



Multifield Quintessence Models and the Swampland

Based on:

[2002.02695](#) [gr-qc] & [2007.11011](#) [hep-th]

with M. Cicoli & Francisco Gil Pedro

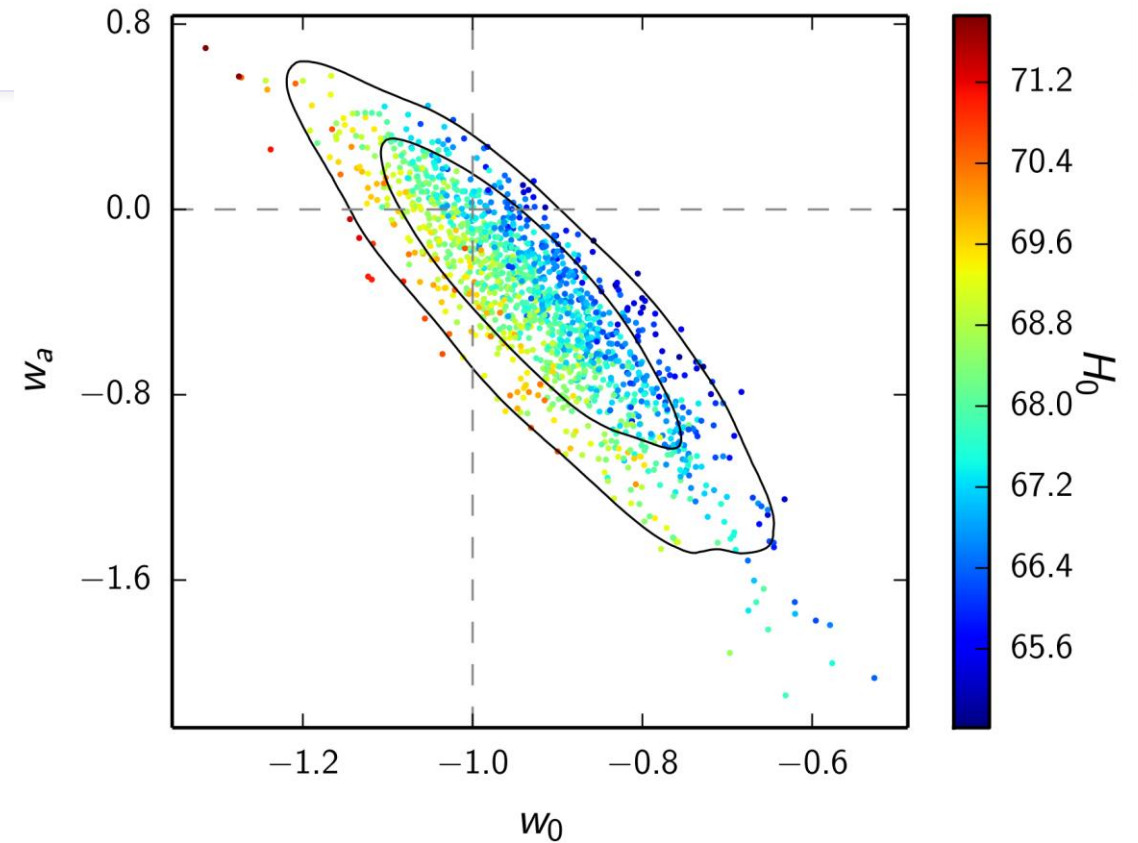
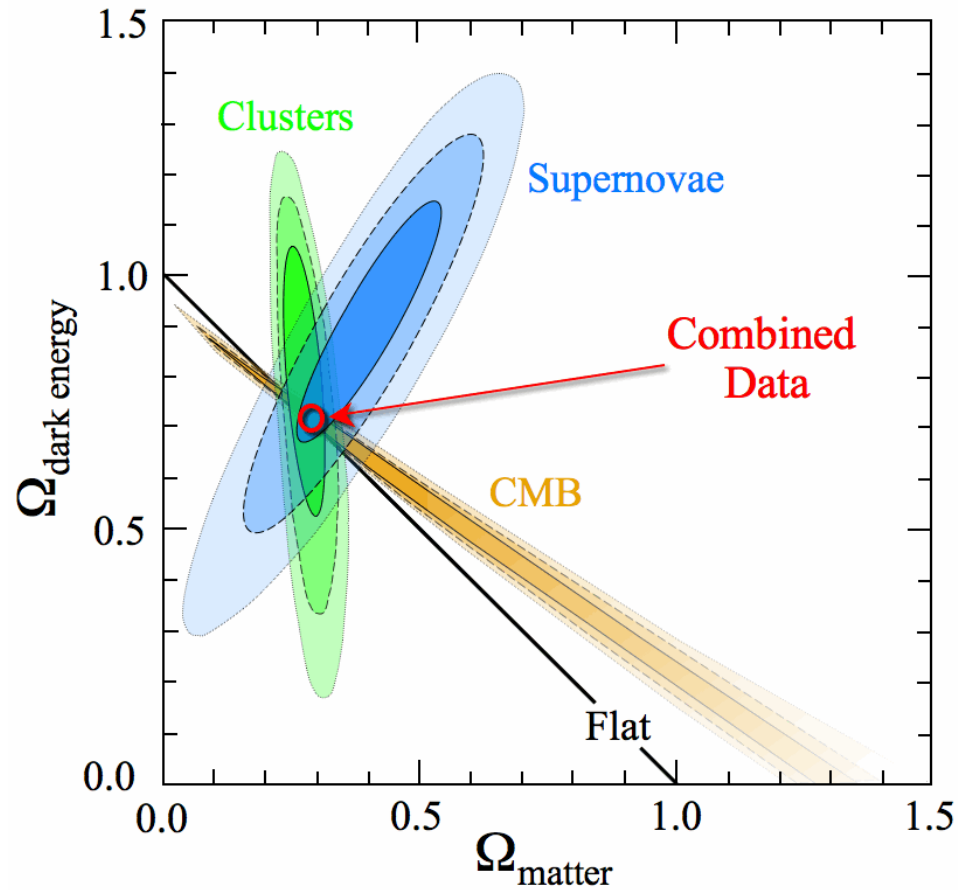




Outline

- Introduction & Motivation
- Single field quintessence and Swampland bounds
- Multifield models: an opportunity!
- Dynamics of 2-field models
- Non-geodesic trajectories and Swampland bounds
- Different scenarios
- Discussion

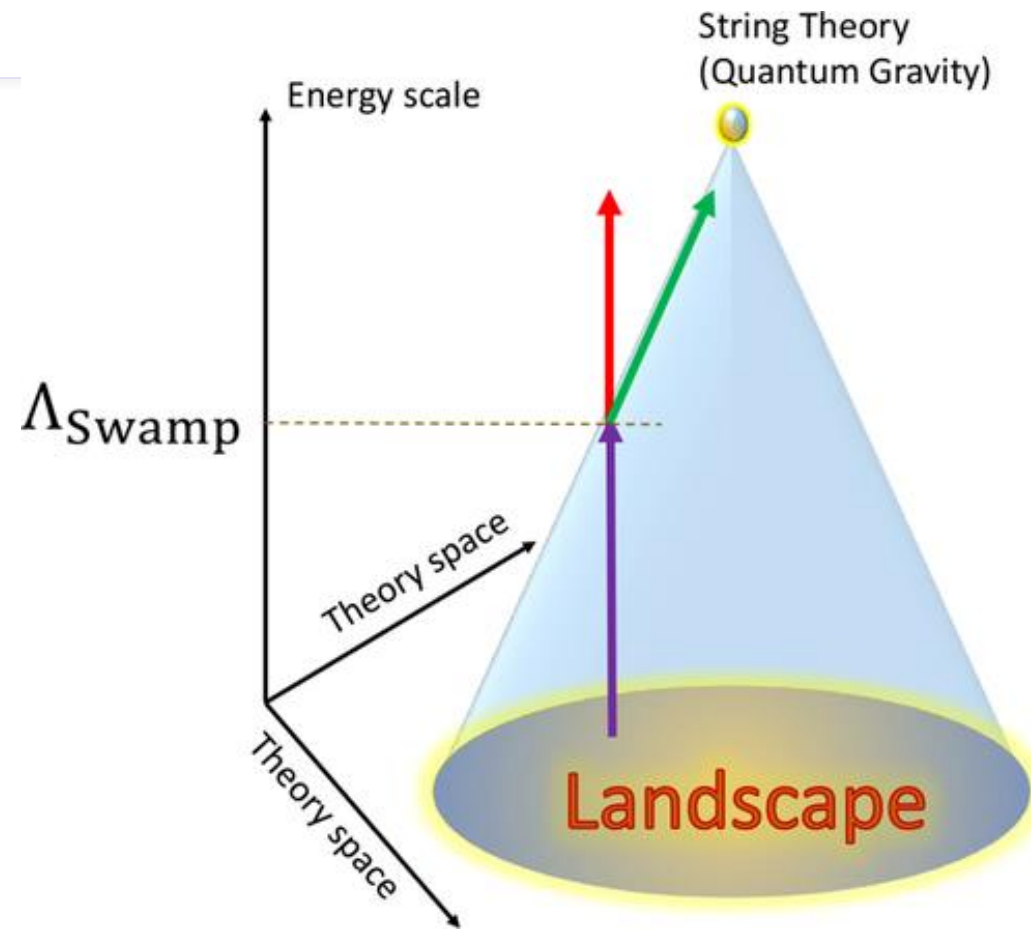
The Λ CDM model



$$\Omega_{\phi} = 0.7 \quad \omega_{\phi} = -1 \pm 0.09 \text{ at } 95\% \text{ CL}$$

[Credits: Planck coll. 1502.01589]

Swampland Conjectures



[Credits: Review by E. Palti, 2019]

The dS Swampland Conjecture

$$\frac{V_{\phi}}{V} \geq \frac{c}{M_{\text{Pl}}}$$

Only models with $c \approx 0,6$ remain compatible with the latest observations!

