

Coalescence plus fragmentation approach for heavy hadrons production: spectra and ratios in AA and pp collisions

Vincenzo Minissale

Dipartimento di Fisica e Astronomia “Ettore Majorana”
Università degli studi di Catania

INFN/LNS

In collaboration with: S. Plumari, V.Greco



UNIVERSITÀ
degli STUDI
di CATANIA



Quinto Incontro Nazionale di Fisica Nucleare INFN 2022
9-11 May 2022

Outline

Hadronization:

- Fragmentation
- Coalescence model

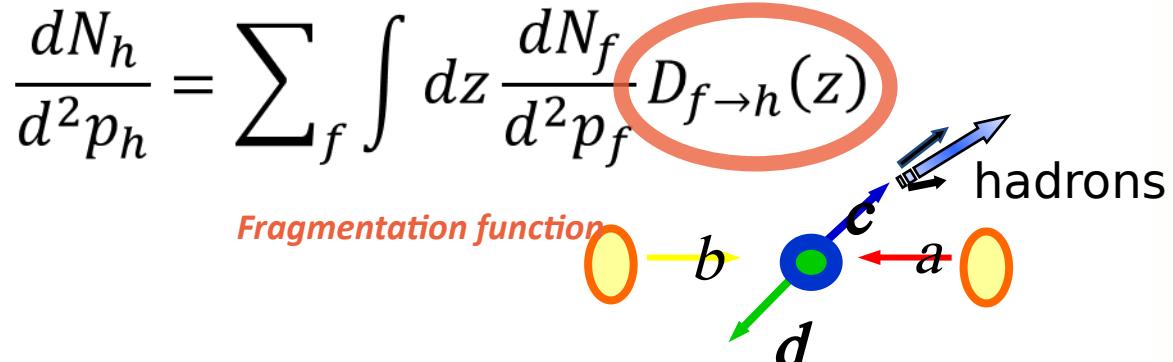
Heavy hadrons in AA collisions:

- Λ_c , D spectra and ratio: RHIC and LHC

Heavy hadrons in small systems (pp @ 5.02 TeV):

- Λ_c/D^0
- Ξ_c/D^0 , Ω_c/D^0

Heavy flavour Hadronization: Fragmentation



The distribution function is evaluated at the Fixed-Order plus Next-to-Leading-Log (FONLL)

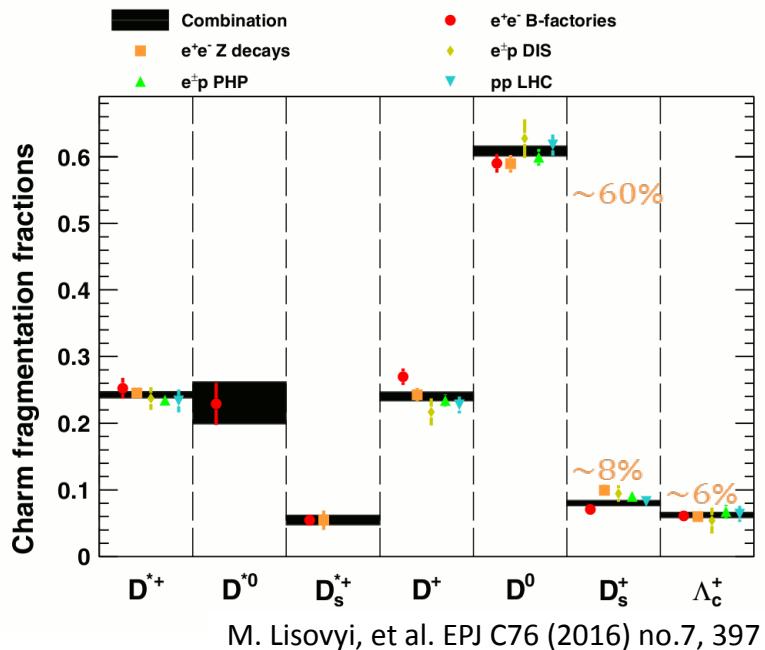
M. Cacciari, P. Nason, R. Vogt, PRL 95 (2005) 122001

In AA: bulk+charm evolution with Relativistic Transport Boltzmann Equation

We use the Peterson fragmentation function

C. Peterson, D. Schalatter, I. Schmitt, P.M. Zerwas PRD 27 (1983) 105

$$D_{f \rightarrow h}(z) \propto \frac{1}{z \left[1 - \frac{1}{z} - \frac{\epsilon}{1-z} \right]^2}$$

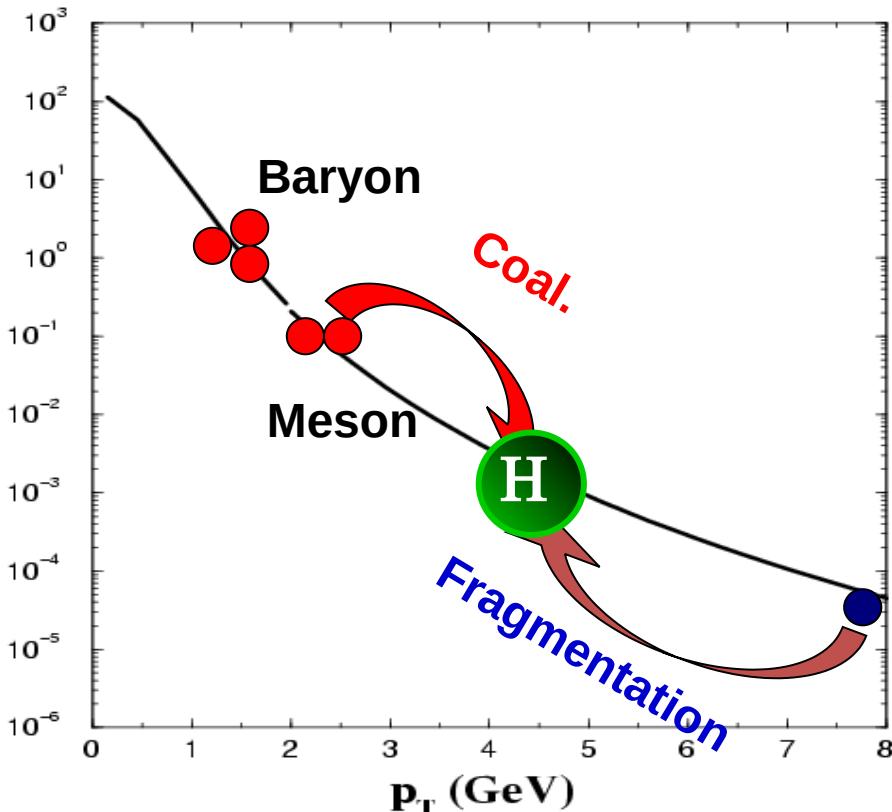


Charm Fragmentation Fraction ($c \rightarrow h$)
Measurement in $e^+ p$, $e^+ e^-$ and old pp data

$$\left(\frac{\Lambda_c^+}{D^0} \right)_{e^+ e^-} \approx 0.1 \quad \left(\frac{D_s^+}{D^0} \right)_{e^+ e^-} \approx 0.13$$

Hadronization: Coalescence

<i>Statistical factor colour-spin-isospin</i> $\frac{dN_{Hadron}}{d^2 p_T} = g_H \int \prod_{i=1}^n p_i \cdot d\sigma_i \frac{d^3 p_i}{(2\pi)^3}$	<i>Parton Distribution function</i> $f_q(x_i, p_i)$	<i>Hadron Wigner function</i> $f_W(x_1, \dots, x_n; p_1, \dots, p_n) \delta(p_T - \sum_i p_{iT})$
--	--	--



Used to describe first observations in light sector for the baryon/meson ratio and elliptic flow splitting at RHIC, more than a decade ago.

V. Greco, C.M. Ko, P. Levai PRL 90, 202302 (2003)

Hadronization: Coalescence

$$\frac{dN_{Hadron}}{d^2 p_T} = g_H \int \prod_{i=1}^n p_i \cdot d\sigma_i \frac{d^3 p_i}{(2\pi)^3} f_q(x_i, p_i) f_W(x_1, \dots, x_n; p_1, \dots, p_n) \delta(p_T - \sum_i p_{iT})$$

*Statistical factor
colour-spin-isospin*

*Parton Distribution
function*

*Hadron Wigner
function*

LIGHT

Thermal+flow for u,d,s ($p_T < 3$ GeV)

$$\frac{dN_{q,\bar{q}}}{d^2 p_T} \sim \exp\left(-\frac{\gamma_T - p_T \cdot \beta_T \mp \mu_q}{T}\right)$$

$$\beta(r) = \frac{r}{R} \beta_{max}$$

$$V = \pi R^2 \tau \cosh(y_z), R(\tau_f) = R_0(1 + \beta_{max} \tau_f)$$

$$\text{PbPb@5ATeV(0-10%)}: \tau_f = 8.4 \frac{fm}{c} \rightarrow V_{|y|<0.5} = 4500 fm^3$$

+quenched minijets for u,d,s ($p_T > 3$ GeV)

