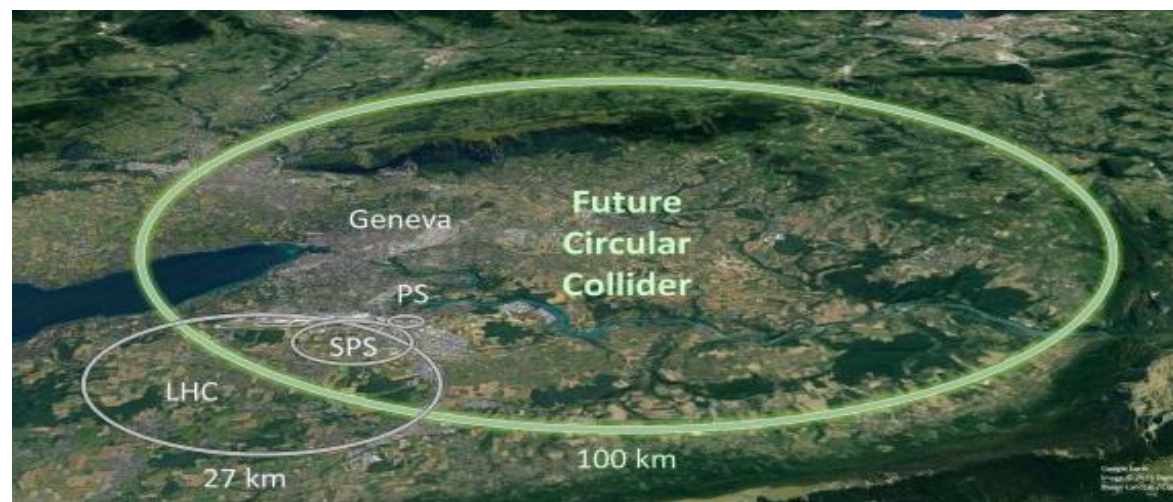


1. New Future Circular Hadron Collider (FCC-hh) Experiments

Possible location for the FCC-hh



- Probing of Higgs-Boson properties, searches for physics beyond the Standard Model
- 100 km circumference with 16 T dipole magnets
- Instantaneous luminosities: $30 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- Integrated luminosity: 20 ab^{-1}
- Center-of-mass energy: $\sqrt{s} = 100 \text{ TeV}$.

The Future Circular Collider (FCC) study is developing designs for a new research infrastructure to host the next generation of higher performance particle colliders to extend the research currently being conducted at the LHC, once the High-Luminosity phase (HL-LHC) reaches its conclusion in 2040 [1].

A muon system will surround the solenoid. It will consist of four parts:

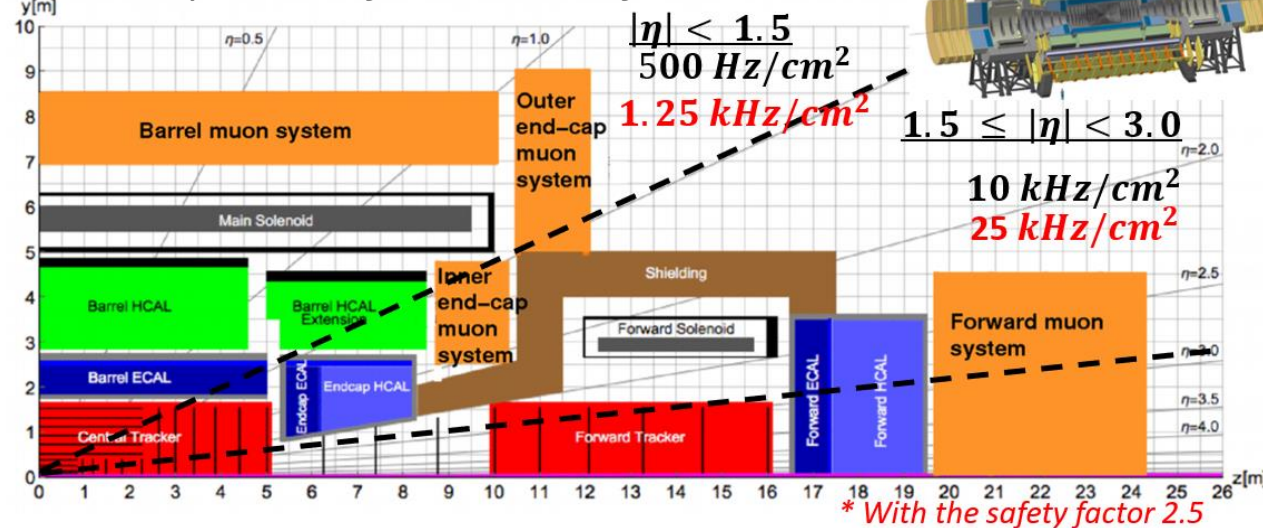
- $|\eta| < 1.0$: barrel;
- $1.0 \leq |\eta| < 1.5$: outer endcap;
- $1.5 \leq |\eta| < 2.2$: inner endcap;
- $2.2 \leq |\eta| < 3.0$: forward.

