



LUND UNIVERSITY

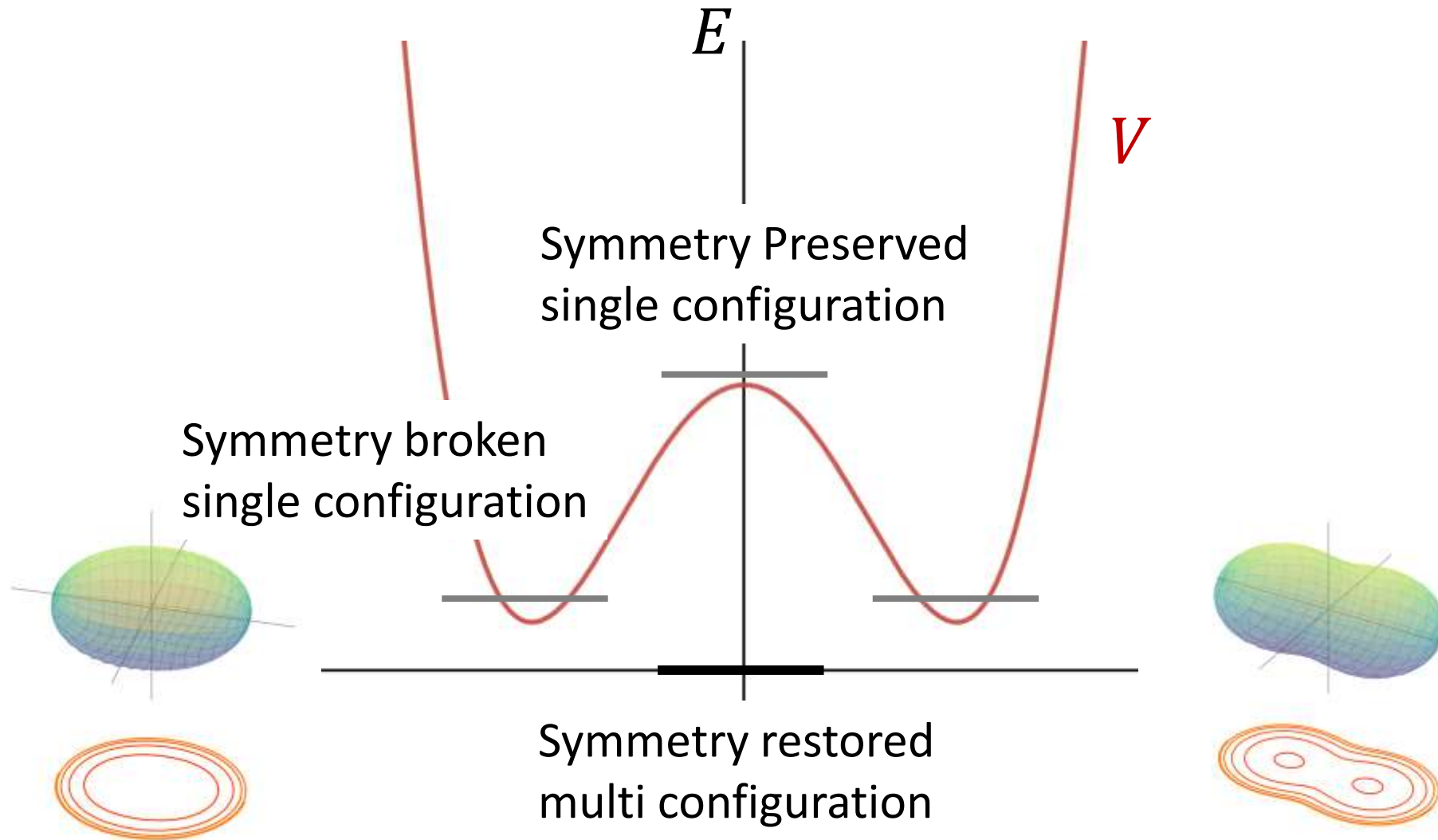
# Properties of rotational bands from symmetry breaking and restoration

**Andrea Idini**

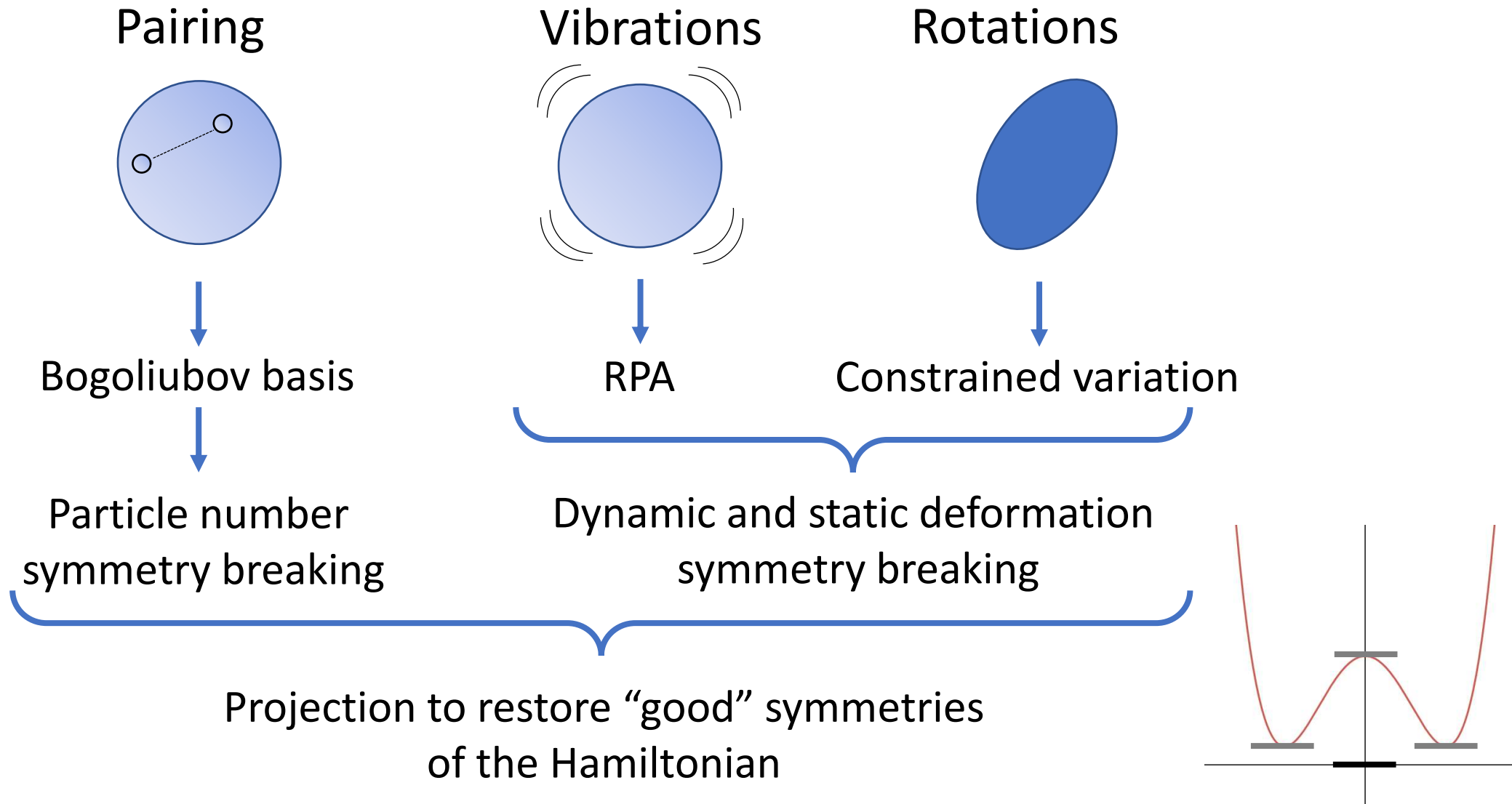
**COMEX7**

**Catania, 12 June 2023**

# Symmetry Breaking



# Nuclear Correlations



# Generator Coordinate Method

$$|\Psi\rangle = \int da f(a) |\Phi(a)\rangle$$

weight function  
generating states  
generator Coordinate

In our case

$$|\psi\rangle = \sum f(\cdot) |\Phi(\beta, \gamma, g_n, g_p, j_x)\rangle$$

Generating states as constrained HFB

Project to good number of particles and angular momentum

$$|\Psi\rangle = P^N P^Z P_{MK}^I (h_1 |\Phi(\text{circle})\rangle + h_2 |\Phi(\text{oval})\rangle + h_3 |\Phi(\text{oval})\rangle + \dots)$$

# Hill Wheeler Equation

$$H|\psi_i\rangle = E_i|\psi_i\rangle \quad \text{Schroedinger Equation}$$



However,  $|\psi_i\rangle$  are described in a non-orthogonal  $|\Phi(\cdot)\rangle$

$$Hh = EOh$$

Hill-Wheeler  $O_{IJ}^{KK'} = \langle \phi_I | \hat{P}_{KK'}^I \hat{P}_Z \hat{P}_N | \phi_J \rangle$

$$H_{IJ}^{KK'} = \langle \phi_I | \hat{H} \hat{P}_{KK'}^I \hat{P}_Z \hat{P}_N | \phi_J \rangle$$

$$|\Psi(I)\rangle = \sum_{KM} P^N P^Z P_{MK}^I (h_1 |\Phi(\text{light blue circle})\rangle + h_2 |\Phi(\text{dark blue oval})\rangle + h_3 |\Phi(\text{dark blue oval})\rangle + \dots)$$

# Construction of the effective Hamiltonian

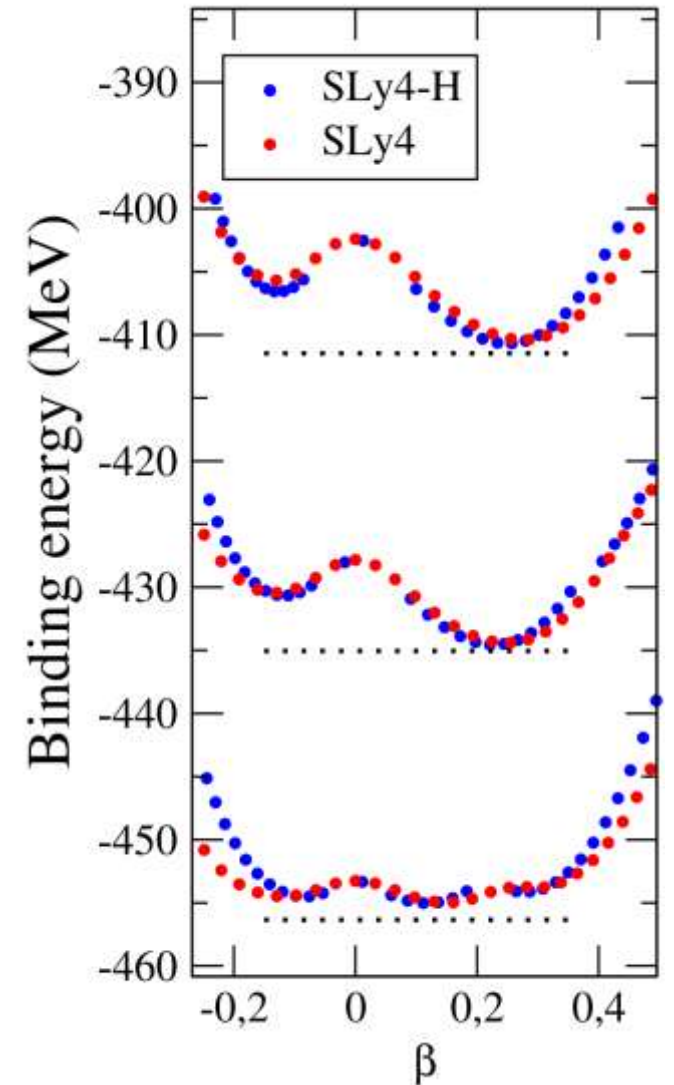
$$\hat{H} = \hat{H}_0 + \hat{H}_Q + \hat{H}_P$$

