



ALMA MATER STUDIORUM  
UNIVERSITÀ DI BOLOGNA

**Beyond Iron peak:  
focus trasversale sulle sorgenti  
di neutroni  $^{13}\text{C}(\alpha, n)$  e  $^{22}\text{Ne}(\alpha, n)$ ,  
e sulla  $^{12}\text{C}+^{12}\text{C}$**

**Cristian Massimi**

Dipartimento di Fisica e Astronomia

Outline:

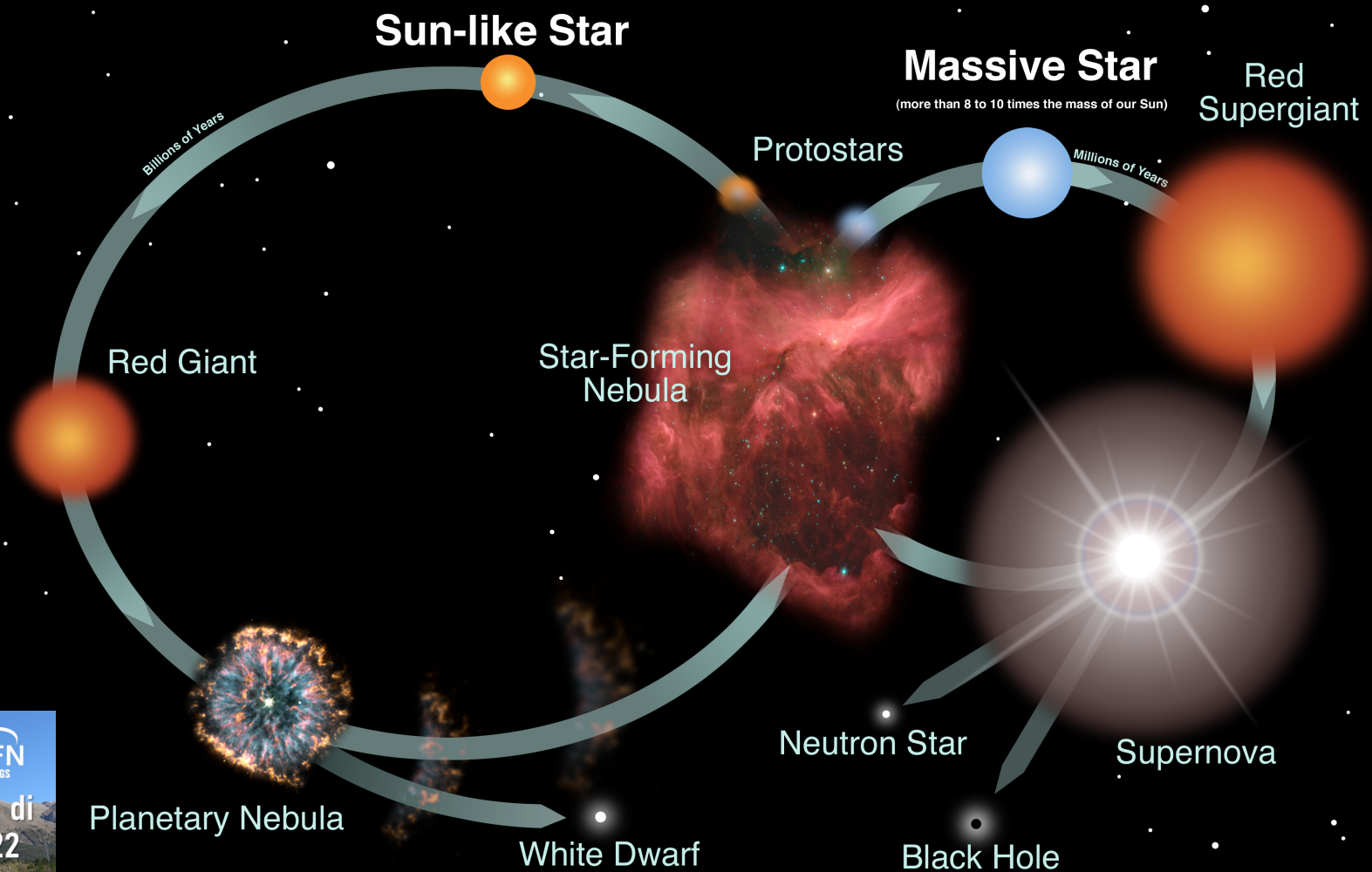
- $^{13}\text{C}(\alpha, n)$
- $^{22}\text{Ne}(\alpha, n)$
- $^{12}\text{C} + ^{12}\text{C}$



Outline:

- $^{13}\text{C}(\alpha,n)$
- $^{22}\text{Ne}(\alpha,n)$
- $^{12}\text{C} + ^{12}\text{C}$

In the cosmos,



Sun-like Star

Massive Star  
(more than 8 to 10 times the mass of our Sun)

Red Supergiant

Protostars

Millions of Years

Billions of Years

Red Giant

Star-Forming Nebula

Neutron Star

Supernova

Planetary Nebula

White Dwarf

Black Hole

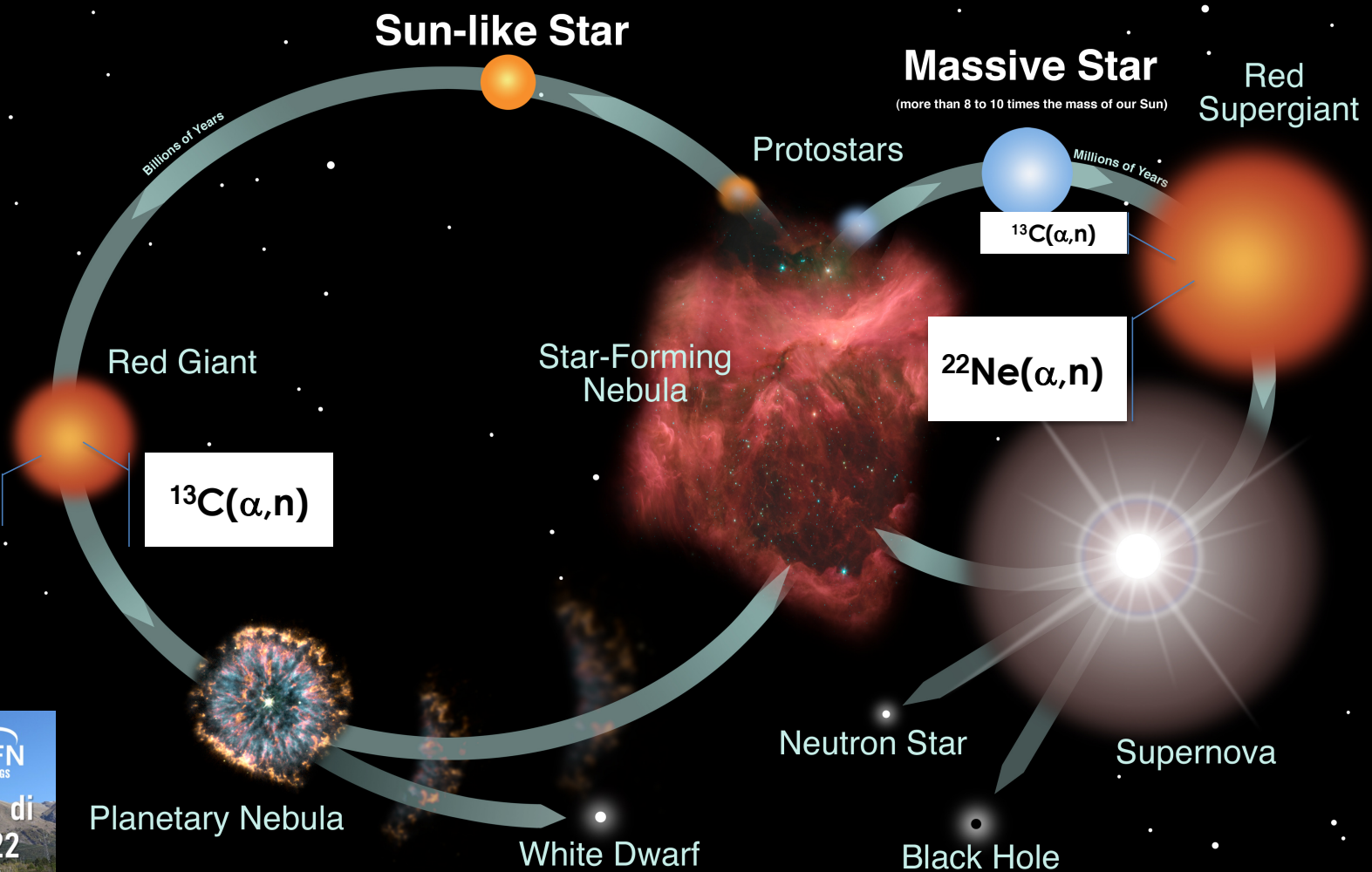


Credit: NASA and the Night Sky Network

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In the cosmos,

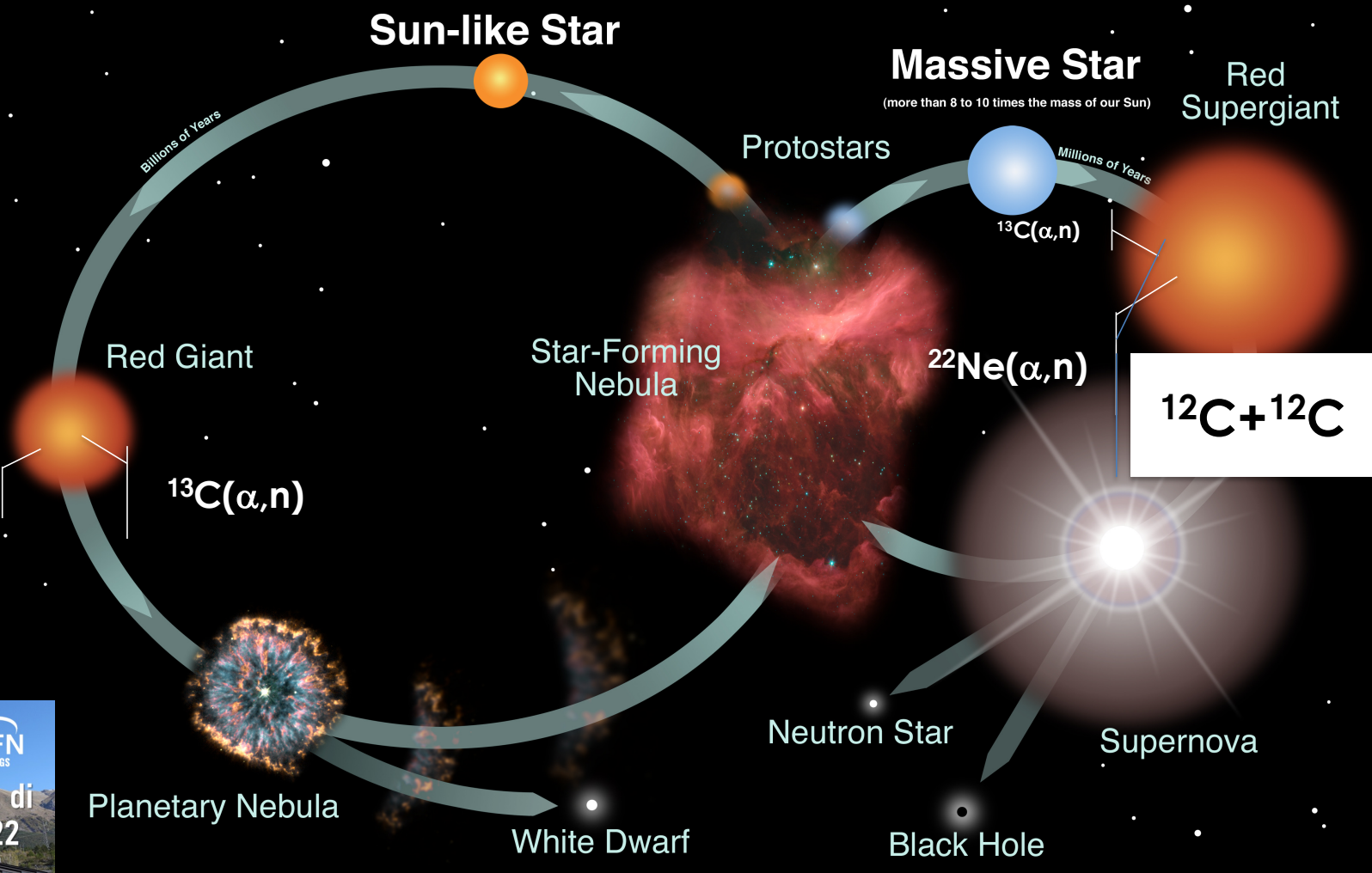


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In the cosmos,



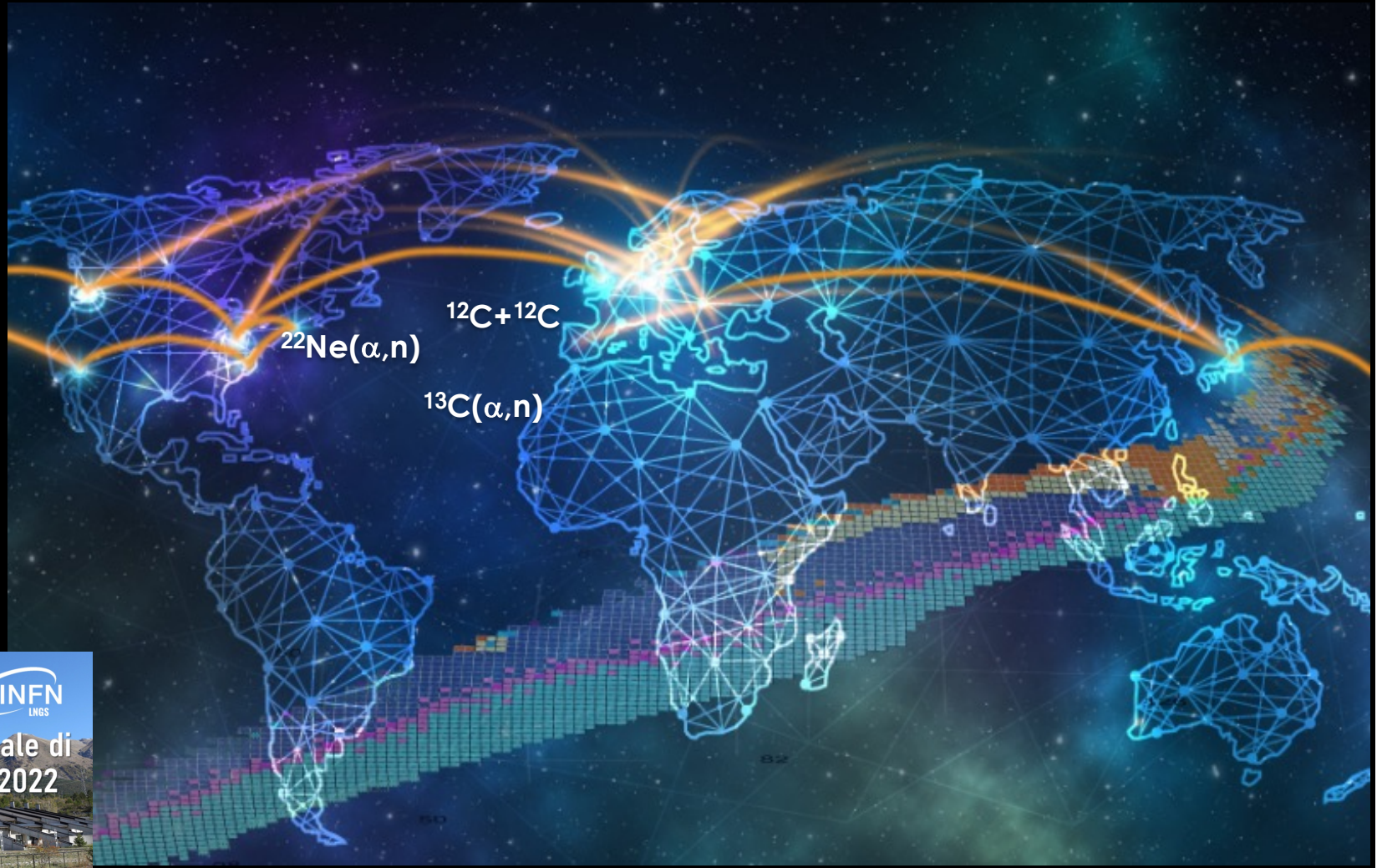
Credit: NASA and the Night Sky Network

Credit: Facility for Rare Isotope Beams

Outline:

- $^{13}\text{C}(\alpha,n)$
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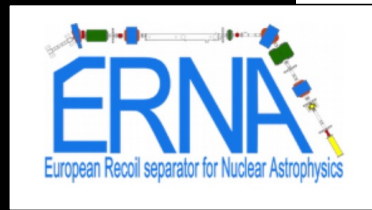
In the cosmos,  
On Earth (labs),



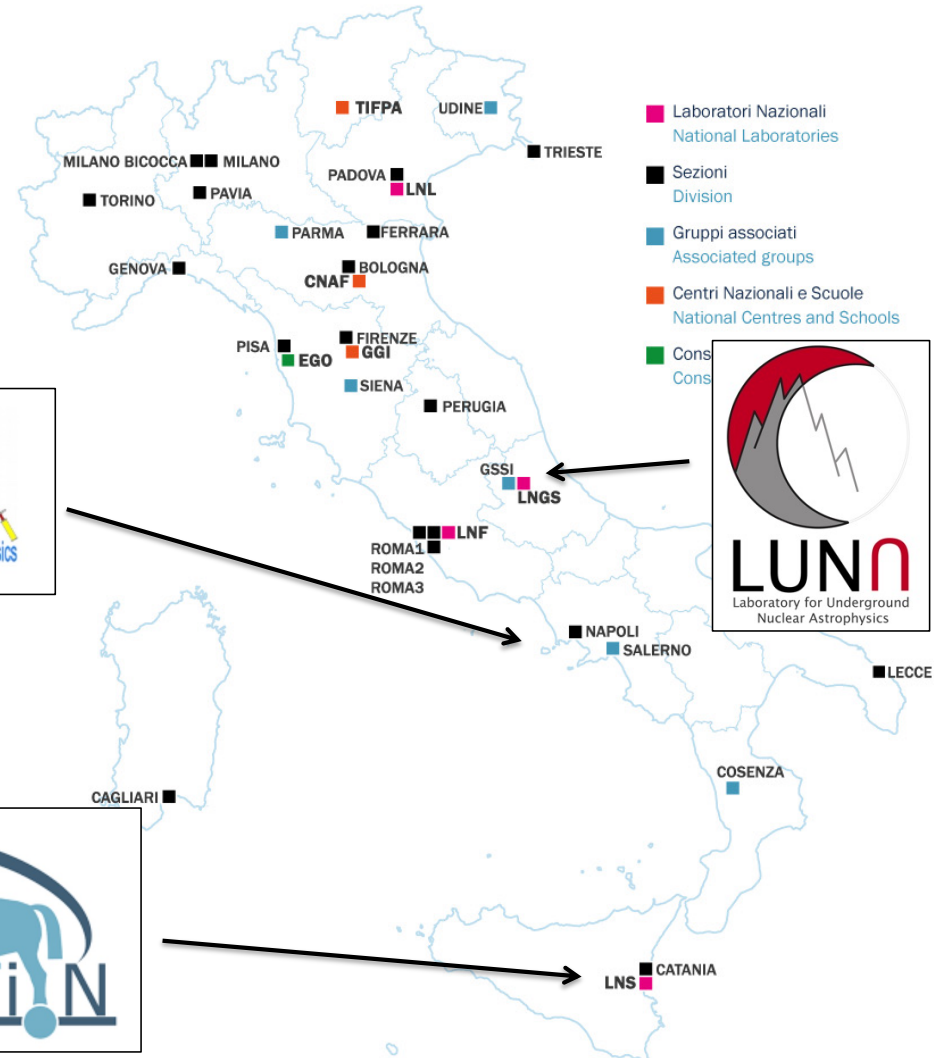
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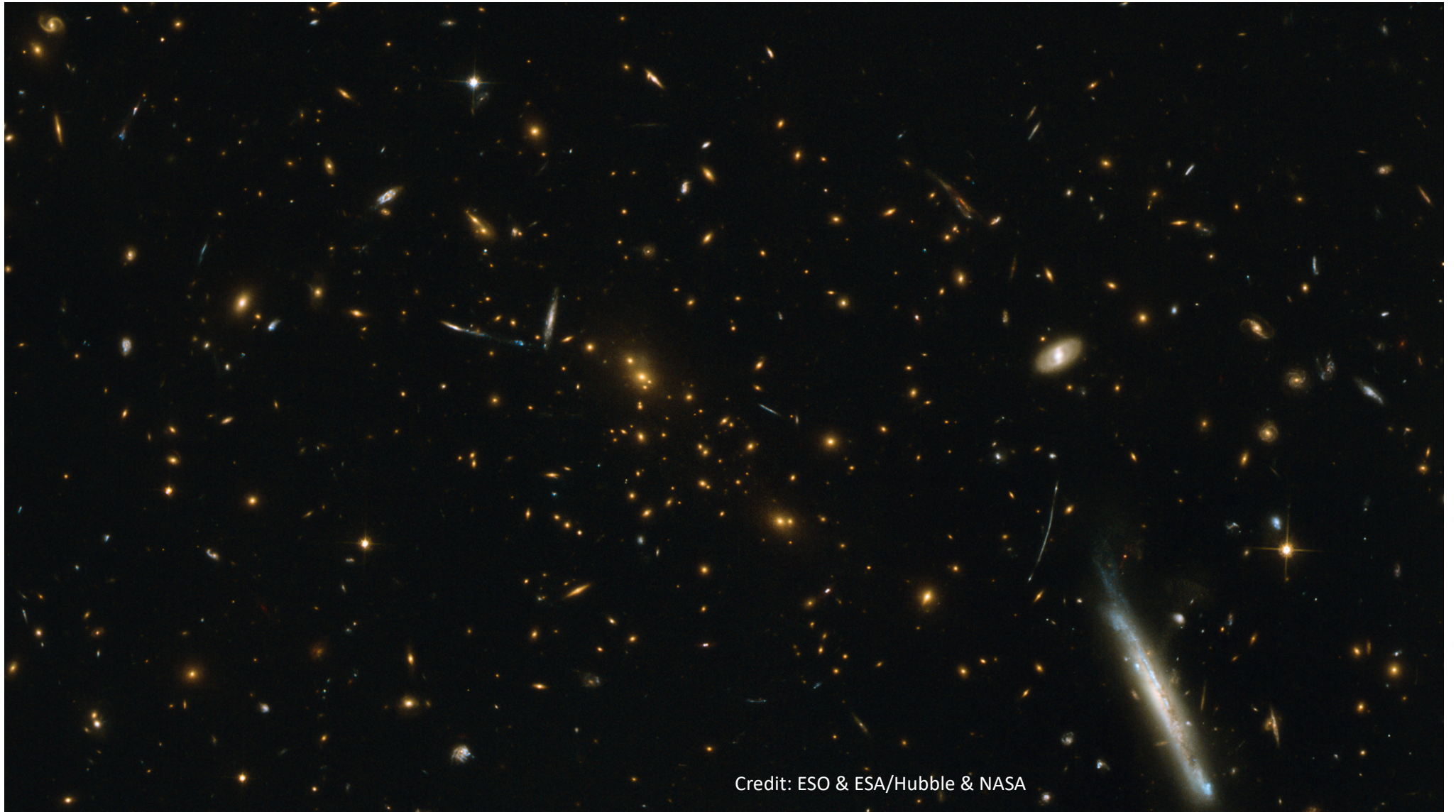
In the cosmos,  
On Earth (labs),  
@ INFN



