



THE HENRYK NIEWODNICZAŃSKI
INSTITUTE OF NUCLEAR PHYSICS
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“Spin” in relativistic heavy-ion collisions

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Based on: arXiv:2103.02592, arXiv:2011.14907 and arXiv:1901.09655.

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Quantum Chromodynamics (QCD) pushed to extreme

Early Universe

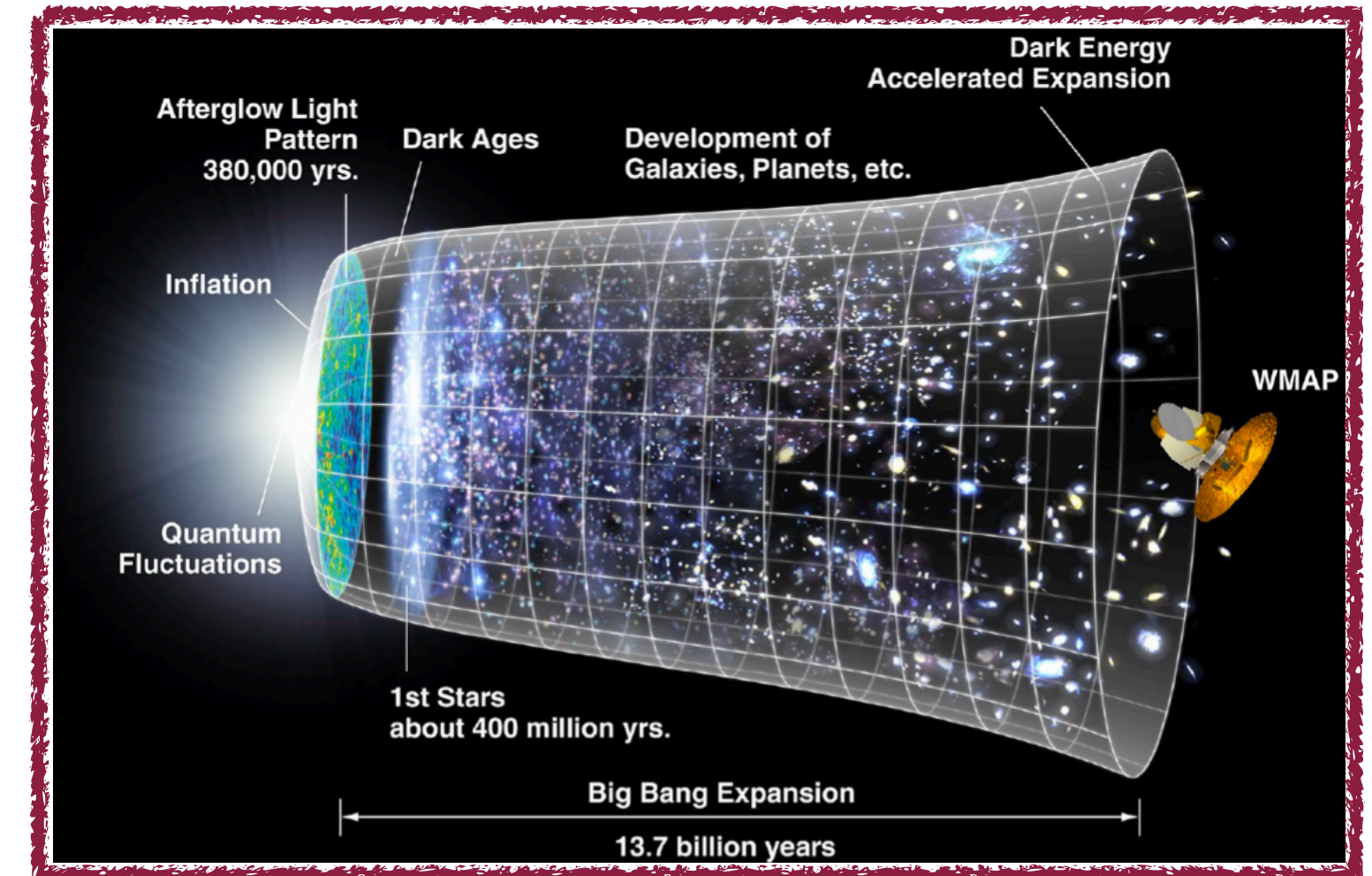


figure: NASA

Cores of neutron stars

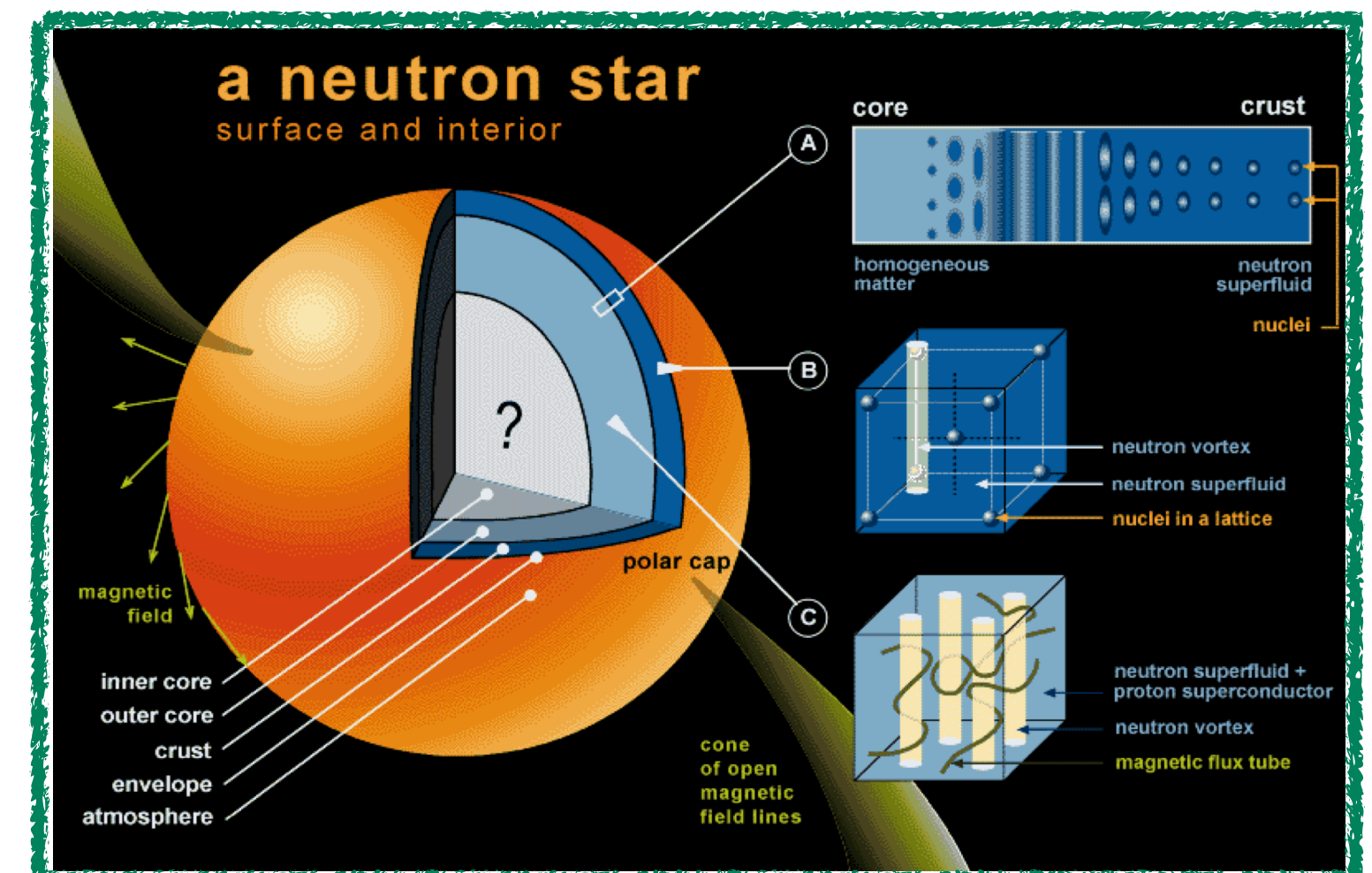
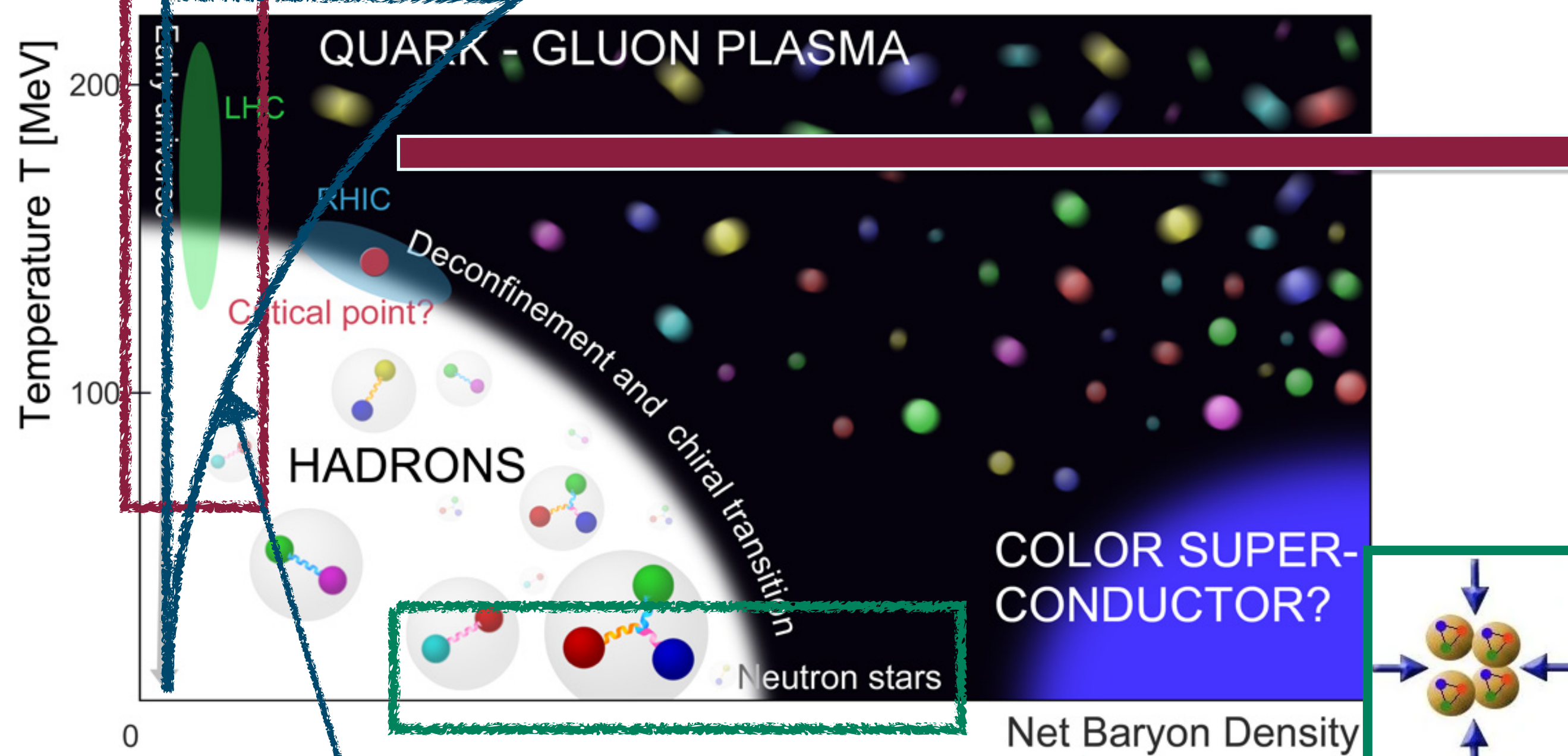