Single particle “mean field”

Quadrupole-Quadrupole

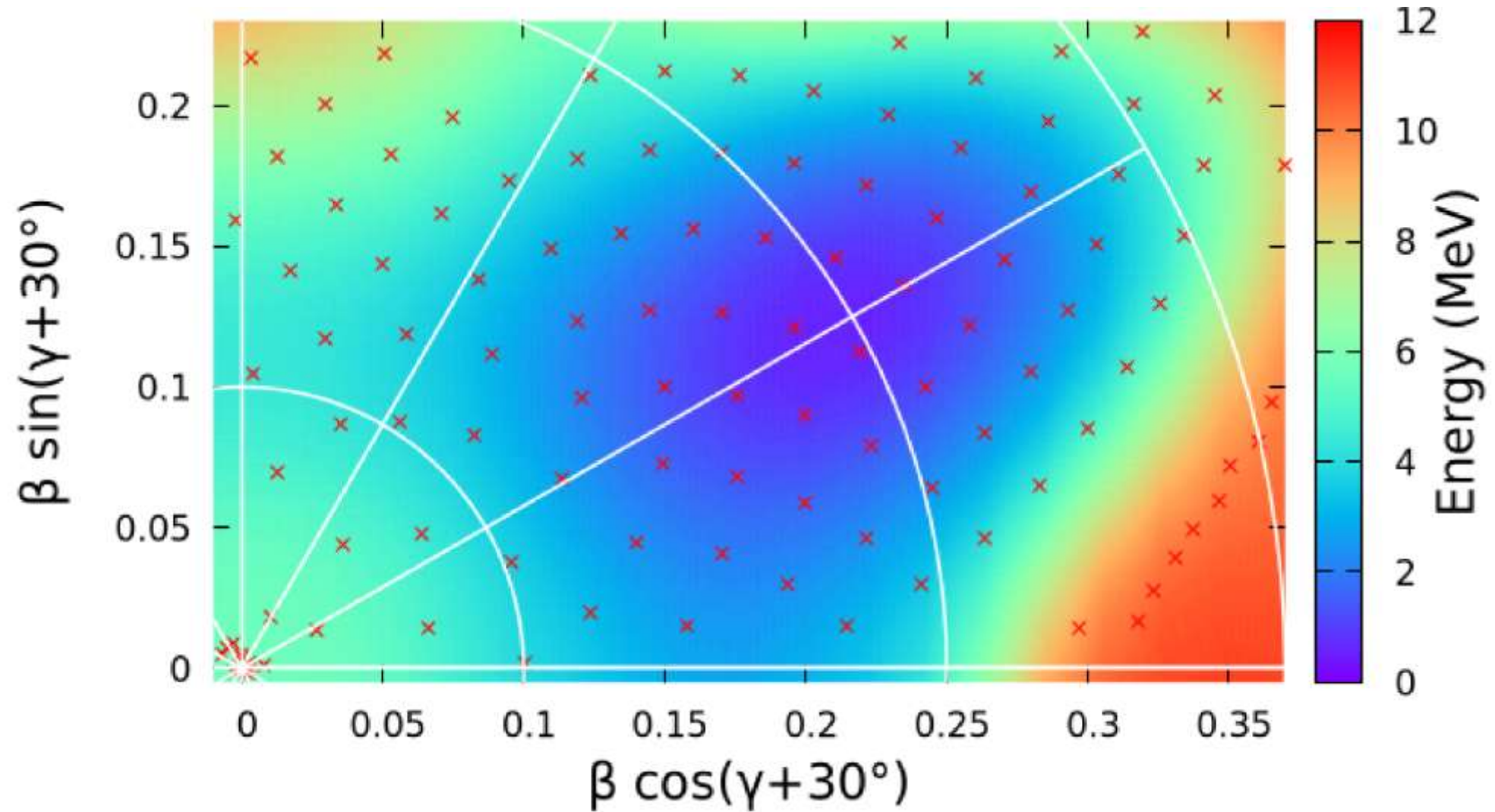
Seniority Pairing

$$H^{eff} = \sum_i \epsilon_i a_i^\dagger a_i - \frac{1}{4} \chi \sum_{\mu,ijkl} [Q_{ij}^{2\mu} Q_{kl}^{2\mu*} - Q_{ik}^{2\mu} Q_{jl}^{2\mu*}] a_i^\dagger a_j^\dagger a_k a_l + G \sum_{ijkl} P_{ij} P_{kl} a_i^\dagger a_j^\dagger a_k a_l$$



# Construction of the basis states

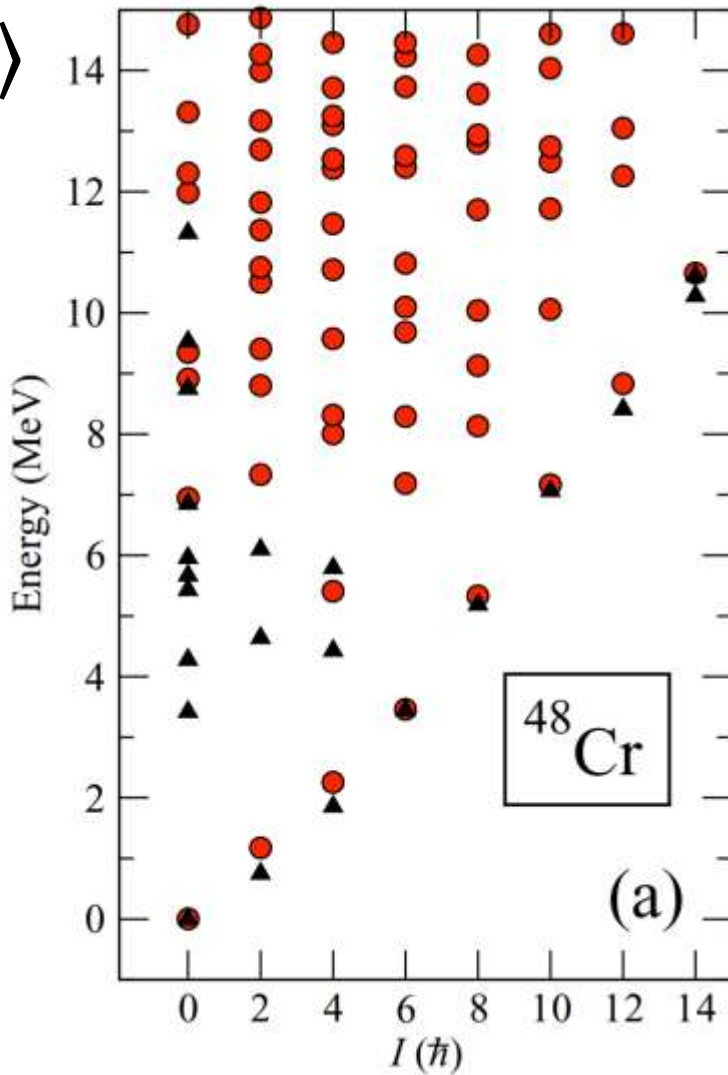
$^{48}\text{Cr}$  SLy4-H,  $N_{\text{max}}=11$



