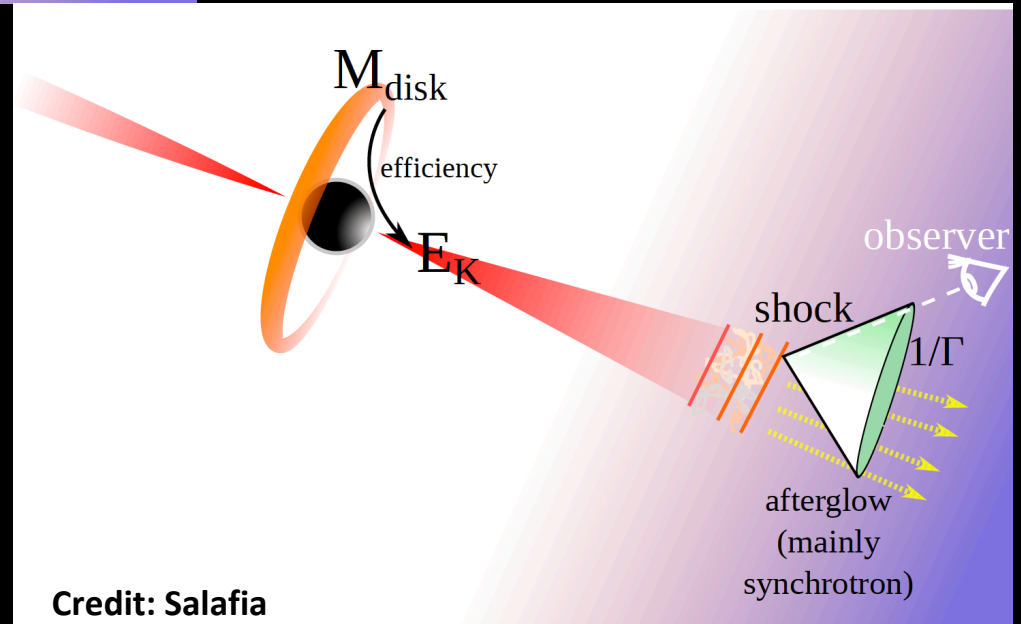
Relativistic beaming angle

$$\theta_b \propto 1/\Gamma$$

(within which ejecta radiates) increases due to the ejecta deceleration

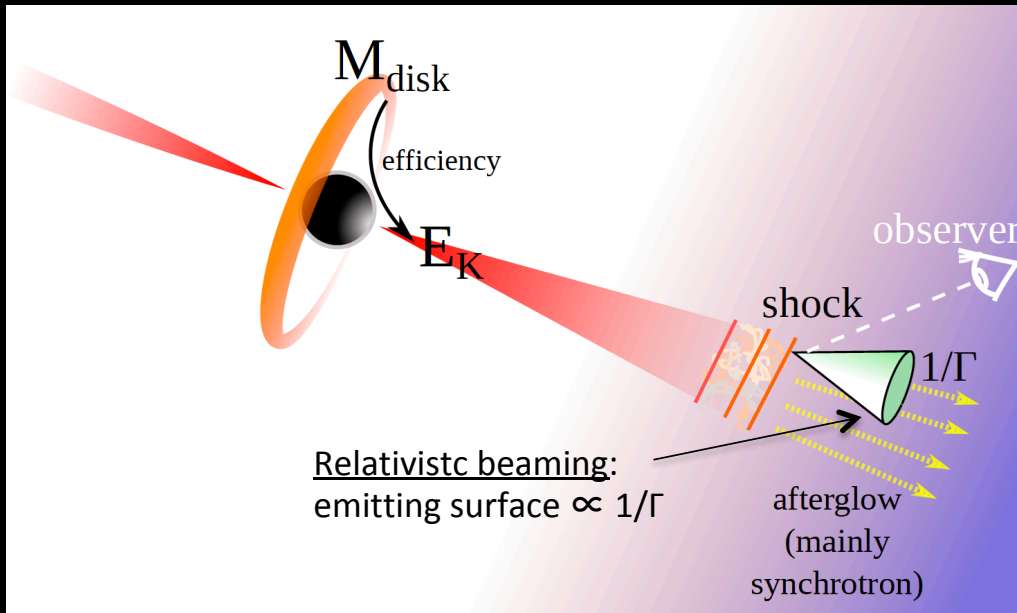
$$\Gamma = \frac{1}{\sqrt{1-\beta^2}} \quad \beta = \frac{v}{c}$$

