

Geometric Approaches to QFT

The Eisenhart Lift



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with Apostolos Pilaftsis and Sotirios Karamitsos

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arXiv:1806.02431 (PRD 2018)

arXiv:2012.15288 (PRD 2021)

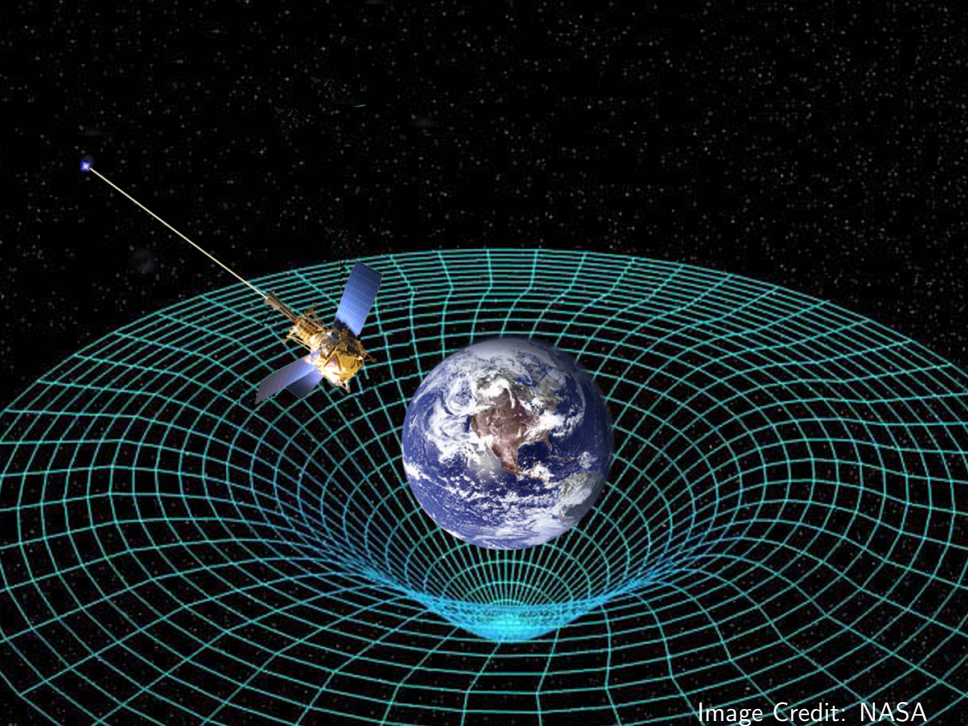


Image Credit: NASA

The Classical Eisenhart Lift

L. Eisenhart 1928

Particle Subject to Conservative Force

$$\blacksquare L = \frac{1}{2} m g_{ij} \dot{x}^i \dot{x}^j - V(\mathbf{x})$$

$$\blacksquare m \ddot{x}^i + \Gamma_{jk}^i \dot{x}^j \dot{x}^k = -V_{,i}$$

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Free Particle on Higher Dimensional Manifold

$$\blacksquare \text{Coordinates: } X^I = \{x^i, y\}$$

$$\blacksquare \text{Metric: } G_{IJ} = \begin{pmatrix} g_{ij} & 0 \\ 0 & \frac{M^2}{mV} \end{pmatrix}$$

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$$\blacksquare L = \frac{1}{2} m g_{ij} \dot{x}^i \dot{x}^j + \frac{1}{2} \frac{M^2}{V(\mathbf{x})} \dot{y}^2$$

$$\blacksquare m \ddot{x}^i + \Gamma_{jk}^i \dot{x}^j \dot{x}^k = -\frac{1}{2} \left(\frac{M \dot{y}}{V(\mathbf{x})} \right)^2 V_{,i}$$

$$\blacksquare \frac{d}{dt} \left(\frac{M \dot{y}}{V(\mathbf{x})} \right) = 0 \Rightarrow \frac{M \dot{y}}{V(\mathbf{x})} = A$$