*L'insocievole socievolezza  
... di GIANTS*



Uff. Comunicazione F. Cuicchio Powered by Multimedia Servic



Credit: ESO & ESA/Hubble & NASA



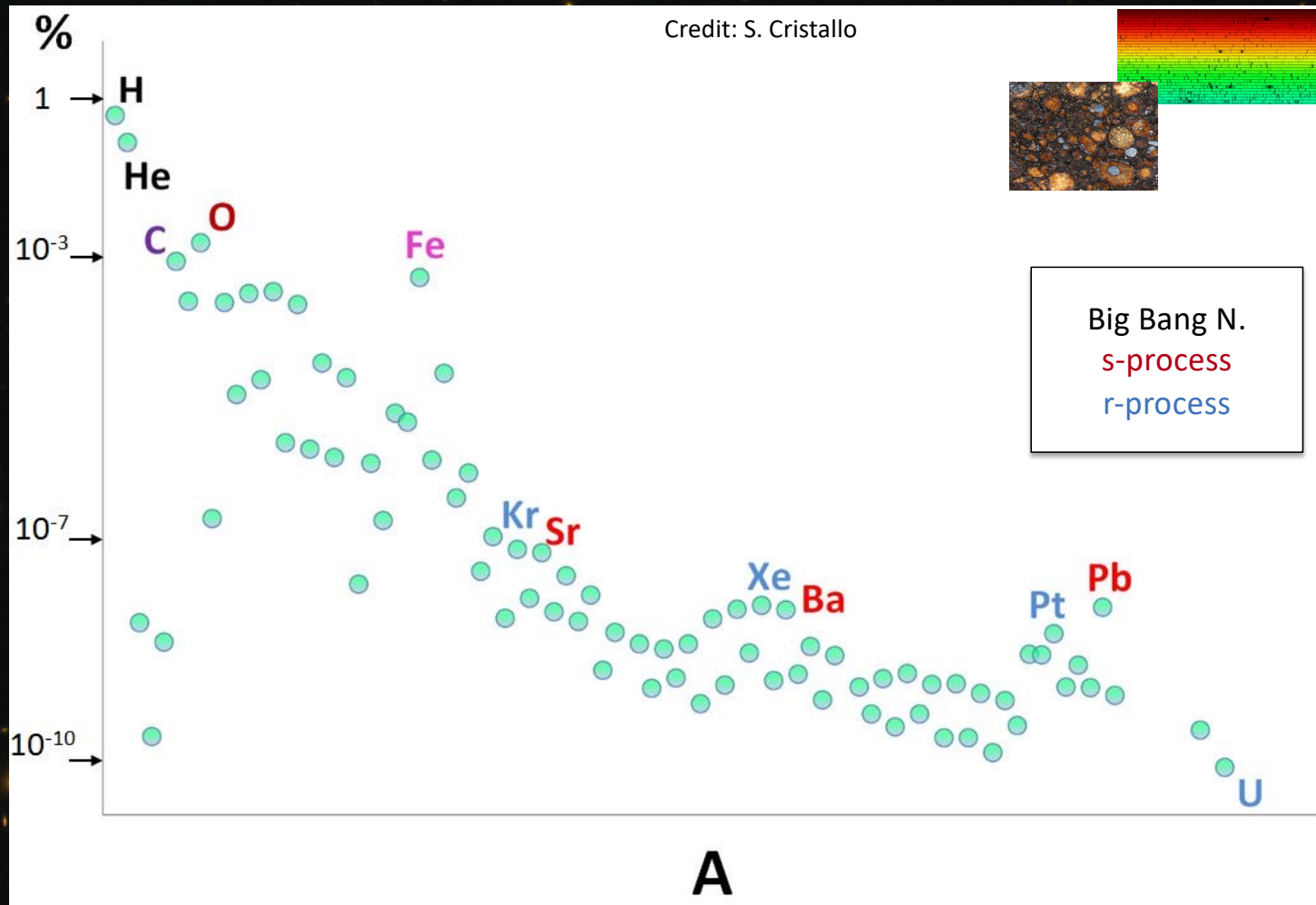
A deep field image of the universe, showing a vast field of galaxies and stars. The background is dark, with numerous bright points of light and some faint, elongated structures. The text is overlaid on the left side of the image.

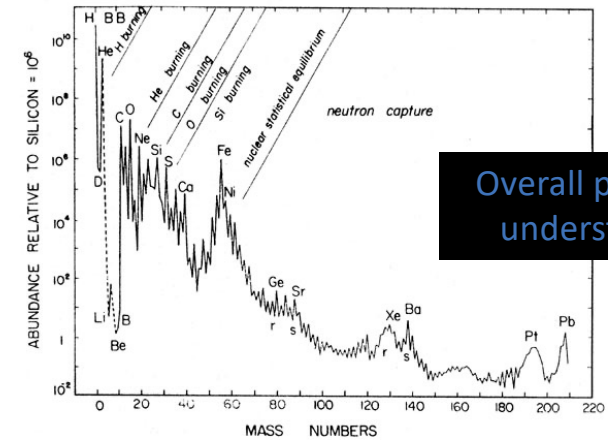
All visible objects in the night sky are powered by nuclear reactions

Nuclear physics governs the evolution of stars from birth to their final fate

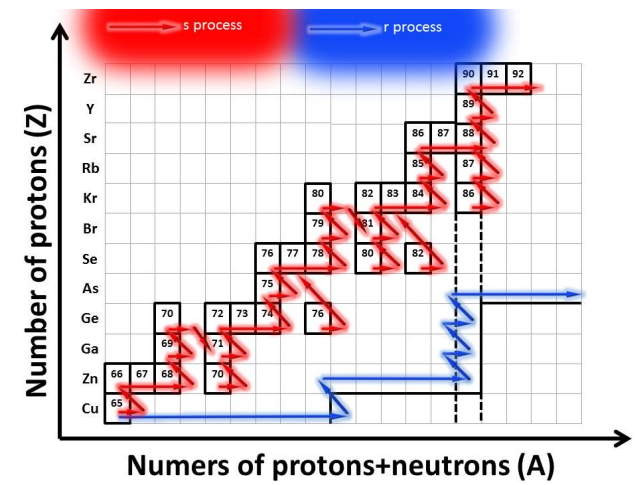
Nuclear reactions in the Big Bang, in stars and in stellar explosions have created every single chemical element found in nature today

Credit: ESO & ESA/Hubble & NASA



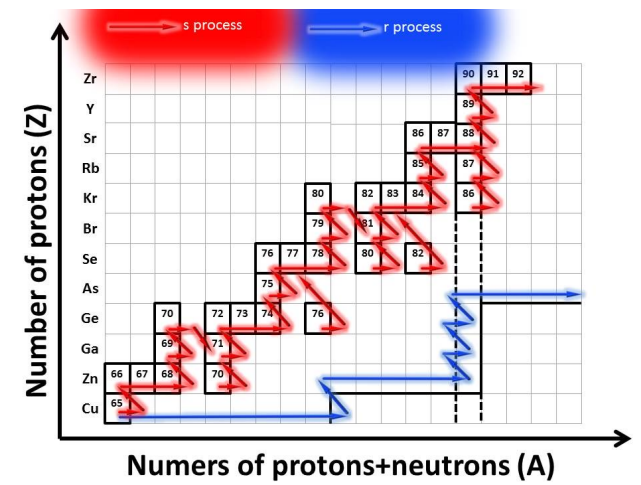
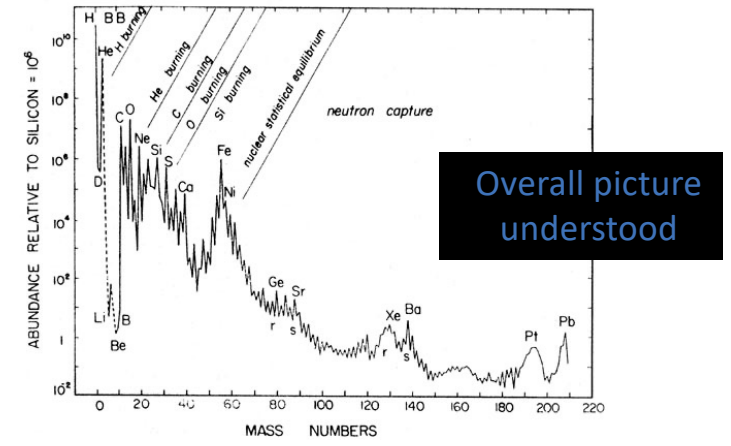


Overall picture understood



## (Some) open questions

- ❖  ${}^7\text{Li}$  produced in Big Bang Nucleosynthesis
- ❖ Elements in the **Sun** beyond He (**metallicity**)
- ❖ **r process** (Main & Weak ?)
- ❖ **r process site** (Supernovae or Neutron Star mergers, ...)
- ❖ Stellar explosion (basic mechanism)
- ❖ X-ray bursts
- ❖ ...

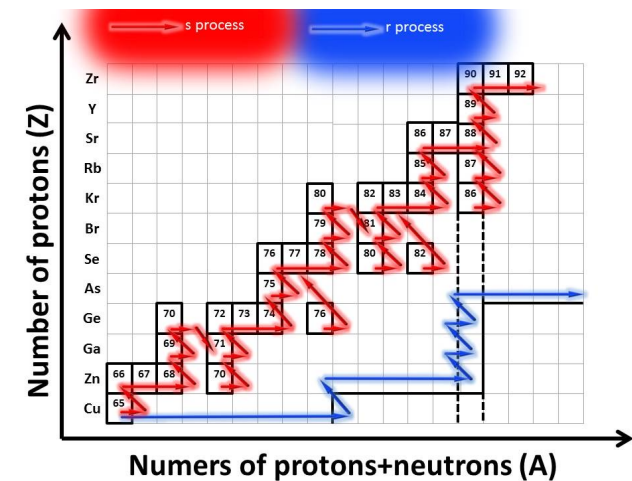
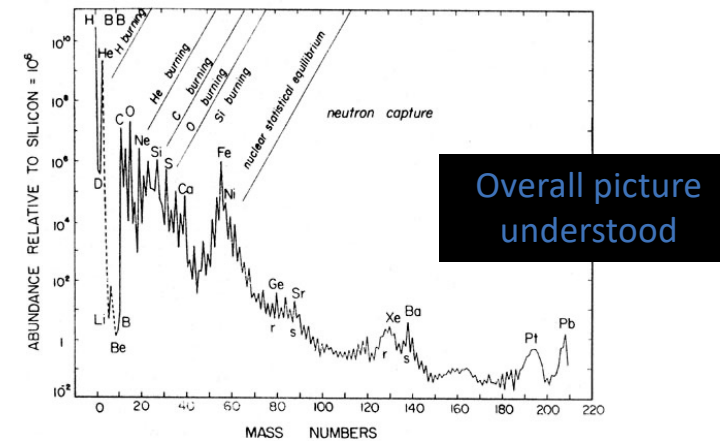


## (Some) open questions

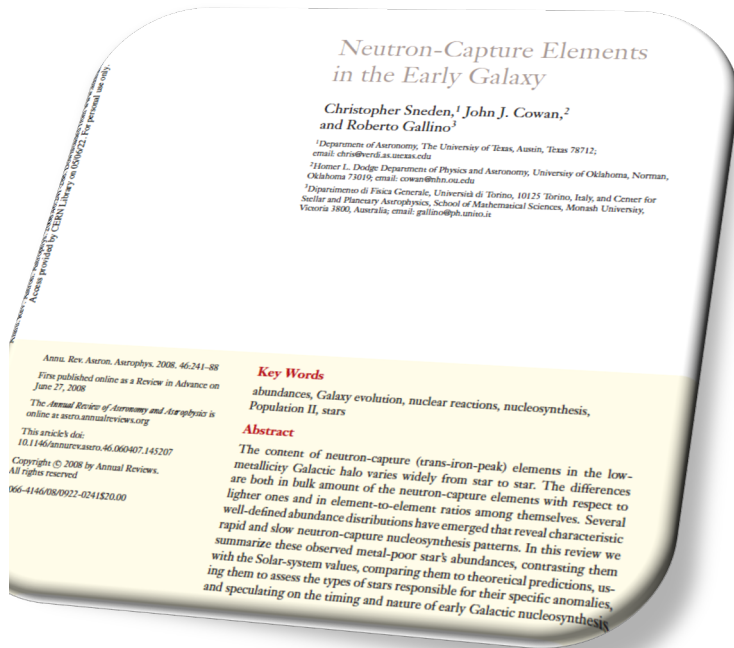
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- ❖ ...

### 3 main drivers pushing the frontiers:

- **Multi-messenger observations** force rethinking our paradigms (e.g. detection of surprisingly massive neutron stars)
- **Multi-facility nuclear experiments** and advances in nuclear theory
- Expansion of **computational capabilities** enables 3D simulations of stellar interiors



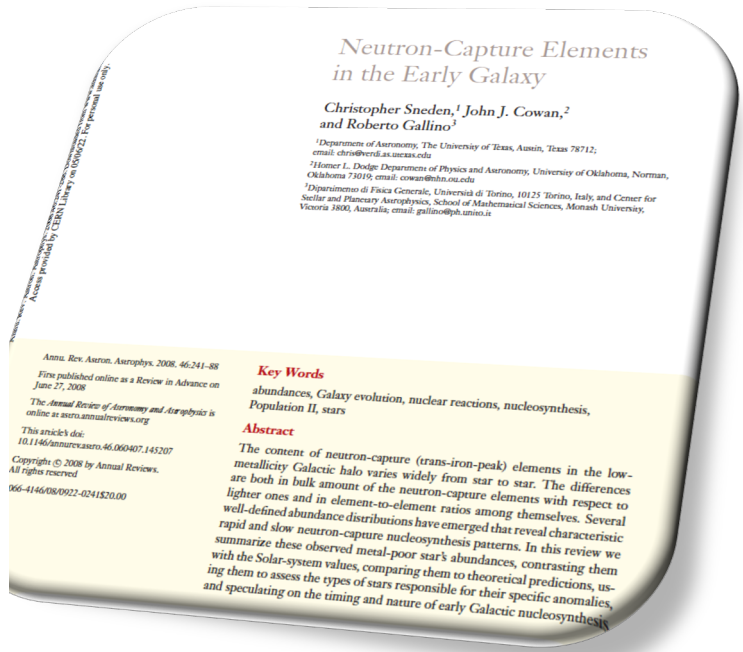
# The role of s-process, an example



$$r = 1 - s$$

↑  
Solar

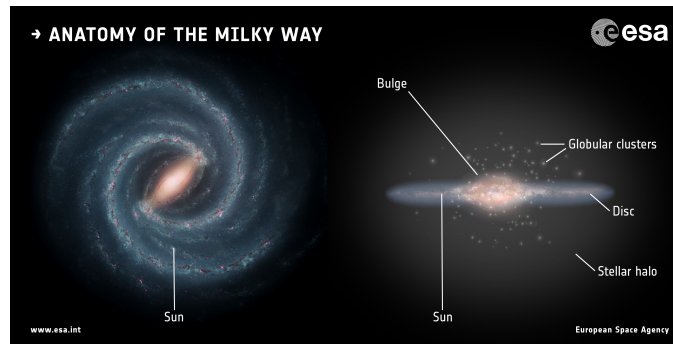
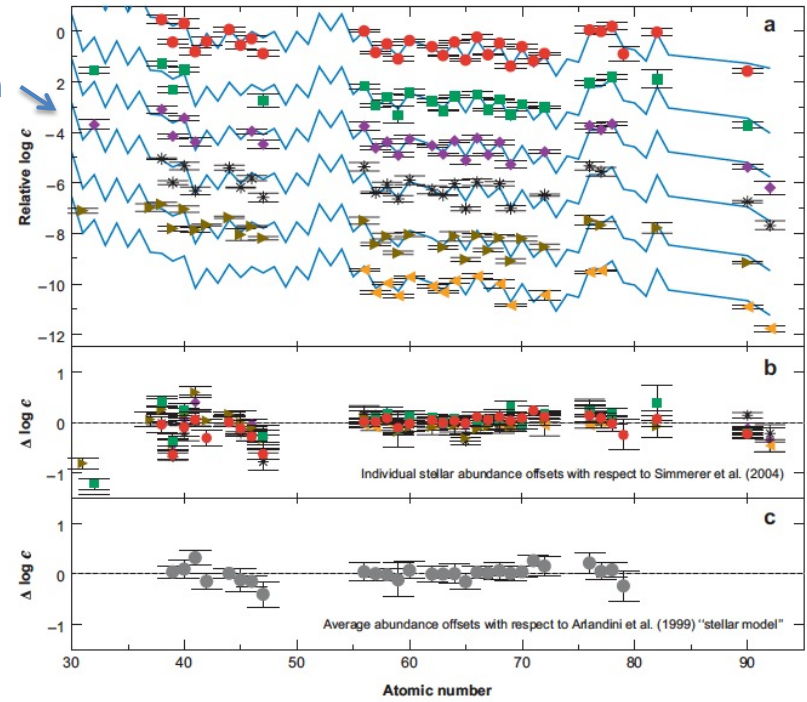
# The role of s-process, an example



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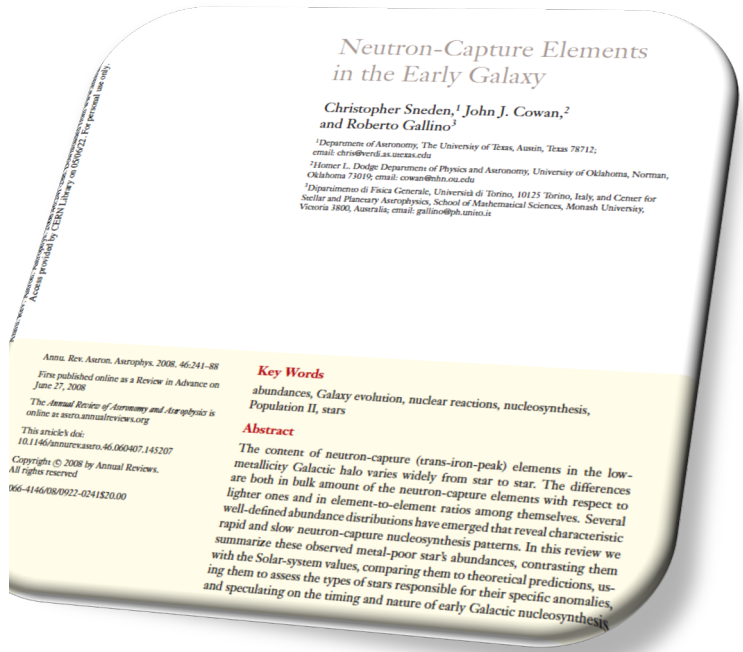
Solar-system r-only abundance distribution



r-process rich Galactic halo stars

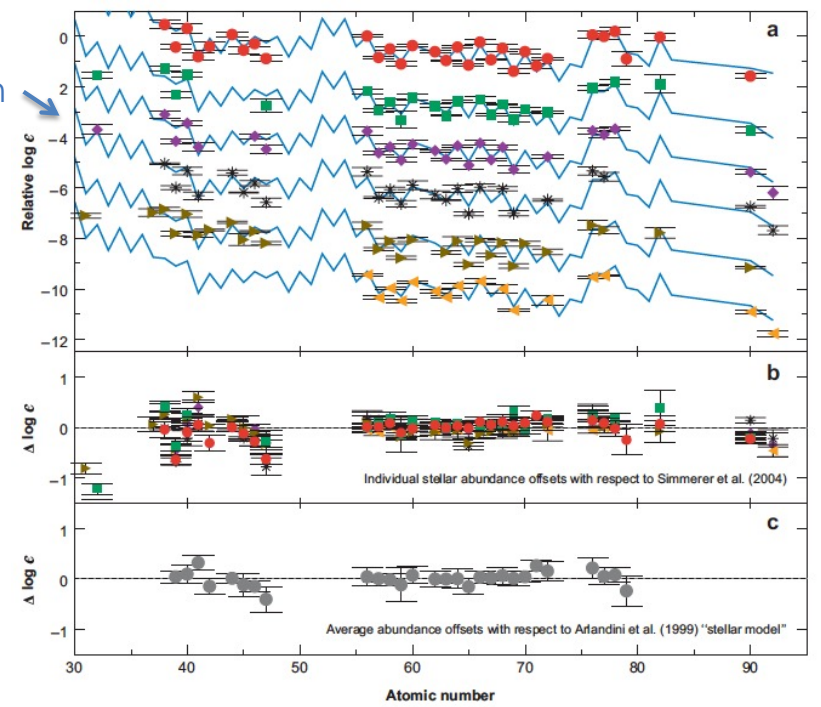
- CS 22892-052: Sneden et al. (2003)
- HD 115444: Westin et al. (2000)
- ★ BD+17°324817: Cowan et al. (2002)
- \* CS 31082-001: Hill et al. (2002)
- ▲ HD 221170: Ivans et al. (2006)
- ◆ HE 1523-0901: Frebel et al. (2007)

# The role of s-process, an example

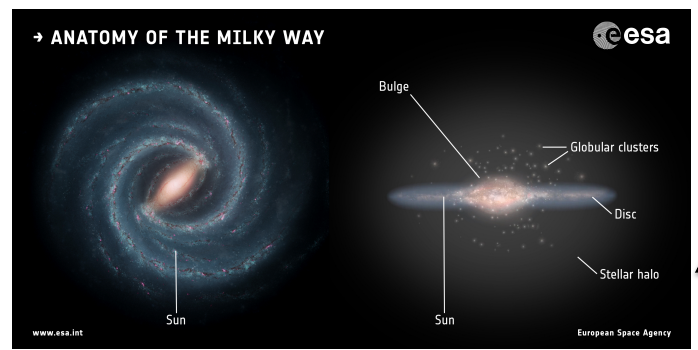


Solar-system r-only abundance distribution

Detailed knowledge of the s-process nucleosynthesis is the prerequisite for the study of the chemical evolution of the Galaxy



$r = 1 - s$   
 ↑  
 Solar



r-process rich Galactic halo stars

- CS 22892-052: Sneden et al. (2003)
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massimi@bo.infn.it

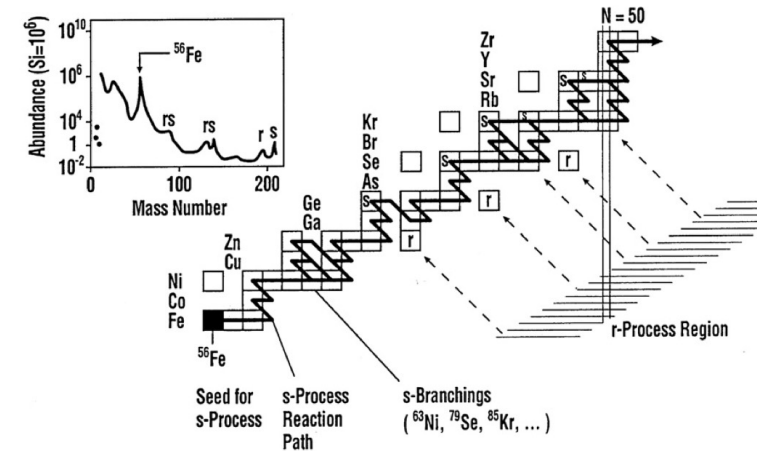


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## s process, what matters?

