

Quantifying the efficiency of state preparation via quantum variational eigensolvers

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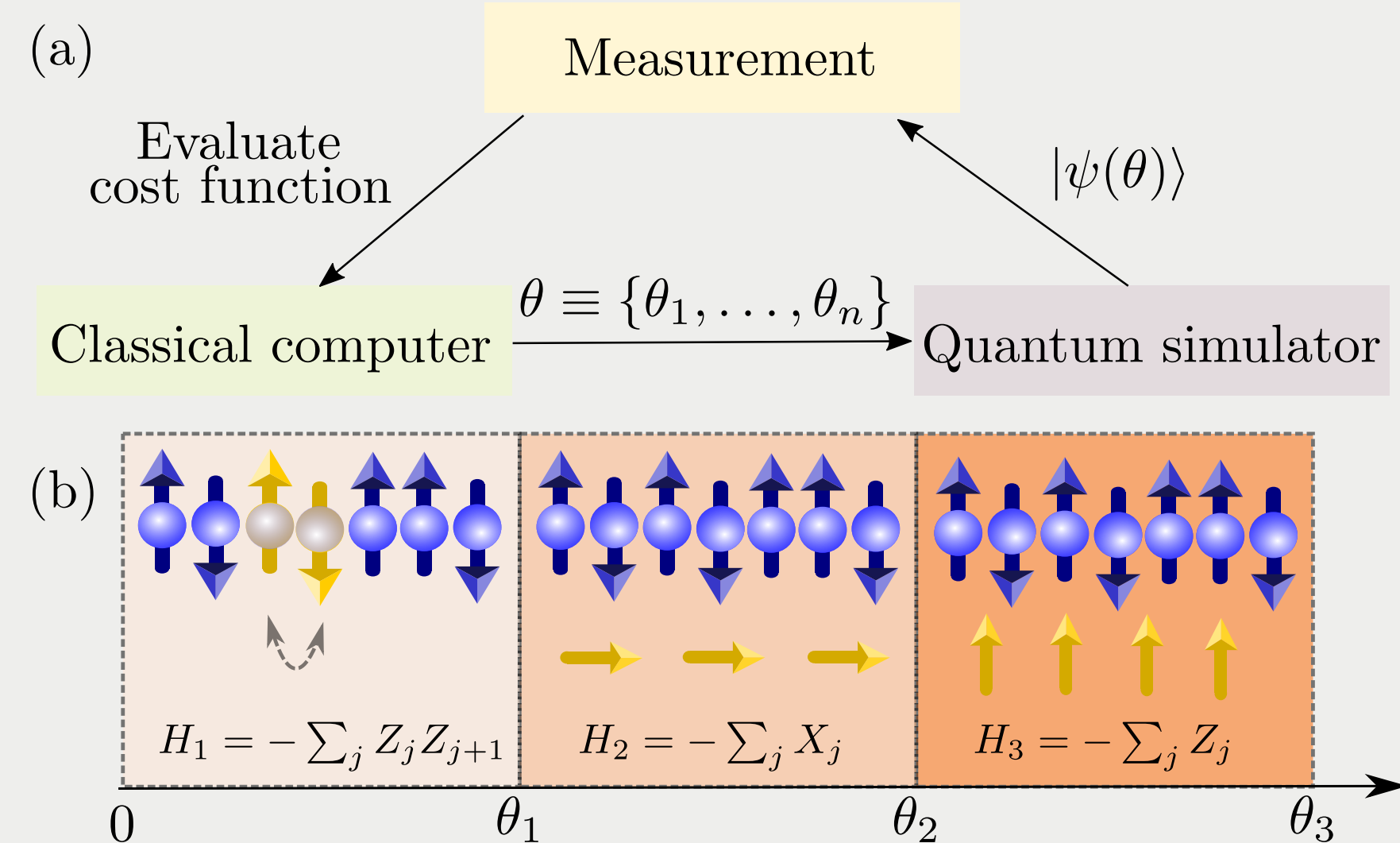
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Quantum Approximate Optimization Algorithm (QAOA)

QAOA is a variational algorithm for preparing quantum states using low-depth quantum circuits [1]. It optimizes over a set of angles θ in order to minimize some cost function, e.g. the energy $\langle \psi(\theta) | H | \psi(\theta) \rangle$, where

$$|\psi(\theta)\rangle = e^{-i\theta_{p,1}H_1} \dots e^{-i\theta_{p,M}H_M} \dots e^{-i\theta_{1,1}H_1} \dots e^{-i\theta_{1,M}H_M} |\psi_0\rangle, \quad (1)$$

$|\psi_0\rangle$ is some easily preparable initial state, $\{H_j\}_{j \in \{1, \dots, M\}}$ is some set of Hamiltonians, and p is the number of steps.



