



Peter Schuck  
(1940-2022)

## ***Solving nuclear structure puzzles with the relativistic nuclear field theory***

*“Ab-initio”  
Equation of Motion  
(EOM) framework*

*Nuclear Field Theory  
(NFT)*



Ricardo Broglio  
(1939-2022)

*Elena Litvinova*

*Western Michigan University*



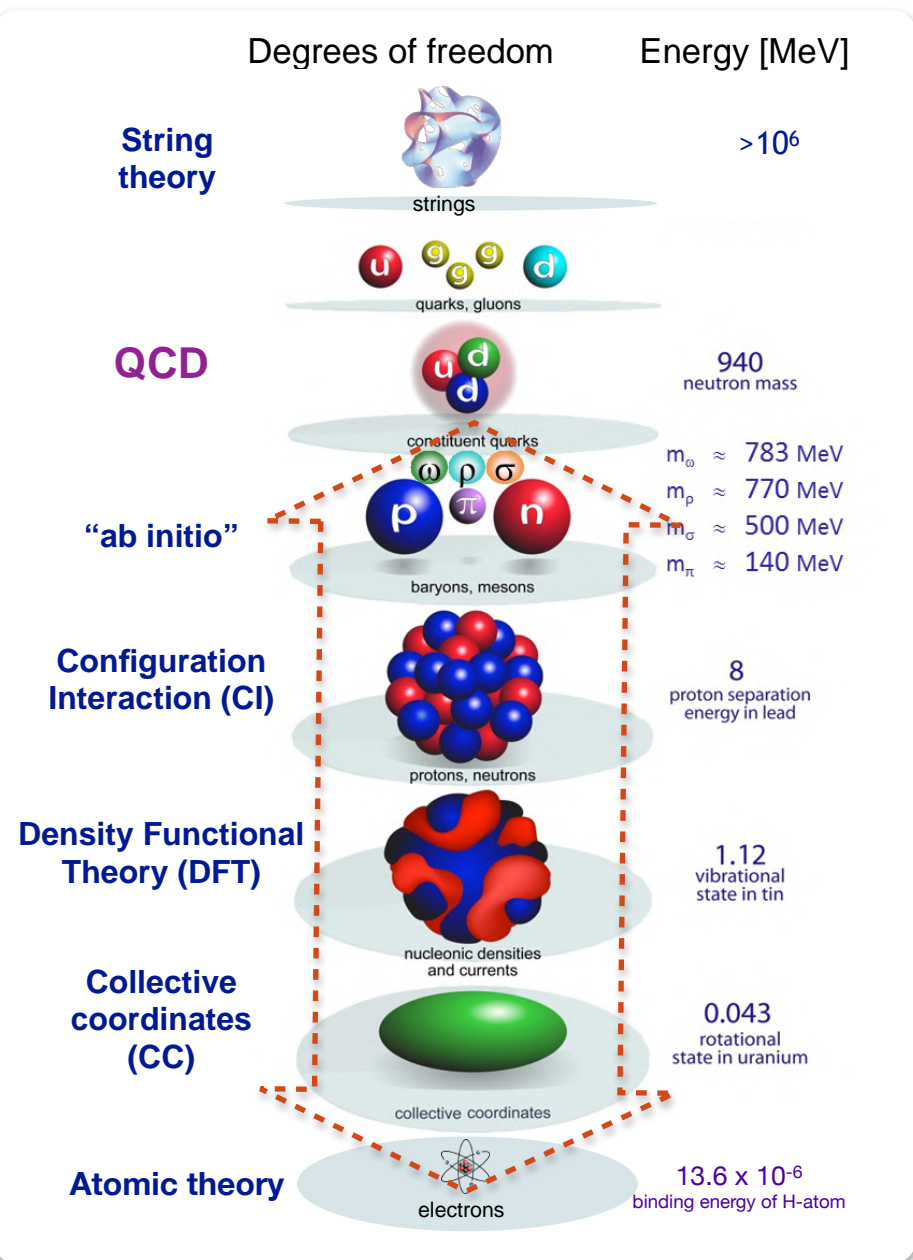
**MICHIGAN STATE  
UNIVERSITY**



***Collaborators:*** Peter Schuck, Peter Ring, Yinu Zhang, Caroline Robin, Herlik Wibowo

***Outline:*** (i) Theory (ii) Results (iii) Outlook

# Hierarchy of energy scales and nuclear many-body problem



• **The major conflict:**

Separation of energy scales => effective field theories

VS

The physics on a certain scale is governed by the next higher-energy scale

**Hamiltonian:**

$$H = K + V$$

center of mass

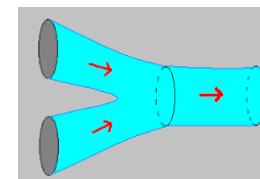
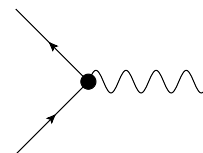
internal degrees of freedom:  
next energy scale

**Standard Model:**

free propagation and interaction, singularities & renormalizations

**String theory:**

merging strings  
NO “Interaction”

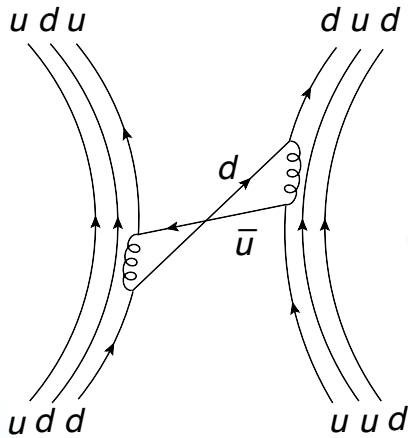


• **Possible solution:**

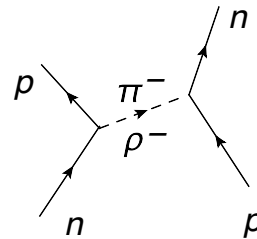
- Keep/establish connections between the scales via emergent phenomena
- A universal approach to the strongly-coupled QMBP?

# The underlying mechanism of NN-interaction:

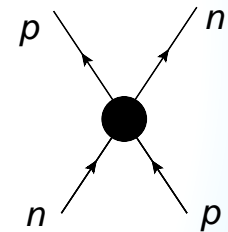
Quantum Chromodynamics (QCD, high energy)



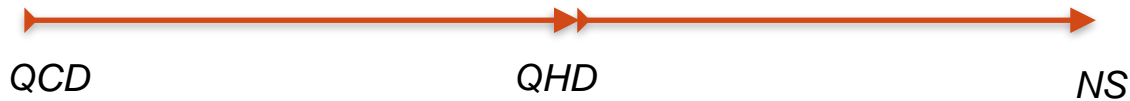
Quantum Hadrodynamics (QHD, intermediate energy)



Nuclear Structure (NS, low energy)

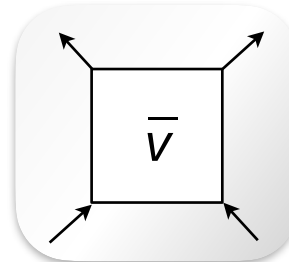


Relay of EFTs:



## Formalism:

- Generic bare "interaction": model-independent, all channels included
- Higher-orders are treated via **in-medium propagators**
- No perturbation theory



## In implementations:

- Meson-exchange (ME) at leading order
- Effective coupling constants/masses (adjusted on the mean-field (MF) level, NL3(\*)) + subtraction of qPVC
- Bare ME + subtraction of MF artifacts (in progress)

# A strongly-correlated many body system: single-fermion propagator, particle-hole propagator and related observables

$$H = \sum_{12} \bar{\psi}_1 (-i\gamma \cdot \nabla + M)_{12} \psi_2 + \frac{1}{4} \sum_{1234} \bar{\psi}_1 \bar{\psi}_2 \bar{v}_{1234} \psi_4 \psi_3 = T + V(2)$$

**Hamiltonian,**  
extendable to 3B forces  
(3BFs are minimized in  
covariant theories)

$$G_{11'}(t - t') = -i \langle T \psi(1) \bar{\psi}(1') \rangle \quad 1 = \{\xi_1, t\}$$

**Single-particle  
propagator**

Fourier transform:  
Spectral  
expansion

$$G_{11'}(\varepsilon) = \sum_n \frac{\eta_1^n \bar{\eta}_{1'}^{n*}}{\varepsilon - \varepsilon_n^+ + i\delta} + \sum_m \frac{\chi_1^m \bar{\chi}_{1'}^{m*}}{\varepsilon + \varepsilon_m^- - i\delta}$$

**Residues** - spectroscopic  
(occupation) factors

$$\eta_1^n = \langle 0^{(N)} | \psi_1 | n^{(N+1)} \rangle \quad \chi_1^m = \langle m^{(N-1)} | \psi_1 | 0^{(N)} \rangle$$

Ground state of  
N particles

(Excited) state  
of (N+1) particles

**Poles** - single-particle  
energies

$$R_{12,1'2'}(t - t') = -i \langle T (\bar{\psi}_1 \psi_2)(t) (\bar{\psi}_{2'} \psi_{1'})(t') \rangle$$

**Particle-hole response  
function**

Fourier transform: Spectral  
expansion

$$R_{12,1'2'}(\omega) = \sum_{\nu > 0} \left[ \frac{\rho_{21}^\nu \bar{\rho}_{2'1'}^{\nu*}}{\omega - \omega_\nu + i\delta} - \frac{\bar{\rho}_{12}^{\nu*} \rho_{1'2'}}{\omega + \omega_\nu - i\delta} \right]$$

**Residues** - transition  
densities

Excitation  
energies

$$\rho_{12}^\nu = \langle 0 | \bar{\psi}_2 \psi_1 | \nu \rangle$$

**Poles** - excitation energies



# Exact equations of motion (EOM) for binary interactions: one-body problem

One-fermion propagator

$$G_{11'}(t - t') = -i \langle T \psi(1) \bar{\psi}(1') \rangle$$

**EOM: Dyson Eq.**

$$G(\omega) = G^{(0)}(\omega) + G^{(0)}(\omega) \Sigma(\omega) G(\omega) \quad (*) \quad \Sigma(\omega) = \Sigma^{(0)} + \Sigma^{(r)}(\omega)$$

**Irreducible kernel (Self-energy, exact):**

Instantaneous term (Hartree-Fock incl. "tadpole")  
**Short-range correlations**

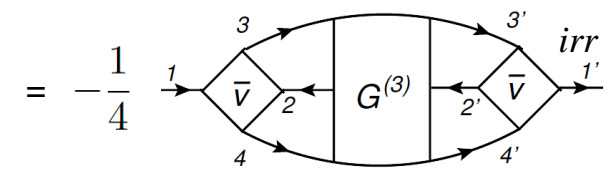
$$\Sigma_{11'}^{(0)} = - \langle \gamma^0 \{ [V, \psi_1], \bar{\psi}_{1'} \} \gamma^0 \rangle$$

$$= \sum_{22'} \bar{v}_{121'2'} \langle \bar{\psi}_2 \psi_{2'} \rangle = \text{Diagram: } \begin{array}{c} \text{circle with } \rho_{2'2} \text{ and arrow} \\ \text{square with } \bar{v} \end{array}$$

t-dependent (dynamical) term (symmetric version): **Long-range correlations**

$$\Sigma_{11'}^{(r)} = i \langle T \gamma^0 [V, \psi_1](t) [V, \bar{\psi}_{1'}](t') \gamma^0 \rangle^{irr}$$

$$= -\frac{1}{4} \sum_{234} \sum_{2'3'4'} \bar{v}_{1234} G_{432', 23'4'}^{(3)irr}(t - t') \bar{v}_{4'3'2'1'}$$



$$\rho_{11'} = -i \lim_{t=t'-0} G_{11'}(t - t')$$

is the full solution of (\*):  
**includes the dynamical term!**

Koltun-Migdal-Galitsky sum rule: **the binding energy**

**"Ab-initio DFT":**

$$E_0 = \frac{1}{2\pi} \int_{-\infty}^{\bar{\epsilon}_F} d\epsilon \sum_{12} (T_{12} + \epsilon \delta_{12}) \text{Im} G_{21}(\epsilon)$$

# Equation of motion (EOM) for the particle-hole response

Particle-hole propagator  
(response function):

$$R_{12,1'2'}(t - t') = -i \langle T(\bar{\psi}_1 \psi_2)(t) (\bar{\psi}_{2'} \psi_{1'})(t') \rangle$$

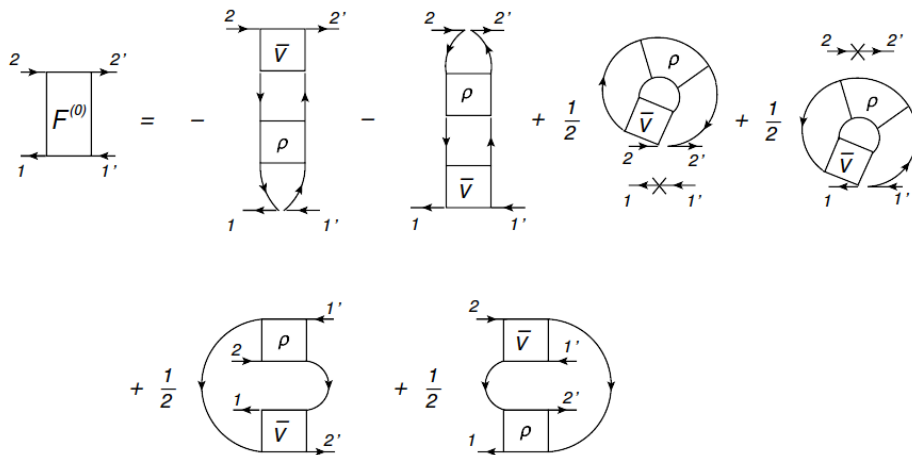
spectra of excitations,  
masses, decays, ...

