

Calculation of ρ meson decay width from the PACS-CS configurations

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Naruhito Ishizuka for PACS-CS collaboration
University of Tsukuba

We present results of ρ meson decay width
evaluated from SC. phase shift for $I=1$ two-pion system
with PACS-CS configurations generated at

$$N_f = 2 + 1, \text{ Imp. Wilson}$$

$$a = 0.091 \text{ fm}, L = 2.9 \text{ fm}$$

$$m_\pi = 410 \text{ MeV}$$

We use NEW Op. to reduce the computational cost.

1. Introduction

Previous works of ρ meson decay

1) From transition amp.

$$\langle 0 | (\pi\pi)(t) \rho(0) | 0 \rangle \sim \left[1 + \langle \pi\pi | \rho \rangle \cdot t + O(t^2) \right] \times e^{-m_\rho t}$$

\sim calc. of WME.

$$\rightarrow \Gamma = |\langle \pi\pi | \rho \rangle|^2 \cdot L^3 k E / (24\pi) = g_{\rho\pi\pi}^2 / (6\pi) \cdot k^3 / E^2$$

- 1) S. Gottlieb, P.B. Mackenzie, H.B. Thacker, D. Weingarten [PL134B\(84\)346](#).

$$\text{Quench, } m_\pi/m_\rho = 0.84 - 0.91$$

- 2) R.D. Loft, T.A. DeGrand [PRD9\(1989\)2692](#).

$$\text{Quench, } m_\pi/m_\rho = 0.9$$

- 3) C. McNeile, C. Michael + UKQCD [PLB556\(2003\)177](#).

$$N_f = 2, m_\pi/m_\rho = 0.578_{-19}^{+13}$$

$$g_{\rho\pi\pi} = 6.77_{-56}^{+91}$$

- 4) K. Jansen, C. McNeile, C. Michael, C. Urbach (ETMC) [PRD80\(2009\)054510](#).

$$N_f = 2, m_\pi/m_\rho = 0.336 - 0.534$$

$$g_{\rho\pi\pi} = 5.2(1.3)$$

Problem : $\langle \pi\pi | \rho \rangle$ at \sqrt{s} Where is \sqrt{s} ? : **ambiguous**

2) From SC. phase

Energy on lat. \longrightarrow $\tan \delta$ \longrightarrow $g_{\rho\pi\pi}, \Gamma$
by finite size formula

exp. :

$$g_{\rho\pi\pi} = 5.98(2)$$

$$\Gamma = \frac{g_{\rho\pi\pi}^2}{6\pi} \frac{k_\rho^3}{m_\rho^2} = 149.1(8) \text{ MeV}$$

1) CP-PACS [PRD76\(07\)094506](#).

$$N_f = 2, a = 0.21 \text{ fm}, L = 2.5 \text{ fm}, m_\pi = 330 \text{ MeV}$$

Imp.Wilson, Moving frame with $\mathbf{p} = \mathbf{e}_3 \cdot (2\pi/L)$

$$g_{\rho\pi\pi} = 6.25(67)$$

2) M. Göckeler, R. Horsley, Y. Nakamura et.al. (QCDSF) [arXiv:0810.5337 \(Lat08\)](#)

$$N_f = 2, a = 0.072 - 0.084 \text{ fm}, m_\pi = 240 - 810 \text{ MeV}$$

Imp.Wilson, CM. frame

$$g_{\rho\pi\pi} = 5.3_{-1.5}^{+2.1}$$

3) X. Feng, K. Jansen, D.B. Renner (ETMC) [arXiv:0910.4871 \(Lat09\)](#)

$$N_f = 2, a = 0.086 \text{ fm}, L = 2.1 \text{ fm}, m_\pi = 391 \text{ MeV}$$

tmQCD, CM and Moving frame with $\mathbf{p} = \mathbf{e}_3 \cdot (2\pi/L)$

$$g_{\rho\pi\pi} = 6.16(48)$$

4) X. Feng, L. Jansen, D.B. Renner (ETMC) *Just previous talk*

5) J. Frison *Talk on 18(Fri.) 17:40 at Room1*

This work : from SC. phase shift

Gauge conf. :

$$N_f = 2 + 1 \quad , \quad \text{Imp. Wilson}$$

$$\beta = 1.9 \quad , \quad 32^3 \times 64 \quad , \quad K_{ud} = 0.13754 \quad , \quad K_s = 0.13640$$

$$a = 0.091 \text{ fm} \quad , \quad \underline{L = 2.9 \text{ fm}}$$

$$\underline{m_\pi = 410 \text{ MeV}}$$

generated by PACS-CS col.
PRD79(2009)034503.

Calc. points :

tot. mom.		state $[\sqrt{s}/m_\rho]$ (without interaction)		
\mathbf{p}	Γ	ground	1st Ex.	referred by
$(0, 0, 0)$ CM. F.	\mathbf{T}_1^-	$\rho_{1,2,3}(0, 0, 0) [1]$	$\pi(0, 0, 1) \pi(0, 0, -1) [1.3]$	$\mathbf{T}_1(\rho(\mathbf{0}))$
$(0, 0, 1)$ Moving F.	\mathbf{E}^-	$\rho_{1,2}(0, 0, 1) [1]$	$\pi(0, 1, 1) \pi(0, -1, 0) [1.4]$	$\mathbf{E}(\rho_T(\mathbf{p}))$
$(0, 0, 1)$ Moving F.	\mathbf{A}_2^-	$\rho_3(0, 0, 1) [1]$	$\pi(0, 0, 0) \pi(0, 0, 1) [1.02]$	$\mathbf{A}_2(\rho_3(\mathbf{p}), \pi(\mathbf{0})\pi(\mathbf{p}))$ ($\mathbf{p} = (0, 0, 1)$)

(by using variational method)

2. Method

Finite size formula :

M. Lüscher, NPB354(1991)531.

K.Rummannen and S.Gottlib, NPB450(1995)397.