QAOA is able to prepare non-trivial quantum states [2] e.g. the GHZ state.

Interaction distance

Given some density matrix ρ , we define the *interaction distance* of ρ [4] as

$$\mathcal{D}_{\mathcal{F}}(\rho) := \min_{\sigma \in \mathcal{F}} \frac{1}{2} \text{Tr} \left(\sqrt{(\rho - \sigma)^2} \right), \quad (2)$$

where \mathcal{F} is the manifold of Gaussian density matrices, which are of the form $\sigma = e^H$, with H quadratic. Interaction distance measures how far a density matrix is from a Gaussian density matrix.

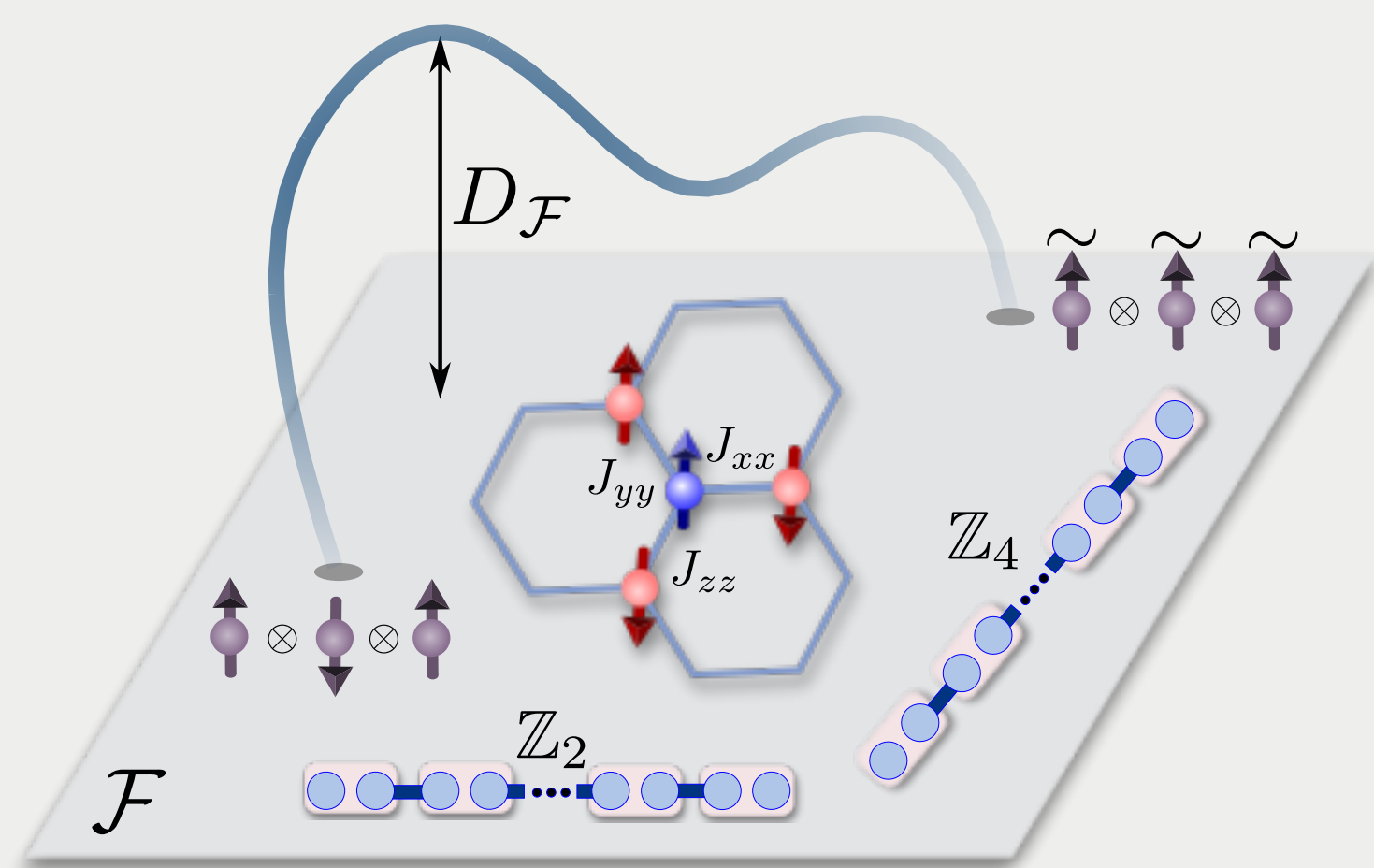


Figure 1: Interaction distance can be interpreted as a distance to the closest free fermion model [3]. The ground states of free fermion models are always Gaussian.

This minimization can be restricted to the Gaussian states simultaneously diagonalizable with ρ , parameterized by the free energies ϵ :

$$\mathcal{D}_{\mathcal{F}}(\rho) = \min_{\epsilon} \frac{1}{2} \sum_k \left| e^{-\beta \rho_k} - \frac{e^{-\beta \sigma_k(\epsilon)}}{\sum_j e^{-\beta \sigma_j(\epsilon)}} \right|, \quad (3)$$

$$\sigma_k(\epsilon) = \sum_{j=0}^{N-1} \epsilon_j n_j^{(k)}, \quad n_j^{(k)} \in \{0, 1\}$$

ρ_k is the entanglement spectrum of ρ in descending order; k also labels the Fock basis such that the σ_k are in descending order, and $n_j^{(k)}$ is the occupancy number of the j th site.