**EOM: Bethe-Salpeter-Dyson Eq.**

$$R(\omega) = R^{(0)}(\omega) + R^{(0)}(\omega)F(\omega)R(\omega) \quad (**) \quad F(t - t') = F^{(0)}\delta(t - t') + F^{(r)}(t - t')$$

**Irreducible kernel (exact):**

Instantaneous term (“bosonic” mean field):  
**Short-range correlations**

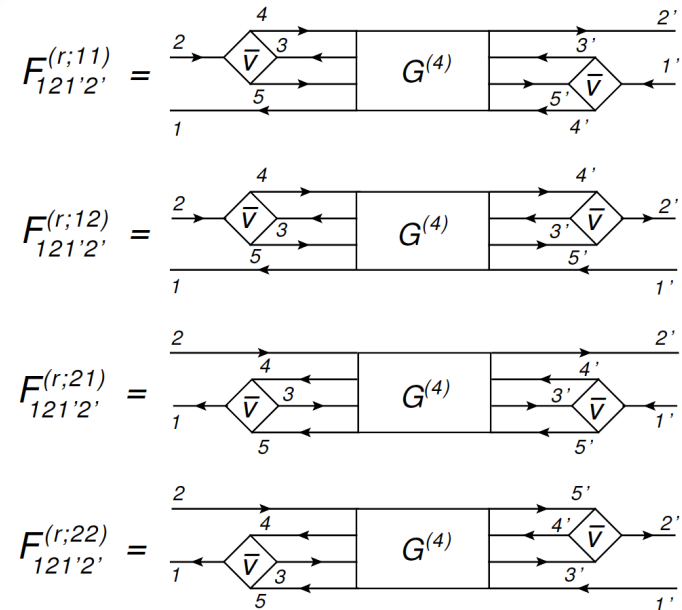


Self-consistent mean field  $F^{(0)}$ , where

$$\rho_{12,1'2'} = \delta_{22'}\rho_{11'} - i \lim_{t' \rightarrow t+0} R_{2'1,21'}(t - t')$$

contains the full solution of (\*\*) including the dynamical term!

$t$ -dependent (dynamical) term:  
**Long-range correlations**



$$F_{12,1'2'}^{(r)}(t - t') = \sum_{ij} F_{12,1'2'}^{(r;ij)}(t - t')$$

# Non-perturbative treatment of two-point $G^{(n)}$ in the dynamical kernels

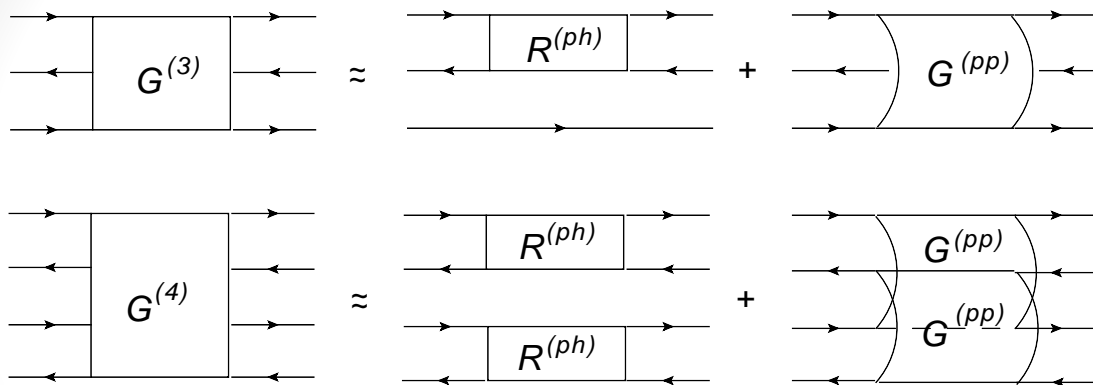
• **Quantum many-body problem in a nutshell:** Direct EOM for  $G^{(n)}$  generates  $G^{(n+2)}$  in the (symmetric) dynamical kernels and further high-rank correlation functions (CFs); an equivalent of the BBGKY hierarchy.  $N_{\text{Equations}} = N_{\text{Particles}} \& \text{ Coupled}$  🙈 !!!

Truncation on two-body level

• **Non-perturbative solutions:** **Cluster decomposition**

$\blacklozenge G^{(3)} = G^{(1)} G^{(1)} G^{(1)} + G^{(2)} G^{(1)} + \Xi^{(3)}$   
“Self-consistent GFs” This work  
“Second RPA” This work

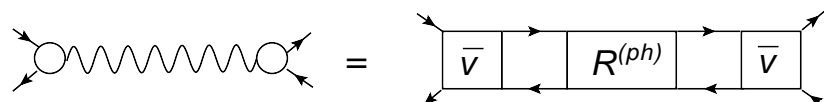
$\blacklozenge G^{(4)} = G^{(1)} G^{(1)} G^{(1)} G^{(1)} + G^{(2)} G^{(2)} + G^{(3)} G^{(1)} + \Xi^{(4)}$



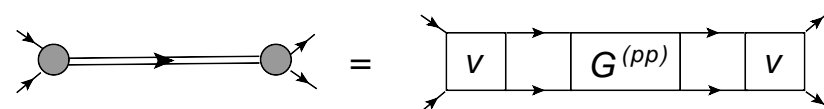
- P. C. Martin and J. S. Schwinger, *Phys. Rev.* 115, 1342 (1959).
- N. Vinh Mau, *Trieste Lectures* 1069, 931 (1970)
- P. Danielewicz and P. Schuck, *Nucl. Phys.* A567, 78 (1994)
- ...

Exact mapping: particle-hole ( $2q$ ) quasibound states

Emergence of effective “particles” (phonons, vibrations):



Emergence of superfluidity:

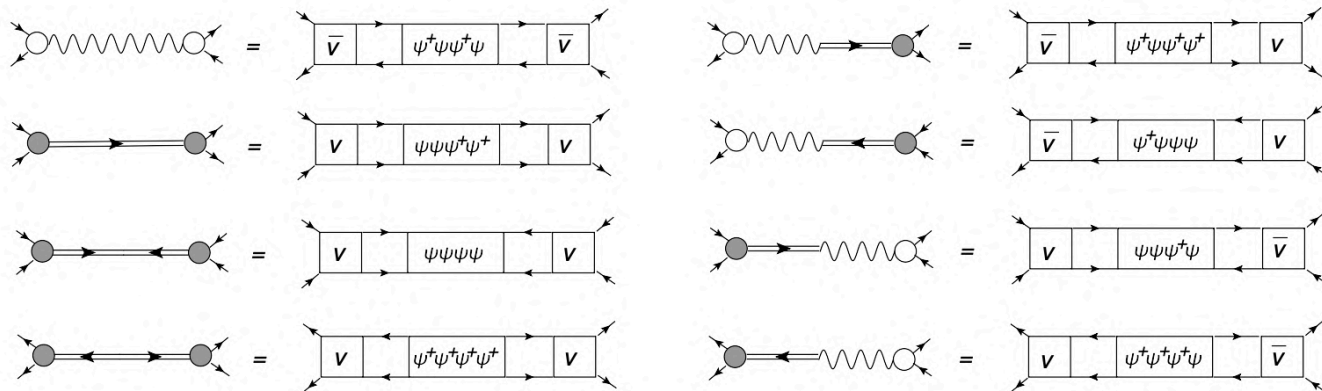


**The genesis**  
of the  
quasiparticle-  
vibration coupling  
(qPVC) in nuclei

# "Ab-initio" qPVC in superfluid systems

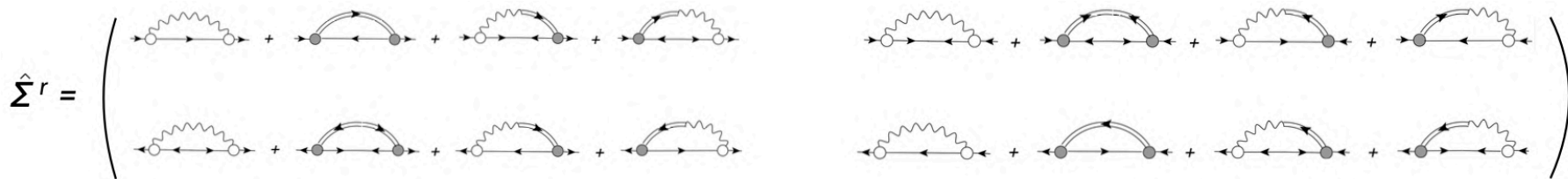
Superfluid dynamical kernel: adding particle-number violating contributions

Mapping on the qPVC in the canonical basis

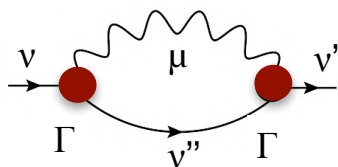


Quasiparticle dynamical self-energy (matrix):

*normal and pairing phonons are unified*



Bogoliubov transformation



Cf.: Quasiparticle static self-energy (matrix) in HFB

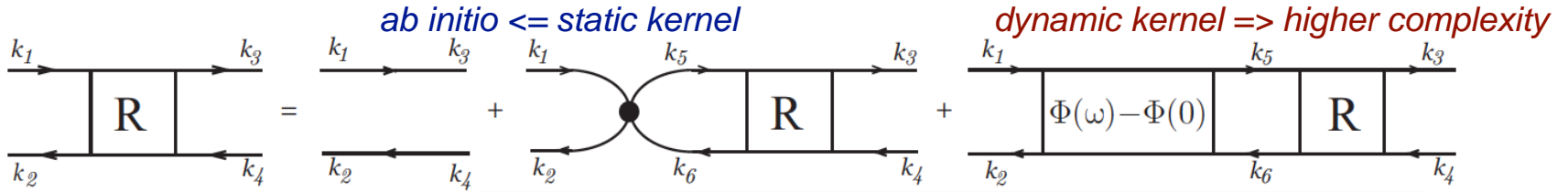
$$\hat{\Sigma}^0 = \begin{pmatrix} \tilde{\Sigma}_{11'} & \Delta_{11'} \\ -\Delta_{11'}^* & -\tilde{\Sigma}_{11'}^T \end{pmatrix}$$

E.L., Y. Zhang, PRC 104, 044303 (2021)

Y. Zhang et al., PRC 105, 044326 (2022)



# Nuclear response: toward a complete theory

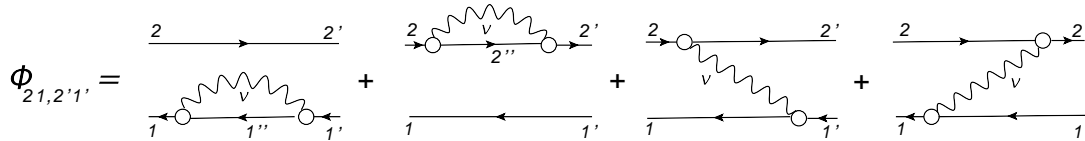


Dyson-Bethe-Salpeter Equation:

$$R(\omega) = R^0(\omega) + R^0(\omega) [V + \Phi(\omega) - \Phi(0)] R(\omega)$$

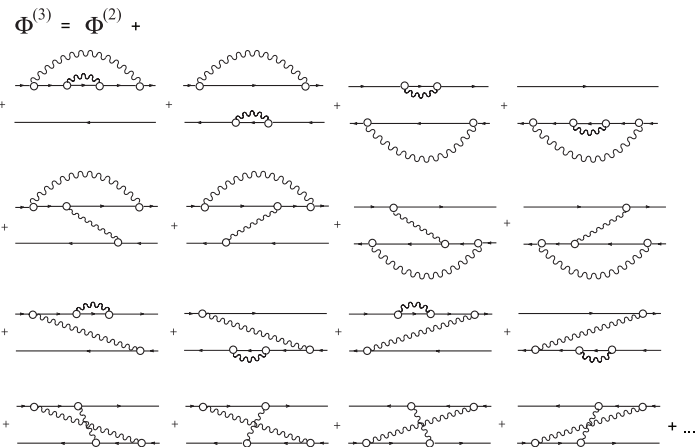
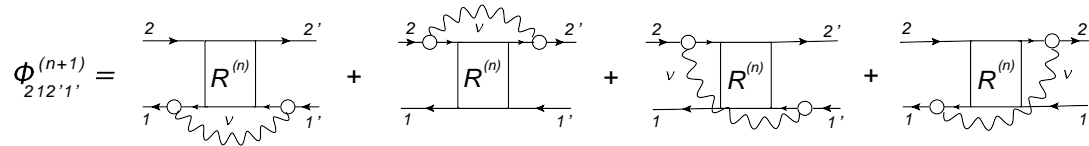
Conventional NFT

Cf. P.-F. Bortignon,  
G. Colò, E. Vigezzi,  
G. Potel, F. Barranco  
&  
V. Tselyaev (t-blocking)



Subtraction  
for effective  
interactions  
(Tselyaev 2013)

Extended NFT:



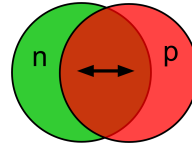
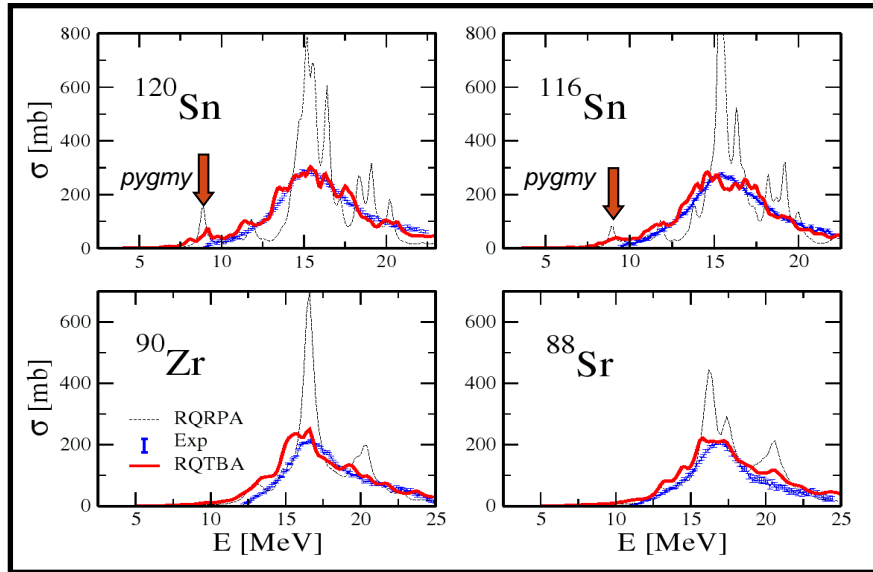
Generalized approach for the correlated propagators

n-th order: E.L. PRC 91, 034332 (2015)

Ab-initio formulation,  
 $\Phi^{(3)}$  implementation; 2q+2phonon correlations:  
E.L., P. Schuck, PRC 100, 064320 (2019)

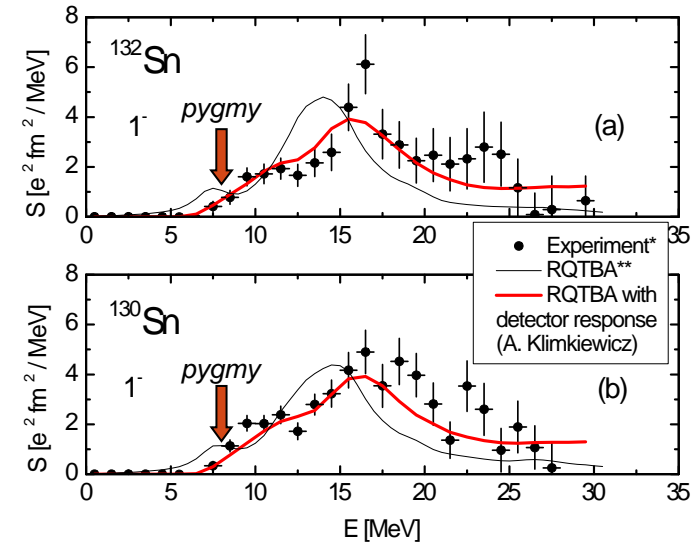
# The giant resonance width puzzle: Relativistic Quasiparticle Time Blocking Approximation (RQTBA)

## Giant dipole resonance (GDR) in stable nuclei



Giant & pygmy dipole resonances

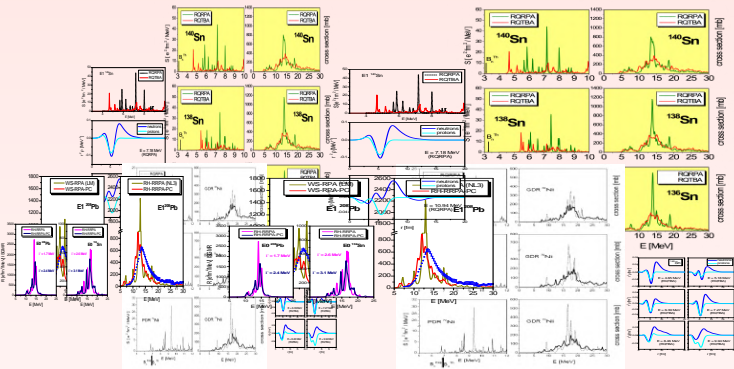
## Neutron-rich Sn



\* P. Adrich et al.,  
PRL 95, 132501 (2005)

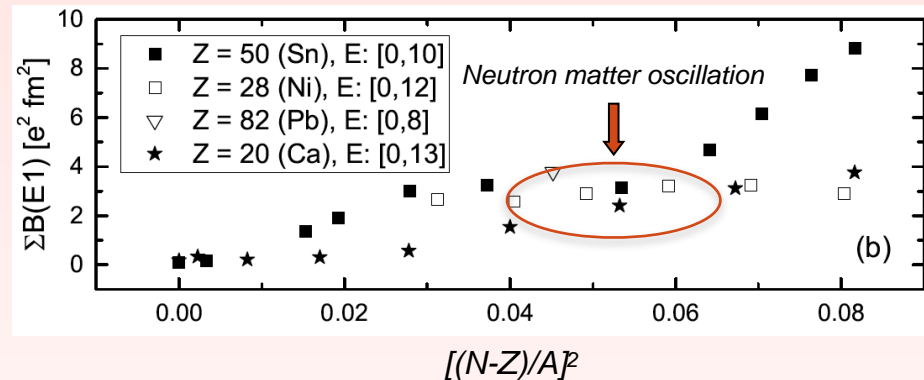
\*\* E. L., P. Ring, and V. Tselyaev,  
Phys. Rev. C 78, 014312 (2008)

## Systematic GMR calculations (various multipoles)



~50+ works on various GMRs

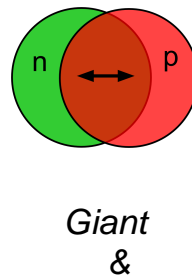
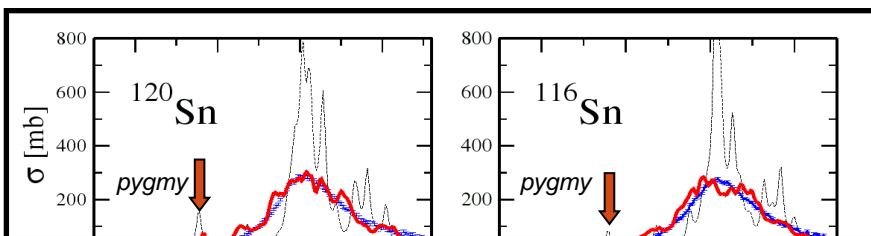
## Pygmy dipole strength systematics (important for EOS)



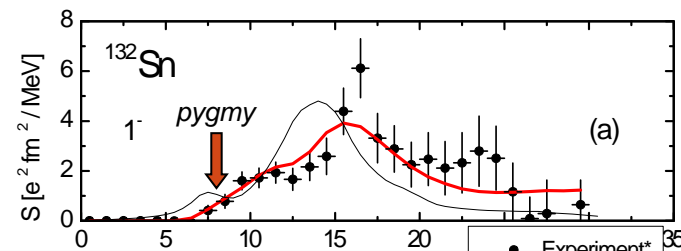
I.A. Egorova, E. Litvinova, Phys. Rev. C 94, 034322 (2016)

# Dipole response in medium-mass and heavy nuclei within Relativistic Quasiparticle Time Blocking Approximation (RQTBA)

## Giant dipole resonance (GDR) in stable nuclei



## Neutron-rich Sn



**Thanks for the collaboration, discussions and references:**

Victor Tselyaev  
Marcello Baldo  
Gianluca Colò  
Armand Bahini  
Luna Pellegrini  
Retief Neveling  
Marine Vandebrouck  
Francisco Barranco  
Enrico Vigezzi  
Yifei Niu  
Achim Richter  
Thomas Aumann

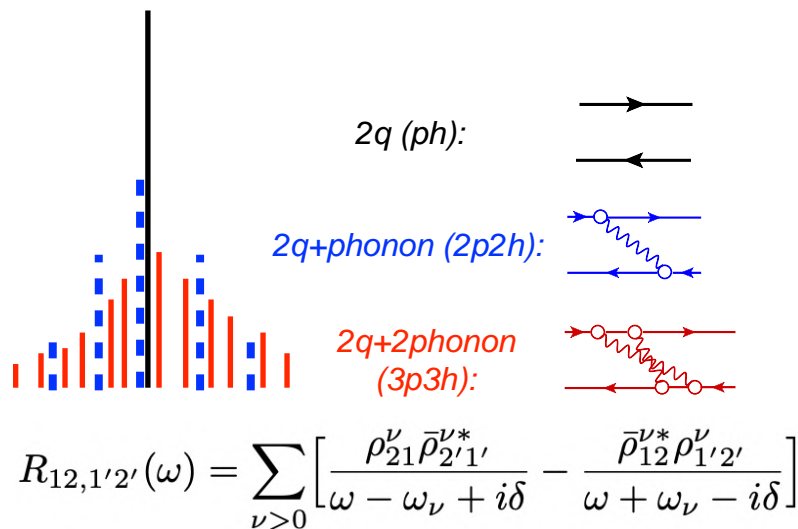
Deniz Savran  
Andreas Zilges  
Michael Weinert  
Maria Markova  
Hiroyuki Sagawa  
Remco Zegers  
Danilo Gambacurta  
Masaki Sasano  
Nils Paar  
Peter von Neumann-  
Cosel

Gregory Potel  
František Knapp  
Javier Roca-Maza  
Oliver Wieland  
Angela Bracco  
Edoardo Lanza  
Mark Spieker  
Adam Maj  
Esra Yüksel  
Atsushi Tamii  
Iolanda Matea

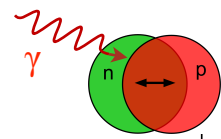
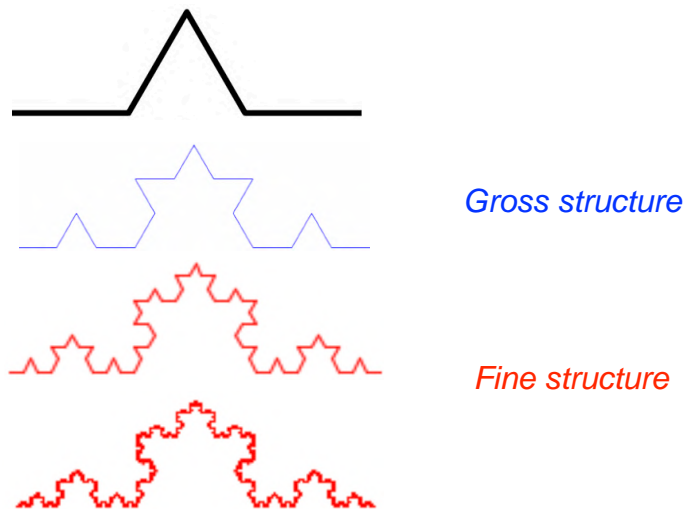
**and Others!**