➡ Tension between theory & observations...

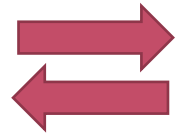
Single-field Quintessence Models

$$S[g_{\mu,\nu}, \phi] = \int d^4x \sqrt{-g} \left(\frac{M_{\text{Pl}}^2}{2} \mathcal{R} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right)$$

$$ds_4^2 = -dt^2 + a(t)^2 ds_{\mathbb{R}^3}^2 \quad \text{and} \quad \phi = \phi(t)$$

$$\epsilon_H \equiv -\frac{\dot{H}}{H^2} < 1$$

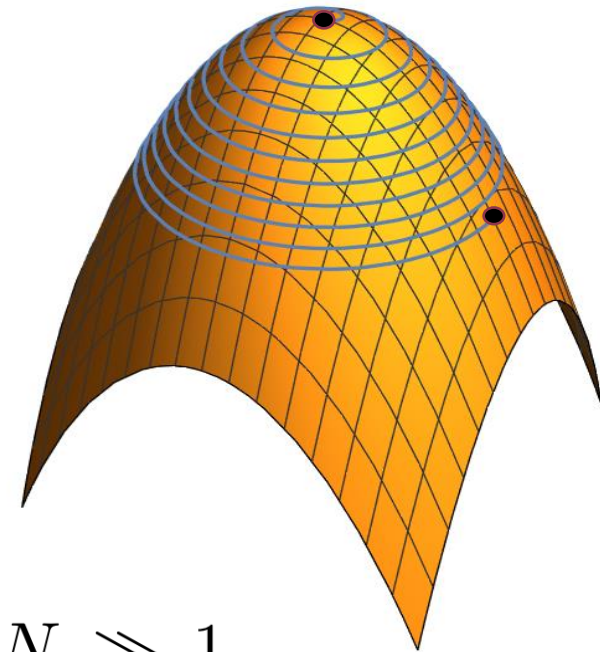
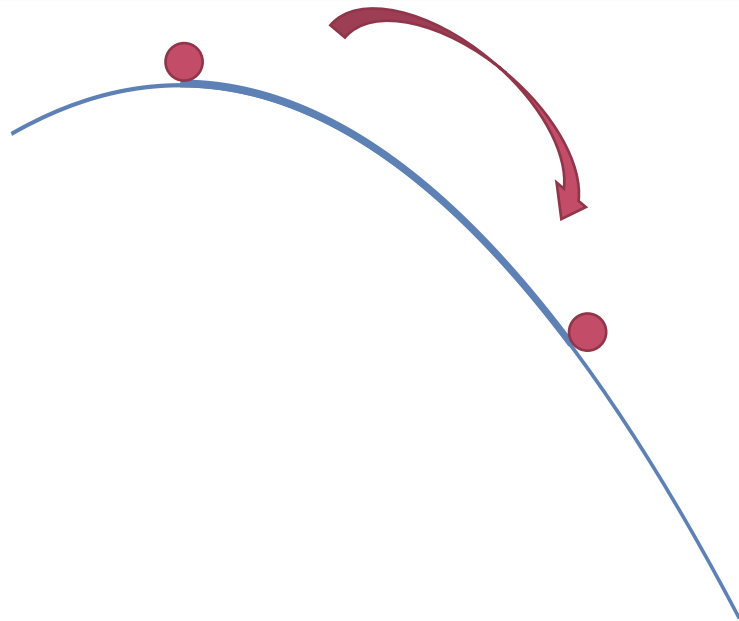
Cosmic acceleration



$$\epsilon_V \equiv \frac{M_{\text{Pl}}^2}{2} \left(\frac{V_\phi}{V} \right)^2 \approx \epsilon_H < 1$$

Flatness of scalar potential V

Multifield Models: an opportunity?



A 2-field Model

$$S[g_{\mu,\nu}, \phi] = \int d^4x \sqrt{-g} \left(\frac{M_{\text{Pl}}^2}{2} \mathcal{R} - \frac{1}{2} (\partial\phi_1)^2 - \frac{1}{2} f(\phi_1)^2 (\partial\phi_2)^2 - V(\phi_1) \right)$$

$$ds_4^2 = -dt^2 + a(t)^2 ds_{\mathbb{R}^3}^2 \quad \text{and} \quad \phi_i = \phi_i(t)$$

The EOM's read

$$\begin{cases} \ddot{\phi}_1 + 3H\dot{\phi}_1 - f f_1 \dot{\phi}_2^2 + V_1 = 0 , \\ \ddot{\phi}_2 + 3H\dot{\phi}_2 + 2\frac{f_1}{f} \dot{\phi}_2 \dot{\phi}_1 = 0 , \end{cases}$$

Together with the Friedman equation

$$H^2 = \frac{1}{3M_{\text{Pl}}^2} \left(\frac{\dot{\phi}_1^2}{2} + \frac{f^2}{2} \dot{\phi}_2^2 + V + \rho_b \right)$$

Now introduce

$$k_1(\phi_1) \equiv -M_{\text{Pl}} \frac{\dot{f}_1}{f}$$

and

$$k_2(\phi_1) \equiv -M_{\text{Pl}} \frac{\dot{V}_1}{V}$$

as well as:

$$x_1 \equiv \frac{\dot{\phi}_1}{\sqrt{6}H M_{\text{Pl}}} , \quad x_2 \equiv \frac{f \dot{\phi}_2}{\sqrt{6}H M_{\text{Pl}}} , \quad y_1 \equiv \frac{\sqrt{V}}{\sqrt{3}H M_{\text{Pl}}}$$

$$x'_1 = 3x_1(x_1^2 + x_2^2 - 1) + \sqrt{\frac{3}{2}}(-2k_1x_2^2 + k_2y_1^2) - \frac{3}{2}\gamma x_1(x_1^2 + x_2^2 + y_1^2 - 1) ,$$

