

# PRECISION PHYSICS AT THE LHC

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INFN Firenze

First Fellini Meeting

Roma  
24.02.2020



# OUTLINE

- 📌 **Past research activity**
  - 📌 Motivation
  - 📌 Brief description (state of the art (QCD) computations for LHC)
- 📌 **Fellini Project** - N<sup>3</sup>LO as the New Standard for Precision Physics at the LHC
  - 📌 The qT-subtraction/resummation formalism
  - 📌 The third order (N<sup>3</sup>LO)
  - 📌 Higgs Phenomenology at N<sup>3</sup>LO
- 📌 **Outlook**

# MOTIVATION

Why precision physics is important at the LHC?

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## Higgs Physics

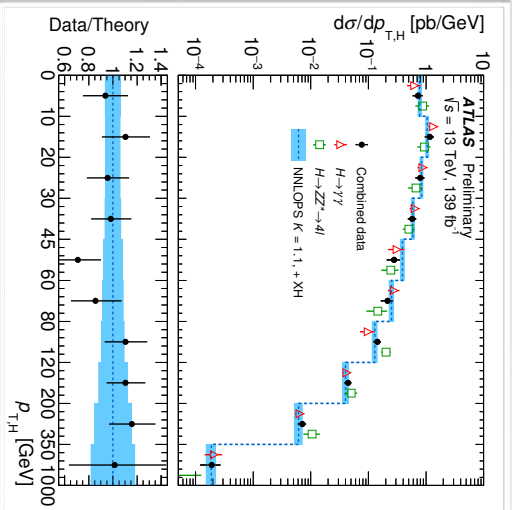
# MOTIVATION

Why precision physics is important at the LHC?

## Higgs Physics

The success in comparisons TH/Data in this slide wouldn't be possible with NLO precision

15 yrs ago



Combined results for  $H \rightarrow \gamma\gamma$  and  $H \rightarrow 4l$

ATLAS-CONF-2019-032

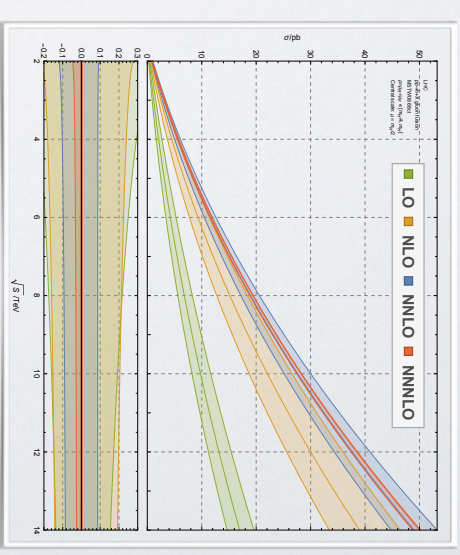
$55.4^{+4.3}_{-4.2}$  pb (  $\pm 3.1$  (stat.)  $^{+3.0}_{-2.8}$  (sys.) )

$55.6 \pm 2.5$  pb

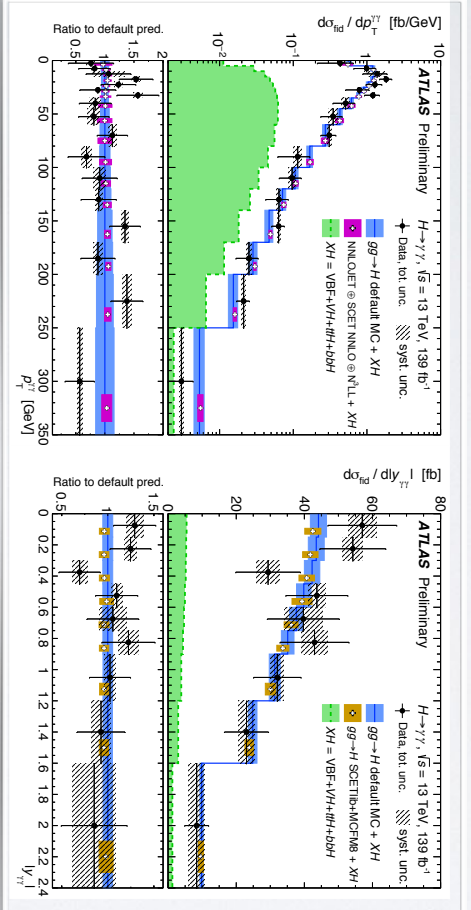
ATLAS

TH

Process	Accuracy	Fraction [%]
$ggF$	N <sup>3</sup> LO in QCD, NLO in EW	87.2
VBF	(approximate) NNLO in QCD, NLO in EW	6.8
VH	qq/gg: NNLO in QCD, NLO in EW; gg: NLO+NLL in QCD	4.0
$i\bar{i}H + tH$	$i\bar{i}H$ : NLO in QCD, NLO in EW; $tH$ : NLO in QCD	1.1
$b\bar{b}H$	NNLO (NLO) in QCD for 5FS (4FS)	0.9



Anastasiou, Duhr, Dulat, Herzog, Mistlberger (2015)



$\sigma_{fid} = 65.2 \pm 4.5$  (stat.)  $\pm 5.6$  (syst.)  $\pm 0.3$  (theo.) fb

$63.6 \pm 3.3$  fb

ATLAS

TH

ATLAS-CONF-2019-029

Results for  $H \rightarrow \gamma\gamma$

# MOTIVATION

Why precision physics is important at the LHC?

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 **Higgs Physics**

 **The 750 GeV excess**

# MOTIVATION

Why precision physics is important at the LHC?

## The 750 GeV excess

Precision Modeling and Model Building working together

Many interpretations in function of New Physics

More than 300 papers

The excess was detected by the ATLAS and CMS collaborations and published in December of 2015. It disappeared in July 2016

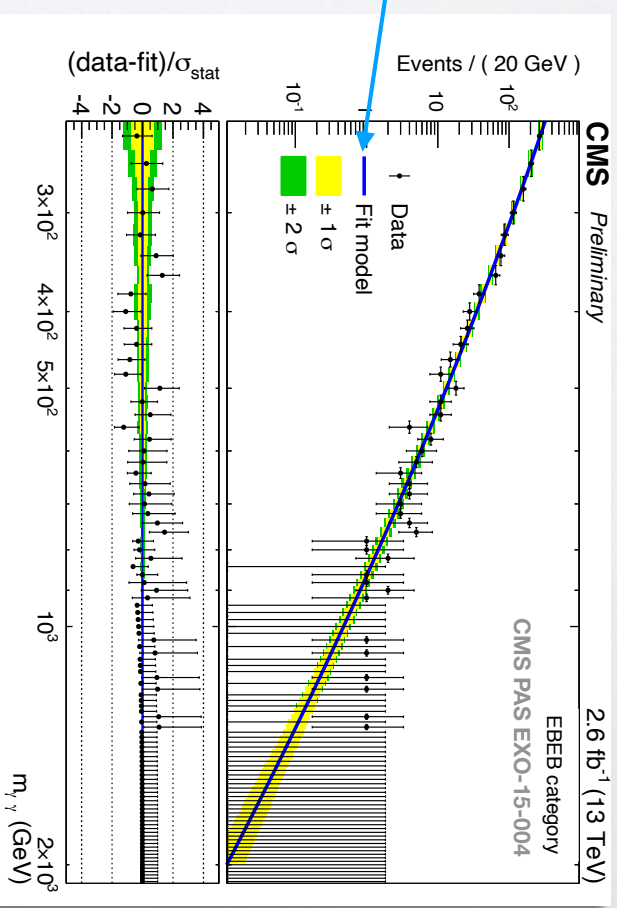
CMS PAS EXO-15-004

ATLAS-CONF-2015-081

More than 500 citations (each one)

Fit model tested with  $2\gamma$ NNLO

- S. Catani, L. C, D. de Florian, G. Ferrera and M. Grazzini (2011)
- S. Catani, L. C, D. de Florian, G. Ferrera and M. Grazzini (2018)
- L. C (2013)
- L. C (2012)
- L. C et. al - Les Houches proceedings (2015)
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# MOTIVATION

Why precision physics is important at the LHC?

## The 750 GeV excess

Precision Modeling and Model Building working together

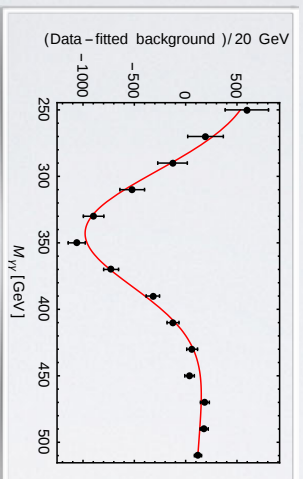
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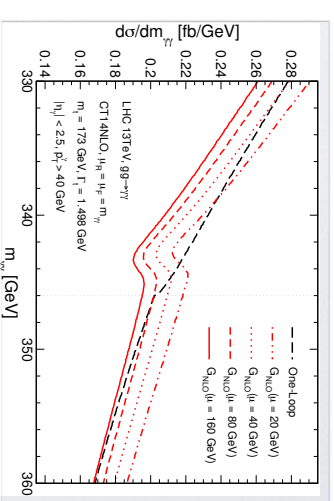
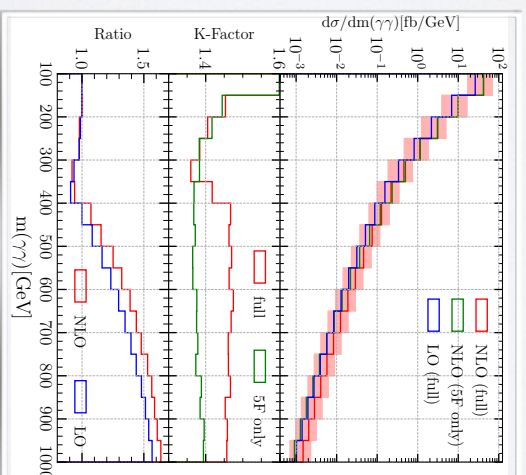
ATLAS-CONF-2015-081

A fitting function is OK, but we have to be aware that backgrounds are not always completely flat

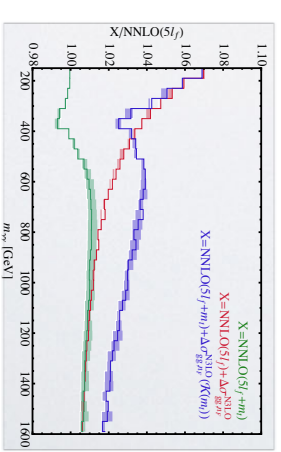


Dugad, Jain, Mitra, Sanyal, Verma (2016)

Maltoni, Mandal, Zhao (2019)



Kawabata, Yokoya (2016)



Campbell, Ellis, Li, Williams (2016)

# MOTIVATION

Why precision physics is important at the LHC?

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- 📌 Higgs Physics
- 📌 The 750 GeV excess
- 📌 Bottom Hadroproduction at the Tevatron



# MOTIVATION

Why precision physics is important at the LHC?

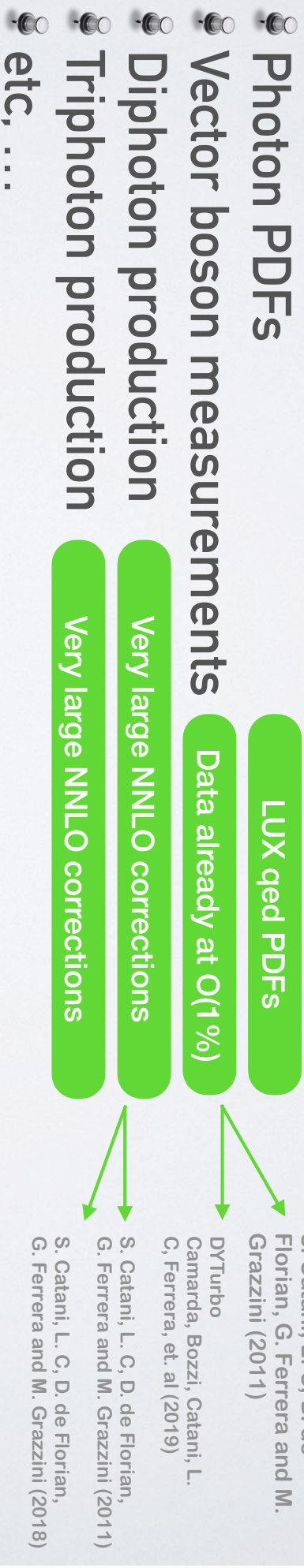
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- **Higgs Physics**
- **The 750 GeV excess**
- **Bottom Hadroproduction at the Tevatron**
- Photon PDFs
- Vector boson measurements
- Diphoton production
- Triphoton production
- etc, ...

# MOTIVATION

Why precision physics is important at the LHC?

- Higgs Physics
- The 750 GeV excess
- Bottom Hadroproduction at the Tevatron



# MOTIVATION

Why precision physics is important at the LHC?

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- **Higgs Physics**
- **The 750 GeV excess**
- **Bottom Hadroproduction at the Tevatron**
- Photon PDFs
- Vector boson measurements
- Diphoton production
- Triphoton production
- etc, ...
- **New interesting tasks are just recently initiated: Yukawas**

# MOTIVATION

Why precision physics is important at the LHC?