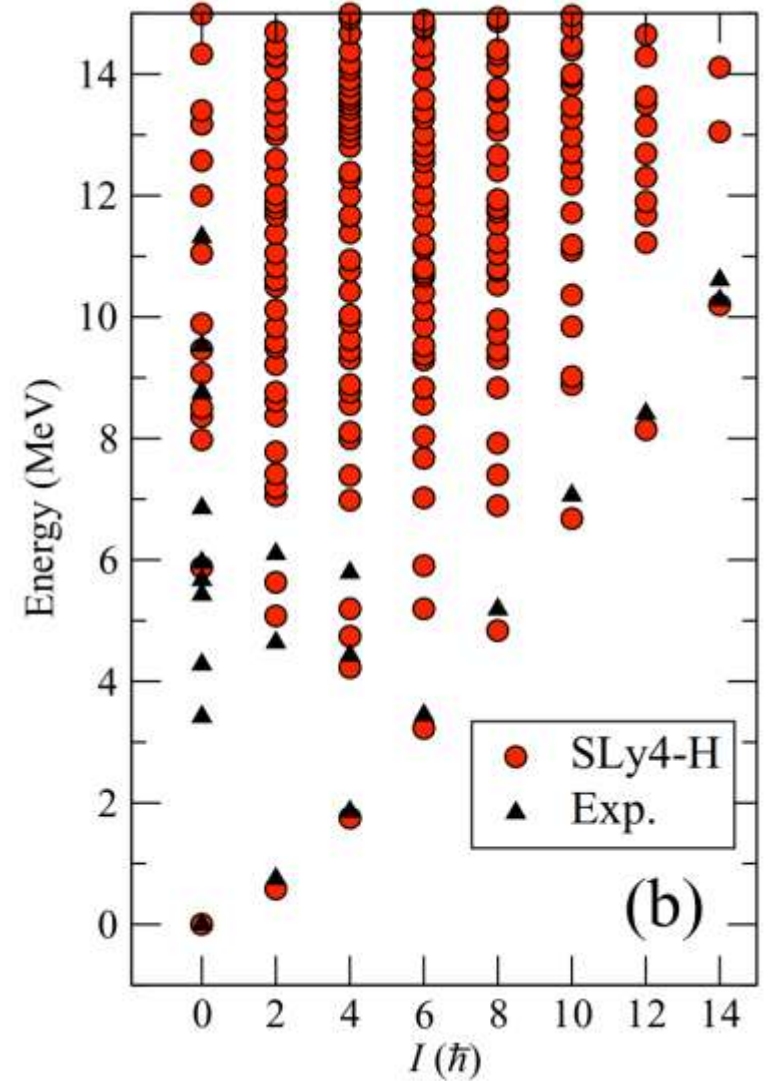


$$|\psi'(\cdot)\rangle = e^{\sum z \beta^+ \beta^-} |\psi(\cdot)\rangle$$

$z$  “temperature like”  
random coefficient

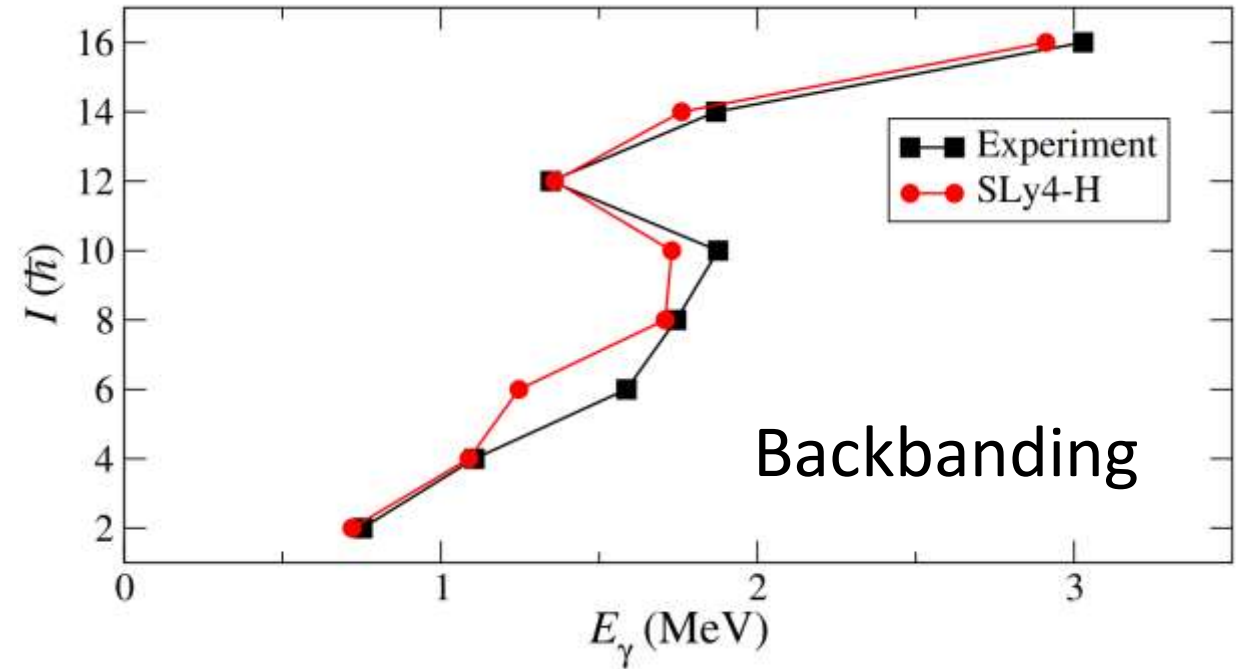
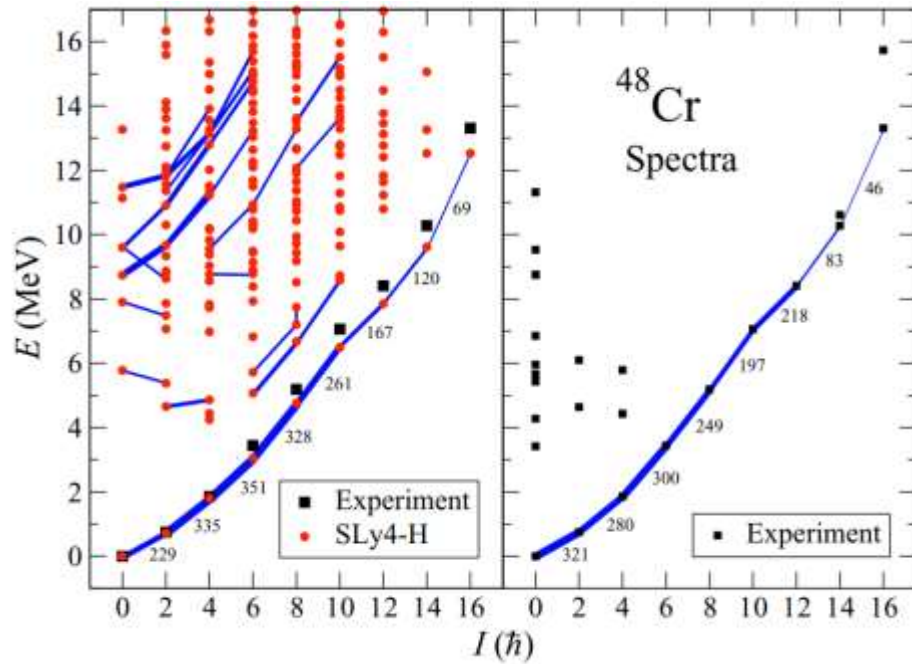
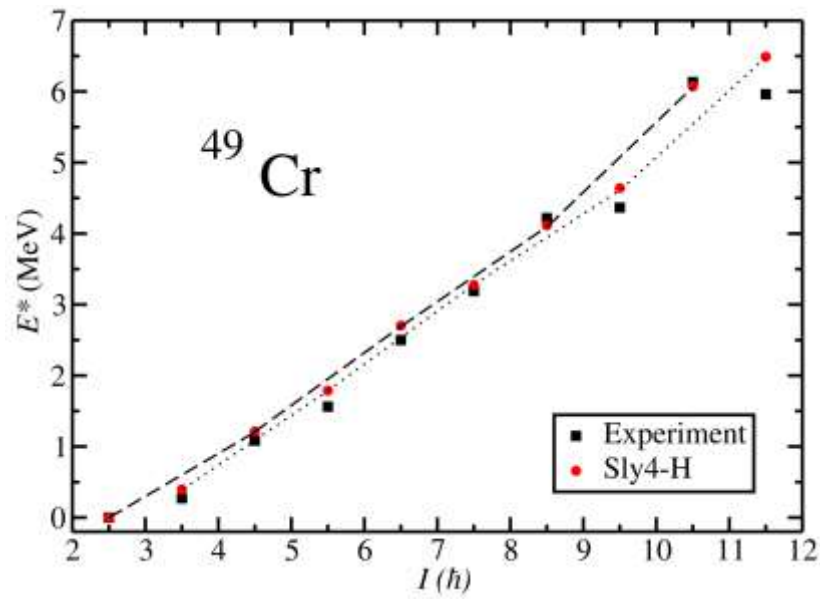


## Excitation mixing

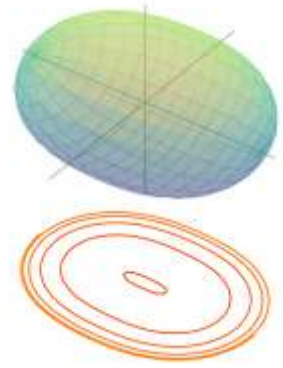
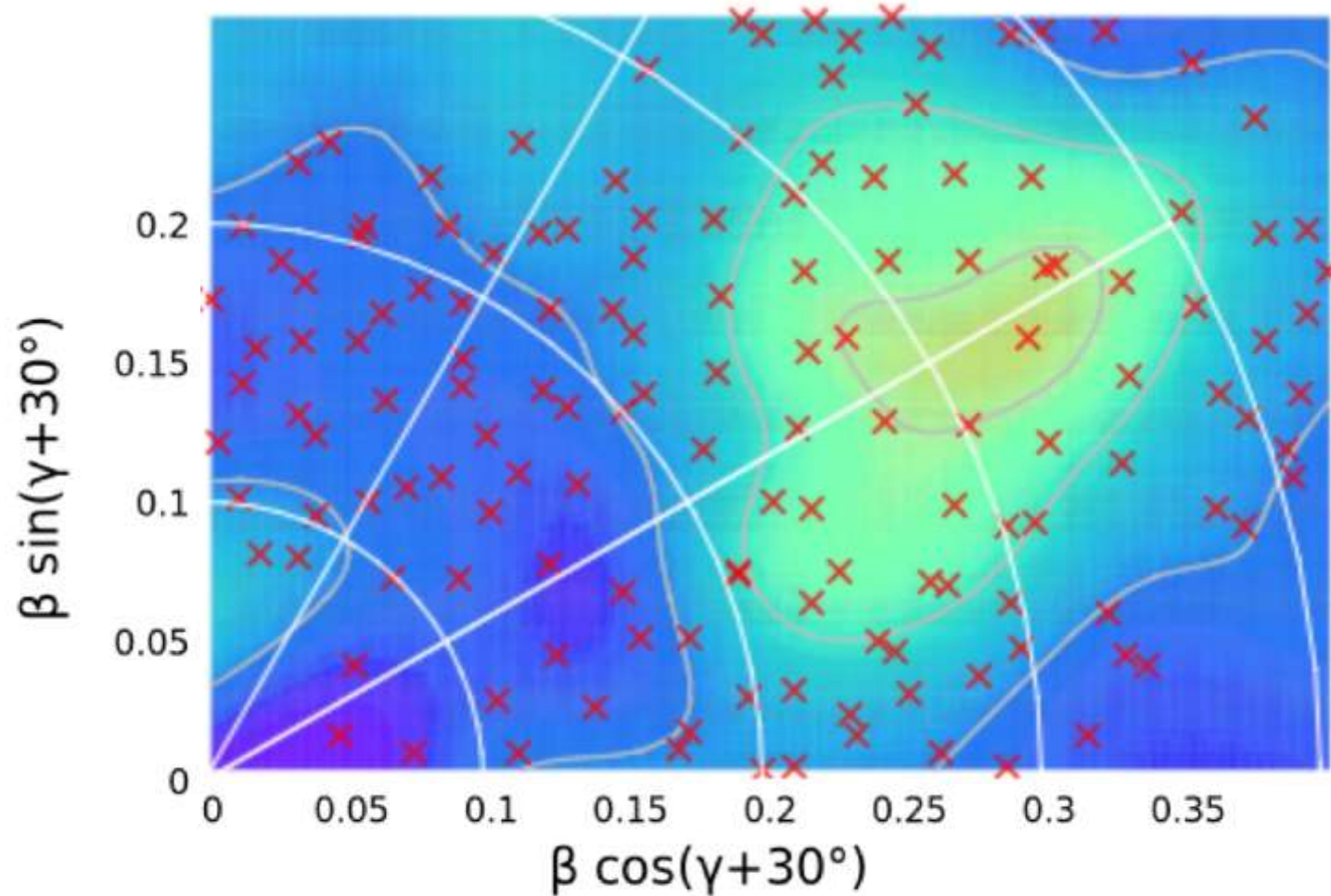