figure: D.E. A. Castillo, talk @RagTime 22

QCD phase diagram

figure: Joint Institute for Computational Fundamental Science



Lattice-QCD simulations

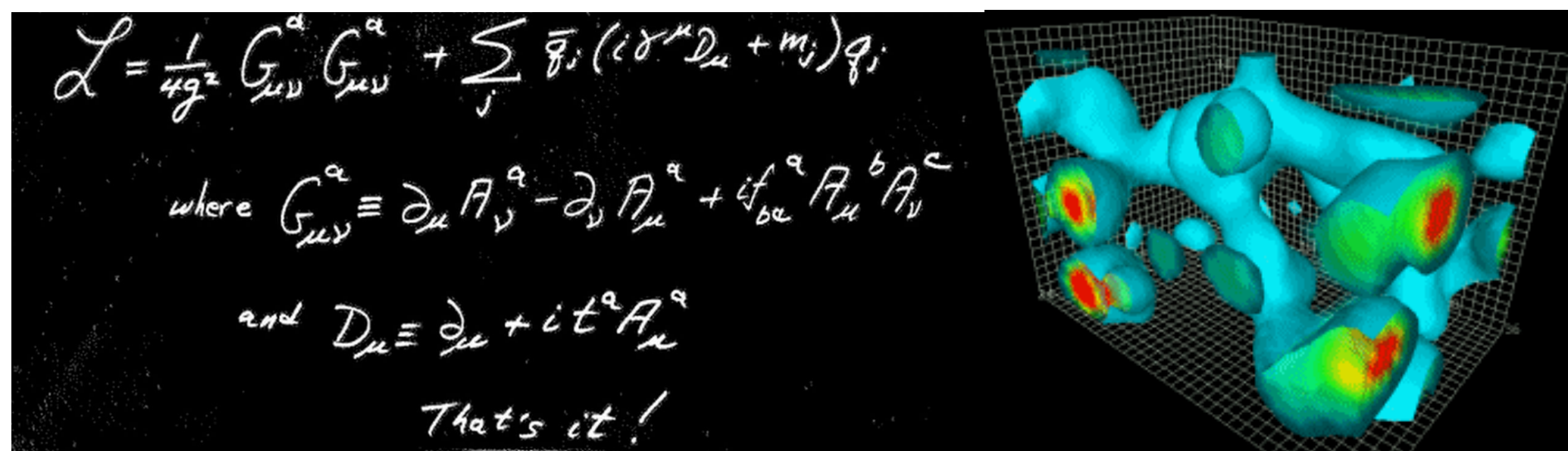
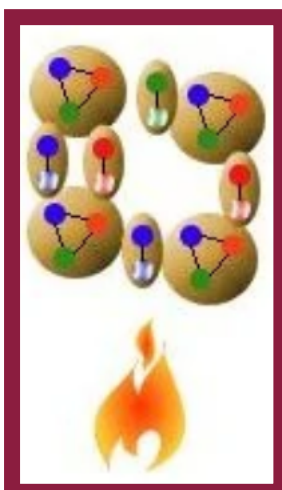
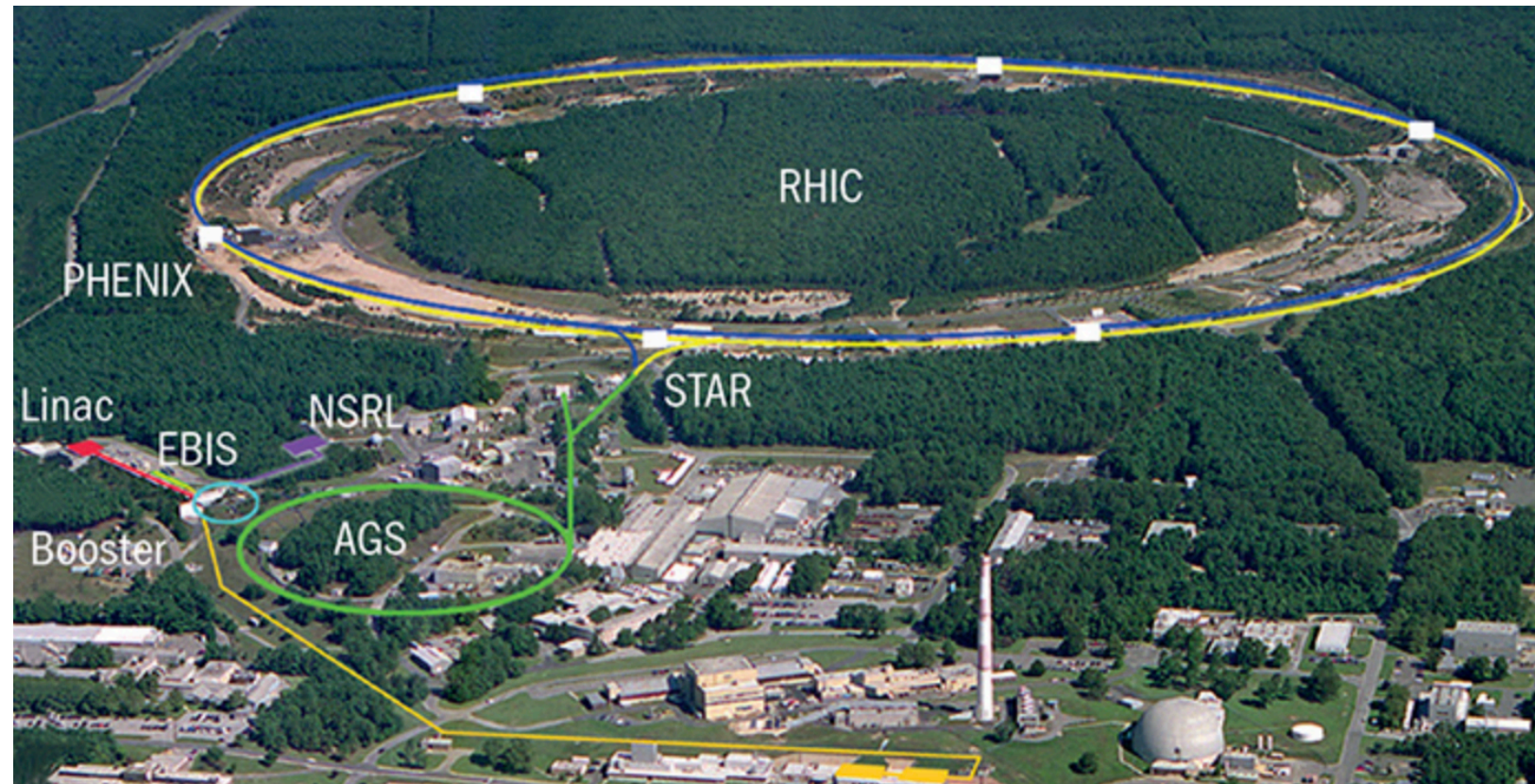