# Leading NFT (q)PVC is insufficient: the “3p3h” configurations

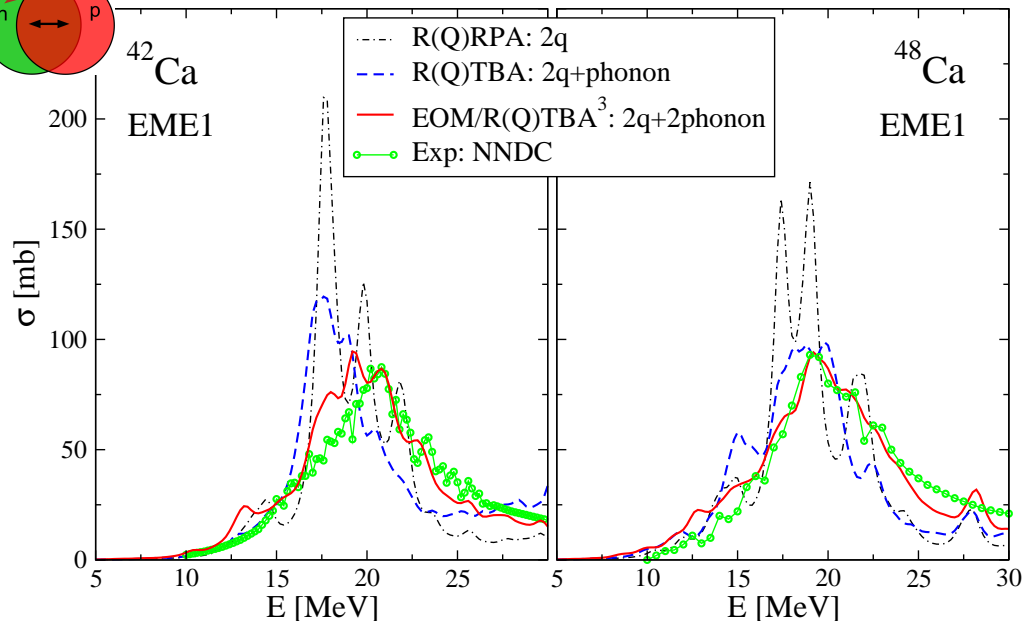
## Fragmentation mechanism



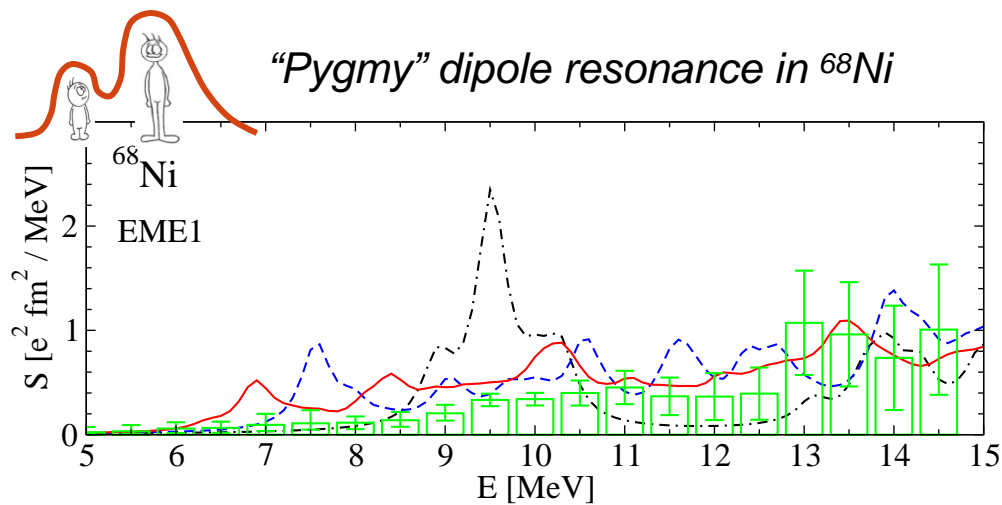
## Fractals: Koch curve



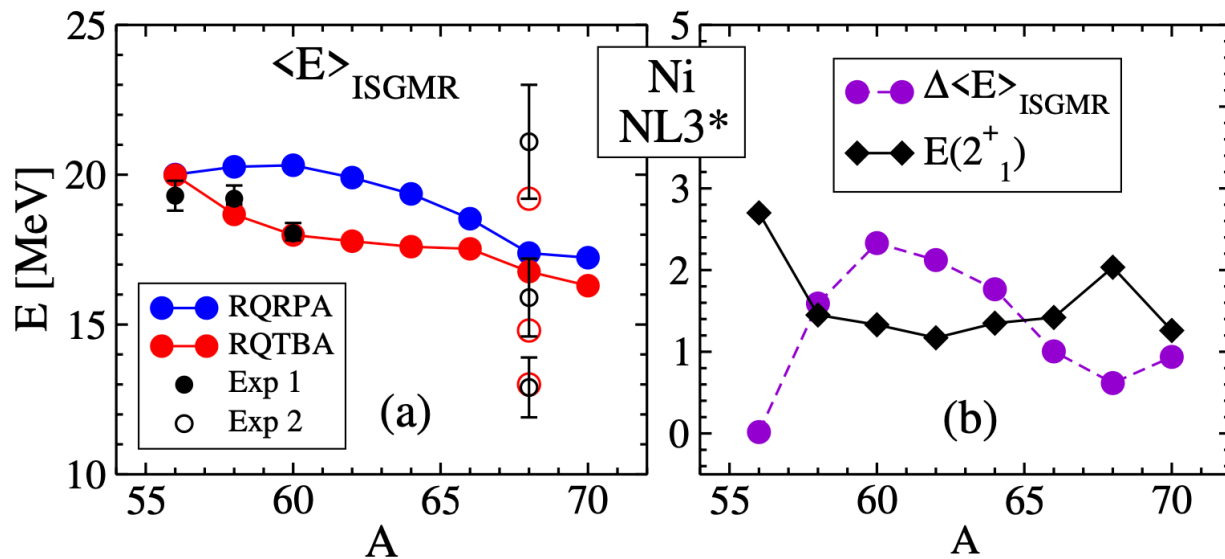
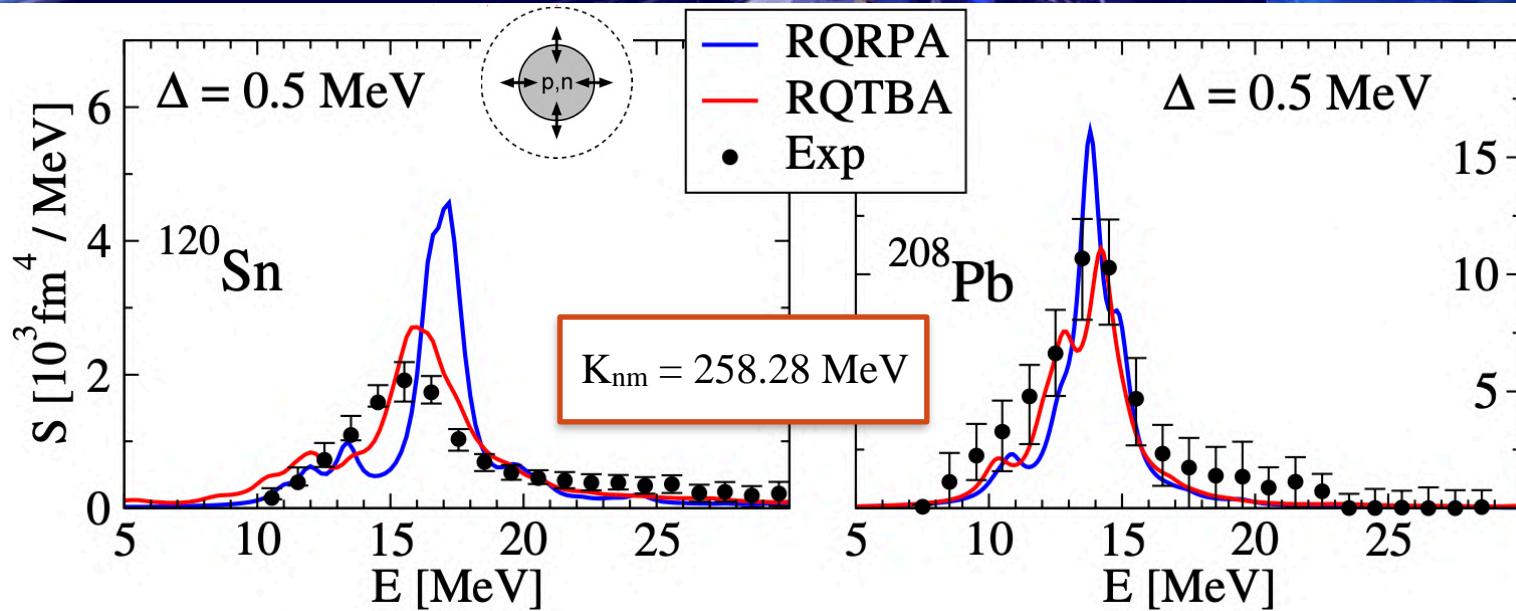
## Giant dipole resonance in $^{42,48}\text{Ca}$



## “Pygmy” dipole resonance in $^{68}\text{Ni}$



# Isoscalar giant monopole resonance (ISGMR): The “fluffiness” puzzle



**“Softness” increases:**

- with the neutron number
- with superfluidity
- with correlations beyond QRPA (q)PVC

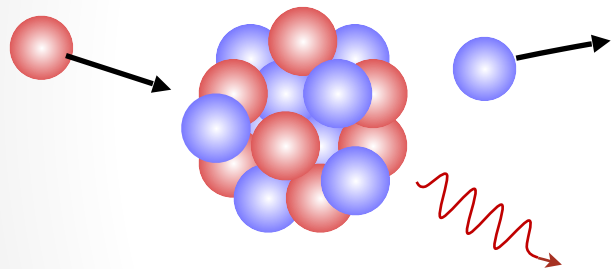
**Stiffer EOS can be used**

*Cf.: Skyrme HFB+qPVC*

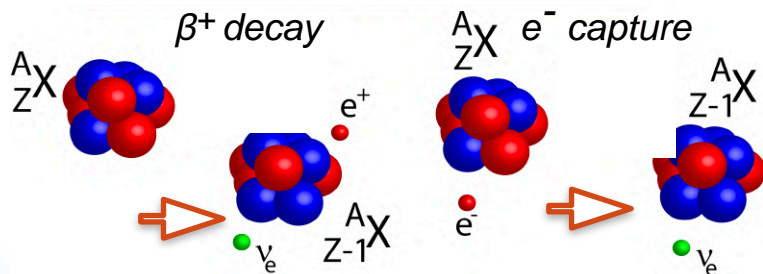
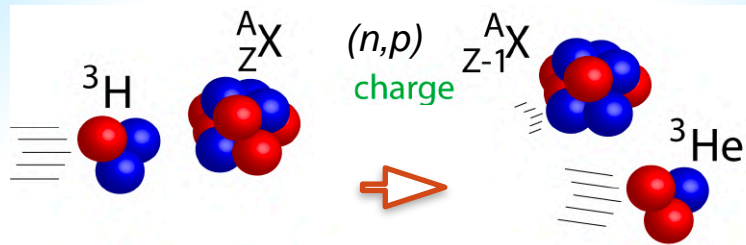
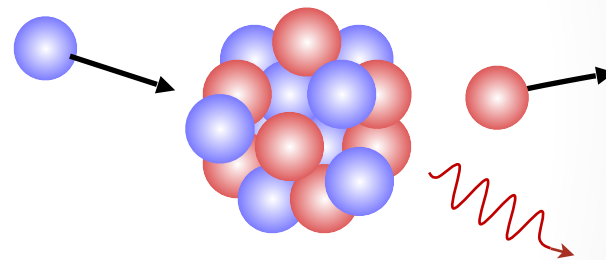
**G. Colò & Y. Niu**

# Spin-isospin excitations

Charge-exchange (p,n) reaction:  $\beta^-$



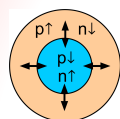
Charge-exchange (n,p) reaction:  $\beta^+$



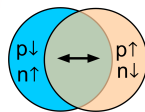
R. Zegers et al.

© NSCL Charge-exchange group:  
[https://groups.nsl.msui.edu/charge\\_exchange/resap.html](https://groups.nsl.msui.edu/charge_exchange/resap.html)

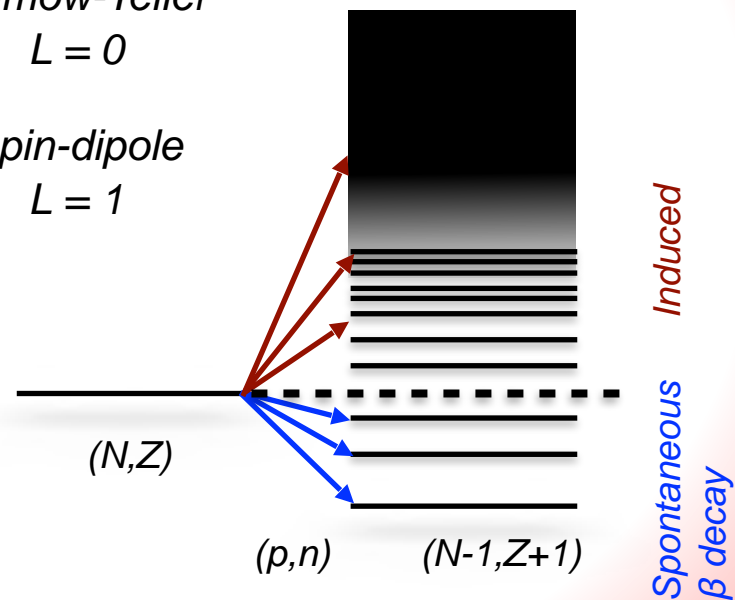
$$P_{\pm}^{\lambda} = \sum_i r_i [\sigma^{(i)} \otimes Y_L(\hat{r}_i)]_{\lambda} \tau_{\pm}^{(i)}$$



Gamow-Teller  
 $L = 0$



Spin-dipole  
 $L = 1$



# Exotic modes: the Isovector Giant Monopole Resonance

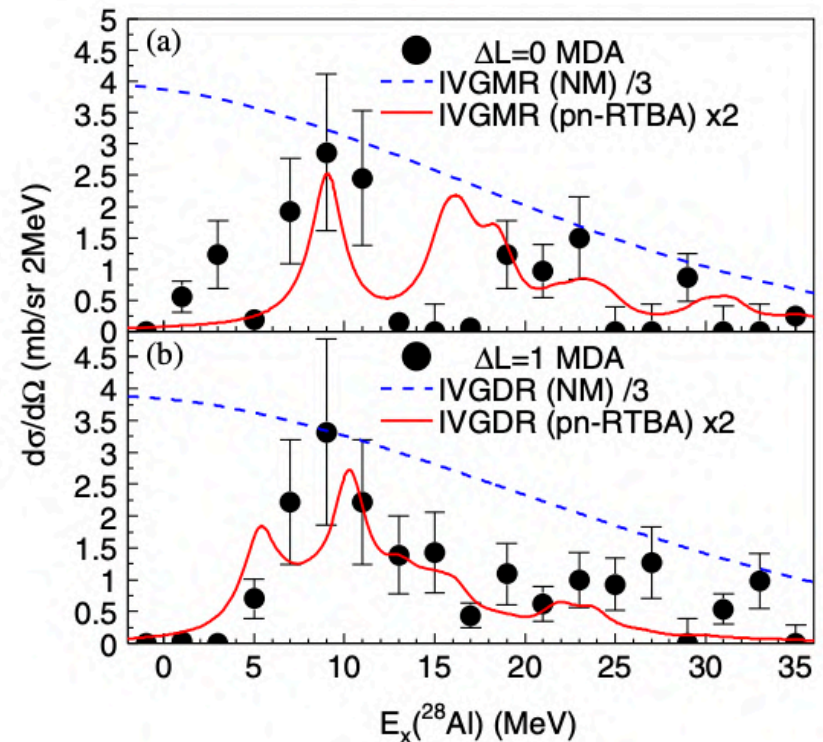
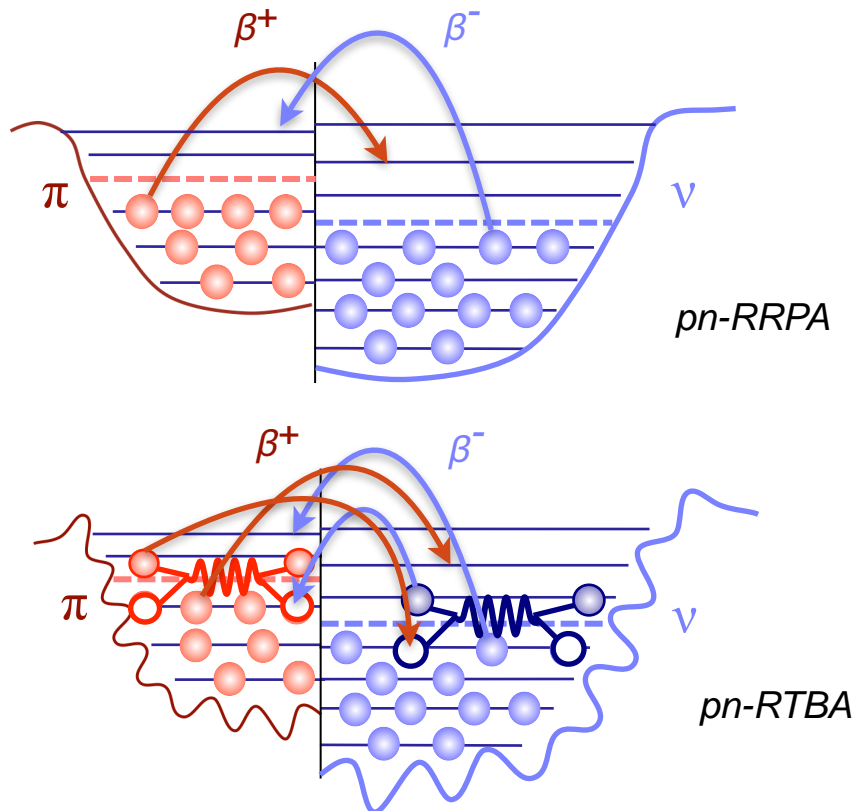
PRL 118, 172501 (2017)