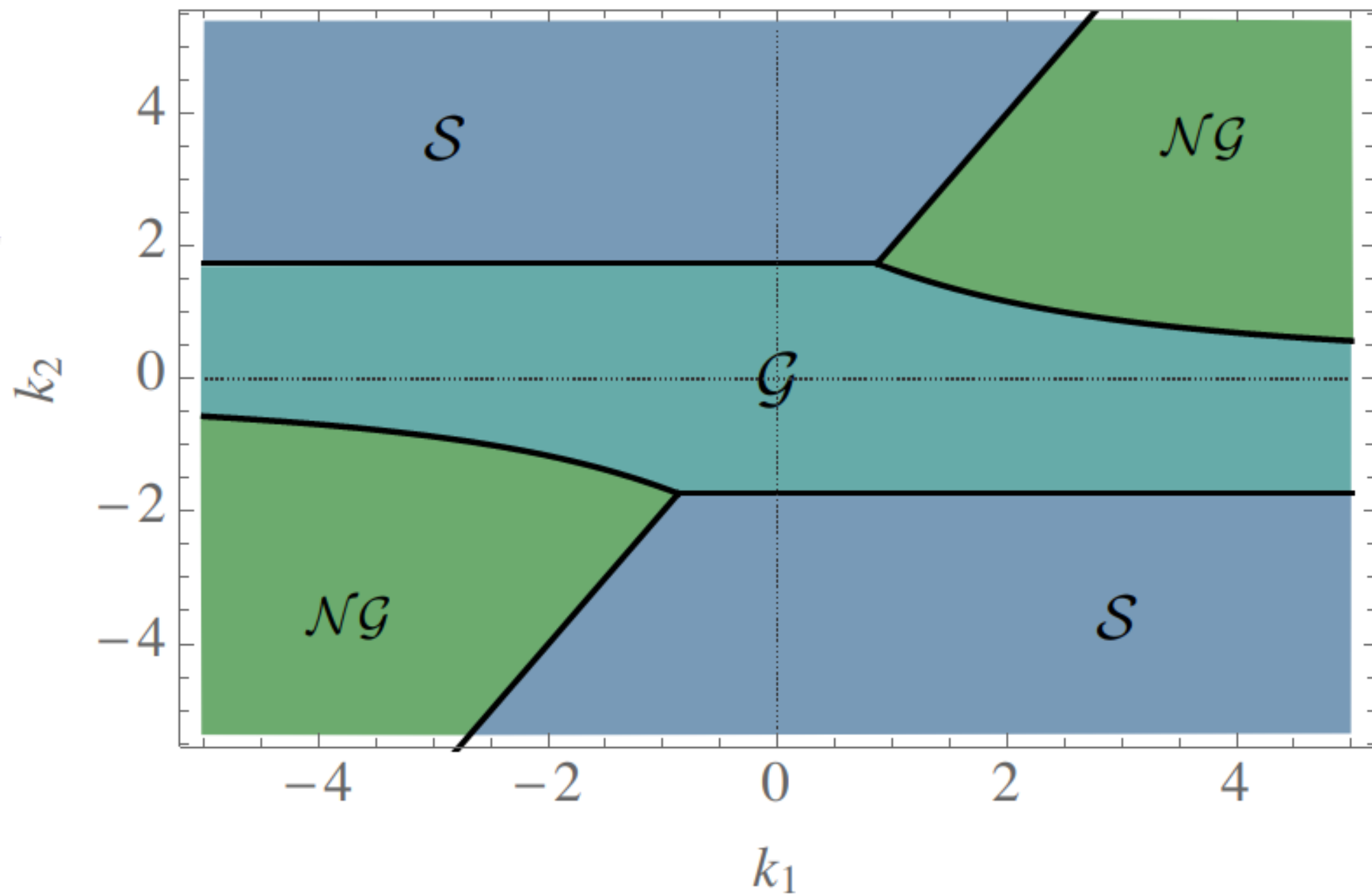
$$x'_2 = 3x_2(x_1^2 + x_2^2 - 1) + \sqrt{6}k_1x_1x_2 - \frac{3}{2}\gamma x_2(x_1^2 + x_2^2 + y_1^2 - 1) ,$$

$$y'_1 = -\sqrt{\frac{3}{2}}k_2x_1y_1 - \frac{3}{2}\gamma y_1(x_1^2 + x_2^2 + y_1^2 - 1) + 3y_1(x_1^2 + x_2^2) ,$$

The Fixed Points Analysis

$$x'_1 = x'_2 = y'_1 = 0$$

	x_1	x_2	y_1	Ω_ϕ	ω_ϕ	existence
\mathcal{K}_+	1	0	0	1	1	all k_1, k_2, γ
\mathcal{K}_-	-1	0	0	1	1	all k_1, k_2, γ
\mathcal{F}	0	0	0	0	undefined	all k_1, k_2, γ
\mathcal{S}	$\frac{\sqrt{3/2\gamma}}{k_2}$	0	$\frac{\sqrt{3/2\gamma(2-\gamma)}}{k_2}$	$\frac{3\gamma}{k_2^2}$	$\gamma - 1$	$0 < \gamma < 2 \wedge k_2^2 \geq 3\gamma$
\mathcal{G}	$\frac{k_2}{\sqrt{6}}$	0	$\sqrt{1 - \frac{k_2^2}{6}}$	1	$-1 + \frac{k_2^2}{3}$	$k_2 < \sqrt{6}$
\mathcal{NG}	$\frac{\sqrt{6}}{(2k_1 + k_2)}$	$\frac{\pm \sqrt{k_2^2 + 2k_2 k_1 - 6}}{2k_1 + k_2}$	$\sqrt{\frac{2k_1}{2k_1 + k_2}}$	1	$\frac{k_2 - 2k_1}{k_2 + 2k_1}$	$k_2 \geq \sqrt{6 + k_1^2} - k_1$



To better connect with **observations**, the scalar dynamics can be formulated in terms of the **density parameter** and the **equation of state parameter**, respectively:

$$\Omega_\phi = x_1^2 + x_2^2 + y_1^2 \quad \text{and} \quad \omega_\phi = \frac{p_\phi}{\rho_\phi} = \frac{x_1^2 + x_2^2 - y_1^2}{x_1^2 + x_2^2 + y_1^2}$$

The time evolution within the $(\omega_\phi, \Omega_\phi)$ plane is determined by the following ODE's

$$\Omega'_\phi = -3(\Omega_\phi - 1)\Omega_\phi(\omega_b - \omega_\phi)$$

$$\omega'_\phi = (\omega_\phi - 1) \left(-k_2 \sqrt{3(\omega_\phi + 1)\Omega_\phi - 6x_2^2} + 3(1 + \omega_\phi) \right)$$

NG Dynamics & Swampland Bounds

The displacement in field space can be expressed as

$$\Delta\phi \equiv \int dt \sqrt{\gamma_{ij} \dot{\phi}^i \dot{\phi}^j}$$

Which in our case, evaluates to:

$$\frac{\Delta\phi}{M_{Pl}} = \int dN \sqrt{-\left(\frac{3}{2k_1}\right)^2 + 3\frac{k_2}{k_1}}$$

$$\frac{\Delta\phi}{M_{Pl}} \sim \frac{\Delta N}{\sqrt{k_1}} < 1$$



This means that we can realize even $O(10)$ e-folds of acceleration
without violating the Distance Swampland Conjecture!!!!

The simplest scenario: Exp-type f & V

$$f = f_0 e^{-k_1 \phi_1}$$

$$V = V_0 e^{-k_2 \phi_1}$$



k_1 and k_2 are **constant**!

$$\omega_\phi = \frac{k_2 - 2k_1}{k_2 + 2k_1} = -1 + \frac{2}{1 + 2\frac{k_1}{k_2}}$$



Need a little
«hierarchy» between
them!