↑
 Early EM emission detectable only by on-axis observers



Credit: Salafia

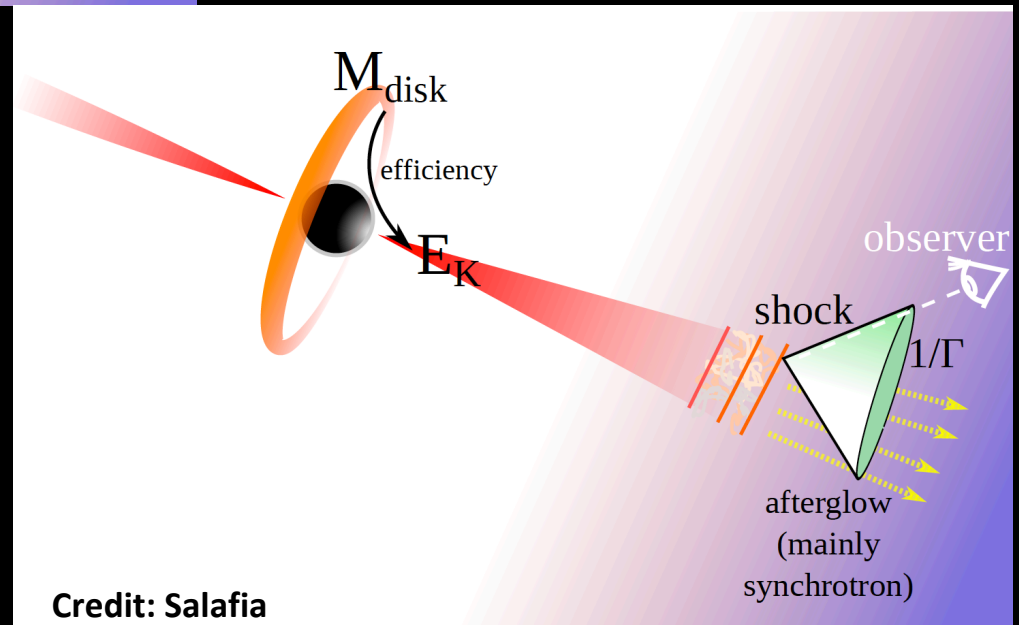


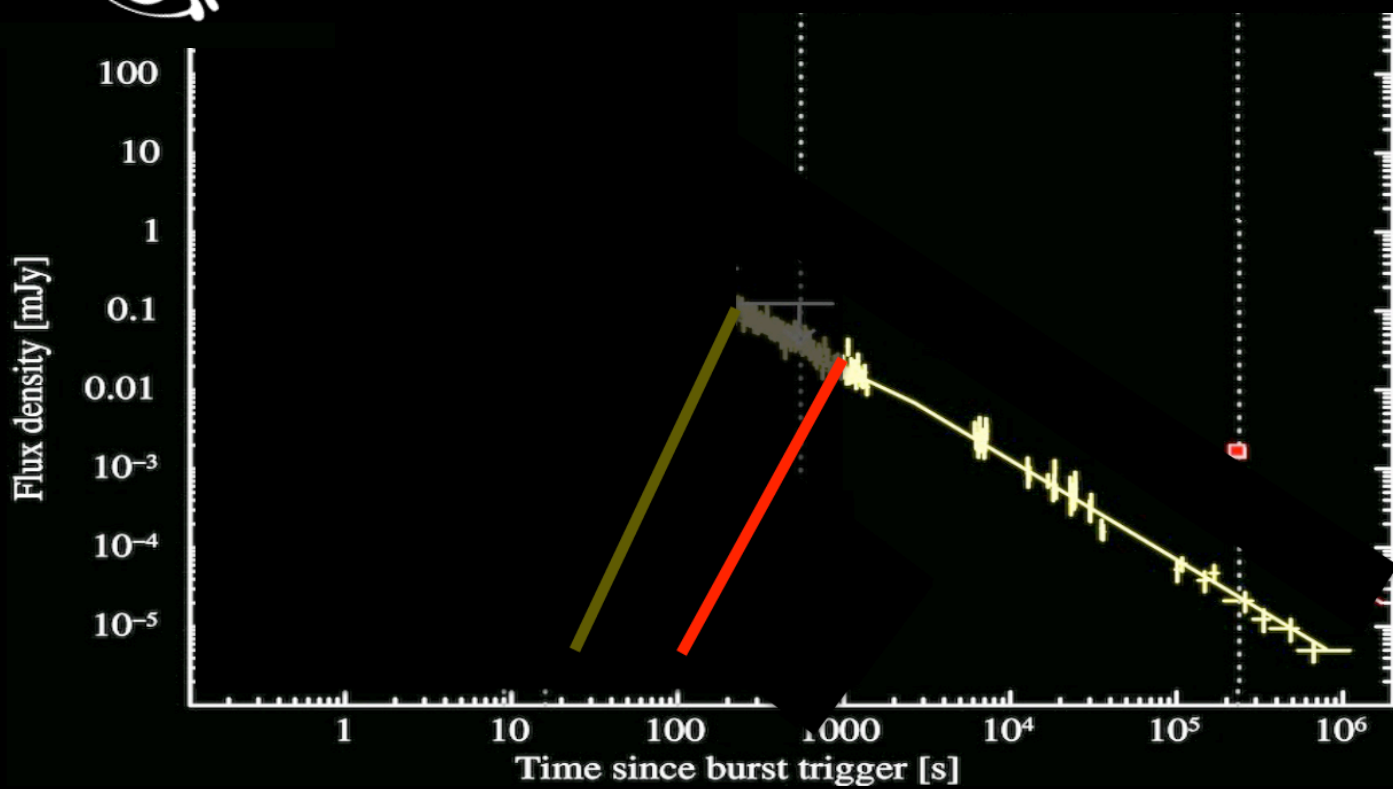
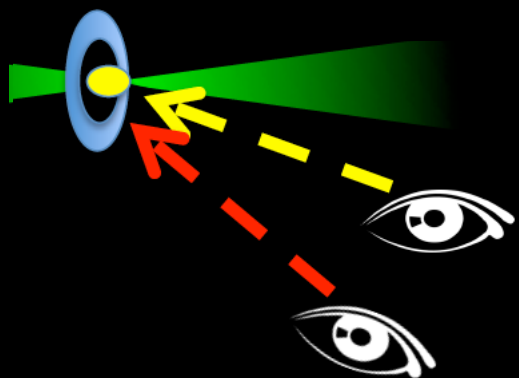