$$\text{for } \mathbf{T}_1(\rho(\mathbf{0})) \quad \left(\begin{array}{l} \sqrt{E^2 - p^2} = \sqrt{s} = 2\sqrt{m_\pi^2 + k^2} \\ q = kL/(2\pi) \end{array} \right)$$
$$\frac{k}{\tan \delta(k)} = Z_{00}(q; \mathbf{0})$$

$$\text{for } \mathbf{E}(\rho_T(\mathbf{p}))$$
$$\frac{k}{\tan \delta(k)} = Z_{00}(q; \mathbf{p}) - \frac{1}{\sqrt{5}} \frac{1}{q^2} Z_{20}(q; \mathbf{p})$$

$$\text{for } \mathbf{A}_2(\rho_3(\mathbf{p}), \pi(\mathbf{0})\pi(\mathbf{p}))$$
$$\frac{k}{\tan \delta(k)} = Z_{00}(q; \mathbf{p}) + \frac{2}{\sqrt{5}} \frac{1}{q^2} Z_{20}(q; \mathbf{p})$$

$$Z_{lm}(q; \mathbf{p}) = 2/(\gamma L \sqrt{\pi}) \cdot \sum_{\mathbf{r} \in D(\mathbf{p})} \mathcal{Y}_{lm}(\mathbf{r}) \cdot (r^2 - q^2)^{-1}$$

: spherical zeta function

$$D(\mathbf{p}) = \{ \mathbf{r} \mid \mathbf{r} = \hat{\gamma}^{-1}(\mathbf{n} + \mathbf{p}/2), \mathbf{n} \in \mathbb{Z}^3 \}$$

$$\gamma = E/\sqrt{s} : \text{Lorentz boost factor}$$

Op. for $\mathbf{A}_2(\rho_3(\mathbf{p}), \pi(\mathbf{0})\pi(\mathbf{p}))$

Energies are extracted by variational method.

For $\mathbf{T}_1(\rho(\mathbf{0}))$ and $\mathbf{E}(\rho_T(\mathbf{p}))$
energies are extracted from time correlation function of ρ meson
as usual calculations of hadron masses.

Source op. :

$$\bar{O}_1(t_s) = \frac{1}{\sqrt{2}} \left(\pi^-(\mathbf{0}, t_s) \pi^+(\mathbf{p}, t_s) - \pi^+(\mathbf{0}, t_s) \pi^-(\mathbf{p}, t_s) \right)$$

$\pi(\mathbf{p}, t_s)$: constructed with U(1) noise

$$\bar{O}_2(t_s) = \sum_{\mathbf{z} \in \Gamma} \frac{1}{\sqrt{2}} \left(\bar{U}(\mathbf{z}, t_s) \gamma_j U(\mathbf{z}, t_s) - \bar{D}(\mathbf{z}, t_s) \gamma_j D(\mathbf{z}, t_s) \right) \cdot e^{i\mathbf{p} \cdot \mathbf{z}}$$

(: smearing OP. for ρ meson with \mathbf{p}) (on Coulomb gauge)

$$U(\mathbf{z}, t_s) = \sum_{\mathbf{x}} u(\mathbf{z} + \mathbf{x}, t_s) \cdot F(\mathbf{x}) \quad (: \text{smearing quark})$$

$$F(\mathbf{x}) = A \cdot e^{-B|\mathbf{x}|} \quad : \text{smearing function}$$

$$\Gamma = \{ \mathbf{z} \mid L/2 \times (n_1, n_2, n_3), n_j = (0 \text{ or } 1) \}$$

Sink op :

$$O_1(t) = \frac{1}{\sqrt{2}} \left(\pi^-(\mathbf{0}, t_1) \pi^+(\mathbf{p}, t) - \pi^+(\mathbf{0}, t_1) \pi^-(\mathbf{p}, t) \right) \cdot e^{m_\pi \cdot (t_1 - t)}$$

with fixed t_1

$$\langle 0 | O_1^\dagger(t) \sim \langle 0 | \pi(\mathbf{0}) | \pi(\mathbf{0}) \rangle \langle \pi(\mathbf{0}) | \pi(\mathbf{p}) | E \rangle \langle E | \times e^{E \cdot t} \quad \text{for } t_1 \gg t$$

$|E\rangle$: energy eigenstate ($\sim |\pi\pi; E\rangle$)

(: same t dependence as usual Heisenberg op.)

$$O_2(t) = \sum_{\mathbf{x}} \frac{1}{\sqrt{2}} \left(\bar{u}(\mathbf{x}, t) \gamma_3 u(\mathbf{x}, t) - \bar{d}(\mathbf{x}, t) \gamma_3 d(\mathbf{x}, t) \right) \cdot e^{i\mathbf{p} \cdot \mathbf{x}}$$

Variational method

2 x 2 correlation matrix :

$$G_{ij}(t) = \langle 0 | O_i^\dagger(t) \bar{O}_j(t_s) | 0 \rangle$$

with

$$\begin{pmatrix} \text{sink} & : & O_1(t) , & O_2(t) \\ \text{source} & : & \bar{O}_1(t) , & \bar{O}_2(t) \end{pmatrix}$$

Assuming 2 state dominant,

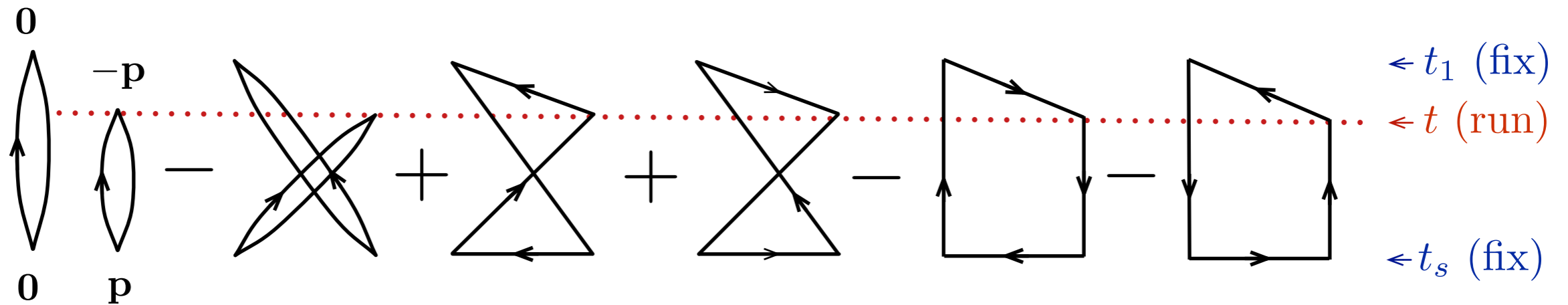
$$\text{Ev}[G(t_R)^{-1} G(t)] = \left\{ \lambda_1(t) = e^{-E_1 \cdot t} , \lambda_2(t) = e^{-E_2 \cdot t} \right\}$$