Ising model

We consider the non-integrable Ising model with PBC

$$H = - \sum_{i=0}^{N-1} \pm \sigma_i^z \sigma_{i+1 \bmod N}^z - h_z \sum_{i=0}^{N-1} \sigma_i^z - h_x \sum_{i=0}^{N-1} \sigma_i^x, \quad (4)$$

In [2], it was argued that the ground states of the transverse field Ising model ($h_x = 0$) can be exactly prepared with $p = \frac{N}{2}$, using $H_1 = - \sum_{i=0}^{N-1} \sigma_i^x$, $H_2 = - \sum_{i=0}^{N-1} \sigma_i^z \sigma_{i+1 \bmod N}^z$.

References

- [1] E. Farhi, J. Goldstone, and S. Gutmann. *arXiv:1411.4028 [quant-ph]*, Nov. 2014.
- [2] W. W. Ho and T. H. Hsieh. *SciPost Phys.*, 6(3):29, 2019.
- [3] J. K. Pachos and Z. Papić. *SciPost Phys. Lect. Notes*, page 4, 2018.
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Results

We want to quantify how efficient QAOA is in preparing the ground state of (4) for $h_x, h_z \in [0, 3]$.

- We use $H'_1 = - \sum_{i=0}^{N-1} \sigma_i^z$, $H'_2 = H_1$, $H'_3 = H_2$.
- We simulate the quantum evolution step classically using $1 - f$ as the cost function, where $f = \langle \psi_{GS} | \psi(\theta) \rangle^2$ and ψ_{GS} is the ground state of H .
- `scipy.minimize.basinhopping` was used to avoid local minima.
- We computed $\mathcal{D}_{\mathcal{F}}$ across the phase diagram of (4) and found good correlation with $1 - f$ (Figures 2 and 3).