figure: D. Leinweber (www.physics.adelaide.edu.au) 2

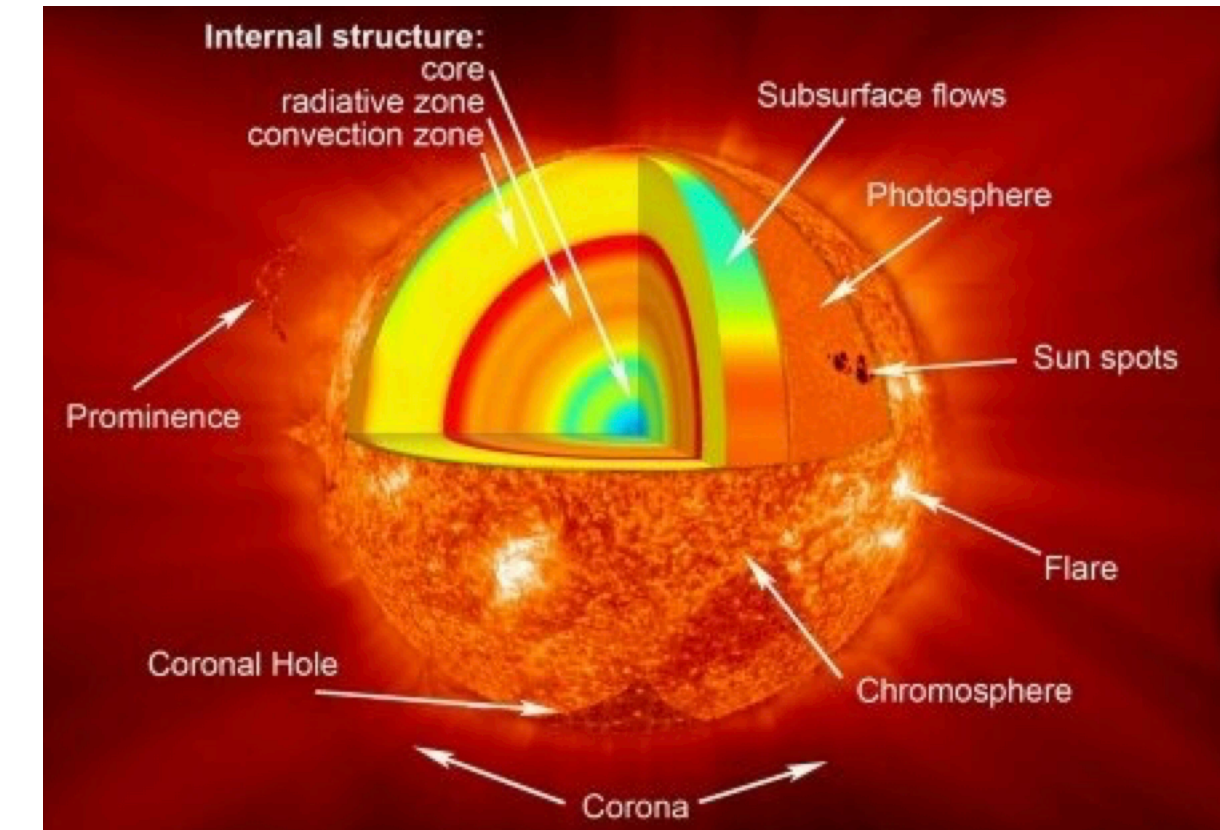


Relativistic heavy-ion collisions - a tool to study QGP

figure: NASA

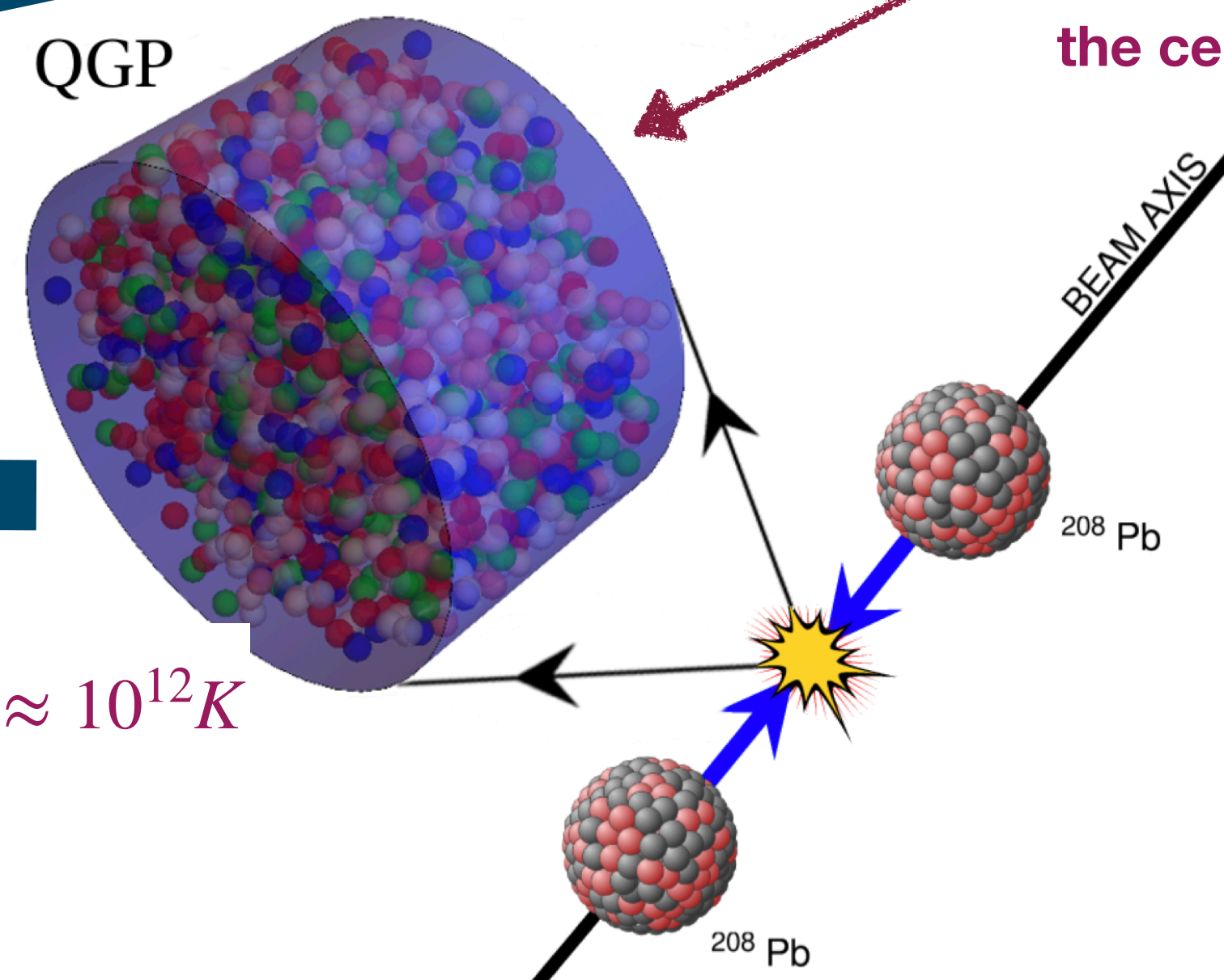
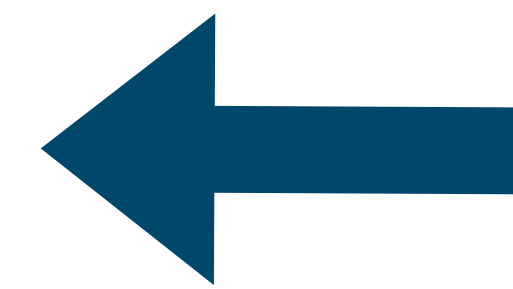
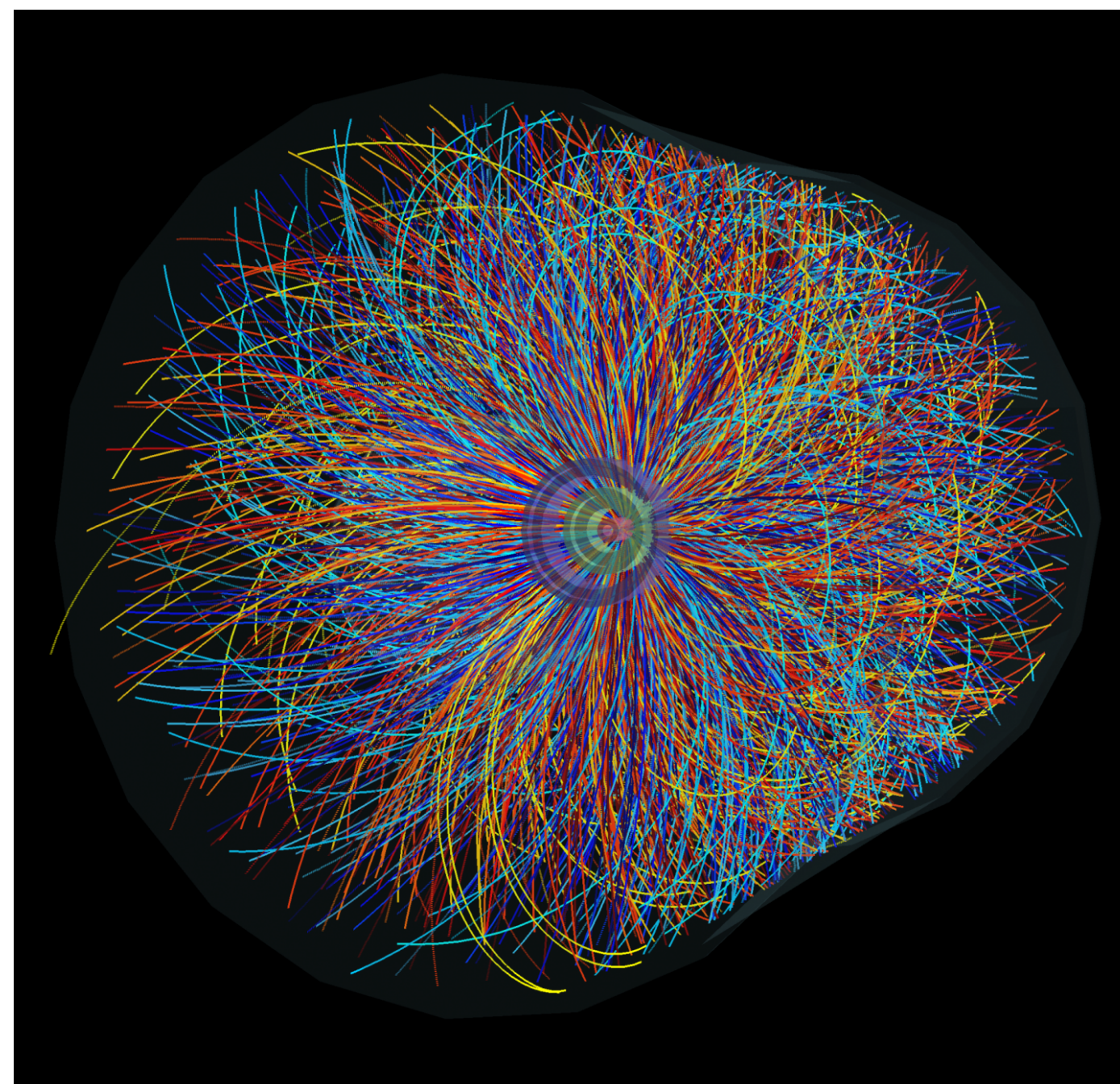