## ESPPU 2013

- g) Europe should support a diverse, vibrant *theoretical physics programme*, ranging from abstract to applied topics, in close collaboration with experiments and extending to neighbouring fields such as *astroparticle physics* and cosmology. Such support should extend also to high-performance computing and software development.
- Great progress since last ESPP in development of (NLO) MC generators and higher-order calculations (e.g. Higgs production at N3LO, NNLO for 2→2 QCD processes) to match precision of LHC (and other) data. Vibrant community in Europe, which also contributed significantly to the study of future facilities. Crucial funding from H2020 (e.g. MChet).  
Note: strong support for theory is crucial for the future of the field: theoretical ideas have led to ground breaking developments in the understanding of nature and have inspired many experimental searches. MC developments and higher-order (signal and background) calculations are fundamental for current and future projects → HEP theory should be supported in Universities and labs and should be an integral part of planning for and funding of future colliders and other projects.

Fabiola Gianotti (Granada 2019)

## physicsworld

PARTICLE AND NUCLEAR | INTERVIEW

### Preparing for a post-LHC future

22 Feb 2019

Richard Blaustein talks to CERN director-general **Fabiola Gianotti** about how the lab is planning the next big experiment in particle physics beyond the Large Hadron Collider



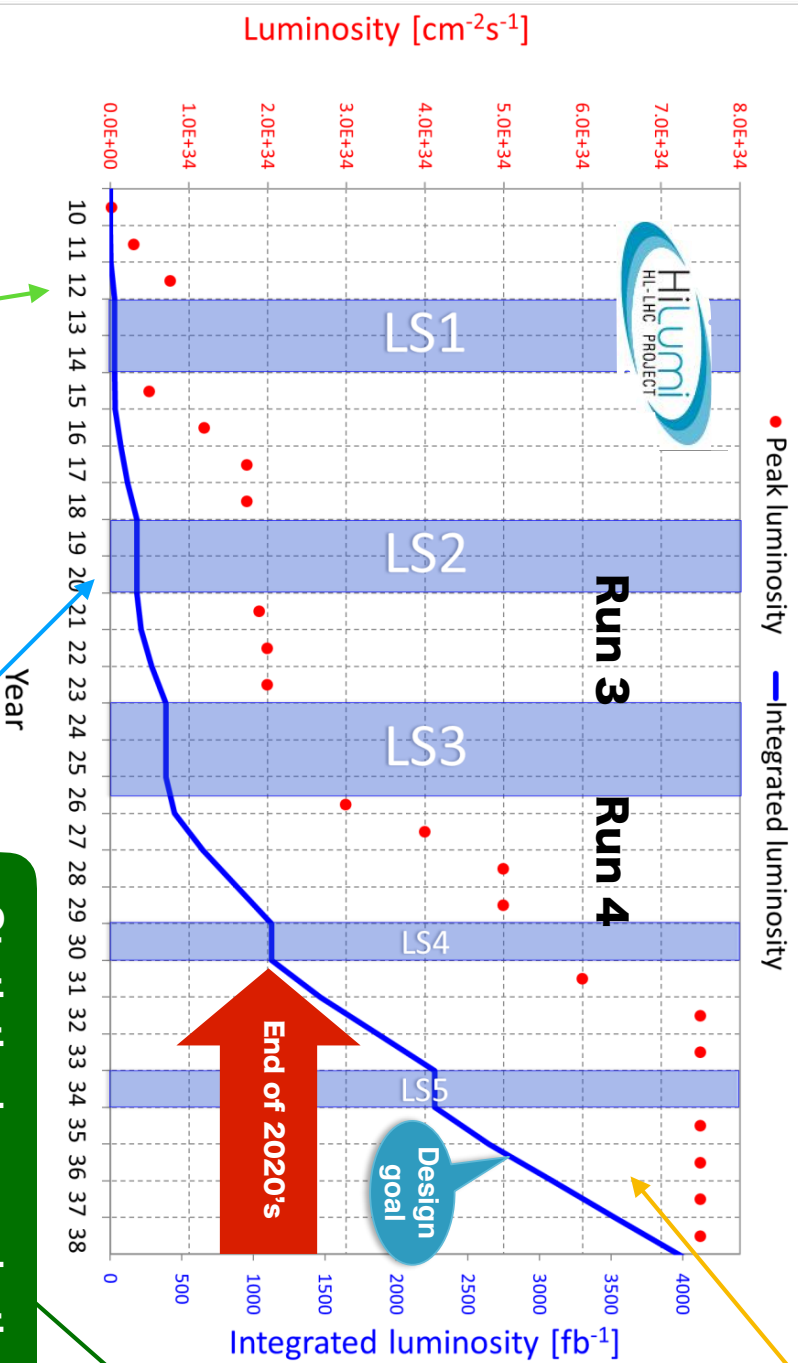
CERN director-general Fabiola Gianotti speaking at the 2019 AAAS meeting in Washington, DC. (Courtesy: Robb Cohen Photography & Video)

“ Precise measurements of known particles and interactions are just as important as finding new particles

# MOTIVATION

Why precision physics is important at the LHC?

HL-LHC lumi: 5-7x today's int.lumi by 2030, 20-30x by 2036



More than 30 yrs of LHC

### ATLAS and CMS

Run 3	Run 4	HL-LHC total
300 fb <sup>-1</sup>	1 ab <sup>-1</sup>	3 – 4 ab <sup>-1</sup>

### LHCb

Run 3	Run 4	HL-LHC total
23 fb <sup>-1</sup>	50 fb <sup>-1</sup>	300 fb <sup>-1</sup>

Statistical errors in the range 1% - 2%  
LHC Physics at % precision?

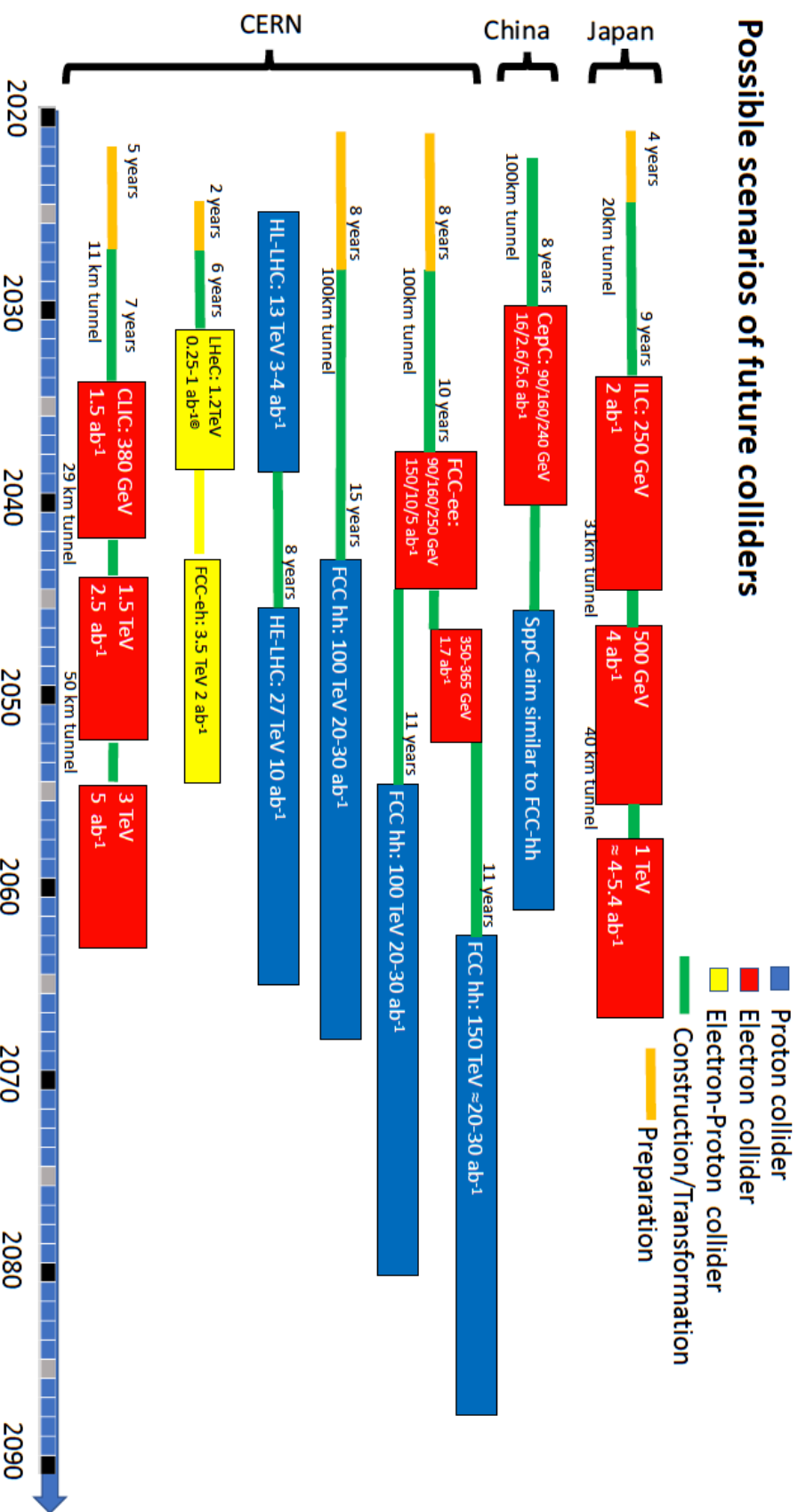
Higgs discovery

We are here

# MOTIVATION

Why precision physics is important at the LHC?

## Possible scenarios of future colliders



Summary of National Inputs

S. Bethke (MPP Munich)

ESPP Symposium, Granada, 15 May 2019

13/05/2019 UB

4

- all options aimed at attobarn<sup>-1</sup> physics
- requires to go far beyond NNLO for theory
- Even conservative estimates not reachable with current techniques

# STATE OF THE ART OF PRECISION PHYSICS AT THE LHC

- QCD perturbative corrections play a crucial role at modern colliders (LHC, Tevatron)
- Until a few years ago, the standard for such calculations was next-to-leading order (NLO) accuracy
- In the past recent years a number of growing next-to-next-to-leading order (NNLO) results were computed

# STATE OF THE ART OF PRECISION PHYSICS AT THE LHC

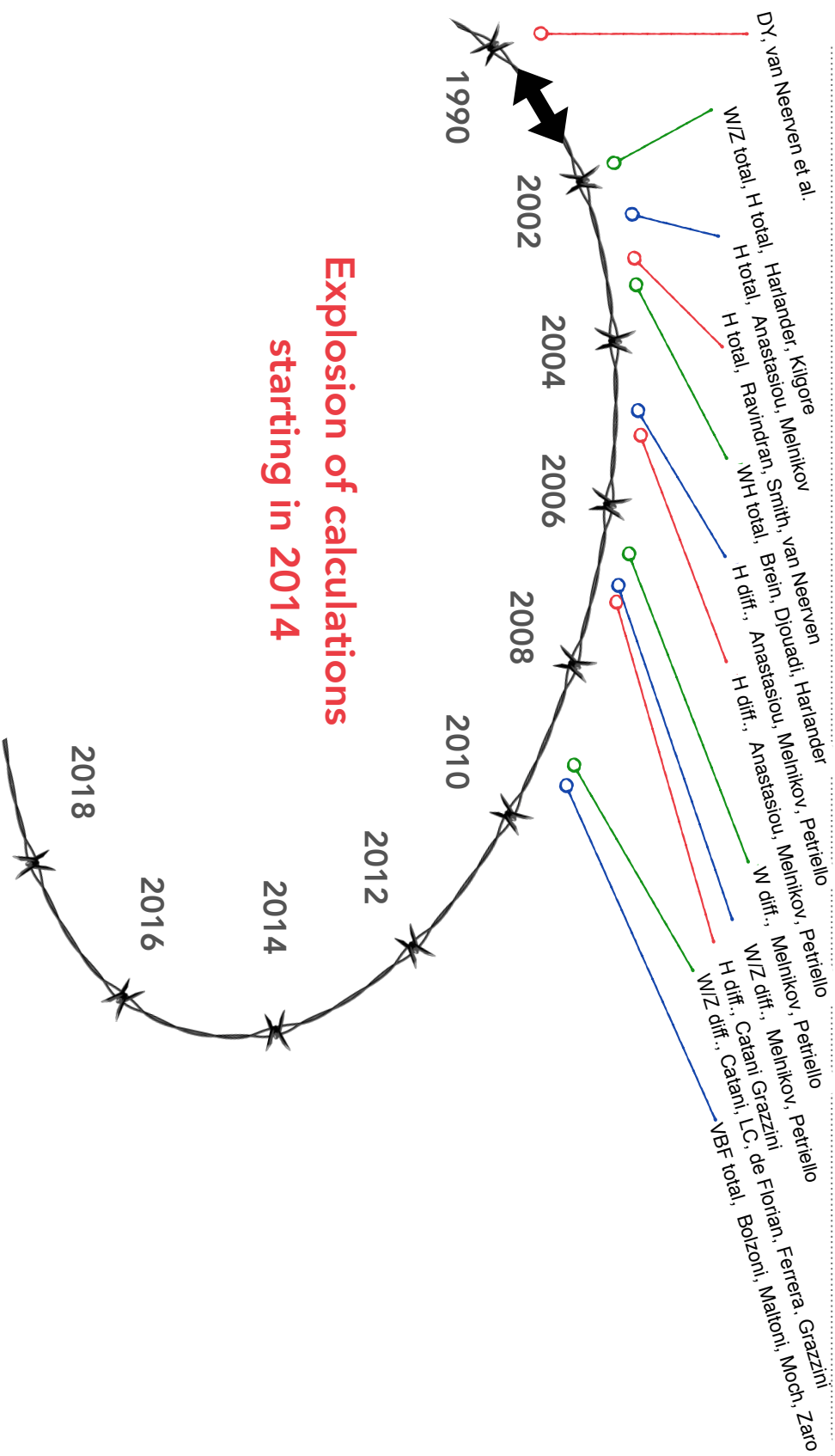
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“The NNLO revolution”