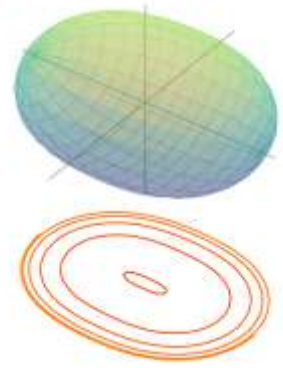
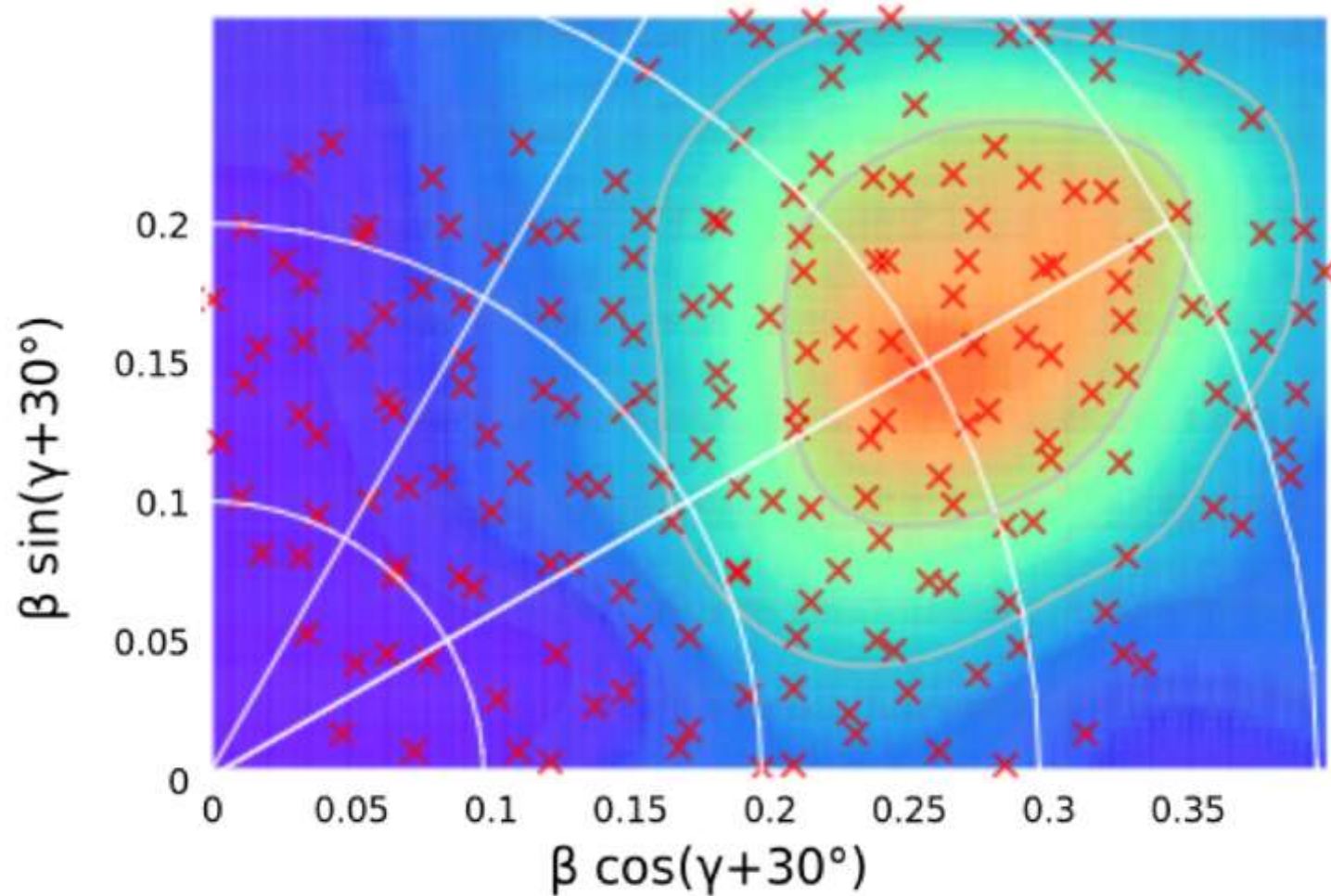
# Results: States and transitions



# Amplitudes of HFB-states for $^{48}\text{Cr}; I = 0$

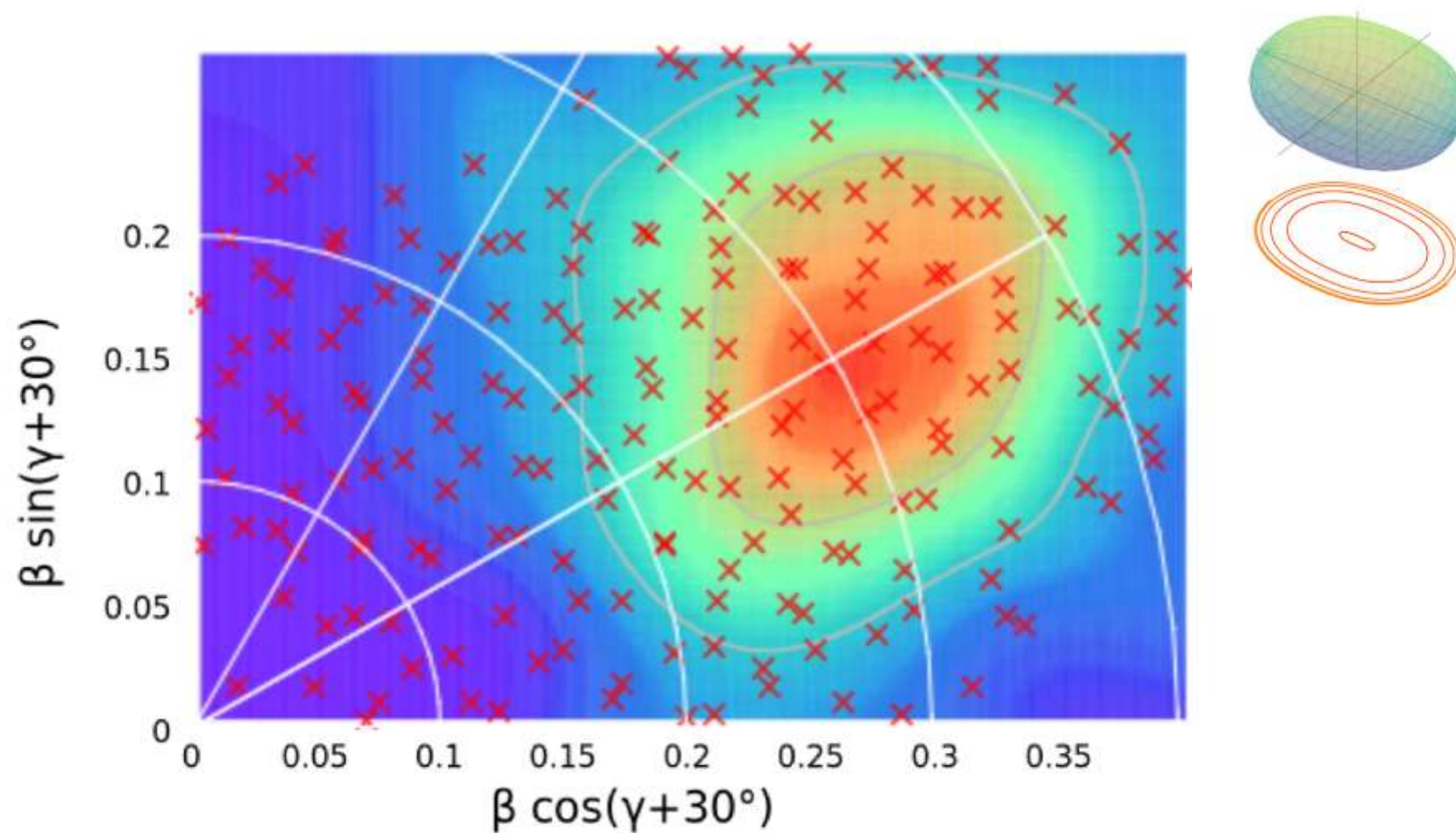


# Amplitudes of HFB-states for $^{48}\text{Cr}; I = 2$

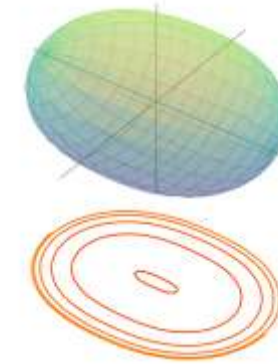
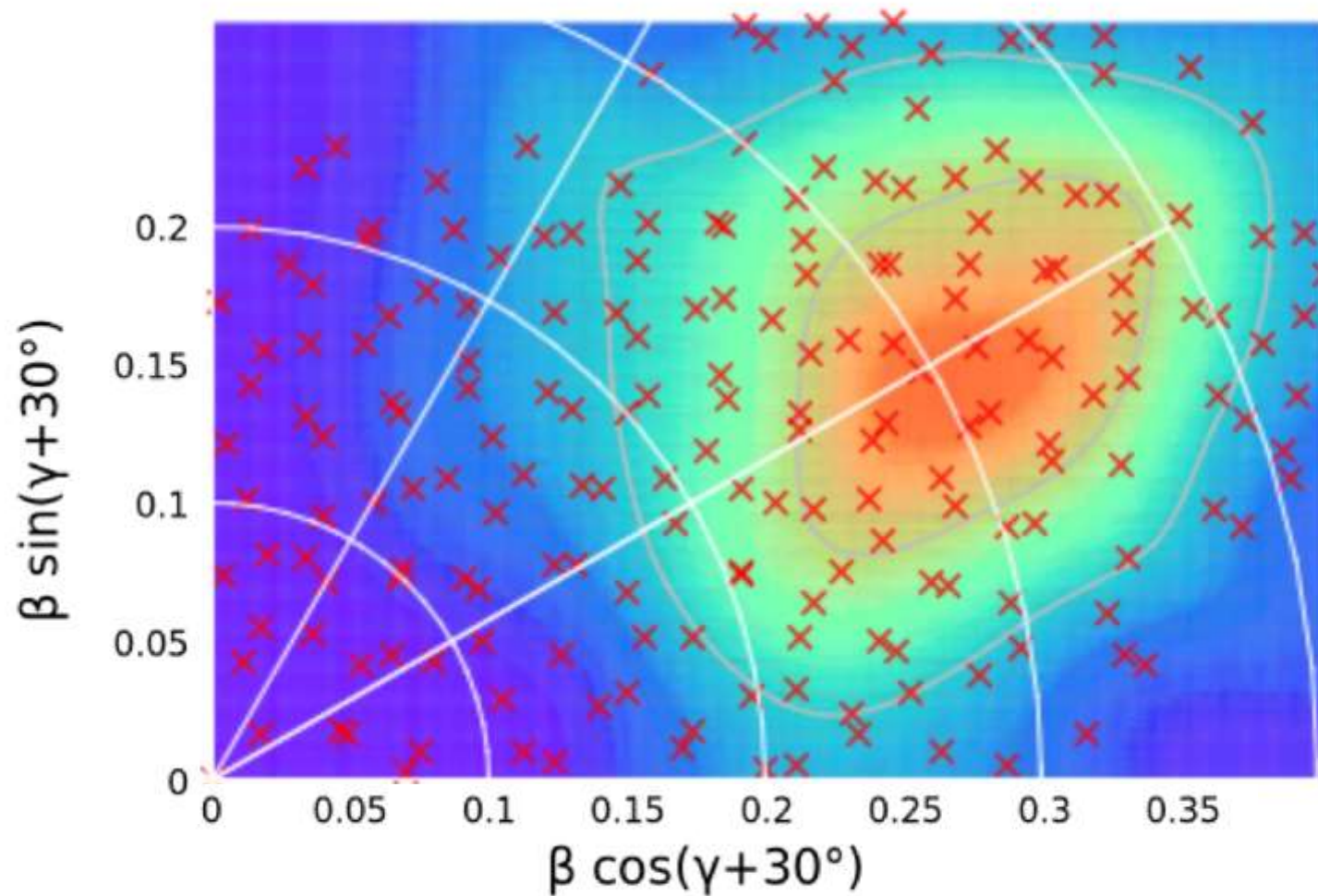