At high beam energies we observe many particles being produced



100 000 times hotter than the center of the Sun

The study of QGP possible only indirectly through the energy and momenta of emitted particles



Spin polarization in heavy-ion collisions - new sensitive probe!

Non-central heavy-ion collisions create fireballs with large global orbital angular momenta

F. Becattini, F. Piccinini, J. Rizzo, PRC 77 (2008) 024906

$$\mathbf{L}_{\text{init}} \sim 10^5 \hbar$$

Part of the angular momentum can be transferred from the orbital to the spin part

$$\mathbf{J}_{\text{init}} = \mathbf{L}_{\text{init}} = \mathbf{L}_{\text{final}} + \mathbf{S}_{\text{final}}$$

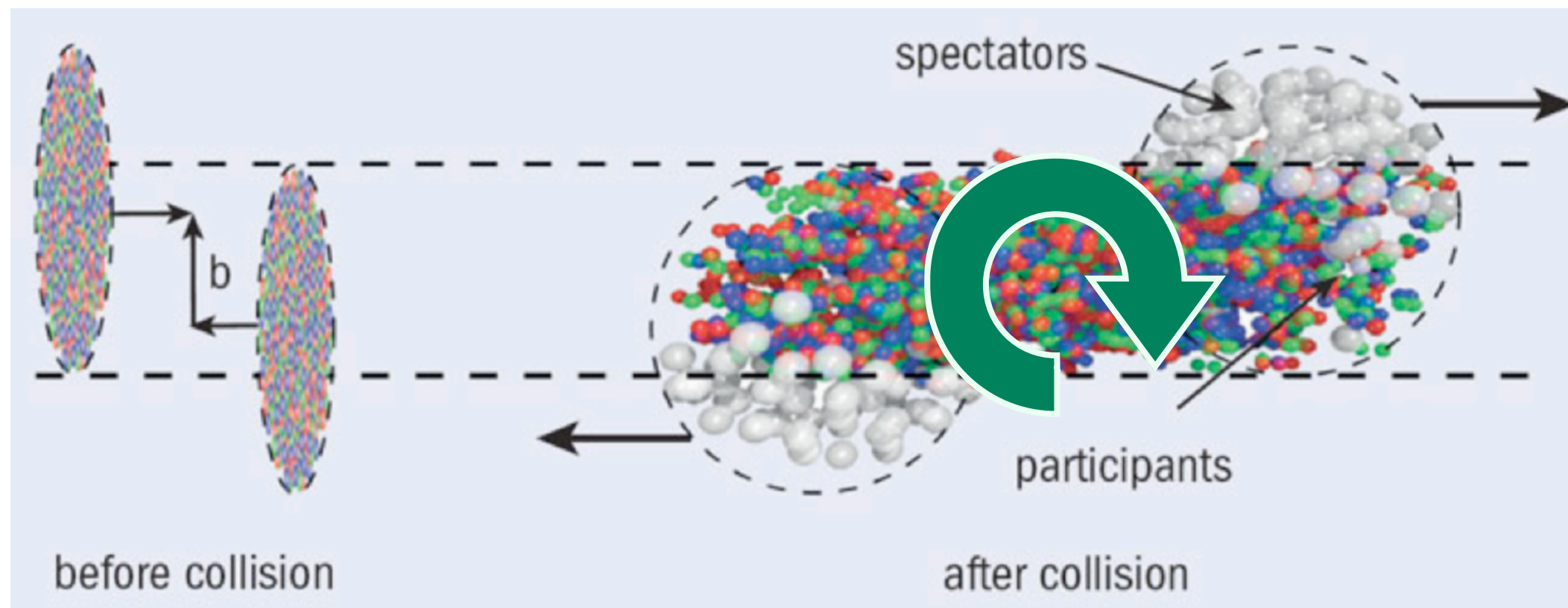


figure: M. Lisa, talk @ "Strangeness in Quark Matter 2016"