CHARM

In AA collisions charm distribution from the studies of R_{AA} and v_2 of **D-meson** to determine the Space Diffusion coefficient:

parton simulations solving relativistic Boltzmann transport equation

In pp collisions the charm distribution are the FONLL distribution

Coalescence simulation in a fireball with radial flow for light quarks \rightarrow dimension set by experimental constraints

Hadronization: Coalescence

<p><i>Statistical factor colour-spin-isospin</i></p> $\frac{dN_{Hadron}}{d^2 p_T} = g_H \int \prod_{i=1}^n p_i \cdot d\sigma_i \frac{d^3 p_i}{(2\pi)^3}$	<p><i>Parton Distribution function</i></p>	<p><i>Hadron Wigner function</i></p> $f_W(x_1, \dots, x_n; p_1, \dots, p_n) \delta(p_T - \sum_i p_{iT})$
--	--	--

Wigner function – Wave function

$$\Phi_M^W(\mathbf{r}, \mathbf{q}) = \int d^3 r' e^{-i\mathbf{q}\cdot\mathbf{r}'} \phi_M\left(\mathbf{r} + \frac{\mathbf{r}'}{2}\right) \phi_M^*\left(\mathbf{r} - \frac{\mathbf{r}'}{2}\right)$$

$\phi_M(\mathbf{r})$ meson wave function

Wigner function width fixed by root-mean-square charge radius from quark model

C.-W. Hwang, EPJ C23, 585 (2002), C. Albertus et al., NPA 740, 333 (2004)

$$\langle r^2 \rangle_{ch} = \frac{3}{2} \frac{m_2^2 Q_1 + m_1^2 Q_2}{(m_1 + m_2)^2} \sigma_{r1}^2 + \frac{3}{2} \frac{m_3^2 (Q_1 + Q_2) + (m_1 + m_2)^2 Q_3}{(m_1 + m_2 + m_3)^2} \sigma_{r2}^2$$

$$\sigma_{ri} = 1/\sqrt(\mu_i \omega) \quad \mu_1 = \frac{m_1 m_2}{m_1 + m_2} \quad \mu_2 = \frac{(m_1 + m_2)m_3}{m_1 + m_2 + m_3}$$

Assuming gaussian wave function

$$f_H(\dots) = \prod_{i=1}^{N_q-1} A_w \exp\left(-\frac{x_{ri}^2}{\sigma_{ri}^2} - p_{ri}^2 \sigma_{ri}^2\right)$$

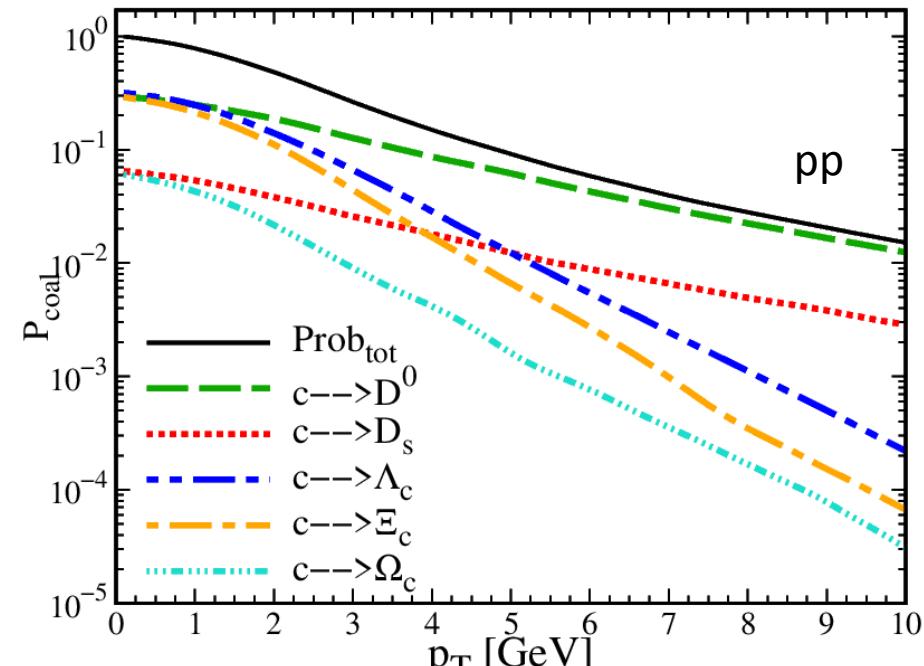
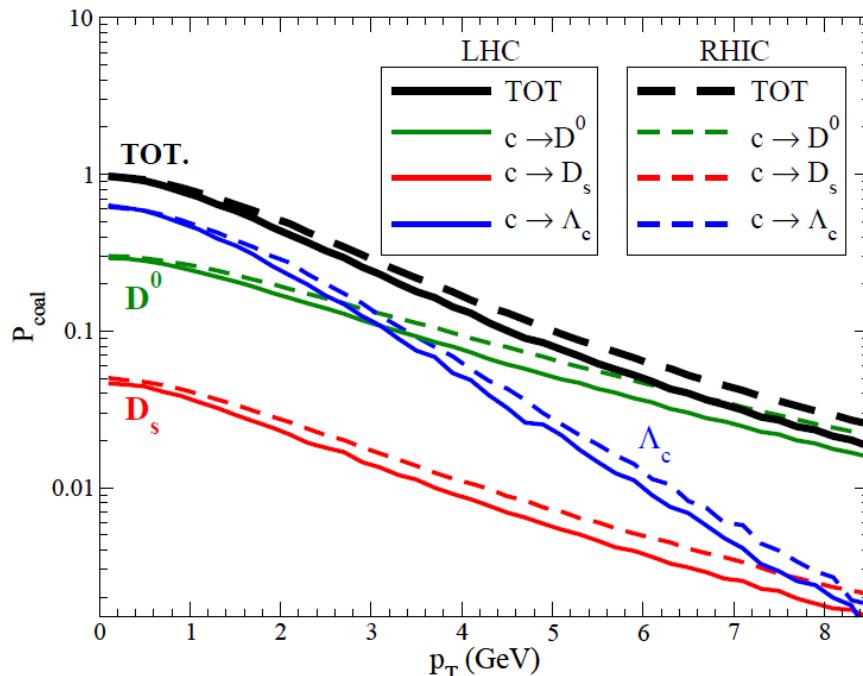
only one width coming from $\phi_M(\mathbf{r})$,
constraint $\sigma_r \sigma_p = 1$

Meson	$\langle r^2 \rangle_{ch}$	σ_{p1}	σ_{p2}
$D^+ = [c\bar{d}]$	0.184	0.282	—
$D_s^+ = [\bar{s}c]$	0.083	0.404	—
Baryon	$\langle r^2 \rangle_{ch}$	σ_{p1}	σ_{p2}
$\Lambda_c^+ = [udc]$	0.15	0.251	0.424
$\Xi_c^+ = [usc]$	0.2	0.242	0.406
$\Omega_c^0 = [ssc]$	-0.12	0.337	0.53

Hadronization: Coalescence

<i>Statistical factor colour-spin-isospin</i> $\frac{dN_{Hadron}}{d^2 p_T} = g_H \int \prod_{i=1}^n p_i \cdot d\sigma_i \frac{d^3 p_i}{(2\pi)^3}$	<i>Parton Distribution function</i> $f_q(x_i, p_i)$	<i>Hadron Wigner function</i> $f_W(x_1, \dots, x_n; p_1, \dots, p_n) \delta(p_T - \sum_i p_{iT})$
--	--	--

- Normalization in $f_W(\dots)$ requiring that $P_{coal}=1$ at $p=0$
- The charm that does not coalesce undergo fragmentation



Heavy flavour: Resonance decay

Meson	Mass(MeV)	$I(J)$	Decay modes	B.R.
$D^+ = \bar{d}c$	1869	$\frac{1}{2}(0)$		
$D^0 = \bar{u}c$	1865	$\frac{1}{2}(0)$		
$D_s^+ = \bar{s}c$	2011	0(0)		
Resonances				
D^{*+}	2010	$\frac{1}{2}(1)$	$D^0\pi^+$; D^+X	68%, 32%
D^{*0}	2007	$\frac{1}{2}(1)$	$D^0\pi^0$; $D^0\gamma$	62%, 38%
D_s^{*+}	2112	0(1)	D_s^+X	100%
Baryon				
$\Lambda_c^+ = udc$	2286	0($\frac{1}{2}$)		
$\Xi_c^+ = usc$	2467	$\frac{1}{2}(\frac{1}{2})$		
$\Xi_c^0 = dsc$	2470	$\frac{1}{2}(\frac{1}{2})$		
$\Omega_c^0 = ssc$	2695	0($\frac{1}{2}$)		
Resonances				
Λ_c^+	2595	0($\frac{1}{2}$)	$\Lambda_c^+\pi^+\pi^-$	100%
Λ_c^+	2625	0($\frac{3}{2}$)	$\Lambda_c^+\pi^+\pi^-$	100%
Σ_c^+	2455	1($\frac{1}{2}$)	$\Lambda_c^+\pi$	100%
Σ_c^+	2520	1($\frac{3}{2}$)	$\Lambda_c^+\pi$	100%
$\Xi_c^{'+,0}$	2578	$\frac{1}{2}(\frac{1}{2})$	$\Xi_c^{+,0}\gamma$	100%
Ξ_c^+	2645	$\frac{1}{2}(\frac{3}{2})$	$\Xi_c^+\pi^-$,	100%
Ξ_c^+	2790	$\frac{1}{2}(\frac{1}{2})$	$\Xi_c'\pi$,	100%
Ξ_c^+	2815	$\frac{1}{2}(\frac{3}{2})$	$\Xi_c'\pi$,	100%
Ω_c^0	2770	0($\frac{3}{2}$)	$\Omega_c^0\gamma$,	100%