for $t_1 \gg t \gg t_s$

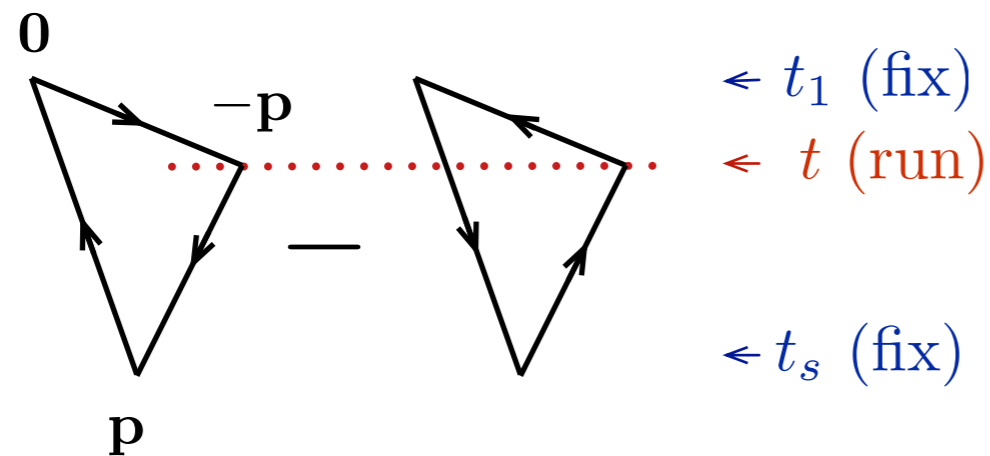
Calc. of $G(t)$ of $\mathbf{A}_2(\rho_3(\mathbf{p}), \pi(\mathbf{0})\pi(\mathbf{p}))$

Diagrams :

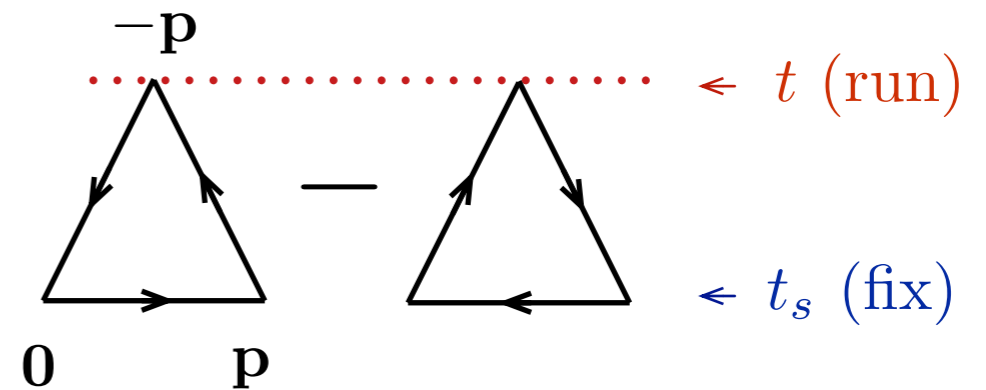
$$G_{\pi\pi \rightarrow \pi\pi}(t) =$$



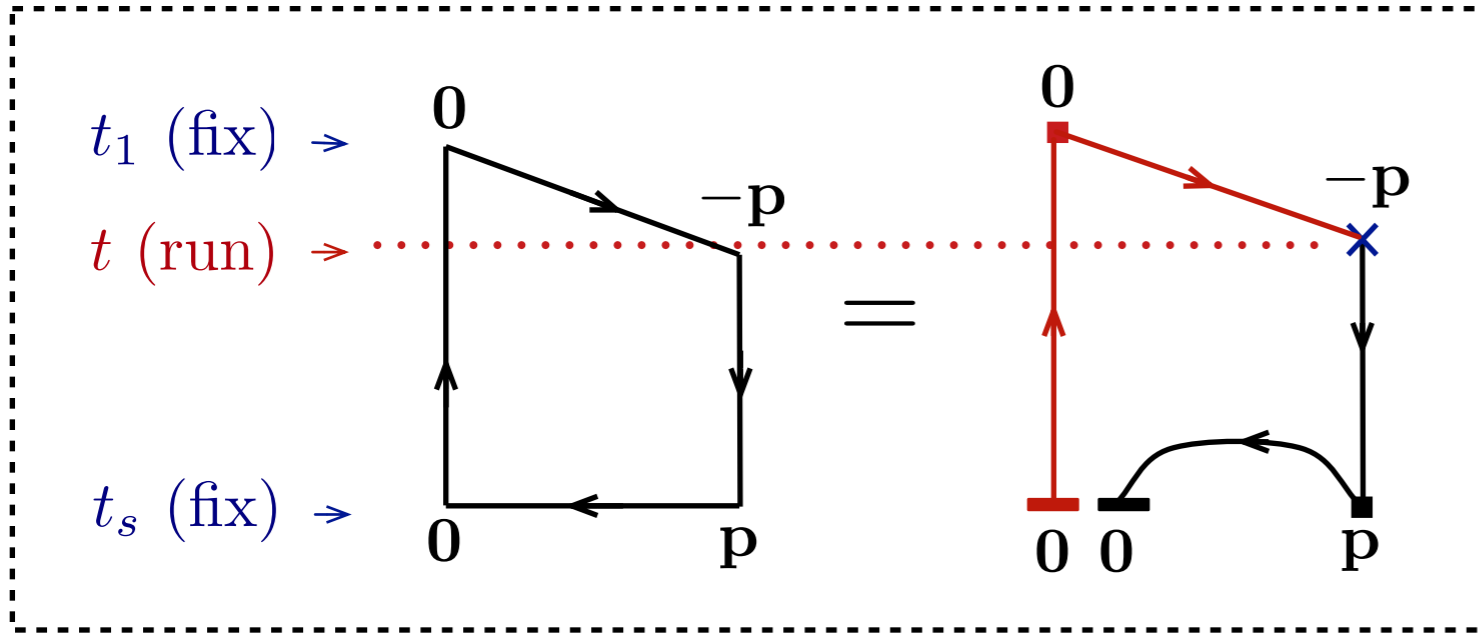
$$G_{\rho \rightarrow \pi\pi}(t) =$$



$$G_{\pi\pi \rightarrow \rho}(t) =$$



Example :



