

# Searching for Dark Photon DM with GW Detectors, O1 and beyond

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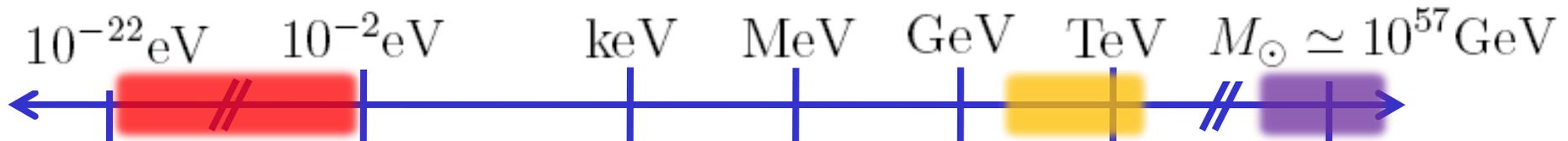
Aaron Pierce, Keith Riles, Yue Zhao  
arXiv:1801.10161 [hep-ph]

Phys.Rev.Lett. 121 (2018) no.6, 061102

Huaike Guo, Keith Riles, F.W. Y., Yue Zhao  
arXiv:1905.04316 [hep-ph]

Nature - Commun.Phys. 2 (2019) 155

# Popular Choices:



- Very light DM particles

Axion and Dark “Photon”

$$10^{-22} \text{ eV} \sim 10^{-2} \text{ eV}$$

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- WIMPs:

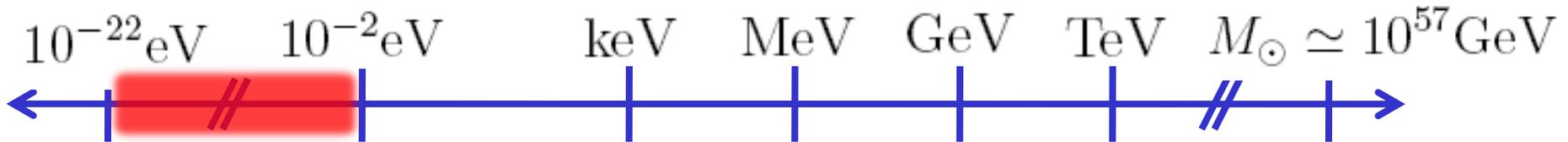
$$100 \text{ GeV} \sim \text{TeV}$$

- Primordial Black Holes:

$$10^{-7} \sim 100 \text{ solar mass}$$

Both ultra-light and ultra-heavy scenarios  
can be proved by GW detectors!

# Popular Choices:



- Very light DM particles

Axion and **Dark “Photon”**

$10^{-22}$  eV  $\sim$   $10^{-2}$  eV

gauge boson of the

**U(1)<sub>B</sub>** or **U(1)<sub>B-L</sub>**

(p+n)

(n)

DM is an oscillating background field.

Dark Photon is dominantly oscillating background dark electric field.

Driving displacements for particles charged under dark gauge group.

# Ultra-light DM – Dark Photon

- Mass

W/Z bosons get masses through the Higgs mechanism.

A dark photon can also get a mass by a dark Higgs,  
or through the **Stueckelberg mechanism.**

  
a special limit of the Higgs mechanism  
unique for U(1) gauge group

- Relic abundance (non-thermal production )

Misalignment mechanism

Light scalar decay

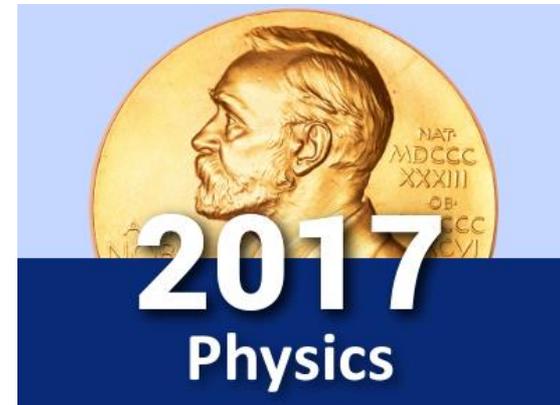
Production from cosmic string

Ultra-light dark photon can be a good candidate of cold dark matter!

# Laser Interferometer Gravitational-Wave Observatory

LIGO (ground-based)

Amazing precision at LIGO:  
 $O(1/1000)$  the radius of a single proton!



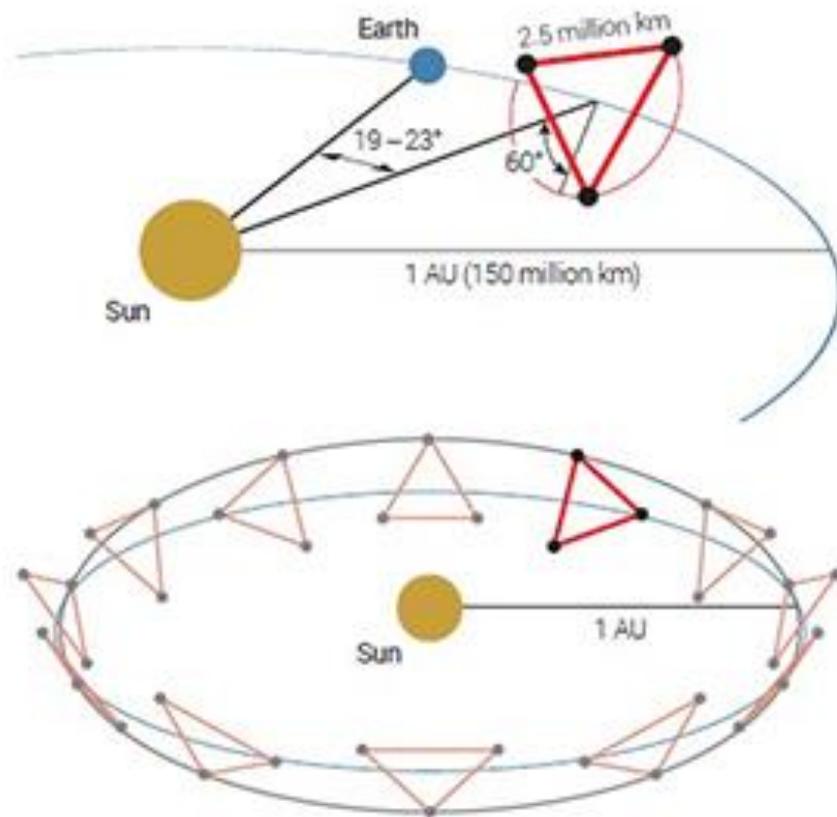
Opened a field:  
Gravitational Wave Astronomy

Enrich our understanding on  
fundamental physics and early  
cosmology.