# Amplitudes of HFB-states for $^{48}\text{Cr}; I = 4$

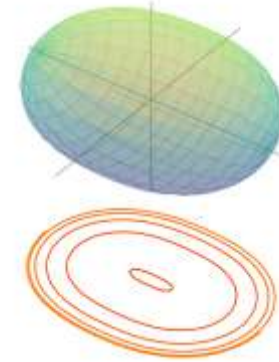
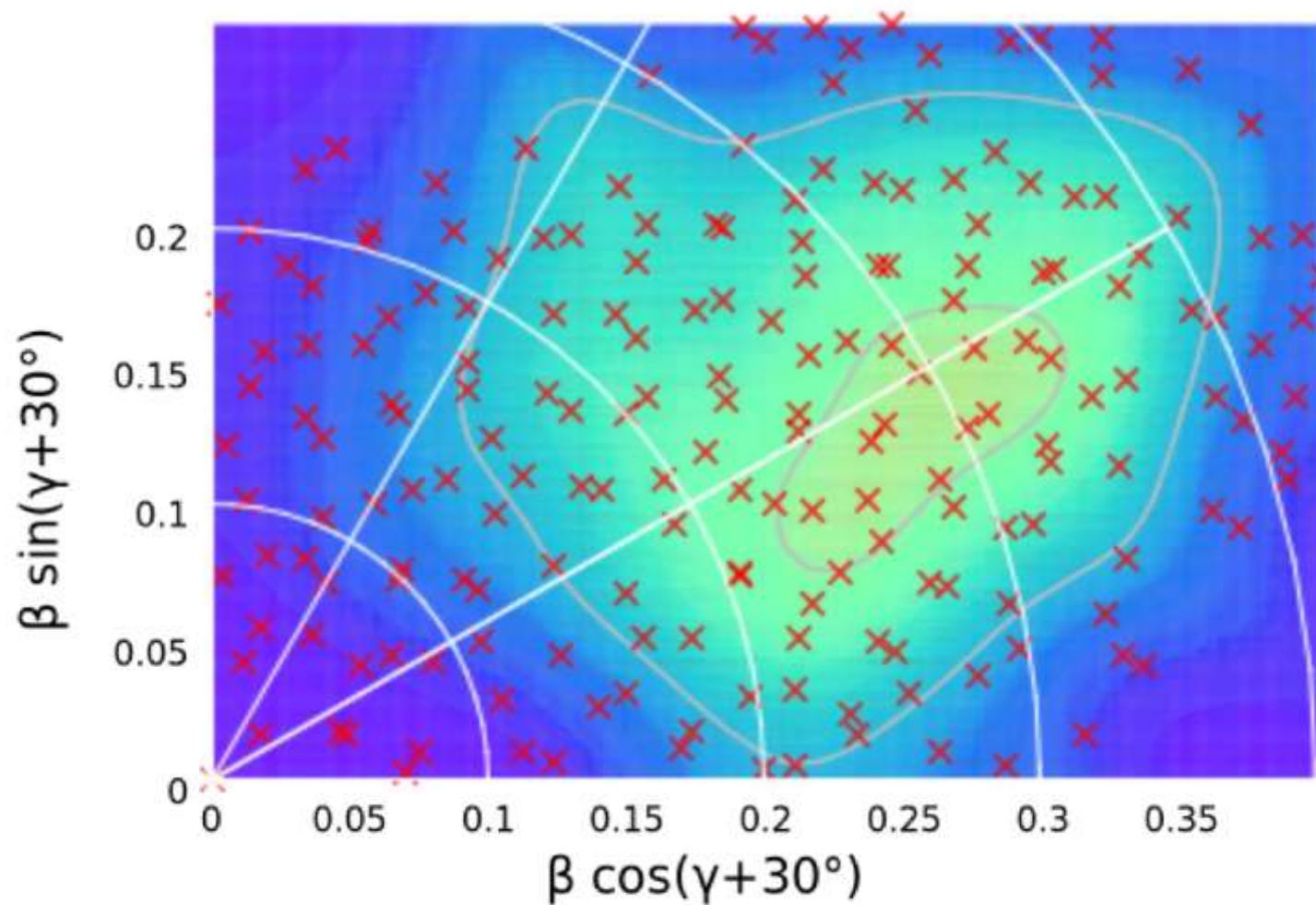


# Amplitudes of HFB-states for $^{48}\text{Cr}; I = 6$

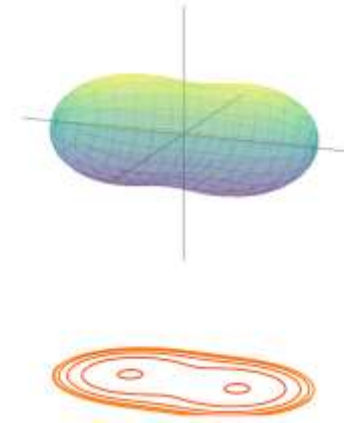
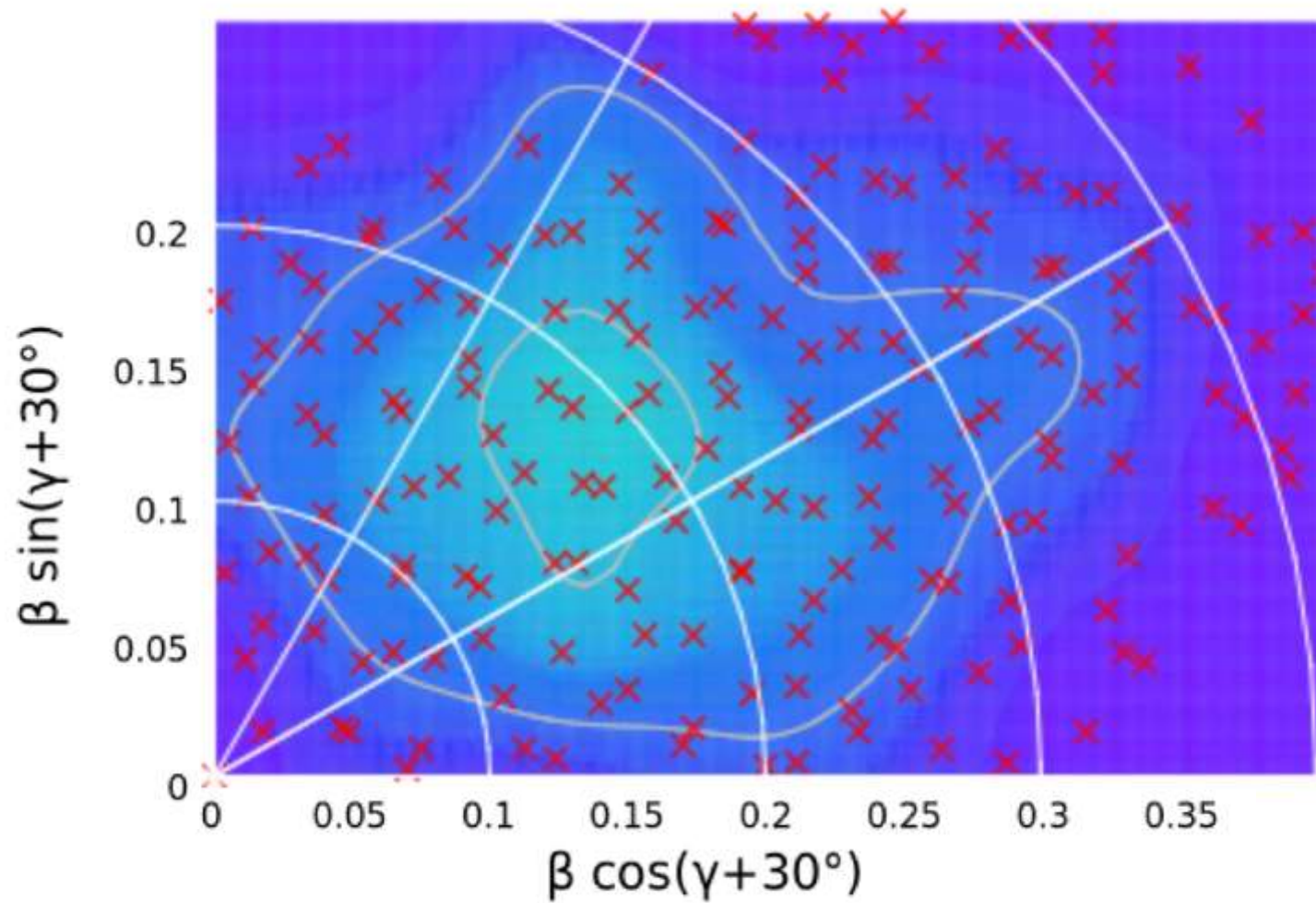