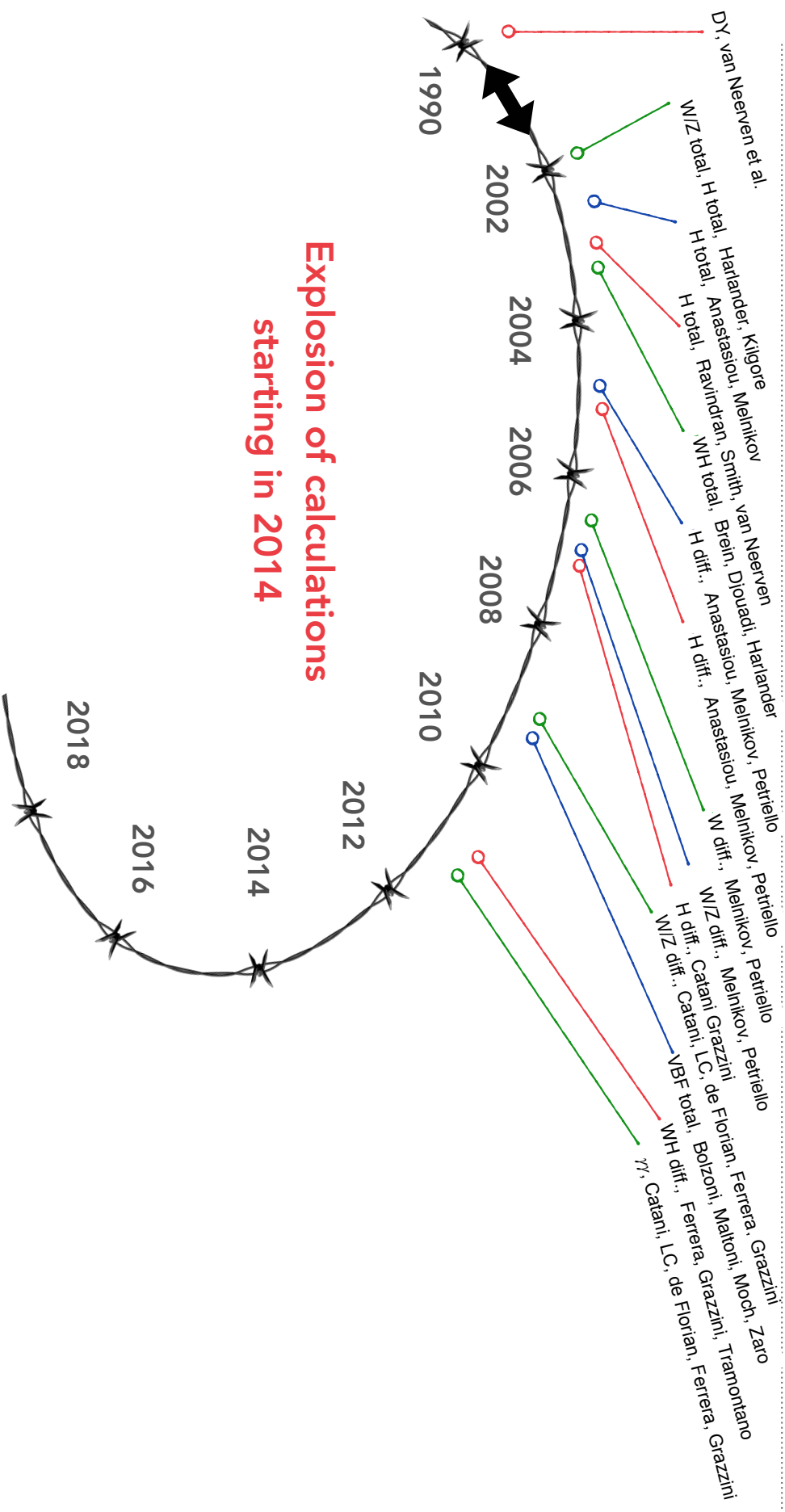
# THE NNLO STANDARD

## NNLO HADRON-COLLIDER CALCULATIONS VS. TIME



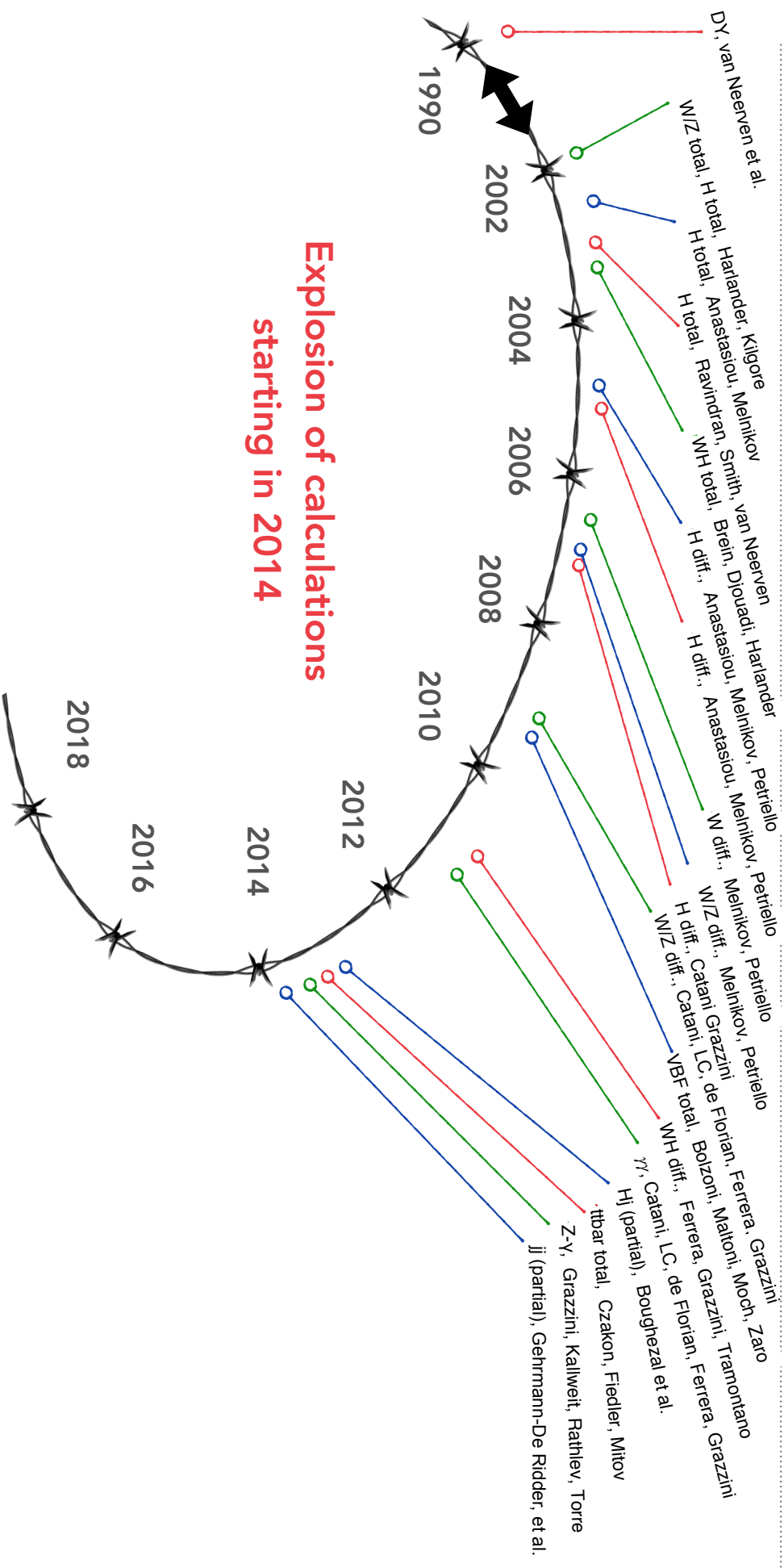
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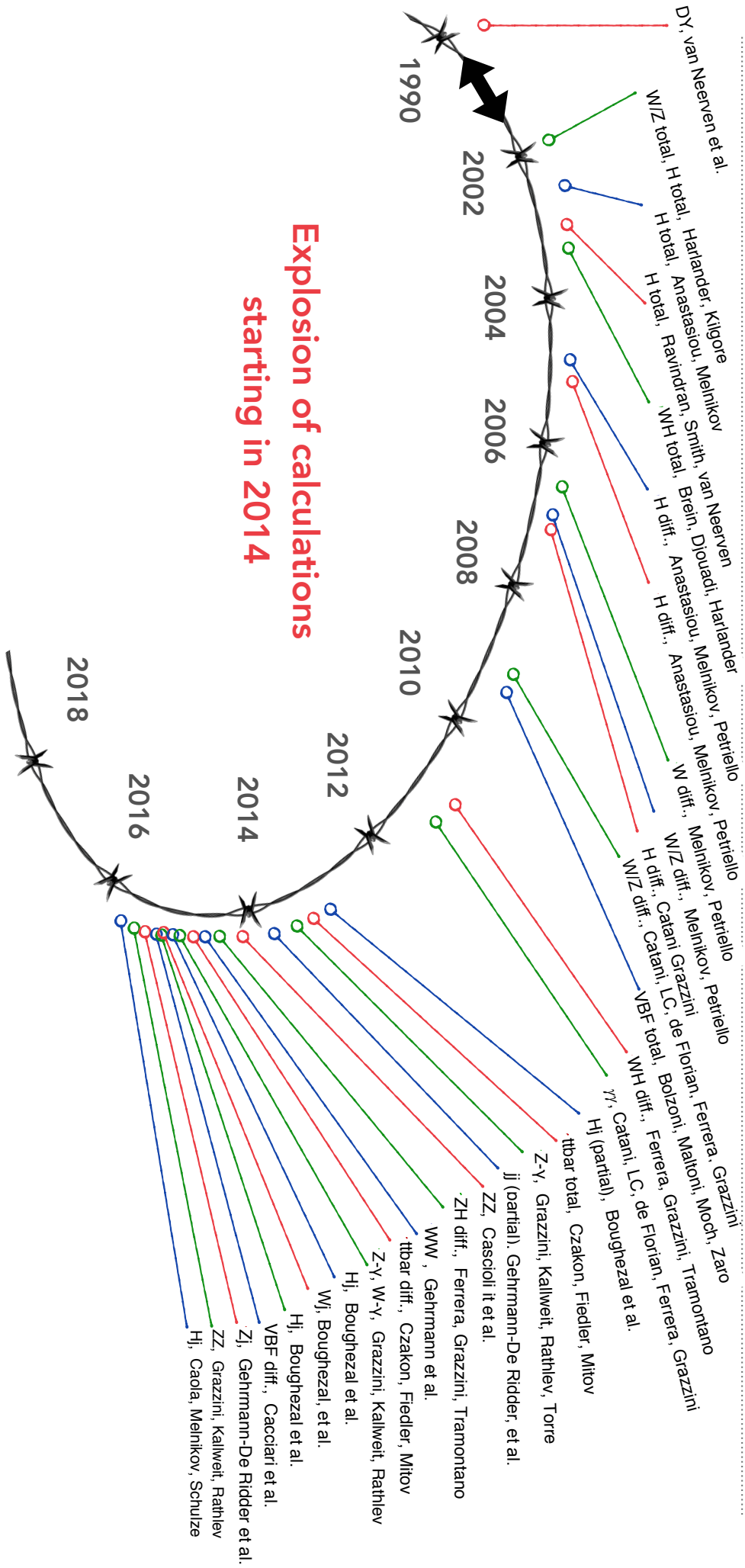
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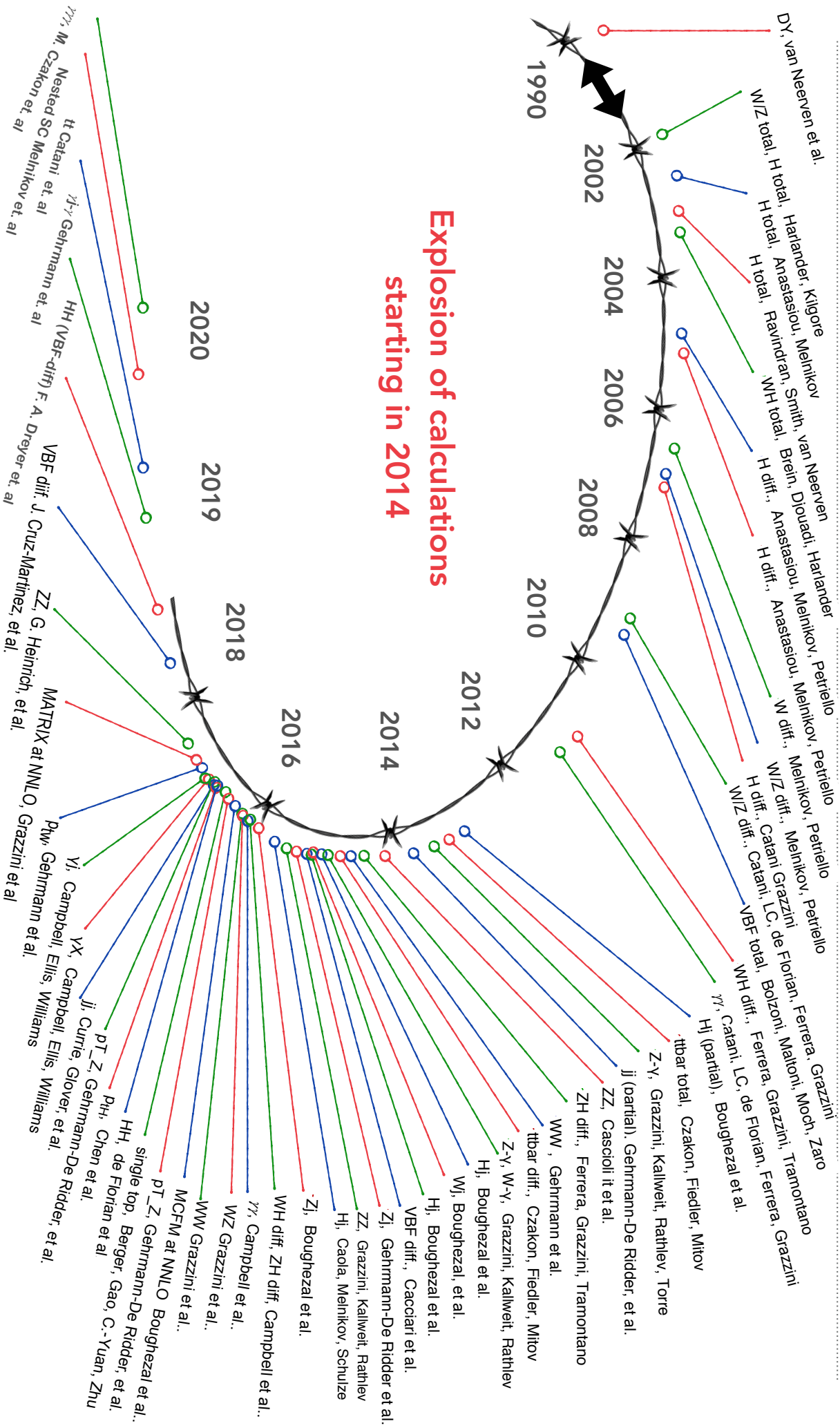
## NNLO HADRON-COLLIDER CALCULATIONS VS. TIME