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- $L = \frac{1}{2} m g_{ij} \dot{x}^i \dot{x}^j + \frac{1}{2} \frac{M^2}{V(\mathbf{x})} \dot{y}^2$

- $m\ddot{x}^i + \Gamma_{jk}^i \dot{x}^j \dot{x}^k = -\frac{A^2}{2} V_{,i}$

- $\frac{d}{dt} \left(\frac{M\dot{y}}{V(\mathbf{x})} \right) = 0 \Rightarrow \frac{M\dot{y}}{V(\mathbf{x})} = A$

Example: Simple Harmonic Motion

Metric:

$$G_{IJ} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{k}{m^5 x^2} \end{pmatrix}$$

The Eisenhart Lift in Quantum Mechanics

KF, S. Karamitsos and A. Pilaftsis 2020 (Arxiv:2012.15288)

Original System

- $H = \frac{g^{ij} p_i p_j}{2m} + V(\mathbf{x})$

- $\dot{x}^i = \frac{g^{ij} p_j}{m}$

- $\dot{p}_i = -mV_{,i}$

Lifted System

- $H = \frac{g^{ij} p_i p_j}{2m} + V(\mathbf{x}) \frac{p_y^2}{2M^2}$

- $\dot{x}^i = \frac{g^{ij} p_j}{m}$

- $\dot{p}_i = -mV_{,i} \frac{p_y^2}{2M^2}$

- $\frac{M\dot{y}}{V(\mathbf{x})} = \frac{p_y}{M} = A$

- $\dot{p}_y = 0$

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Poisson Brackets

- x^i, y, p_i, p_y
- $\{x^i, p_i\} = \{y, p_y\} = 1$
- $\{x^i, p_y\} = \{y, p_i\} = 0$
- $\{x^i, x^j\} = \{y, y\} = 0$
- $\{x^i, y\} = 0$
- $\{p_i, p_j\} = \{p_y, p_y\} = 0$
- $\{p_i, p_y\} = 0$



Commutators

- $\hat{x}^i, \hat{y}, \hat{p}_i, \hat{p}_y$
- $[\hat{x}^i, \hat{p}_i] = [\hat{y}, \hat{p}_y] = i$
- $[\hat{x}^i, \hat{p}_y] = [\hat{y}, \hat{p}_i] = 0$
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Original Hamiltonian

$$\hat{H} = \frac{1}{2m}\hat{p}_x^2 + V(\hat{x})$$

Original Schrödinger Equation

$$-\frac{1}{2m}\psi''(x) + V(x)\psi(x) = E\psi(x)$$

Lifted Hamiltonian

$$\hat{H} = \frac{1}{2m}\hat{p}_x^2 + \frac{V(\hat{x})}{2M^2}\hat{p}_y^2$$

Lifted Schrödinger Equation

$$-\frac{1}{2m}\Psi_{,xx}(\mathbf{x}, y) - \frac{V(\mathbf{x})}{2M^2}\Psi_{,yy}(\mathbf{x}, y) = E\Psi(\mathbf{x}, y)$$

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Lifted Hamiltonian

$$\hat{H} = \frac{1}{2m} \hat{p}_x^2 + \frac{V(\hat{x})}{2M^2} \hat{p}_y^2$$

Lifted Schrödinger Equation

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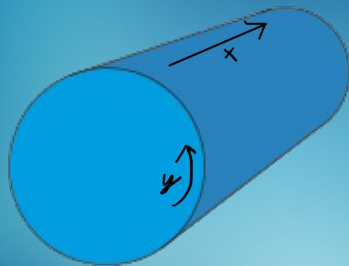
Separation of Variables