# Amplitudes of HFB-states for $^{48}\text{Cr}$ ; $I = 8$

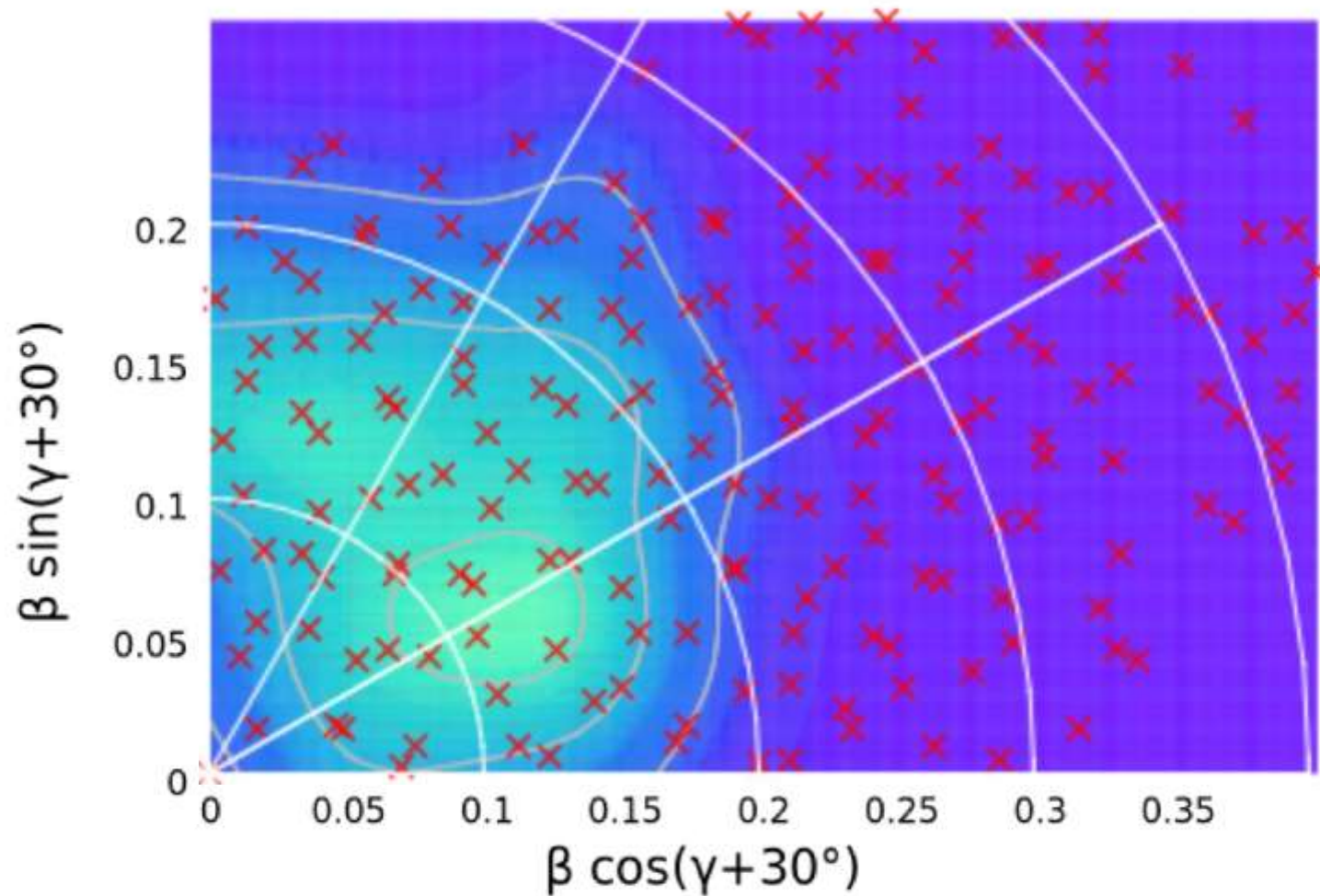


# Amplitudes of HFB-states for $^{48}\text{Cr}$ ; $I = 10$

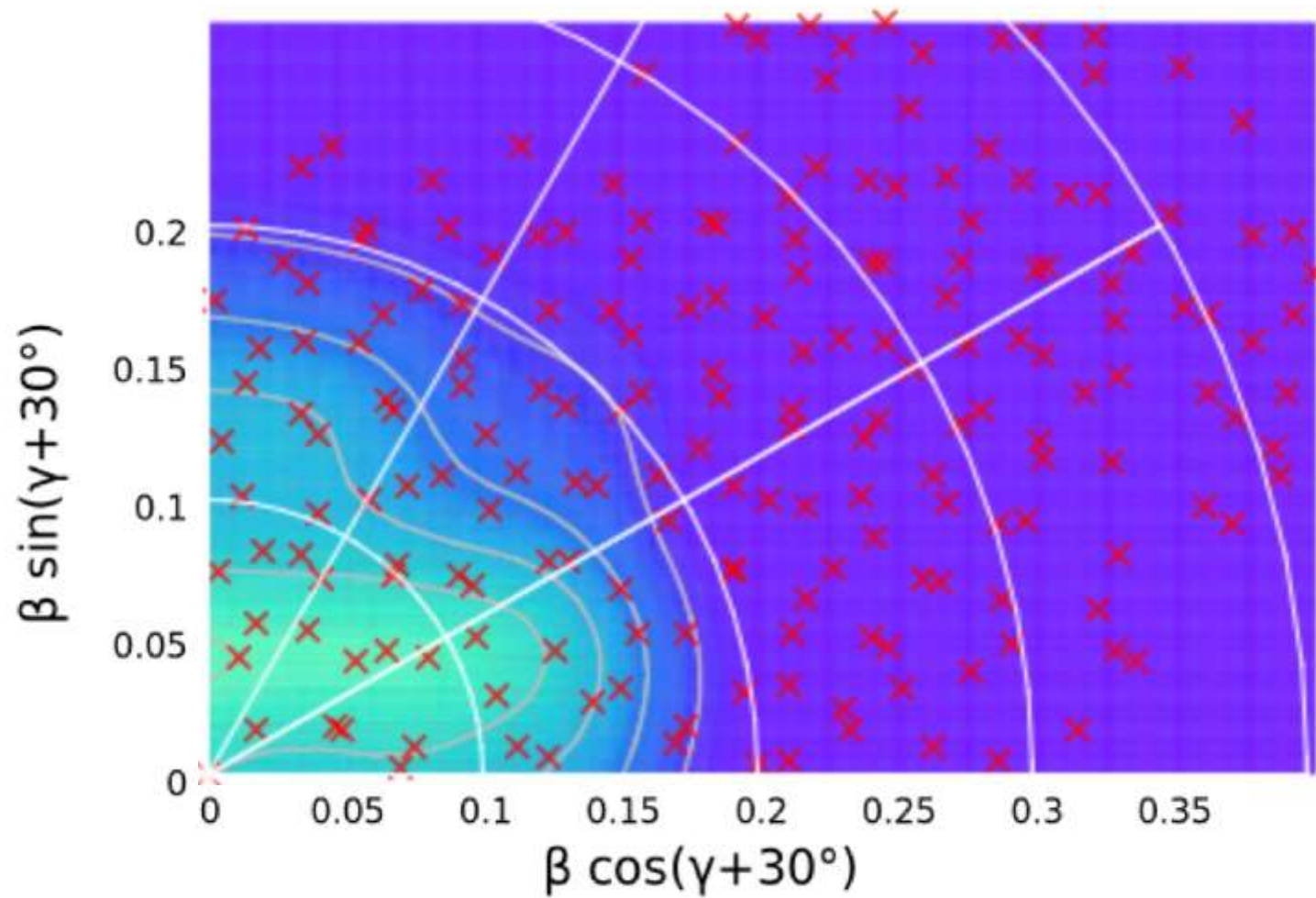




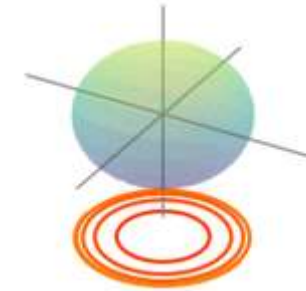
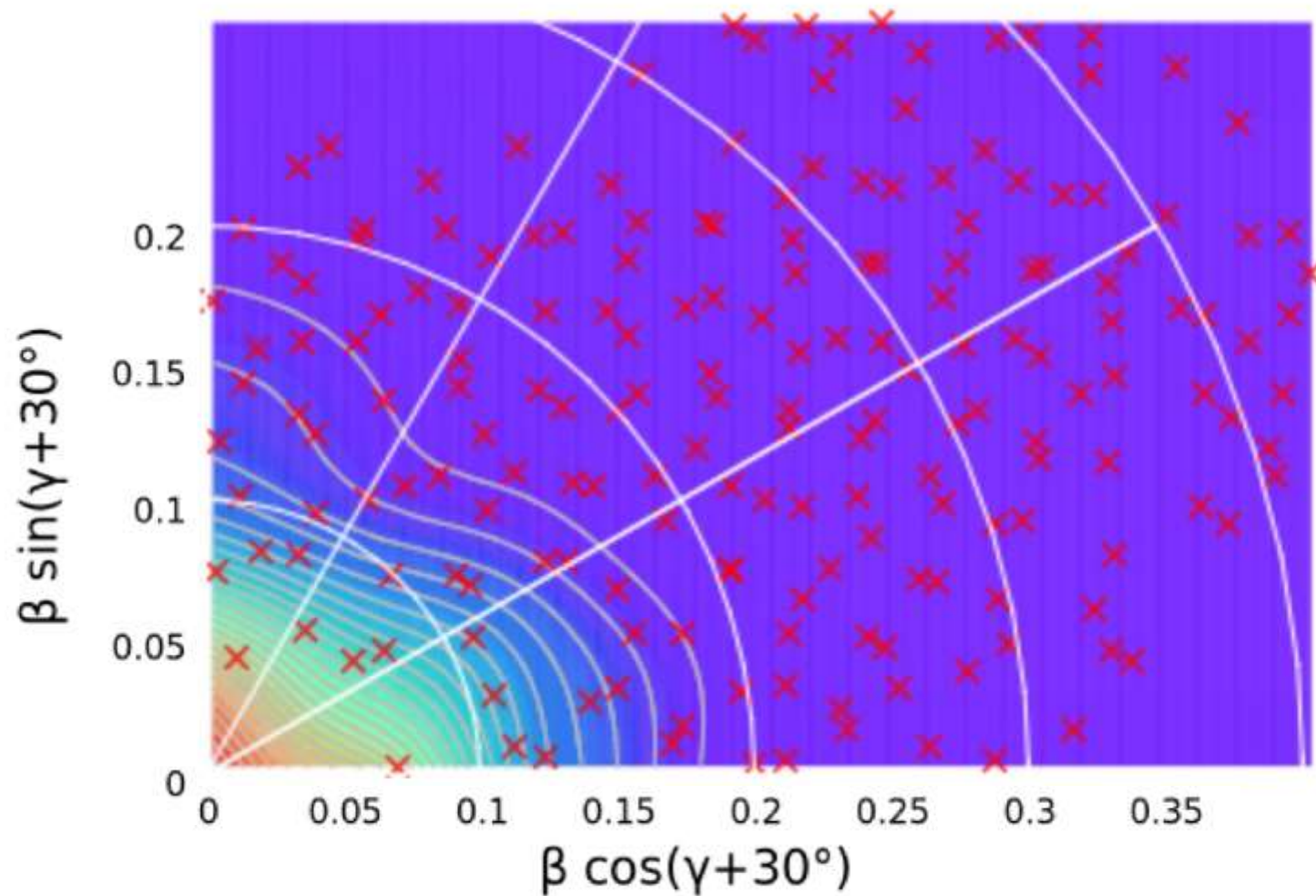
# Amplitudes of HFB-states for $^{48}\text{Cr}$ ; $I = 12$



# Amplitudes of HFB-states for $^{48}\text{Cr}$ ; $I = 14$

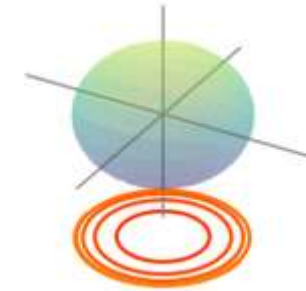
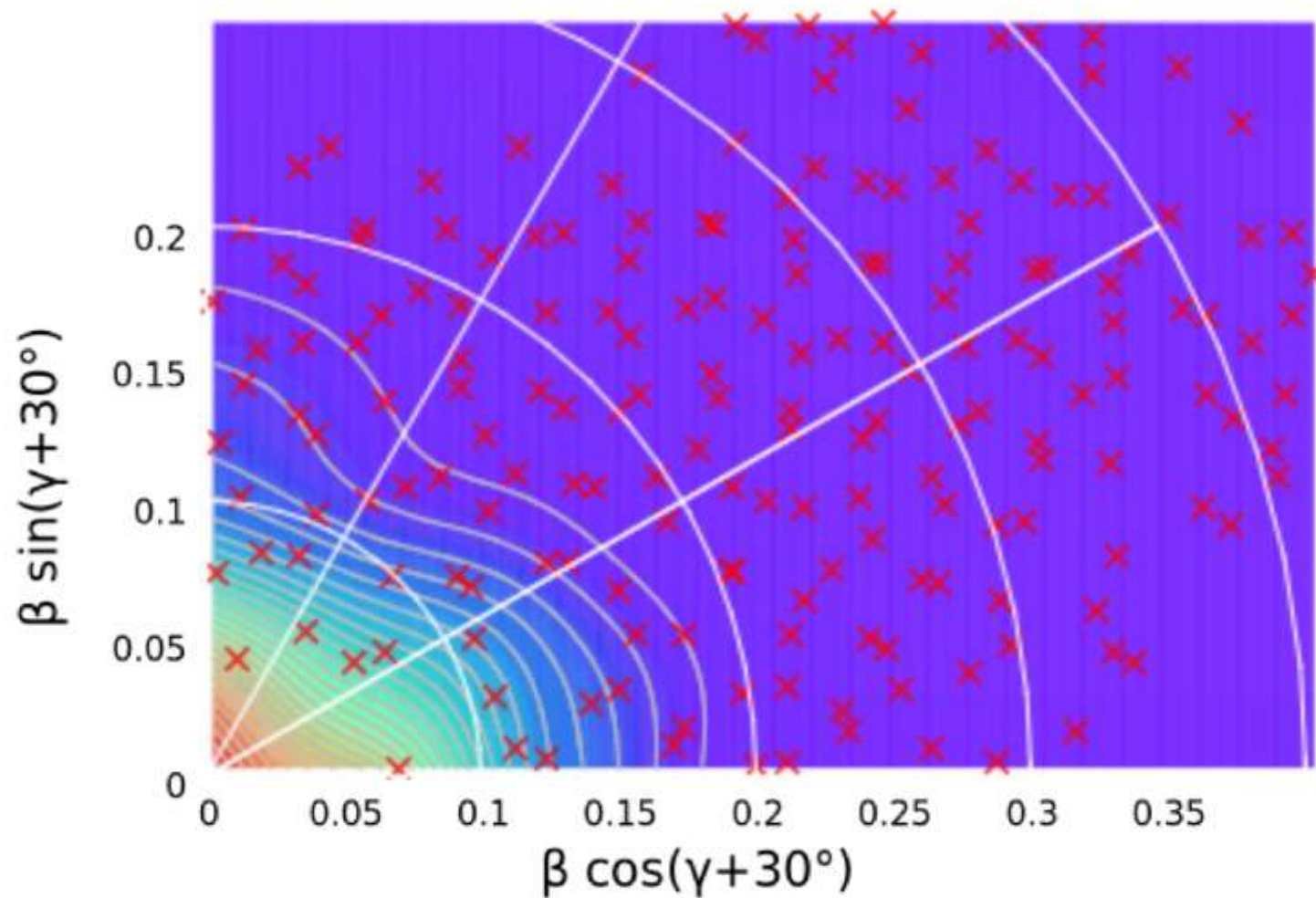