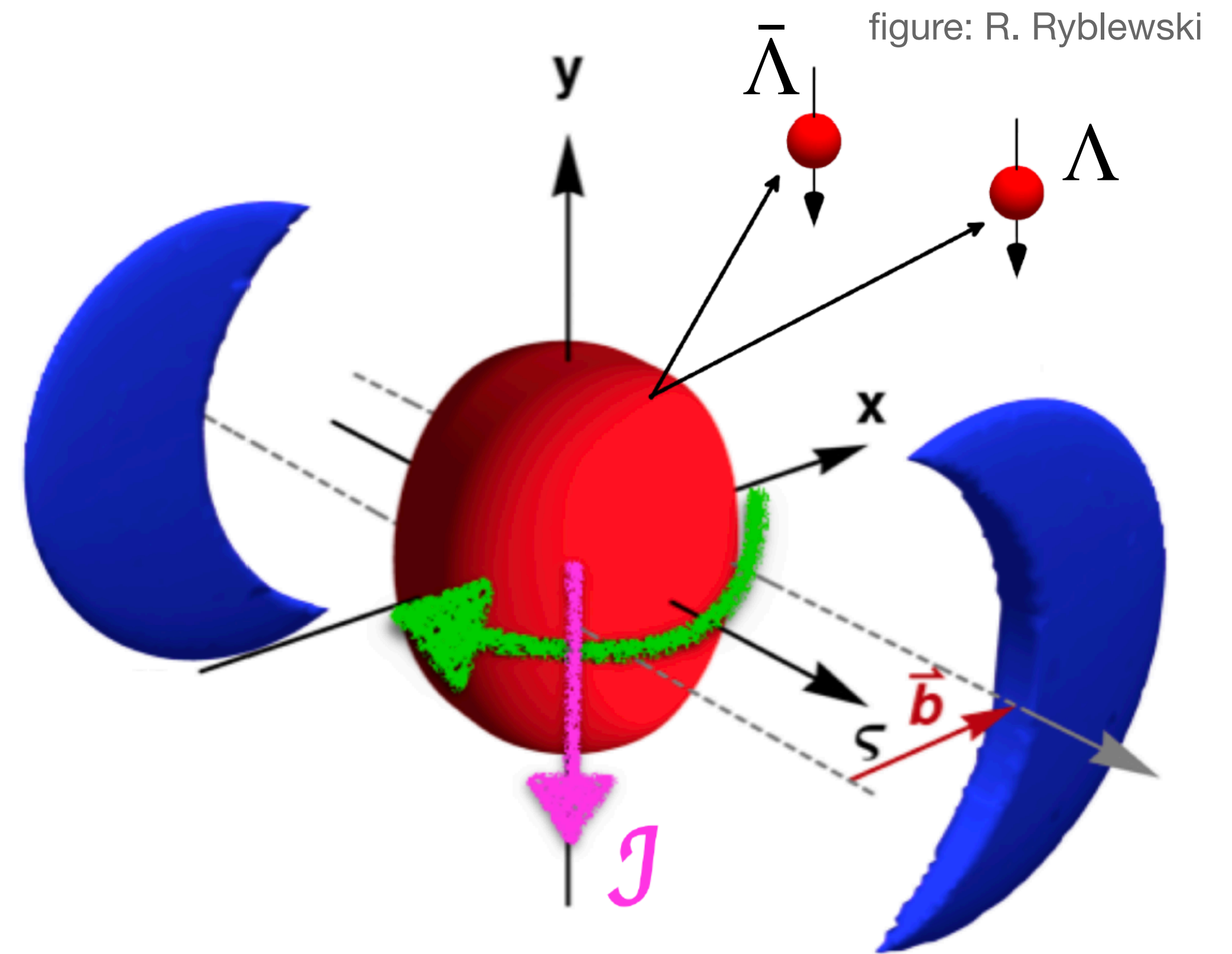
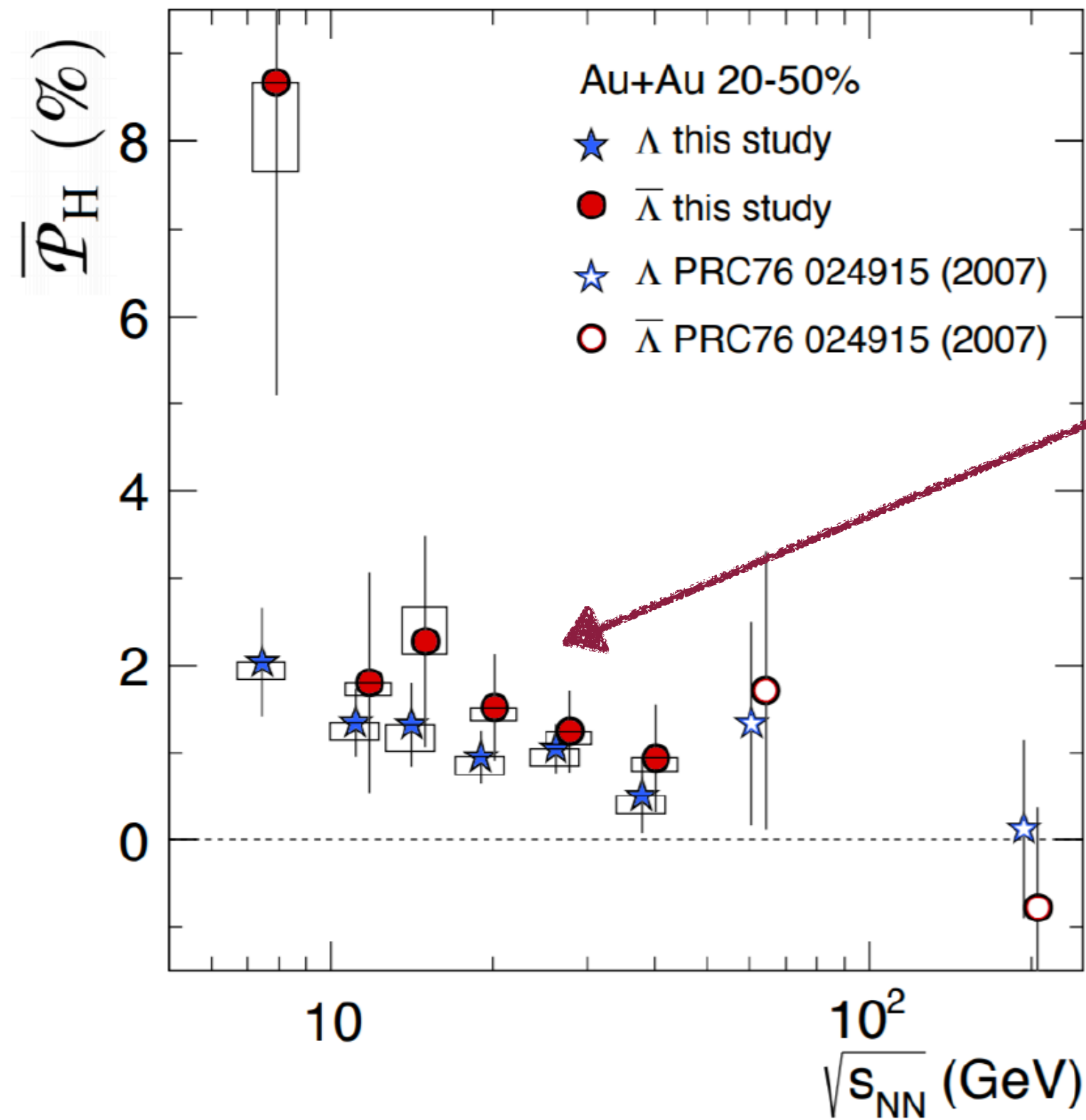


figure: R. Ryblewski

Emitted particles are expected to be globally polarized along the system's angular momentum

Measurement of Λ and $\bar{\Lambda}$ spin polarization in heavy-ion collisions

L. Adamczyk et al. (STAR) (2017), Nature 548 (2017) 62-65



~2% - small but measurable effect

Self-analysing parity-violating hyperon weak decay allows to measure polarization of Λ

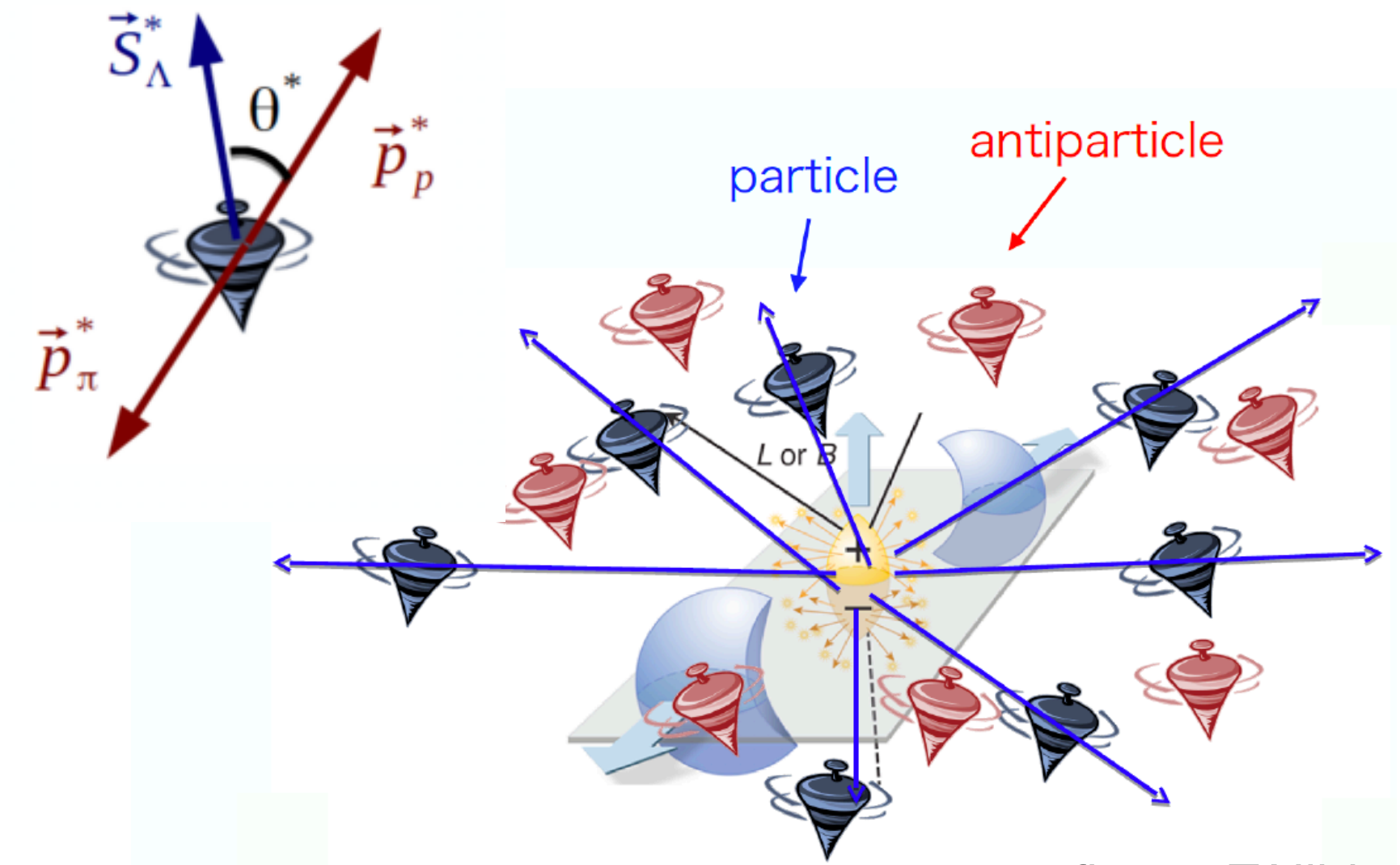


figure: T.Niida

... the hottest, least viscous – and now, most vortical – fluid produced in the laboratory ...