- ❖ Reaction rate of the two main neutron sources  $^{13}\text{C}(\alpha,n)^{16}\text{O}$  and  $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$
- ❖  $(n,\gamma)$  cross sections with small (5%) uncertainty
- ❖  $(n,\gamma)$  cross sections of branching-point isotopes and  $\beta$ -decay rates
- ❖ Stellar models
- ❖ Stellar enhancement factors



REVIEWS OF MODERN PHYSICS The s process: Nuclear physics, stellar models, and observations, Vol. 83, 2011

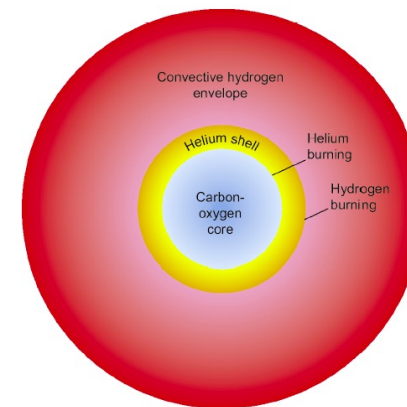
Sample	Half-life (yr)	Q value (MeV)	Comment
$^{63}\text{Ni}$	100.1	$\beta^-$ , 0.066	TOF work in progress (Couture, 2009), sample with low enrichment. PRL 110, 022501 (2013)
$^{79}\text{Se}$	$2.95 \times 10^5$	$\beta^-$ , 0.159	Important branching, constrains s-process temperature in massive stars
$^{81}\text{Kr}$	$2.29 \times 10^5$	EC, 0.322	Part of $^{79}\text{Se}$ branching
$^{85}\text{Kr}$	10.73	$\beta^-$ , 0.687	Important branching, constrains neutron density in massive stars
$^{95}\text{Zr}$	64.02 d	$\beta^-$ , 1.125	Not feasible in near future, but important for neutron density low-mass AGB stars
$^{134}\text{Cs}$	2.0652	$\beta^-$ , 2.059	Important branching at $A = 134, 135$ , sensitive to s-process temperature in low-mass AGB stars, measurement not feasible in near future
$^{135}\text{Cs}$	$2.3 \times 10^6$	$\beta^-$ , 0.269	So far only activation measurement at $kT = 25$ keV by Patronis <i>et al.</i> (2004)
$^{147}\text{Nd}$	10.981 d	$\beta^-$ , 0.896	Important branching at $A = 147/148$ , constrains neutron density in low-mass AGB stars
$^{149}\text{Pm}$	2.6234	$\beta^-$ , 0.225	Part of branching at $A = 147/148$
$^{150}\text{Pm}$	5.368 d	$\beta^-$ , 2.464	Not feasible in the near future
$^{151}\text{Sm}$	90	$\beta^-$ , 0.076	Existing TOF measurements, full set of MACS data available (Abbondanno <i>et al.</i> , 2004a; Wisshak <i>et al.</i> , 2006c) PRL 93, 161103 (2004)
$^{154}\text{Eu}$	8.593	$\beta^-$ , 1.978	Complex branching at $A = 154, 155$ , sensitive to temperature and density
$^{155}\text{Eu}$	4.753	$\beta^-$ , 0.246	So far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1995)
$^{153}\text{Gd}$	0.658	EC, 0.244	Part of branching at $A = 154, 155$
$^{160}\text{Tb}$	0.198	$\beta^-$ , 1.833	Weak temperature-sensitive branching, very challenging experiment
$^{163}\text{Ho}$	4570	EC, 0.0026	Branching at $A = 163$ sensitive to mass density during s process, so far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1996b)
$^{170}\text{Tm}$	0.352	$\beta^-$ , 0.968	Important branching, constrains neutron density in low-mass AGB stars
$^{171}\text{Tm}$	1.921	$\beta^-$ , 0.098	Part of branching at $A = 170, 171$
$^{179}\text{Ta}$	1.82	EC, 0.115	Crucial for s-process contribution to $^{180}\text{Ta}$ , nature's rarest stable isotope
$^{185}\text{W}$	0.206	$\beta^-$ , 0.432	Important branching, sensitive to neutron density and s-process temperature in low-mass AGB stars
$^{204}\text{Tl}$	3.78	$\beta^-$ , 0.763	Determines $^{205}\text{Pb}/^{205}\text{Tl}$ clock for dating of early Solar System





❖ Main neutron source in low-mass AGB stars at temperature  $\sim 90\text{-}100\text{ MK}$

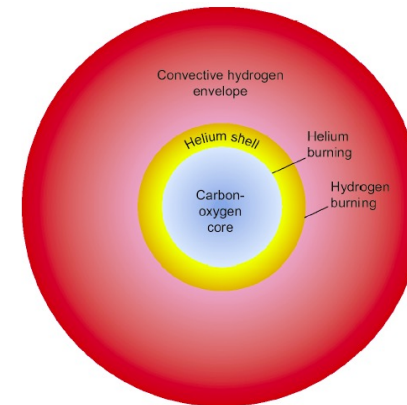
❖ Gamow window  $\sim 150\text{-}230\text{ keV}$



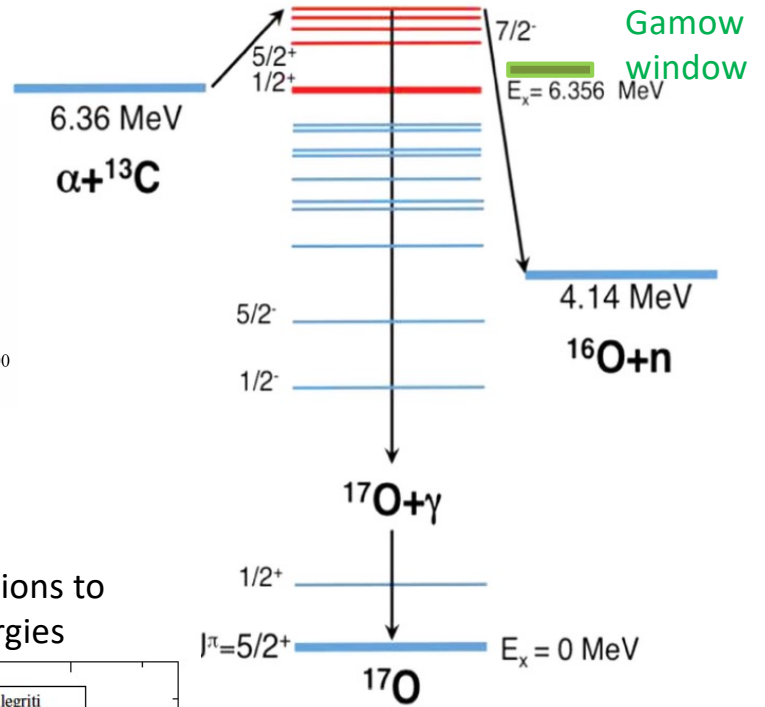
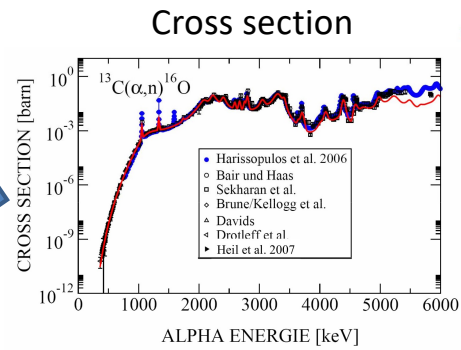
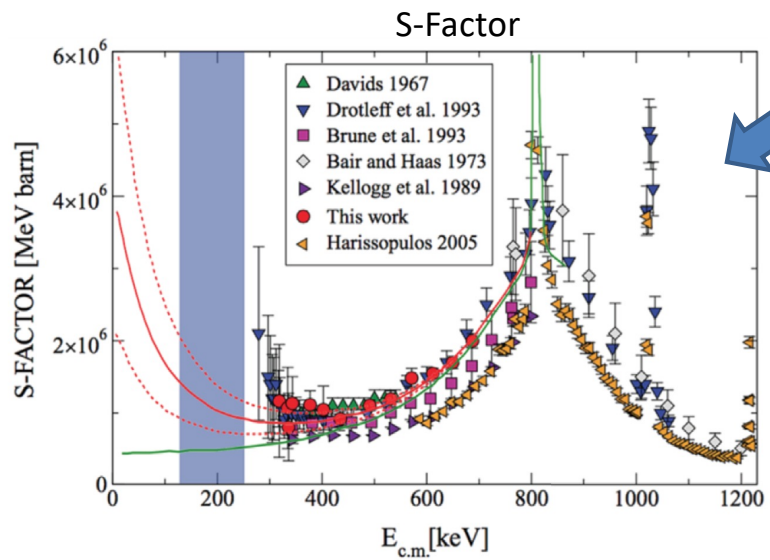


❖ Main neutron source in low-mass AGB stars at temperature  $\sim 90\text{-}100\text{ MK}$

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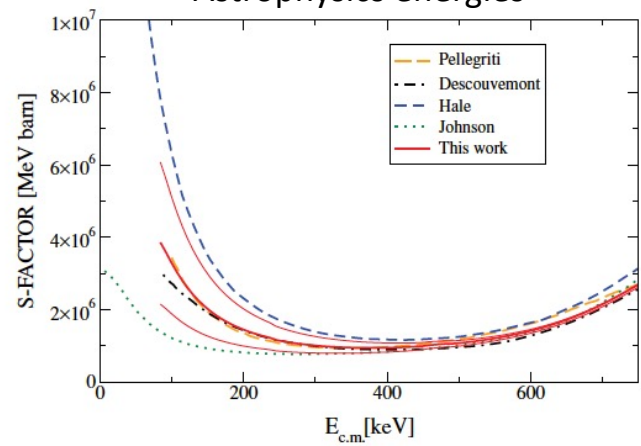


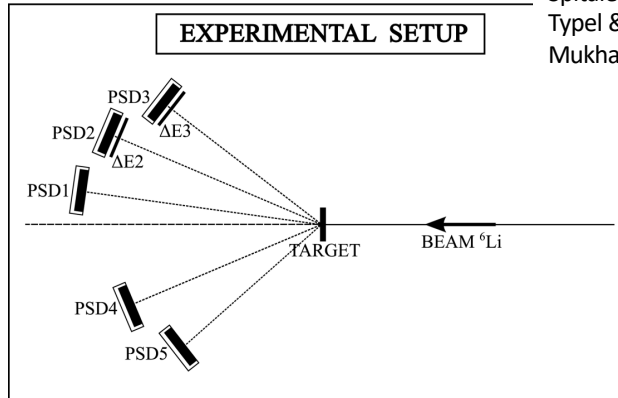
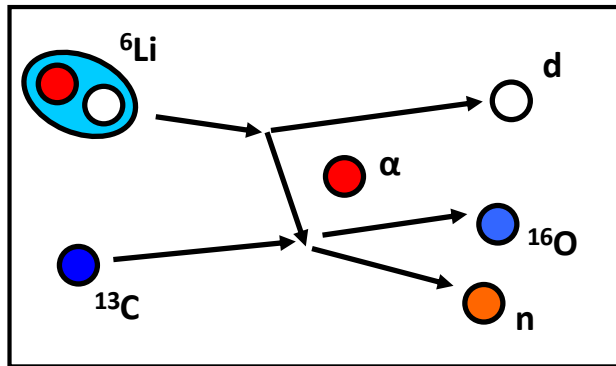
# $^{13}\text{C}(\alpha, n)^{16}\text{O}$ "BEFORE"



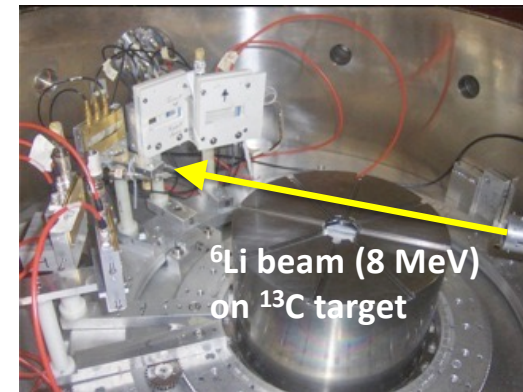
The broad  $1/2+$  state at 6.356 keV plays a major role

## Tentative extrapolations to Astrophysics energies





Baur PLB 178 (1986) 135  
 Spitaleri, in *Problems of Fundamental Modern Physics II*, World Sci. (1991) 21  
 Spitaleri et al., PAN 74 (2011) 1763  
 Typel & Baur, Ann. Phys. 305 (2003) 228  
 Mukhamedzhanov et al., JPG 35 (2008) 014016



Measurement of the the sub-Coulomb  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction through the  $^{13}\text{C}(^6\text{Li}, n^{16}\text{O})d$  reaction in the quasi-free kinematics regime @ Florida State University.



- Charged particle detection (instead of neutron)
- Measurement at  $\sim 10 - 60$  MeV kinetic energy ( $\gg$  Coulomb barrier)
- No electron screening



- Model dependent (from indirect to direct)
- Normalization needed

Measurement of the  $-3$  keV Resonance in the Reaction  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  of Importance in the s-Process

M. La Cognata,<sup>1,4</sup> C. Spitaleri,<sup>1,2</sup> O. Trippella,<sup>1,3</sup> G. G. Kiss,<sup>1,4</sup> G. V. Rogachev,<sup>5</sup> A. M. Mukhamedzhanov,<sup>6</sup> M. Avila,<sup>5</sup> G. L. Guardo,<sup>1,2</sup> E. Koshchiy,<sup>5</sup> A. Kuchera,<sup>5</sup> L. Lamia,<sup>2</sup> S. M. R. Puglia,<sup>1,2</sup> S. Romano,<sup>1,2</sup> D. Santiago,<sup>5</sup> and R. Sparta<sup>1,2</sup>  
<sup>1</sup>Istituto Nazionale di Fisica Nucleare, Laboratori Nazionali del Sud, 95123 Catania, Italy  
<sup>2</sup>Dipartimento di Fisica e Astronomia, Università di Catania, 95123 Catania, Italy  
<sup>3</sup>Istituto Nazionale di Fisica Nucleare, Sezione di Perugia, Perugia, Italy and Dipartimento di Fisica, Università di Perugia, 06123 Perugia, Italy  
<sup>4</sup>Institute of Nuclear Research (ATOMKI), 4026 Debrecen, Hungary  
<sup>5</sup>Department of Physics, Florida State University, Tallahassee, Florida 32306, USA  
<sup>6</sup>Cyclotron Institute, Texas A&M University, College Station, Texas 77843, USA  
 (Received 3 July 2012; revised manuscript received 10 August 2012; published 4 December 2012)



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# $^{13}\text{C}(\alpha, n)^{16}\text{O}$



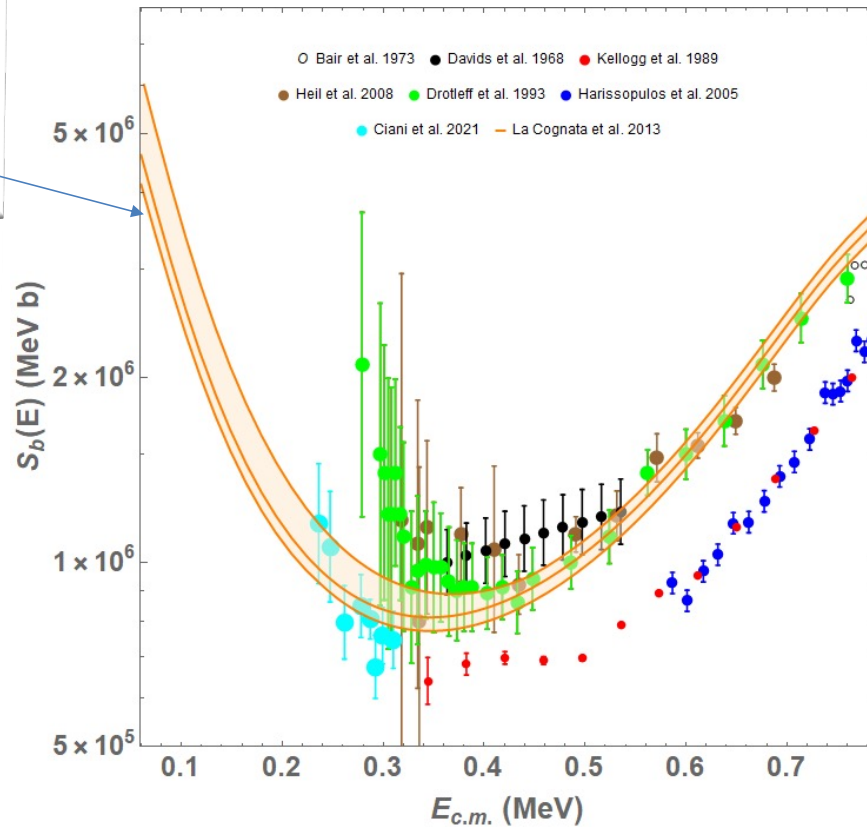
THE ASTROPHYSICAL JOURNAL, 837:41 (12pp), 2017 March 1  
 © 2017, The American Astronomical Society. All rights reserved. <https://doi.org/10.3847/1538-4357/aa5cb5>

**Concurrent Application of ANC and THM to assess the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  Absolute Cross Section at Astrophysical Energies and Possible Consequences for Neutron Production in Low-mass AGB Stars**

O. Trippella<sup>1,2</sup> and M. La Cognata<sup>3</sup>  
<sup>1</sup> Department of Physics and Geology, University of Perugia, Perugia, Italy; [oscar.trippella@pg.infn.it](mailto:oscar.trippella@pg.infn.it)  
<sup>2</sup> Istituto Nazionale di Fisica Nucleare, Section of Perugia, Perugia, Italy  
<sup>3</sup> Istituto Nazionale di Fisica Nucleare, Laboratori Nazionali del Sud, Catania, Italy; [lacognata@lns.infn.it](mailto:lacognata@lns.infn.it)  
 Received 2016 December 16; revised 2017 February 3; accepted 2017 February 3; published 2017 March 1

Good agreement with direct data from the LUNA collaboration (and some previous extrapolations).

Towards a concordance scenario for the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  S(E)



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[massimi@bo.infn.it](mailto:massimi@bo.infn.it)

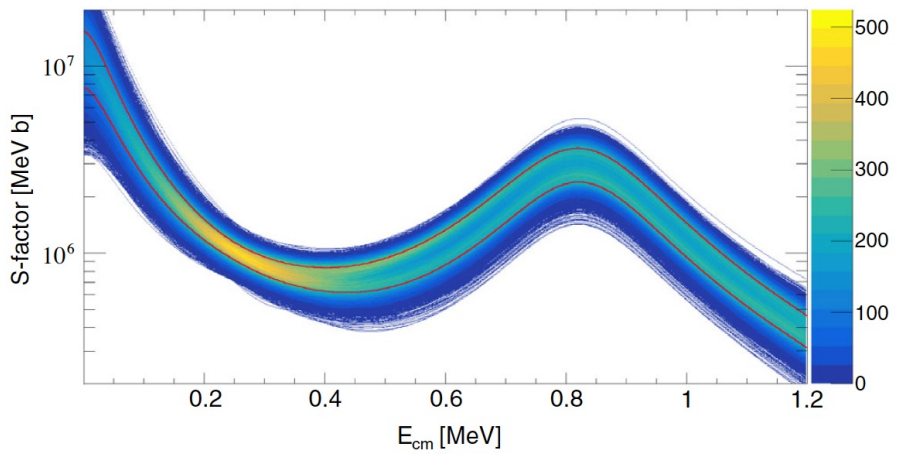
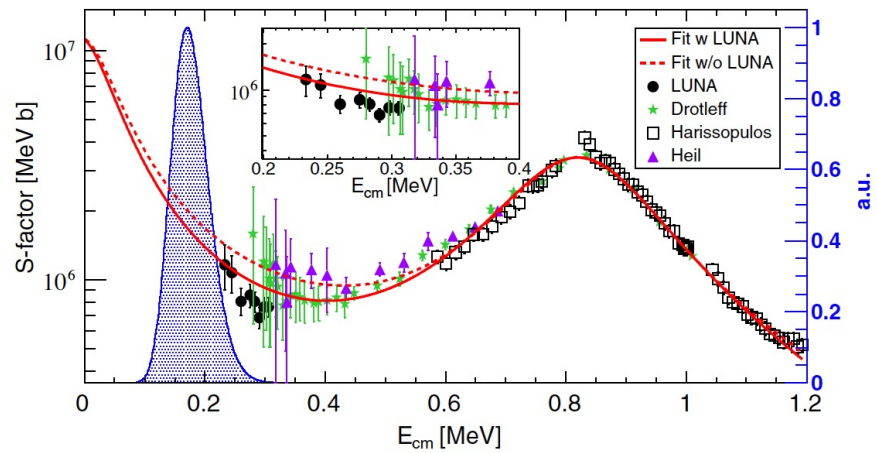


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# $^{13}\text{C}(\alpha, n)^{16}\text{O}$



PHYSICAL REVIEW LETTERS 127, 152701 (2021)

**Direct Measurement of the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  Cross Section into the s-Process Gamow Peak**

G. F. Ciani,<sup>1,2,3</sup> L. Csedreki,<sup>1,2,3</sup> D. Rapagnani,<sup>4,5</sup> M. Aliotta,<sup>6</sup> J. Balibrea-Correa,<sup>4,5</sup> F. Barile,<sup>7,8</sup> D. Bemmerer,<sup>12</sup> A. Best,<sup>4,5,\*</sup> A. Boeltzig,<sup>4,5</sup> C. Broggini,<sup>10</sup> C. G. Bruno,<sup>6</sup> A. Cacioli,<sup>10,11</sup> F. Cavanna,<sup>12</sup> T. Chillery,<sup>6</sup> P. Colombetti,<sup>12</sup> P. Corvisiero,<sup>13,14</sup> S. Cristallo,<sup>15,16</sup> T. Davinson,<sup>6</sup> R. Depalo,<sup>11,10</sup> A. Di Leva,<sup>4,5</sup> Z. Elekes,<sup>3</sup> F. Ferraro,<sup>13,14</sup> E. Fiore,<sup>7,8</sup> A. Formicola,<sup>2,†</sup> Zs. Fülöp,<sup>3</sup> G. Gervino,<sup>17,18</sup> A. Guglielmetti,<sup>19,20</sup> C. Gustavino,<sup>21</sup> Gy. Gyürky,<sup>3</sup> G. Imbriani,<sup>4,5</sup> M. Junker,<sup>2</sup> M. Lugaro,<sup>22,23</sup> P. Marigo,<sup>10,11</sup> E. Masha,<sup>19,20</sup> R. Menegazzo,<sup>10</sup> V. Mossa,<sup>8</sup> F.R. Pantaleo,<sup>7,8</sup> V. Paticchio,<sup>8</sup> R. Perrino,<sup>8,‡</sup> D. Piatti,<sup>10,11</sup> P. Prati,<sup>13,14</sup> L. Schiavulli,<sup>7,8</sup> K. Stöckel,<sup>9,24</sup> O. Straniero,<sup>15,2</sup> T. Szücs,<sup>3</sup> M.P. Takács,<sup>9,24</sup> F. Terrasi,<sup>25,5</sup> D. Vescovi,<sup>16,26</sup> and S. Zavatarelli<sup>14</sup>

(LUNA Collaboration)



- ❖ Deep-underground measurement
- ❖ Energy range 230 –300 keV
- ❖ Drastically reduced uncertainties
- ❖ First time providing data directly inside the s-process Gamow peak

[massimi@bo.infn.it](mailto:massimi@bo.infn.it)



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# $^{13}\text{C}(\alpha, n)^{16}\text{O}$

Eur. Phys. J. A (2020) 56:75  
 https://doi.org/10.1140/epja/s10050-020-00077-0

THE EUROPEAN PHYSICAL JOURNAL A

Special Article - New Tools and Techniques

## A new approach to monitor $^{13}\text{C}$ -targets degradation in situ for $^{13}\text{C}(\alpha, n)^{16}\text{O}$ cross-section measurements at LUNA