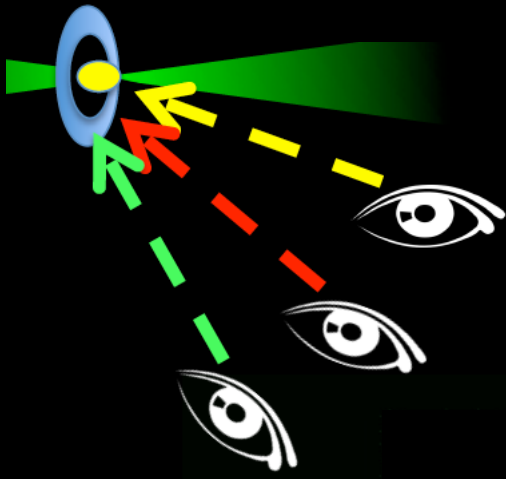
EM emission
detectable also by
off-axis observers



Early EM emission
detectable only by
on-axis observers



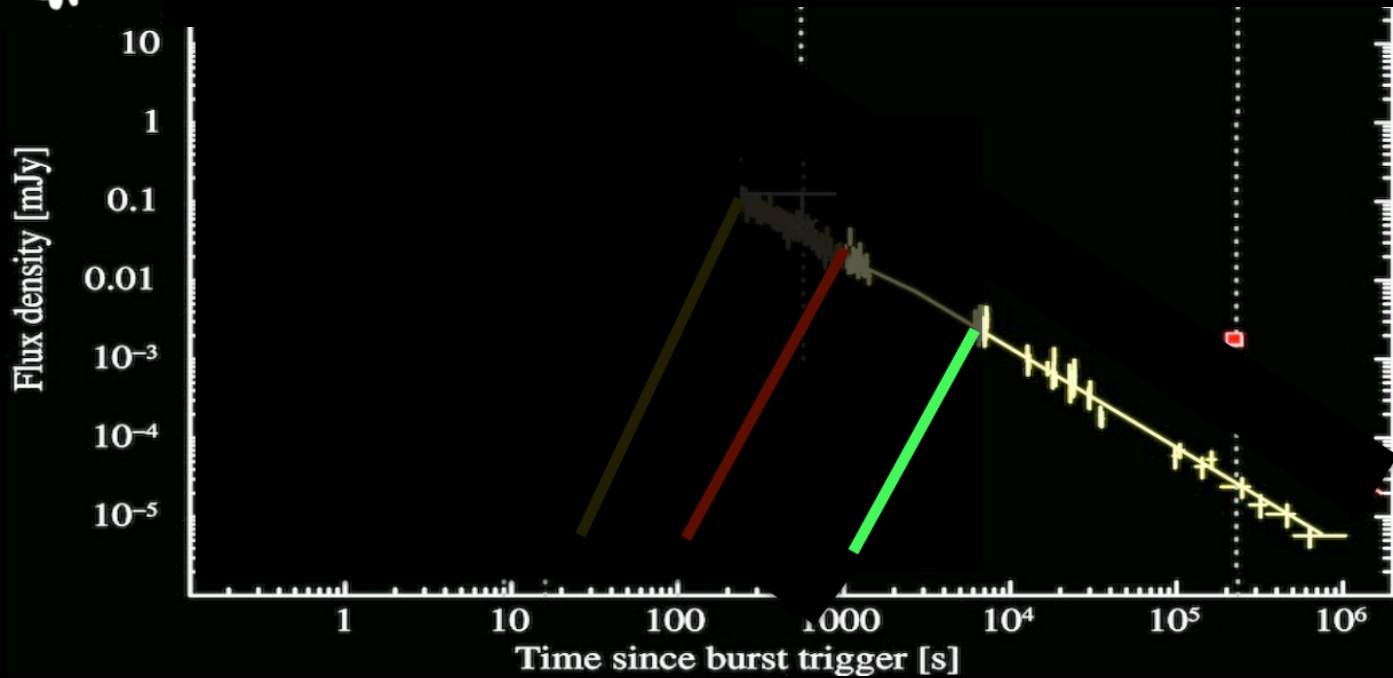




More off—axis:

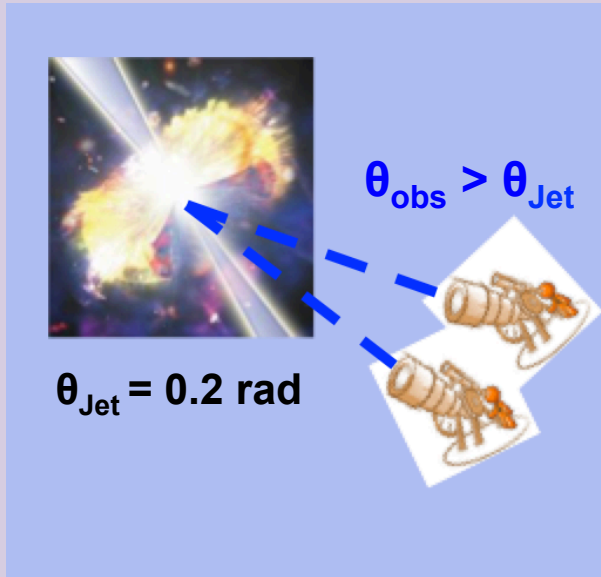
(a) the emission peaks later

(b) the flux at peak is fainter



Optical afterglows of Off-axis GRBs

Off-axis GRB



LONG bright GRB

$E_{\text{jet}} = 2e51 \text{ erg}$, $n = 1 \text{ cm}^{-3}$