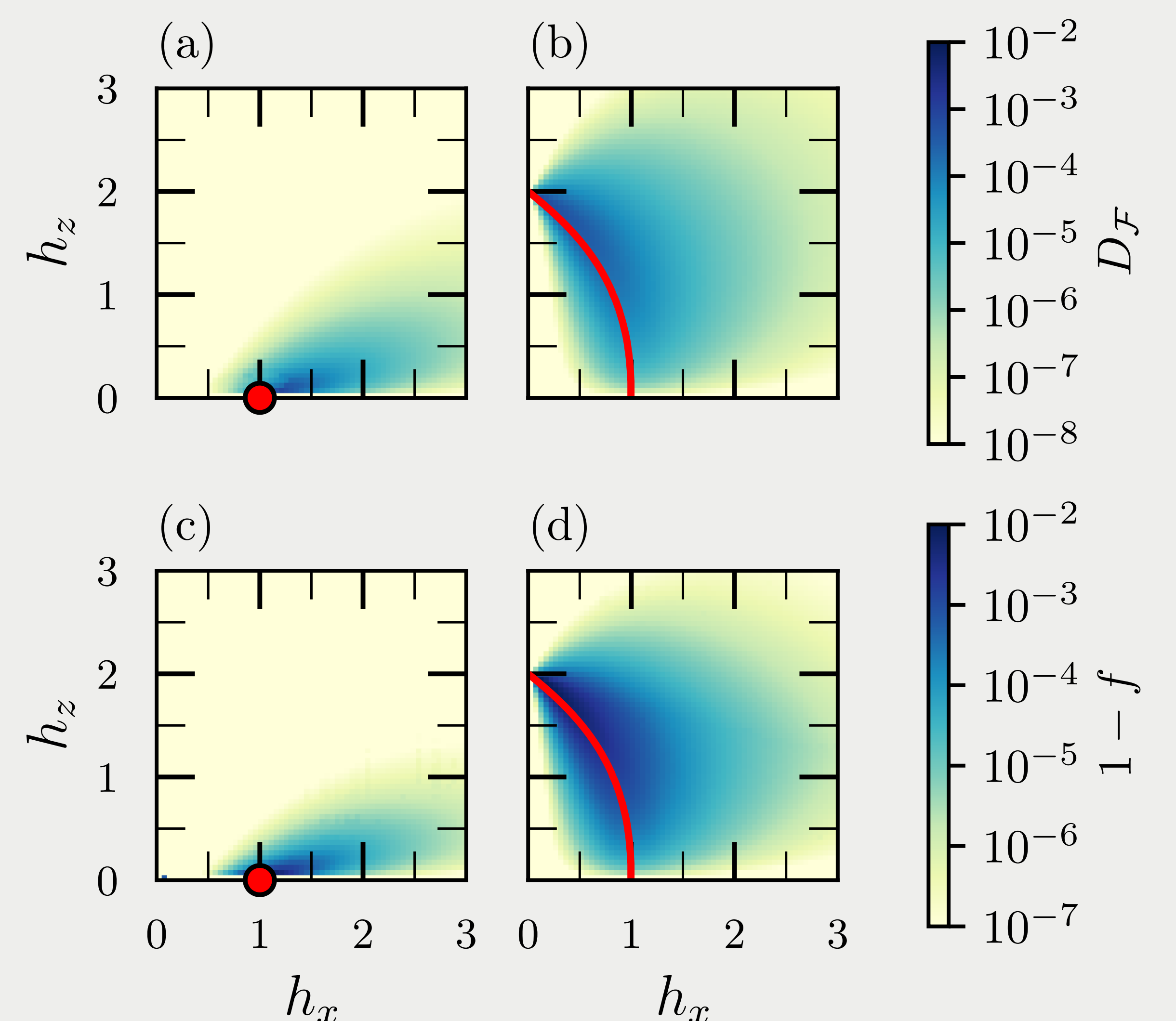


Figure 2: Ferromagnetic (left) and antiferromagnetic (right) Ising model.