G. F. Ciani<sup>1,2,3,a</sup>, L. Csedreki<sup>1,2,b</sup>, J. Balibrea-Correa<sup>4,5</sup>, A. Best<sup>4,5</sup>, M. Aliotta<sup>6</sup>, F. Barile<sup>7</sup>, D. Bemmerer<sup>8</sup>, A. Boeltzig<sup>1,2</sup>, C. Broggini<sup>9</sup>, C. G. Bruno<sup>6</sup>, A. Cacioli<sup>10</sup>, F. Cavanna<sup>11</sup>, T. Chillery<sup>6</sup>, P. Colombetti<sup>12,13</sup>, P. Corvisiero<sup>11,14</sup>, T. Davinson<sup>6</sup>, R. Depalo<sup>6</sup>, A. Di Leva<sup>4,5</sup>, L. Di Paolo<sup>2</sup>, Z. Elekes<sup>6</sup>, F. Ferraro<sup>11,14</sup>, E. M. Fiore<sup>15</sup>, A. Formicola<sup>6</sup>, Zs. Fülöp<sup>3</sup>, G. Gervino<sup>12,13</sup>, A. Guglielmetti<sup>16,17</sup>, C. Gustavino<sup>18</sup>, Gy. Gyürky<sup>3</sup>, G. Imbriani<sup>4,5</sup>, M. Junker<sup>2</sup>, I. Kochanek<sup>19</sup>, M. Lugaro<sup>19</sup>, P. Marigo<sup>9,10</sup>, E. Masha<sup>16,17</sup>, R. Menegazzo<sup>6</sup>, V. Mossa<sup>1</sup>, F. R. Pantaleo<sup>7,20</sup>, V. Paticchio<sup>7</sup>, R. Perrino<sup>7,24</sup>, D. Piatti<sup>9,10</sup>, P. Prati<sup>11,14</sup>, L. Schiavulli<sup>7,19</sup>, K. Stöckel<sup>21</sup>, O. Straniero<sup>7,22</sup>, T. Szics<sup>6</sup>, M. P. Takács<sup>8,21,25</sup>, F. Terrasi<sup>23</sup>, D. Trezzi<sup>16,17</sup>, S. Zavatarelli<sup>11</sup>



Nuclear Inst. and Methods in Physics Research, A 994 (2021) 165081

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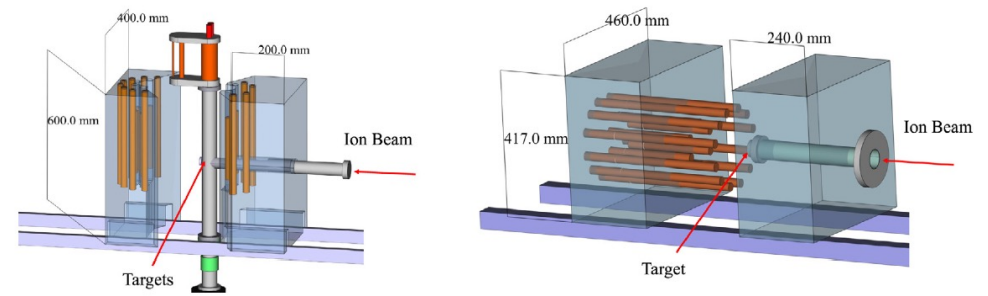
Nuclear Inst. and Methods in Physics Research, A

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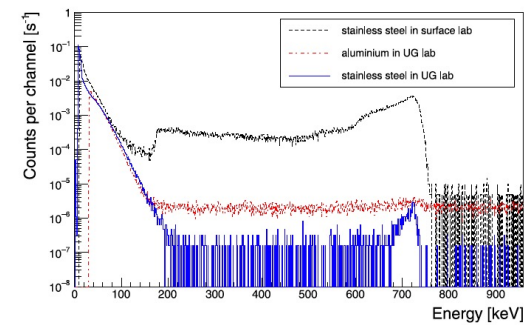
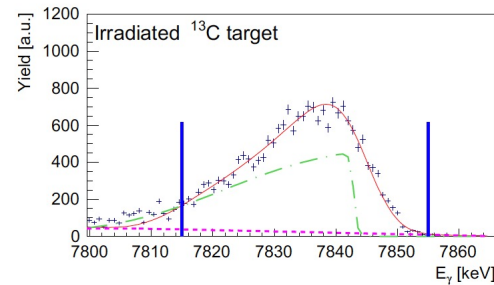
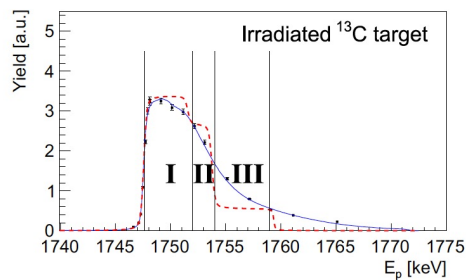
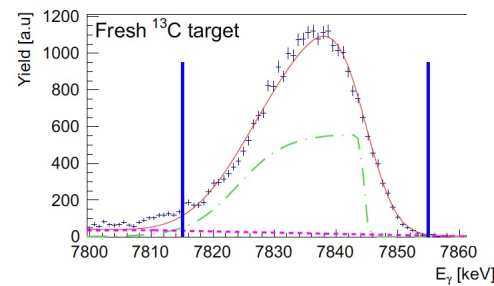
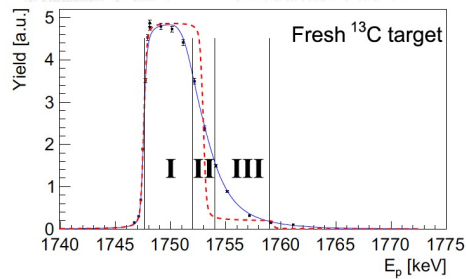
## Characterization of the LUNA neutron detector array for the measurement of the $^{13}\text{C}(\alpha, n)^{16}\text{O}$ reaction

L. Csedreki<sup>a,b,c</sup>, G.F. Ciani<sup>a,b,c</sup>, J. Balibrea-Correa<sup>d</sup>, A. Best<sup>d</sup>, M. Aliotta<sup>e</sup>, F. Barile<sup>f</sup>, D. Bemmerer<sup>g</sup>, A. Boeltzig<sup>h</sup>, C. Broggini<sup>h</sup>, C.G. Bruno<sup>e</sup>, A. Cacioli<sup>h,i</sup>, F. Cavanna<sup>j</sup>, T. Chillery<sup>e</sup>, P. Colombetti<sup>k,l</sup>, P. Corvisiero<sup>l,m</sup>, T. Davinson<sup>e</sup>, R. Depalo<sup>h,i</sup>, A. Di Leva<sup>h,i</sup>, F. Elekes<sup>e</sup>, F. Ferraro<sup>l,m</sup>, E.M. Fiore<sup>l,n</sup>, A. Formicola<sup>o</sup>, Zs. Fülöp<sup>q</sup>, G. Gervino<sup>k,l</sup>, A. Guglielmetti<sup>o</sup>, C. Gustavino<sup>o</sup>, Gy. Gyürky<sup>q</sup>, G. Imbriani<sup>d</sup>, Z. Janas<sup>r</sup>, M. Junker<sup>s</sup>, I. Kochanek<sup>t</sup>, M. Lugaro<sup>u,v</sup>, P. Marigo<sup>h,i</sup>, E. Masha<sup>o</sup>, C. Mazzocchi<sup>h</sup>, R. Menegazzo<sup>h</sup>, V. Mossa<sup>f</sup>, F.R. Pantaleo<sup>h,s</sup>, V. Paticchio<sup>f</sup>, R. Perrino<sup>h</sup>, D. Piatti<sup>h,i</sup>, P. Prati<sup>l,m</sup>, L. Schiavulli<sup>h</sup>, K. Stöckel<sup>g,t</sup>, O. Straniero<sup>h,u</sup>

## Neutron detector based on $^3\text{He}$ counters



## Collaboration with Atomki & HZDR



ALMA MATER STUDIORUM  
UNIVERSITÀ DI BOLOGNA

# $^{13}\text{C}(\alpha, n)^{16}\text{O}$



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<https://doi.org/10.3847/1538-4357/aae177>

## The Importance of the $^{13}\text{C}(\alpha, n)^{16}\text{O}$ Reaction in Asymptotic Giant Branch Stars

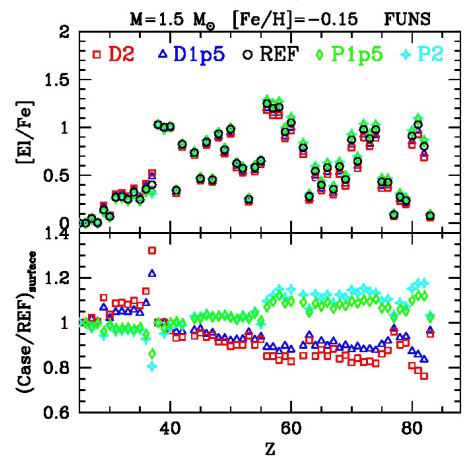
S. Cristallo<sup>1,2</sup>, M. La Cognata<sup>3</sup>, C. Massimi<sup>4,5</sup>, A. Best<sup>6,7</sup>, S. Palmerini<sup>2,8</sup>, O. Straniero<sup>1,9</sup>, O. Trippella<sup>2</sup>, M. Busso<sup>2,8</sup>, G. F. Ciani<sup>10</sup>, F. Mingrone<sup>11</sup>, L. Piersanti<sup>1,2</sup>, and D. Vescovi<sup>2,10</sup>

<sup>1</sup> INAF—Osservatorio Astronomico d’Abruzzo, via M. Maggini snc, Teramo, Italy; [sergio.cristallo@inaf.it](mailto:sergio.cristallo@inaf.it)  
<sup>2</sup> INFN—Sezione di Perugia, Via A. Pascoli snc, Perugia, Italy  
<sup>3</sup> INFN—Laboratori Nazionali del Sud, Via S. Sofia 62, Catania, Italy  
<sup>4</sup> Dipartimento di Fisica e Astronomia, Università di Bologna, Via Irnerio 46, Bologna, Italy  
<sup>5</sup> INFN—Sezione di Bologna, Viale Berti Pichat 6/2, Bologna, Italy  
<sup>6</sup> Università degli Studi di Napoli “Federico II,” Via Cintia, Napoli, Italy  
<sup>7</sup> INFN—Sezione di Napoli, Via Cintia, Napoli, Italy  
<sup>8</sup> Università degli Studi di Perugia, Via A. Pascoli snc, Perugia, Italy  
<sup>9</sup> INFN—Laboratori Nazionali del Gran Sasso, Via G. Acitelli 22, Assergi, Italy  
<sup>10</sup> Gran Sasso Science Institute, Viale Francesco Crispi 7, L’Aquila, Italy  
<sup>11</sup> CERN, Route de Meyrin 1211, Genève, France

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**Abstract**  
 Low-mass asymptotic giant branch stars are among the most important polluters of the interstellar medium. In their interiors, the main component ( $A \gtrsim 90$ ) of the slow neutron capture process (the s-process) is synthesized, the most important neutron source being the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction. In this paper, we review its current experimental status, discussing possible future synergies between some experiments. Moreover, in order to determine the level of precision needed to sensitivity study, carried out with the FUNS evolutionary stella the rate up to a factor of 2 with respect to a reference case, appreciably affect s-process distributions for masses above  $3 M_{\odot}$  the differences are always below 5%. The situation is not environment: this occurs in FUNS models with  $M < 3 M_{\odot}$ .  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction rate leads to nonnegligible variations of larger peaks for some elements (such as rubidium) and neutron are found in low-mass, low-metallicity models if protons are in the surface abundances of the heavier elements may vary by

**Key words:** nuclear reactions, nucleosynthesis, abundances



There is some interest in measuring the inverse reaction at n\_TOF (a first attempt with gas detector failed, and a new setup is being developed).

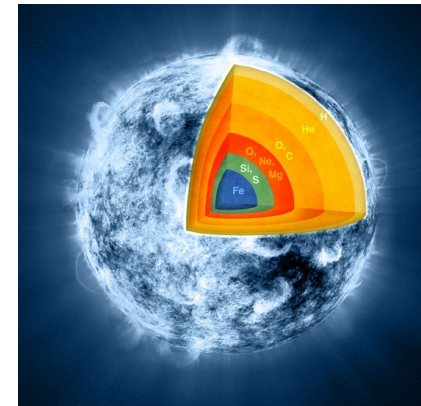
By using the detailed balance (i.e., time-reversal invariance theorem) the reaction cross section of  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  is deduced from the measurement in the reverse direction  $^{16}\text{O}(n, \alpha)^{13}\text{C}$ .

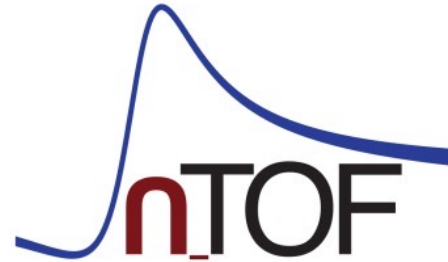
A single measurement could provide data in a large energy region.



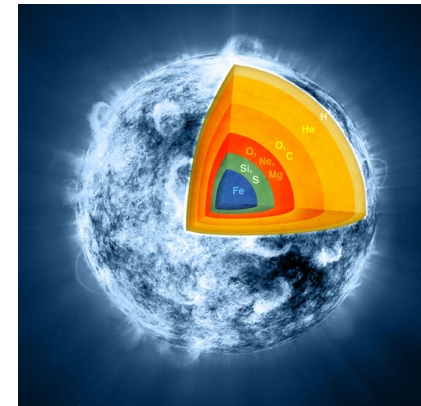


- ❖ Main neutron source in massive stars
- ❖ Gamow window: threshold  $\sim 300$  keV

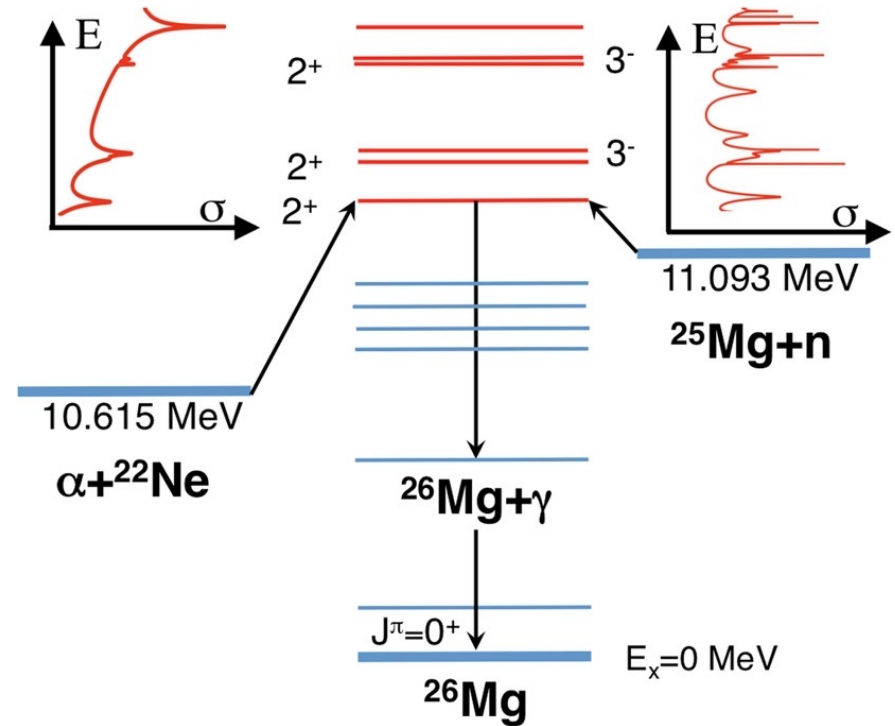
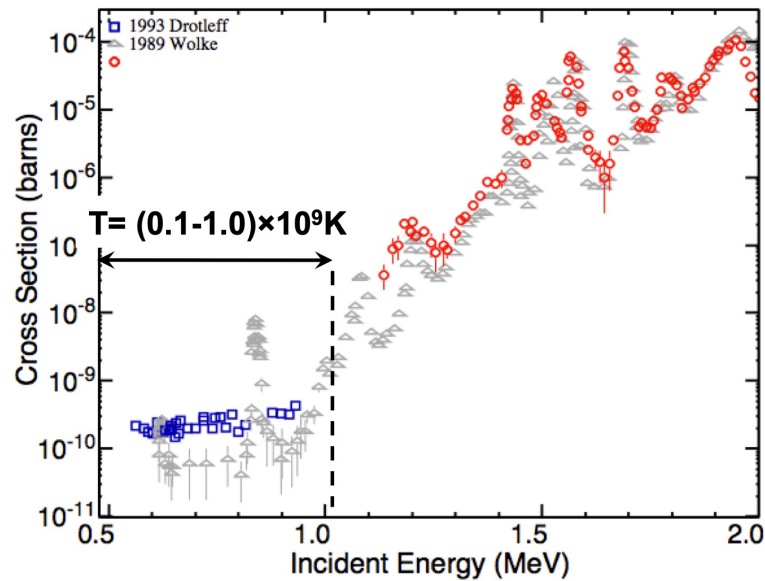




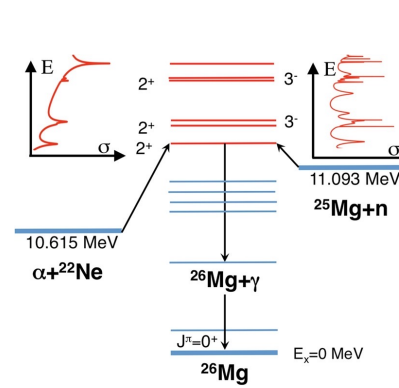
- ❖ Main neutron source in massive star
- ❖ Gamow window: threshold  $\sim 300$  keV



# $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$



- ❖ Possible presence of resonances in the Gamow window (unmeasured)
- ❖ Large uncertainty in extrapolations at low energy
- ❖ Background on surface labs. too high (only upper limit in Gamow window)
- ❖ Non negligible  $^{22}\text{Ne}(\alpha, \gamma)$  branch





# $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$

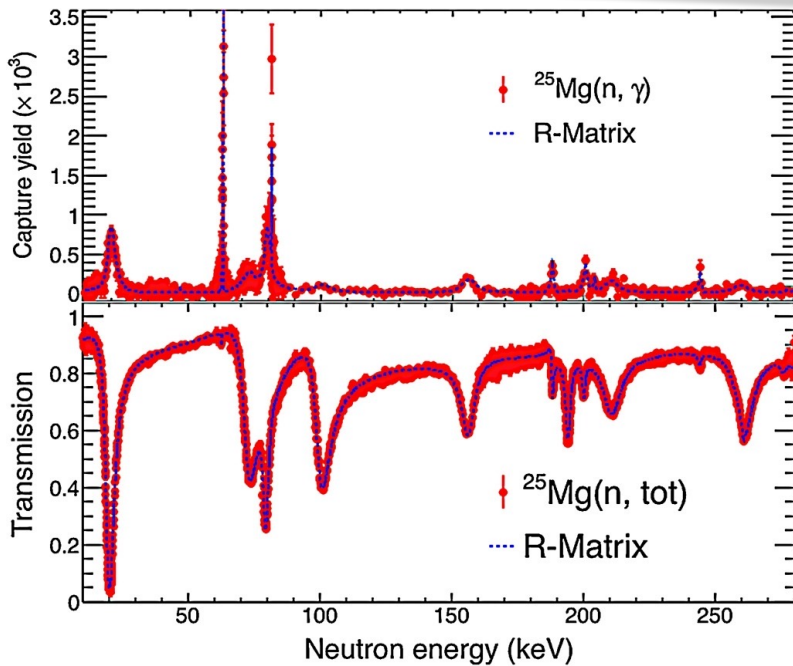
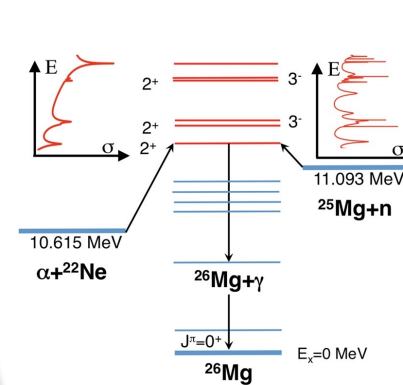
Physics Letters B 768 (2017) 1–6  
 Contents lists available at ScienceDirect  
 Physics Letters B  
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Neutron spectroscopy of  $^{26}\text{Mg}$  states: Constraining the stellar neutron source  $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$

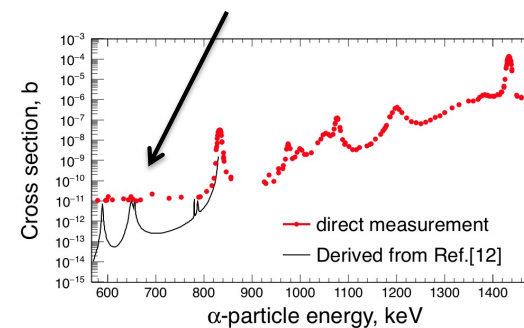
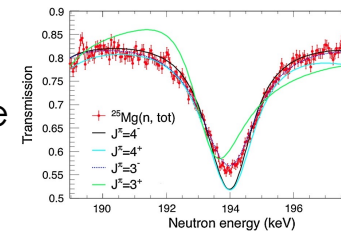
C. Massimi<sup>a,b,\*</sup>, S. Altstadt<sup>c</sup>, I. Andriejowski<sup>d</sup>, I. Audouin<sup>e</sup>, M. Barbagallo<sup>f</sup>, V. Récaries<sup>g</sup>

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- ❖ Indirect approach, resonances above  $n$ -threshold
- ❖ R-Matrix parametrization ( $E_R, \Gamma_\gamma, \Gamma_n, J^\pi$ )
- ❖ Deduced  $^{26}\text{Mg}$  states with natural parity
- ❖ No experimental  $\Gamma_\alpha$



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# $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$

# international effort ChETEC

- $n+^{25}\text{Mg}$
- $^{26}\text{Mg}(\alpha, \alpha')^{26}\text{Mg}$ ,  $^{26}\text{Mg}(p, p')^{26}\text{Mg}$ ,  $^{26}\text{Mg}(d, d')^{26}\text{Mg}$
- $^{22}\text{Ne}(^6\text{Li}, d)^{26}\text{Mg}$
- $^{11}\text{B}(^{16}\text{O}, p)^{26}\text{Mg}$