# Laser Interferometer Space Antenna

LISA (space-based)



Recently approved by the European Space Agency.

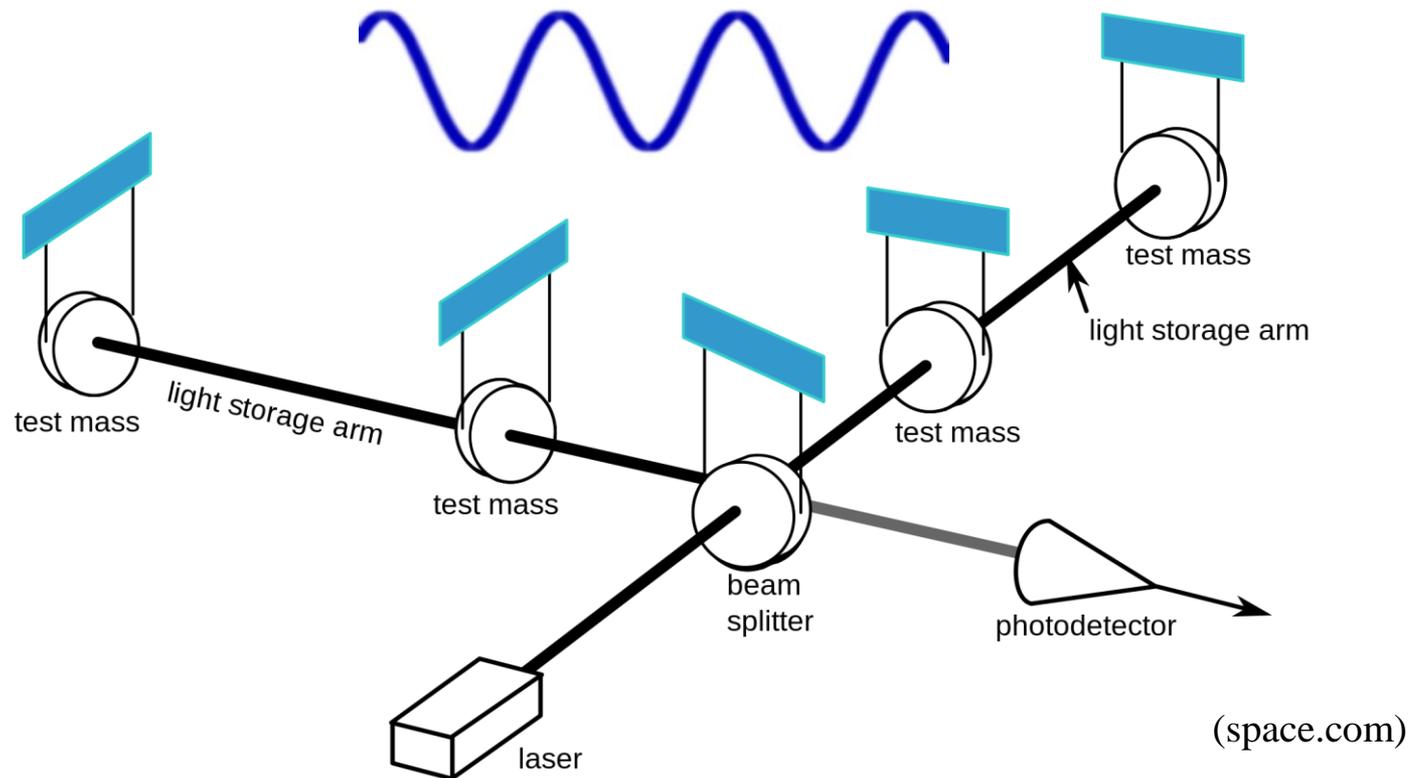
U.S. (NASA) just rejoined the program.

LISA PathFinder is a great success!

(LISA Mission Consortium)