- The **Muon System** will have to provide a muon trigger.

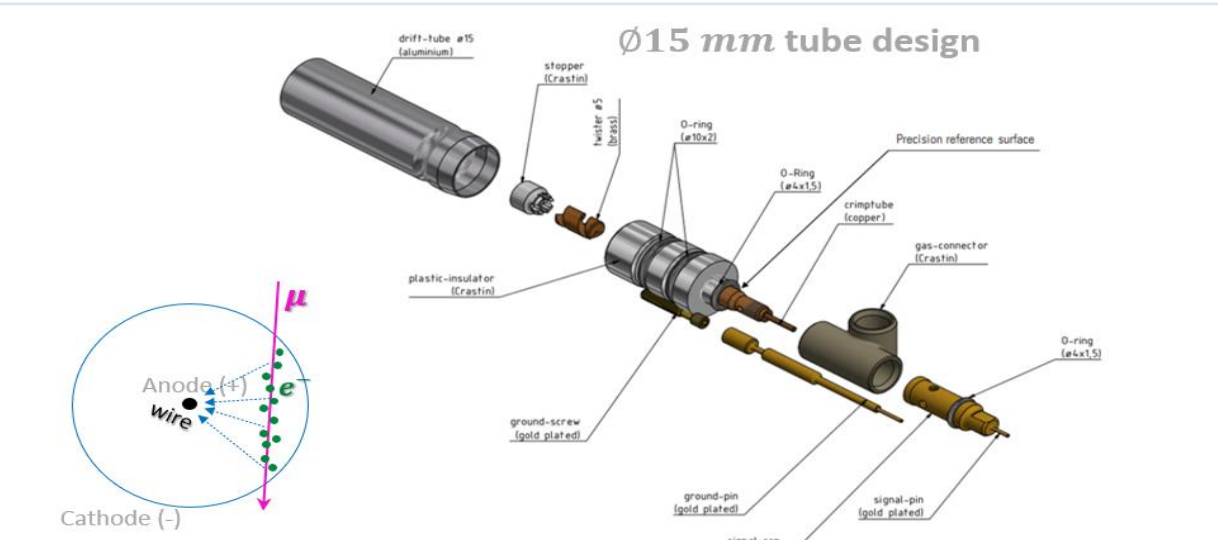
- Deflection angle α measurement of muon p_T .

- An angular resolution of $70 \mu\text{rad}$ is required for a high momentum resolution up to $p_T = 1 \text{ TeV}$.