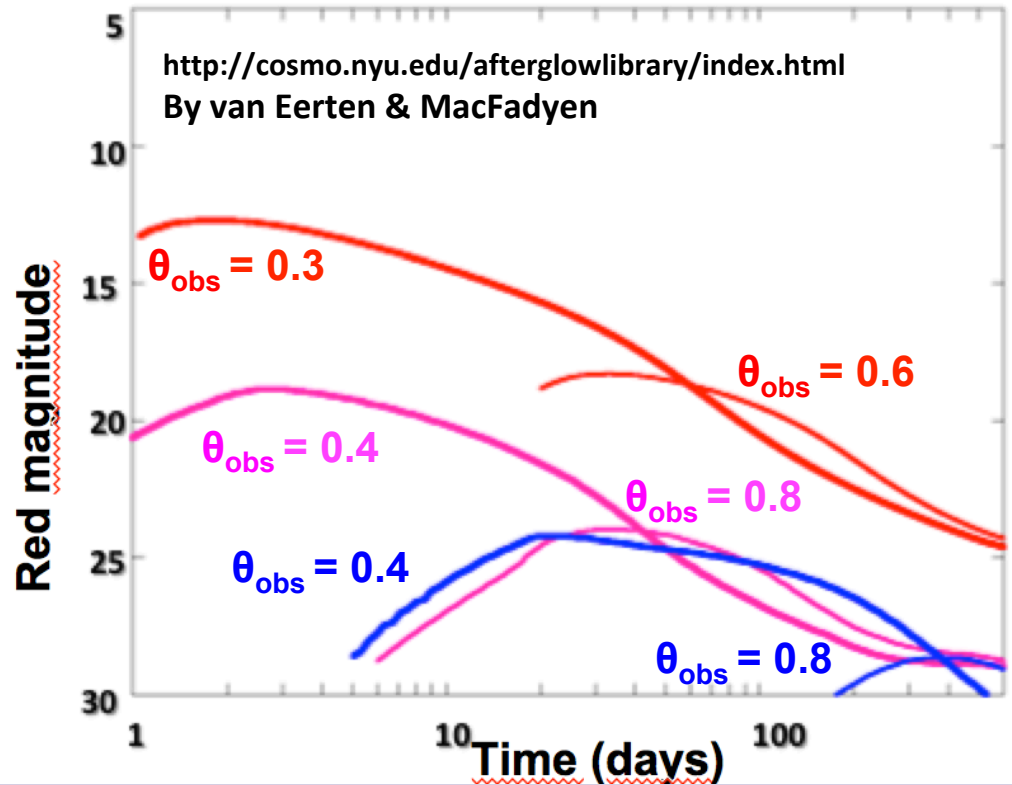
LONG faint/ SHORT bright GRB

$E_{\text{jet}} = 1e50 \text{ erg}$, $n = 1 \text{ cm}^{-3}$

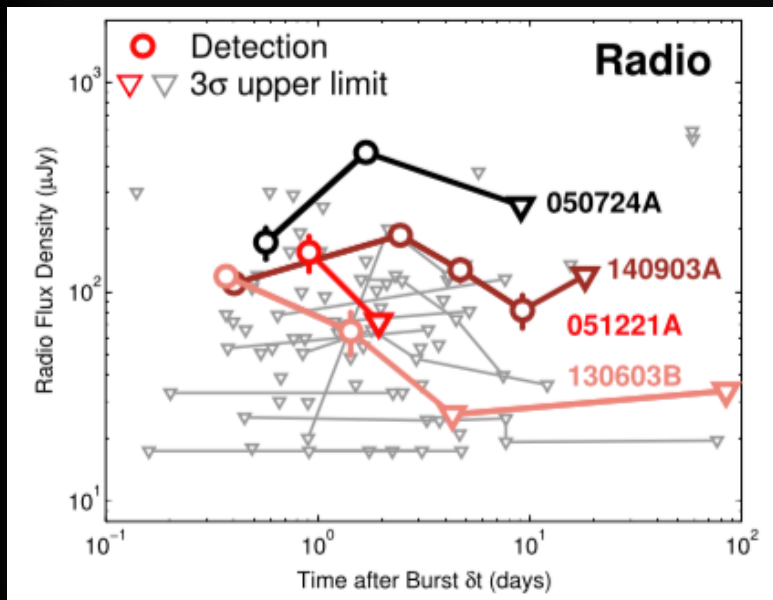
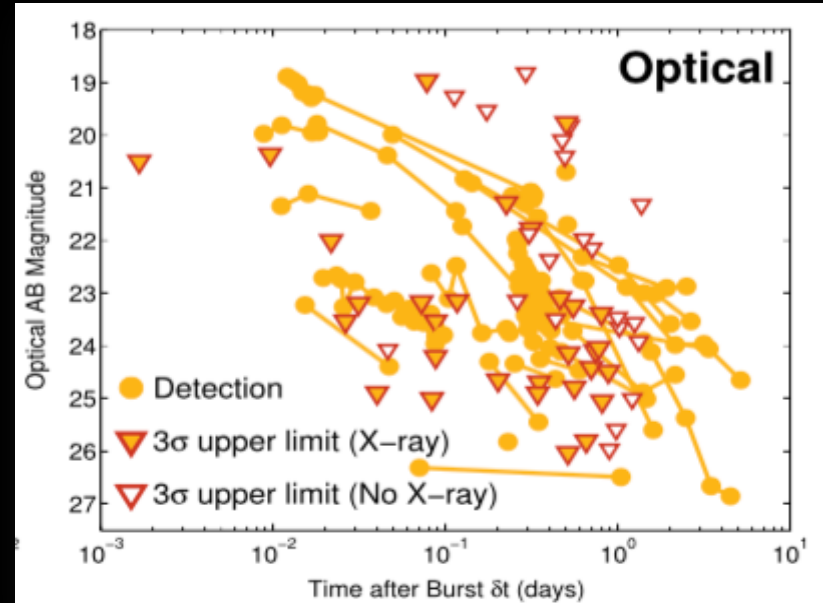
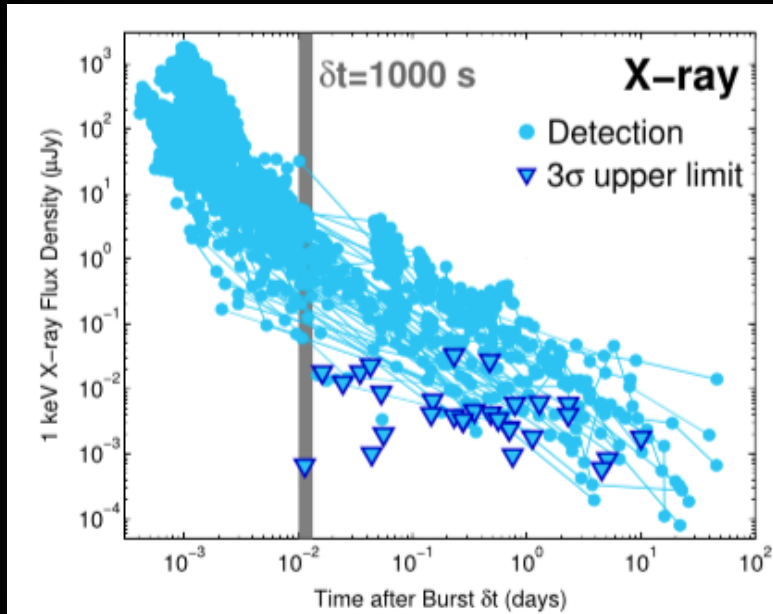
SHORT GRB