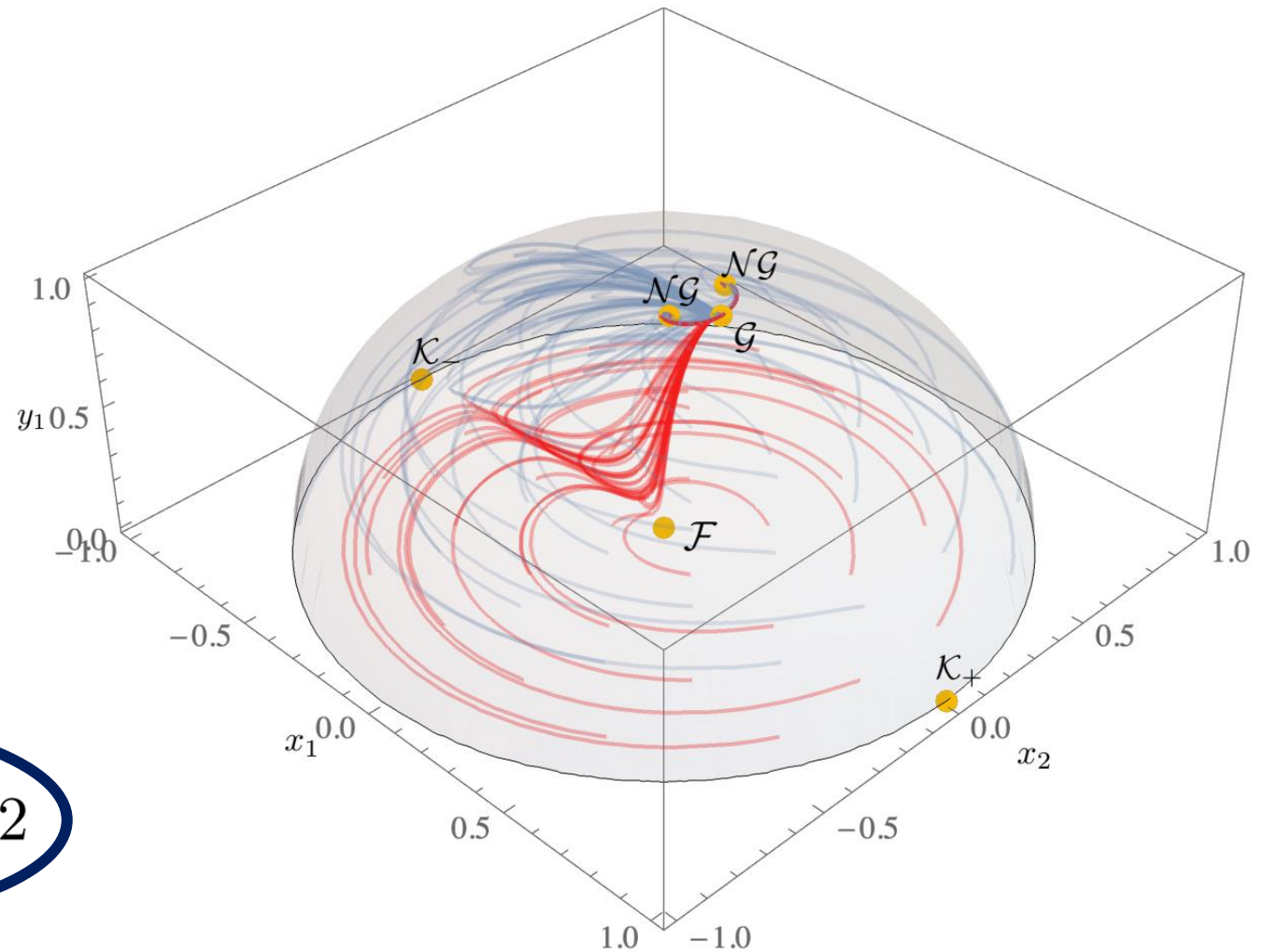
$$k_1 \gg k_2$$

Should be **small**...

The system evolves from matter domination towards \mathcal{G} (saddle), before settling into \mathcal{NG} (stable).

Plot:

$$k_1 = 10, k_2 = 1/2$$



Stringy Examples...

Stabilize everything but **ONE** complex modulus

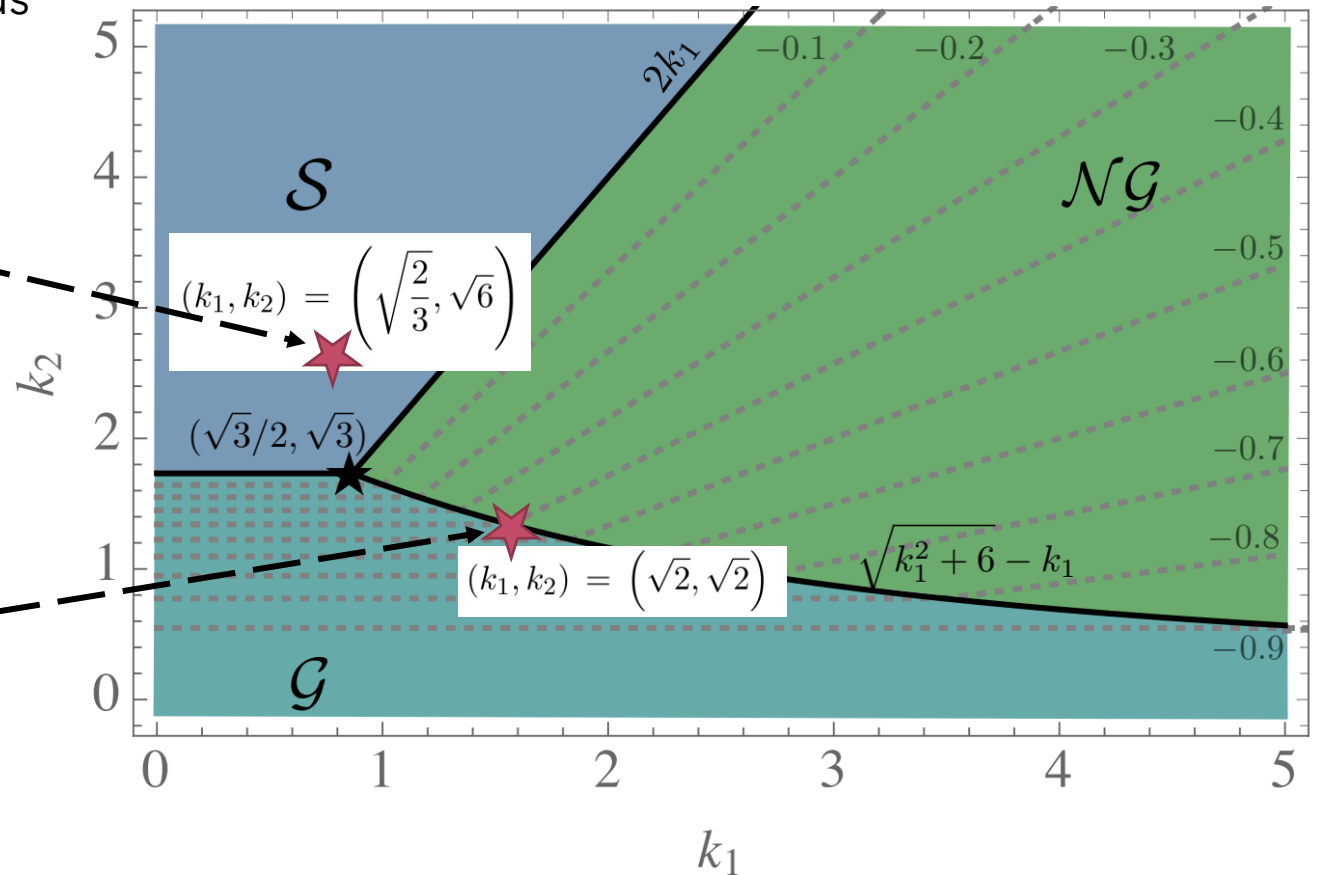
- Type IIB on a CY 3-fold with O3 planes

$$T = C_{(4)} + i \text{vol}_6$$

[Saltman, Silverstein, 2002]

- Heterotic on CY 3-fold

$$S = B_{(6)} + i e^{-\Phi}$$



The case of non-Exp f 's

k_1 and k_2 are no longer **constant**!