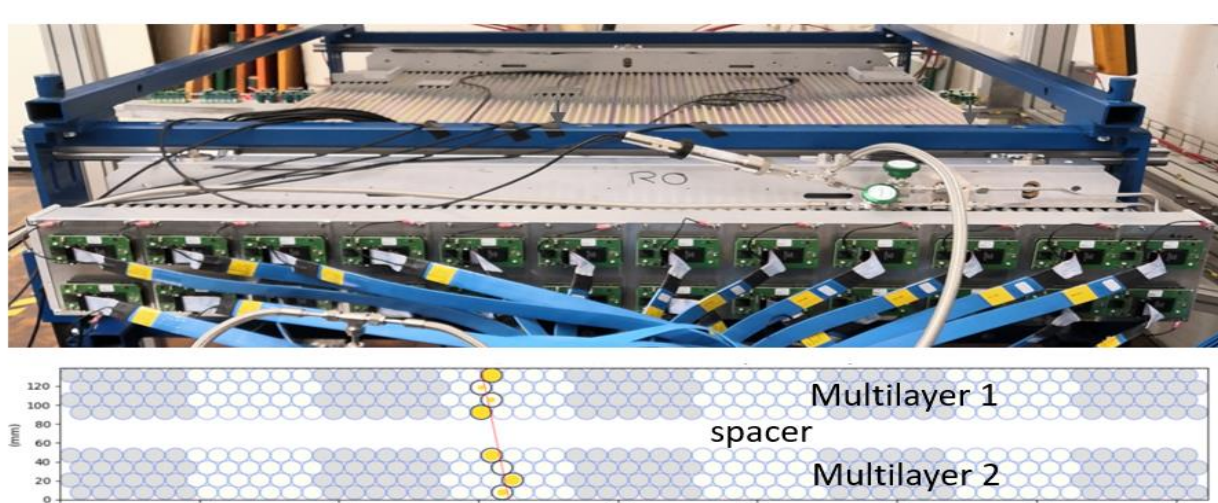
One quadrant of the FCC-hh reference detector



3. New small Drift Tube Detector Technology



New sMDT for the ATLAS Muon System upgrade



- Faster Max. Drift times $\sim 175 \text{ ns}$.
- High mechanical accuracy ($\sim 5 \mu\text{m}$ sense wire positioning accuracy).
- Spatial resolution ($< 40 \mu\text{m}$) over large areas.
- 10 times high-rate capability.
- Aging effect measured and didn't observe with Ar:CO_2 (93:7) drift gas at 3 bar up to 9 C/cm charge accumulation on wire.

Properties	New sMDT
Tube diameter, Wall thickness	15 mm, 400 μm
Anode wire diameter	50 μm
Number of tube layers	8
Operating Gas Mixture	Ar: CO ₂ (93:7)
Operating Pressure	3 bar
Operating HV working point	2730 V
Gas gain	2×10^4
Max. Drift time	$\sim 175 \text{ ns}$
Single tube Space resolution at 500 Hz/cm ² background rate	$110 \pm 2 \mu\text{m}$