PHYSICAL REVIEW LETTERS

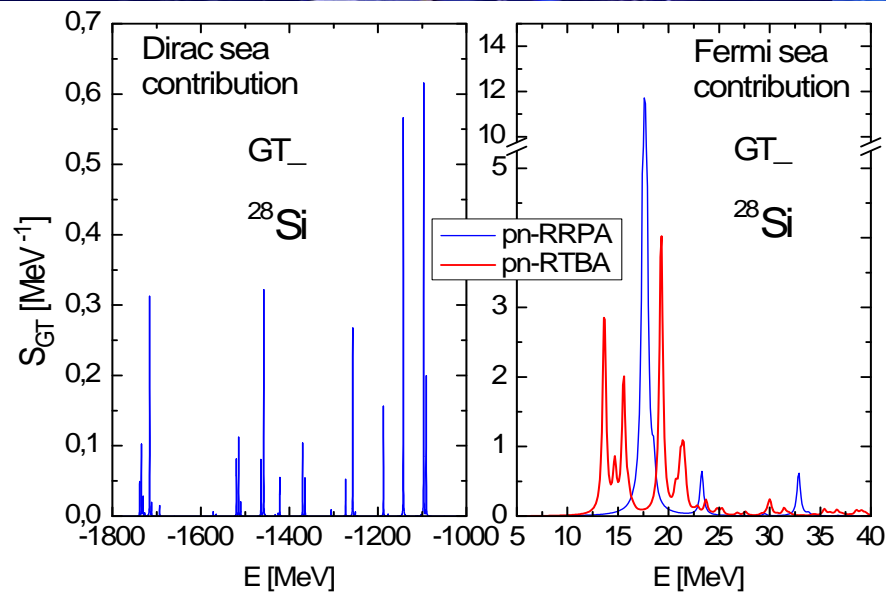
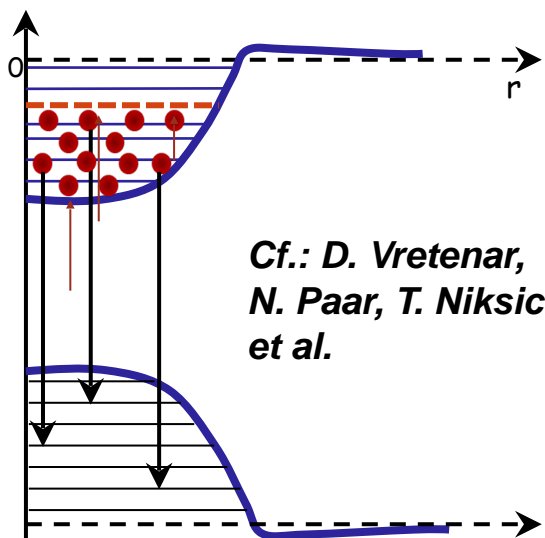
week ending  
28 APRIL 2017

## Observation of the Isovector Giant Monopole Resonance via the $^{28}\text{Si}(^{10}\text{Be}, ^{10}\text{B}^* [1.74 \text{ MeV}])$ Reaction at 100 AMeV

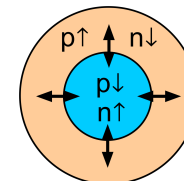
M. Scott,<sup>1,2,3</sup> R. G. T. Zegers,<sup>1,2,3,\*</sup> R. Almus,<sup>4</sup> Sam M. Austin,<sup>1,2,3</sup> D. Bazin,<sup>1</sup> B. A. Brown,<sup>1,2,3</sup> C. Campbell,<sup>5</sup> A. Gade,<sup>1,3</sup> M. Bowry,<sup>1</sup> S. Galès,<sup>6,7</sup> U. Garg,<sup>8</sup> M. N. Harakeh,<sup>9</sup> E. Kwan,<sup>1</sup> C. Langer,<sup>1,2</sup> C. Loelius,<sup>1,2,3</sup> S. Lipschutz,<sup>1,2,3</sup> E. Litvinova,<sup>10,1,2</sup> E. Lunderberg,<sup>1,3</sup> C. Morse,<sup>1,3</sup> S. Noji,<sup>1,2</sup> G. Perdikakis,<sup>4,1,2</sup> T. Redpath,<sup>4</sup> C. Robin,<sup>10,2</sup> H. Sakai,<sup>11</sup> Y. Sasamoto,<sup>12,11</sup> M. Sasano,<sup>11</sup> C. Sullivan,<sup>1,2,3</sup> J. A. Tostevin,<sup>13,1</sup> T. Uesaka,<sup>11</sup> and D. Weisshaar<sup>1</sup>



# The Gamow-Teller resonance quenching puzzle: $^{28}\text{Si}$



$$P = \sum_i \sigma^{(i)} \tau_{\pm}^{(i)}$$



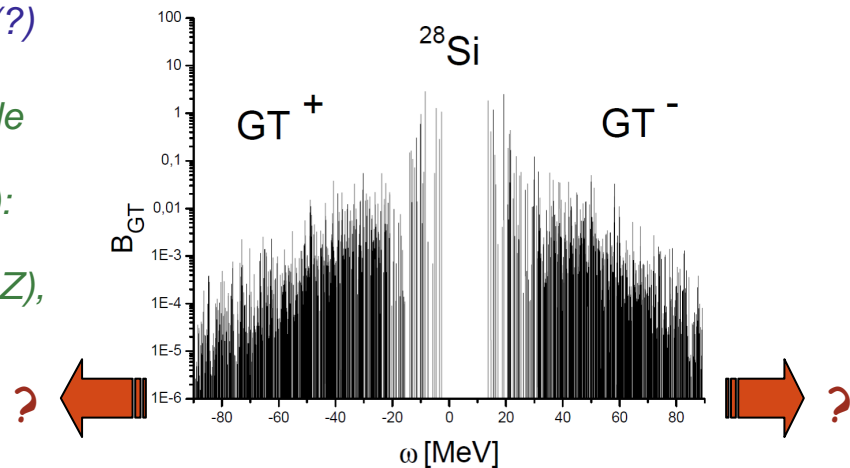
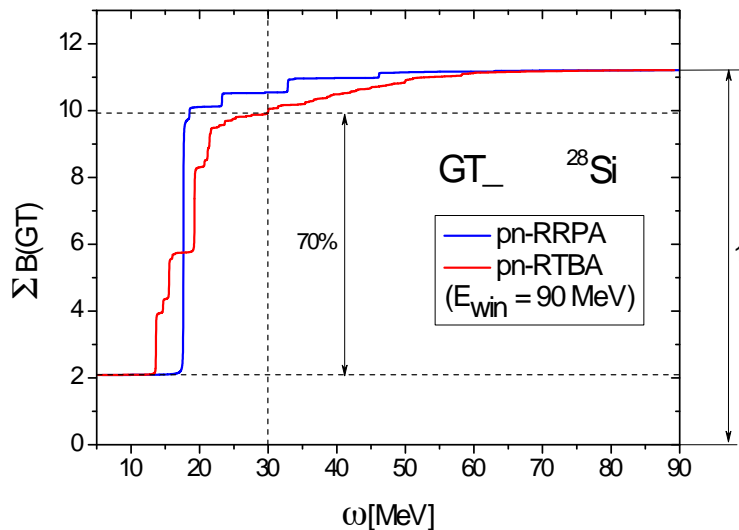
$$\begin{aligned} \Delta L &= 0 \\ \Delta T &= 1 \\ \Delta S &= 1 \end{aligned}$$

„Microscopic“ quenching of  $B(GT)$ :  
 (i) relativistic effects,  
 (ii)  $ph$ +phonon and higher configurations  
 (iii)  $2b$  currents (?)

Cf.: H. Sagawa, Y. Niu, D. Gambacurta

Ikeda Sum rule  
 (model independent):

$$S_- - S_+ = 3(N - Z),$$

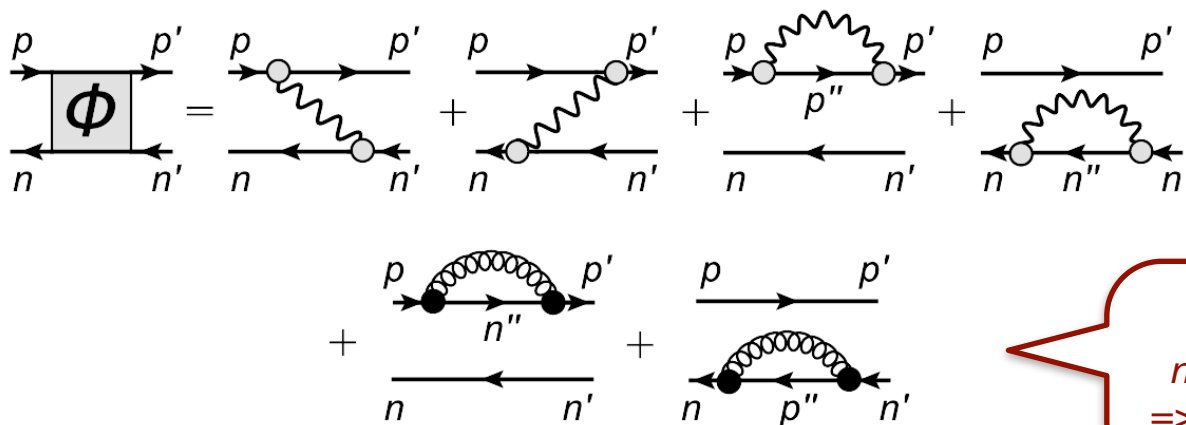


Problem: finite basis



# More on quenching: Coupling to charge-exchange (CE) phonons

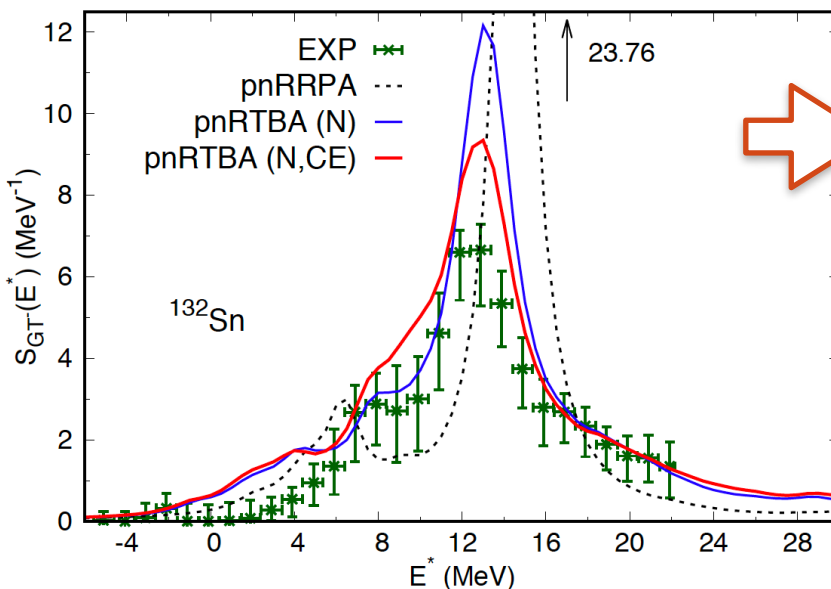
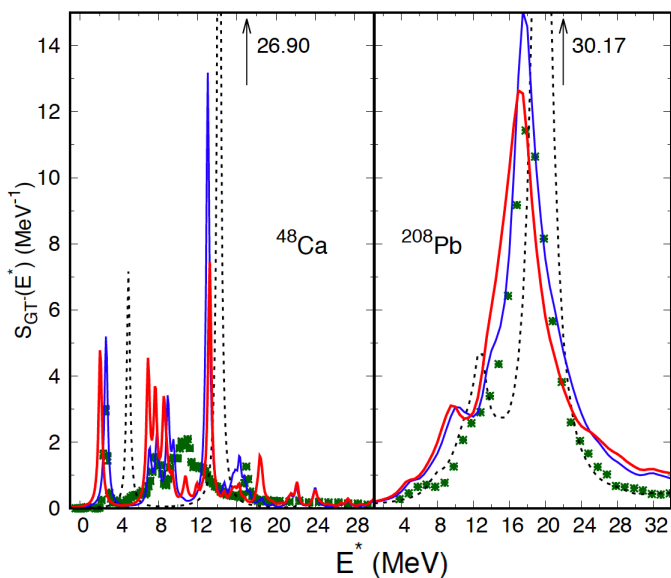
## The role of coupling to charge-exchange (CE) vibrations



Neutral phonons:  
Both phonon-exchange  
and self-energy  
contributions

Charge-exchange phonons:  
no phonon-exchange counterparts  
=> larger than expected contribution!