(1)

$$\begin{aligned}
 & \mathbf{q}, t_s \quad \text{---} \quad \mathbf{x}, t \\
 &= \sum_{\mathbf{y}} D^{-1}(\mathbf{x}, t; \mathbf{y}, t_s) \cdot \xi(\mathbf{y}) e^{i\mathbf{q} \cdot \mathbf{y}} \\
 &\equiv Q(\mathbf{x}, t; \mathbf{q}, t_s)
 \end{aligned}$$

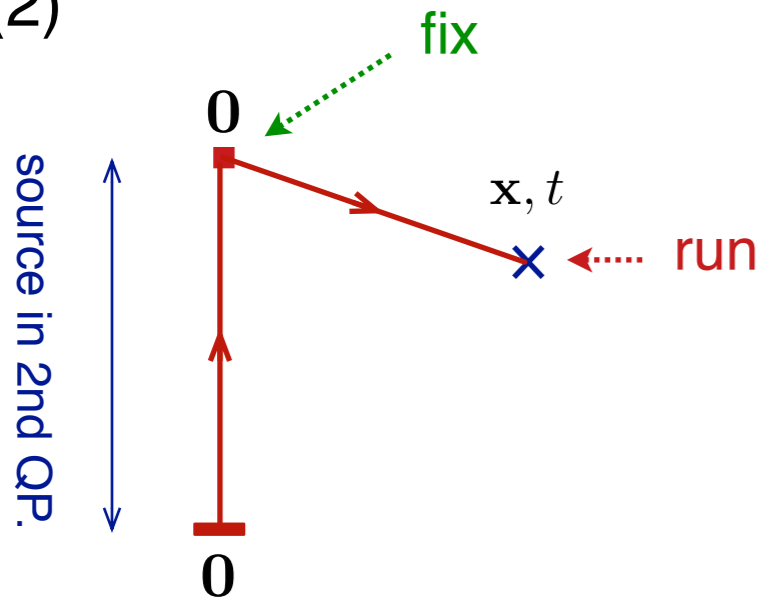
$\xi(\mathbf{x})$: U(1) noise

$$\sum_{j=1}^{N_R} \xi_j^\dagger(\mathbf{x}) \xi_j(\mathbf{y}) = \delta^3(\mathbf{x} - \mathbf{y})$$

for $N_R \rightarrow \infty$

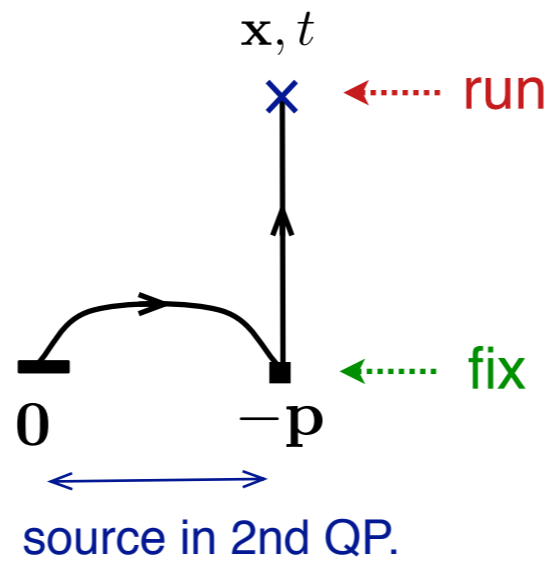
($N_R = 10$ (in present work))

(2)



$$= \sum_{\mathbf{y}} D^{-1}(\mathbf{x}, t; \mathbf{y}, t_s) \cdot \gamma_5 Q(\mathbf{y}, t_1; \mathbf{0}, t_s)$$