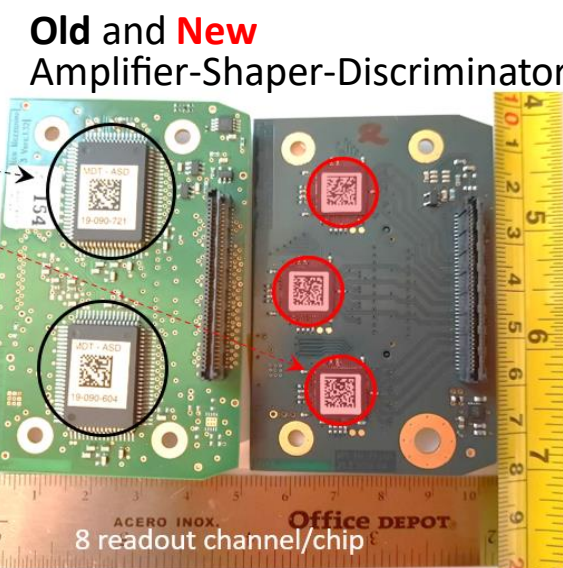
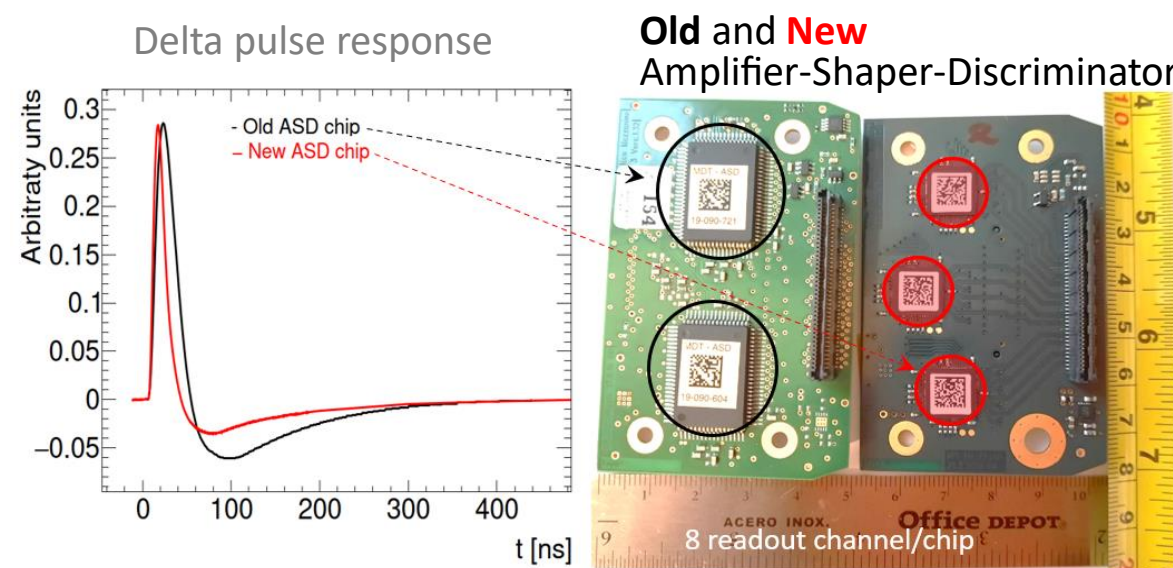
New sMDT stations already installed in the ATLAS Muon System during LS2 2019-2022



Design, materials and construction method optimised for large scale production. It allows for high-precision monolithic two-multilayer detectors for precise track angle measurement [3].