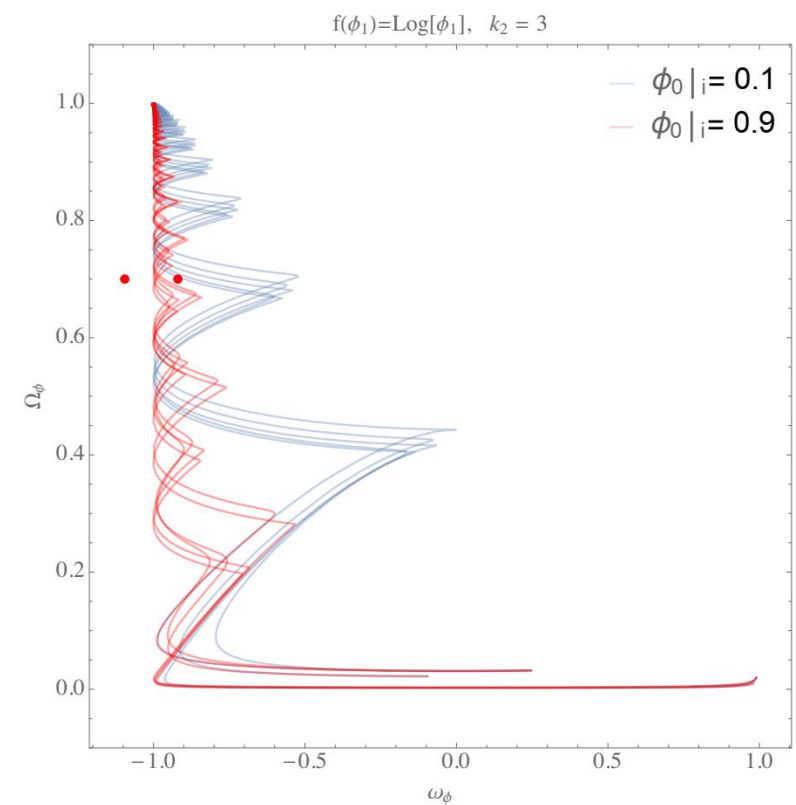
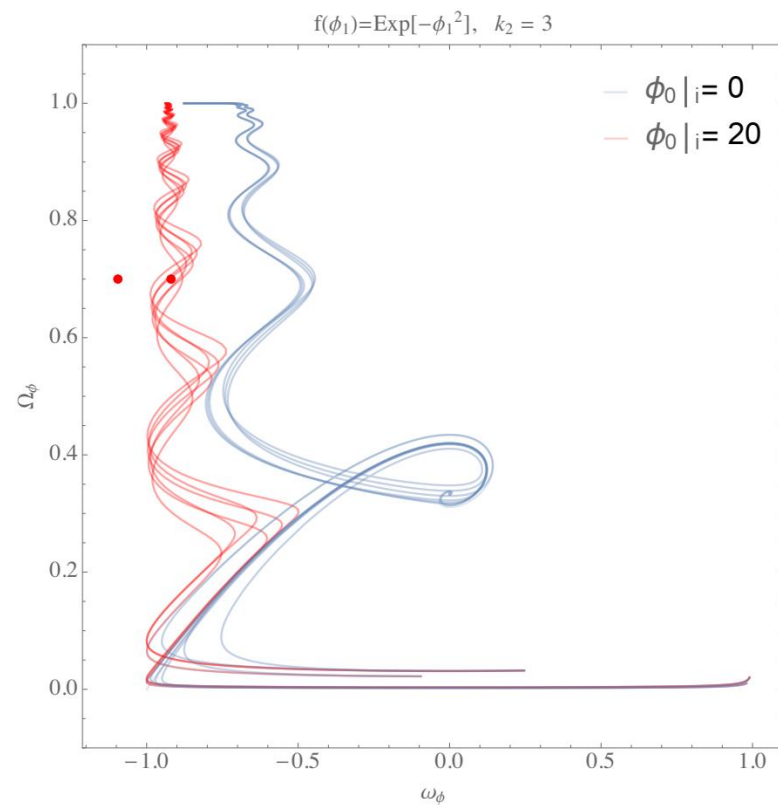
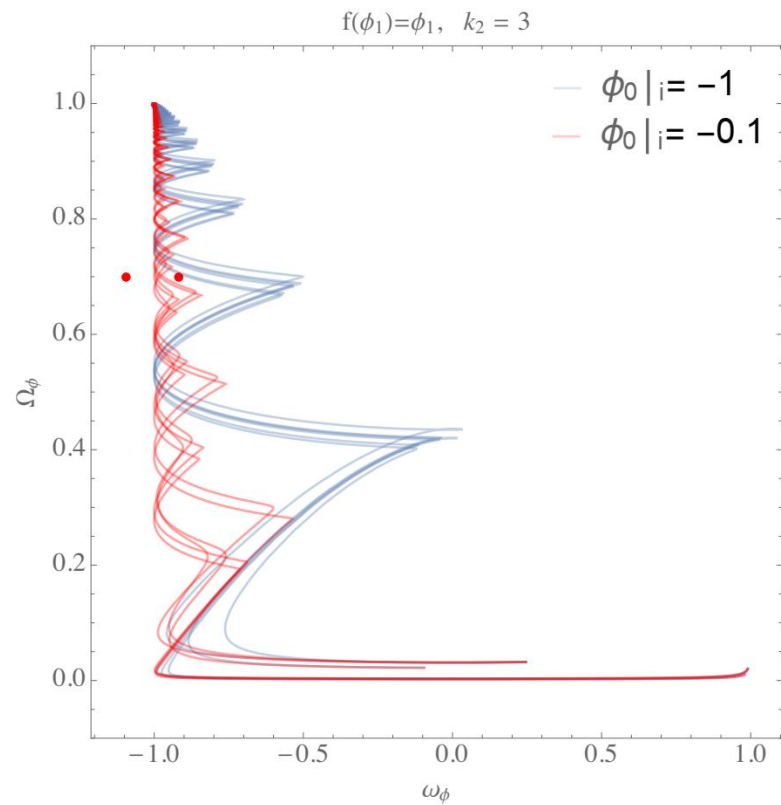
Their time evolution is given by

$$k'_i = \sqrt{6}k_i^2 x_1 (1 - \Gamma_i)$$

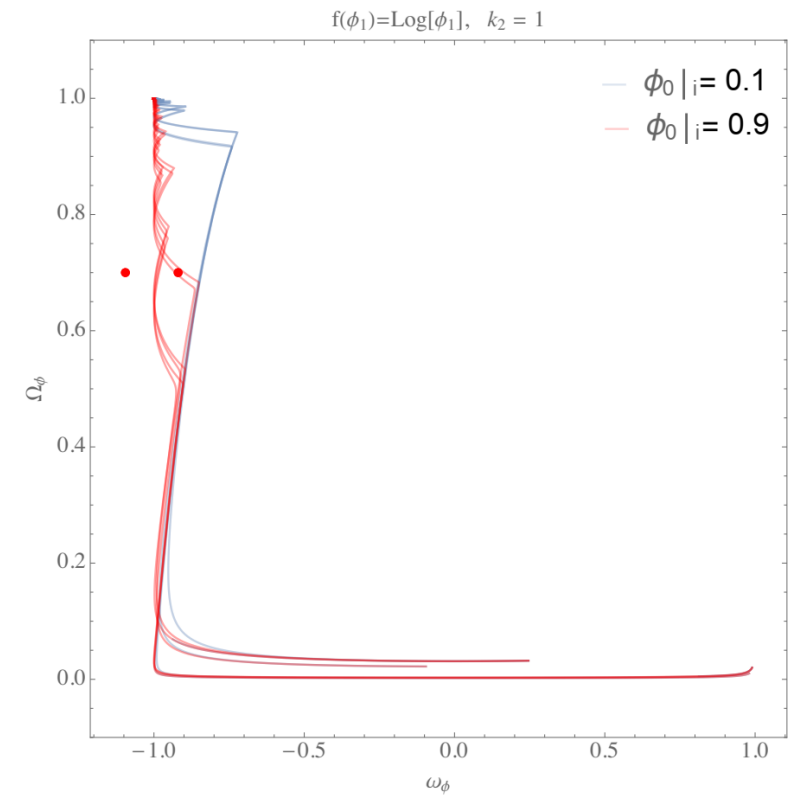
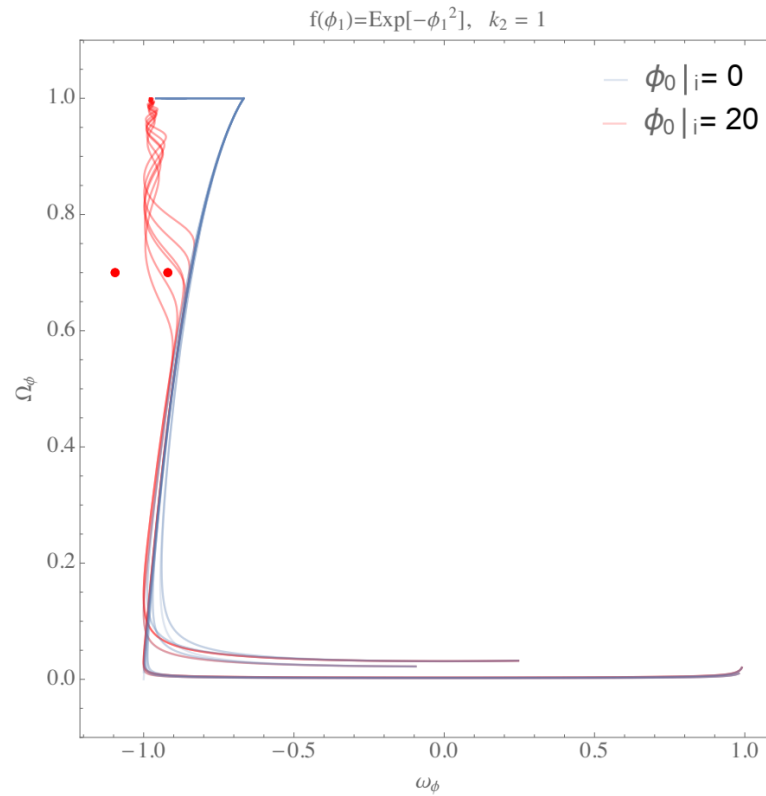
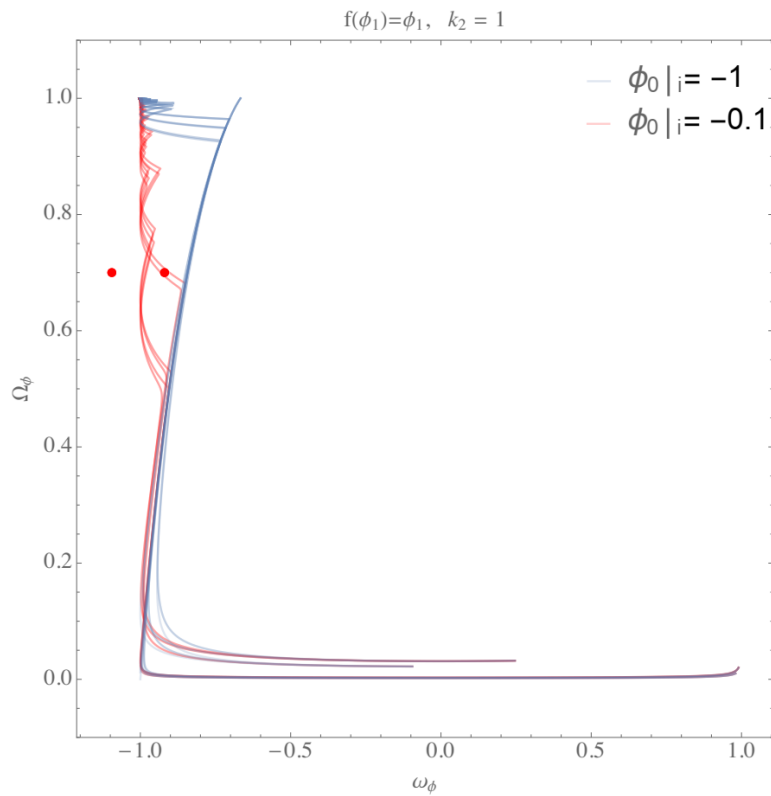
Where:

$$\Gamma_1 \equiv \frac{f_{11}f}{f_1^2} \quad \text{and} \quad \Gamma_2 \equiv \frac{V_{11}V}{V_1^2}$$