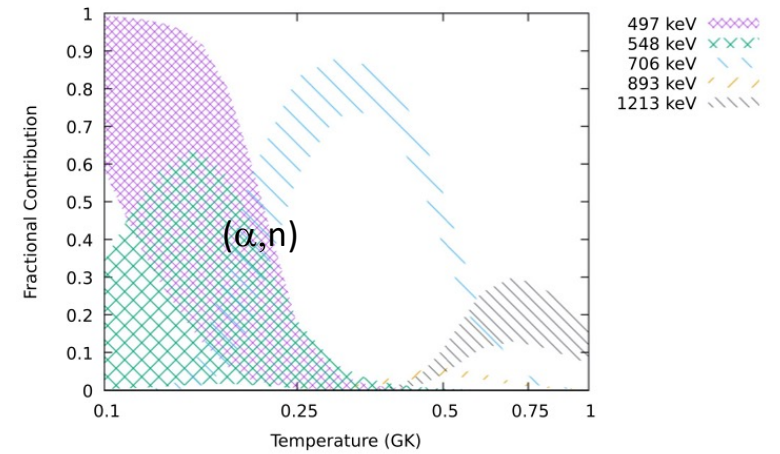
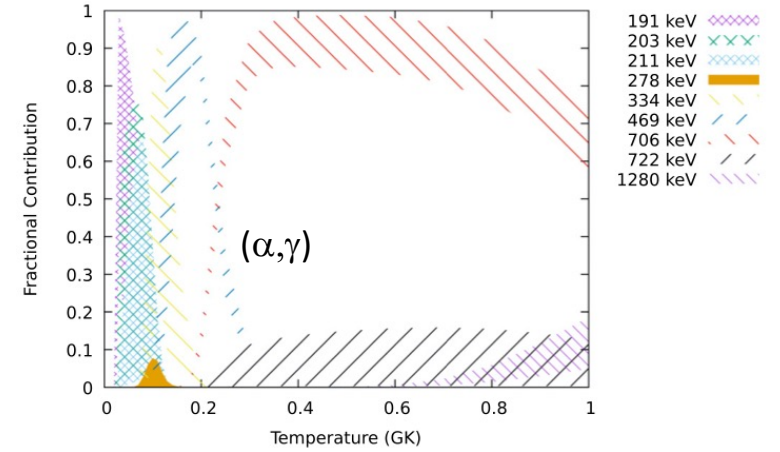
PHYSICAL REVIEW C **103**, 015805 (2021)

## Reevaluation of the $^{22}\text{Ne}(\alpha, \gamma)^{26}\text{Mg}$ and $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction rates

Philip Adsley<sup>1,2,3,\*</sup>, Umberto Battino<sup>4,†</sup>, Andreas Best<sup>5,6</sup>, Antonio Cacioli<sup>7,8</sup>, Alessandra Guglielmetti<sup>9</sup>,  
 Gianluca Imbriani<sup>5,6</sup>, Heshani Jayatissa<sup>10</sup>, Marco La Cognata<sup>11</sup>, Livio Lamia<sup>12,11,13</sup>, Eliana Masha<sup>9</sup>,  
 Cristian Massimi<sup>14,15</sup>, Sara Palmerini<sup>16,17</sup>, Ashley Tattersall<sup>4,†</sup> and Raphael Hirschi<sup>18,19,†</sup>

$E_x$ (MeV)	$E_x^{\text{CM}}$ (keV)	$J^\pi$	$\omega\gamma_{(\alpha,\gamma)}$ (eV)	$\omega\gamma_{(\alpha,n)}$ (eV)	$\Gamma_\alpha$ (eV)	$\Gamma_\gamma$ (eV)	$\Gamma_n$ (eV)	Integrate resonance?
10.6963(4)	81.6(4)	4 <sup>+</sup>			$3.5(18) \times 10^{-46}$	3.0(15)	0	No
11.084(1)	469(1)	2 <sup>+</sup>			$5.7(1.5) \times 10^{-11}$	3.0(15)	0	No
11.321(1)	706(1) <sup>a</sup>	0 <sup>+</sup> /1 <sup>-</sup>	$3.7(4) \times 10^{-5}$	$4.2(11) \times 10^{-5}$				No
11.44120(4)	826.46(5)	3 <sup>-</sup>		$3.9(10) \times 10^{-5}$	$5.50(14) \times 10^{-6}$	3.0(15)	$1.47(8) \times 10^3$	Yes
11.46574(6)	851.00(6)	3 <sup>-</sup>		$5.5(17) \times 10^{-5}$	$7.9(2.4) \times 10^{-6}$	3.0(15)	$6.55(9) \times 10^3$	Yes
11.5080(9)	893.3(9)	1 <sup>-</sup>		$3.5(6) \times 10^{-4}$	$1.2(4) \times 10^{-4}$	3.0(15)	$1.27(25) \times 10^3$	Yes
11.5260(15)	911.3(15)	1 <sup>-</sup>		$1.3(4) \times 10^{-3}$	$4.3(11) \times 10^{-4}$	3.0(15)	$1.80(25) \times 10^3$	Yes
11.630(1)	1015.3(14)	1 <sup>-</sup>		$7.1(15) \times 10^{-3}$	$2.4(5) \times 10^{-3}$	3.0(15)	$13.5(17) \times 10^3$	Yes
11.749(5)	1133(6)	1 <sup>-</sup>		$5.9(8) \times 10^{-2}$	$2.0(3) \times 10^{-2}$	3.0(15)	$64(9) \times 10^3$	Yes
11.787(3)	1172(3)	1 <sup>-</sup>		$2.5(9) \times 10^{-2}$	$8(3) \times 10^{-3}$	3.0(15)	$24.5(24) \times 10^3$	Yes
11.828(1)	1213(1)	2 <sup>+</sup>		$2.5(3) \times 10^{-4}$	$1.8(1) \times 10^{-1}$	3.0(15)	$1.10(25) \times 10^3$	Yes
11.863(3)	1248(3)	1 <sup>-</sup>			$1.5(10) \times 10^{-2}$	3.0(15)	$2.45(34) \times 10^4$	Yes
11.880(3)	1265(3)	1 <sup>-</sup>		$1.9(19) \times 10^{-1}$	$6.30(63) \times 10^{-2}$	3.0(15)	$3.0(15) \times 10^3$	No
11.895(4)	1280(4)	1 <sup>-</sup>	$2.0(2) \times 10^{-3}$	$4.1(4) \times 10^{-1}$				No
11.911(1)	1297(3)	1 <sup>-</sup>	$3.4(4) \times 10^{-3}$	1.4(1)	1.9(9.8)	3.0(15)	$5(2) \times 10^3$	Yes
11.953(3)	1338(3)	2 <sup>+</sup>	$3.4(4) \times 10^{-3}$	1.60(13)	$3.2(1.7) \times 10^{-1}$	3.0(15)	$2(1) \times 10^3$	Yes
12.050(1)	1436(3)	2 <sup>+</sup>	$6.0(7) \times 10^{-3}$	4.7(3)	$1.1(3) \times 10^{-1}$	3.0(15)	$4(1) \times 10^3$	Yes
12.141(1)	1526(3)	1 <sup>-</sup>	$1.0(2) \times 10^{-3}$	2.4(2)	1.7(5)	3.0(15)	$1.5(2) \times 10^4$	Yes
12.184(5)	1569(7)	0 <sup>+</sup>	$1.1(2) \times 10^{-3}$	$1.21(29) \times 10^1$	0.90(11)	3.0(15)	$3.3(5) \times 10^4$	Yes
12.270(5)	1658(7)	0 <sup>+</sup>	$8.9(1) \times 10^{-3}$	$2.1(2) \times 10^1$	$2.2(4) \times 10^2$	3.0(15)	$7.3(9) \times 10^4$	Yes
12.344(2)	1728(4)	0 <sup>+</sup>	$5.4(7) \times 10^{-2}$	$1.57(10) \times 10^2$	$6.30(12) \times 10^2$	3.0(15)	$3.5(5) \times 10^4$	Yes

<sup>a</sup>Resonance energy is taken using the state observed in Ref. [17] and assuming a single resonance.



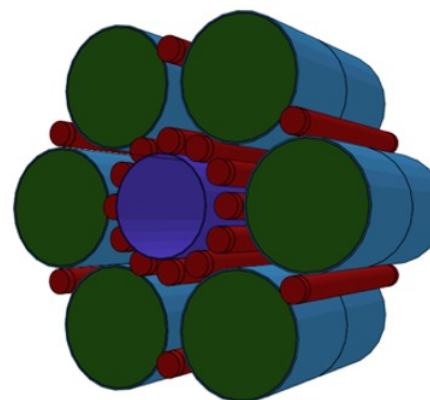
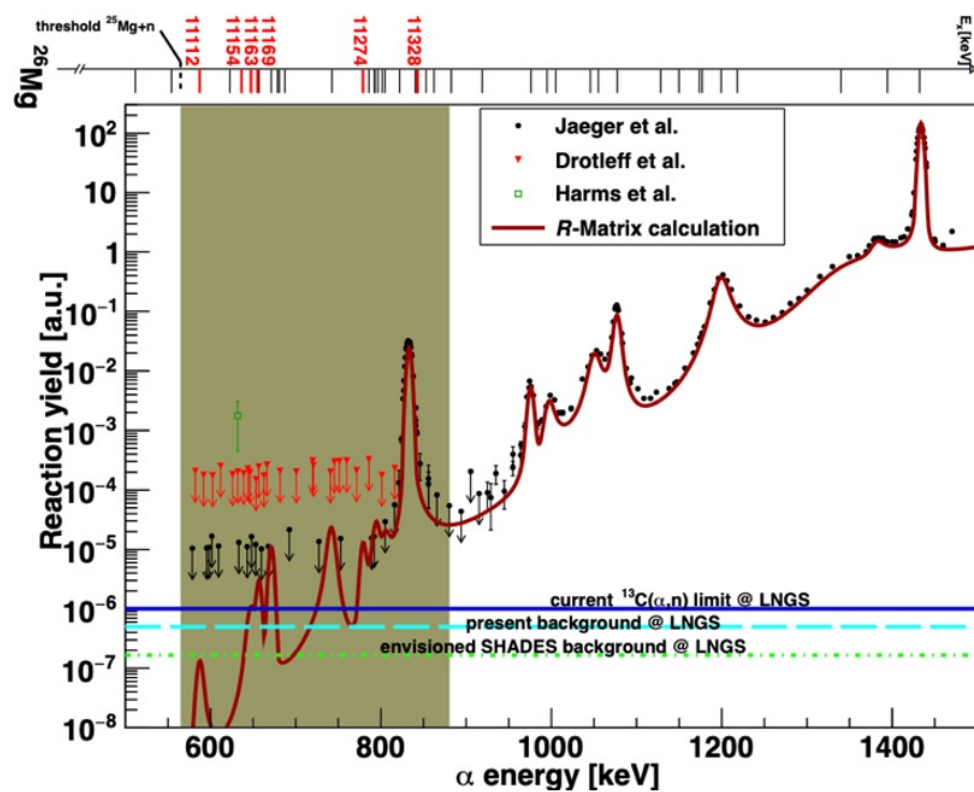
Only natural-parity states in  $^{26}\text{Mg}$  contribute



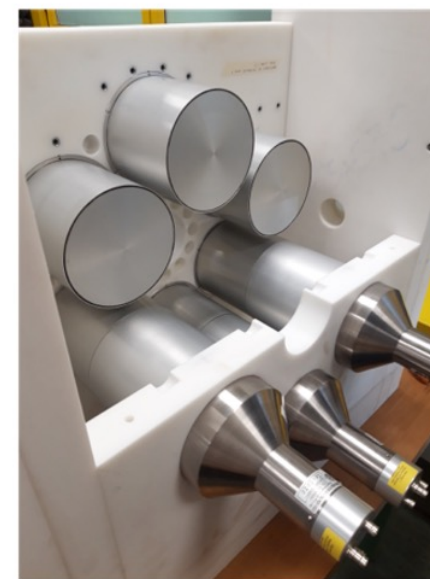
INFN2022, 9–11 May 2022, LNGS

## $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$

- ❖ LUNA-MV scientific proposal: **started in 2020, first measurements in 2022**
- ❖ Novel high-efficiency, energy sensitive detector array ( $^3\text{He}$  + liquid scintillator)

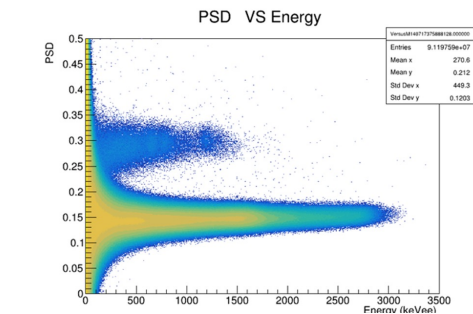
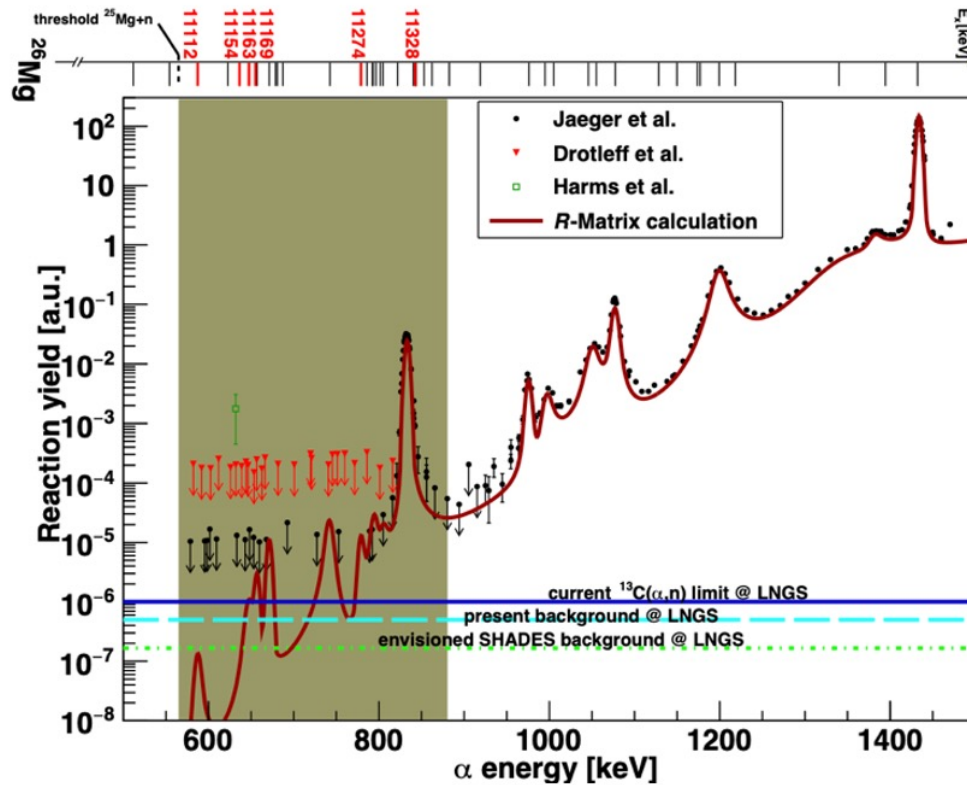


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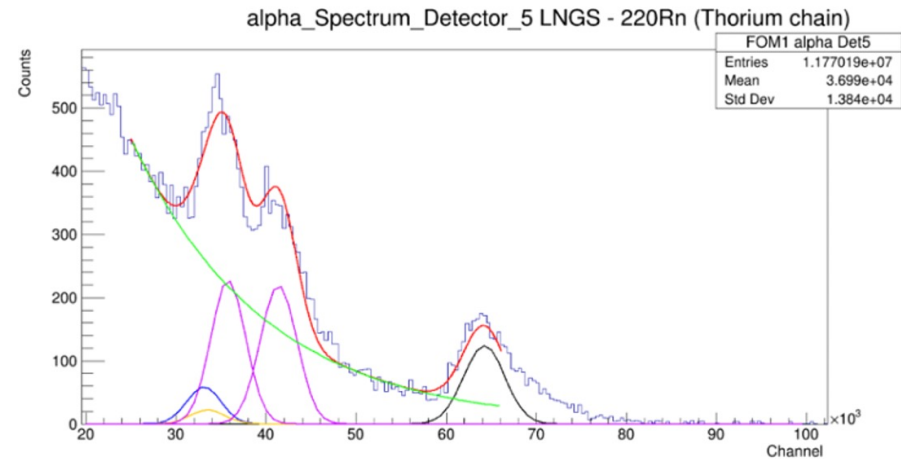


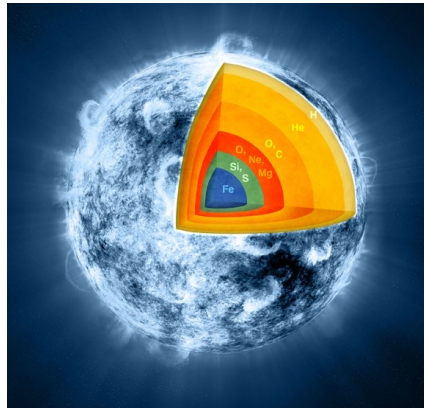
# $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$

- ❖ LUNA-MV scientific proposal: **started in 2020, first measurements in 2022**
- ❖ Novel high-efficiency, energy sensitive detector array ( $^3\text{He}$  + liquid scintillator)



Characterization of detectors and gas target ongoing

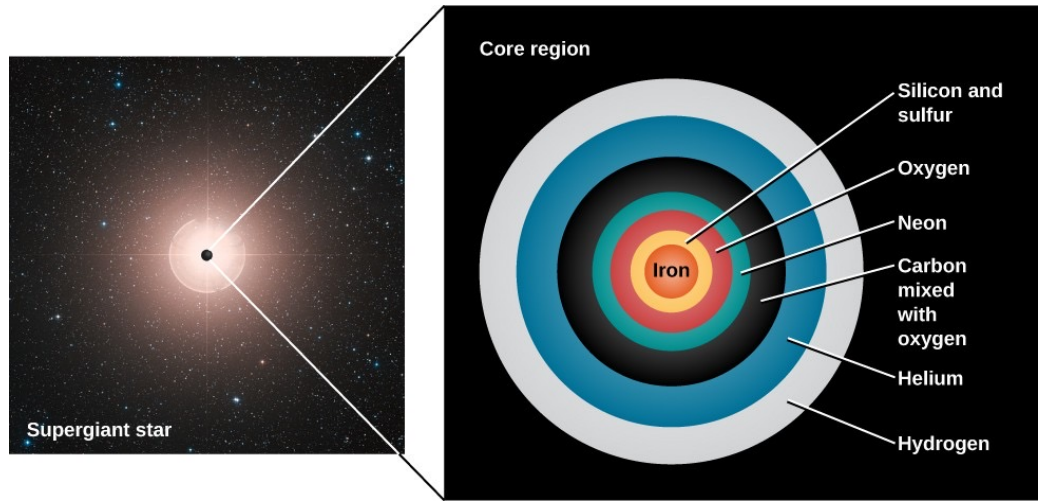




ASFIN aims at measuring the inverse reaction  $^{25}\text{Mg}(n, \alpha)^{22}\text{Ne}$  with an indirect method based on Trojan Horse Method (feasibility studies ongoing)



# Explosive burning (massive stars)



Quiescent  $^4\text{He}$  and heavy ion burnings result in massive stars with convective shells at different densities: onion structure

Quiescent and explosive nucleosynthesis in massive stars are constrained by a few nuclear reactions:

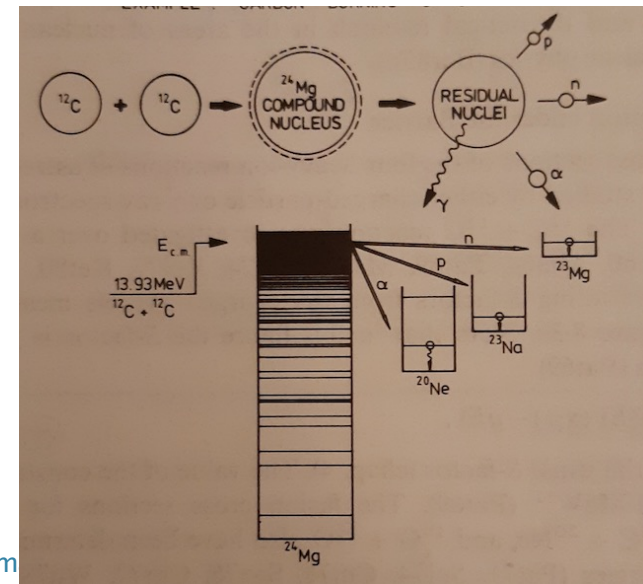
- $^{12}\text{C}(\alpha, n)$
- $^{12}\text{C}(^{12}\text{C}, p)^{23}\text{Na}$
- $^{12}\text{C}(^{12}\text{C}, \alpha)^{20}\text{Ne}$
- $^{12}\text{C} + ^{16}\text{O}$
- ... and others



Type Ia Supernovae

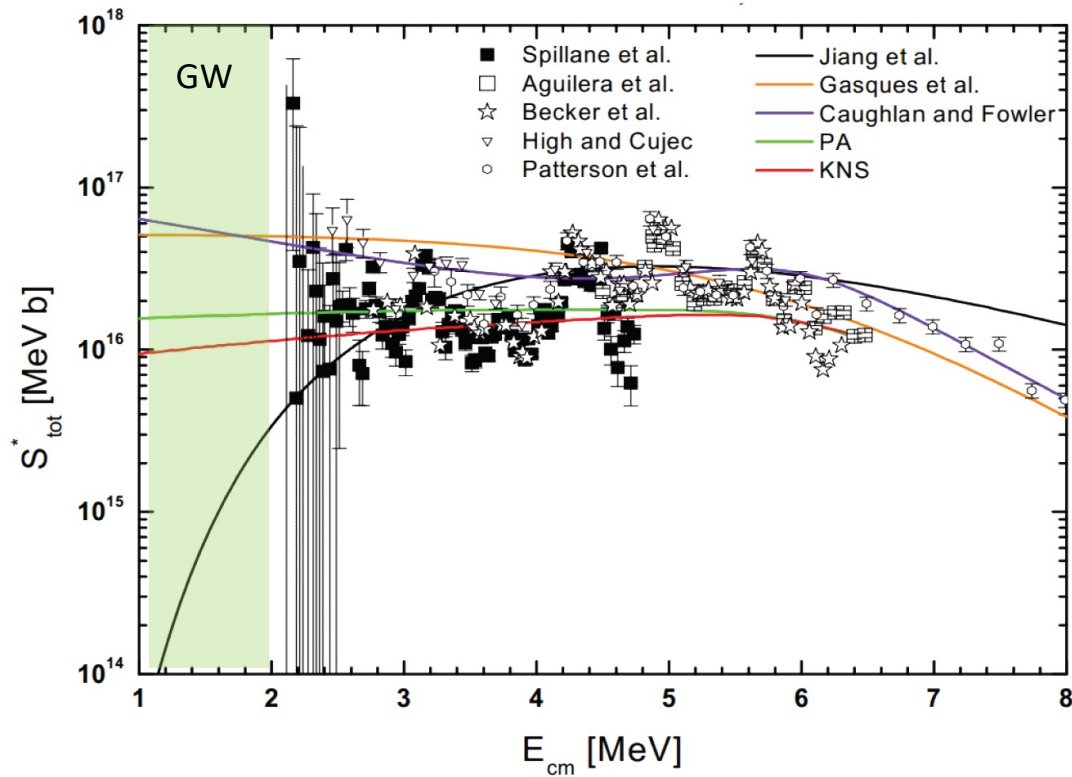


- ❖ Reaction rate determines the late phase of stellar evolution
- ❖ Gamow window below Coulomb barrier  $\sim 1\text{--}2\text{ MeV}$





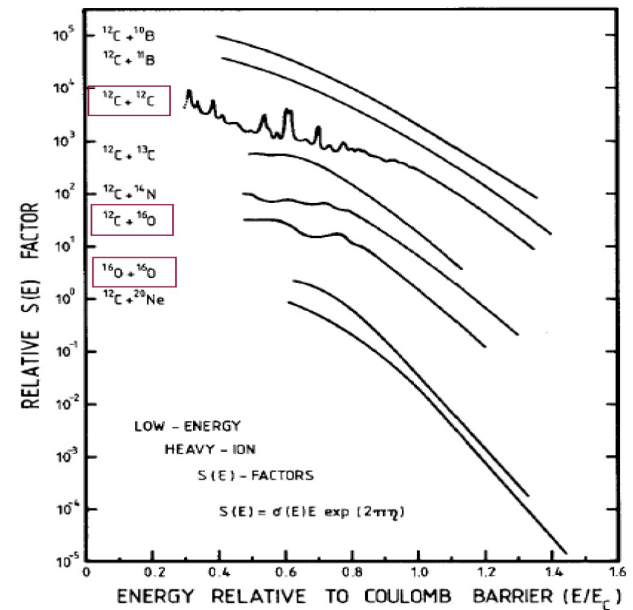
# $^{12}\text{C}+^{12}\text{C}$ "BEFORE"