5. Improved pulse response of the new readout chip

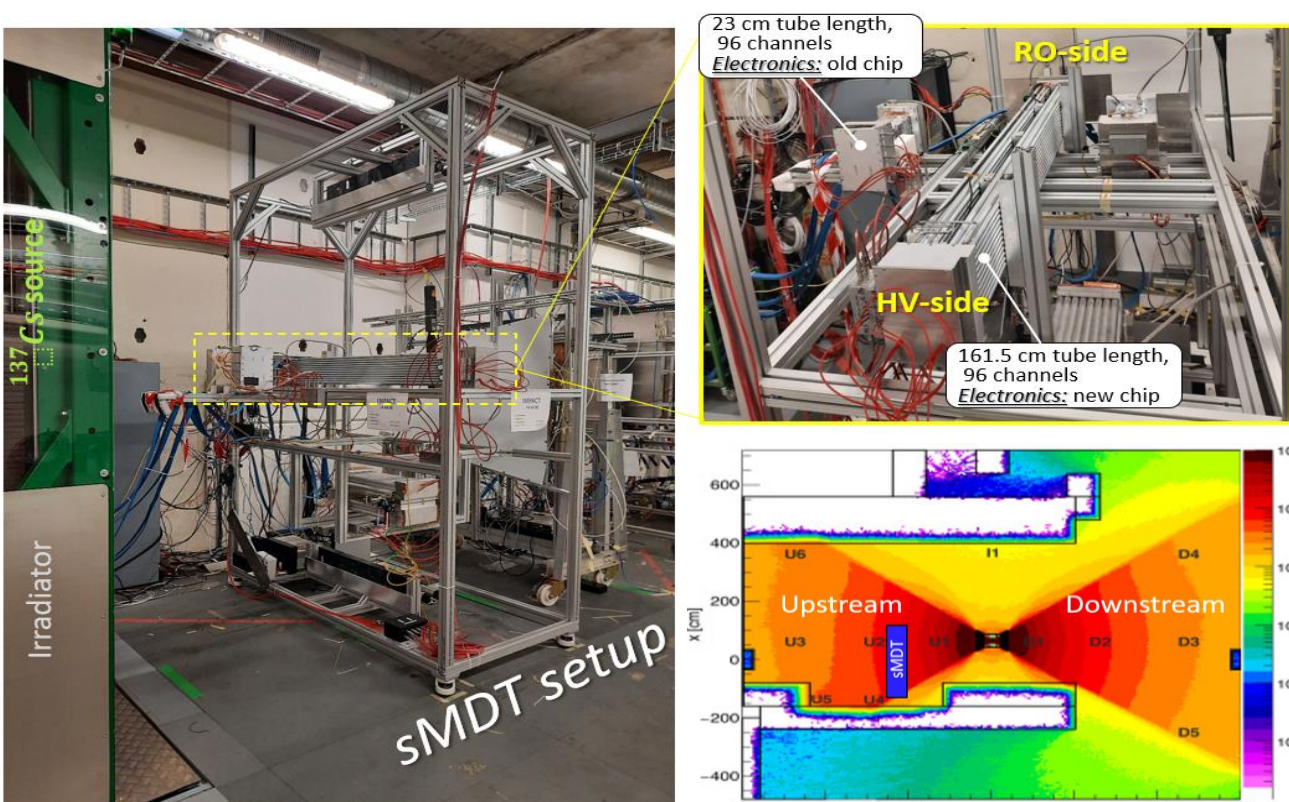
To improve the performance of the new sMDT chambers, as well as the existing MDT detectors installed at the ATLAS muon spectrometer new readout ASD chip has been developed at the Max-Planck Institute in Munich, and used for the HL-LHC era, as well as can be used for FCC-hh experiments.



The **main advantage** is the **shorter peaking time**:

- Reducing discriminator threshold crossing time jitter.
- Faster baseline recovery, reducing the signal pile-up effects at high counting rates.
- Improved chip design leads to an overall noise reduction, which makes the discriminator threshold adjustment essentially unnecessary.

6. sMDT Setup at CERN Gamma Irradiation Facility (GIF++)



The new ASD chips were tested on a sMDT chamber with 1.6 m long in CERN's Gamma Irradiation Facility (GIF++).

The sMDT detector under test is placed at $\sim 2.0 \text{ m}$ from the source in the upstream field:

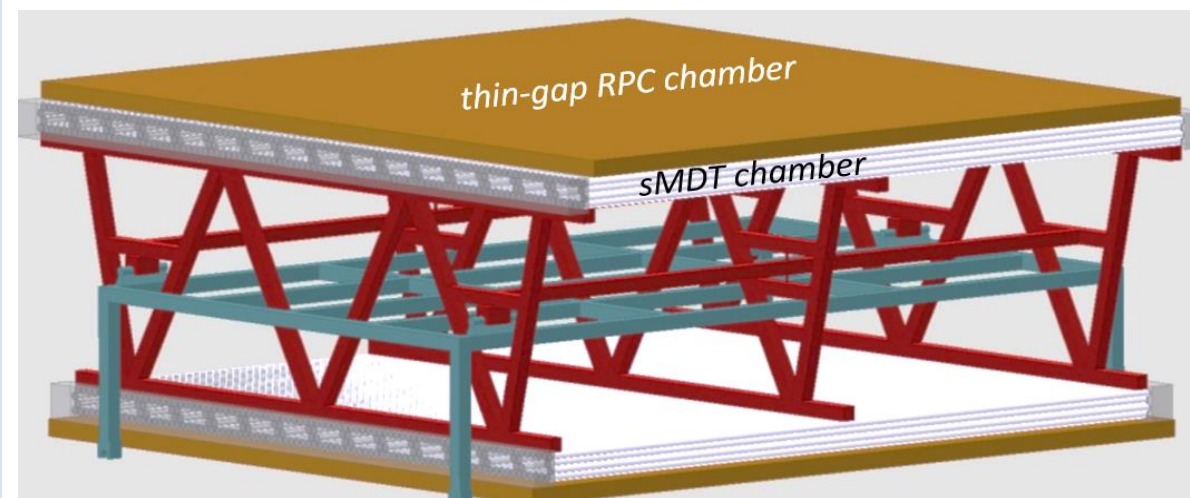
- a high-energy muon beam (100 GeV/c , $\sim 10^4 \text{ muon/spill}$), from secondary SPS beam line H4 in EHN1.
- a gamma photons (662 keV), from ^{137}Cs source with flux from 0 (source OFF) up to $\sim 10^7 \gamma/\text{cm}^2 \cdot \text{s}$ by using different absorber (ABS) settings [4].