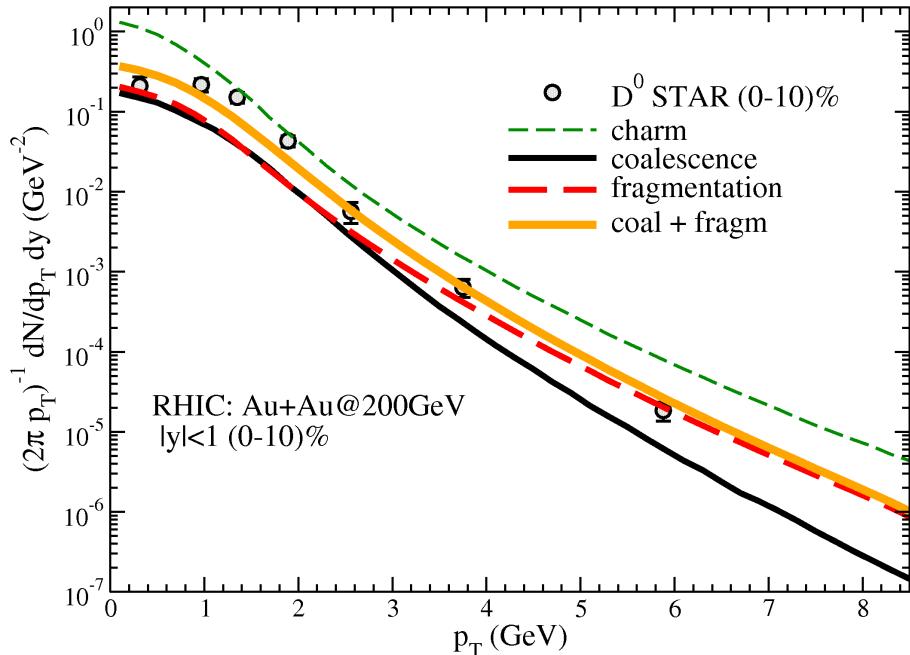
In our calculations we take into account hadronic channels including the ground states + first excited states

Statistical factor suppression for resonances

$$\frac{[(2J+1)(2I+1)]_{H^*}}{[(2J+1)(2I+1)]_H} \left(\frac{m_{H^*}}{m_H} \right)^{3/2} e^{-(m_{H^*} - m_H)/T}$$

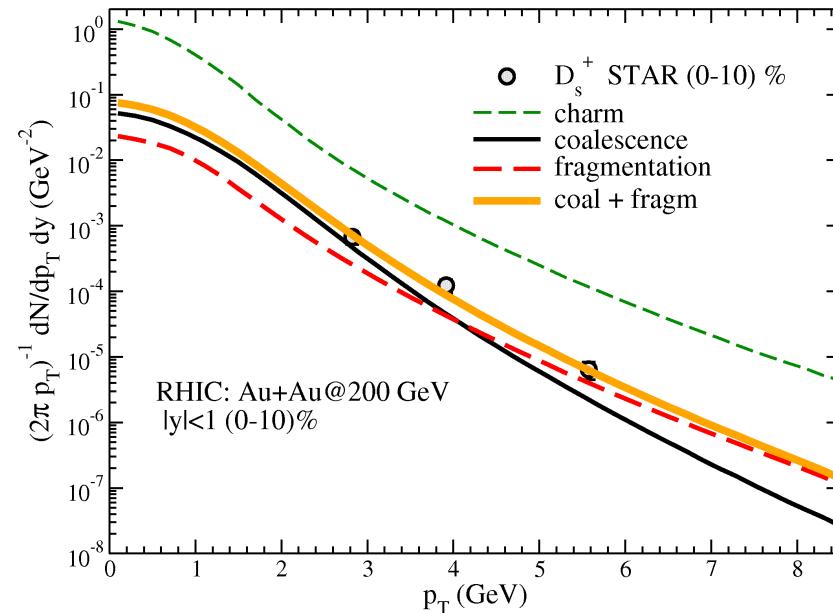
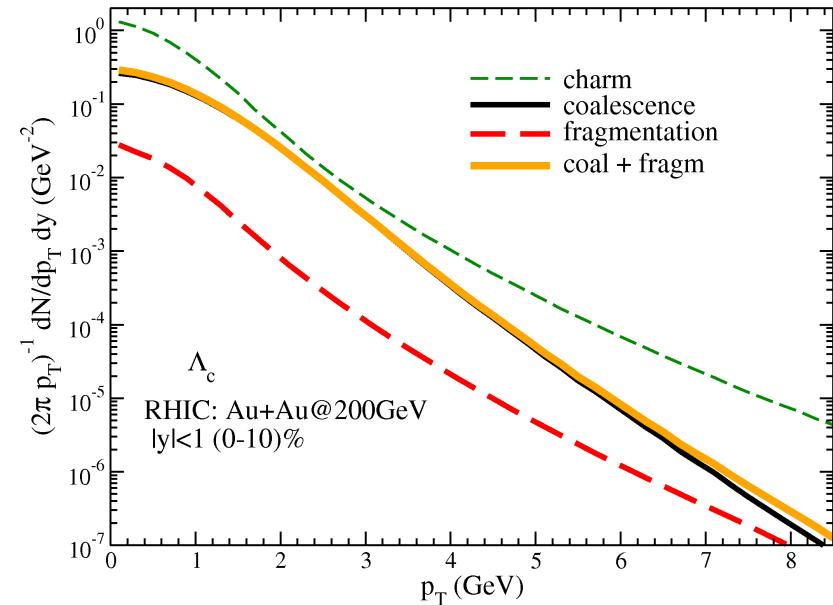
RHIC: results

S. Plumari, V. Minissale et al., Eur. Phys. J. **C78** no. 4, (2018) 348



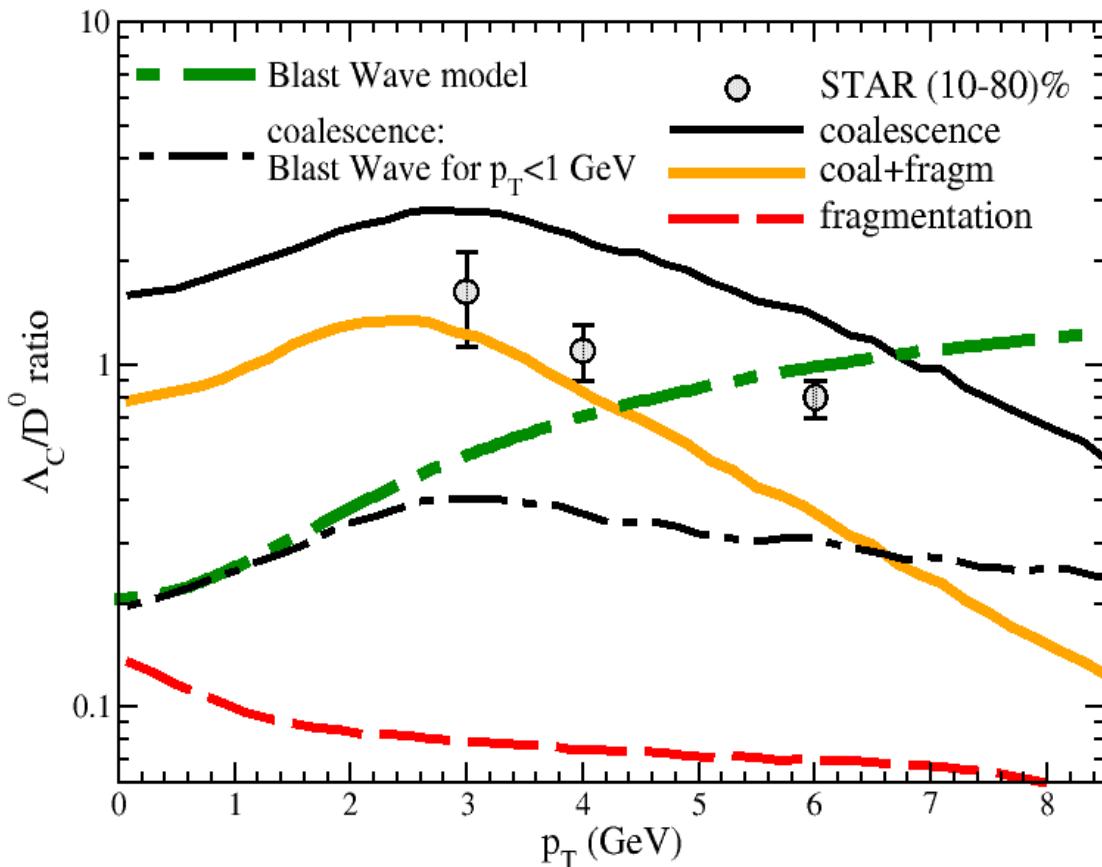
Data from STAR Coll. PRL **113** (2014) no.14, 142301