Explosion of calculations starting in 2014

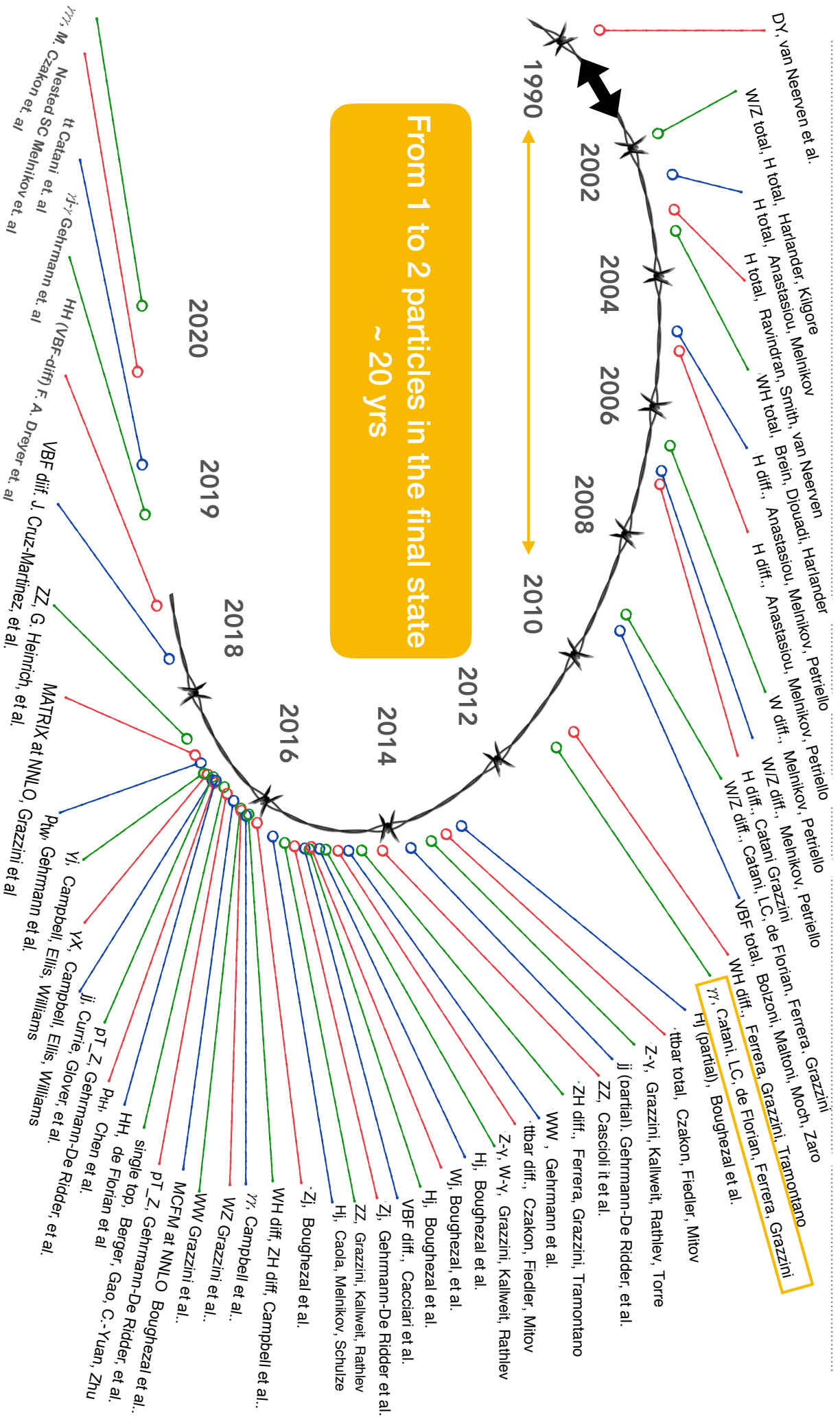
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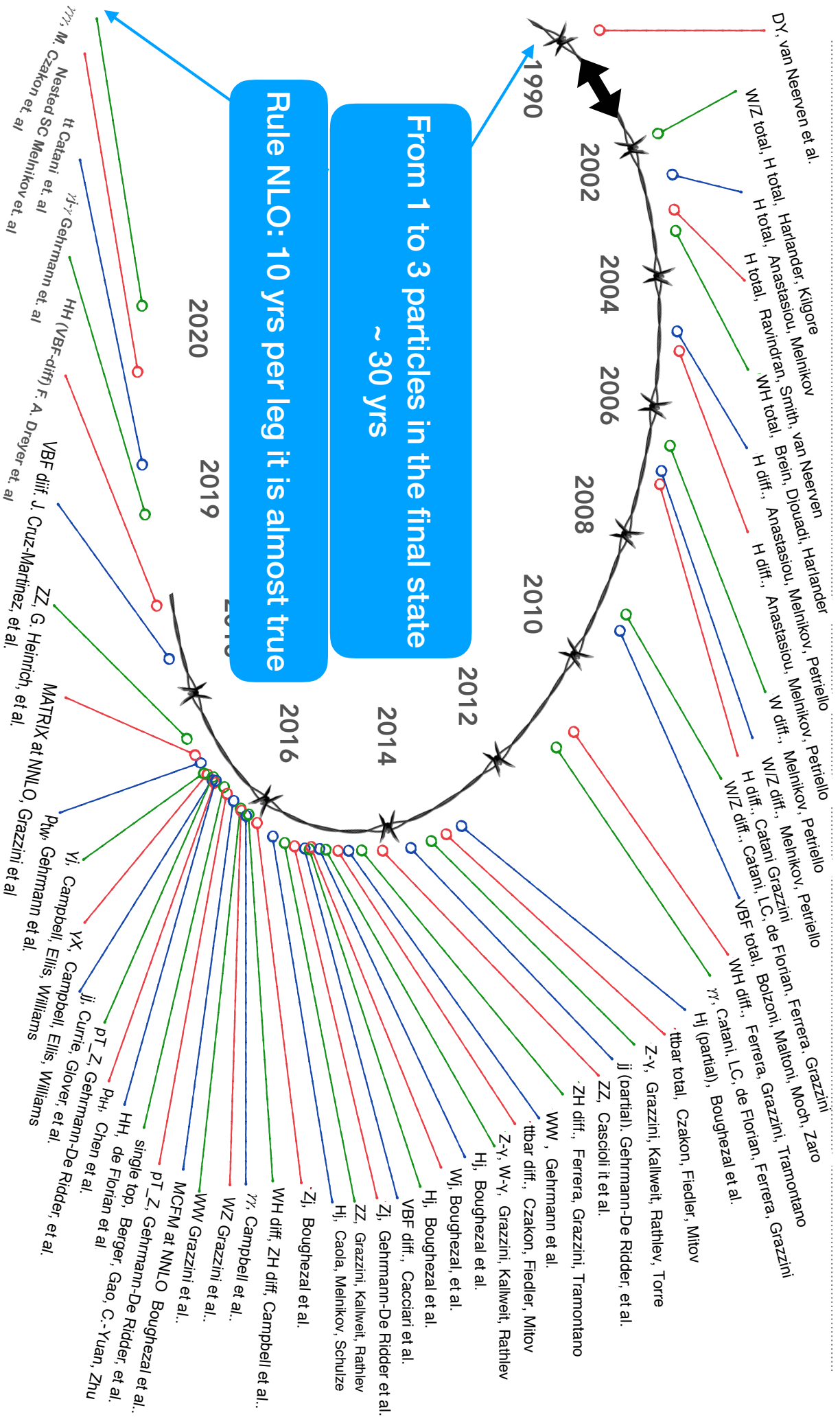
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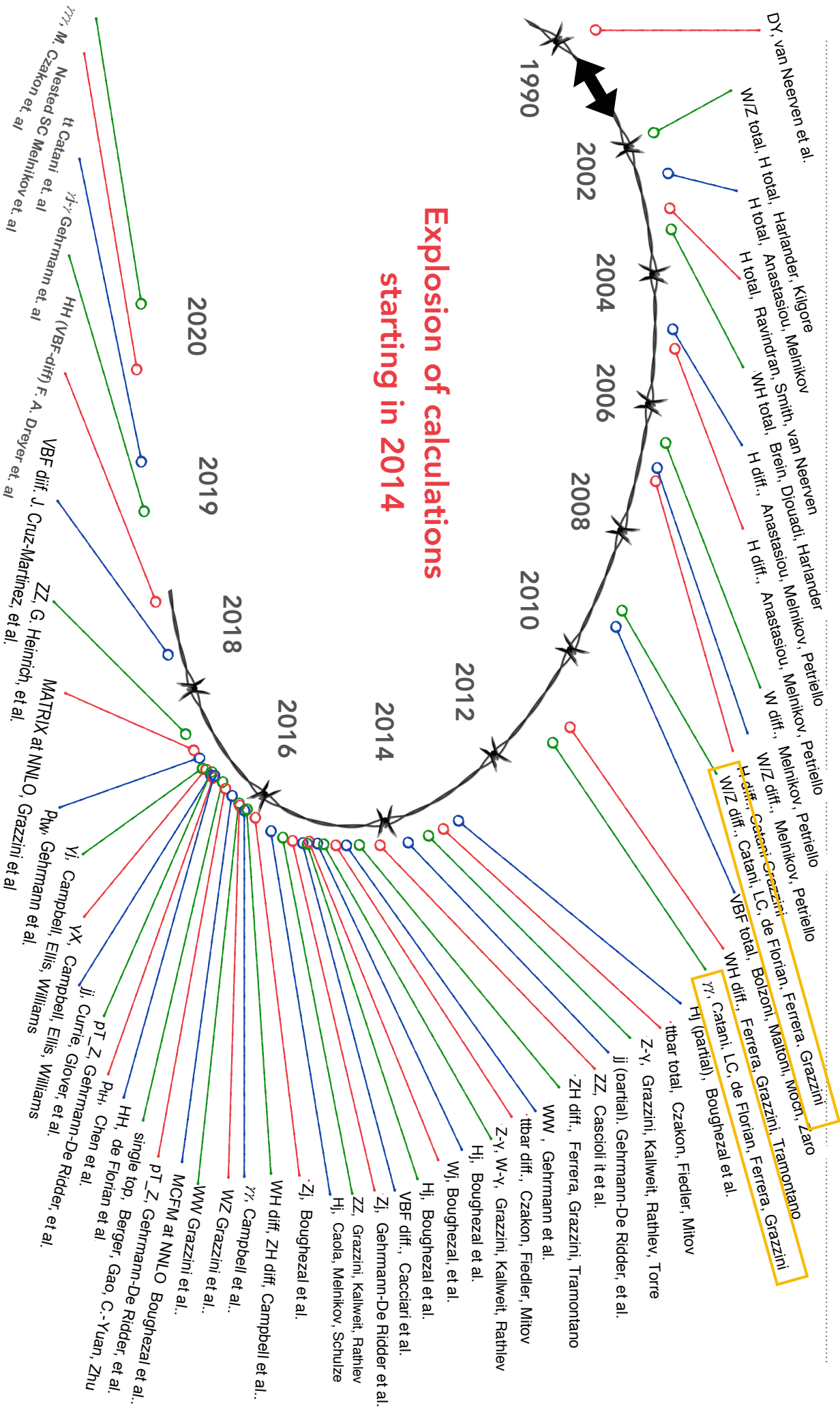
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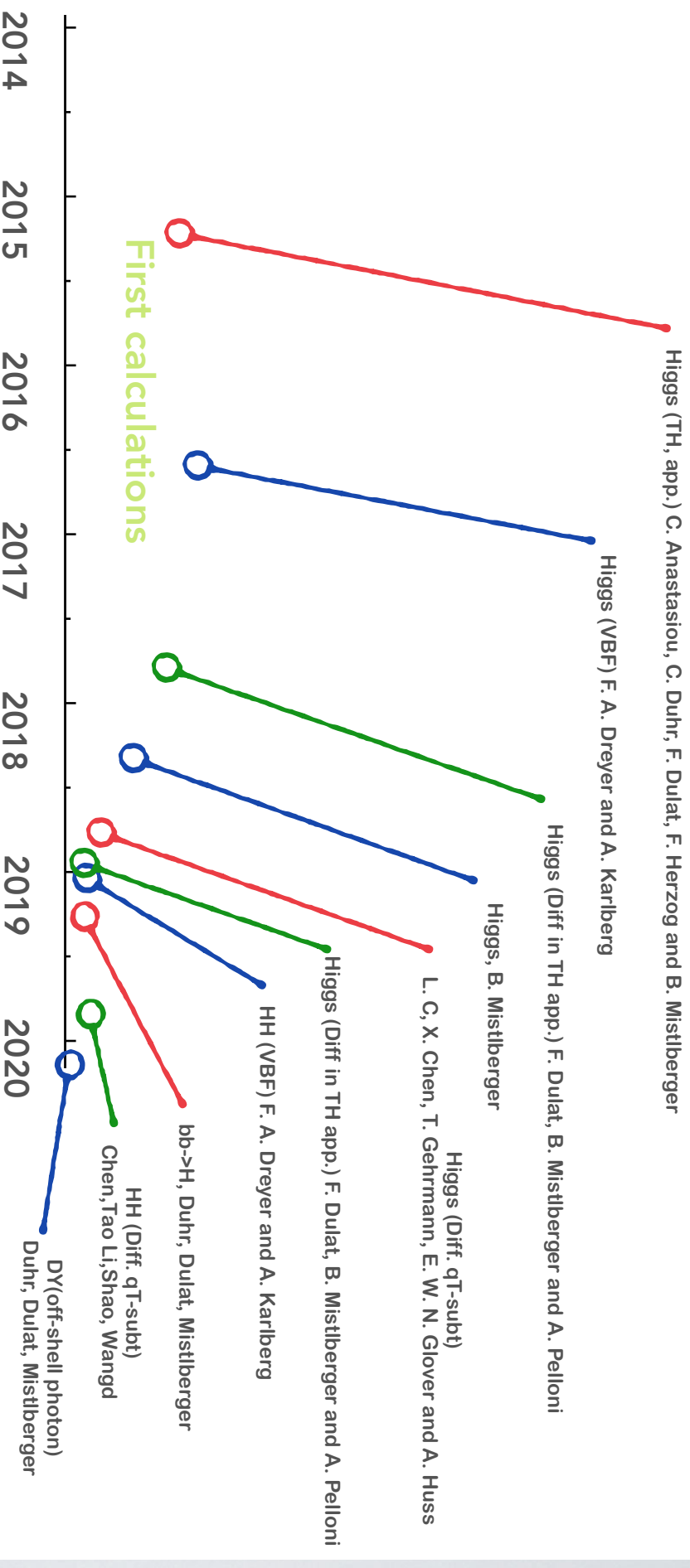
## NNLO HADRON-COLLIDER CALCULATIONS VS. TIME