sMDT's operating condition:

- Gas mixture: Ar: CO₂ (93:7) @ 3 bar
- Operating working point: 2730 V ($G = 2 \times 10^4$)

2. Precision sMDT chambers for FCC-hh Muon System

An angular resolution of $70 \mu\text{rad}$ can be achieved with small-diameter Monitored Drift-Tube chambers, so-called "sMDT chambers". The muon detectors will be exposed to a large background of neutrons and γ rays, which will lead to counting rates of up to $500 \text{ Hz} \times \text{cm}^{-2}$ in the barrel and the outer end-cap muon system and up to $10 \text{ kHz} \times \text{cm}^{-2}$ in the inner endcap and forward muon system. sMDT chambers can be operated at counting rates of up to $30 \text{ kHz} \times \text{cm}^{-2}$ and background occupancies of up to 30% [2].



2×4 layers of 15 mm diameter drift tubes at 1.5 m multilayer distance provide the required angular resolution.

Main Goal is to achieve the required angular resolution of $70 \mu\text{rad}$ with an sMDT chamber with two quadruple layers at 1.5 m distance, the single tube resolution must not exceed $150 \mu\text{m}$.

Advantages of the sMDT chambers:

- The small-diameter MDT detector technology developed for the upgrade of the ATLAS detector can operate in an environment up to $\sim 30 \text{ kHz/cm}^2$.
- High reliability and robustness.
- Cost-effective way solution for precise track point.
- Angle measurement over large areas (about $\sim 1200 \text{ m}^2$).
- Can achieve resolution of better than $50 \mu\text{rad}$ and $70 \mu\text{rad}$.

4. High-Rate Phenomena in sMDT Chambers

Spatial resolution and gas gain loss in drift tubes due to space charge effect

- Altered electric field inside the tube due to a constant positive space charge ρ_V :

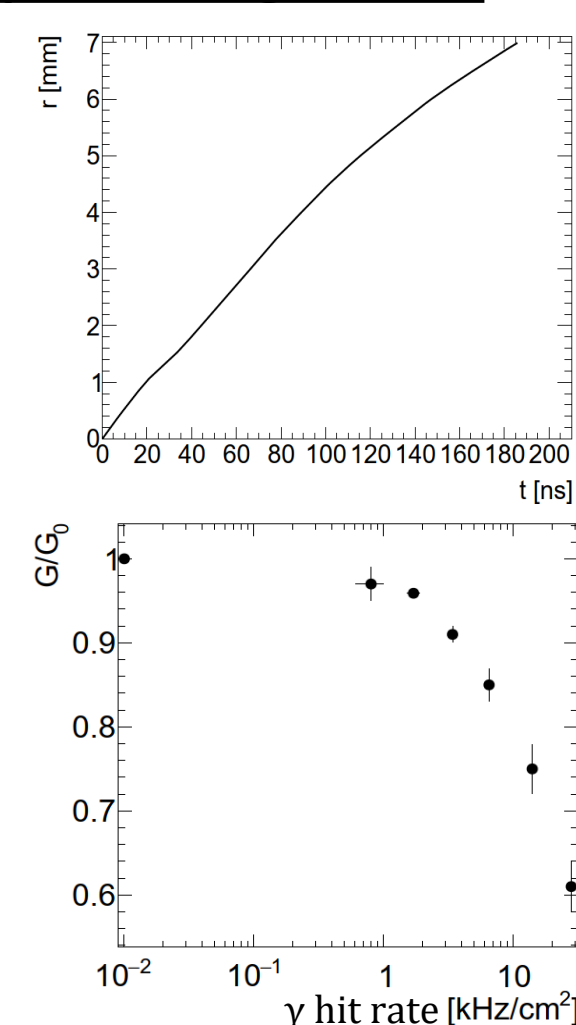
$$E(r) = \frac{U_0 - \frac{\rho_V R^2}{4\epsilon_0}}{\ln \frac{R}{r_0}} \times \frac{1}{r} + \frac{1}{2} \frac{\rho_V}{\epsilon_0} r$$