$$\omega = (P_\Lambda + P_{\bar{\Lambda}})k_B T / \hbar \sim 0.6 - 2.7 \times 10^{22} \text{ s}^{-1}$$

$$P_\Lambda \approx \frac{1}{2} \frac{\omega}{T} + \frac{\mu_\Lambda B}{T} \quad P_{\bar{\Lambda}} \approx \frac{1}{2} \frac{\omega}{T} - \frac{\mu_\Lambda B}{T}$$

$$\frac{dN}{d\Omega^*} = \frac{1}{4\pi} (1 + \alpha_H \mathbf{P}_H \cdot \mathbf{p}_p^*)$$

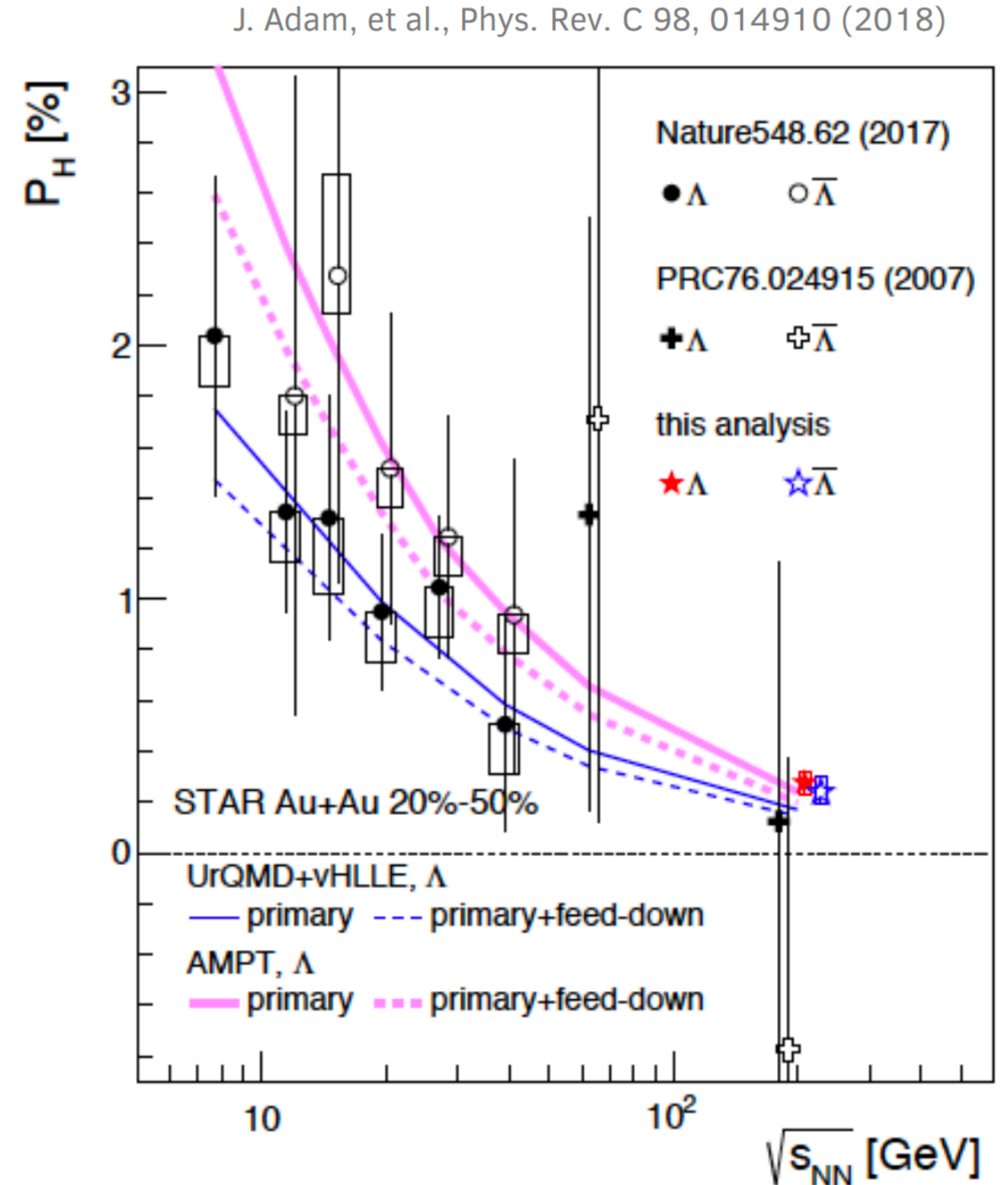
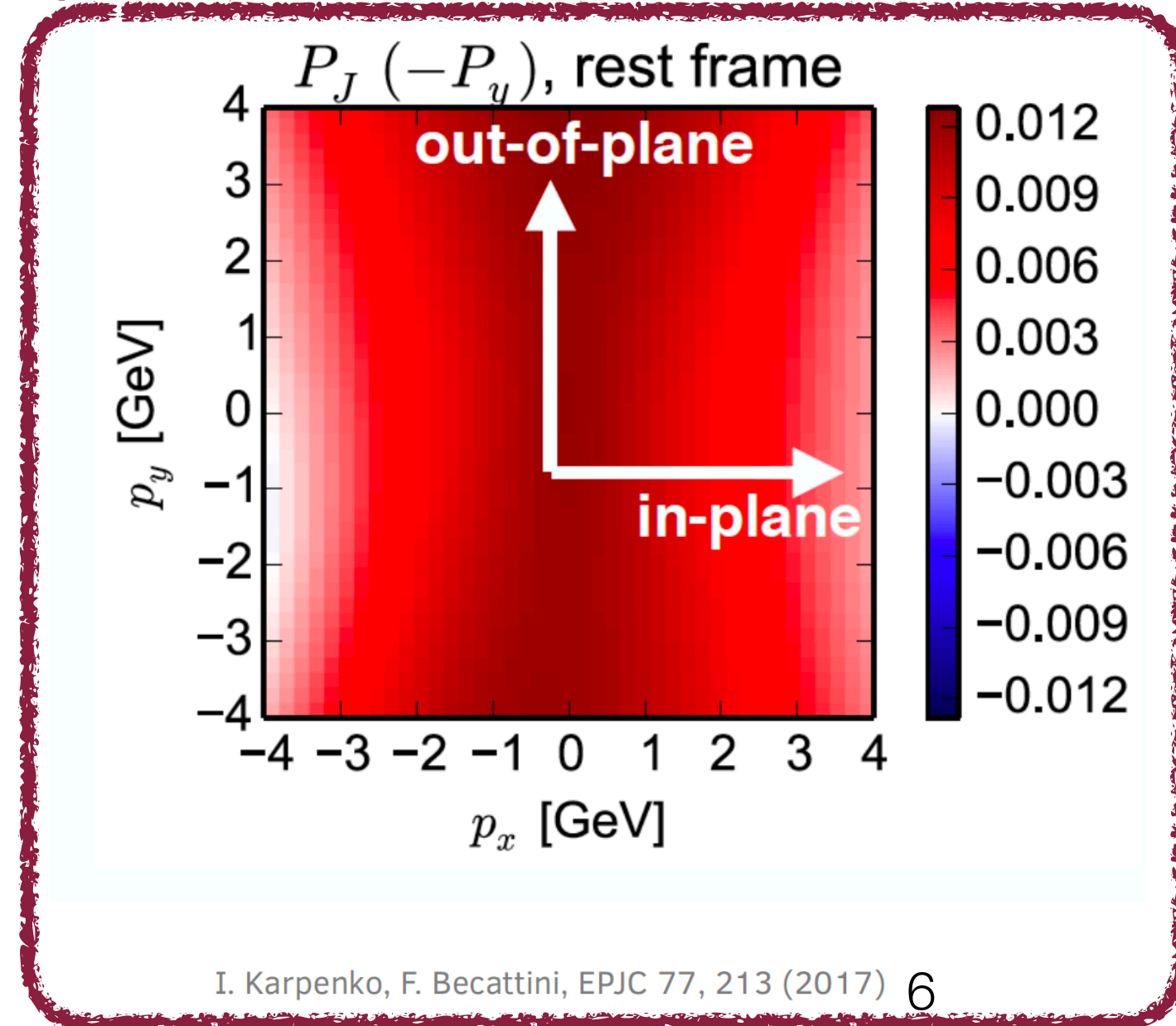
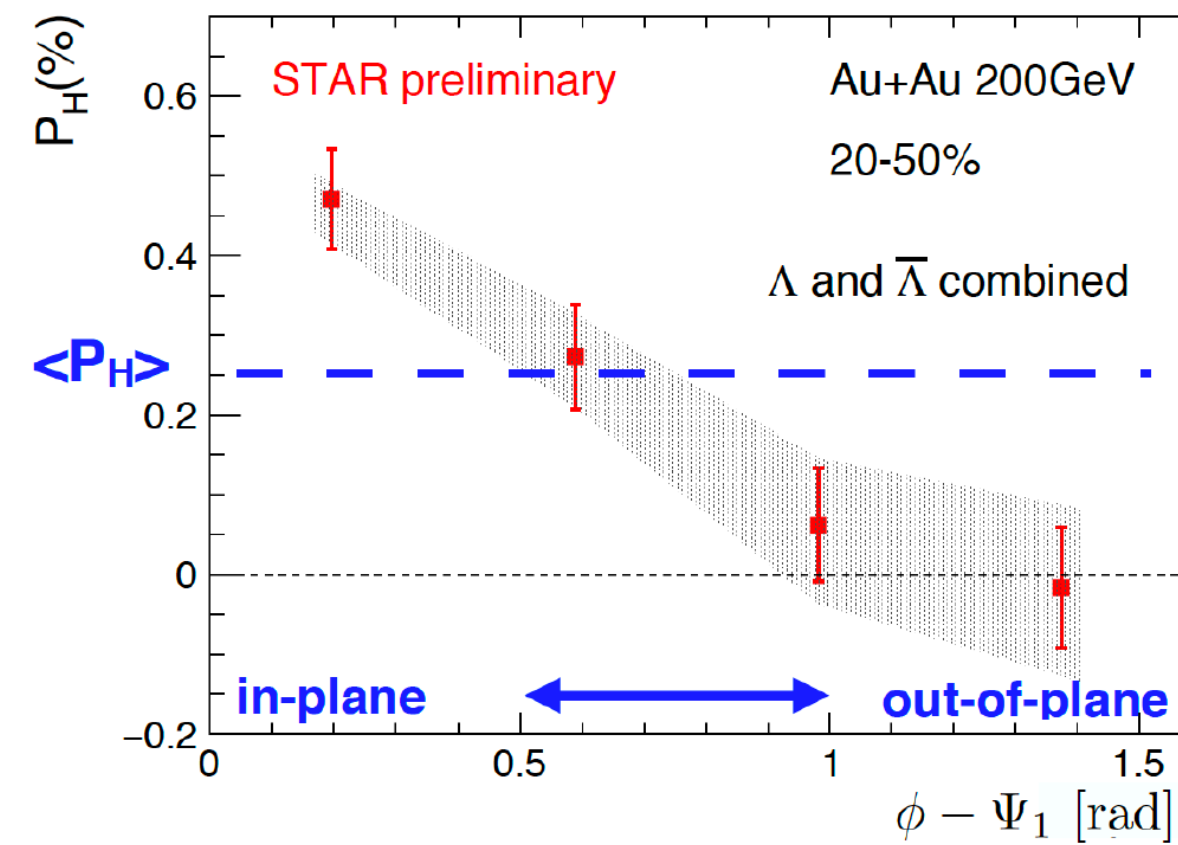
$P_\Lambda \approx P_{\bar{\Lambda}}$ → first direct observation of spin