# THE N<sup>3</sup>LO ERA

## N<sup>3</sup>LO HADRON-COLLIDER CALCULATIONS VS. TIME

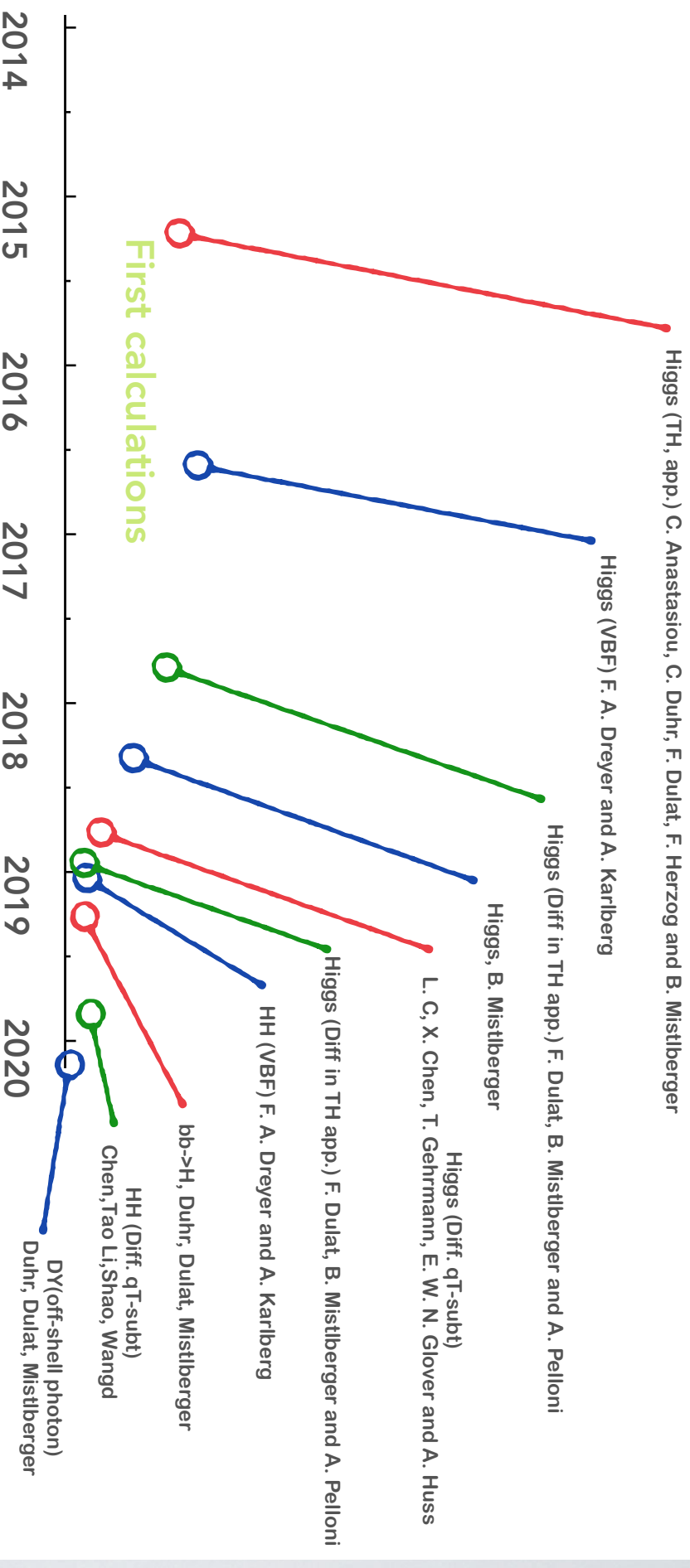
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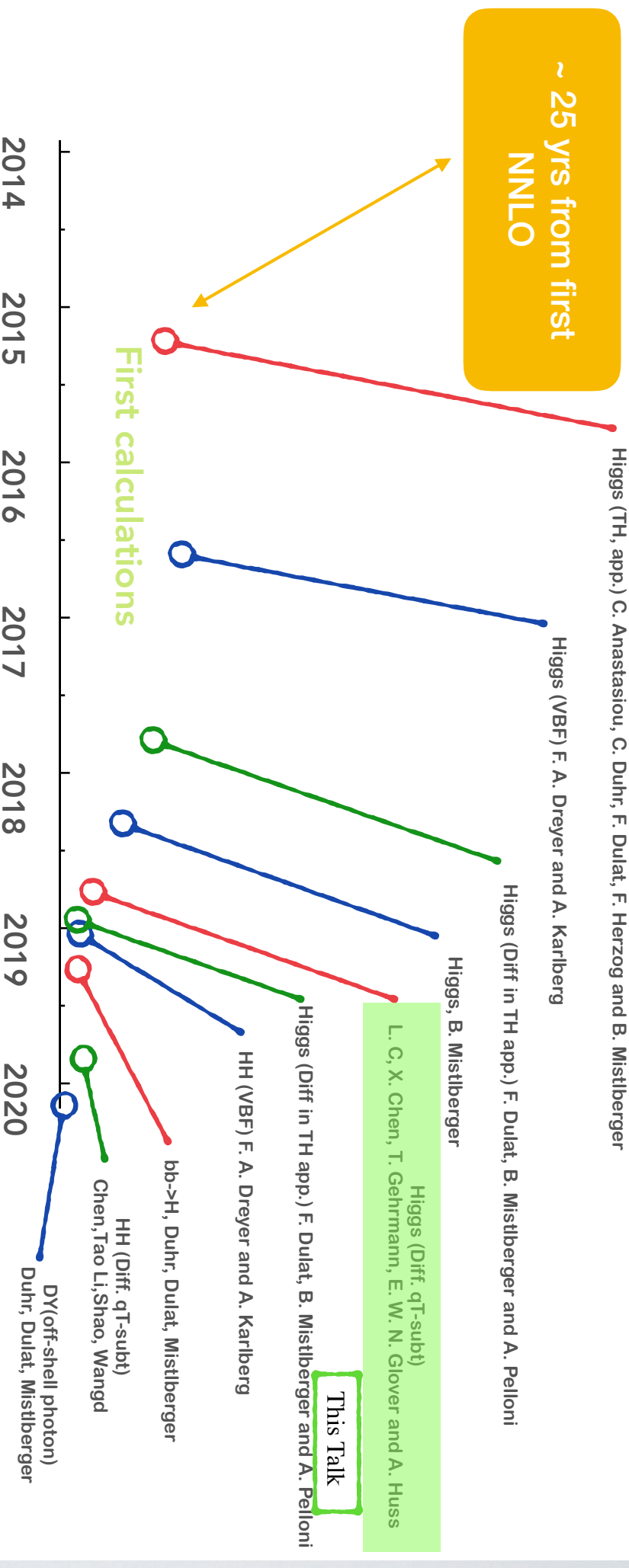
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This line starts to filling quickly

# THE N<sup>3</sup>LO ERA

## N<sup>3</sup>LO HADRON-COLLIDER CALCULATIONS VS. TIME



# THE QT-SUBTRACTION/RESUMMATION FORMALISM

## 📍 NNLO+NNLL QCD

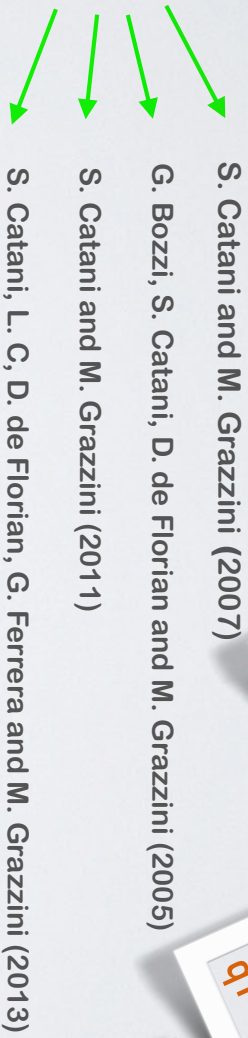
The method

Explicit form  
qT-subt. Ingredients

Helicity-Flip  
Contributions

Universal relation between  
hard-virtual factors and  
two-loop scattering amplitudes