- Due to the linearity of the space drift-time (r-t) relationship the electron drift velocity is unaffected by the modification of the electric field.
- Only consequence of the altered $E(r)$: gas gain drop due reduced field at the anode wire.

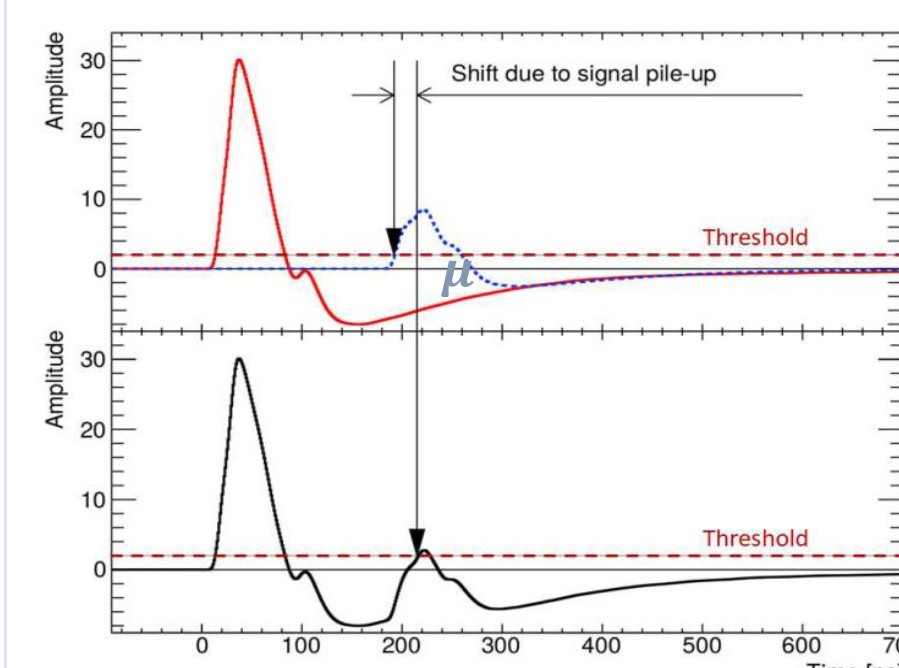
→ **Negative Results:** deterioration of the tube spatial resolution and muon detection efficiency.

→ **Solution:** need to adjust the operating voltage U_0 to compensate for the drop in the gas gain.