# Amplitudes of HFB-states for $^{48}\text{Cr}$ ; $I = 16$





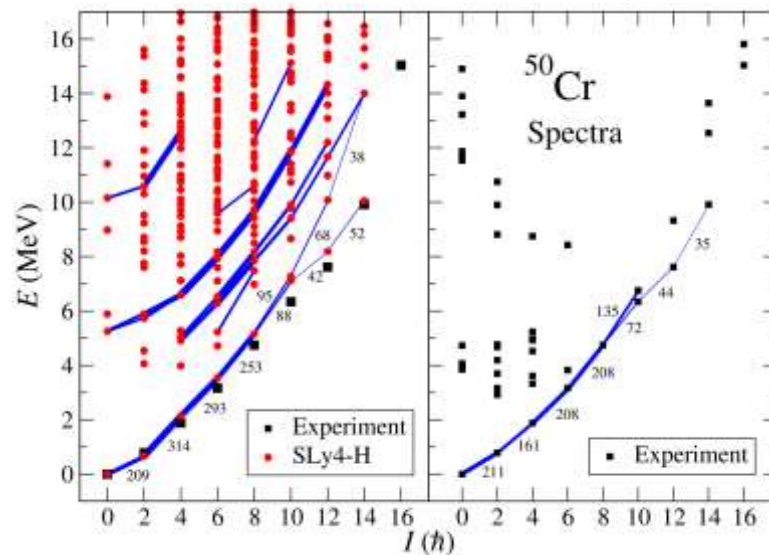
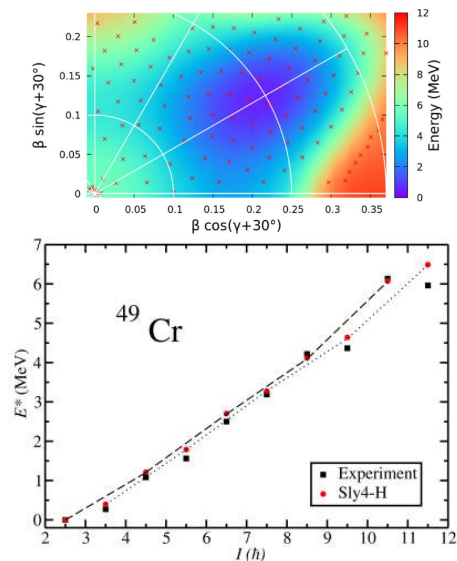
# Amplitudes of HFB-states for $^{48}\text{Cr}$ ; $I = 16$



4 x 2 particles  
in  $f_{7/2}$

# Conclusions & Outlook

- GCM can describe low energy states with a great degree of collectivity within a consistent elegant framework.
- Working towards:
  - Generalizing the effective Hamiltonian
  - Additional observables (e.g. scattering)



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arXiv:2211.14263  
arXiv:2212.07673

# GCM

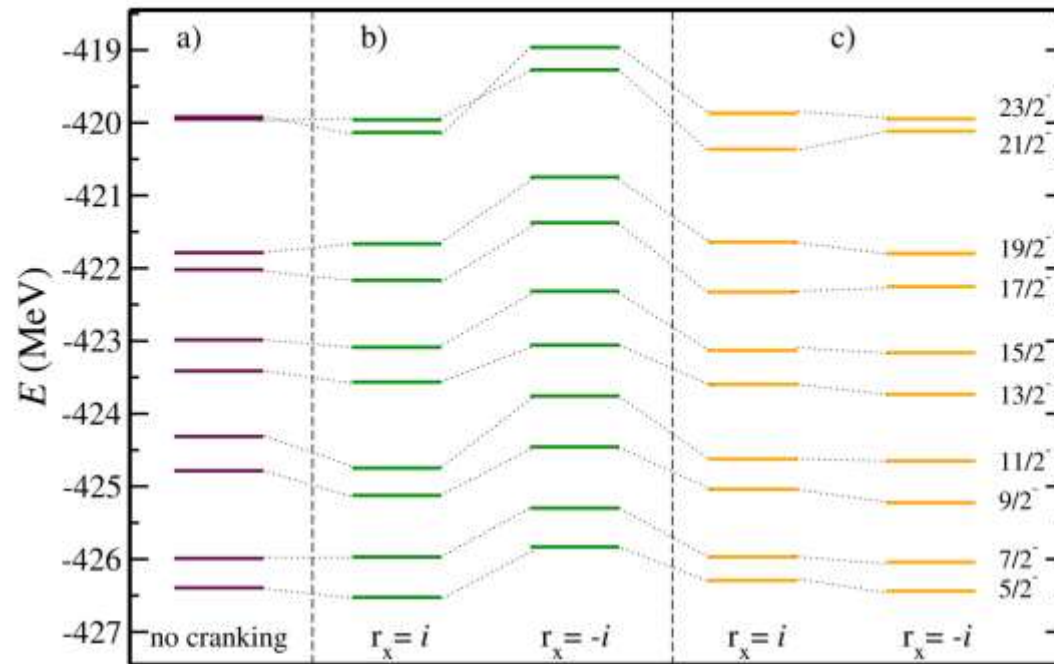


Figure 7: Energies of the g.s. rotational band in  $^{49}\text{Cr}$  computed in different bases. In the panel a) the basis include 116 states and the cranking is switched off. In the panels b) and c), respectively, 114 and 156 states are included in the basis and the cranking is switched on. Results on the left (right) side of panels b) and c) are obtained in a basis with signature