$E_{\text{jet}} = 1e50 \text{ erg}$, $n = 10^{-3} \text{ cm}^{-3}$

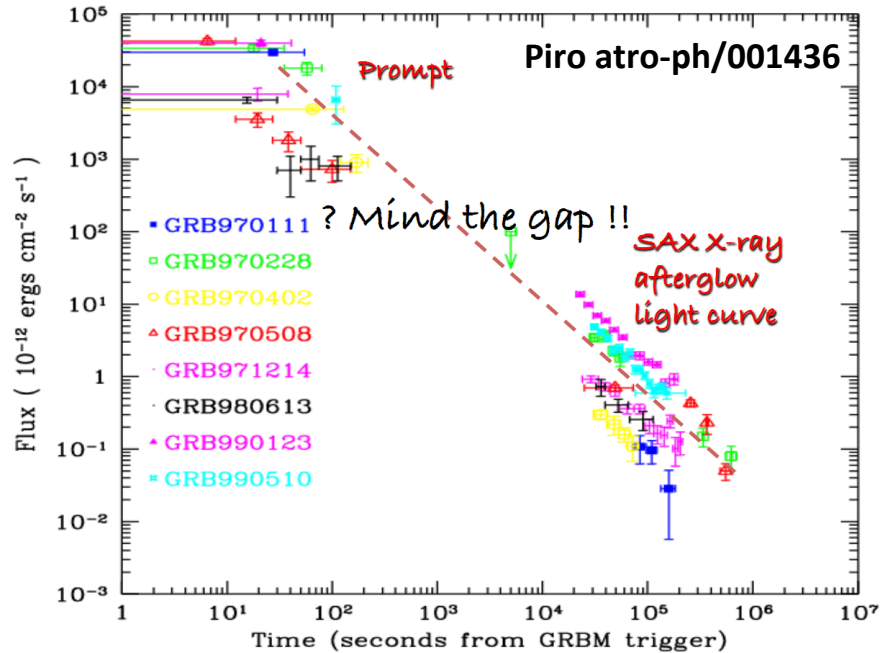
Modelled afterglows - Source at 200 Mpc



Short GRB afterglows in numbers



- About 140 SGRBs detected since 2005
- Afterglow detection percentage :
 - 90% in X-rays
 - 40% in opt
 - 7% in radio
- About 30 with redshift
- $z_{\min} = 0.12 \rightarrow 560$ Mpc
- Energy = 10^{48-52} erg



← X-ray afterglow before Swift

X-ray afterglow NOW ←