Growing kinetic coupling, steep potential

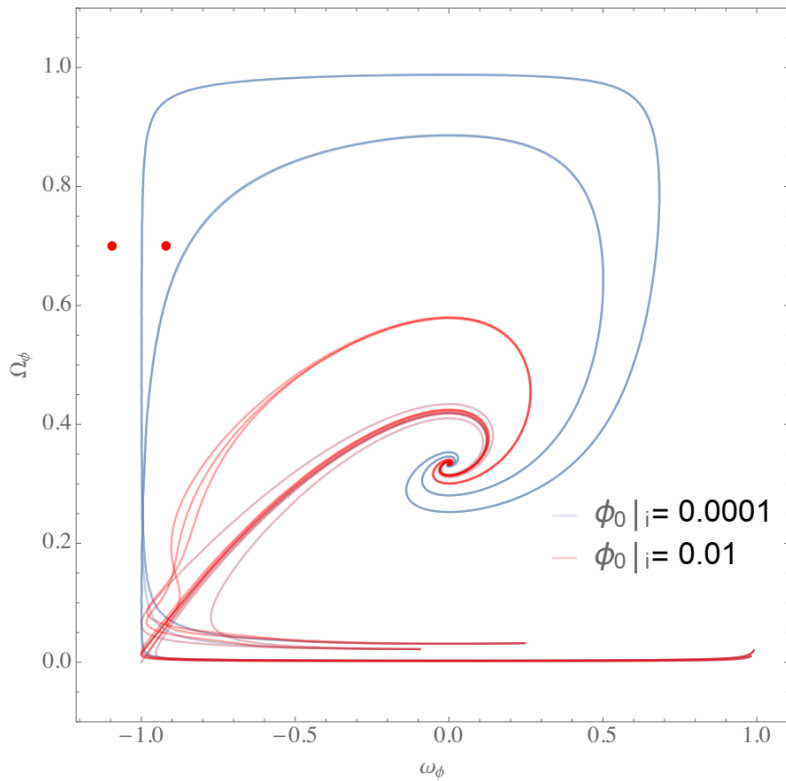


Growing kinetic coupling, shallow potential

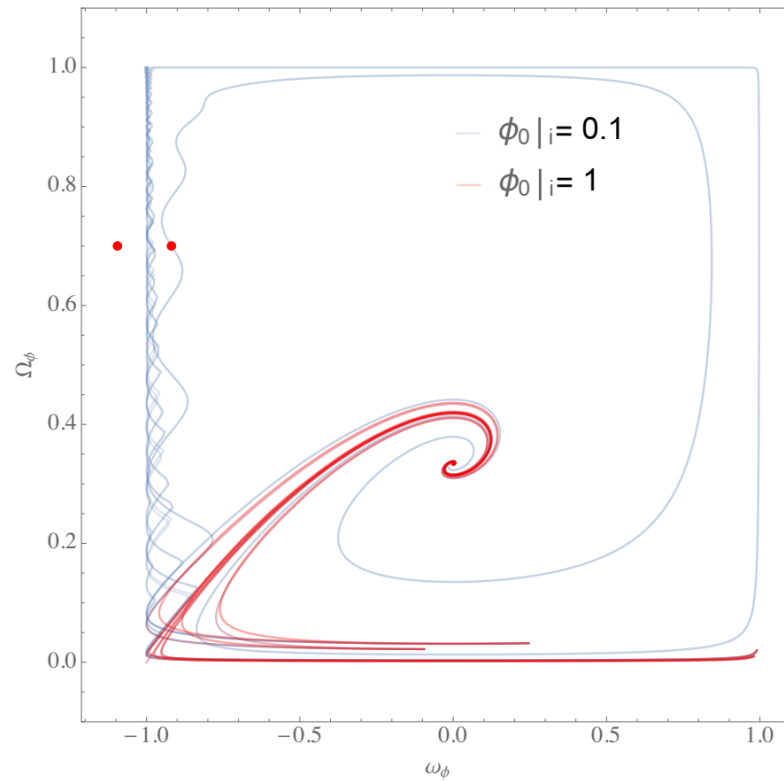


Decaying kinetic coupling, steep potential

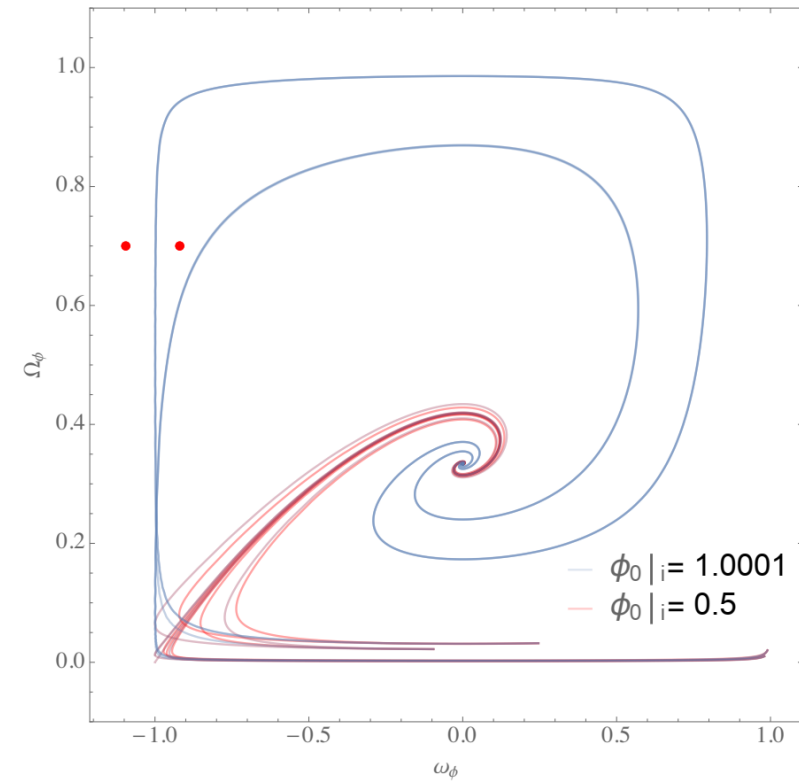
$$f(\phi_1)=1/\phi_1, \quad k_2 = 3$$



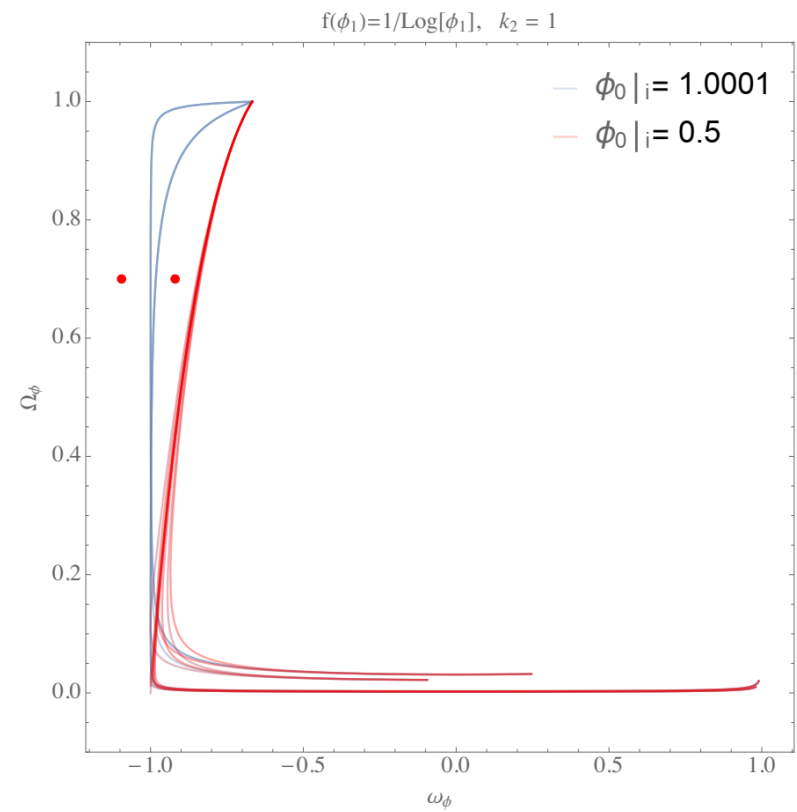