# General Picture:

## LIGO/LISA: advanced Michelson–Morley interferometer

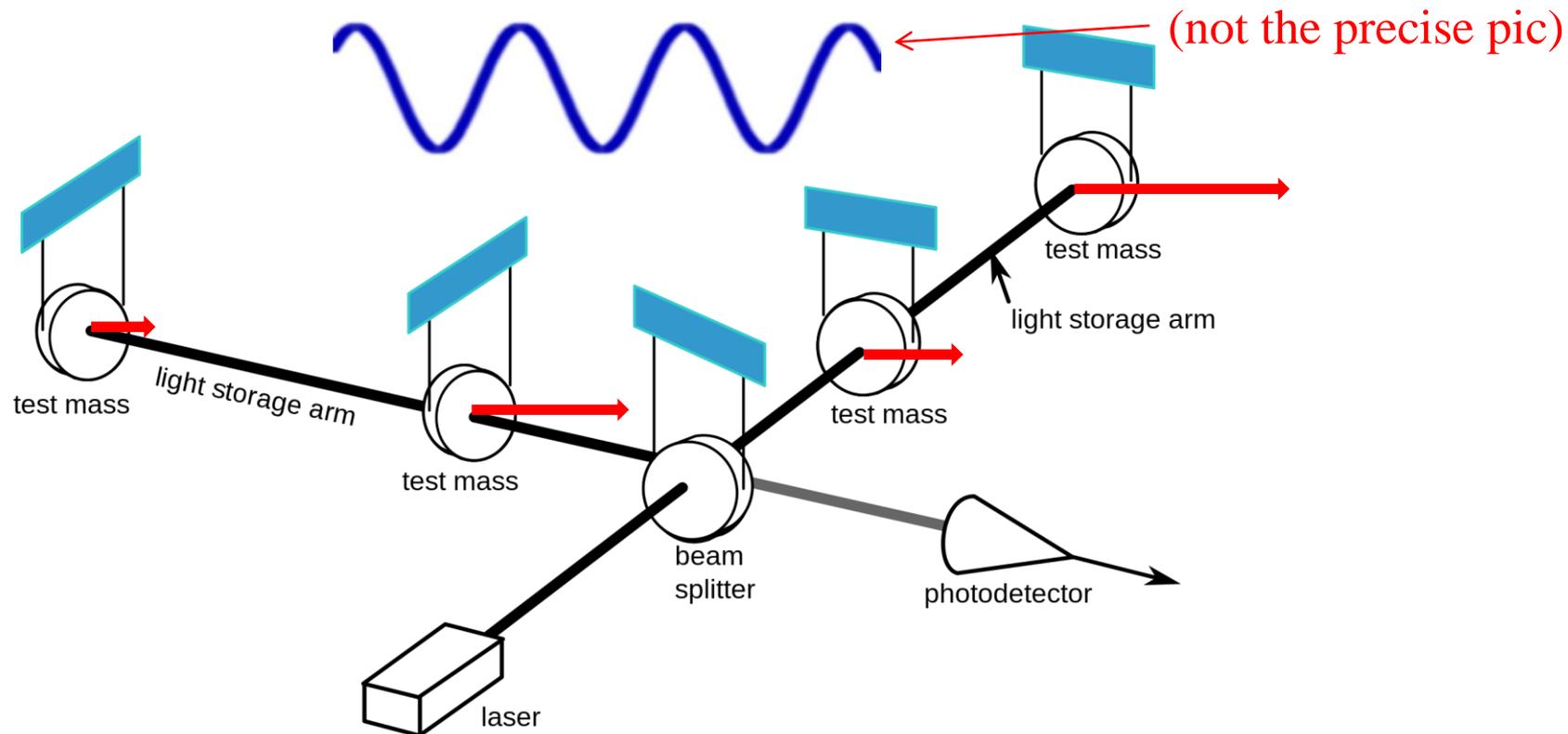


Gravitational wave changes the distance between mirrors.

⇒ Change photon propagation time between mirrors. ⇒ interferometer pattern

# General Picture:

Ultra-light DM  $\Rightarrow$  classical oscillating background field



Dark photon dark matter moves mirrors.  $\Rightarrow$  Change photon propagation time between mirrors.  $\Rightarrow$  interferometer pattern

# Maximal Displacement:

Local DM energy density:

$$\frac{1}{2} m_A^2 A_{\mu,0} A_0^\mu \simeq 0.4 \text{ GeV/cm}^3$$

local field strength of DP

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$\partial^\mu A_\mu = 0$$

$$E_i \sim m_A A_i \quad \gg \quad B^i \sim m_A v_j A_k \epsilon^{ijk}$$

# Maximal Displacement:

$$\vec{a}_i(t) = \frac{\vec{F}_i(t)}{M_i} \simeq \underbrace{\epsilon e}_{\text{dark photon coupling}} \underbrace{\frac{q_{D,i}}{M_i}}_{\text{charge mass ratio of the test object}} \underbrace{\partial_t \vec{A}(t, \vec{x}_i)}_{\text{dark electric field}}$$

charge mass ratio of the test object

Silicon mirror:

$$U(1)_B : 1/\text{GeV}$$

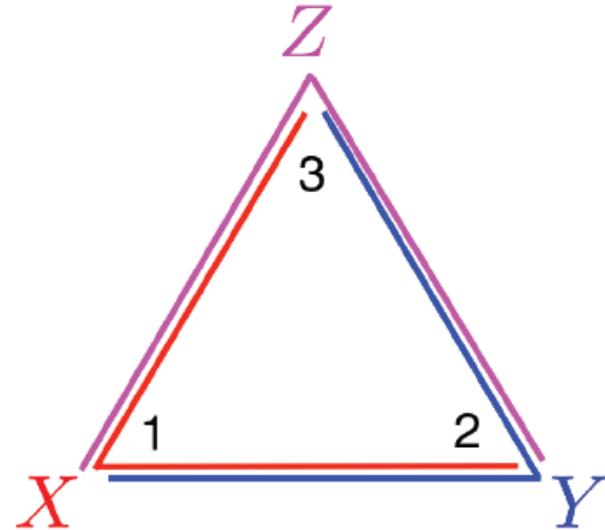
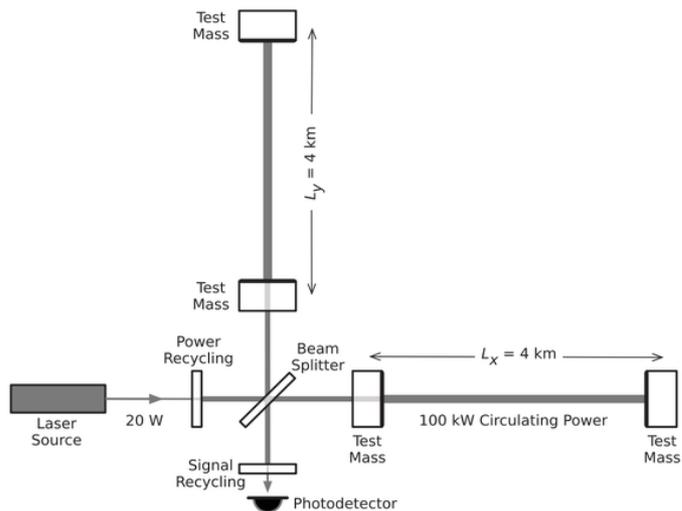
$$U(1)_{B-L} : 1/(2\text{GeV})$$

$$\Delta s_{\parallel,i} = \int dt \int dt a_{\parallel,i}(t)$$

projected along the arm direction

# Maximal GW-like Displacement:

$$\Delta L[t] = (x_1[t] - x_2[t]) - (y_1[t] - y_2[t])$$



$$\sqrt{\langle \Delta L^2 \rangle}_{LIGO} |_{max} = \frac{\sqrt{2}}{3} \frac{|a||k|L}{m_A^2}$$

Averaging on directions of acceleration and momentum vectors.

$$\sqrt{\langle \Delta L^2 \rangle}_{LISA} |_{max} = \frac{1}{\sqrt{6}} \frac{|a||k|L}{m_A^2}$$

$v_{vir} = 0$  gives same force to all test objects, not observable.  
Net effect is proportional to velocity.

# Properties of DPDM Signals:

Signal:

- almost monochromatic

$$f \simeq \frac{m_A}{2\pi}$$

- very long coherence time

$$\Delta f / f = v_{vir}^2 \simeq 10^{-6}$$

DM velocity dispersion.  
Determined by gravitational  
potential of our galaxy.

⇒ A bump hunting search in frequency space.

Can be further refined as a detailed template search,  
assuming Boltzmann distribution for DM velocity.

Once measured, we know great details of the local DM properties!