$$\Psi(x, y) = \psi(x)\chi(y)$$

$$\chi''(y) = -P^2\chi(y)$$

$$-\frac{1}{2m} \psi''(x) + \frac{P^2}{2M^2} V(x)\psi(x) = E\psi(x)$$

The Eisenhart Lift in Quantum Mechanics

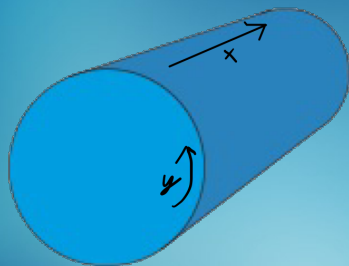


- Compactify fictitious dimension y so that $0 \leq y < \ell$

- $P = \frac{2\pi k}{\ell}, \quad k \in \mathbb{Z}$

- $-\frac{1}{2m}\psi''(x) + \frac{2\pi^2 k^2}{M^2 \ell^2} V(x)\psi(x) = E\psi(x)$

The Eisenhart Lift in Quantum Mechanics



- Compactify fictitious dimension y so that $0 \leq y < \ell$
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Example: Simple Harmonic Oscillator

- Hamiltonian: $\hat{H} = \frac{\hat{p}_x^2}{2m} + m\omega^2 \hat{x}^2 \frac{\hat{p}_y^2}{2M^2}$
- Energy Eigenstates: $\Psi(x, y) = \langle \mathbf{x} | k, n \rangle = \exp\left(\frac{2\pi i k y}{\ell}\right) \psi_n\left(\frac{x}{k}\right)$
- Energy Eigenvalues: $k\omega\left(n + \frac{1}{2}\right)$
- Orthonormality: $\langle k', n' | k, n \rangle = \delta_{k'k} \delta_{n'n}$

The Eisenhart Lift for Classical Field Theory

KF, S. Karamitsos and A. Pilaftsis 2018 (arXiv:1806.02431)

Original Theory

- $\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi)$
- $\square \phi = -V'(\phi)$

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Lifted Theory

- $\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{1}{2} \frac{M^4}{V(\phi)} \partial_\mu B^\mu \partial_\nu B^\nu$
- $\square \phi = -\frac{1}{2} \left(\frac{M^2 \partial_\nu B^\nu}{V(\phi)} \right)^2 V'(\phi)$
- $\partial_\mu \left(\frac{M^2 \partial_\nu B^\nu}{V(\phi)} \right) = 0 \Rightarrow \frac{M^2 \partial_\nu B^\nu}{V(\phi)} = A$

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Hamiltonian Formulation

Lifted Hamiltonian

$$\mathcal{H}_{\text{lift}} = \frac{\pi_\phi^2 + (\nabla\phi)^2}{2} + \frac{V(\phi)}{M^4} \pi_0^2 - \pi_0 \partial_i B^i + \pi_i \dot{B}^i$$

Hamilton's Equations

$$\dot{\phi} = \pi_\phi$$

$$\dot{\pi}_\phi = \nabla^2 \phi - \frac{1}{2M^4} \pi_0^2 V'(\phi)$$

$$\pi_0 = M^4 \frac{\partial_\mu B^\mu}{V(\phi)} = M^2 A$$

$$\partial_\mu \pi_0 = 0$$

$$\pi_i = 0$$

The Eisenhart Lift for Quantum Field Theory

Lifted Hamiltonian

$$\widehat{\mathcal{H}}_{\text{lift}} = \frac{\widehat{\pi}_\phi^2 + (\nabla \widehat{\phi})^2}{2} + \frac{m^2 \widehat{\phi}^2}{2M^4} \widehat{\pi}_0^2 - \widehat{\pi}_0 \partial_i \widehat{B}^i + \widehat{\pi}_i \dot{\widehat{B}}^i$$

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Creation and Annihilation Operators

$$a_{\mathbf{k}}, a_{\mathbf{k}}^\dagger = \int d^3\mathbf{x} e^{i\mathbf{k}\cdot\mathbf{x}} \left[\frac{\hat{\pi}_0(\mathbf{x})}{\sqrt{2M^2}} \sqrt{\frac{\omega_{\mathbf{k}}}{2}} \hat{\phi}(\mathbf{x}) \pm \frac{i}{\sqrt{2\omega_{\mathbf{k}}}} \hat{\pi}_\phi(\mathbf{x}) \right]$$