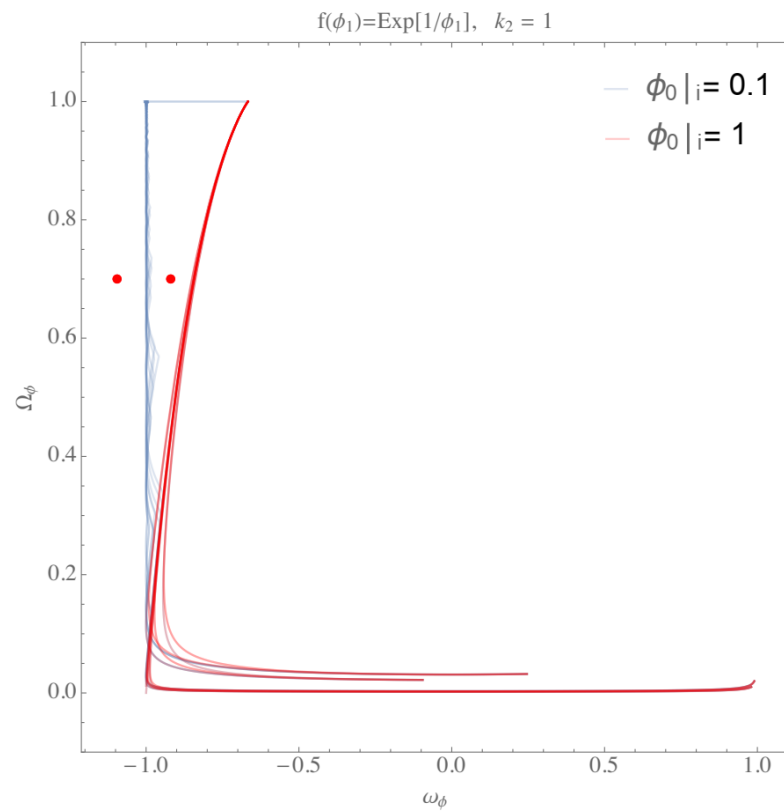
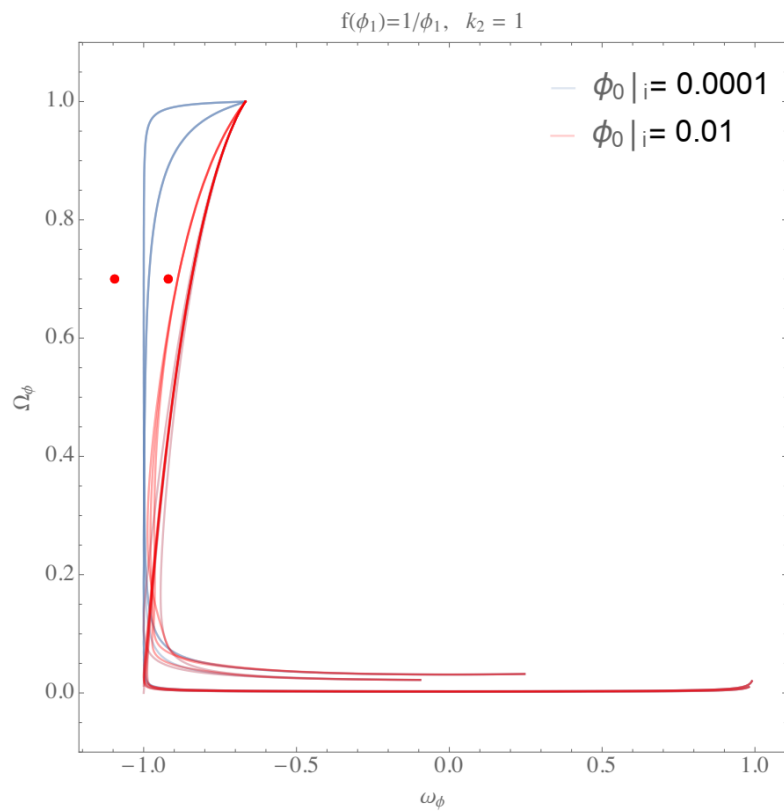
$$f(\phi_1)=\text{Exp}[1/\phi_1], \quad k_2 = 3$$



$$f(\phi_1)=1/\text{Log}[\phi_1], \quad k_2 = 3$$



Decaying kinetic coupling, shallow potential



Towards more general scenarios...

It's **hard to be exhaustive** and the equations get much more complicated...