# Properties of DPDM Signals:

Signal:

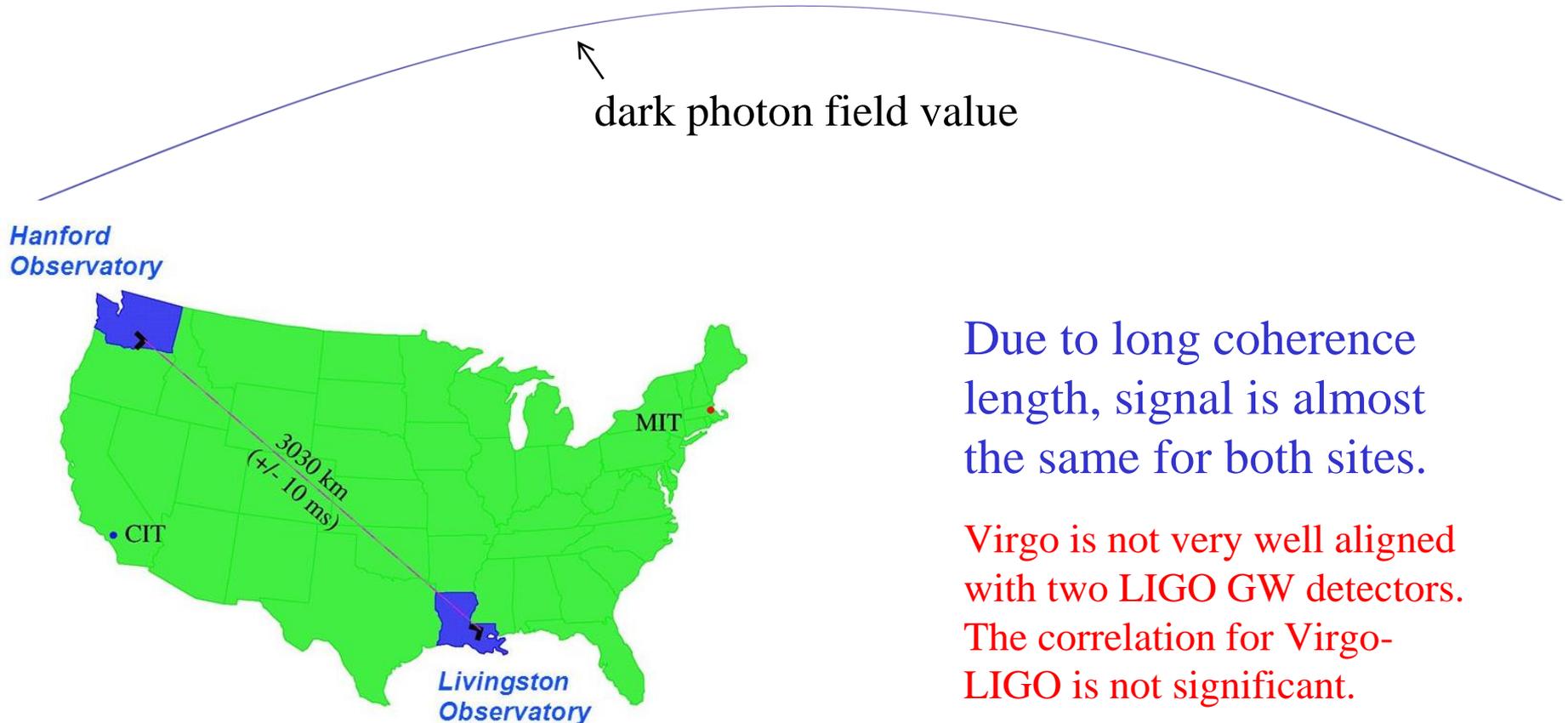
- very long coherent distance

$$l_{coh} \simeq \frac{1}{m_A v_{vir}} \simeq 3 \times 10^9 \text{m} \left( \frac{100 \text{Hz}}{f} \right)$$

Within the coherent distance, the detectors will experience almost the same DPDM background field.

# Properties of DPDM Signals:

Correlation between two sites is important to reduce background!



# Sensitivity to DPDM signal of GW detectors:

First we estimate the sensitivity in terms of **GW strain**.

(Allen & Romano, Phys.Rev.D59:102001,1999)

One-sided power spectrum function:

later map to  $\Delta L/L$

$$S_{GW}(f) = \frac{3H_0^2}{2\pi^2} f^{-3} \Omega_{GW}(f)$$

energy density carried by  
a GW planewave  $\rho_{GW}(f) = \frac{\langle \dot{h}^2 \rangle}{16\pi G}$

$$\Omega_{GW}(f) \equiv \frac{f}{\rho_c} \frac{d\rho_{GW}}{df} = \frac{f}{\rho_c} \frac{\rho_{GW}(f)}{\Delta f}$$

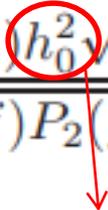
$$\Delta f/f = v_{vir}^2 \simeq 10^{-6}$$

Concretely predicted by  
Maxwell–Boltzmann distribution!

A template search is possible,  
and a better reach is expected!

# Sensitivity to DPDM signal of GW detectors:

Translate strain sensitivity to parameters of DPDM:

$$\text{SNR} = \frac{\gamma(|f|)h_0^2\sqrt{T}}{2\sqrt{P_1(f)P_2(f)\Delta f}}.$$


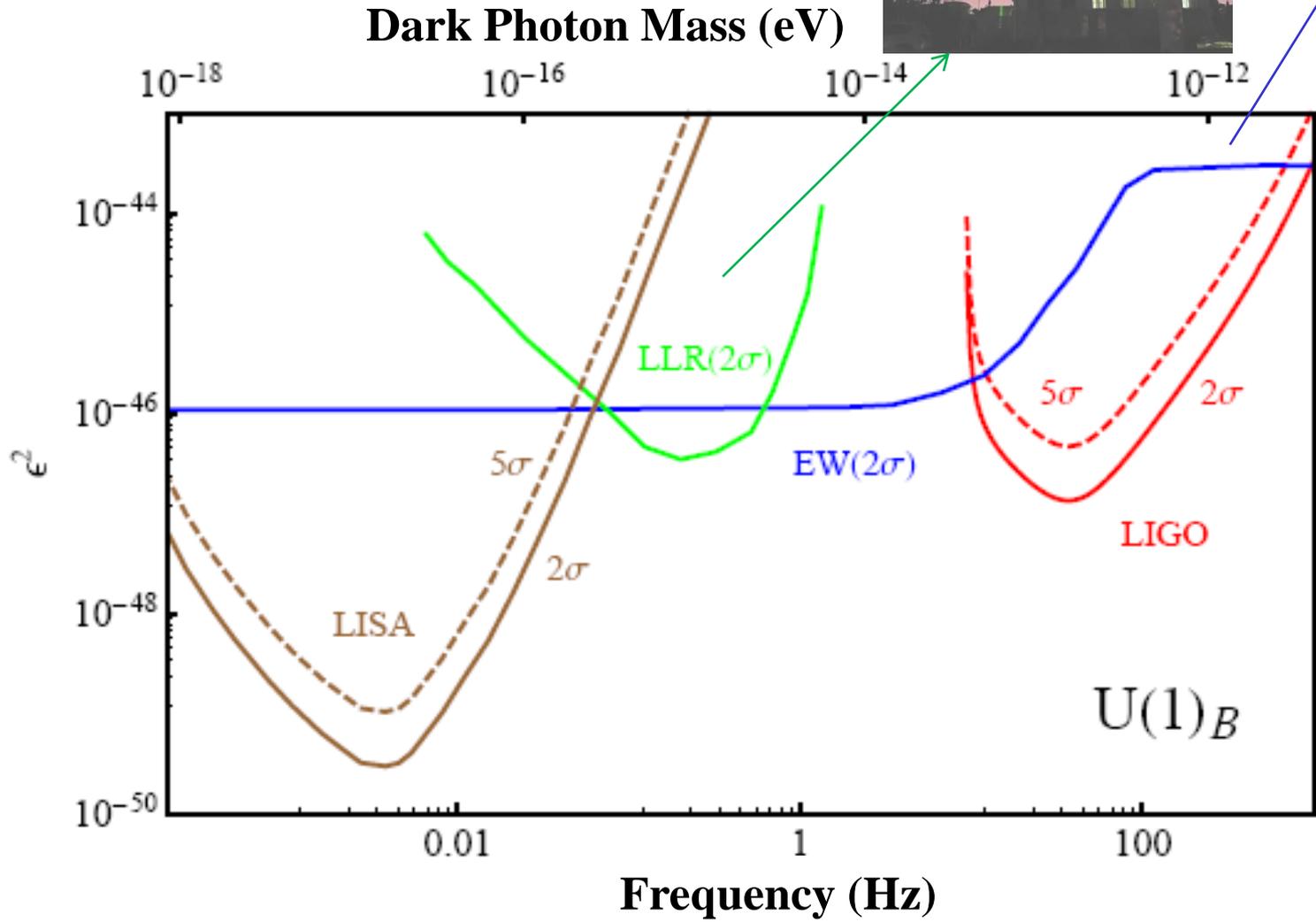
effectively the max differential displacement of two arms

a GW with strain  $h$   $\Rightarrow$  change of relative displacement as  $h$

$\Rightarrow \sqrt{\langle \Delta L^2 \rangle}_{LIGO|_{max}}$

$\Rightarrow$  sensitivity of DPDM parameters (mass, coupling)

# Sensitivity Plot:



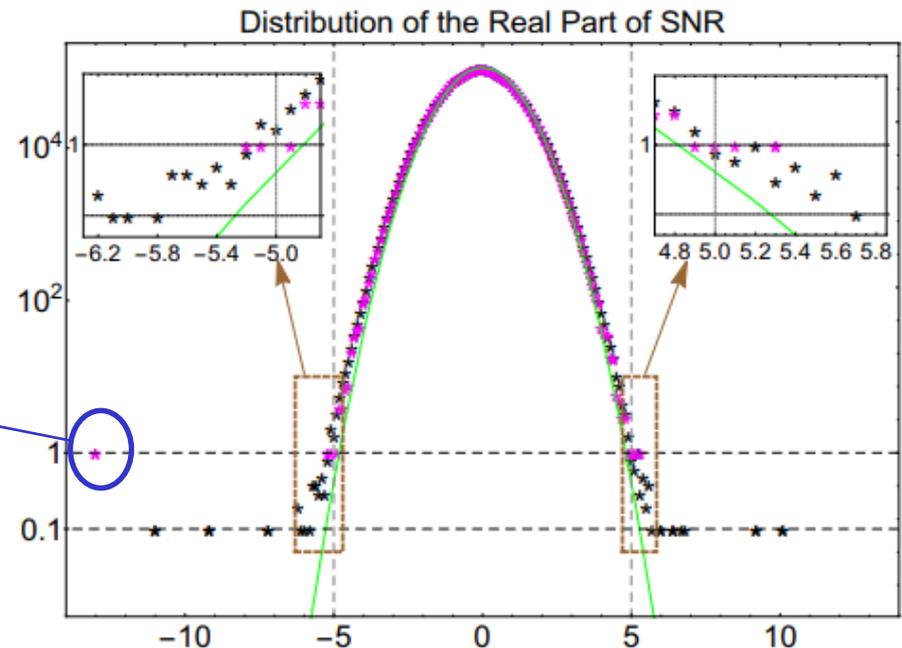
(Eöt-Wash web)  
 Loránd Eötvös  
 → Eöt-Wash

design sensitivities, 2 yrs

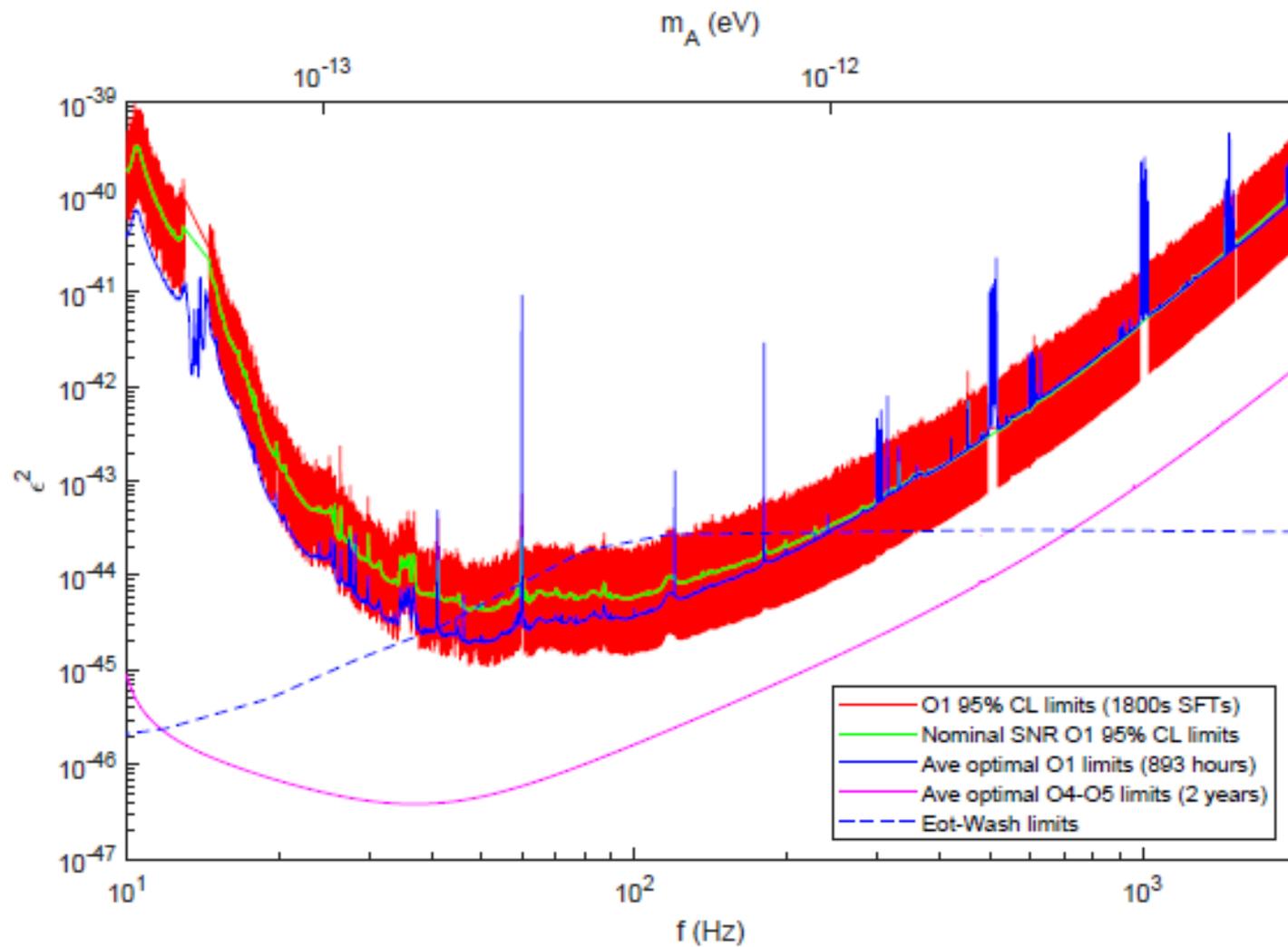
# O1 Result:

- 1800s FT: optimized for a signal at  $f \sim 500$  Hz
- Remove known noise bins and their neighbor bins
- Within 10-2000 Hz frequency band, require  $\text{Re}[\text{SNR}] < -5.8$   
 $\sim 1\%$  false alarm probability after including trial factors.
- Frequency lags: to deal with non-Gaussian noise  
offset bins ( $-50, -40, \dots, -10, +10, \dots, +50$ )  
Remove single interferometer artifacts and broadband correlated artifacts

known continuous wave  
“hardware injections”  
with random phase

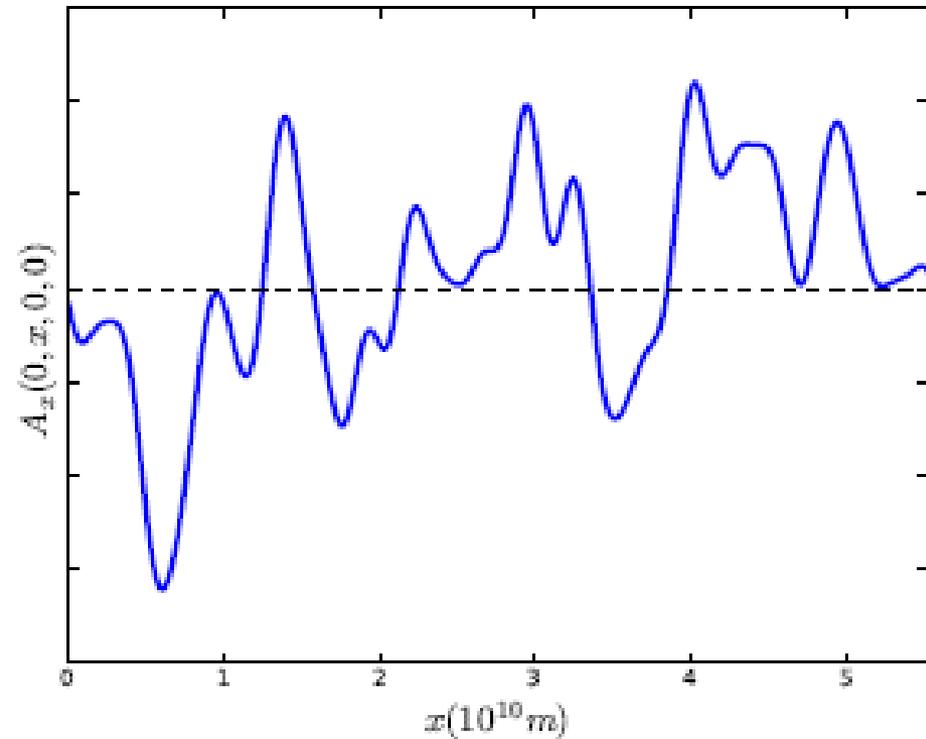
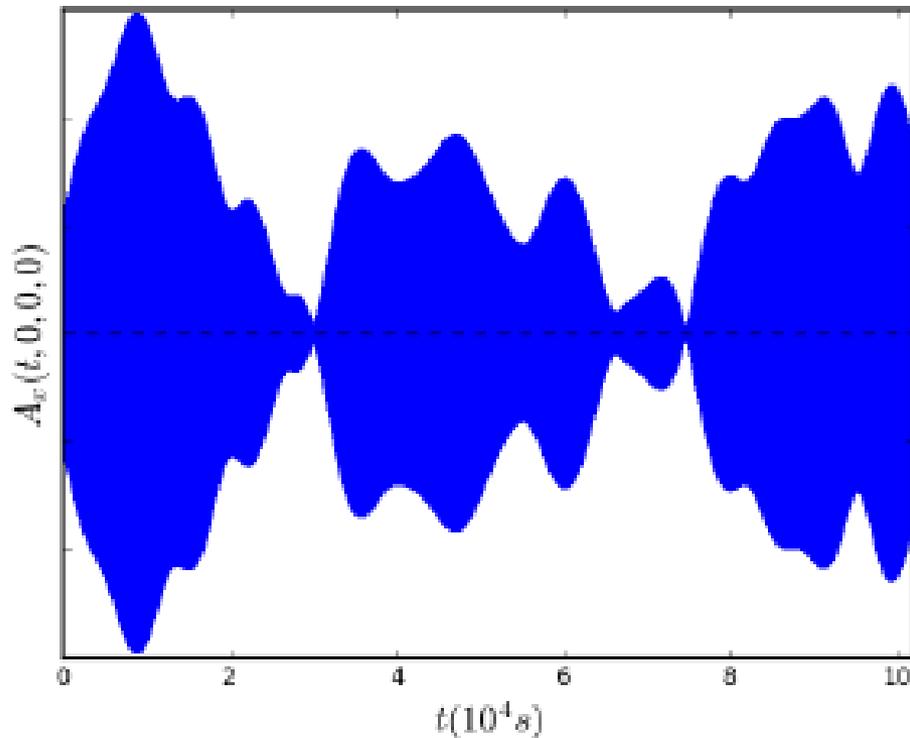