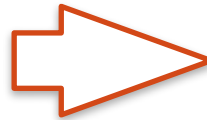
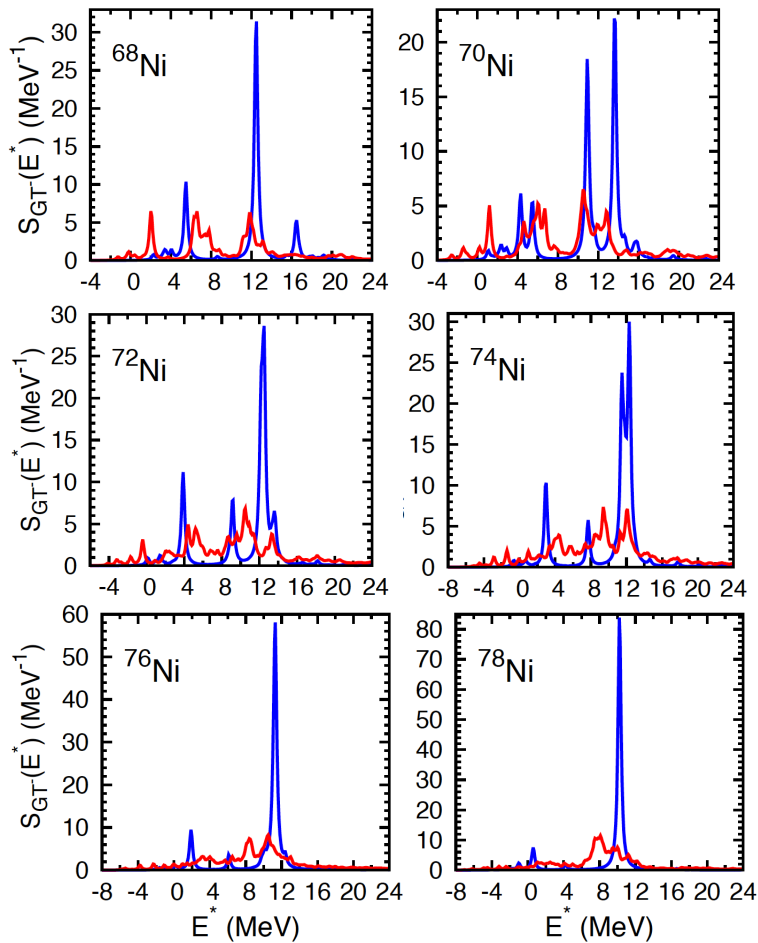
## Gamow-Teller response



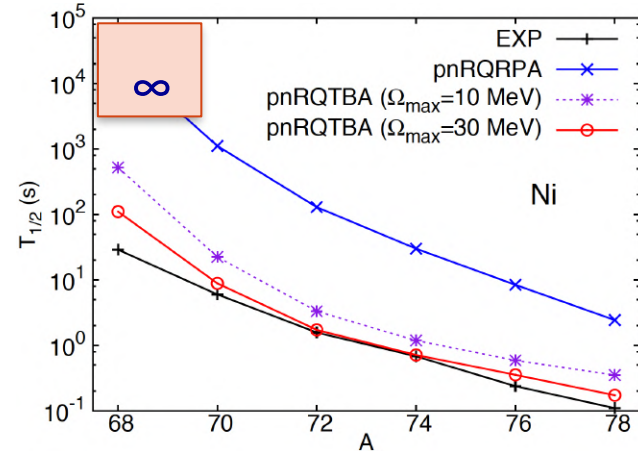
3p3h-  
configurations  
are needed  
for a consistent  
description  
without applying  
quenching  
factors  
(in progress)

# The beta decay puzzle: Gamow-Teller resonance and $T_{1/2}$

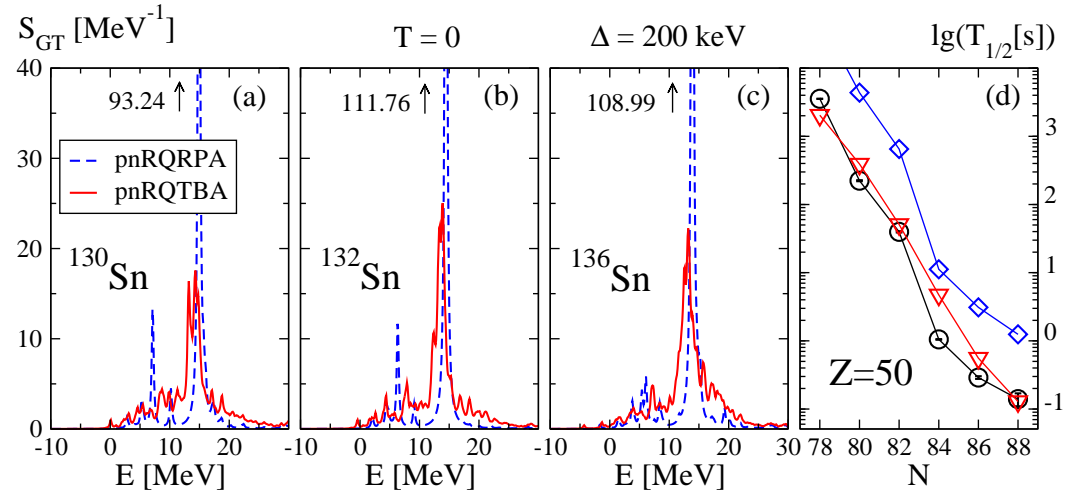
Overall strength  
68-78Ni isotopes:



Beta decay  $T_{1/2}$



128-138Sn isotopes:



—  $pn$ -RQRPA  
—  $pn$ -RQTBA

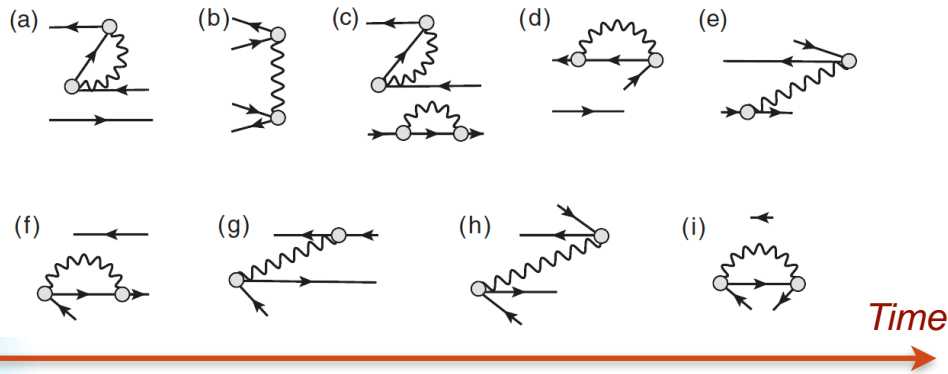
C. Robin, E.L., *Eur. Phys. J. A* 52, 205 (2016)

E.L., C. Robin, H. Wibowo, *PLB* 800, 135134 (2020)

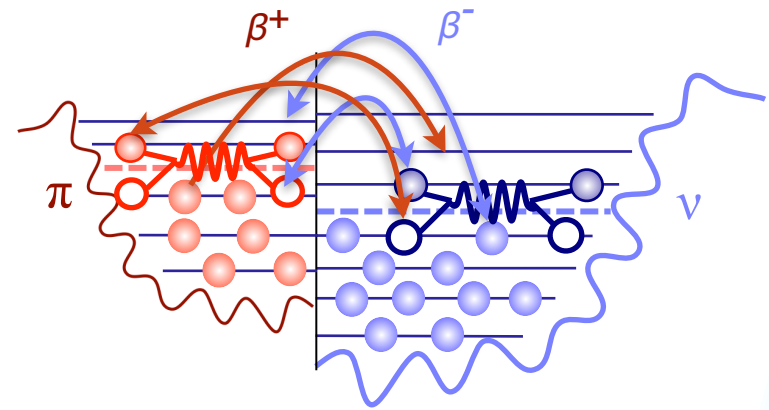
*No fits, no artificial quenching,  
no adjustable proton-neutron pairing*

# $\beta^+$ strength in neutron-rich nuclei: absent in the leading qPVC approach

Ground state correlations induced by qPVC:  
backward-going diagrams



New unblocking mechanism:

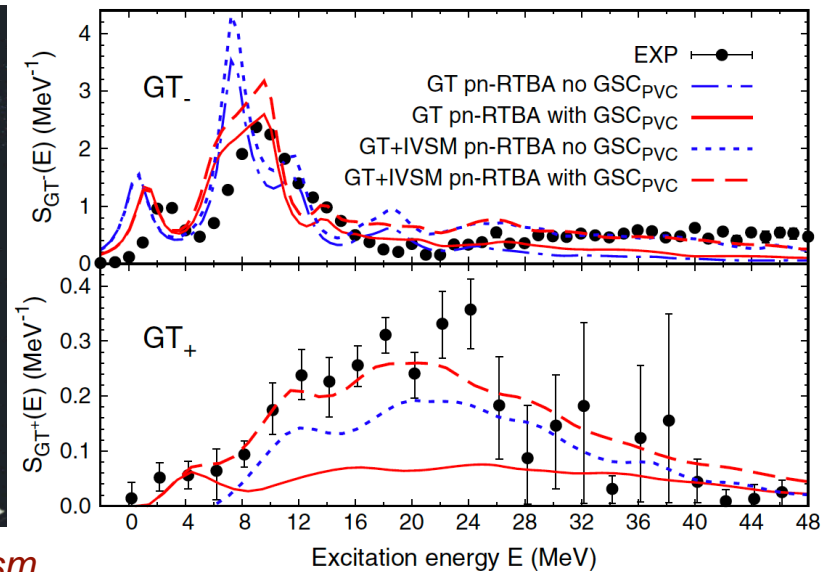


Emergent “time machine”



The backward-going diagrams are the leading mechanism of the  $\beta^+$  strength formation in neutron-rich nuclei

Gamow-Teller +IVSM strength in  $^{90}\text{Zr}$ :



C. Robin, E.L., *Phys. Rev. Lett.* 123, 202501 (2019)

# The "Astro" puzzle: Finite-temperature response with the ph+phonon dynamical kernel

$$R_{12,1'2'}(t-t') = -i\langle \mathcal{T}(\bar{\psi}_1\psi_2)(t)(\bar{\psi}_{2'}\psi_{1'})(t') \rangle \rightarrow -i\langle \mathcal{T}(\bar{\psi}_1\psi_2)(t)(\bar{\psi}_{2'}\psi_{1'})(t') \rangle_T$$

$$\langle \dots \rangle \equiv \langle 0|\dots|0 \rangle \rightarrow \langle \dots \rangle_T \equiv \sum_n \exp\left(\frac{\Omega - E_n - \mu N}{T}\right) \langle n|\dots|n \rangle$$

averages



thermal averages

$$\begin{aligned} \mathcal{R}_{14,23}(\omega, T) &= \tilde{\mathcal{R}}_{14,23}^0(\omega, T) + \\ &+ \sum_{1'2'3'4'} \tilde{\mathcal{R}}_{12',21'}^0(\omega, T) [\tilde{V}_{1'4',2'3'}(T) + \delta\Phi_{1'4',2'3'}(\omega, T)] \mathcal{R}_{3'4,4'3}(\omega, T) \\ \delta\Phi_{1'4',2'3'}(\omega, T) &= \Phi_{1'4',2'3'}(\omega, T) - \Phi_{1'4',2'3'}(0, T) \end{aligned}$$

**Method: EOM  
for Matsubara  
Green's functions**

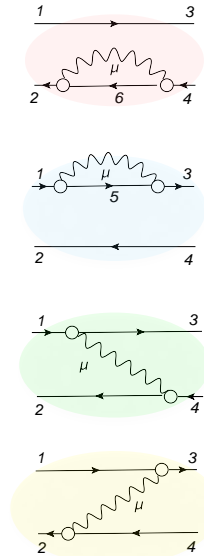


See talk of H. Wibowo

$T > 0$ :

$$\begin{aligned} \Phi_{14,23}^{(ph)}(\omega, T) &= \frac{1}{n_{43}(T)} \sum_{\mu; \eta_\mu = \pm 1} \eta_\mu \left[ \delta_{13} \sum_6 \gamma_{\mu;62}^{\eta_\mu} \gamma_{\mu;64}^{\eta_\mu*} \times \right. \\ &\times \frac{(N(\eta_\mu\Omega_\mu) + n_6(T))(n(\varepsilon_6 - \eta_\mu\Omega_\mu, T) - n_1(T))}{\omega - \varepsilon_1 + \varepsilon_6 - \eta_\mu\Omega_\mu} + \\ &+ \delta_{24} \sum_5 \gamma_{\mu;15}^{\eta_\mu} \gamma_{\mu;35}^{\eta_\mu*} \times \\ &\times \frac{(N(\eta_\mu\Omega_\mu) + n_2(T))(n(\varepsilon_2 - \eta_\mu\Omega_\mu, T) - n_5(T))}{\omega - \varepsilon_5 + \varepsilon_2 - \eta_\mu\Omega_\mu} - \\ &- \gamma_{\mu;13}^{\eta_\mu} \gamma_{\mu;24}^{\eta_\mu*} \times \\ &\times \frac{(N(\eta_\mu\Omega_\mu) + n_2(T))(n(\varepsilon_2 - \eta_\mu\Omega_\mu, T) - n_3(T))}{\omega - \varepsilon_3 + \varepsilon_2 - \eta_\mu\Omega_\mu} - \\ &- \gamma_{\mu;31}^{\eta_\mu*} \gamma_{\mu;42}^{\eta_\mu} \times \\ &\left. \times \frac{(N(\eta_\mu\Omega_\mu) + n_4(T))(n(\varepsilon_4 - \eta_\mu\Omega_\mu, T) - n_1(T))}{\omega - \varepsilon_1 + \varepsilon_4 - \eta_\mu\Omega_\mu} \right], \end{aligned}$$