- For D^0 coalescence and fragmentation comparable at 2 GeV
- fragmentation fraction for D_s^+ are small and less than about 8% of produced total heavy hadrons
- Λ_c^+ fragmentation is even more smaller, coalescence gives the dominant contribution



RHIC: Baryon/meson

STAR, Phys.Rev.Lett. 124 (2020) 17, 172301



Compared to light baryon/meson ratio
the Λ_c/D^0 ratio has a larger width
(flatter)

More flatter \rightarrow should coalescence
extend to higher p_T ? Indication also in
light sector

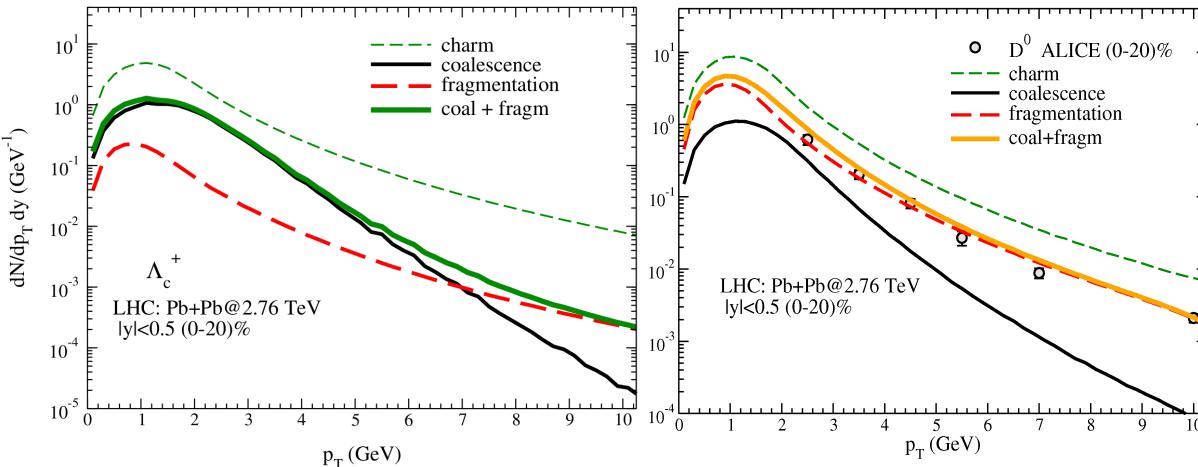
V. Minissale, F. Scardina, V. Greco PRC 92,054904 (2015)
Cho, Sun, Ko et al., PRC 101 (2020) 2, 024909

Needed data at low p_T

LHC: results

wave function widths σ_p of baryon and mesons are the same at RHIC and LHC!

Data from ALICE Coll. JHEP 1209 (2012) 112



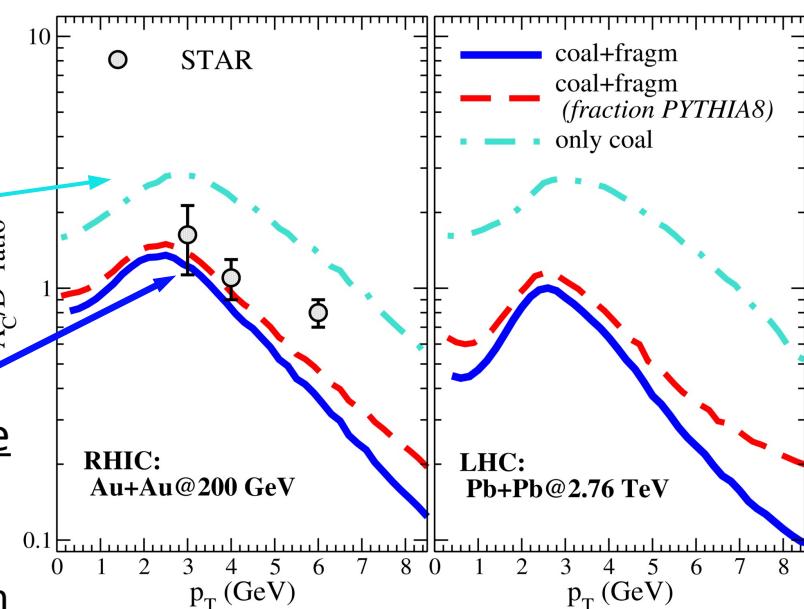
Only Coalescence ratio is similar at both energies.

Fragmentation ~ 0.1 at both energies.

the **combined ratio is different** because the coalescence over fragmentation ratio at LHC is smaller than at RHIC

Therefore at LHC the larger contribution in particle production from fragmentation leads to a final ratio that is smaller than at RHIC.

Coalescence lower than at RHIC \rightarrow main contribution from Fragmentation



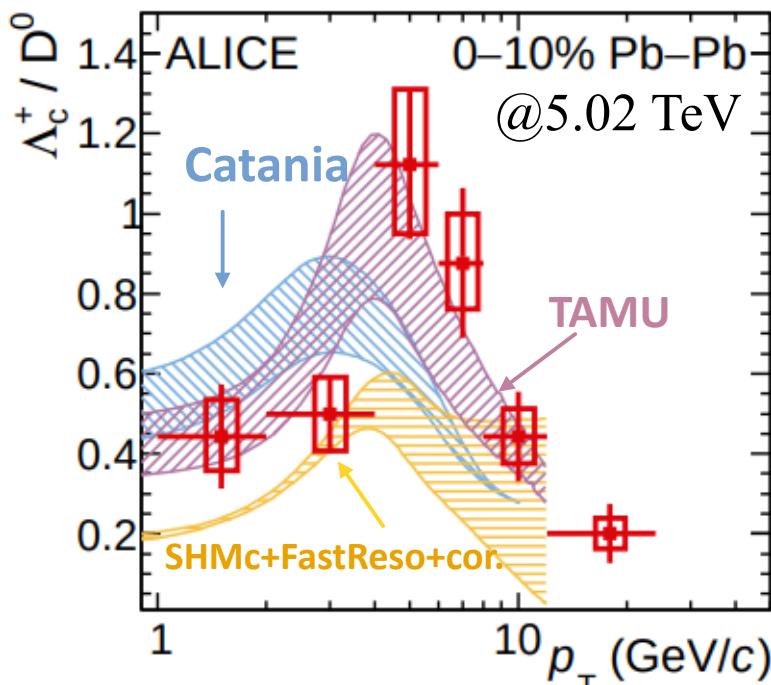
LHC: results

wave function widths σ_p of baryon and mesons are the same at RHIC and LHC!

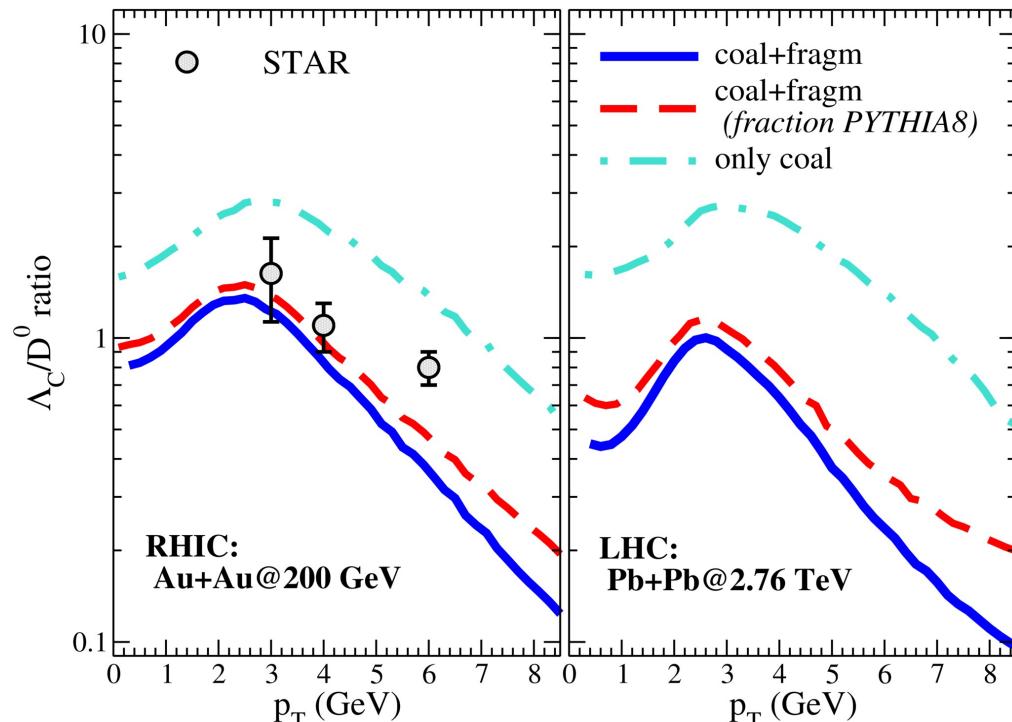
Results for 0-10% in PbPb @5.02TeV:

Consistent with the trend shown at RHIC and LHC @2.76TeV

Available data at low p_T → differences recombination vs SHM



ALICE Coll. arXiv:2112.08156v1



S. Plumari, V. Minissale et al., Eur. Phys. J. C78 no. 4, (2018) 348

Small systems: Coalescence in pp?