### Colourless final states



### Heavy quark production



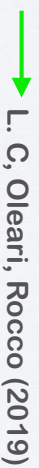
### Extension at NLO+NLL QED+QCD



### Extension at NLO QED FSR



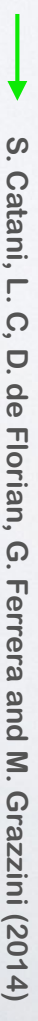
### Sub leading power corrections NLO



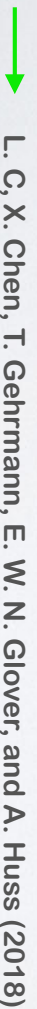
## 📍 N3LO+N3LL QCD

**Leonardo Vernazza's Talk**

### Soft virtual cross sections and threshold resummation



### N3LO QCD



# QT-SUBTRACTION AT M3LO

L. C. Chen, T. Gehrmann, E. W. N. Glover, and A. Huss (2018)

# N<sup>3</sup>L0 DIFFERENTIAL DISTRIBUTIONS

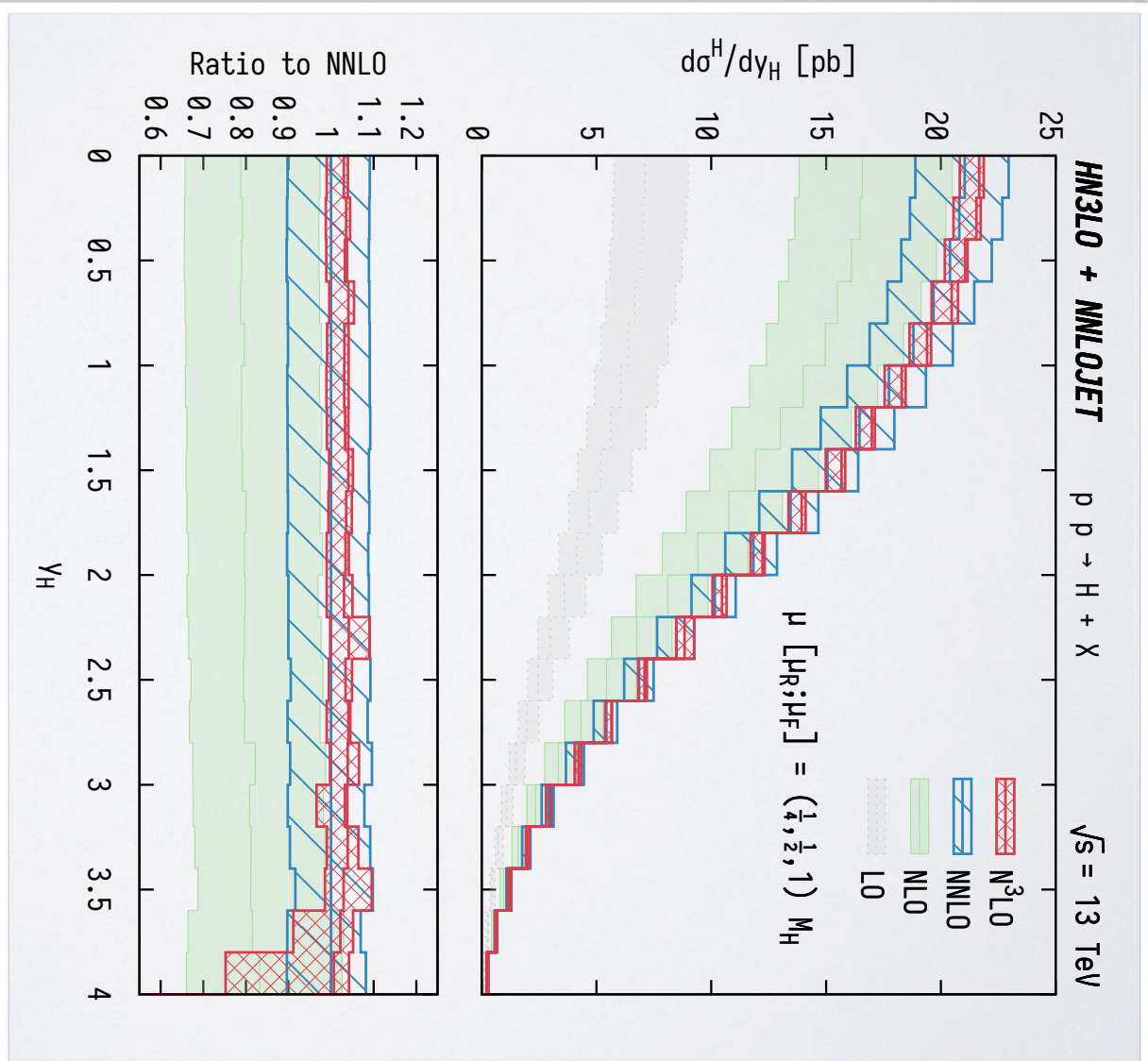
L. C. X. Chen, T. Gehrmann, E. W. N. Glover, and A. Huss (2018)

## Rapidity distribution

We calculate the  $y_H$  distribution at N<sup>3</sup>L0 employing a seven-point scale variation and carefully assess systematic errors arising from different  $q_T$  cut and CN3 values.

The combined theoretical uncertainty at N<sup>3</sup>L0 is at most of  $\pm 5\%$  level with respect to the central scale choice

N<sup>3</sup>L0 prediction at  $q_T$  cut=2GeV +  $q_T$  cut uncertainties + systematic uncertainties

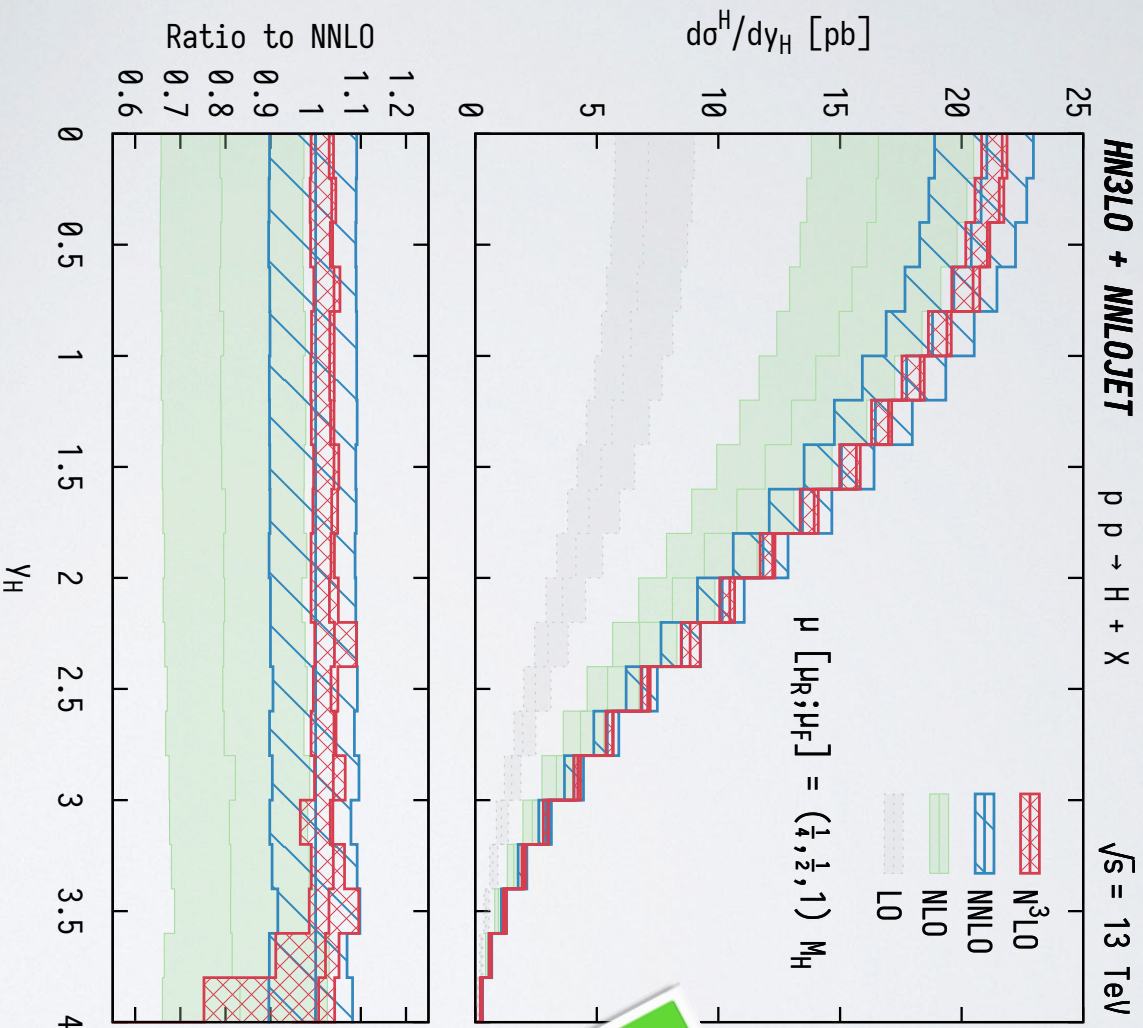


Compared to the NNLO  $y_H$  distributions, we observe a large reduction of theory uncertainties by more than 50% at N<sup>3</sup>L0. The scale variation band at N<sup>3</sup>L0 stays within the NNLO band with a flat K-factor of about 1.034 in the central rapidity region ( $|y_H| \leq 3.6$ ).

# N<sup>3</sup>LO DIFFERENTIAL DISTRIBUTIONS

L. C. X. Chen, T. Gehrmann, E. W. N. Glover, and A. Huss (2018)

## Rapidity distribution



In agreement with  
Dulat, Mistlberger and Pelloni (2018)  
[Higgs production in TH expansion  
approximation]  
Which is far from a trivial test!!



# N<sup>3</sup>LO DIFFERENTIAL DISTRIBUTIONS

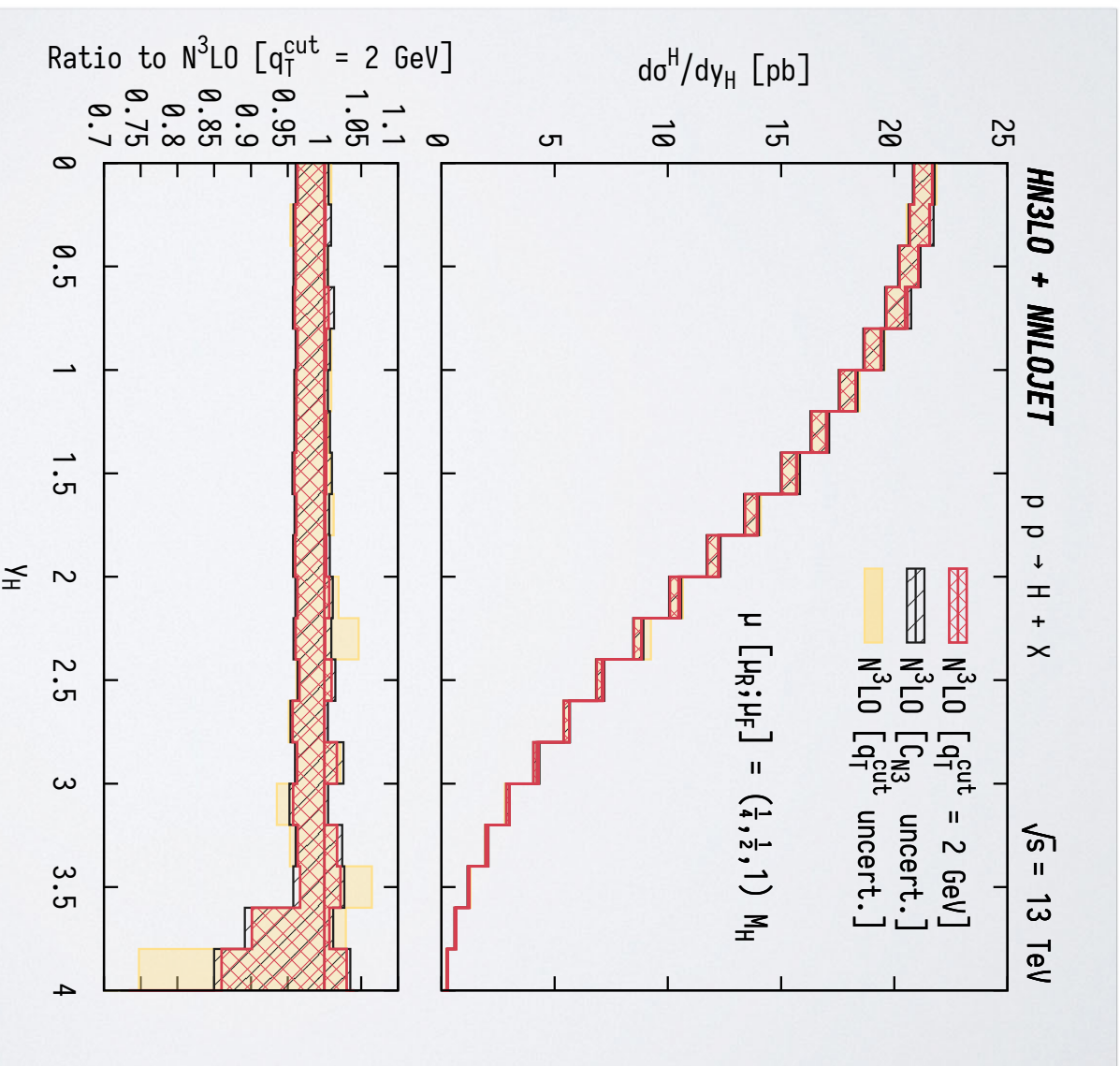
L. C. X. Chen, T. Gehrmann, E. W. N. Glover, and A. Huss (2018)

## Uncertainties at N<sup>3</sup>LO

N<sup>3</sup>LO prediction at  
 $q_T^{\text{cut}}=2\text{GeV}$   
 +  $q_T^{\text{cut}}$  uncertainties +  
 systematic uncertainties

$q_T^{\text{cut}}$  uncertainties  
 Calculated taking three  
 different predictions for  
 $q_T^{\text{cut}} = (2 \pm 1) \text{ GeV}$

The black band is  
 obtained as the envelope  
 of the seven-point scale  
 variation at  $q_T^{\text{cut}} = 2$   
 GeV considering for each  
 scale the two extremal  
 CN3 coefficients  
 corresponding to its  
 maximum and minimum  
 statistical deviations  
 $\text{CN3} = \{-1165, -721\}$



In order to produce  
 smooth histograms the  
 code has to run for a  
 couple of months in  
 2000 cores approx.