Global polarization

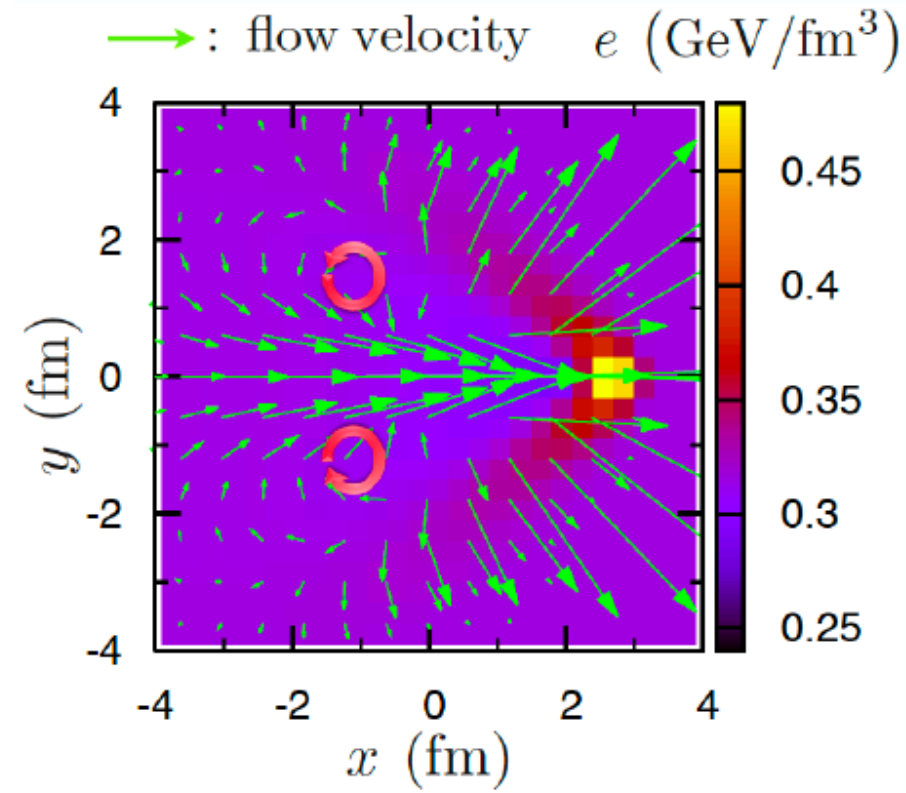
Global polarization data supports the spin-thermal approach

Signal is pretty robust and agrees for both multiphase transport model (AMPT) and viscous hydrodynamics (UrQMD+vHLLE)

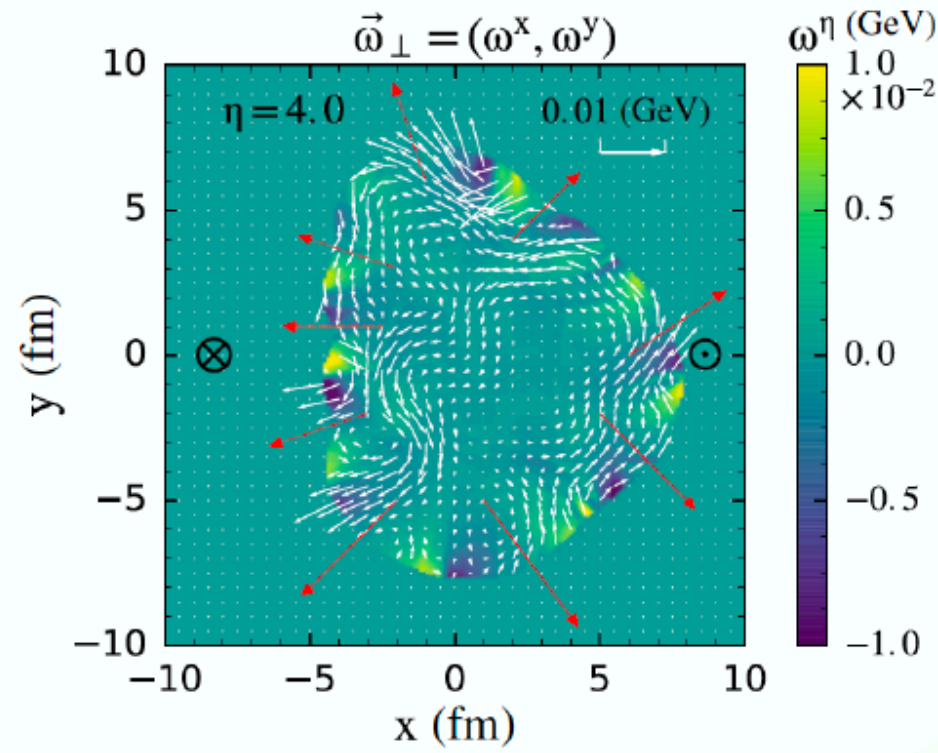
Azimuthal modulation is not captured



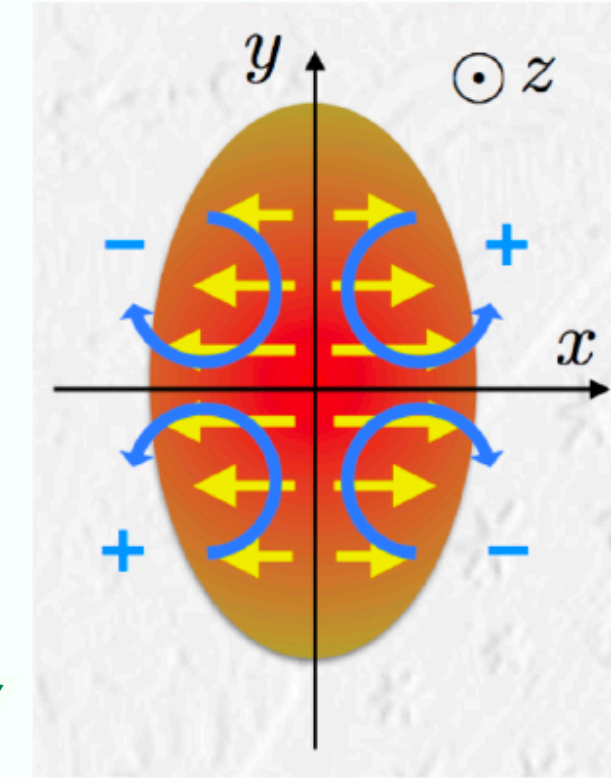
Local (momentum-differential) polarization