However, an **interesting** case is given by the following setup

$$V = V_0 \phi_1^{p_2}$$

and

$$f = f_0 \phi_1^{p_1}$$

Even though the k 's are not constant, **their ratio still is!**

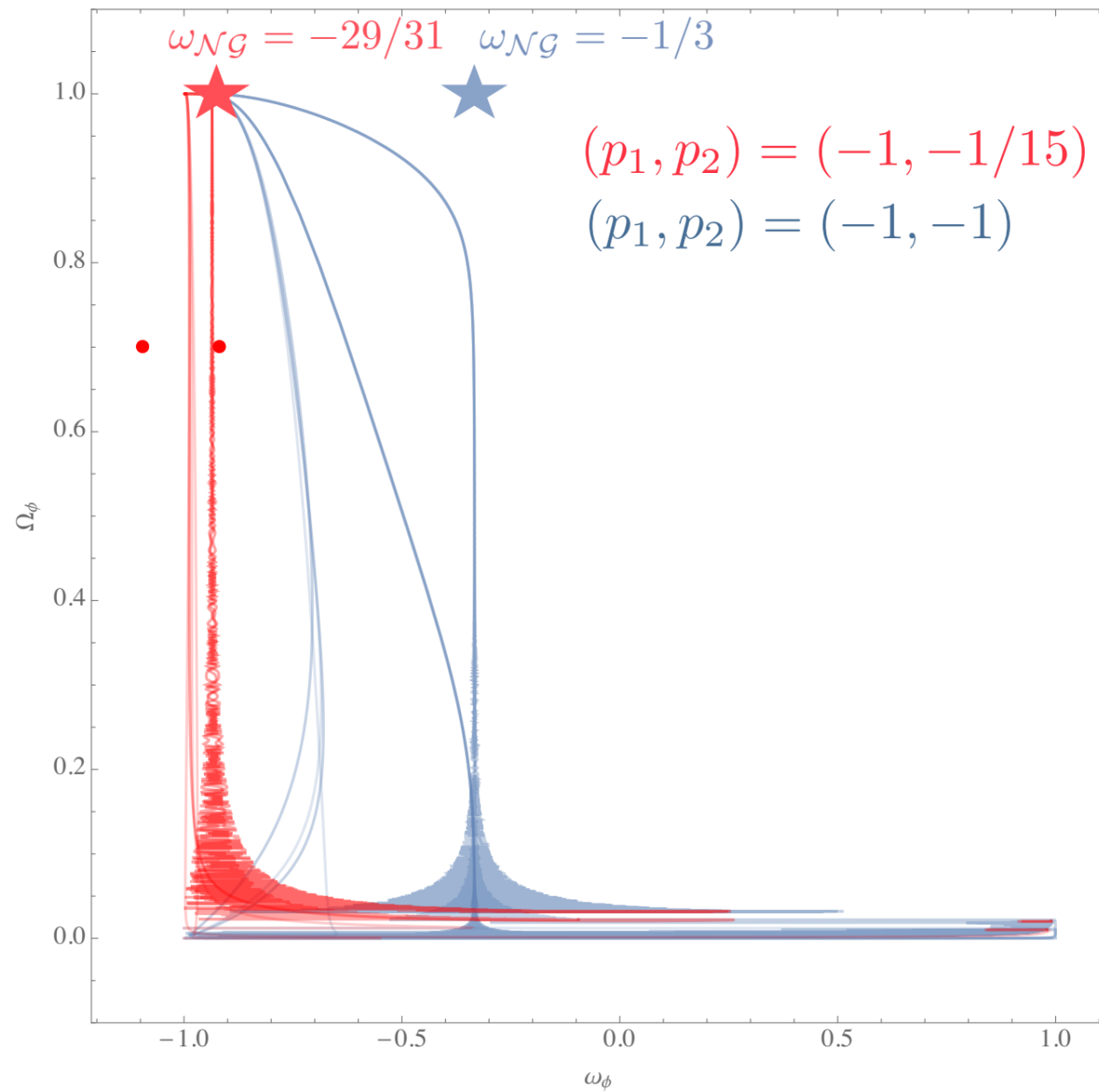
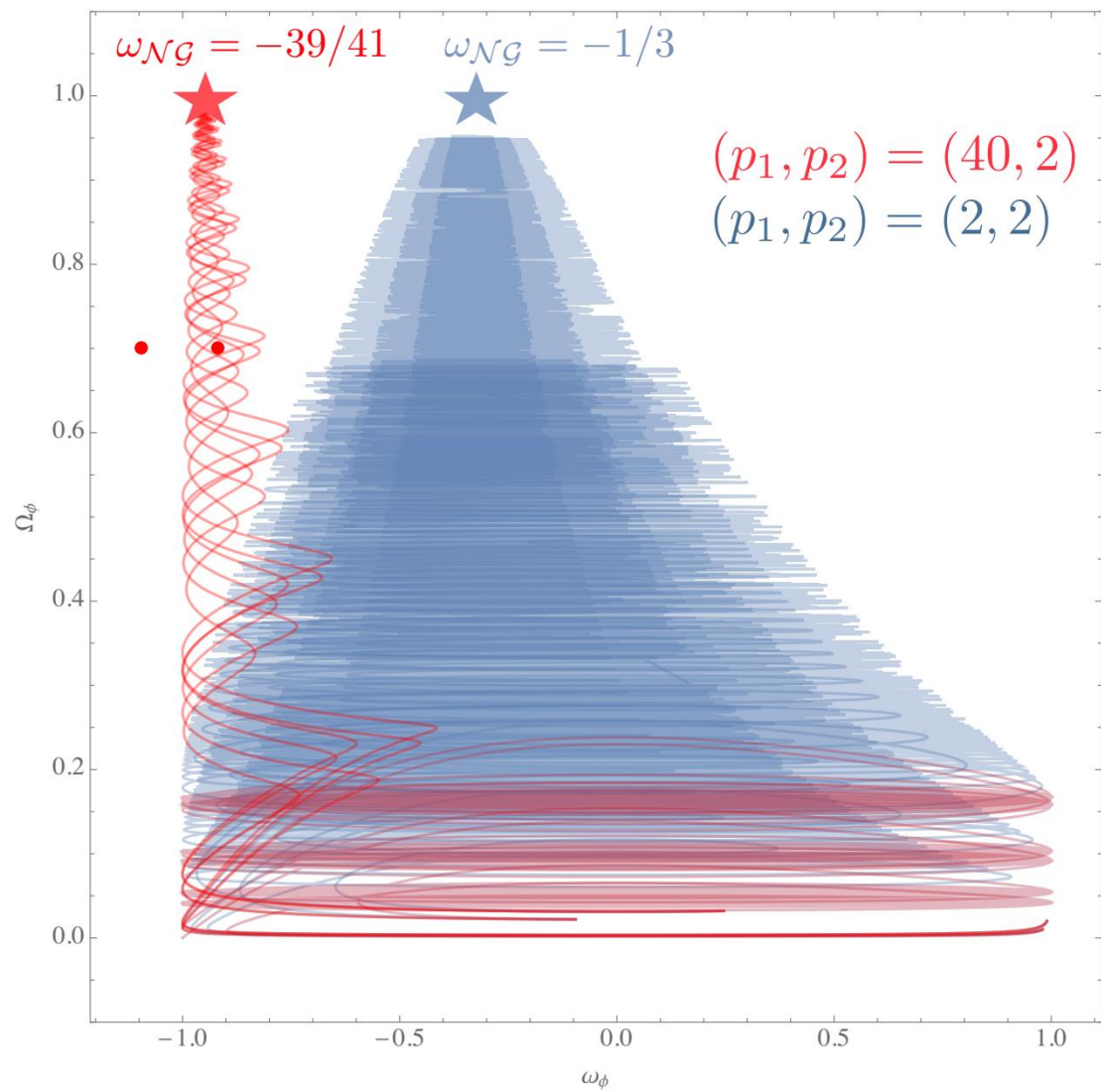


$$\frac{k_1}{k_2} = \frac{p_1}{p_2}$$

Need the same hierarchy but this time for the p 's!



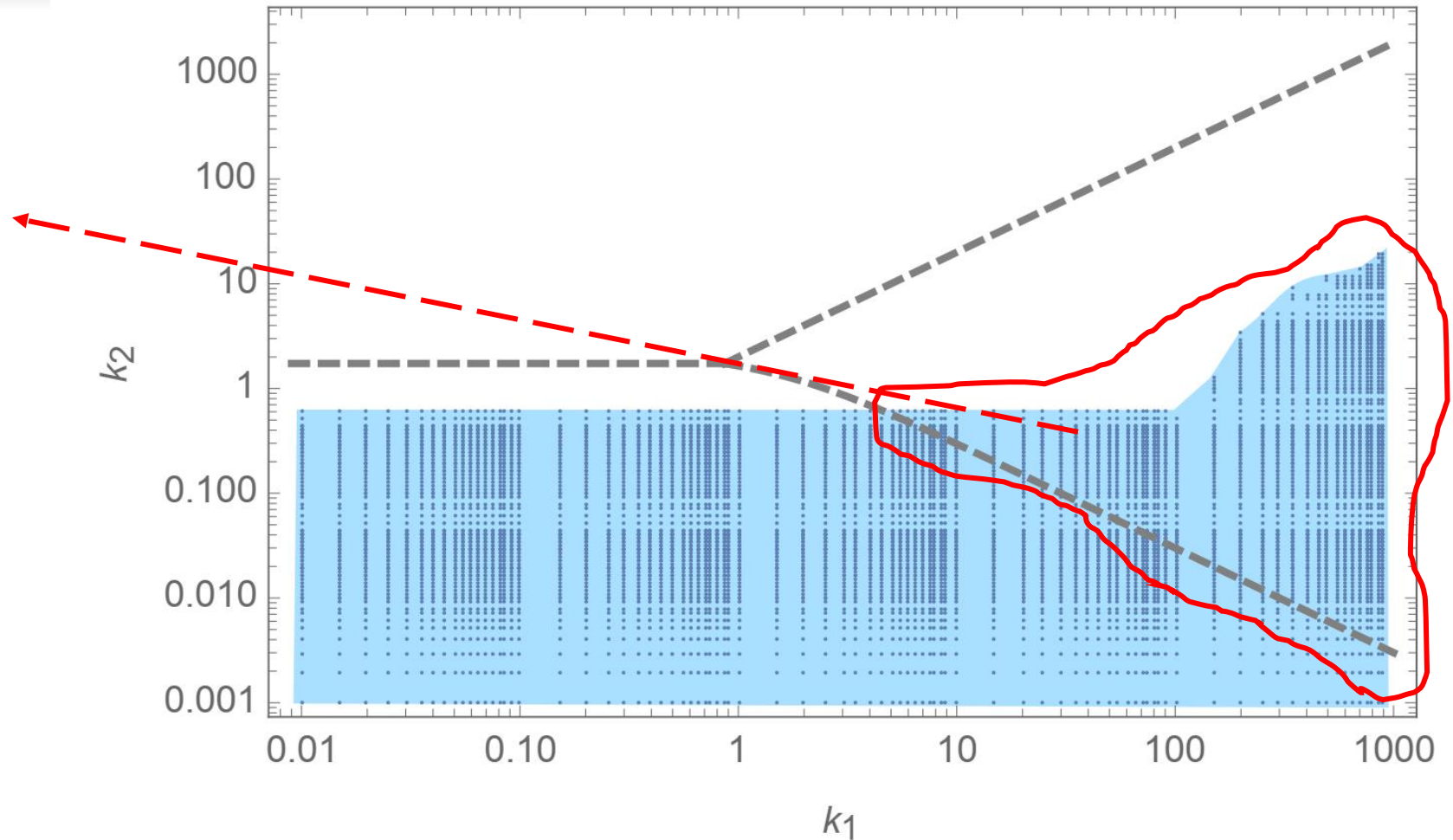
$$\omega_{\mathcal{NG}} = \frac{p_2 - 2p_1}{2p_1 + p_2}$$



What about string theory?

Here's where
you would like
your model to
be!

BUT...



Conclusions & Outlook

- Cosmic acceleration at late times is a **challenge** for QG
- A **dS** phase has been conjectured to be in the **Swampland**
- What about **Quintessence**? Single-field needs flat potentials and hence has the same fate as dS
- **Multifield** Quintessence offers a way out cause it may realize a quasi-dS phase with **NO need for flat potentials**
- A simple setup though, still needs some **hierarchy** which is not very natural for stringy constructions



We are working on this, so stay tuned...

Thank you for your attention!