1p1h+phonon dynamical kernel:



$T = 0$ :

$$\begin{aligned} \Phi_{14,23}^{(ph,ph)}(\omega) &= \sum_{\mu} \times \\ &\times \left[ \delta_{13} \sum_6 \frac{\gamma_{62}^{\mu} \gamma_{64}^{\mu*}}{\omega - \varepsilon_1 + \varepsilon_6 - \Omega_{\mu}} + \right. \\ &+ \delta_{24} \sum_5 \frac{\gamma_{15}^{\mu} \gamma_{35}^{\mu*}}{\omega - \varepsilon_5 + \varepsilon_2 - \Omega_{\mu}} - \\ &- \frac{\gamma_{13}^{\mu} \gamma_{24}^{\mu*}}{\omega - \varepsilon_3 + \varepsilon_2 - \Omega_{\mu}} - \\ &\left. - \frac{\gamma_{31}^{\mu*} \gamma_{42}^{\mu}}{\omega - \varepsilon_1 + \varepsilon_4 - \Omega_{\mu}} \right] \end{aligned}$$

# The “upbend” puzzle: understanding the Oslo data

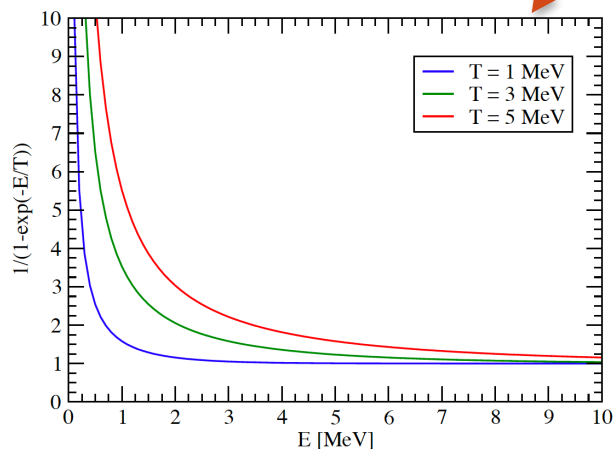
$$S(E, T) = -\frac{1}{\pi} \lim_{\Delta \rightarrow +0} \text{Im} \langle V^{0\dagger} \mathcal{R}(E + i\Delta, T) V^0 \rangle$$

The final strength function at  $T > 0$ :

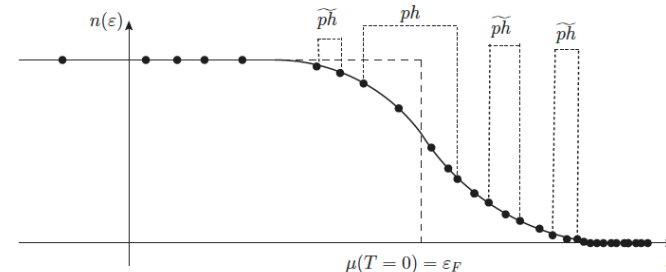
$$\tilde{S}(E) = \frac{1}{1 - e^{-E/T}} S(E)$$

$$\lim_{E \rightarrow 0} S(E, T) = 0$$

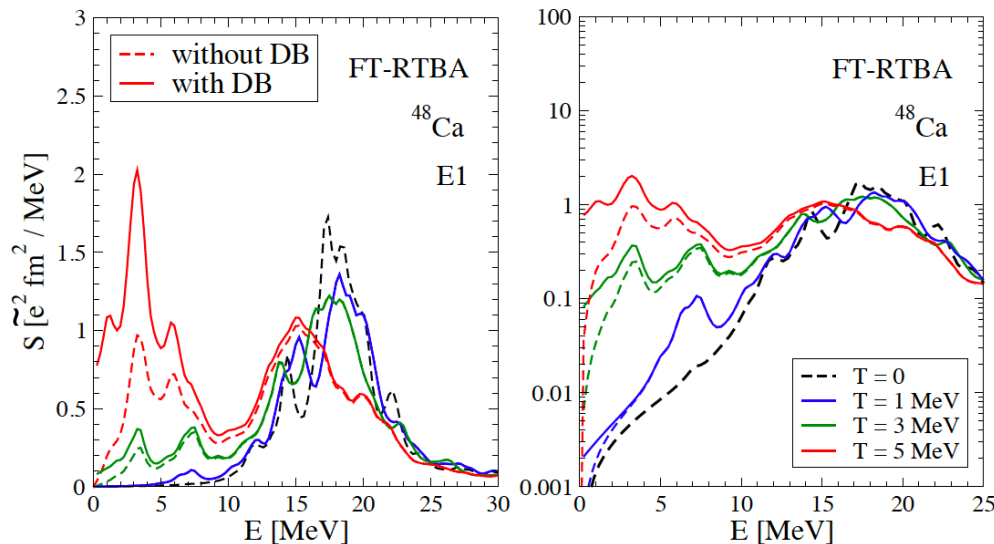
The **generic** exponential factor:



Thermal unblocking:



Dipole strength: absorption at  $T > 0$ :



• The exponential factor brings an additional enhancement in  $E < T$  energy region and provides the finite zero-energy limit of the strength (regardless its spin-parity)

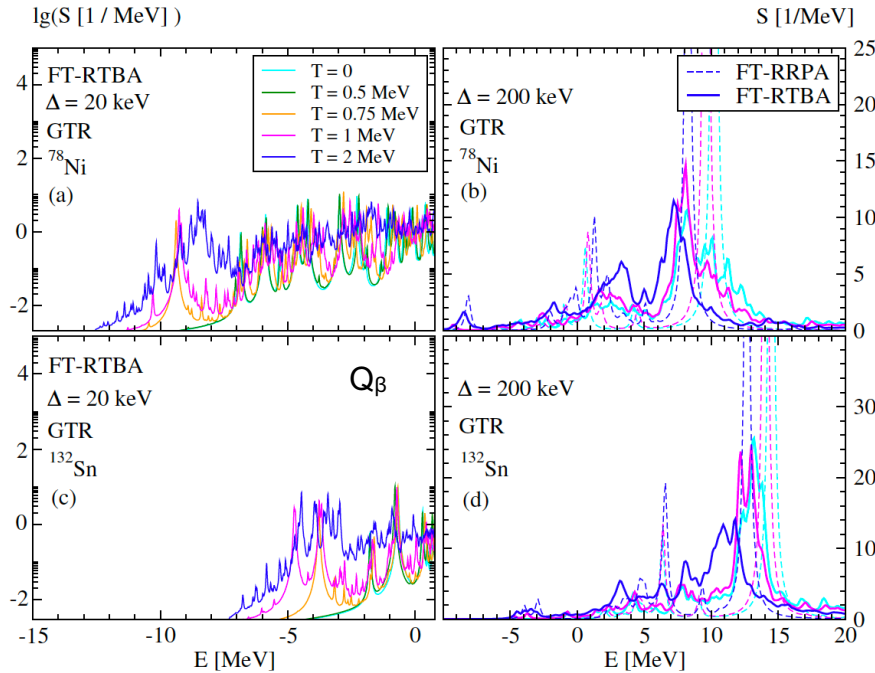
See talk of H. Wibowo

E.L., H. Wibowo, *Phys. Rev. Lett.* 121, 082501 (2018)

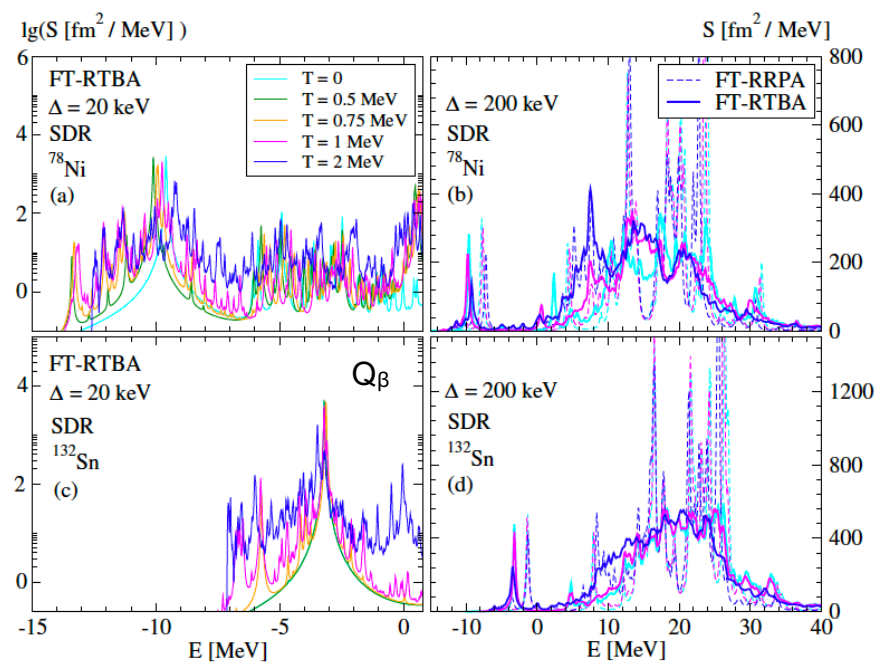
H. Wibowo, E.L., *Phys. Rev. C* 100, 024307 (2019)

# Spin-Isospin response and beta decay in hot stellar environments

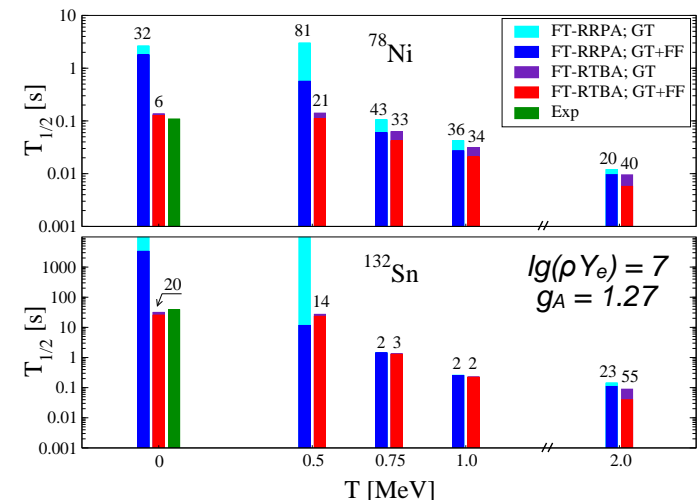
## Gamow-Teller $GT_-$ response of $^{78}\text{Ni}$ and $^{132}\text{Sn}$



## Spin Dipole response of $^{78}\text{Ni}$ and $^{132}\text{Sn}$



## Beta decay half-lives in a hot stellar environment



**An advanced microscopic description of the beta decay beyond (Q)RPA is available for the r-process modeling**

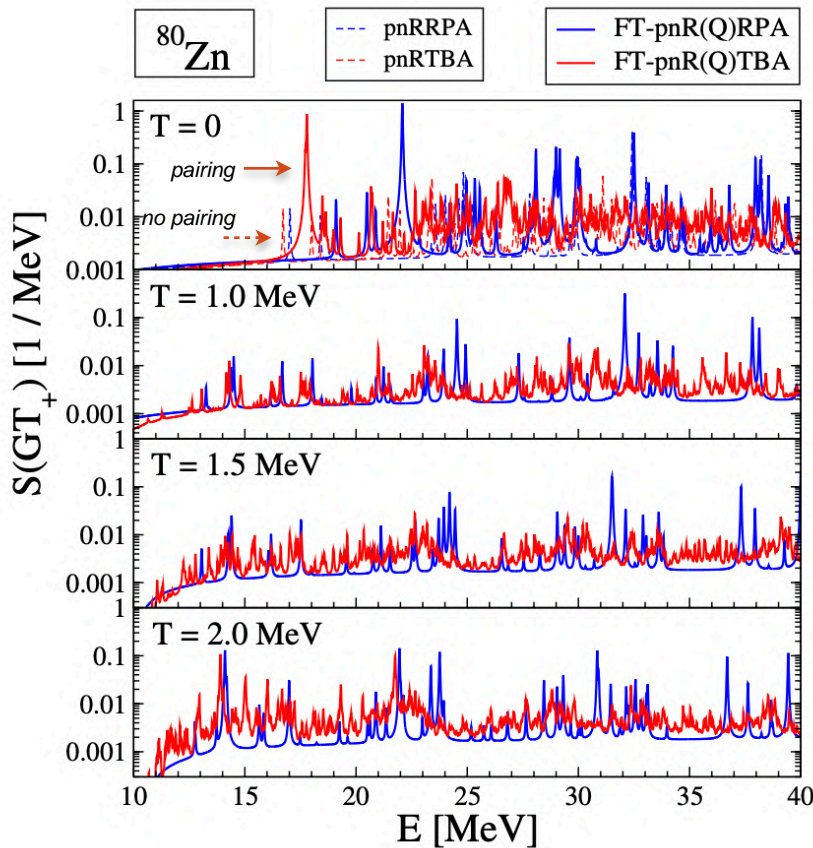
E. Litvinova, H. Wibowo, *Phys. Rev. Lett.* 121, 082501 (2018)