Standard *Diethorn formula* can be used to predict the necessary adjustment



Readout electronics signal pile-up at high counting rates

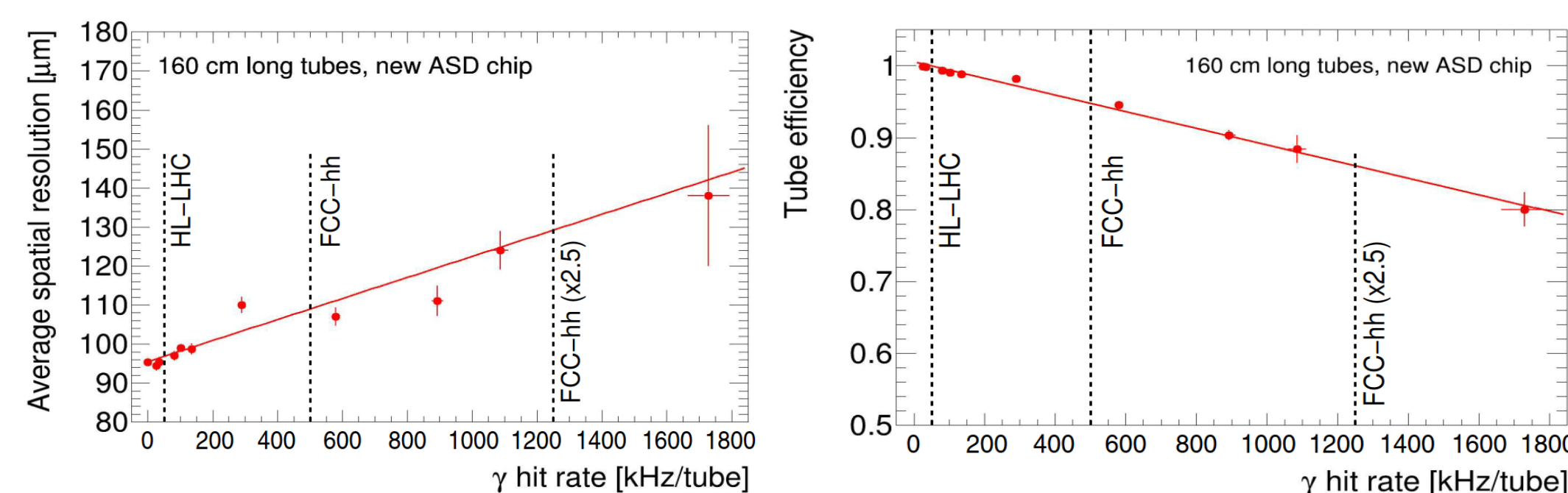


- A bipolar shaping technique is used by ATLAS sMDT in order to avoid baseline shift associated with unipolar shaping.

Disadvantage: the large and long undershoot of equal area after the positive signal pulse.

- Muon hits falling into the negative tail of a preceding γ hit get deteriorated leading to a degradation of the spatial resolution.

7. Spatial resolution and muon detection efficiency as a function of the γ background rate



- By using the operating voltage adjustment technique, the good spatial resolution result of $\sim 135 \mu\text{m}$ at the highest rate of about 1.7 MHz/tube has been achieved which correspond to the 30% occupancy of the sMDT detector.
- Tube muon detection efficiency of more than 80% at the 1.7 MHz/tube counting rate (30% occupancy) has been achieved.

→ This result demonstrate that we can ensure that more 99% muon track reconstruction efficiency with 8-layer drift tube of sMDT detector.

→ This result achieved thanks to the short shaping time implemented in the new ASD chip that also minimized the pile-up effect.

- Small-diameter Drift Tube Detector (sMDT) technology is the ideal choice for the cost-effective instrumentation of large-area muon systems when high spatial resolution is required.
- sMDT chambers with new read-out electronics can be operated with the high spatial resolution and muon detection efficiency up to the highest background rates expected at future hadron collider experiments.
- It is still necessary to further optimize and improve the performance of the readout electronics to fully exploit the sMDT rate capability.

[1] M. Benedikt et al., FCC-hh: The Hadron Collider: Future Circular Collider Conceptual Design Report Volume 3.
[2] O. Kortner et al., Design of the FCC-hh muon detector and trigger system, <https://doi.org/10.1016/j.nima.2018.10.013>
[3] H. Kroha et al., Design and construction of integrated small-diameter drift tube and thin-gap resistive plate chambers for the phase-1 upgrade of the ATLAS muon spectrometer, <https://doi.org/10.1016/j.nima.2018.10.139>
[4] D. Pfeiffer et al., The radiation field in the Gamma Irradiation Facility GIF++ at CERN, <https://doi.org/10.1016/j.nima.2017.05.045>