$$\equiv R(\mathbf{x}, t; \mathbf{0}, t_s)$$



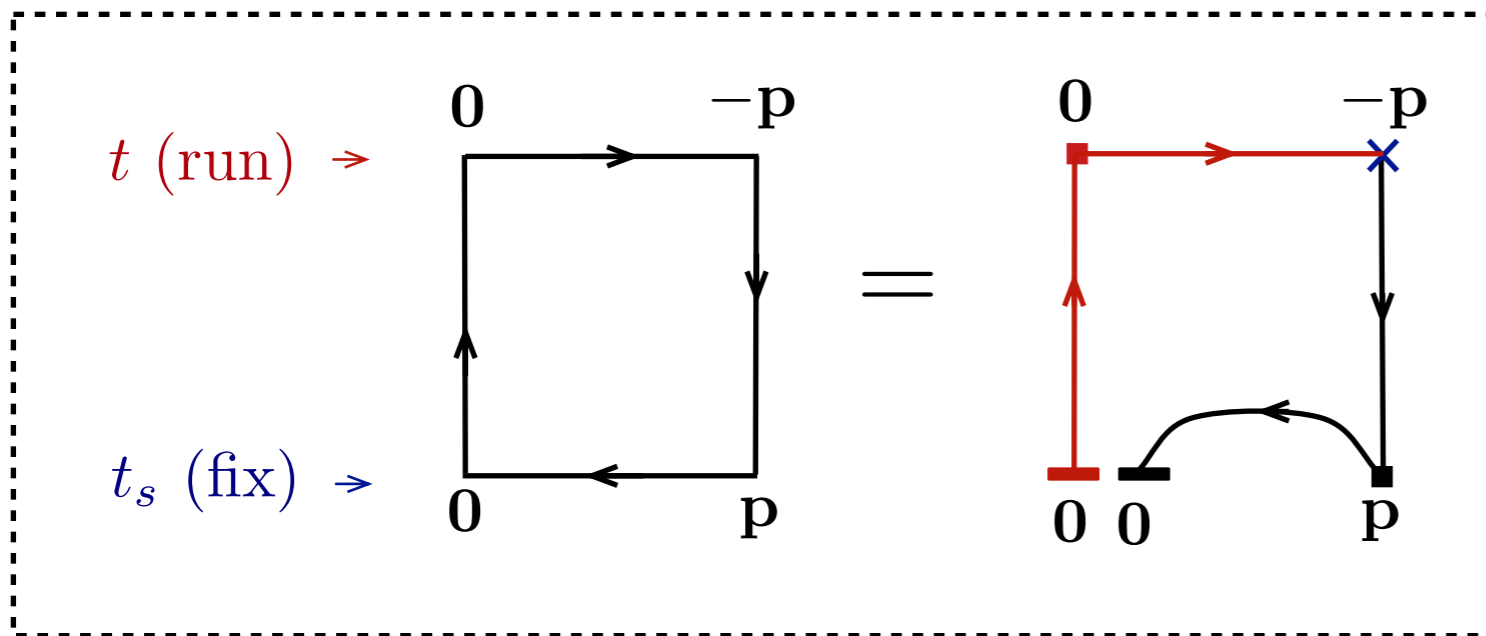
$$= \sum_{\mathbf{y}} D^{-1}(\mathbf{x}, t; \mathbf{y}, t_s) \cdot \gamma_5 e^{-i\mathbf{p} \cdot \mathbf{y}} Q(\mathbf{y}, t_s; \mathbf{0}, t_s)$$

$$\equiv B(\mathbf{x}, t; \mathbf{0}, t_s)$$

(3)

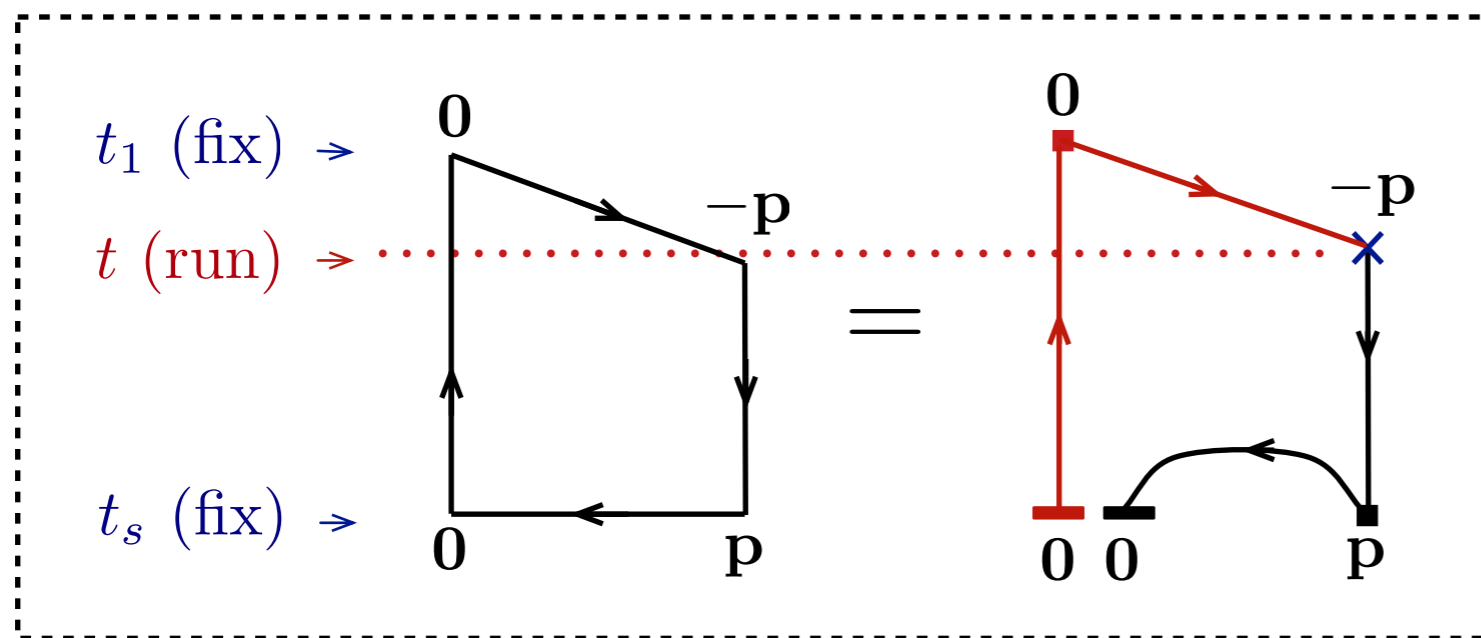
$$\sum_{\mathbf{y}} \text{Tr} \left[R(\mathbf{x}, t; \mathbf{0}, t_s) B^\dagger(\mathbf{x}, t; \mathbf{0}, t_s) \right] \cdot e^{-i\mathbf{p} \cdot \mathbf{x}}$$