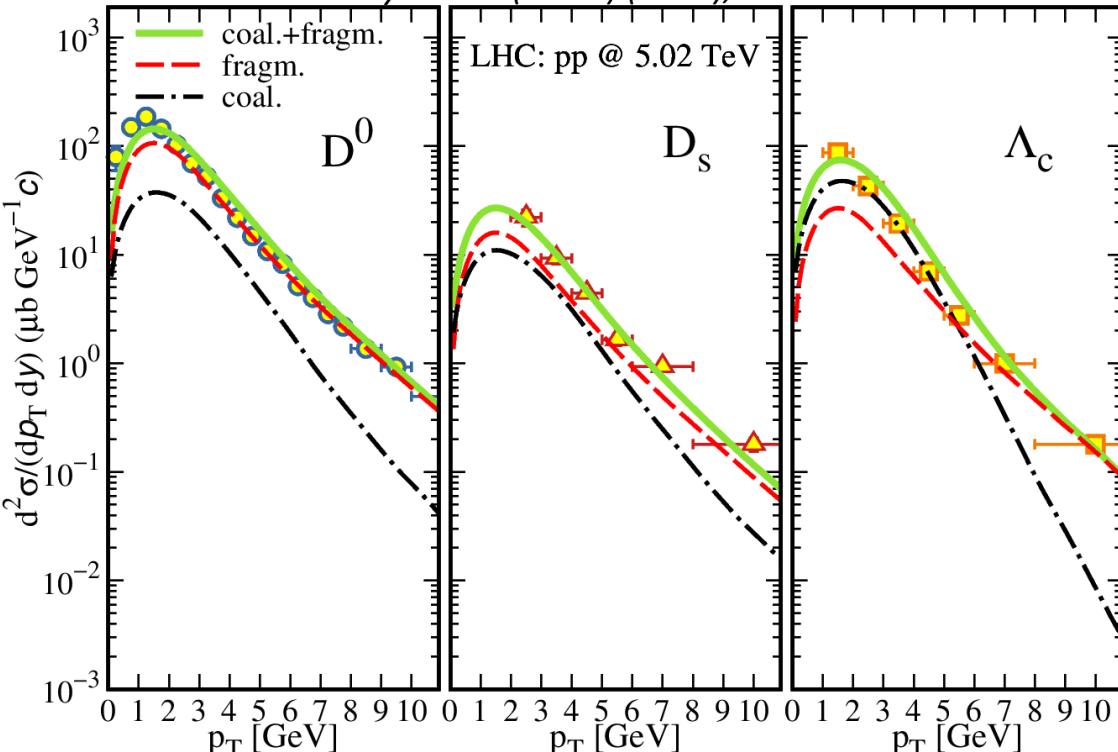
Common consensus of possible presence of QGP in smaller system.

If we assume in p+pp @ 5 TeV a medium similar to the one simulated in hydro:

What if:

- Assuming QGP formation also in pp?
- What coalescence+fragmentation predicts in this case?

S. Acharya et al. (ALICE), Eur. Phys. J. C 79, 388 (2019)
S. Acharya et al. (ALICE) (2020), 2011.06078 - 2011.06079



p+p @ 5 TeV

- $\tau_{pp}=2 \text{ fm}/c$
- $\beta_0=0.4$
- $R=2.5 \text{ fm}$
- $V \sim 30 \text{ fm}^3$

■ Thermal Distribution ($p_T < 2 \text{ GeV}$) LIGHT

$$\frac{dN_q}{d^2 r_T d^2 p_T} = \frac{g_g \tau m_T}{(2\pi)^3} \exp\left(-\frac{\gamma_T(m_T - p_T \cdot \beta_T)}{T}\right)$$

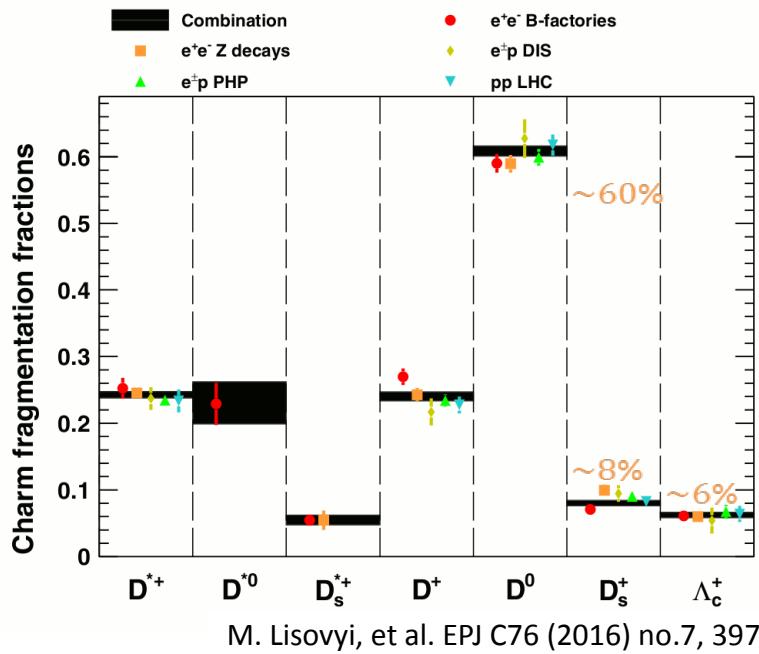
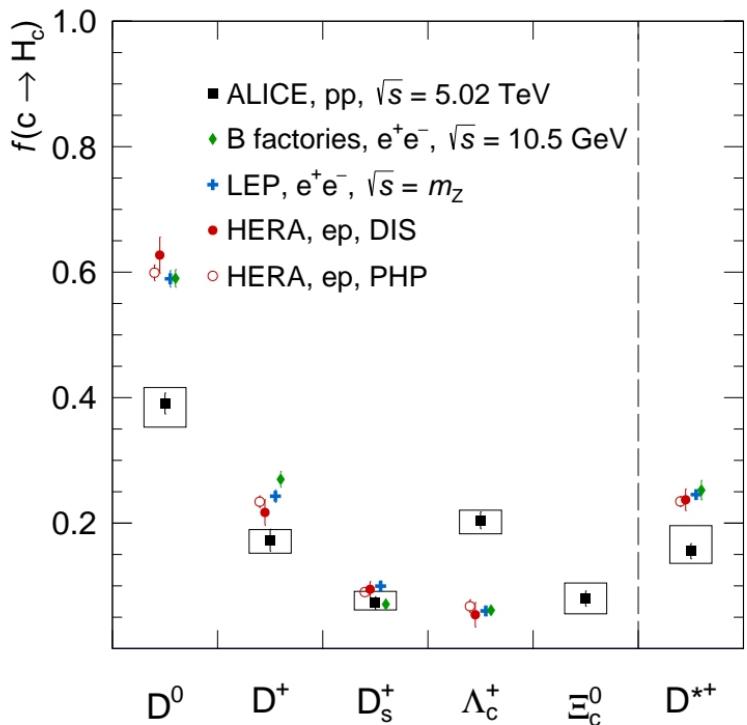
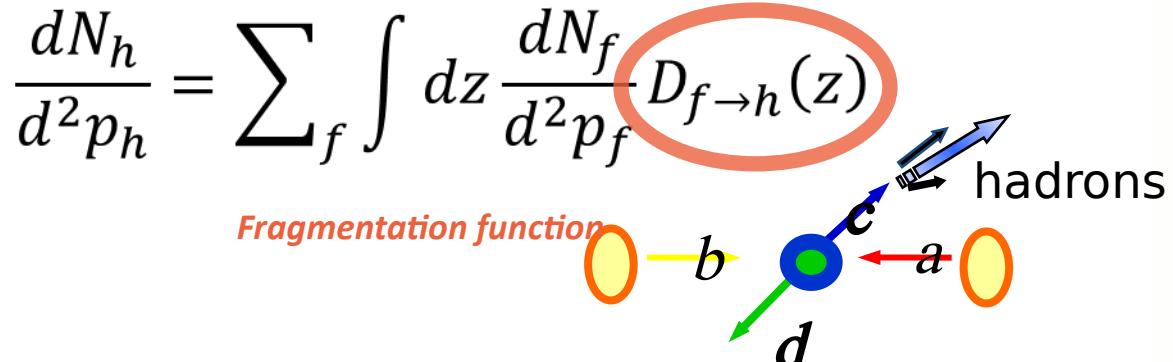
■ Minijet Distribution ($p_T > 2 \text{ GeV}$)
NO QUENCHING