Measurements based on  $\gamma$ -ray or particle spectroscopy



- ❖ Presence of resonances in the Gamow window (unmeasured)
- ❖ Large uncertainty in extrapolations at low energy
- ❖ Microscopic model?
- ❖ Impact of the background



R. Stokstad et al., Phys.Rev.Lett. 37 (1976)

$^{12}\text{C}+^{12}\text{C}$



# $^{12}\text{C}+^{12}\text{C}$

Eur. Phys. J. A (2018) 54: 132  
DOI 10.1140/epja/i2018-12564-8

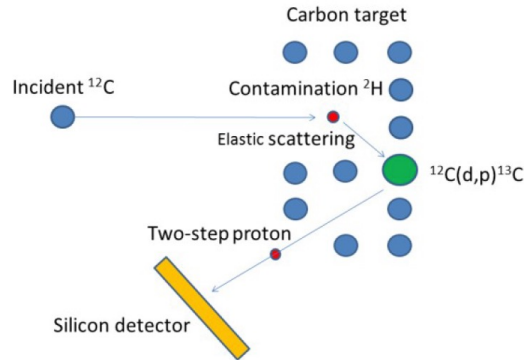
THE EUROPEAN PHYSICAL JOURNAL A

Regular Article – Experimental Physics

**Reduction of deuterium content in carbon targets for  $^{12}\text{C} + ^{12}\text{C}$  reaction studies of astrophysical interest**

L. Morales-Gallegos<sup>1,2</sup>, M. Aliotta<sup>1,3,a</sup>, C.G. Bruno<sup>1</sup>, R. Buompane<sup>3,2</sup>, T. Davinson<sup>1</sup>, M. De Cesare<sup>1,2</sup>, A. Di Leva<sup>5,2</sup>, A. D'Onofrio<sup>3,2</sup>, J.G. Duarte<sup>3,2</sup>, L.R. Gasques<sup>6,3,2</sup>, L. Gialanella<sup>3,2,b</sup>, G. Imbriani<sup>2,2</sup>, G. Porzio<sup>2</sup>, D. Rapanotti<sup>7,8</sup>, M. Romoli<sup>2</sup>, D. Schürmann<sup>3,2</sup>, F. Terrasi<sup>3,2</sup>, and L.Y. Zhang<sup>8,1</sup>

## Beam-induced background background



### Low energy beam induced background studies for a $^{12}\text{C}(^{12}\text{C},p)^{23}\text{Na}$ reaction cross section measurement

James Zickefoose, Jeffrey Schweitzer, Timothy Spillane  
University of Connecticut  
Storrs CT, USA  
zickefoose@phys.uconn.edu

Frank Strieder<sup>1</sup>, Hans-Werner Becker, Claus Rolfs  
Ruhr-Universität Bochum  
Bochum, Germany  
strieder@ep3.rub.de

Antonino Di Leva, Mario De Cesare, Nicola De Cesare, Filippo Terrasi  
Seconda Università di Napoli and INFN Napoli  
Caserta, Italy

Lucio Gialanella, Daniel Schürmann  
INFN Napoli  
Naples, Italy

Yongjing Guan  
Guangxi University  
Nanning, China

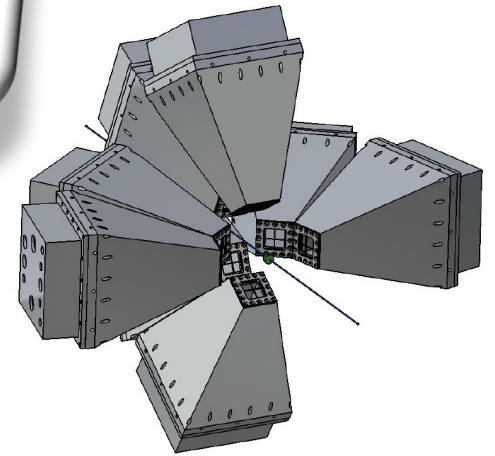
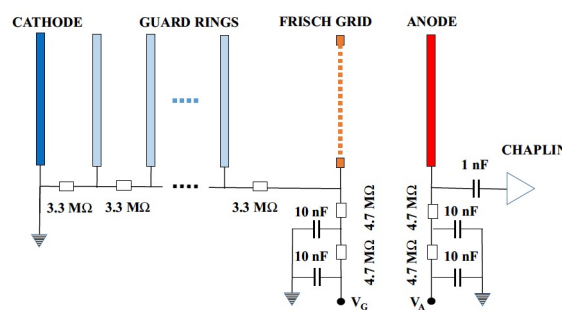
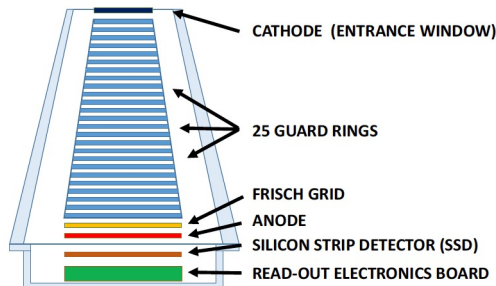
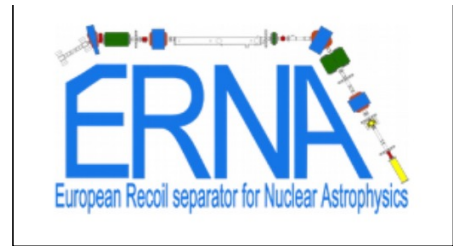
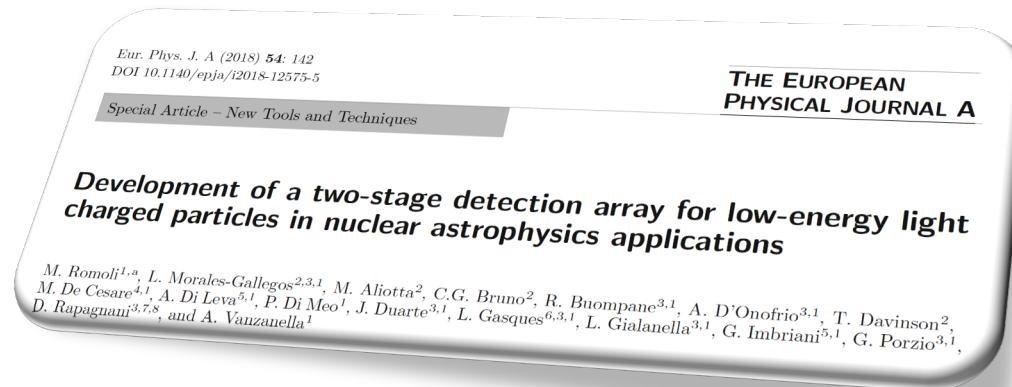
Gianluca Imbriani, Benedicita Limata  
Università di Napoli Federico II and INFN Napoli  
Naples, Italy

Using thick  $^{12}\text{C}$  target methods (Taking the difference of two measurements at different energies) experimental evidences are protons and  $\alpha$  particle from the reactions  $^{12}\text{C}(^{12}\text{C},p)^{23}\text{Na}$  &  $^{12}\text{C}(^{12}\text{C},\alpha)^{20}\text{Ne}$



- ❖ Detailed study of the background
- ❖ Study of suitable  $^{12}\text{C}$  carbon targets → Morales-Gallegos et al 2015 J. Phys.: Conf. Ser. 578 012002

$^{12}\text{C}+^{12}\text{C}$



Gastly -> particle spectroscopy



- ❖ Detailed study of the background
- ❖ Study of suitable  $^{12}\text{C}$  carbon targets → Morales-Gallegos et al 2015 J. Phys.: Conf. Ser. 578 012002
- ❖ Development of detection setup for p and  $\alpha$

$^{12}\text{C}+^{12}\text{C}$

@ CIRCE  
 $2 < E_{\text{CM}} < 4 \text{ MeV}$



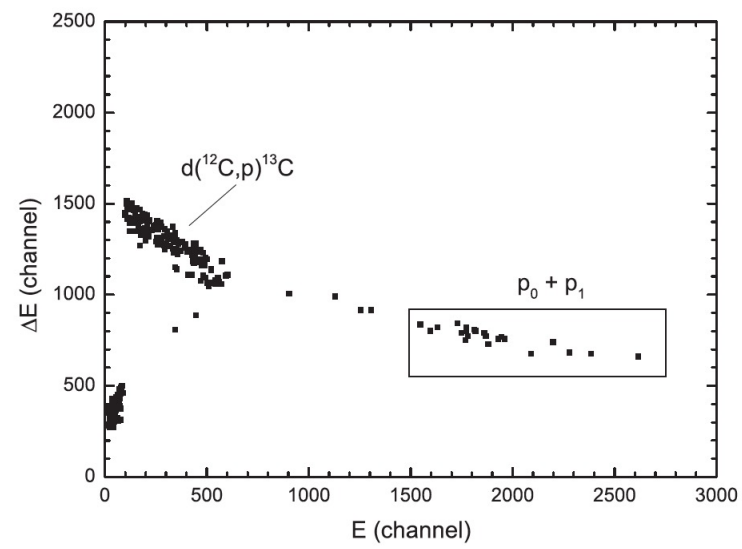
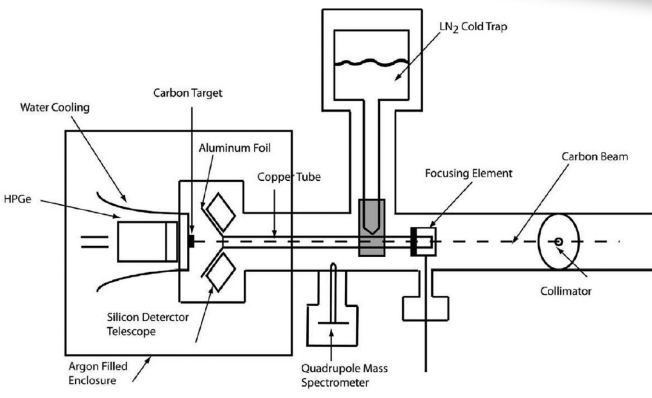
PHYSICAL REVIEW C 97, 065806 (2018)

**Measurement of the  $^{12}\text{C}(^{12}\text{C}, p)^{23}\text{Na}$  cross section near the Gamow energy**

J. Zickefoose,<sup>1,\*</sup> A. Di Leva,<sup>2,3</sup> F. Strieder,<sup>4</sup> L. Gialanella,<sup>5,3</sup> G. Imbriani,<sup>2,3</sup> N. De Cesare,<sup>5,3,†</sup> C. Rolfs,<sup>6</sup> J. Schweitzer,<sup>1</sup> T. Spillane,<sup>1</sup> O. Straniero,<sup>7,8</sup> and F. Terrasi<sup>5,3</sup>



Transition	$J^\pi$	$E_x$ (MeV)	$Q_x$ (MeV)	$E_p$ (MeV)
$p_0$	$3/2^+$	0	2.240	3.71
$p_1$	$5/2^+$	0.440	1.801	3.33
$p_2$	$7/2^+$	2.076	0.165	1.95
$p_3$	$1/2^+$	2.391	-0.150	1.69
$p_4$	$1/2^-$	2.640	-0.399	1.48
$p_5$	$9/2^+$	2.704	-0.463	1.43
$p_6$	$3/2^+$	2.982	-0.741	1.20
$n_7$	$3/2^-$	3.678	-1.437	0.67
	$5/2^-$	3.848	-1.607	0.53
	$5/2^+$	3.914	-1.673	0.48
	$1/2^+$	4.430	-2.189	0.13



# $^{12}\text{C}+^{12}\text{C}$

Eur. Phys. J. A (2022) 58:65  
<https://doi.org/10.1140/epja/s10050-022-00717-7>

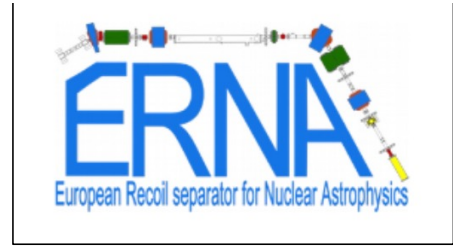
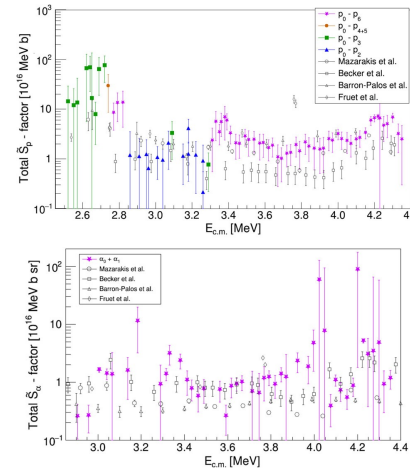
Regular Article - Experimental Physics

THE EUROPEAN PHYSICAL JOURNAL A

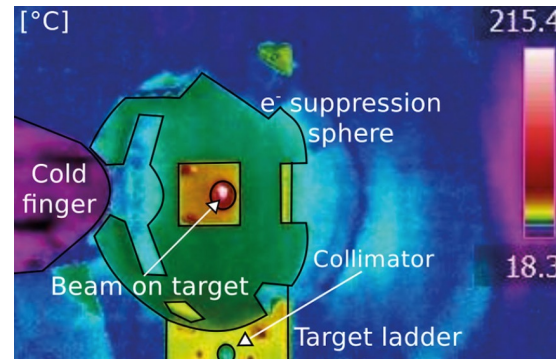
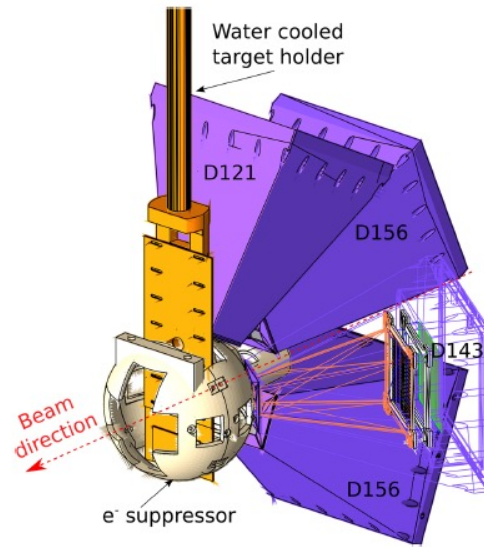
Check for updates

**Direct measurements of the  $^{12}\text{C}+^{12}\text{C}$  reactions cross-sections towards astrophysical energies**

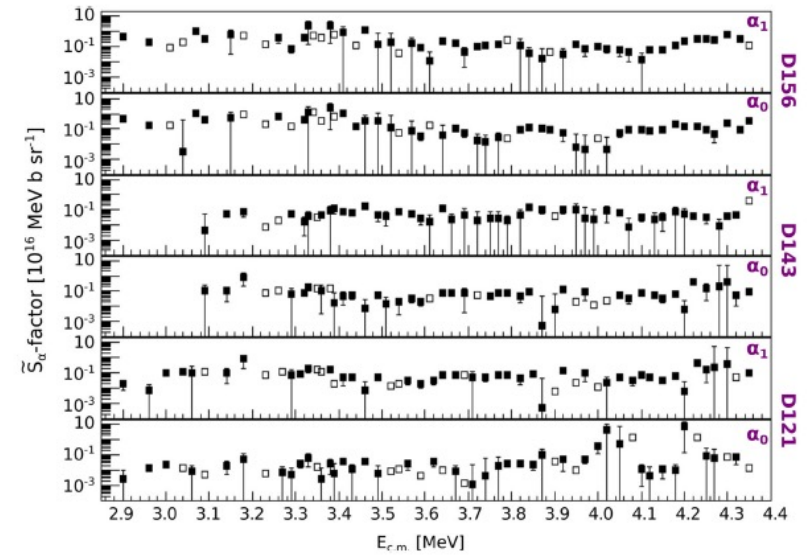
L. Morales-Gallegos<sup>1,2,3,a</sup>, M. Aliotta<sup>1,2</sup>, L. Gialanella<sup>1,3</sup>, A. Best<sup>3,5</sup>, C. G. Bruno<sup>2</sup>, R. Buompane<sup>1,3</sup>, T. Davinson<sup>2</sup>, M. De Cesare<sup>3,4</sup>, A. Di Leva<sup>3,5</sup>, A. D'Onofrio<sup>1,3</sup>, J. G. Duarte<sup>1,3,6</sup>, L. R. Gasques<sup>1,3,7</sup>, G. Imbriani<sup>3,5</sup>, G. Porzio<sup>1</sup>, D. Rapagnani<sup>1,3</sup>, M. Romoli<sup>3</sup>, F. Terrasi<sup>1,3</sup>



To be completed ...  
 data analysis ongoing

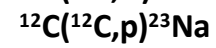
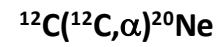


Temperature monitor

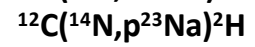
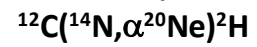




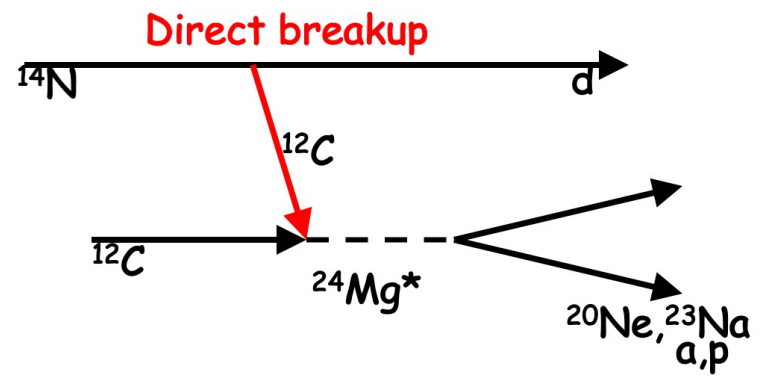
## $^{12}\text{C}+^{12}\text{C}$



via the Trojan Horse Method applied to

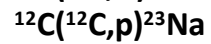
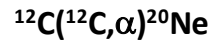


three-body processes

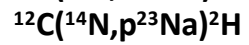
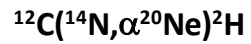




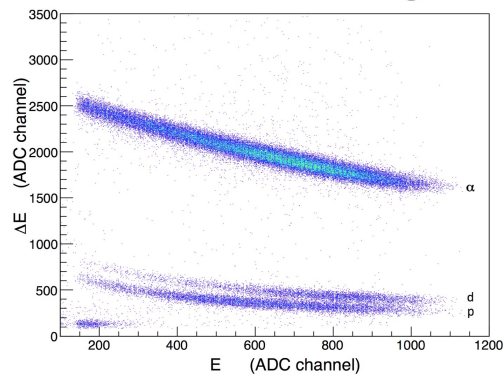
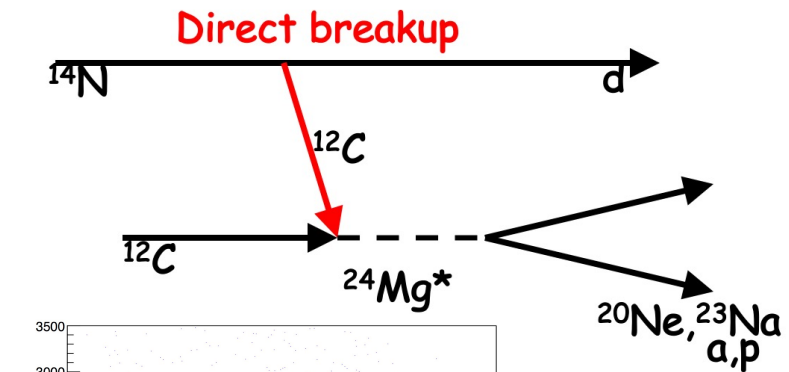
## $^{12}\text{C}+^{12}\text{C}$



via the Trojan Horse Method applied to



three-body processes



- ❖ Normalization to direct data
- ❖ Model to extract direct cross section







# $^{12}\text{C}+^{12}\text{C}$

$^{12}\text{C}(^{12}\text{C},\alpha)^{20}\text{Ne}$

$^{12}\text{C}(^{12}\text{C},\text{p})^{23}\text{Na}$

via the Trojan Horse Method applied to

$^{12}\text{C}(^{14}\text{N},\alpha)^{20}\text{Ne}+^2\text{H}$

$^{12}\text{C}(^{14}\text{N},\text{p})^{23}\text{Na}+^2\text{H}$

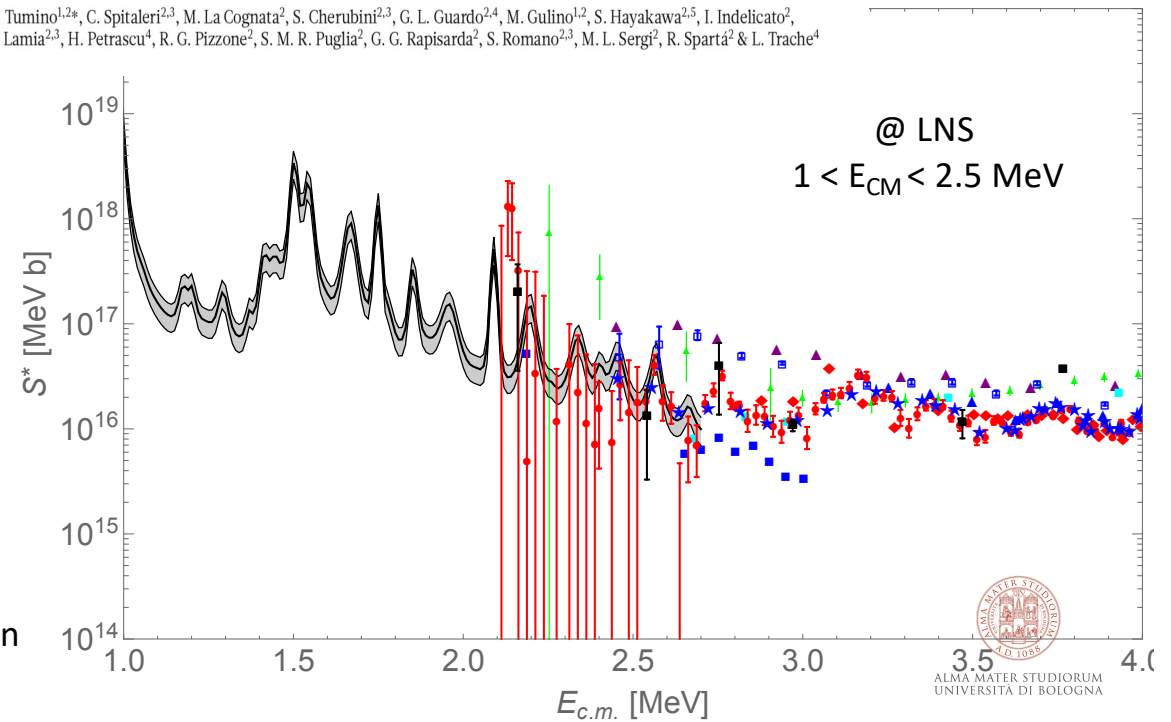
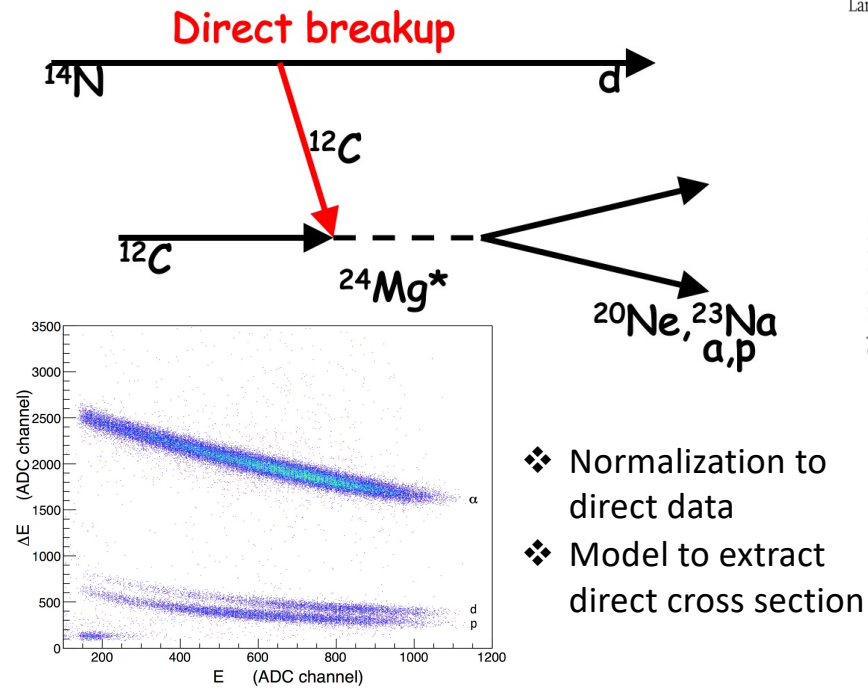
three-body processes

## LETTER

<https://doi.org/10.1038/s41586-018-0149-4>

### An increase in the $^{12}\text{C} + ^{12}\text{C}$ fusion rate from resonances at astrophysical energies

<sup>A</sup> Tumino<sup>1,2\*</sup>, C. Spitaleri<sup>2,3</sup>, M. La Cognata<sup>2</sup>, S. Cherubini<sup>2,3</sup>, G. L. Guardo<sup>2,4</sup>, M. Gulino<sup>1,2</sup>, S. Hayakawa<sup>2,5</sup>, I. Indelicato<sup>2</sup>, Lamia<sup>2,3</sup>, H. Petruscu<sup>4</sup>, R. G. Pizzone<sup>2</sup>, S. M. R. Puglia<sup>2</sup>, G. G. Rapisarda<sup>2</sup>, S. Romano<sup>2,3</sup>, M. L. Sergi<sup>2</sup>, R. Spart<sup>4</sup> & L. Trache<sup>4</sup>



$^{12}\text{C}+^{12}\text{C}$

international effort { STELLA @ IPN ORSAY  
NOTRE DAME Univ.