CP-PACS(2007) $\pi\pi$ op. : $\pi(\mathbf{0}, t)\pi(\mathbf{p}, t)$



[# of Red type of QP.]
 \propto [# of time slice : T]
 $\Rightarrow 1$

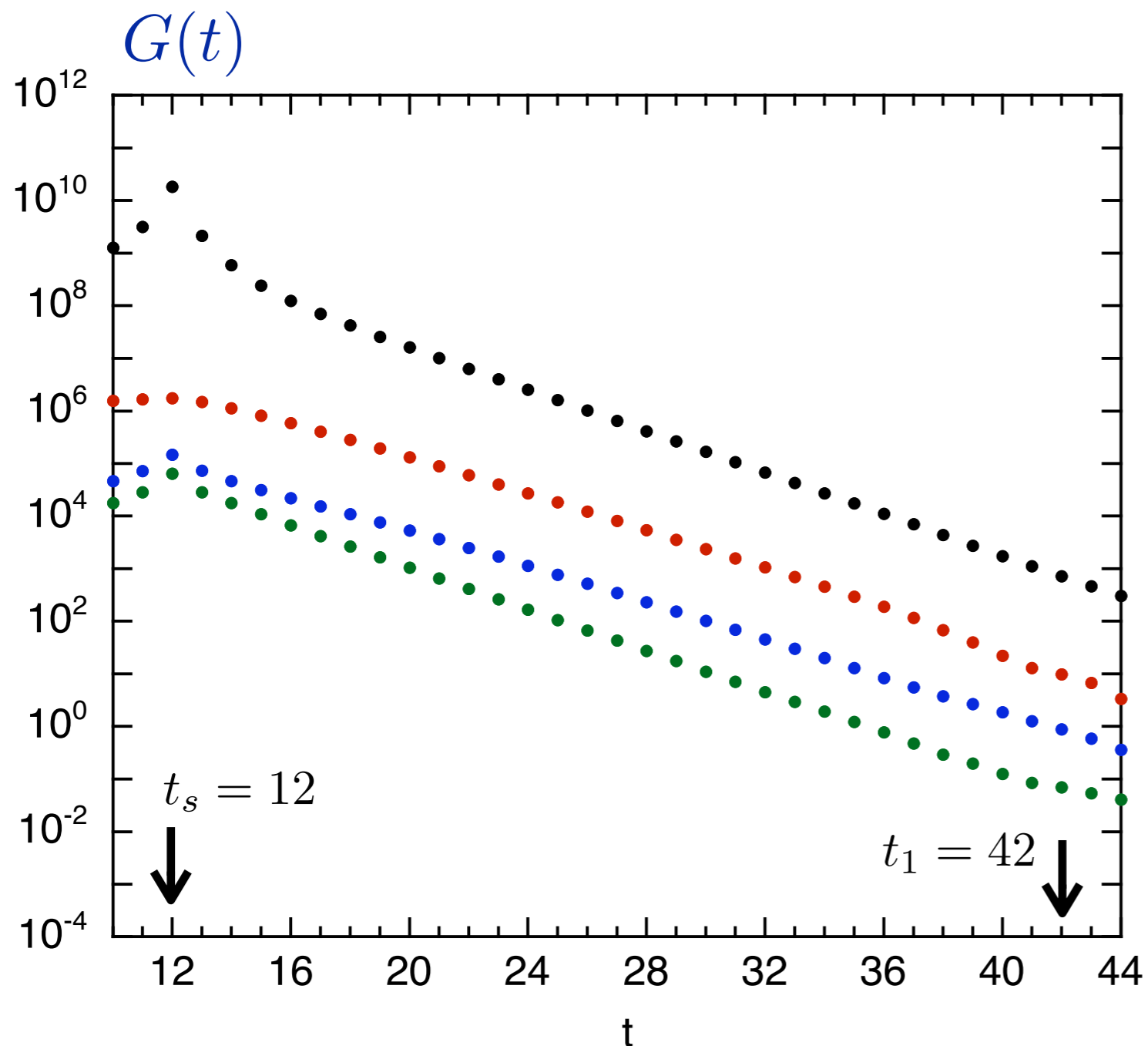
This work $\pi\pi$ op. : $\pi(\mathbf{0}, t_1)\pi(\mathbf{p}, t)$



Computational cost
 is reduced by T .

3. Results

1) Results of $A_2(\rho_3(\mathbf{p}), \pi(\mathbf{0})\pi(\mathbf{p}))$



- $\text{Re}[\pi\pi \rightarrow \pi\pi]$
- $\text{Im}[\rho \rightarrow \pi\pi]$
- $-\text{Im}[\pi\pi \rightarrow \rho]$
- $\text{Re}[\rho \rightarrow \rho]$

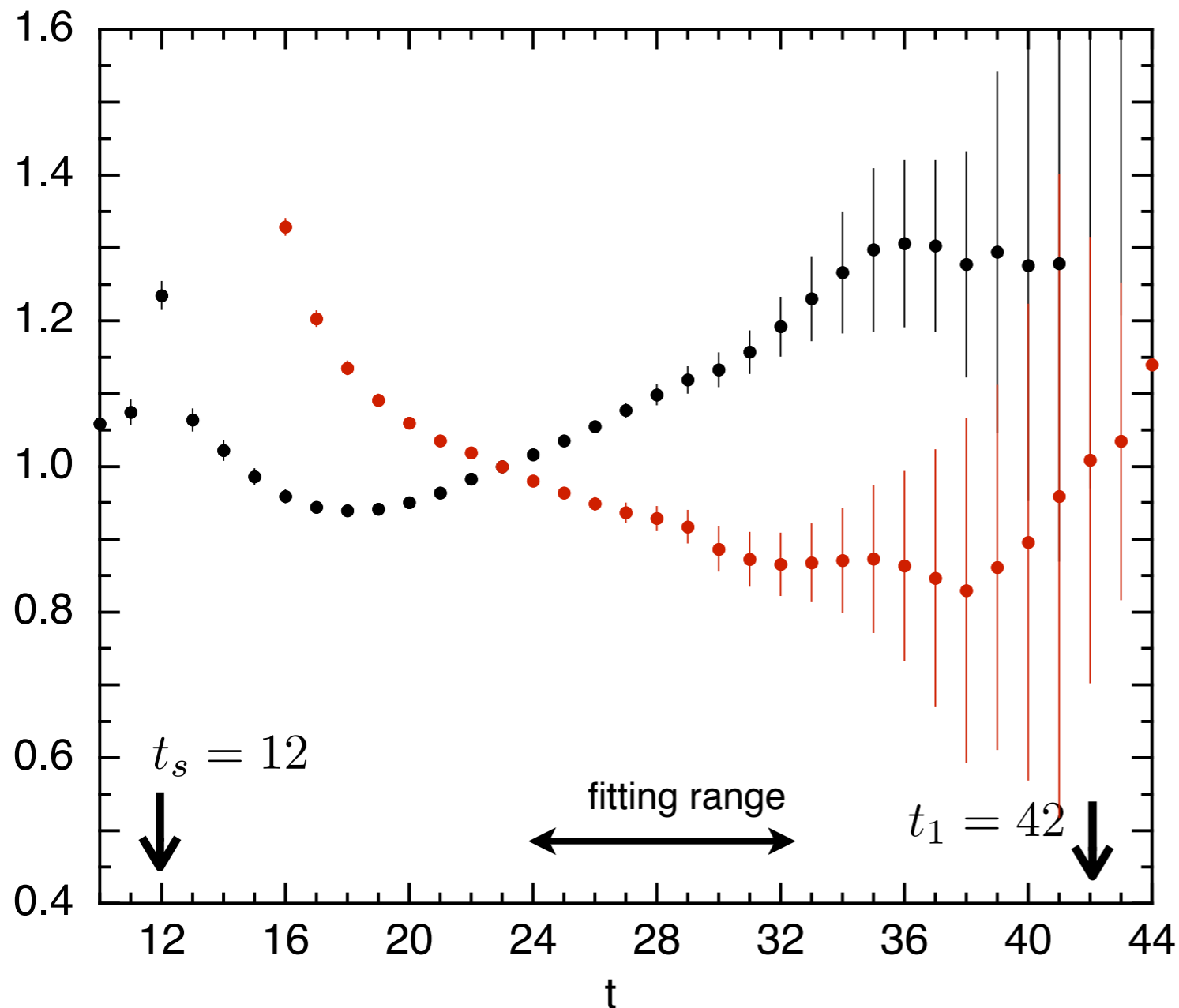
From CP and P sym.