CHARM

FONLL Distribution

wave function widths σ_p of baryon and mesons kept the same from AA to pp

Heavy flavour Hadronization: Fragmentation



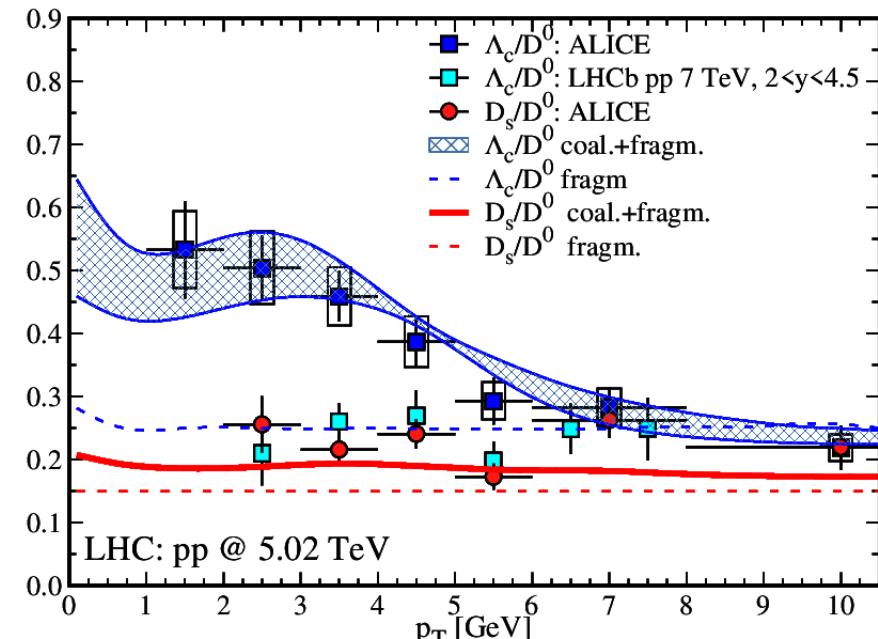
Updated fractions after experimental evidence in pp@5TeV:

Fragmentation fractions ($c \rightarrow h$) depends on collision system...and QGP presence?

No more Universality?

Small systems: Coalescence in pp?

V. Minissale, S. Plumari, V. Greco, Phys.Lett.B821 (2021) 136622



Error band correspond to $\langle r^2 \rangle$ uncertainty in quark model

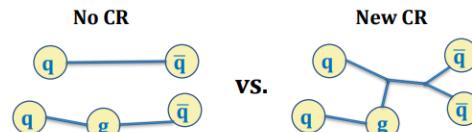
Other models:

He-Rapp, PLB795(2019): Increase ≈ 2 to Λ_c production:

SHM with resonance not present in PDG

PYTHIA8 + color reconnection

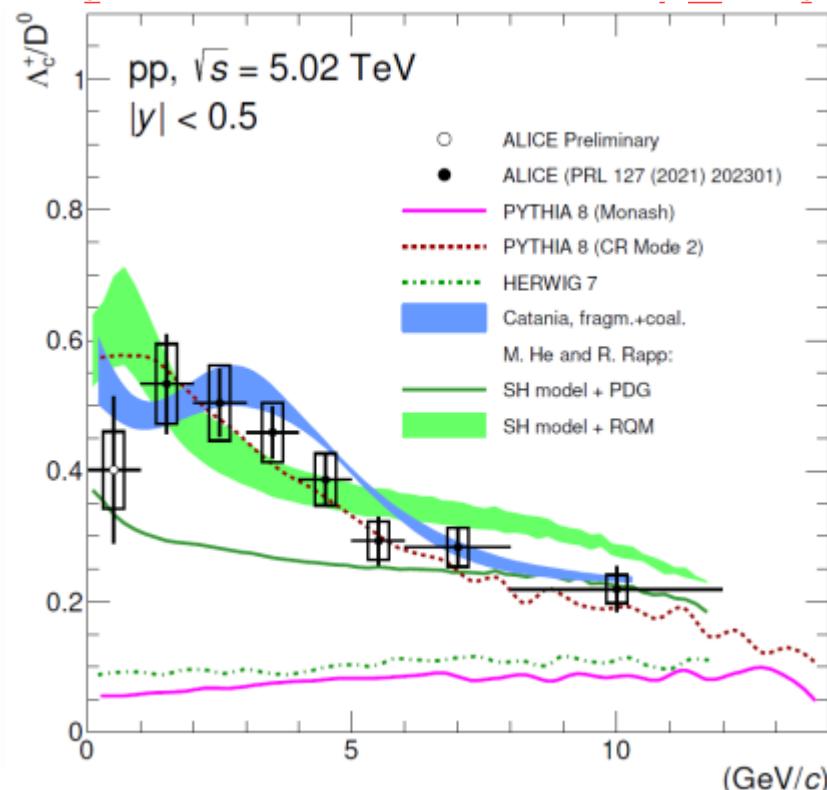
CR with SU(3) weights and string length minimization



Reduction of rise-and-fall behaviour in Λ_c^+ / D^0 ratio:

- Confronting with AA: Coal. contribution smaller w.r.t. Fragm.
- FONLL distribution flatter w/o evolution trough QGP
- Volume size effect

The increase of Λ_c production in pp have effect on R_{AA} of Λ_c

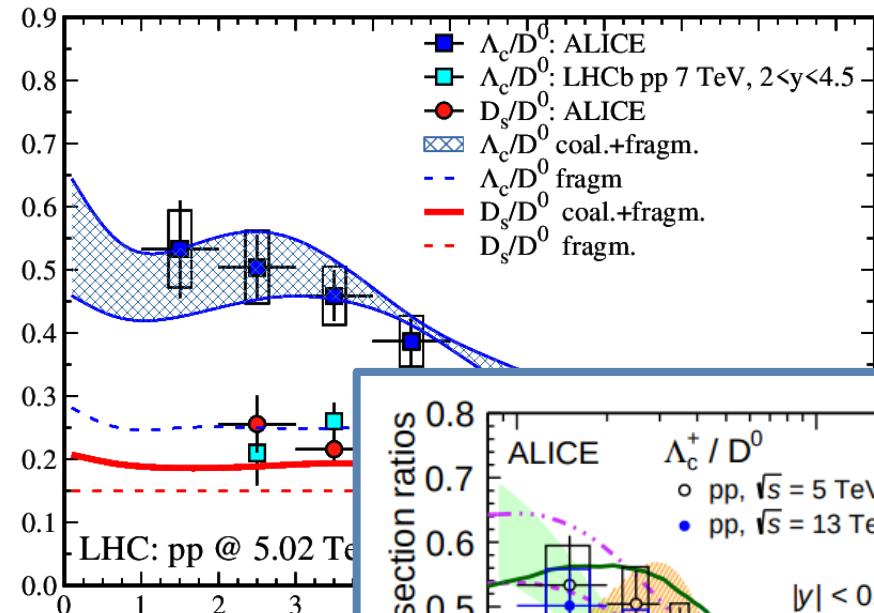


ALICE, Phys.Rev.Lett. 127 (2021) 20, 202301

Small systems: Coalescence in pp?