Y. Tachibana and T. Hirano, NPA904-905 (2013) 1023

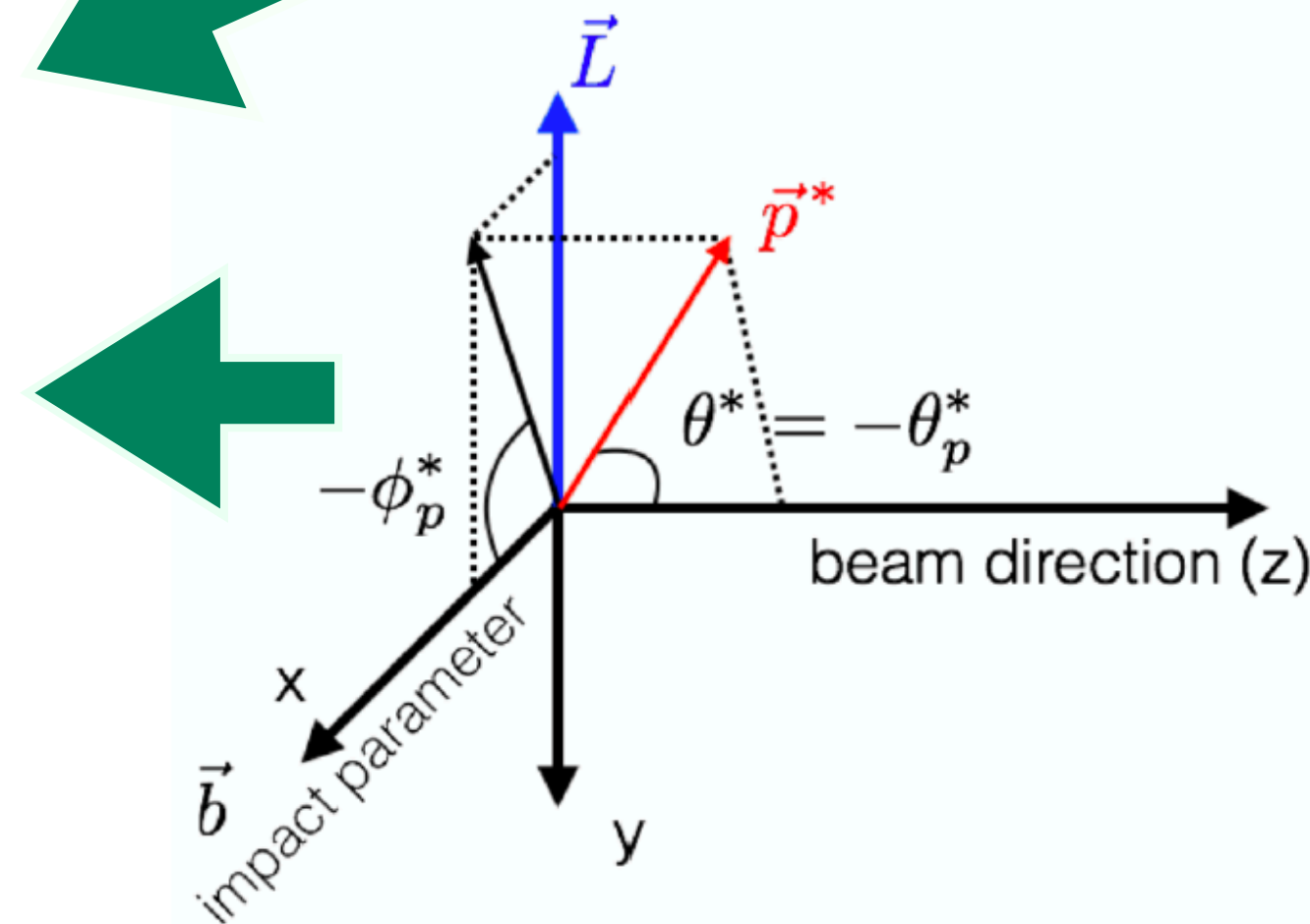
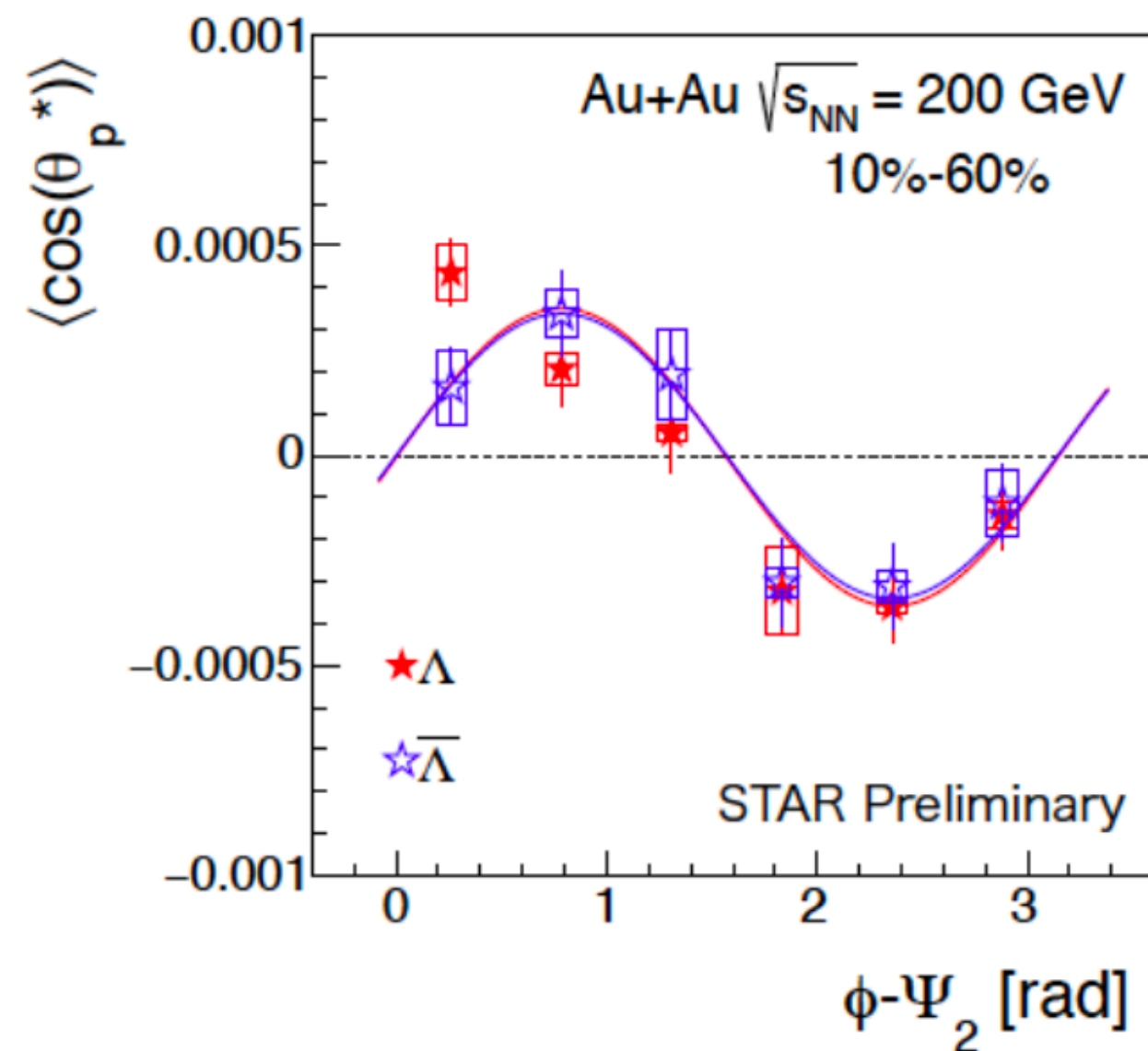


L.-G. Pang, H. Peterson, Q. Wang, ... Wang PRL117, 192001 (2016)



Flow structure in the transverse plane (jet, vbe fluctuations etc.) may generate longitudinal polarization

F. Becattini and I. Karpenko, PRL120.012302 (2018)
S. Voloshin, EPJ Web Conf.171, 07002 (2018)



$$\frac{dN}{d\Omega^*} = \frac{1}{4\pi} (1 + \alpha_H \mathbf{P}_H \cdot \mathbf{p}_p^*)$$

$$\langle \cos \theta_p^* \rangle = \int \frac{dN}{d\Omega^*} \cos \theta_p^* d\Omega^*$$

$$= \alpha_H P_z \langle (\cos \theta_p^*)^2 \rangle$$

$$\therefore P_z = \frac{\langle \cos \theta_p^* \rangle}{\alpha_H \langle (\cos \theta_p^*)^2 \rangle}$$

$$= \frac{3 \langle \cos \theta_p^* \rangle}{\alpha_H} \quad (\text{if perfect detector})$$

α_H : hyperon decay parameter

θ_p^* : θ of daughter proton in Λ rest frame