# O1 Result:

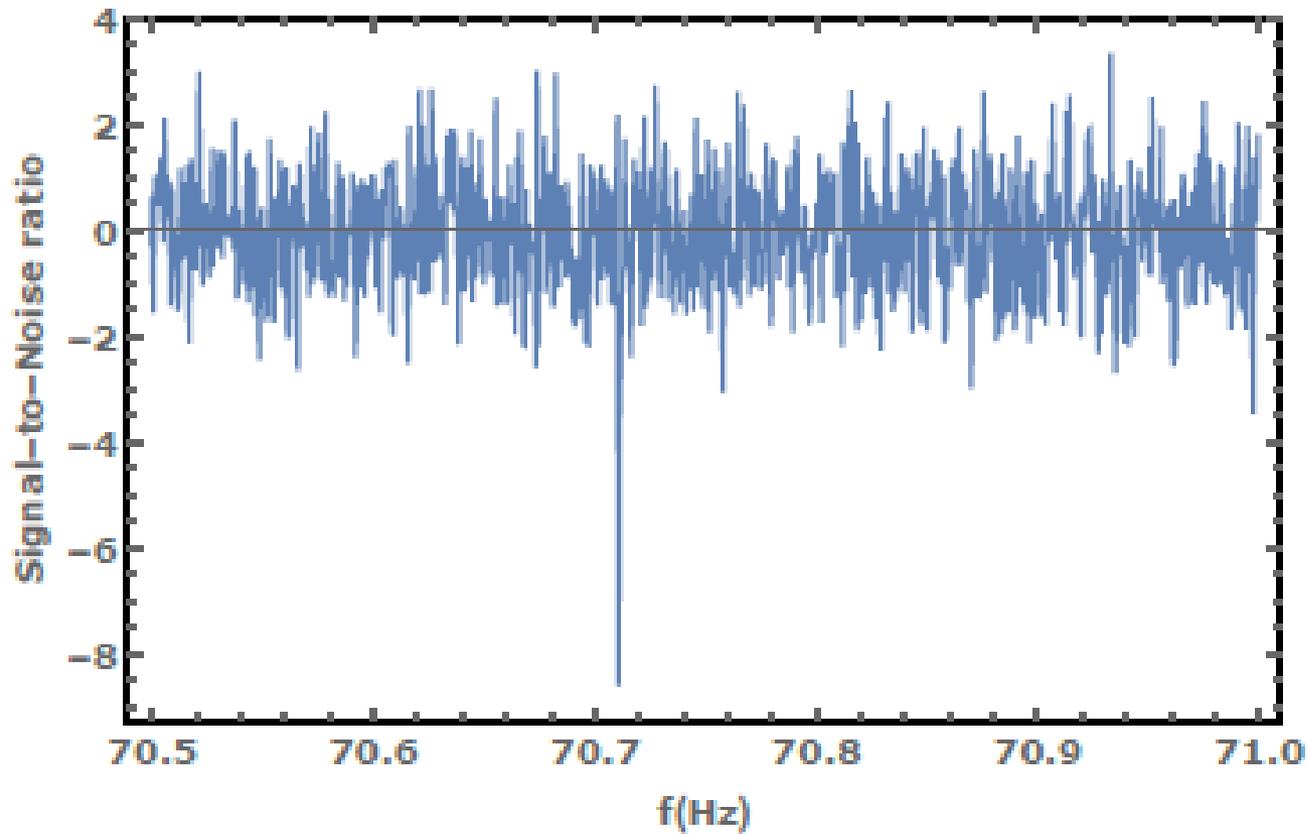


## Modeling DPDM background:

$$\vec{A}_{total}(t, \mathbf{x}) = \sum_{i=1}^N \vec{A}_{i,0} \sin(\omega_i t - \vec{k}_i \cdot \vec{x} + \phi_i)$$