V. Minissale, S. Plumari, V. Greco, Phys.Lett.B821 (2021)

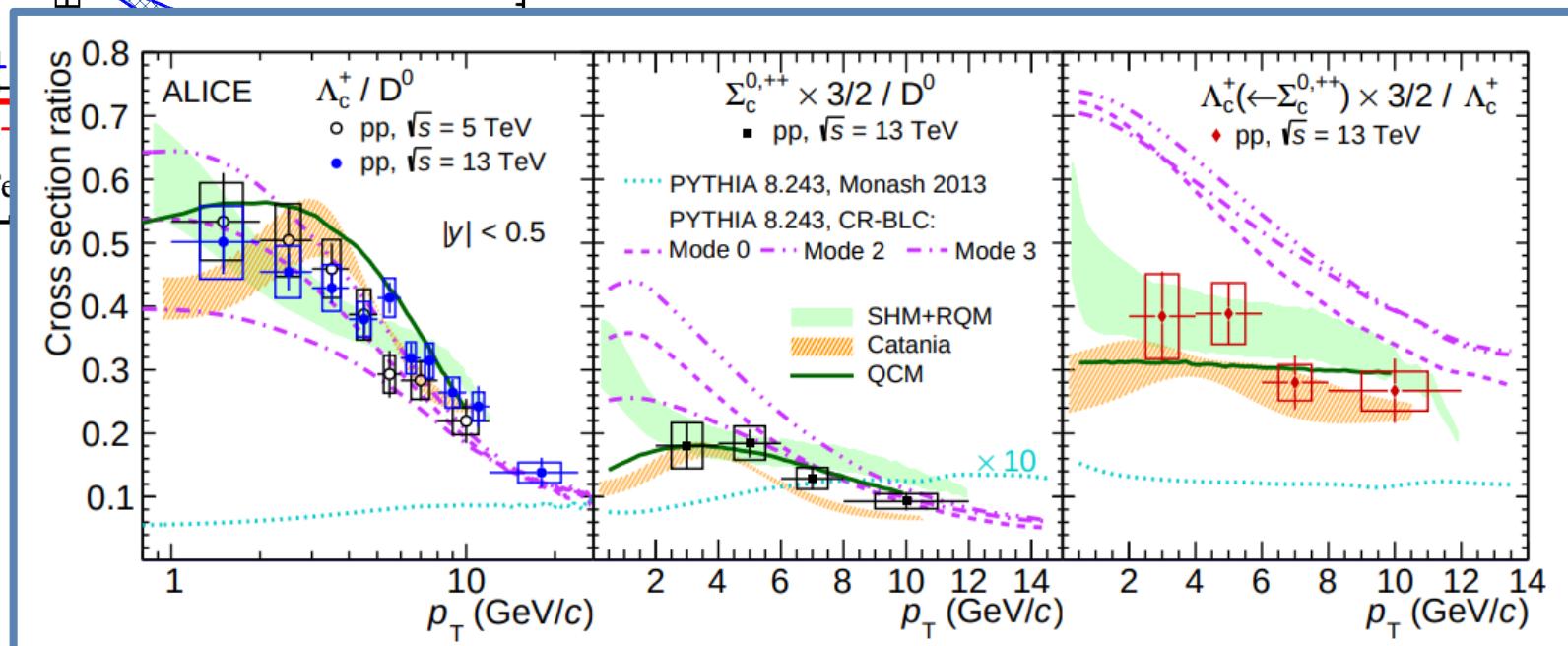
136622



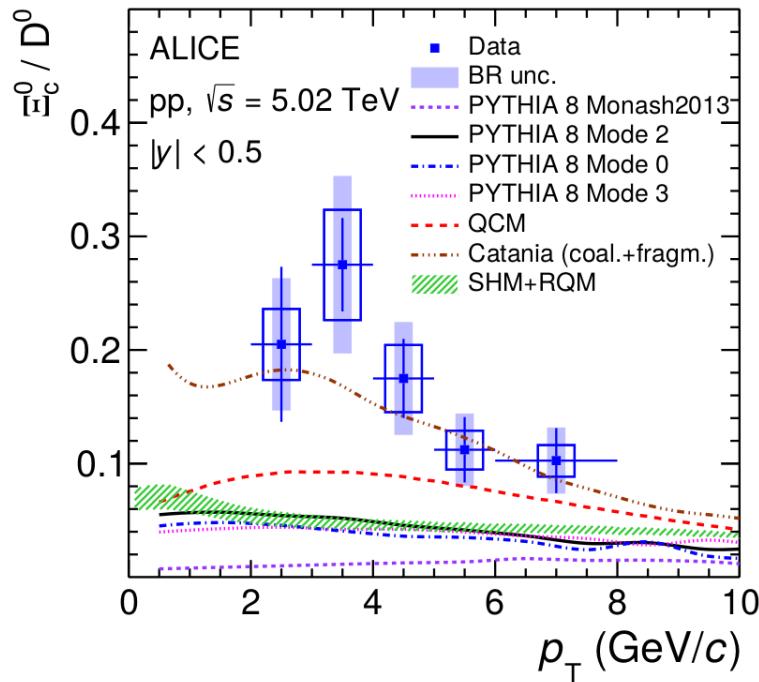
Reduction of rise-and-fall behaviour in Λ_c / D^0 ratio:

- Confronting with AA: Coal. contribution smaller w.r.t. Fragm.
- FONLL distribution flatter w/o evolution trough QGP
- Volume size effect

The increase of Λ_c production in pp have effect on R_{AA} of Λ_c



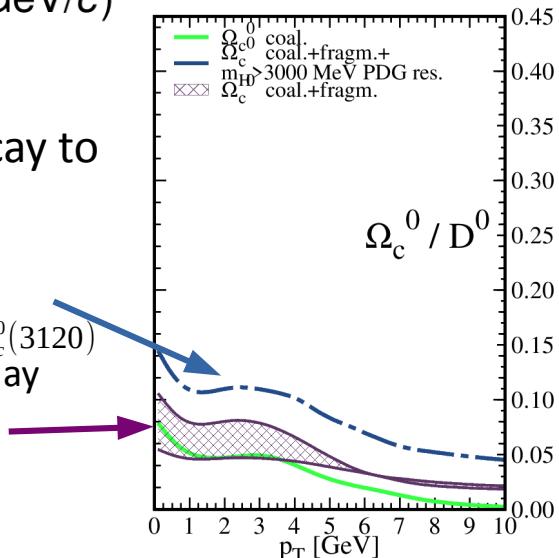
Small systems: Coalescence in pp?



Assuming additional PDG
resonances with $J=3/2$ and decay to
 Ω_c^0 additional to $\Omega_c^0(2770)$

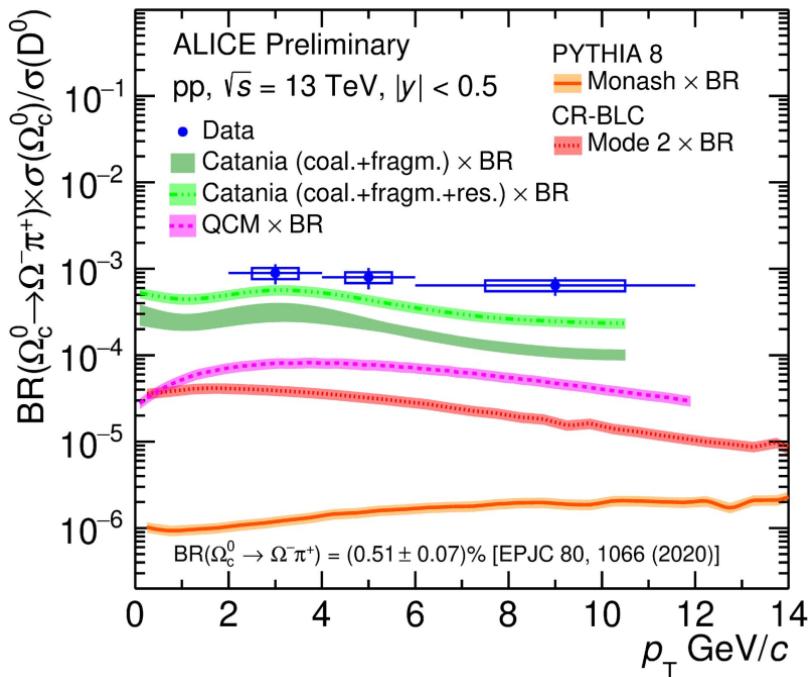
$\Omega_c^0(3000), \Omega_c^0(3005), \Omega_c^0(3065), \Omega_c^0(3090), \Omega_c^0(3120)$
supply an idea of how these states may

affect the ratio



New measurements of heavy hadrons at ALICE:

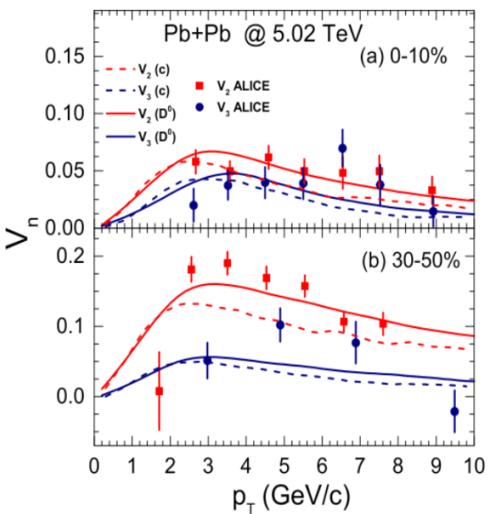
- Ξ_c^0 / D^0 ratio, same order of Λ_c^+ / D^0 : coalescence gives enhancement
- very large Ω_c^0 / D^0 ratio, our model does not get the big enhancement



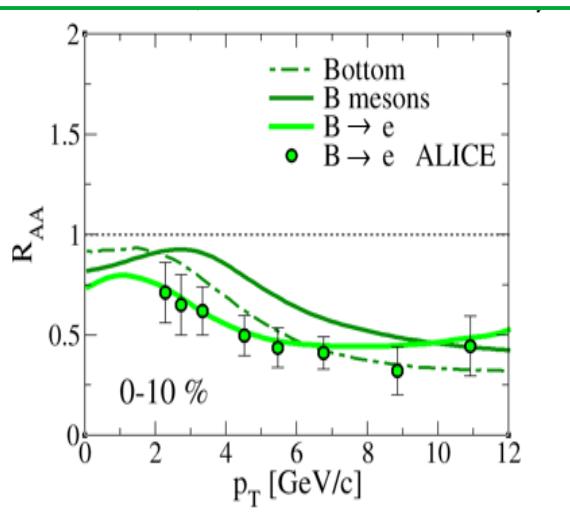
ALICE Coll. JHEP 10 (2021) 159

V. Minissale, S. Plumari, V. Greco,
Phys.Lett.B821 (2021) 136622

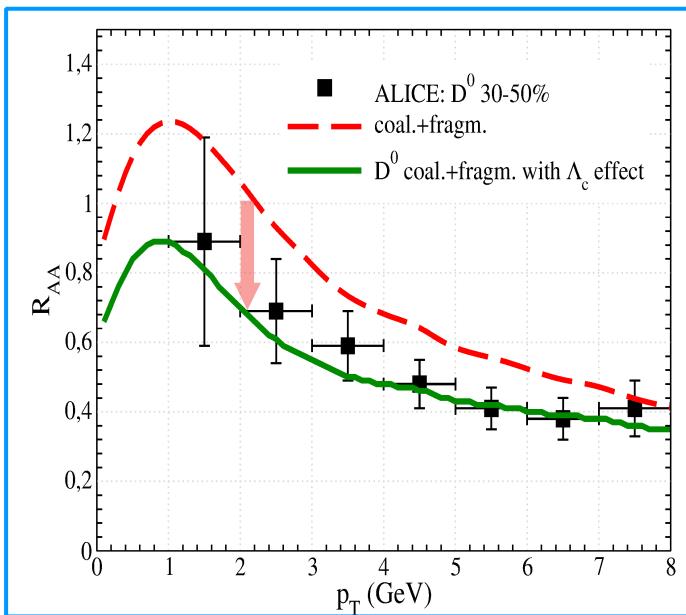
Implications, developments and outlooks:



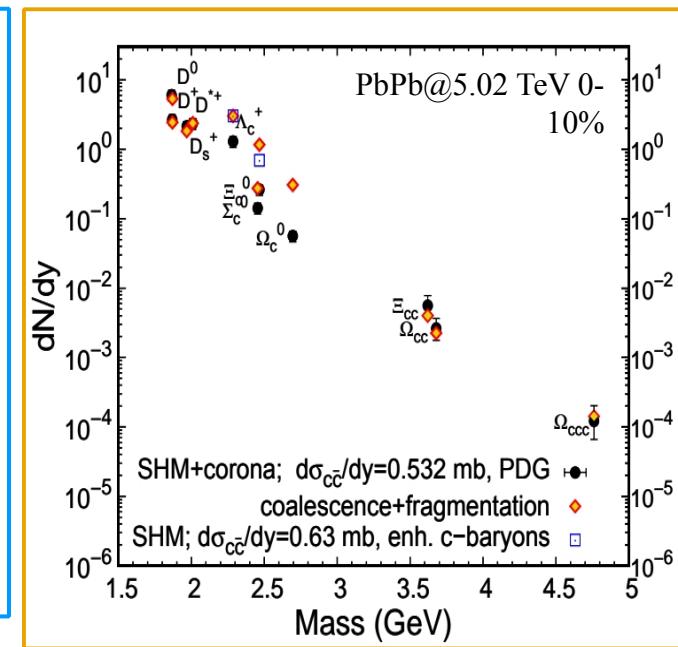
Coalescence give an enhancement to the $v_n(p_T)$ of final hadrons compared to the charm $v_n(p_T)$.
Sambataro,Sun,Minissale,Plumari,Greco(in preparation)



Electrons from semileptonic B meson decay with a coal + fragm model for B meson production
Sambataro, Minissale et al.(in preparation)



The large Λ_c production has effects on the R_{AA} of D^0 , because of the charm conservation



Yields of multicharm hadrons similar to SHM in PbPb, possible effects with system size and on the distribution vs p_T are under investigation.
A baryon like Ω_{ccc} formed only by heavy quarks can give insights on the baryon wave function and the potential $V_{cc}(r,T)$ between quarks.
Enhancement for single charmed baryons as seen in pp collisions.

Minissale, Plumari, Greco (in preparation)

MULTICCHARM

Conclusions

- *Good agreement with experimental data of hadrons spectra in AA collisions from RHIC to LHC*
- *Extension to pp: description of D mesons and Λ_c spectra*
- *Coalescence plus fragmentation gives peculiar enhancement in baryon/ meson ratio for all heavy hadrons $\Lambda_c, \Xi_c, \Omega_c$*
- *Outlook: multicharm hadrons production*

Backup Slides

Elliptic Flow – Quark Number Scaling

Fourier expansion of the azimuthal distribution

$$f(\varphi, p_T) = 1 + 2 \sum_{n=1}^{\infty} v_n(p_T) \cos n\varphi$$

momentum anisotropy in the transverse plane

$n=2$ Elliptic flow

coalescence brings to

$$\begin{aligned} v_{2,M}(p_T) &\approx 2v_{2,q}(p_T/2) \\ v_{2,B}(p_T) &\approx 3v_{2,q}(p_T/3) \end{aligned}$$

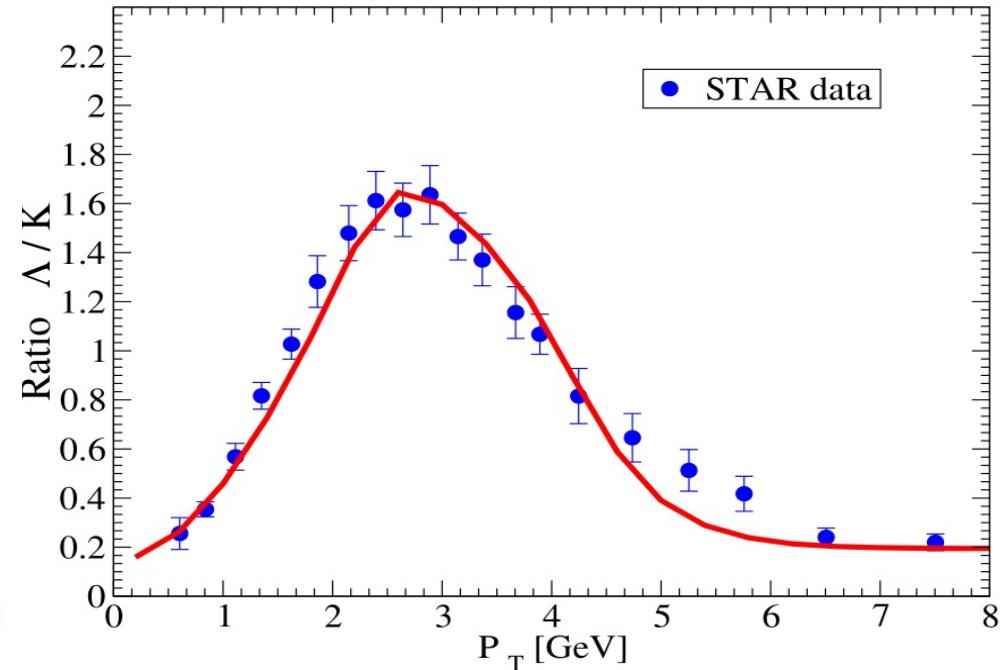
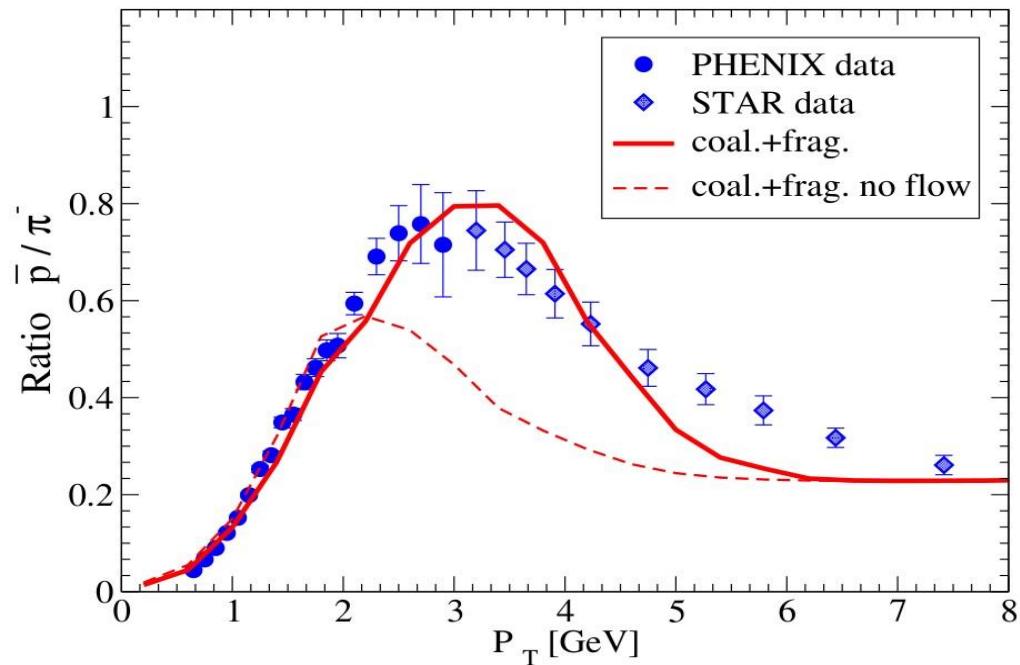
Partonic
elliptic flow

Hadronic
elliptic flow

Assumption

- one dimensional
- Dirac delta for Wigner function
- isotropic radial flow
- not including resonance effect

Baryon to meson ratio at RHIC



- coalescence naturally predict a baryon/meson enhancement in the region $p_T \simeq 2-4\text{GeV}$ with respect to pp collisions
- Lack of baryon yield in the region $p_T \simeq 5-7\text{GeV}$

Relativistic Boltzmann transport at finite η/s

Bulk evolution

$$p^\mu \partial_\mu f_{q,g}(x, p) + M(x) \partial_x^\mu M(x) \partial_p^\mu f_{q,g}(x, p) = C_{22}[f_{q,g}]$$

free-streaming
field interaction
 $\varepsilon-3p\neq 0$
collisions
 $\eta\neq 0$

Heavy quark evolution

$$p^\mu \partial_\mu f_Q(x, p) = C[f_q, f_g, f_Q]$$

- Describes the evolution of the one body distribution function $f(x,p)$
 - It is valid to study the evolution of both bulk and Heavy quarks
 - Possible to include $f(x,p)$ out of equilibrium