H. Wibowo, E. Litvinova, *Phys. Rev. C* 100, 024307 (2019)

E. Litvinova, C. Robin, H. Wibowo, *Phys. Lett. B* 800, 135134 (2020)

# GT+ response and electron capture (EC) rates at $T > 0$ : the neighborhood of $^{78}\text{Ni}$

## GT+ response



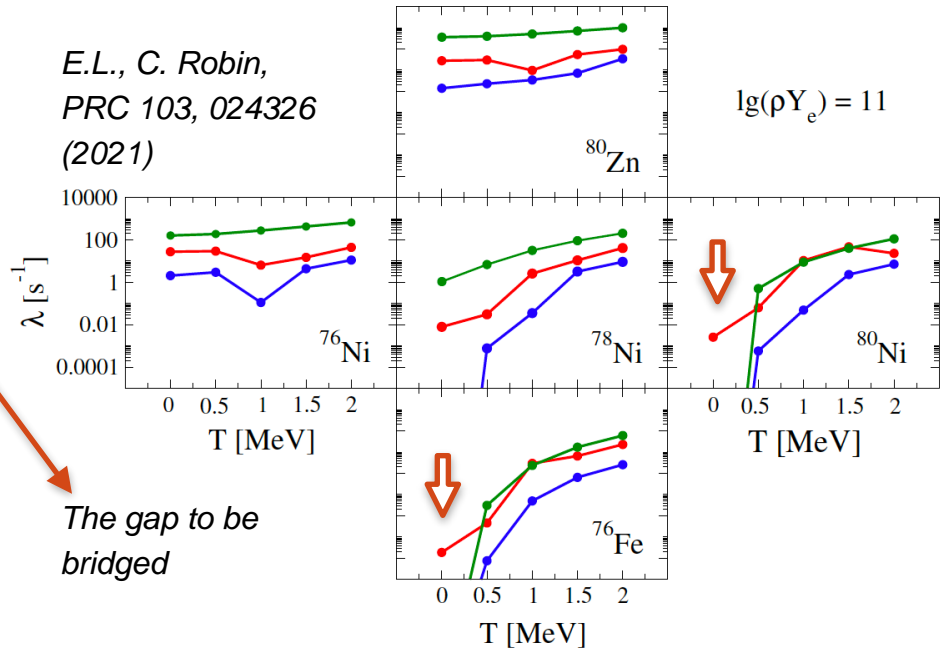
Superfluid

$T_c$

Non-superfluid

## Electron capture rates around $^{78}\text{Ni}$

E.L., C. Robin,  
PRC 103, 024326  
(2021)



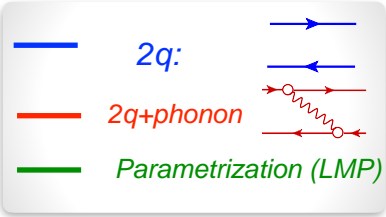
The gap to be bridged

**Interplay** of superfluidity and collective effects in the core-collapse supernovae:

- Amplifies the EC rates and, consequently,
- Reduces the electron-to-baryon ratio leading to lower pressure
- Promotes the gravitational collapse
- Increases the neutrino flux and effective cooling
- Allows heavy nuclei to survive the collapse

**An advanced microscopic description of the EC in the supernova is available**

beyond (R)QRPA [E. Yüksel, N. Paar, Y. Niu, R. Zegers et al.]



# Towards complete formalism at $T>0$ : the pairing channel

Averages redefined:

$$G_{12,1'2'}(t-t') = -i\langle \mathcal{T}(\psi_1\psi_2)(t)(\bar{\psi}_{2'}\bar{\psi}_{1'})(t') \rangle \rightarrow -i\langle \mathcal{T}(\psi_1\psi_2)(t)(\bar{\psi}_{2'}\bar{\psi}_{1'})(t') \rangle_T$$

**Grand Canonical average:**  $\langle \dots \rangle \equiv \langle 0|\dots|0 \rangle \rightarrow \langle \dots \rangle_T \equiv \sum_n \exp\left(\frac{\Omega - E_n - \mu N}{T}\right) \langle n|\dots|n \rangle$

Matsubara imaginary-time formalism: temperature-dependent dynamical kernel

Direct:

$$\begin{aligned} \mathcal{K}_{121'2'}^{(r;11)}(\omega_n) &= - \sum_{\nu'\nu''} w_{\nu'} w_{\nu''} \\ &\times \left[ \sum_{\nu\mu} \frac{\Theta_{121'2'}^{\mu\nu;\nu'\nu''(+)} }{i\omega_n - \omega_{\nu\nu'} - \omega_{\mu\nu''}^{(++)}} (e^{-(\omega_{\nu\nu'} + \omega_{\mu\nu''}^{(++)})/T} - 1) \right. \\ &\left. - \sum_{\nu\kappa} \frac{\Theta_{121'2'}^{\kappa\nu;\nu'\nu''(-)} }{i\omega_n + \omega_{\nu\nu'} + \omega_{\kappa\nu''}^{(--)}} (e^{-(\omega_{\nu\nu'} + \omega_{\kappa\nu''}^{(--)})/T} - 1) \right] \end{aligned}$$

Exchange:

$$\begin{aligned} \mathcal{K}_{121'2'}^{(r;12)}(\omega_n) &= \sum_{\nu'\nu''} w_{\nu'} w_{\nu''} \\ &\times \left[ \sum_{\nu\mu} \frac{\Sigma_{121'2'}^{\mu\nu;\nu'\nu''(+)} }{i\omega_n - \omega_{\nu\nu'} - \omega_{\mu\nu''}^{(++)}} (e^{-(\omega_{\nu\nu'} + \omega_{\mu\nu''}^{(++)})/T} - 1) \right. \\ &\left. - \sum_{\nu\kappa} \frac{\Sigma_{121'2'}^{\kappa\nu;\nu'\nu''(-)} }{i\omega_n + \omega_{\nu\nu'} + \omega_{\kappa\nu''}^{(--)}} (e^{-(\omega_{\nu\nu'} + \omega_{\kappa\nu''}^{(--)})/T} - 1) \right], \end{aligned}$$

E.L., P.Schuck, Phys. Rev. C 104, 044330 (2021)

BCS-like gap Eq., but with non-trivial  $T$ -dependence in  $K^{(r)}$ :

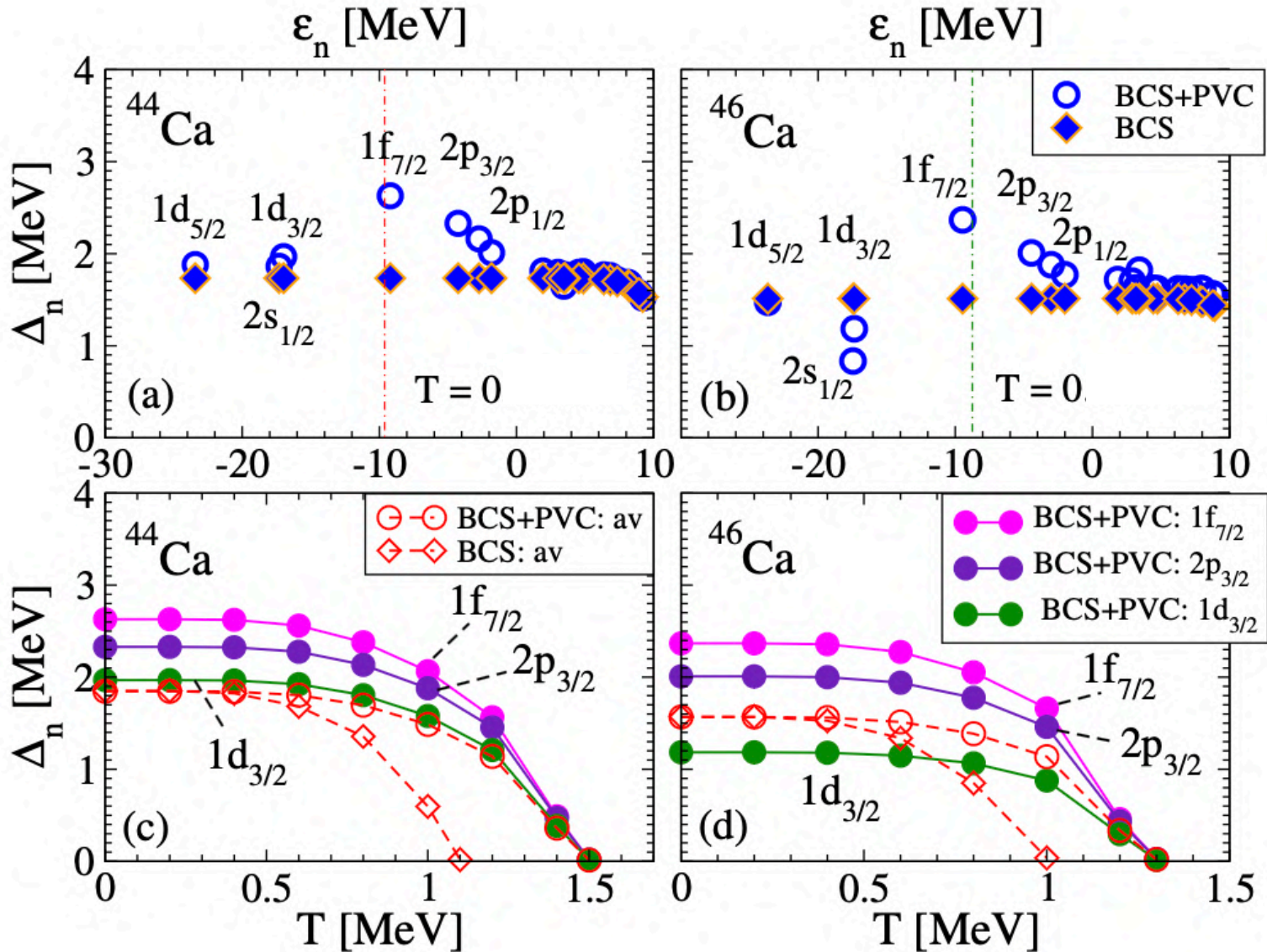
$$\Delta_1(T) = - \sum_2 \nu_{1\bar{1}2\bar{2}} \frac{\Delta_2(T)(1 - 2f_2(T))}{2E_2}$$

$$f_1(T) = \frac{1}{\exp(E_1/T) + 1}$$

$$\mathcal{V}_{121'2'} = \frac{1}{2} \left( K_{121'2'}^{(0)} + K_{121'2'}^{(r)}(2\lambda) \right)$$



Pairing gap at  $T = 0$ ,  $T > 0$  and critical temperature





# Outlook

## Summary:

- The relativistic nuclear field theory (RNFT) is formulated and advanced in the Equation of Motion (EOM) framework, with the emphasis on **emergent collectivity**, solving burning nuclear structure issues
- The **emergent collective effects** renormalize interactions in correlated media, underly the spectral fragmentation mechanisms, affect superfluidity and weak decay rates.
- Relativistic NFT is **generalized to finite temperature** and applied to nuclear superfluidity.
- Weak rates at astrophysical conditions are extracted: **the correlations beyond QRPA are found significant.**

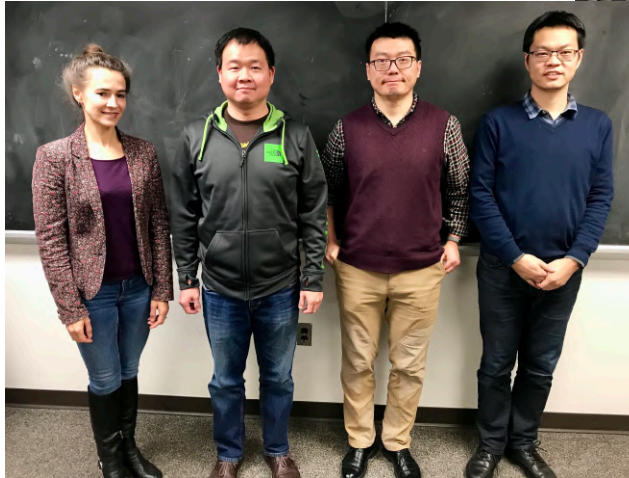
## Current and future developments:

- **Deformed nuclei:** correlations vs shapes; first results just released (Yinu Zhang et al.);
- Efficient algorithms; **quantum computing** (Manqoba Hlatshwayo et al.);
- Implementation of the EC rates into the **core-collapse supernovae simulations**;
- Toward an “*ab initio*” description: implementations with bare NN-interactions;
- **Superfluid pairing at  $T>0$**  to extend the application range (*r*-process);
- **Relativistic EOM’s, bosonic EOM’s, beyond Standard Model, ...**

# Many thanks for collaboration and support:

*Yinu Zhang (WMU)*  
*Manqoba Hlatshwayo (WMU)*  
*Herlik Wibowo (AS Taipei)*  
*Caroline Robin (U. Bielefeld & GSI)*  
*Peter Schuck (IPN Orsay)*  
*Peter Ring (TU München)*  
*Tamara Niksic (U Zagreb)*

2018:



2017:



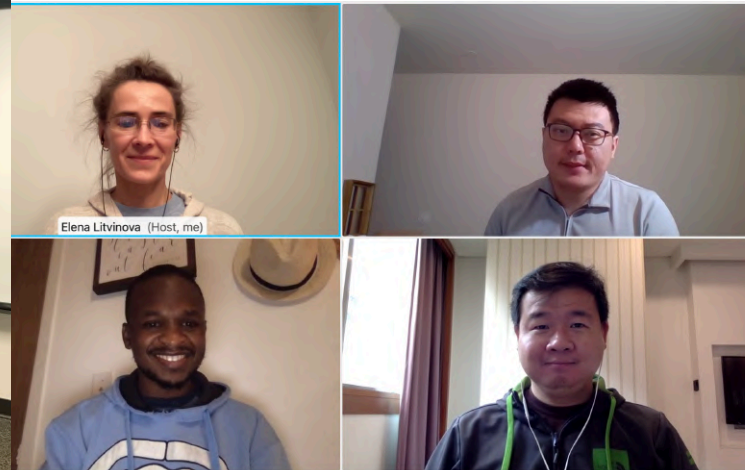
**US-NSF CAREER**  
**PHY-1654379 (2017-2023)**

**US-NSF PHY-2209376**  
**(2022-2025)**

2019-2020:



2020-2022:





*Thank you!*