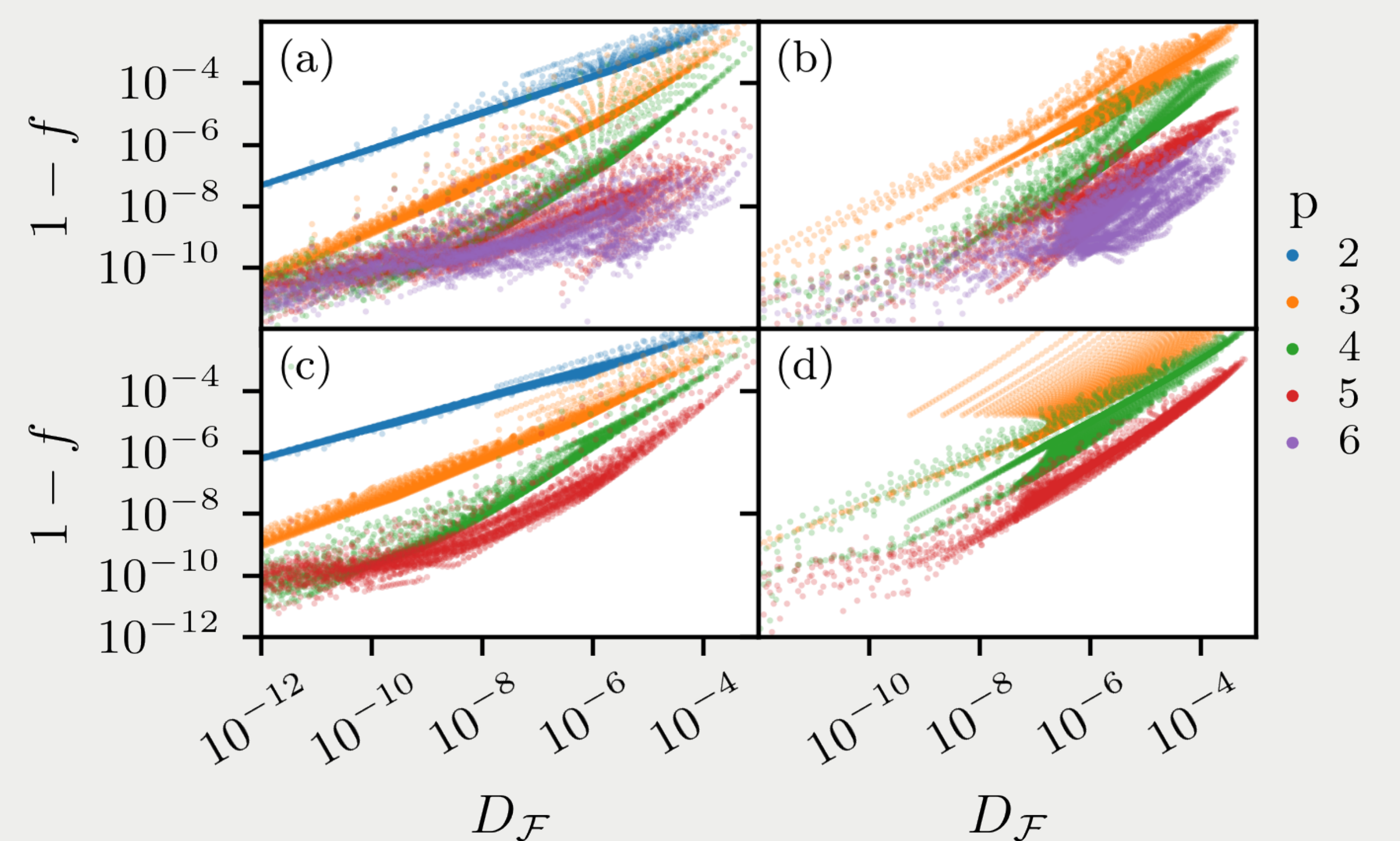


Figure 3: Correlation between $\mathcal{D}_{\mathcal{F}}$ and $1 - f$ in ferromagnetic (left) and antiferromagnetic (right) Ising model, for $N = 6$ (top) and $N = 8$ (bottom) for the data in Figure 2.

Analysis

- We have found a correlation between $\mathcal{D}_{\mathcal{F}}$ and QAOA around $p = \frac{N}{2}$ in the Ising model (4).
- This can be explained by the perturbatively small angles associated with H'_1 (ask me about this!).
- We relate how hard it is to prepare a ground state with how free it is, as measured by $\mathcal{D}_{\mathcal{F}}$
- This correlation is also present in the three-spin Ising model.
- This is a hard minimization problem. The initial guess and the classical simulation of the quantum evolution can be improved.

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