$\rho \rightarrow \pi\pi$ and $\pi\pi \rightarrow \rho$
: pure imaginary

$\pi\pi \rightarrow \pi\pi$ and $\rho \rightarrow \rho$
: real

Eigenvalue $\lambda(t) = \text{Ev}[G(t_R)^{-1} G(t)] \quad (t_R = 23)$

$$\lambda(t) / \langle \pi(\mathbf{0}, t) \pi(\mathbf{0}, t_s) \rangle / \langle \pi(\mathbf{p}, t) \pi(\mathbf{p}, t_s) \rangle \sim e^{-\Delta E \cdot t}$$



- ground state

$$\Delta E = -0.0191(27)$$

$$E = 0.4413(36)$$

$$\sqrt{s} = 0.3952(40)$$

$$\tan \delta(k) = +0.0942(47)$$

attractive

- 1st Ex. state

$$\Delta E = 0.0150(42)$$

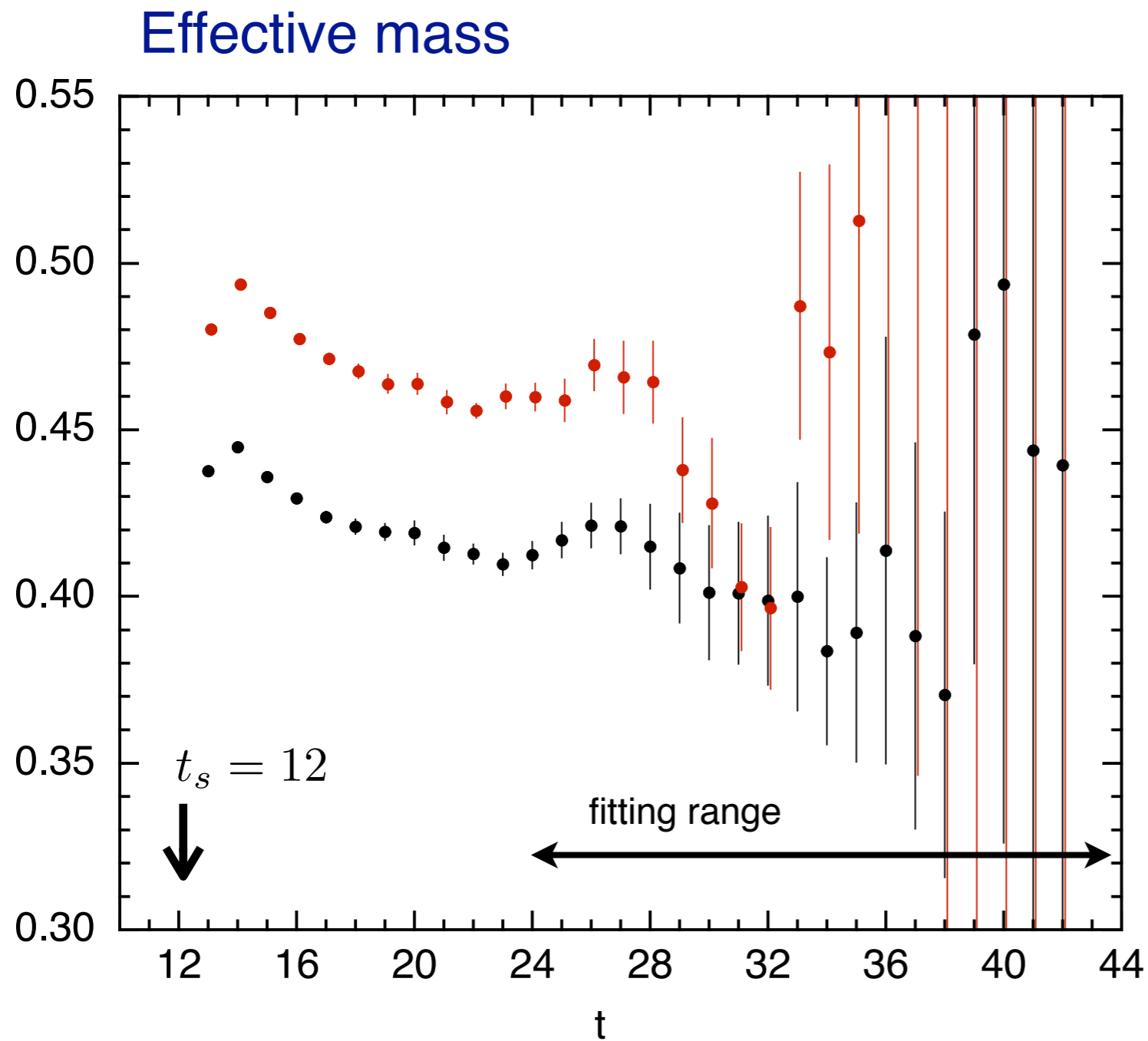
$$E = 0.4755(44)$$

$$\sqrt{s} = 0.4330(48)$$

$$\tan \delta(k) = -0.165(50)$$

repulsive

2) Results of $\mathbf{T}_1(\rho(\mathbf{0}))$ and $\mathbf{E}(\rho_T(\mathbf{p}))$