# N<sup>3</sup>LO DIFFERENTIAL DISTRIBUTIONS

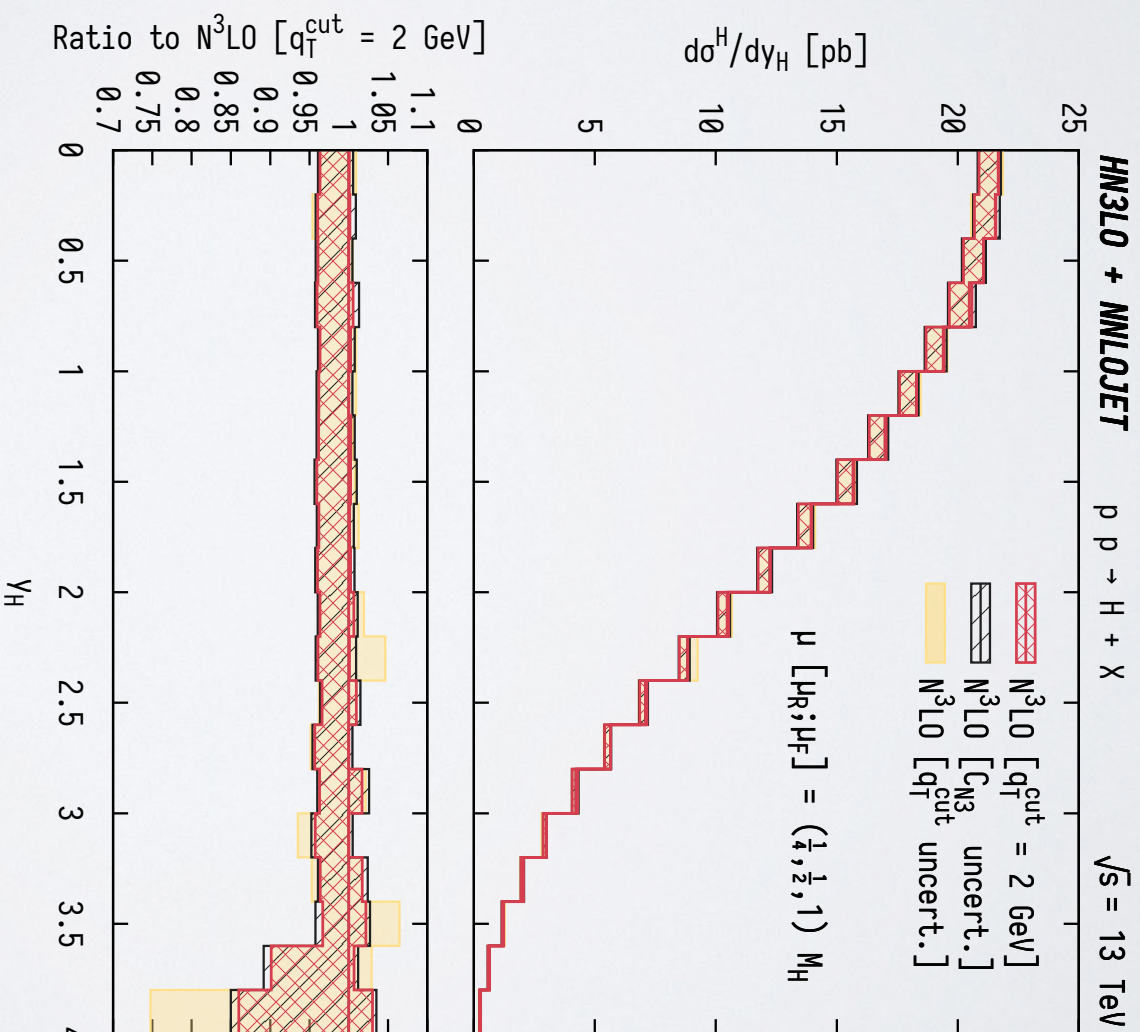
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For the  
 inclusive: ~2  
 min in one  
 core!

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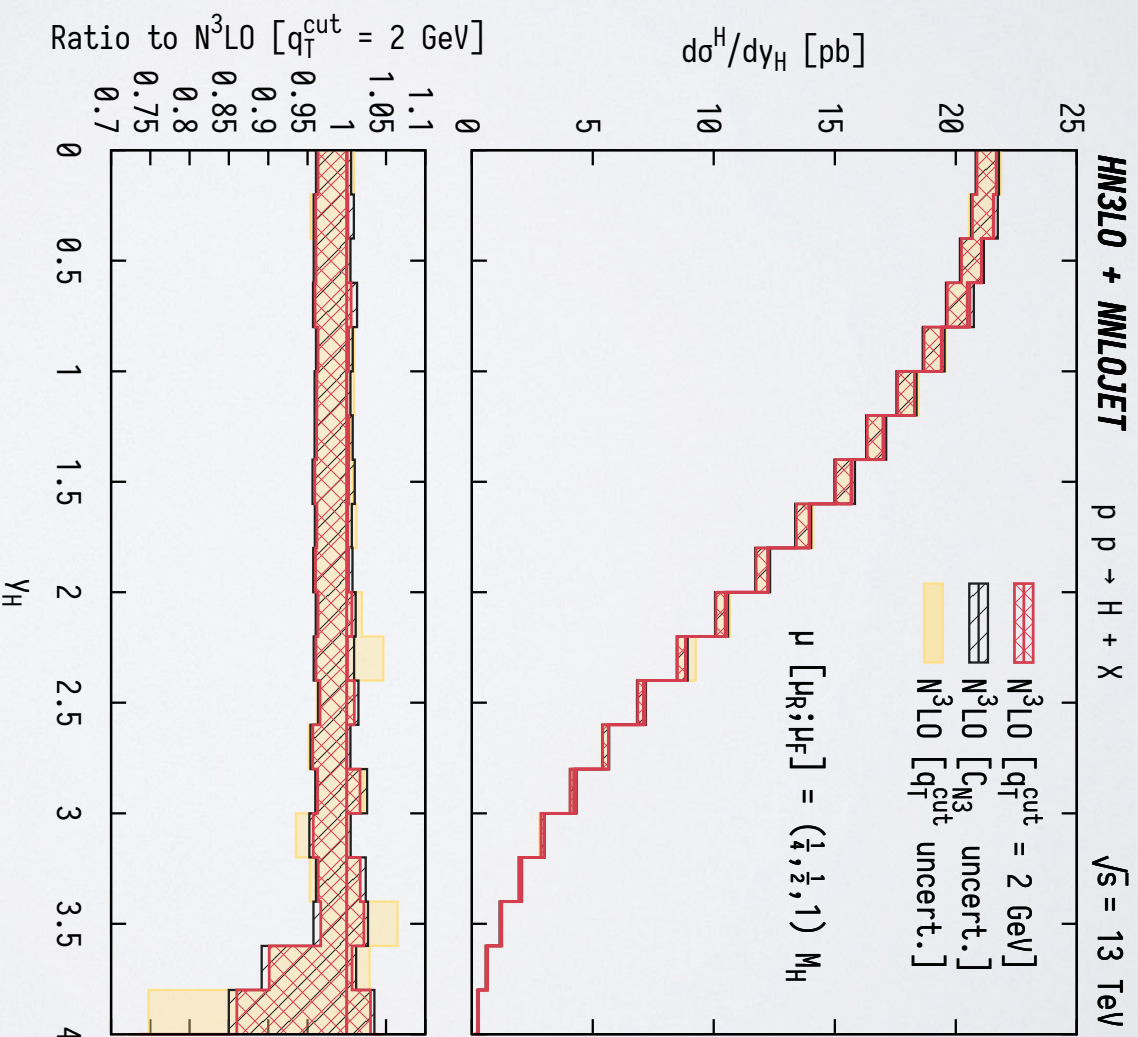
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For the  
 inclusive: ~2  
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 smooth histograms the  
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 couple of months in  
 2000 cores approx.

~ 100 yrs in a  
 desktop computer

We need new  
 techniques, new  
 approaches

# N3LO AS THE NEW STANDARD FOR PRECISION PHYSICS AT THE LHC

Next steps following the research plan

---

Drell-Yan at N3LO differential

All the qT-subtraction ingredients are known for qqbar  
M. Luo, T. Yang, H. Xing Zhu, Y. Jiao Zhu (2019)

# N3LO AS THE NEW STANDARD FOR PRECISION PHYSICS AT THE LHC

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Total cross section not known yet

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qT-resummation at the same  
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DYTurbo  
Camarda, Bozzi, Catani, L. C, Ferrera, et. al (2019)

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Sub-leading power corrections

L. C, Oleari, Rocco (2019)

They are not known even at NNLO

Leonardo Vernazza's Talk

Hints on local subtraction

Setup general framework (analytical/numerical) to apply to any colourless final state



# CONCLUSIONS AND OUTLOOK

- We presented results at N<sup>3</sup>LO for completely differential distributions for observables measured at the LHC  
(in particular for Higgs boson production in gluon fusion)
- The qT-subtraction method at N<sup>3</sup>LO, presented here, can be applied to other processes initiated by gluon fusion (if the three-loop scattering amplitudes are known)
- In the future, two elements are necessary in order to complete analytically the qT-subtraction method at N<sup>3</sup>LO
  - The universal relation between the three-loop virtual scattering amplitudes and the hard virtual factor
  - The third order collinear functions (gluon-gluon channel)
- Our results can be applied to perform qT-resummation at the corresponding level of logarithmic accuracy. **i.e without reweighing by the total cross section**
- We have all the elements in the case of the qqbar channel

# CONCLUSIONS AND OUTLOOK

- **The completion of this research project will produce**
  - new complementary use of different techniques, analytic and semi-analytic techniques; new ideas related to the introduction of more complex mathematical structures related to the calculation of Feynman diagrams
  - new subtraction scheme developed
  - new developed Monte Carlo code at N<sup>3</sup>LO (one of the few existing)
  - innovative studies of processes that test the SM with very high precision. Such high-precision can eventually give hints about NP contributions

# CONCLUSIONS AND OUTLOOK

• I expect that our fellowship serve us to produce “state of the art computations” that can be used to compare with the most precise LHC data

• I particularly expect that our fellowship, and the work that we will perform along our research plans, will benefit us at the time of application to other prestigious fellowships and fixed positions

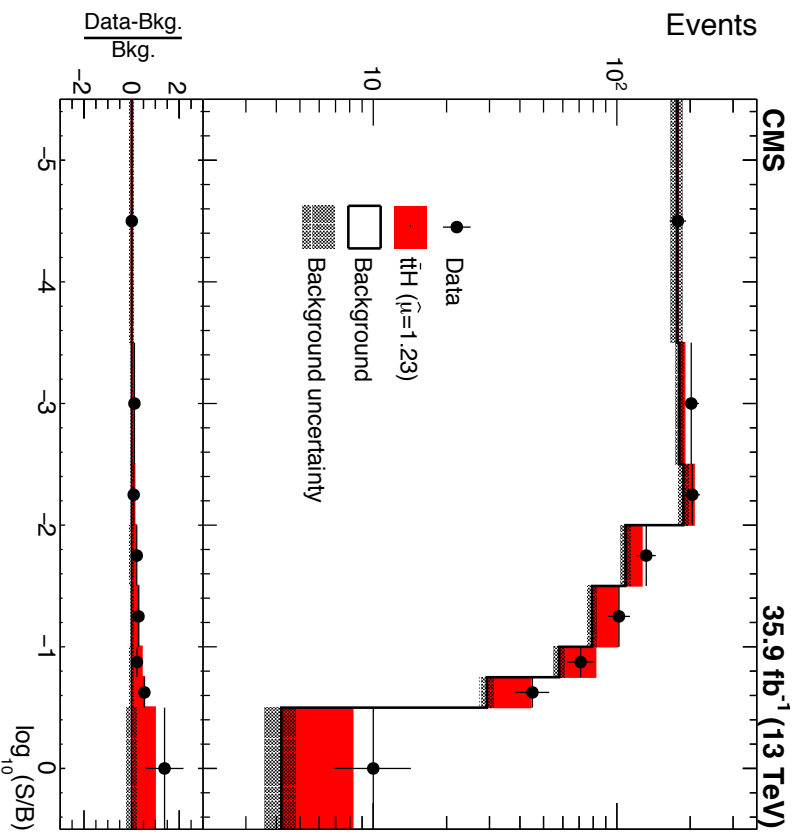
**THANK YOU!**

# BACKUP SLIDES

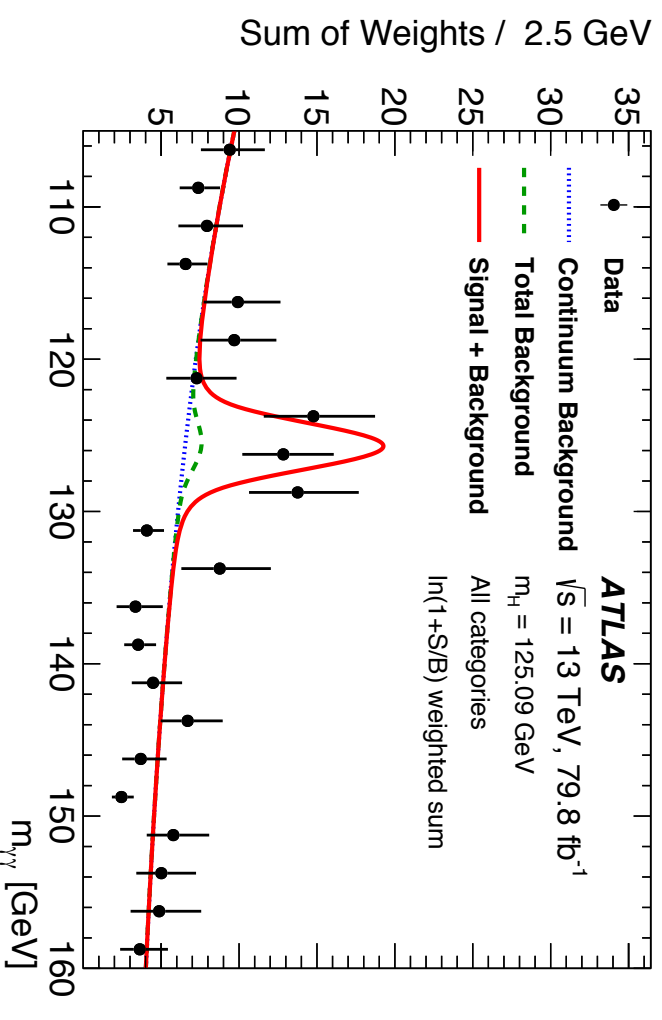
# New interesting tasks are just recently initiated: Yukawas

the news of the past 12 months

A few weeks ago:  
CMS >5-sigma ttH



This week:  
ATLAS >5-sigma ttH



$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{\partial}\psi + \bar{\psi} \gamma_5 \psi \phi + h.c. + |D_\mu\phi|^2 - V(\phi)$$

This equation neatly sums up our current understanding of fundamental particles and forces.