INFN2022, 9–11 May 2022, LNGS

[massimi@bo.infn.it](mailto:massimi@bo.infn.it)

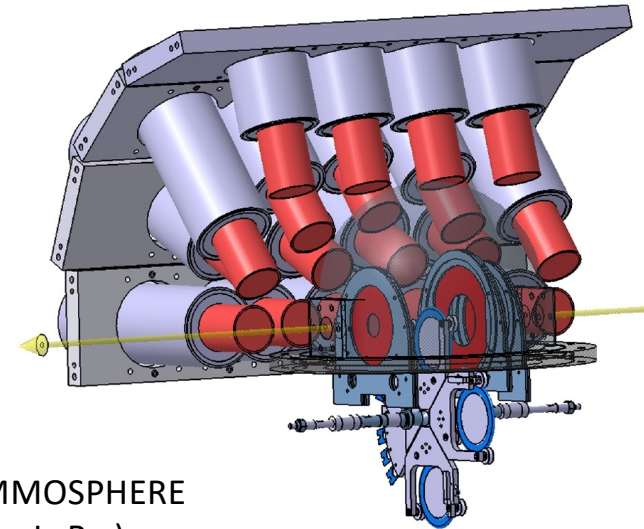
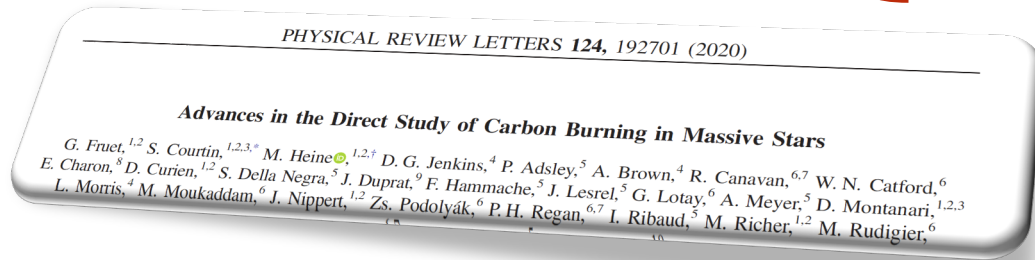


ALMA MATER STUDIORUM  
UNIVERSITÀ DI BOLOGNA

$^{12}\text{C}+^{12}\text{C}$

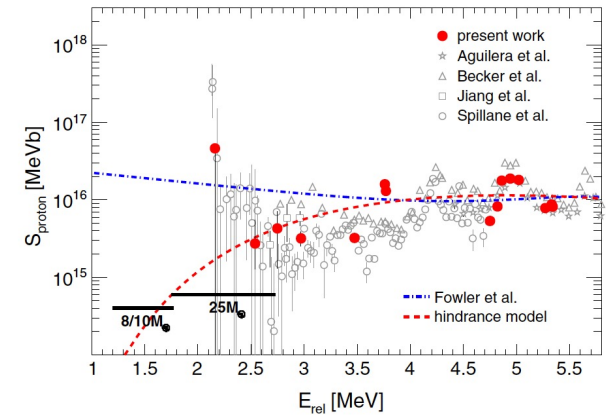
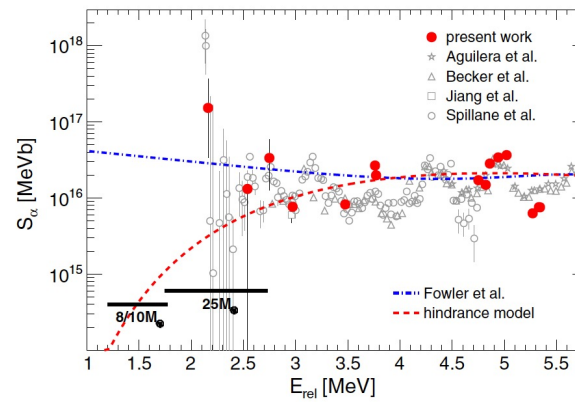
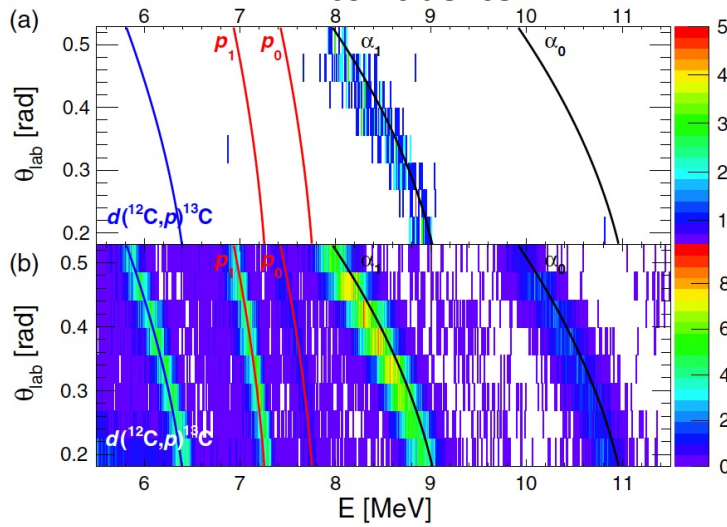
international effort

STELLA @ IPN ORSAY  
NOTRE DAME Univ.



Measurements based on  
 $\gamma$ -ray + particle spectroscopy  
In coincidence

STELLA + GAMMOSPHERE  
(Anular Si + LaBr<sub>3</sub>)

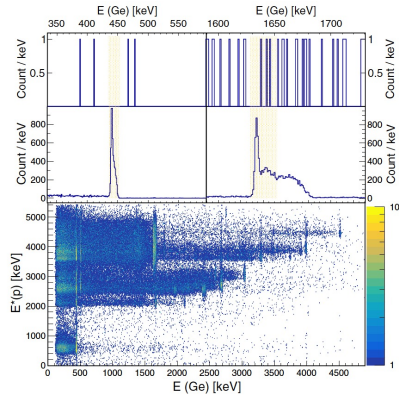


$^{12}\text{C}+^{12}\text{C}$

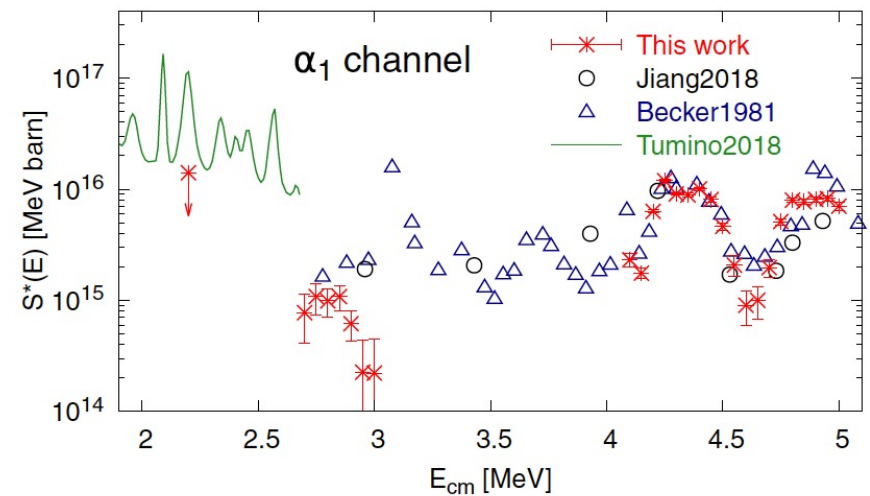
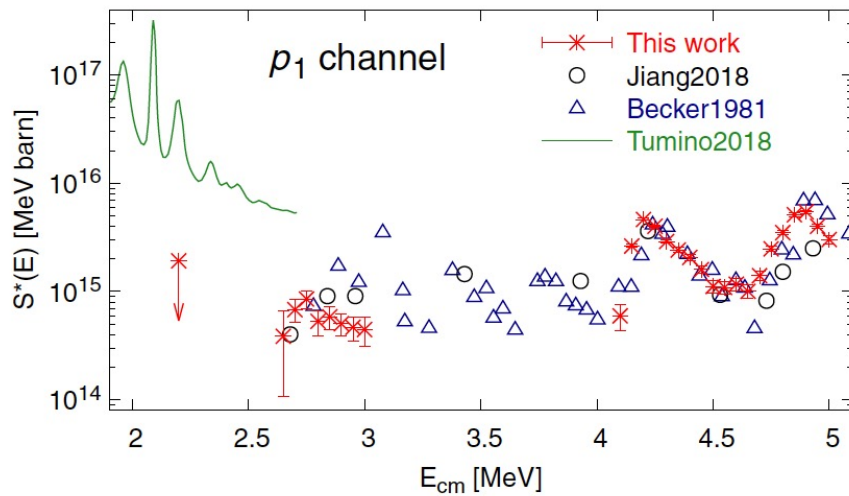
international effort

STELLA @ IPN ORSAY  
NOTRE DAME Univ.

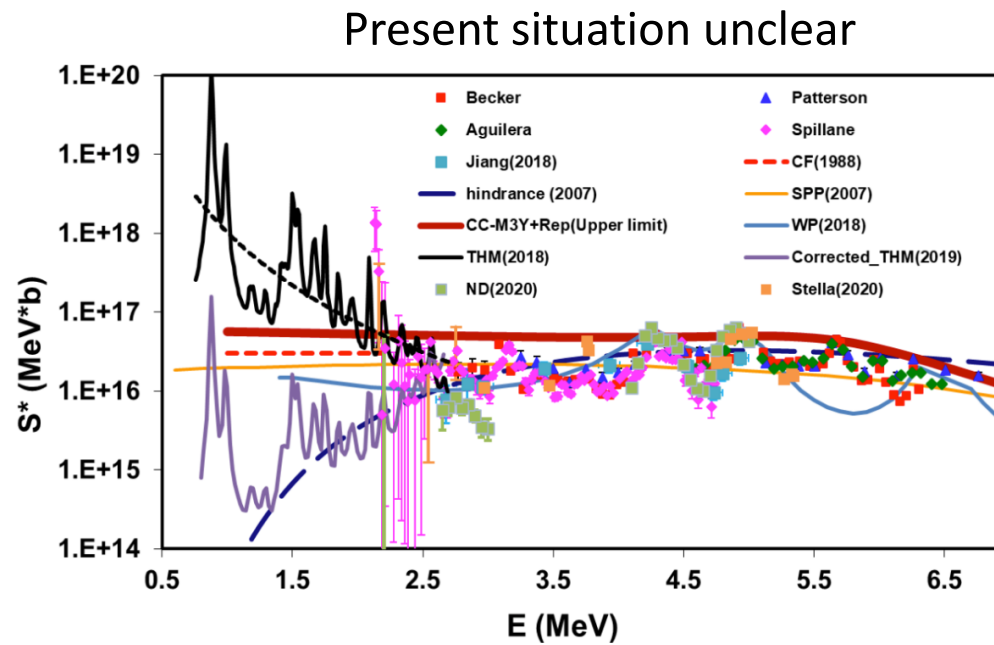
Silicon detector array + HPGe



Measurements based on  
 $\gamma$ -ray + particle spectroscopy  
in coincidence



$^{12}\text{C}+^{12}\text{C}$



Measurement to be included in the LUNA-MV scientific program

# $^{12}\text{C}+^{12}\text{C}$

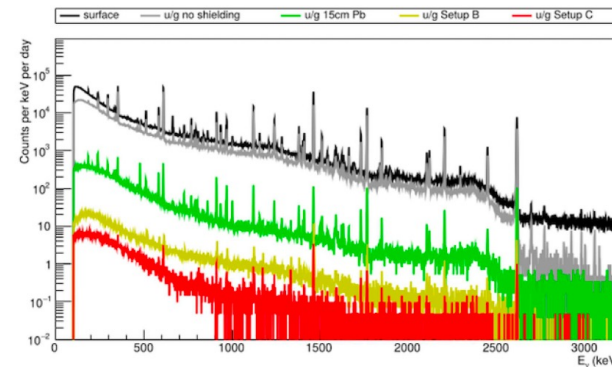
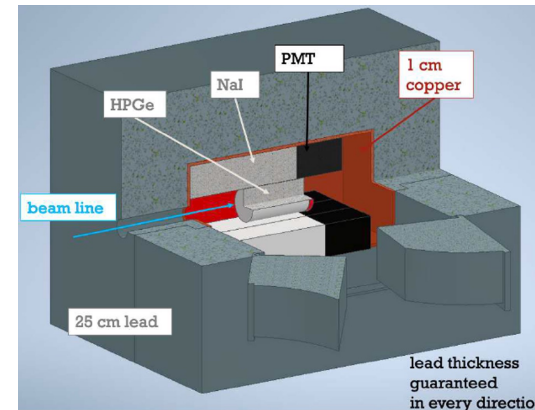
LUNA aims at measuring down to  $E=1500$  keV through detection of de-excitation  $\gamma$ -rays of residual nuclei ( $^{20}\text{Ne}$  or  $^{23}\text{Na}$ ) and/or detection of emitted particles (GASTLY ?)

Several tasks are being afforded:

1)  $^{12}\text{C}$  target production and characterization, with particular emphasis on the strategy to reduce and quantify residual H and D impurities that cause beam induced background

2) Setup for gamma measurement: a new ultra pure HPGe with 158% efficiency is being tested at LNGS

2-a) A proper shielding is being developed and the use of ancillary NaI detectors is under evaluation



## Conclusions

Many **intriguing open questions** in Nuclear Astrophysics.

**Common international effort** required to address some of these puzzles.

**Neutron source** in Red Giant Stars:

$^{13}\text{C}(\alpha, n)$  **well characterised, possible** some future activities @ n\_TOF & perhaps LUNA)

$^{22}\text{Ne}(\alpha, n)$  fair knowledge, ongoing activities (ASFIN & **LUNA**)

Heavy-Ion burning in Red Supergiants stars:

$^{12}\text{C} + ^{12}\text{C}$  still far from being accurately constrained, ongoing activities (ERNA, **LUNA**).

After Normalization to LUNA data, ASFIN data will provide S(E) in the whole Gamow window

## Acknowledgments

**Thank you for your attention**

**Thanks to the organizers**

**Many Thanks to  
L. Lamia & G. Pizzone (RL ASFIN)  
A. Di Leva (RN ERNA)  
G. Imbriani (RN LUNA)**





ALMA MATER STUDIORUM  
UNIVERSITÀ DI BOLOGNA

**Cristian Massimi**

Dipartimento di Fisica e Astronomia

[cristian.massimi@unibo.it](mailto:cristian.massimi@unibo.it)

[www.unibo.it](http://www.unibo.it)

# Backup

Indirect study of the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction via the  $^{13}\text{C}(^7\text{Li}, t)^{17}\text{O}$  transfer reaction

M. G. Pellegrini,<sup>1,\*</sup> F. Hammache,<sup>1,†</sup> P. Roussel,<sup>1</sup> L. Audouin,<sup>1</sup> D. Beausse,<sup>1</sup> P. Decouvremont,<sup>1</sup> S. Fortier,<sup>1</sup> L. Gaudetroy,<sup>1</sup> J. Kienzer,<sup>1</sup> A. Lefèvre-Suhlert,<sup>1</sup> M. Szafraniec,<sup>1</sup> V. Tatischeff,<sup>1</sup> and M. Vilmsay<sup>1</sup>  
<sup>1</sup>IPN-Orsay, IN2P3-CNRS, Université Paris XI, F-91406 Orsay, France  
<sup>2</sup>Physique Nucléaire Théorique et Physique Mathématique, ULB CP229, B-1050 Brussels, Belgium  
<sup>3</sup>GANIL, IN2P3-CNRS, Caen, France  
<sup>4</sup>CNSM, IN2P3-CNRS, Université Paris XI, F-91405 Orsay, France  
<sup>5</sup>GSI, Postfach 110552, D-64220 Darmstadt, Germany  
 (Received 31 May 2007; revised manuscript received 21 December 2007; published 22 April 2008)

The  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction is considered the main neutron source for the  $s$ -process in low mass asymptotic giant branch (AGB) stars. In the Gamow peak, the cross section sensitively depends on the  $1/2^+$  subthreshold state of  $^{16}\text{O}$  ( $E_x = 6.356$  MeV). In this work, we determined the astrophysical  $S$  factor through an evaluation of the  $\alpha$ -spectroscopic factor and the corresponding asymptotic normalization factor (ANC) of the 6.356 MeV state using the transfer reaction  $^{13}\text{C}(^7\text{Li}, t)^{17}\text{O}$  at two different incident energies. Our result confirms that the contribution of the  $1/2^+$  state is dominant at astrophysical energies. Our reaction rate at  $T = 0.09$  GK is slightly lower than the value adopted in the Nuclear Astrophysics Compilation of Reaction Rates (NACRE), but two times larger than the one obtained in a recent ANC measurement.

DOI: 10.1103/PhysRevC.77.042801

PACS number(s): 25.70.Hi, 25.55.Hp, 26.20.-f, 27.20.+n

Nearly half of the heavy elements observed in the universe are produced by a sequence of slow neutron capture reactions, the so-called  $s$ -process nucleosynthesis. In AGB stars of 1–3 solar masses at low temperatures, the description of this process critically depends on the neutron flux from the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction [1].

Direct measurements of the  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  cross section have been performed down to 270 keV [2], whereas in asymptotic giant branch (AGB) stars at temperatures around  $10^8$  K, the Gamow peak is at  $E_{\text{cm}} \sim 190$  keV.  $R$ -matrix extrapolations [3] of the cross sections measured at higher energies then have to be performed and have to include the contribution of the  $1/2^+$  state of  $^{16}\text{O}$  which lies at 6.356 MeV (3 keV below the  $\alpha$ - $^{16}\text{O}$  threshold). This contribution strongly depends on the  $\alpha$ -spectroscopic factor  $S_\alpha$  of this state. With values  $S_\alpha \sim 0.3$ – $0.7$  deduced from a nuclear model [4] and considered in the  $s$ -process modeling, a rise of the astrophysical  $S$  factor is expected when the energy decreases. This rise is compatible with the experimental data of Drotleff *et al.* [2] but their error bars are too large to derive definite conclusions.

From the analyses of the experimental results of Kubono *et al.* [5] on the  $^{13}\text{C}(^7\text{Li}, t)^{17}\text{O}$  transfer reaction, performed at the incident energy of 60 MeV, two different results are deduced: one by Kubono *et al.* [5],  $S_\alpha \sim 0.011$ , and the other one by Keeley *et al.* [6],  $S_\alpha \sim 0.4$ , which lies within the theoretical values [4]. In addition to these controversial analyses, the new study of the contribution of the  $1/2^+$  state to the total  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  reaction rate by a  $^{13}\text{C}(^7\text{Li}, t)^{17}\text{O}$  asymptotic normalization factor (ANC) measurement [7] led to an astrophysical  $S$  factor at the energy of 190 keV ten times smaller than the value adopted in the Nuclear Astrophysics

Compilation of Reaction Rates (NACRE) [8] but five times larger than the value deduced by Kubono *et al.* Consequently, it appeared highly desirable to perform a new determination of this  $S_\alpha$  factor.

In this article, we report on a study of the 6.356 MeV state and of three other excited states at 3.055 MeV ( $1/2^-$ ), 4.553 MeV ( $3/2^-$ ), and 7.380 MeV ( $5/2^-$ ) by means of the  $\alpha$  transfer reaction  $^{13}\text{C}(^7\text{Li}, t)^{17}\text{O}$ . Two incident energies, 28 and 34 MeV, were used to check the direct mechanism character of our transfer reaction. The use of  $^7\text{Li}$  instead of  $^6\text{Li}$  reduces possible multiplet effects [9] and the transfer cross sections to low spin states are enhanced because of the nonzero  $\alpha$  angular momentum in  $^7\text{Li}$  as observed when comparing the transfer reactions ( $^6\text{Li}, t$ ) and ( $^7\text{Li}, t$ ) on  $^{13}\text{C}$  [10,11]. Moreover, an estimate of the systematic errors associated with the DWBA analysis can be evaluated from the comparison of the results of different reactions involving different angular momenta and performed at different energies.

The experiment was performed with a  $^7\text{Li}^{3+}$  beam at the Orsay TANDEM. Two self-supporting enriched  $^{13}\text{C}$  targets, with 72(4) and 133(7)  $\mu\text{g}/\text{cm}^2$  of  $^{13}\text{C}$  and an initial purity of 90%, were used. A  $^{12}\text{C}$  target of 80(4)  $\mu\text{g}/\text{cm}^2$  was also used for calibration and background subtraction. Despite a rather low gas pressure in the reaction chamber ( $\sim 10^{-5}$  mb), a  $^{12}\text{C}$  buildup was observed during the whole experiment ending with an amount of  $^{12}\text{C}$  in the enriched  $^{13}\text{C}$  target about twice its initial value. This buildup up was of least consequence as shown in the analysis described below and was however duly monitored. Note that a more enriched target would not have brought a big improvement.