# LIGO simulation output:

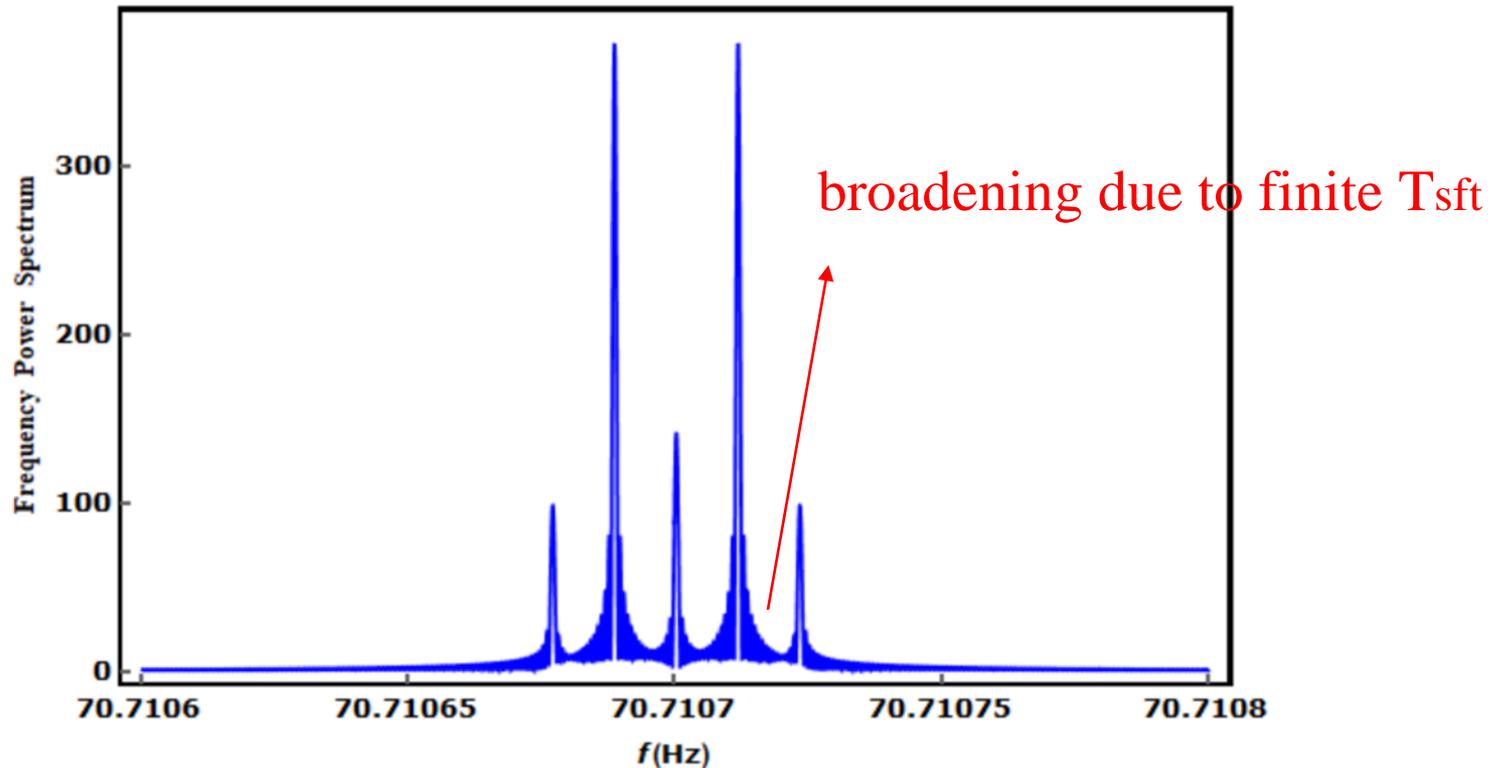


$$\epsilon^2 = 5 \times 10^{-44}, \quad f = 70.71 \text{ Hz} \quad T_{\text{SFT}} = 1800 \text{ s} \quad T_{\text{tot}} = 200 \text{ hr}$$

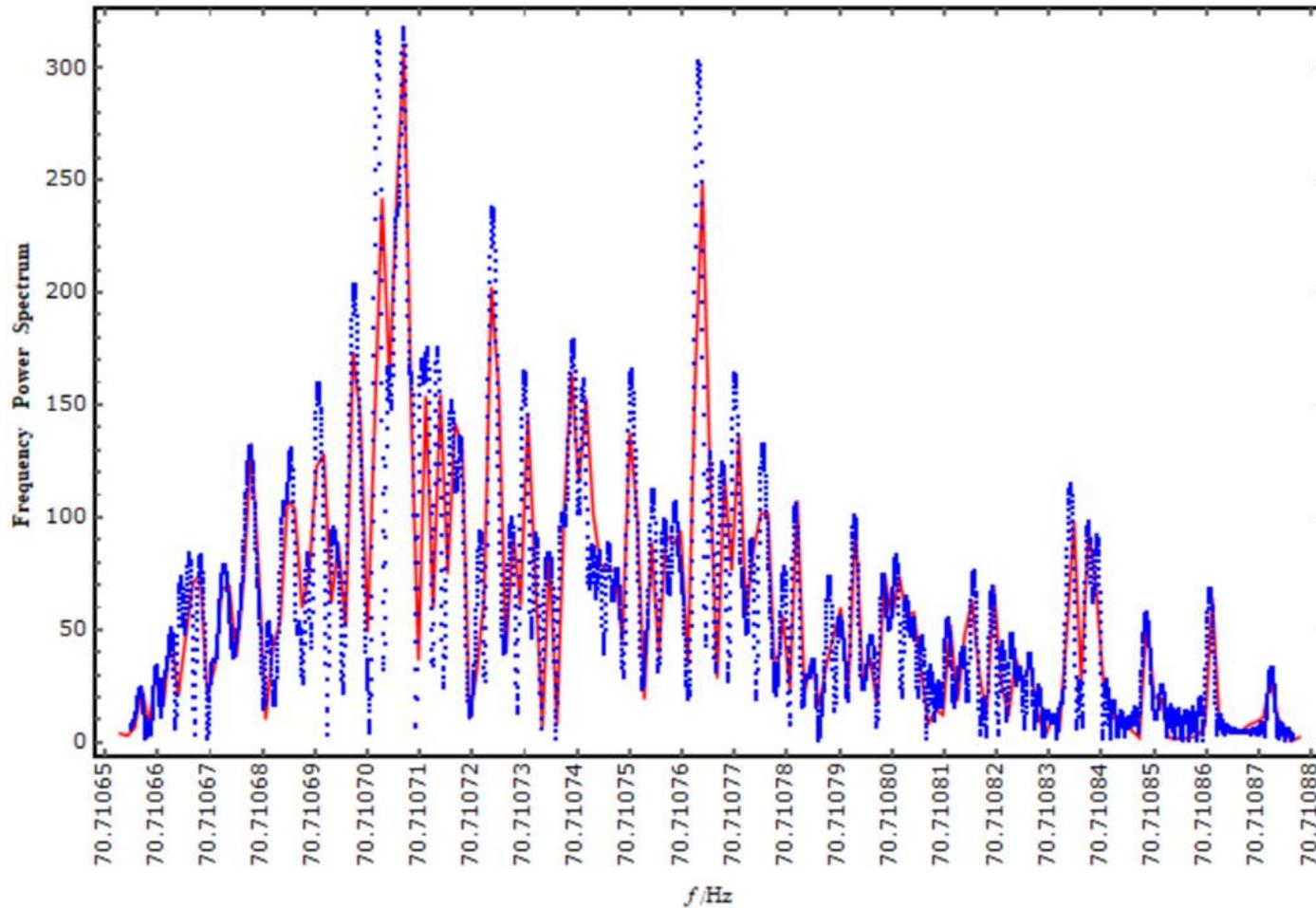
⇒ SNR  $\simeq -8$ ,

# Earth Rotation Effects:

$$R_L \approx - \sum_{i=1}^n \frac{\cos(\omega_i t + \Phi_i)}{\omega_i^2} \left( C_{2,1}^i \cos(2\omega_{Et}) + C_{2,2}^i \sin(2\omega_{Et}) + C_{1,1}^i \cos(\omega_{Et}) + C_{1,2}^i \sin(\omega_{Et}) + C_0^i \right)$$



# Fine structure of the signal:



Analytic understanding matches very well with numerical result!

# Conclusion

The applications of GW experiments can be extended!

⇒ Particularly sensitive to relative displacements.

Coherently oscillating DPDM generates such displacements.

It can be used as a DM direct detection experiment.

The analysis is straightforward!

⇒ Very similar to stochastic GW searches.

Better coherence between separated interferometers than Stochastic GW background.

The sensitivity can be extraordinary!

⇒ O1 data has already beaten existing experimental constraints.

Can achieve 5-sigma discovery at unexplored parameter regimes.

Once measured, great amount of DM information can be extracted!