Why do Yukawa couplings matter?  
 (2) Because, within SM **conjecture**, they're what give masses to all **quarks**

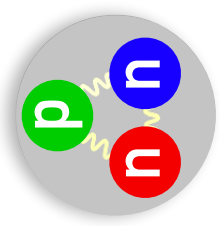
Up quarks (mass  $\sim 2.2$  MeV) are lighter than  
 down quarks (mass  $\sim 4.7$  MeV)

**proton** (up+up+down):  $2.2 + 2.2 + 4.7 + \dots = 938.3$  MeV  
**neutron** (up+down+down):  $2.2 + 4.7 + 4.7 + \dots = 939.6$  MeV

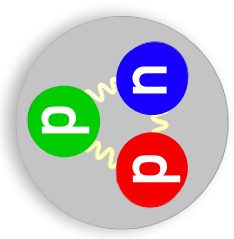
So protons are **lighter** than neutrons,  
 → protons are stable.

Which gives us the hydrogen atom,  
 & chemistry and biology as we know it

*proton*  
 mass = 938.3MeV

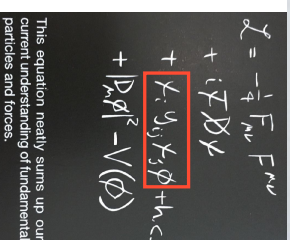


*neutron*  
 mass = 939.6MeV



## Why do Yukawa couplings matter?

(3) Because, within SM **conjecture**, they're what give masses to all **leptons**



This equation neatly sums up our current understanding of fundamental particles and forces.

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{\partial}\psi + \bar{\psi} \gamma_5 \psi \phi + h.c. + \bar{\psi}\not{\partial}^2\psi - V(\phi)$$

**Bohr radius**

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = \frac{\hbar}{m_e c \alpha}$$

electron mass determines size of all atoms

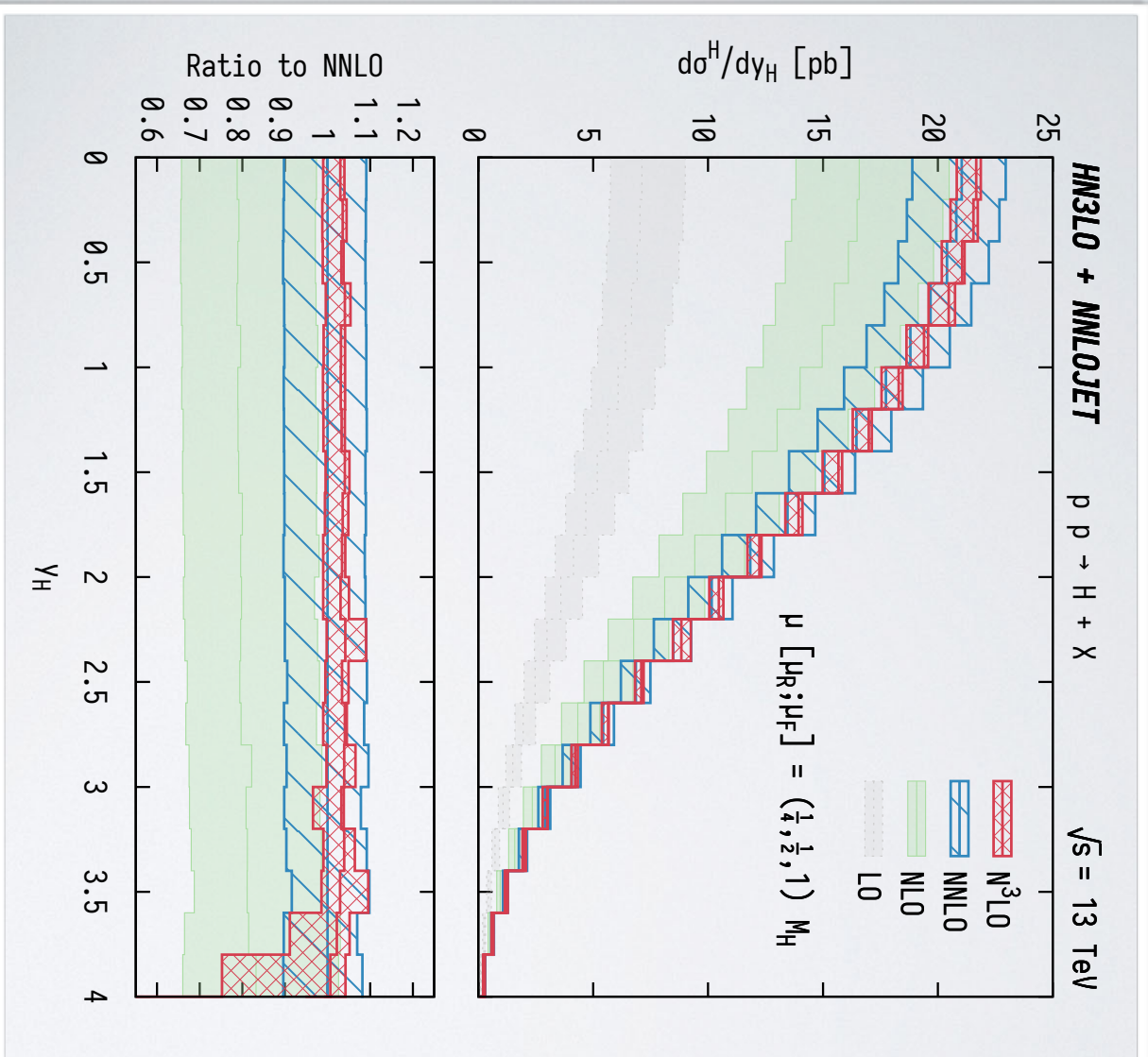
it sets energy levels of all chemical reactions



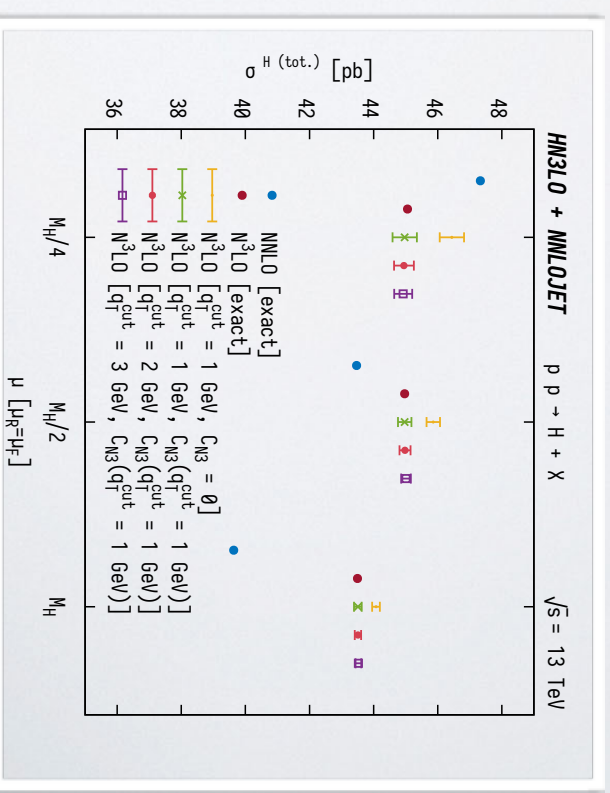
# N<sup>3</sup>LO DIFFERENTIAL DISTRIBUTIONS

L. C. X. Chen, T. Gehrmann, E. W. N. Glover, and A. Huss (2018)

## Rapidity distribution



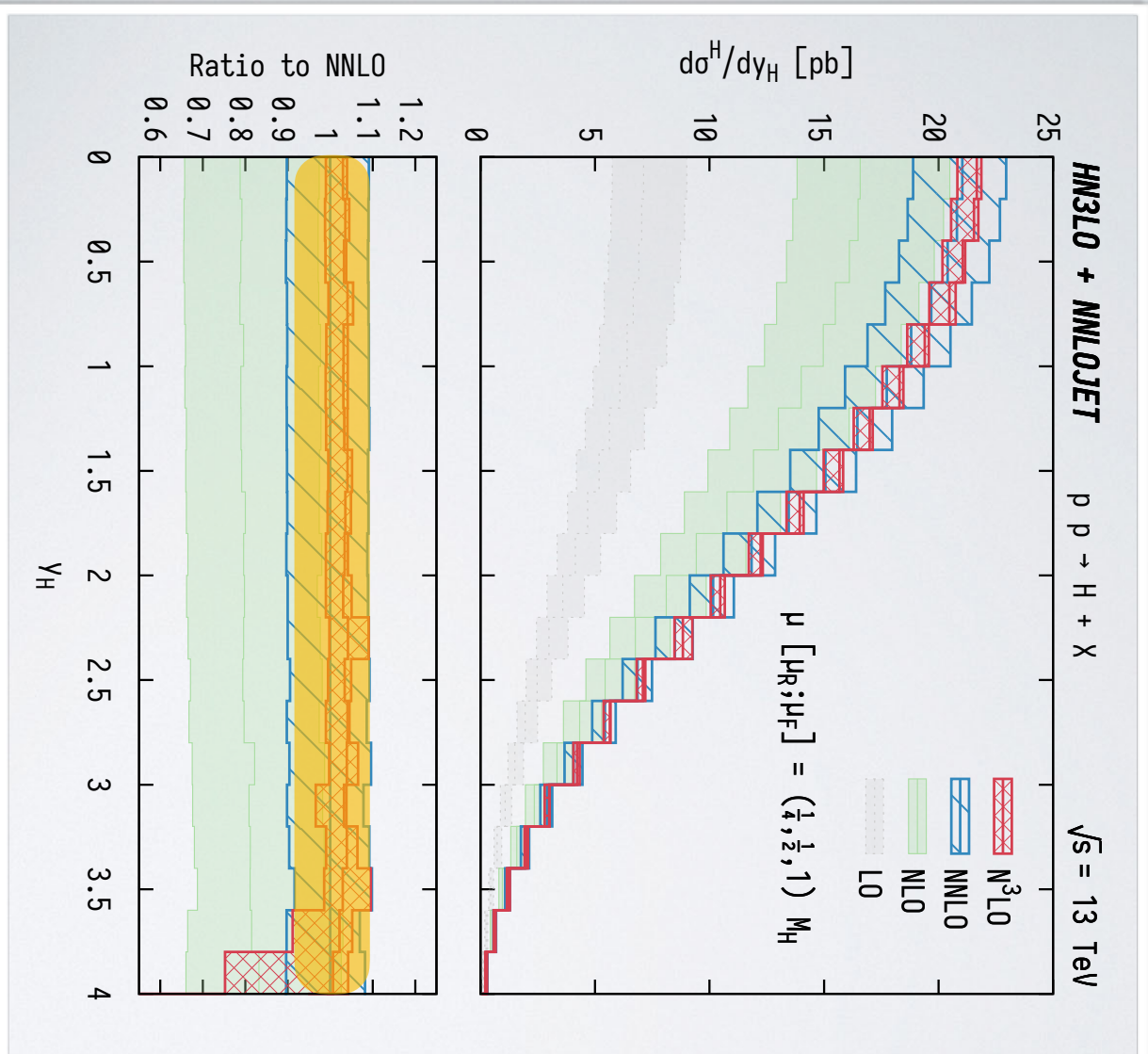
The central prediction at N<sup>3</sup>LO, almost coincides with the upper edge of the band, as was already observed for the total cross section



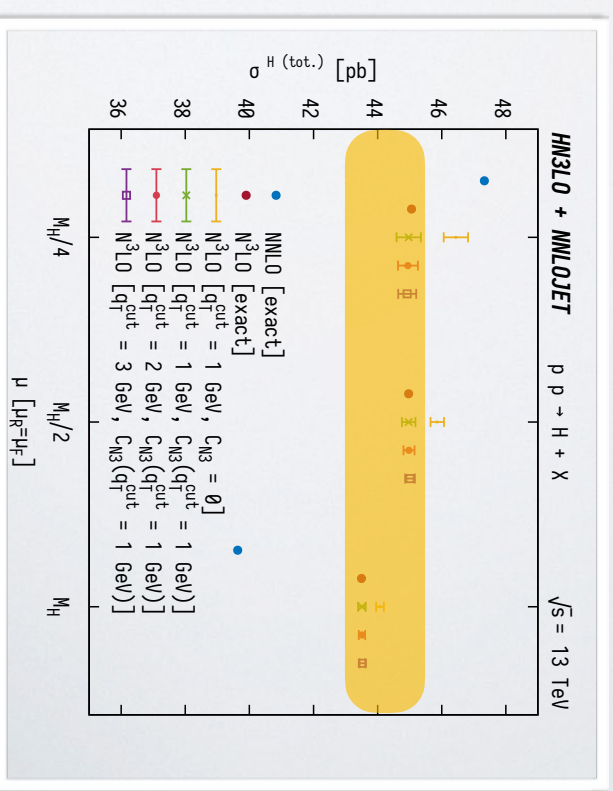
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## Rapidity distribution



The central prediction at N<sup>3</sup>LO, almost coincides with the upper edge of the band, as was already observed for the total cross section



# MOTIVATION

Why precision physics is important at the LHC?

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F_{\mu\nu} + i\bar{\psi} D\psi + Y_i \bar{\psi}_i \chi_j \phi + h.c. + |D_\mu \phi|^2 - V(\phi)$$

## Yukawa couplings

.....  
until 6 years ago this was essentially conjecture

no such term had ever been seen in nature  
hadn't even been probed in electroweak precision tests

Gavin Salam's talk LHCP (2018)

🔗 New interesting tasks are just recently initiated: Yukawas

# MOTIVATION

Why precision physics is important at the LHC?

Yukawa interactions are important not merely because they had never before been directly observed, but also because they are hypothesized to be responsible for the stability of hydrogen, and for determining the size of atoms and the energy scales of chemical reactions.

$$\mathcal{L} = \dots + i\bar{\chi}\chi + \chi_i y_{ij} \chi_j \phi + h.c. + |\partial_\mu \phi|^2 - V(\phi)$$

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