Commutation Relations

- $[a_{\mathbf{k}}, a_{\mathbf{q}}^\dagger] = \frac{\hat{\pi}_0}{\sqrt{2M^2}} (2\pi)^3 \delta^{(3)}(\mathbf{k} - \mathbf{q})$
- $[\hat{H}_{\text{lift}}, a_{\mathbf{k}}^\dagger] = \frac{\hat{\pi}_0}{\sqrt{2M^2}} \omega_{\mathbf{k}} a_{\mathbf{k}}^\dagger$
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Conserved Momentum

- $[\hat{\mathcal{H}}_{\text{lift}}, \hat{\pi}_0] = 0$
- $[\hat{P}_i, \hat{\pi}_0] = 0$
- $\therefore \hat{\pi}_0$ conserved in time and space

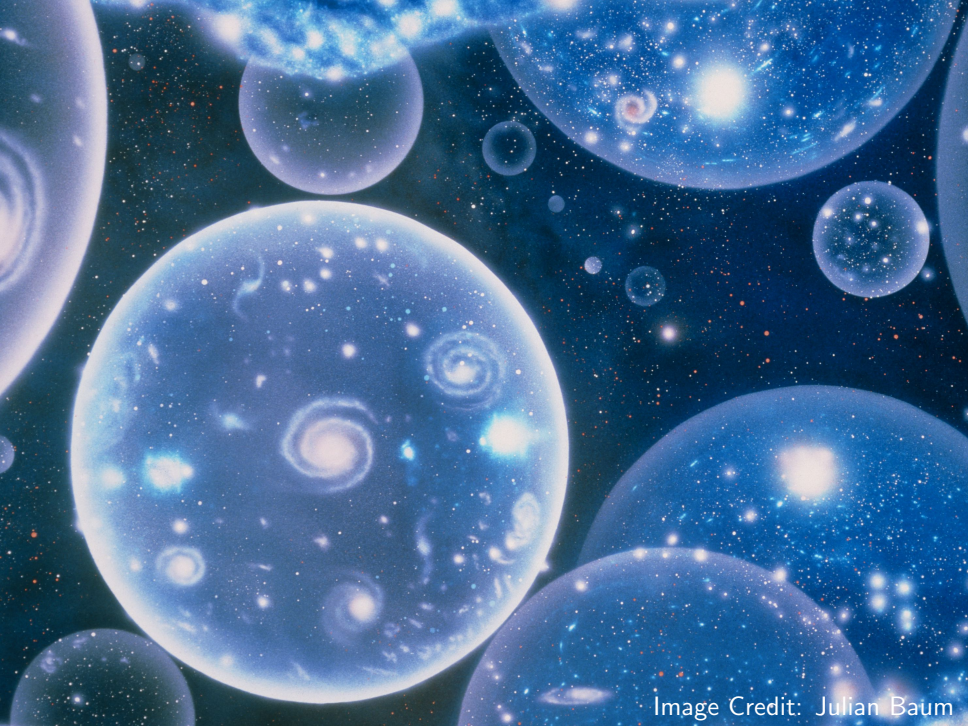


Image Credit: Julian Baum



Possible Applications

- Anthropic explanation of the cosmological constant problem
- Gauge hierarchy problem
- The initial conditions of inflation (see [arXiv:1812.07095](https://arxiv.org/abs/1812.07095))

Summary

- A conservative force can be described by the curvature of a higher dimension
- The potential term of a quantum field theory can be described as the kinetic term of a vector field
- Different values of the conjugate momentum π_0 lead to an ensemble of universes with different physics
- May have applications to cosmological constant, gauge hierarchy and measure problem