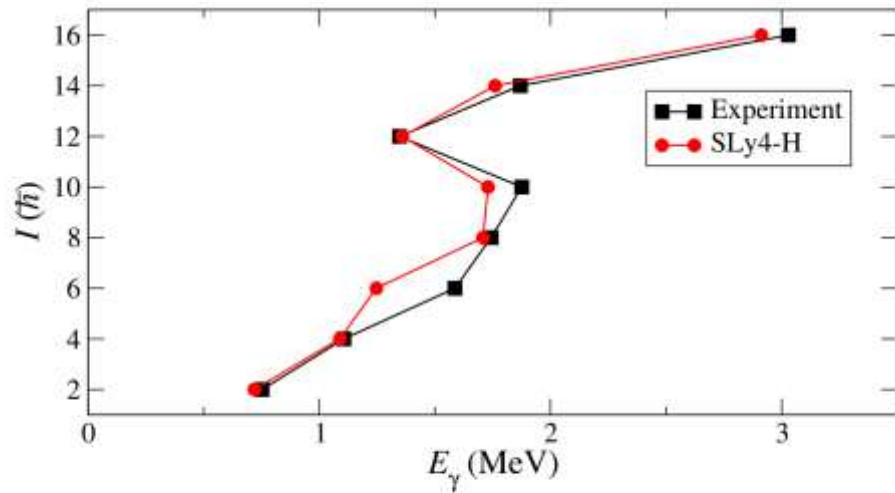
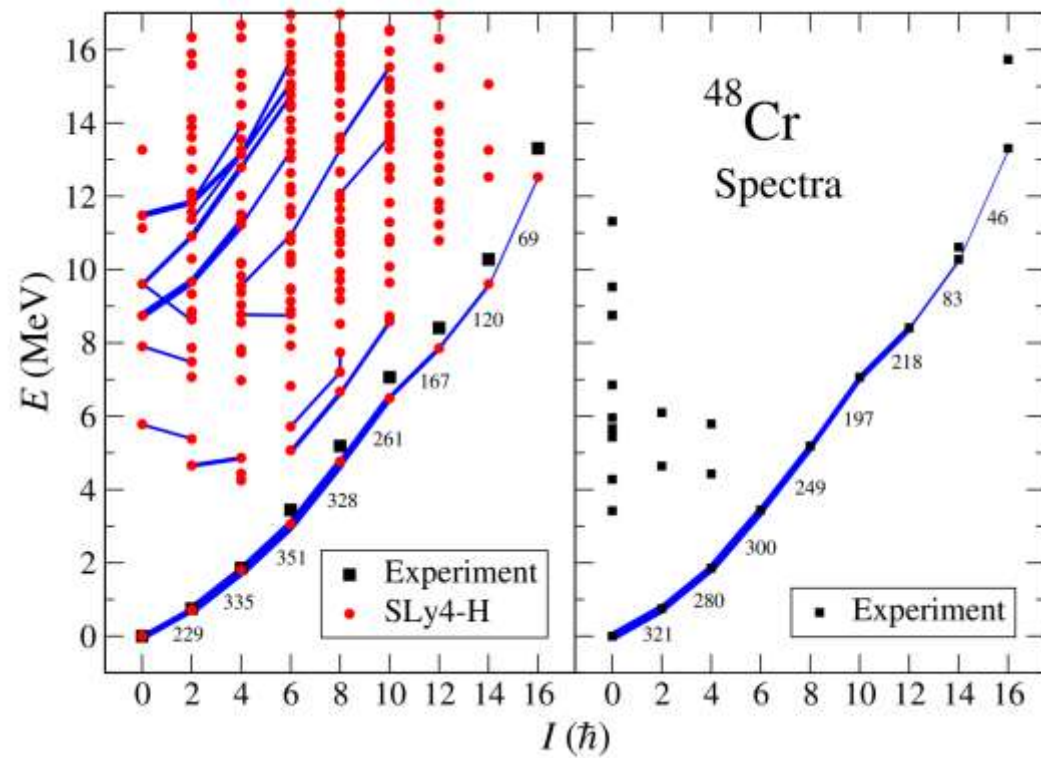
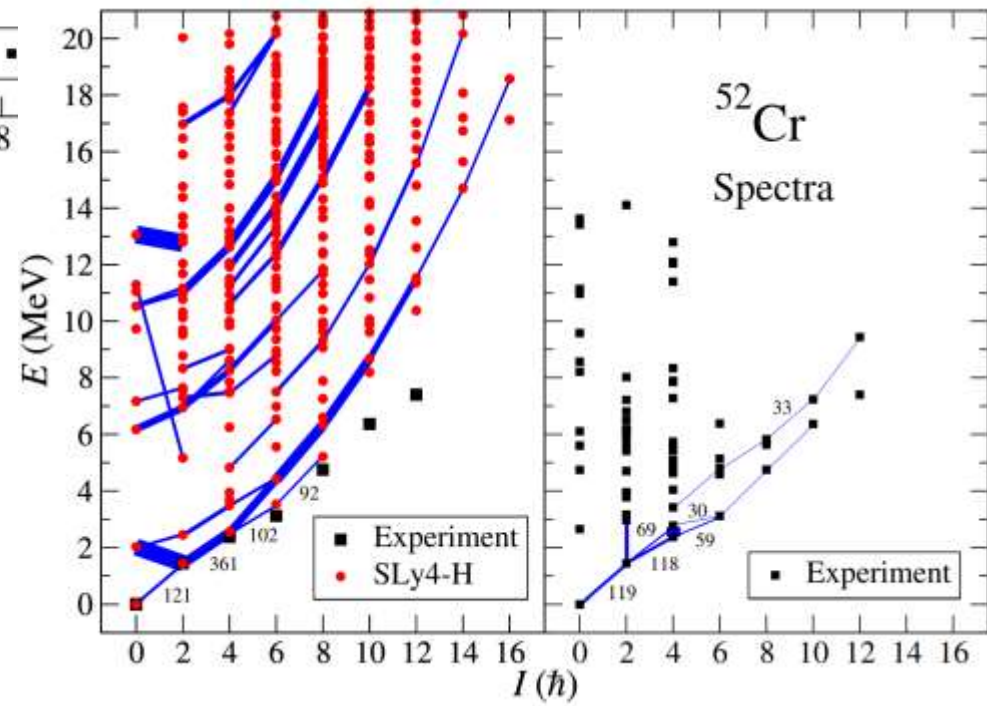
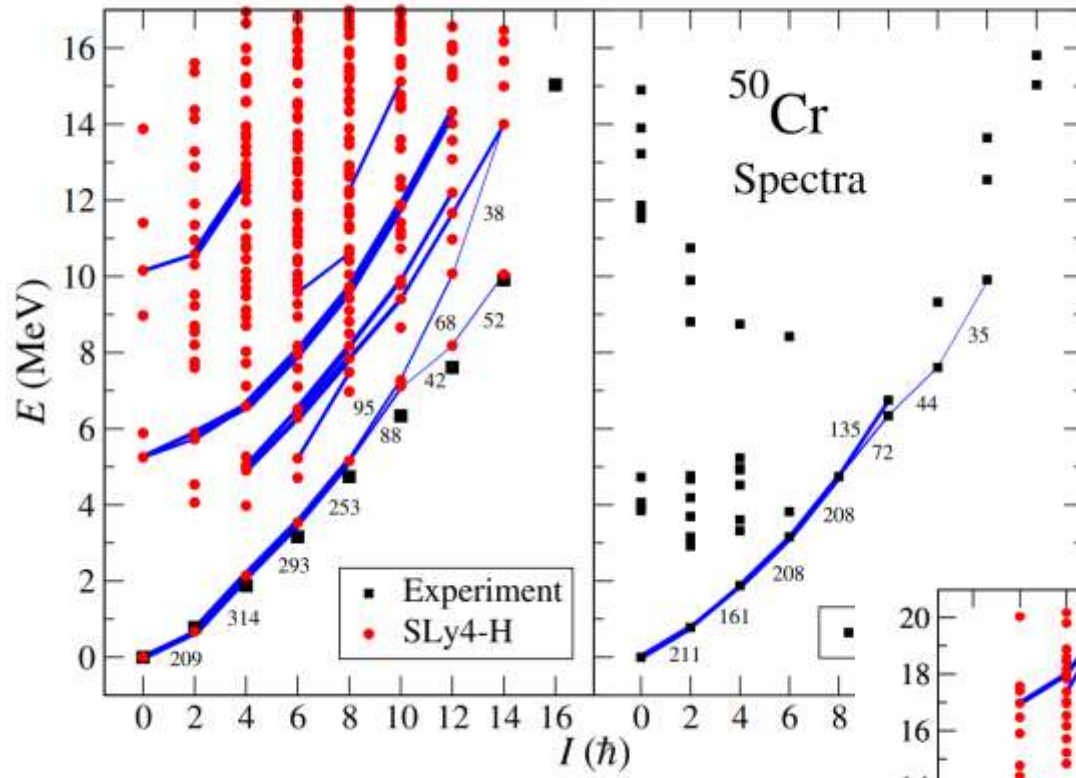
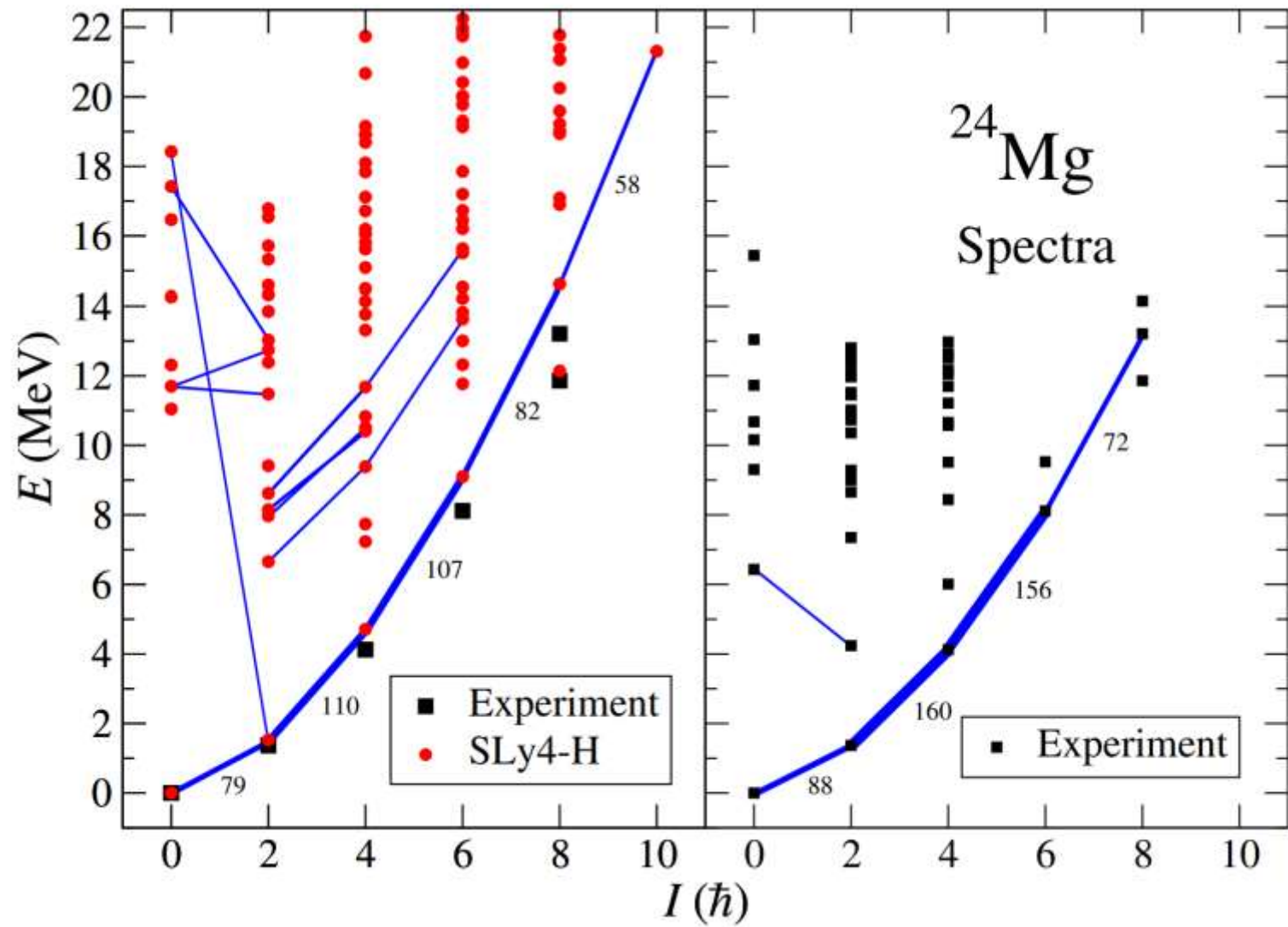
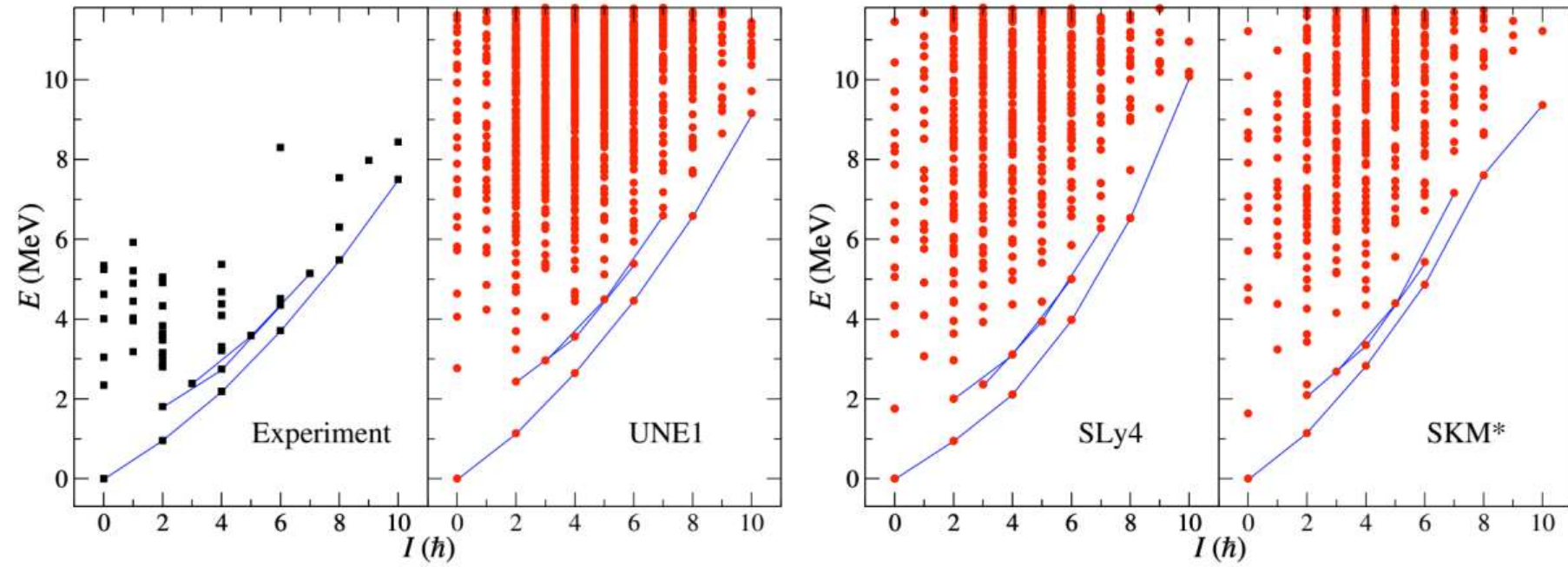


Figure 6: Released gamma ray energies for transit  $I$  for the yrast band. The backbending at  $I = 1$  a change in the internal structure.









$^{62}\text{Zn}$