The absolute (and constant) amount of  $^{13}\text{C}$  and the final amount of  $^{12}\text{C}$  in the  $^{13}\text{C}$  targets were deduced from  $\alpha$  energy loss measurements on all targets and by comparing the ratio between 28 MeV elastically scattered  $^7\text{Li}$  particles at  $21^\circ$  from the  $^{12}\text{C}$  nuclei in  $^{12}\text{C}$  and  $^{13}\text{C}$  targets (see Figs. 1(a) and 1(b)) measured in a run at the end of the experiment.

\* Present address: Dipartimento di Fisica e Astronomia, Università di Catania and Laboratori Nazionali del Sud - INFN, Catania, Italy.  
 † Corresponding author: hammache@ipn.in2p3.fr

Measurement of the  $-3$  keV Resonance in the Reaction  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  of Importance in the  $s$ -Process

M. La Cognata,<sup>1,\*</sup> C. Spitaleri,<sup>1,2</sup> O. Trippella,<sup>1,3</sup> G. G. Kiss,<sup>1,4</sup> G. V. Rogachev,<sup>5</sup> A. M. Mukhamedzhanov,<sup>6</sup> M. Avila,<sup>5</sup> G. L. Guardo,<sup>1,2</sup> E. Koshchiy,<sup>5</sup> A. Kuchera,<sup>5</sup> L. Lamia,<sup>2</sup> S. M. R. Puglia,<sup>1,2</sup> S. Romano,<sup>1,2</sup> D. Santiago,<sup>5</sup> and R. Spartà<sup>1,2</sup>

<sup>1</sup>Istituto Nazionale di Fisica Nucleare, Laboratori Nazionali del Sud, 95123 Catania, Italy

<sup>2</sup>Dipartimento di Fisica e Astronomia, Università di Catania, 95123 Catania, Italy

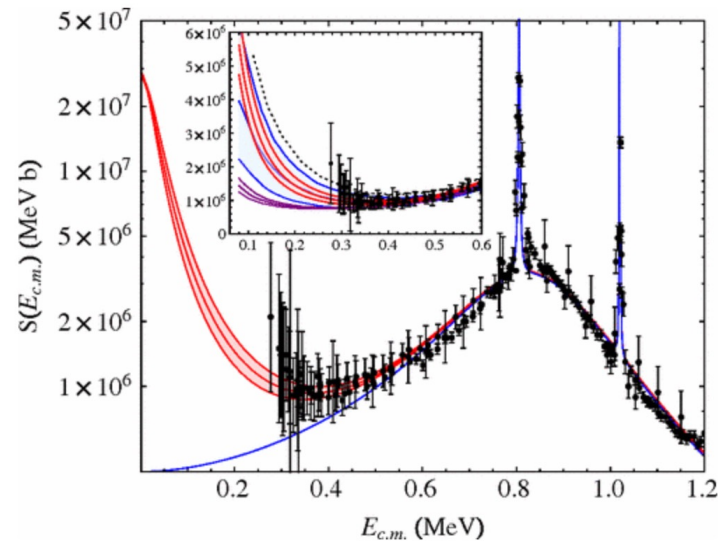
<sup>3</sup>Istituto Nazionale di Fisica Nucleare, Sezione di Perugia, Perugia, Italy and Dipartimento di Fisica, Università di Perugia, 06123 Perugia, Italy

<sup>4</sup>Institute of Nuclear Research (ATOMKI), 4026 Debrecen, Hungary

<sup>5</sup>Department of Physics, Florida State University, Tallahassee, Florida 32306, USA

<sup>6</sup>Cyclotron Institute, Texas A&M University, College Station, Texas 77843, USA

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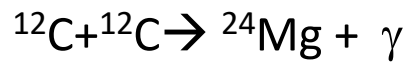
## Backup $^{12}\text{C}+^{12}\text{C}$



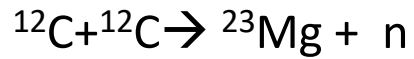
$$Q = 4.62 \text{ MeV}$$



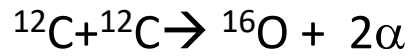
$$Q = 2.24 \text{ MeV}$$



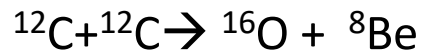
$$Q = 13.93 \text{ MeV}$$



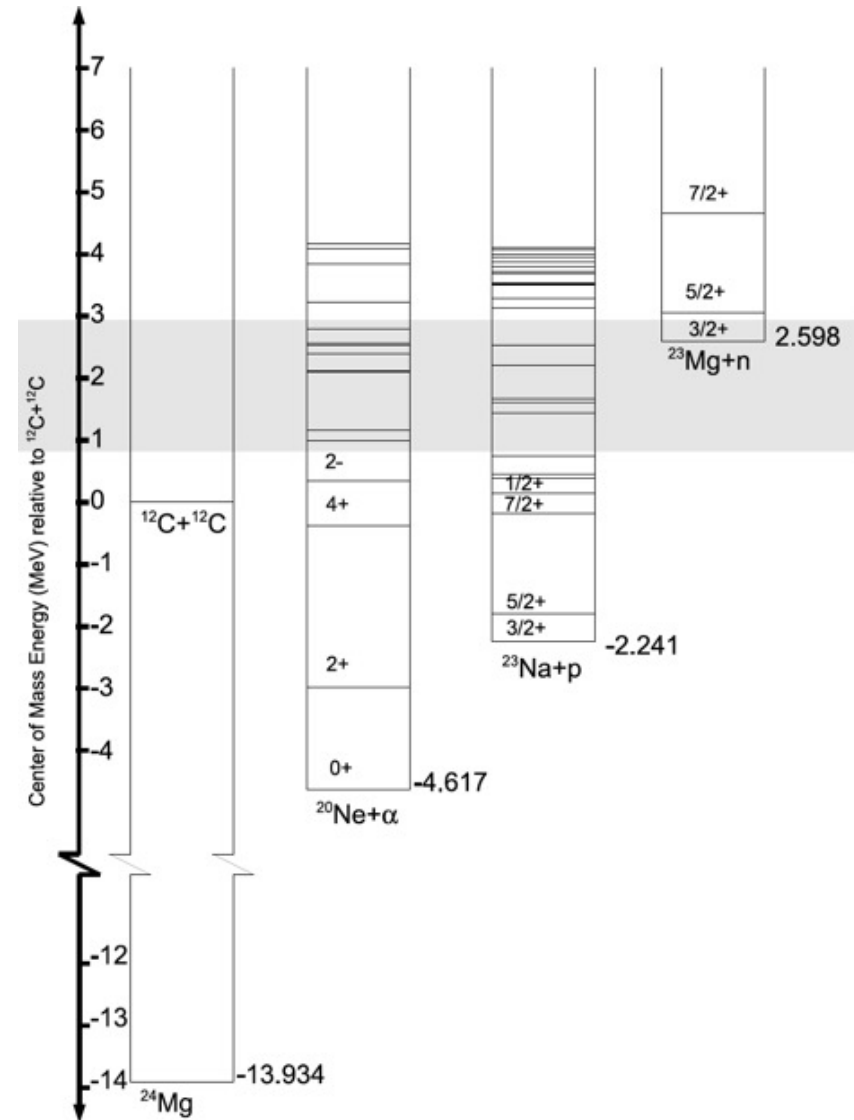
$$Q = -2.62 \text{ MeV}$$



$$Q = -0.12 \text{ MeV}$$



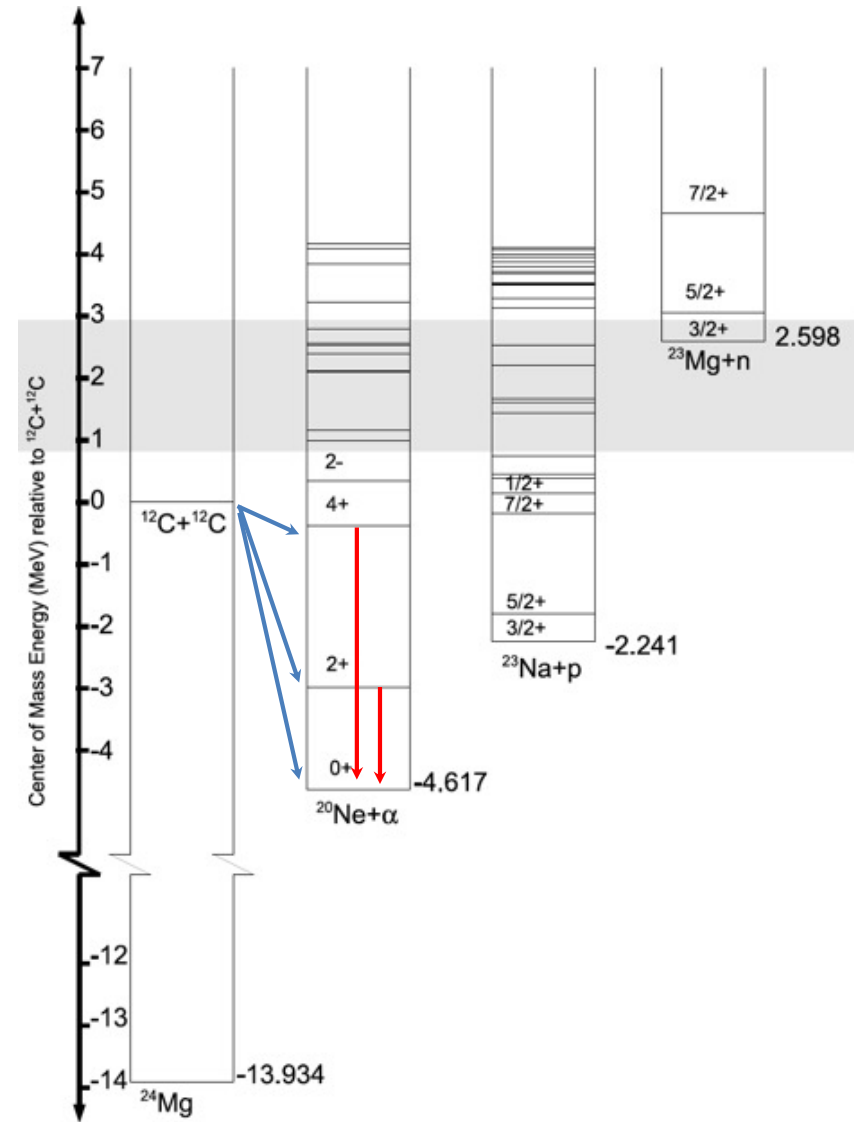
$$Q = -0.21 \text{ MeV}$$



# Backup $^{12}\text{C}+^{12}\text{C}$



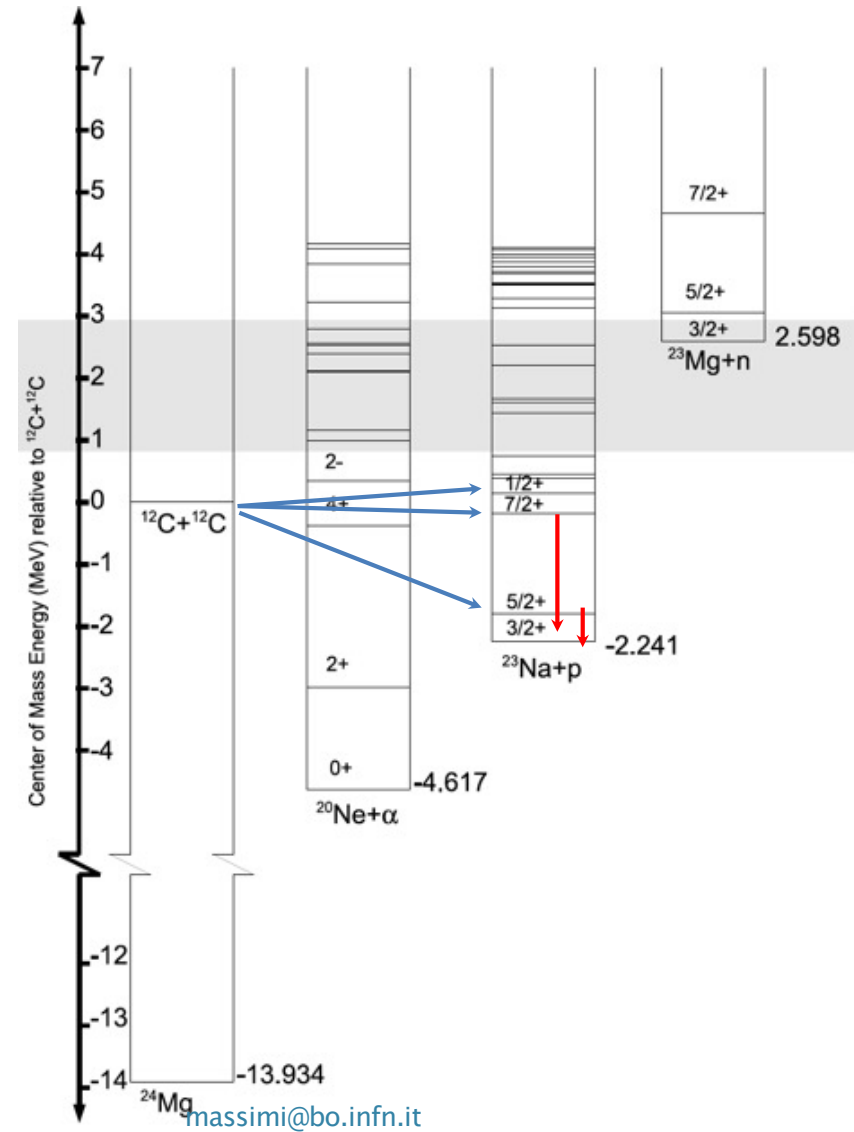
$\gamma$ -rays and $\alpha$ particles energies for excited states for $^{12}\text{C}(^{12}\text{C}, \alpha)^{20}\text{Ne}$ ( $Q = 4.617 \text{ MeV}$ )				
$E_x$ (MeV)	$J^p$	Main $\gamma$ transitions (MeV)	ID	$E_{\alpha\text{-max}}$ (MeV) ( $E^{\text{CM}} = 2 \text{ MeV}$ )
0.0	$0^+$		$\alpha_0$	8.6
1.63	$2^+$	$1.63 \rightarrow 0$ 1.63	$\alpha_1$	6.8
4.24	$4^+$	$4.24 \rightarrow 1.63$ 2.61	$\alpha_2$	3.9
4.96	$2^-$	$4.96 \rightarrow 1.63$ 3.33	$\alpha_3$	3.1
5.62	$3^-$	$5.62 \rightarrow 1.63$ 3.98	$\alpha_4$	2.2
5.78	$1^-$	$5.78 \rightarrow 1.63$ 4.15 $5.78 \rightarrow 0$ 5.78	$\alpha_5$	2.0



# Backup $^{12}\text{C}+^{12}\text{C}$



$\gamma$ -rays and p particles energies for excited states for $^{12}\text{C}(^{12}\text{C}, \text{p})^{23}\text{Na}$ ( $Q = 2.241 \text{ MeV}$ )					
$E_x$ (MeV)	$J^p$	Main $\gamma$ transitions (MeV)		ID	$E_{p\text{-max}}$ (MeV) ( $E_{\text{CM}} = 2 \text{ MeV}$ )
0.0	$3/2^+$			$p_0$	5.3
0.44	$5/2^+$	$0.44 \rightarrow 0$ <b>0.44</b>		$p_1$	4.8
2.07	$7/2^+$	$2.07 \rightarrow 0.44$ <b>1.63</b>		$p_2$	3
2.39	$1/2^+$	$2.39 \rightarrow 0.44$ <b>1.95</b>	$2.39 \rightarrow 0$ <b>2.39</b>	$p_3$	2.6
2.64	$1/2^-$	$2.64 \rightarrow 0$ <b>2.64</b>		$p_4$	2.3
2.70	$9/2^+$	$2.70 \rightarrow 2.07$ <b>0.62</b>	$2.70 \rightarrow 0.44$ <b>2.26</b>	$p_5$	2.3
2.98	$3/2^+$	$2.98 \rightarrow 0.44$ <b>2.54</b>	$2.98 \rightarrow 0$ <b>2.98</b>	$p_6$	1.9



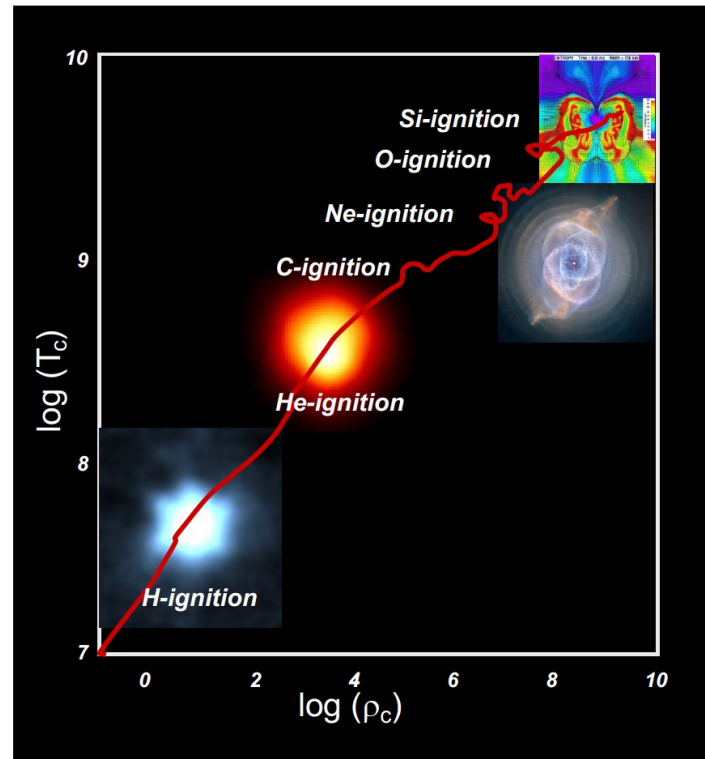
# Backup $^{12}\text{C}+^{12}\text{C}$

different burning phases  
characterize the evolution  
of a „massive“ star

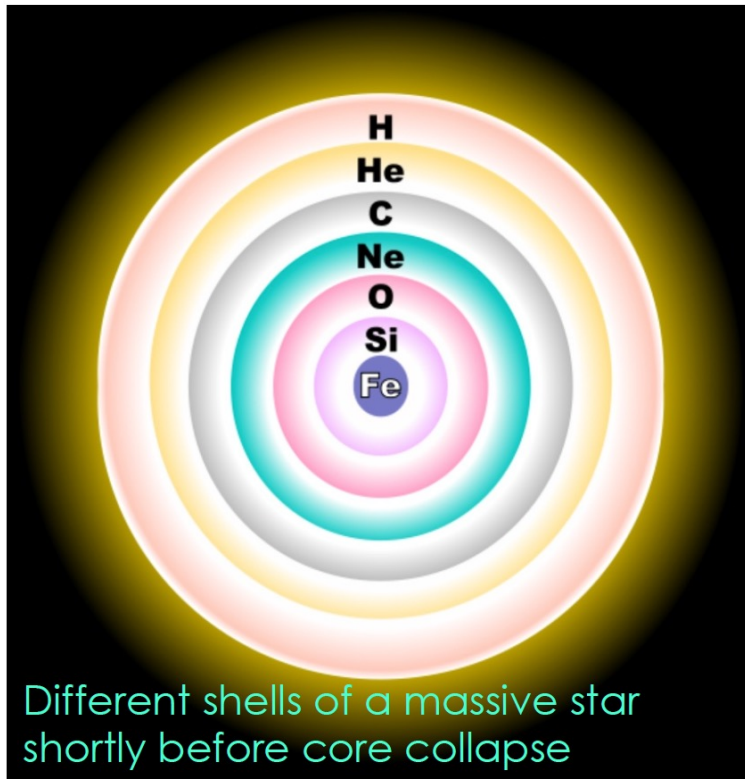


each burning phase is controlled  
by different nuclear reactions,  
which govern the:

- energy production
- time scale
- nucleosynthesis



# Backup $^{12}\text{C}+^{12}\text{C}$



- key reactions at each stage of stellar burning

Fuel	Main Product	Secondary Product	T (10 <sup>9</sup> K)	Time (yr)	Main Reaction
H	He	<sup>14</sup> N	0.02	10 <sup>7</sup>	$4\text{H} \xrightarrow{\text{CNO}} \text{He}$
He	O, C	<sup>18</sup> O, <sup>22</sup> Ne s-process	0.2	10 <sup>6</sup>	$3\text{He}^4 \rightarrow \text{C}$ $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$
C	Ne, Mg	Na	0.8	10 <sup>3</sup>	$^{12}\text{C} + ^{12}\text{C}$
Ne	O, Mg	Al, P	1.5	3	$^{20}\text{Ne}(\gamma, \alpha)^{16}\text{O}$ $^{20}\text{Ne}(\alpha, \gamma)^{24}\text{Mg}$
O	Si, S	Cl, Ar, K, Ca	2.0	0.8	$^{16}\text{O} + ^{16}\text{O}$
Si	Fe	Ti, V, Cr, Mn, Co, Ni	3.5	0.02	$^{28}\text{Si}(\gamma, \alpha)\dots$

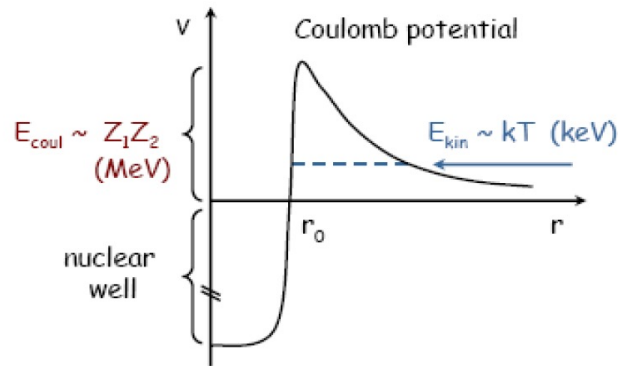
- In a star of 8-11 Solar masses, a carbon flash lasts just milliseconds.
- In a star of 25 Solar masses carbon burning lasts about 600 years.



# Backup

charged particles  $\rightarrow$  Coulomb barrier

energy available: from thermal motion

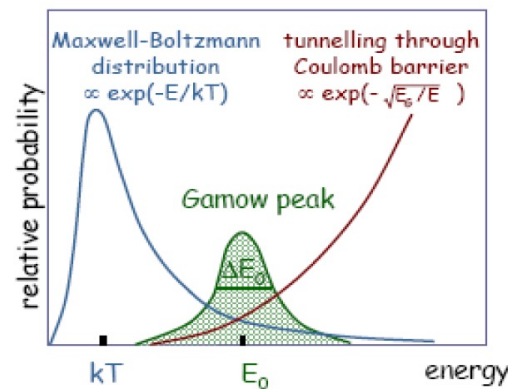


$T \sim 15 \times 10^6 \text{ K}$  (e.g. our Sun)  $\Rightarrow kT \sim 1 \text{ keV}$

during static burnings:  $kT \ll E_{\text{coul}}$

reactions occur through TUNNEL EFFECT

$\rightarrow$  tunneling probability  $P \propto \exp(-2\pi\eta)$



$$\sigma(E) = \frac{1}{E} \exp(-2\pi\eta) S(E)$$

non-nuclear origin  
STRONG energy  
dependence

nuclear origin  
WEAK energy  
dependence

**ASTROPHYSICAL S(E)-FACTOR**