● $\mathbf{T}_1(\rho(\mathbf{0}))$

$$E = 0.4122(84)$$

$$\sqrt{s} = 0.4122(84)$$

$$\tan \delta(k) = -0.74(24)$$

repulsive

● $\mathbf{E}(\rho_T(\mathbf{p}))$

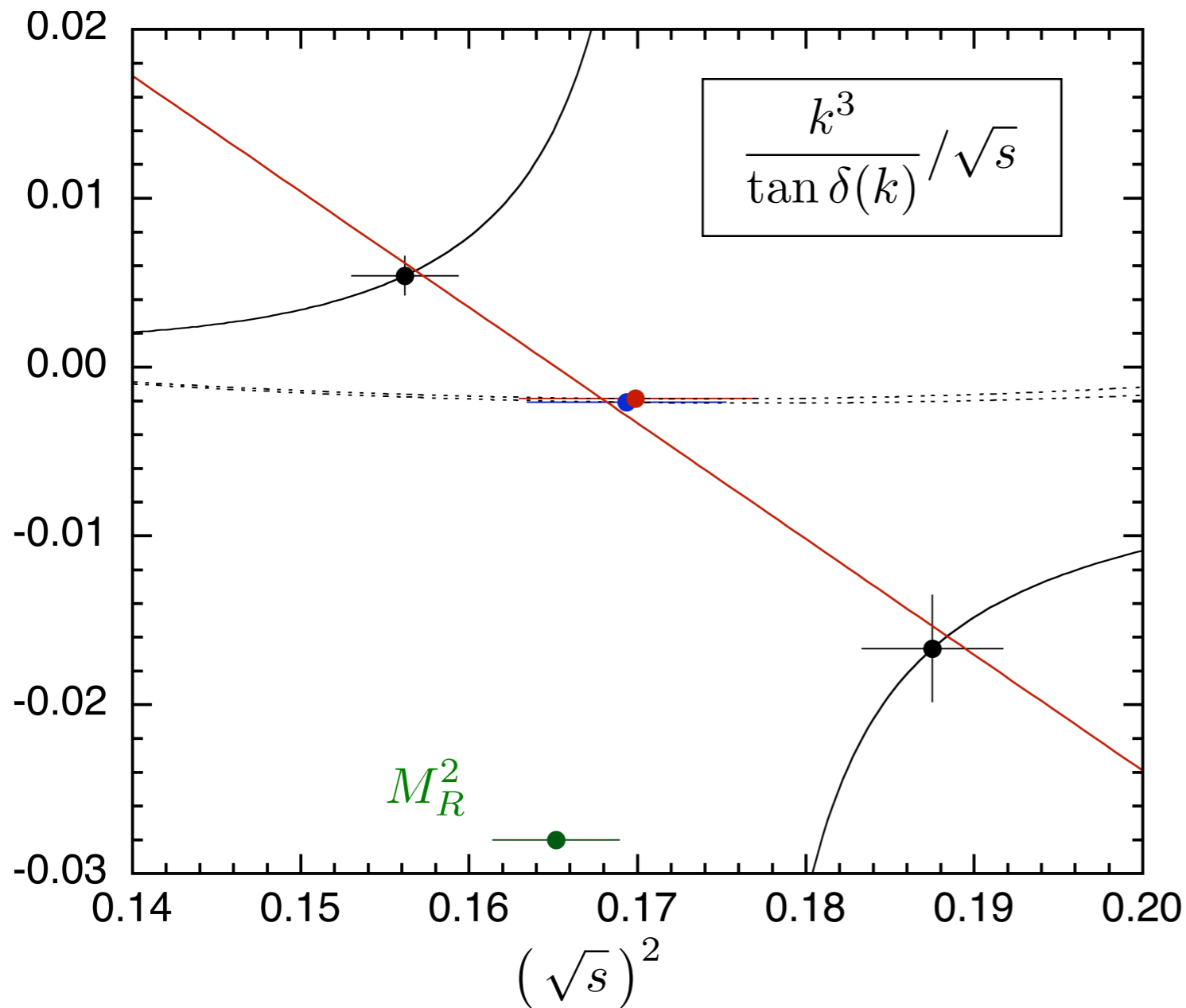
$$E = 0.4560(65)$$

$$\sqrt{s} = 0.4115(72)$$

$$\tan \delta(k) = -0.65(18)$$

repulsive

3) Results of decay width



exp. :

$$g_{\rho\pi\pi} = 5.98(2)$$

$$\Gamma = 149.1(8) \text{ MeV}$$

● $\mathbf{A}_2 (\rho_3(\mathbf{p}), \pi(\mathbf{0})\pi(\mathbf{p}))$

● $\mathbf{T}_1 (\rho(\mathbf{0}))$

● $\mathbf{E} (\rho_T(\mathbf{p}))$

— : finite size formula

— : fitting line

$$\frac{k^3}{\tan \delta(k) / \sqrt{s}} = \frac{6\pi}{g_{\rho\pi\pi}^2} \cdot \left(M_R^2 - (\sqrt{s})^2 \right) \text{ for } \sqrt{s} \sim M_R$$

$$g_{\rho\pi\pi} = 5.24(51)$$

$$M_R = 0.4064(46)$$

$$\Gamma = \frac{g_{\rho\pi\pi}^2 k_\rho^3}{6\pi m_\rho^2} = 113(22) \text{ MeV}$$

4. Summary

We calculate SC. phase shift for $I=1$ two-pion system and evaluate $\rho\pi\pi$ coupling constant with PACS-CS configurations generated at

$$N_f = 2 + 1, \text{ Imp. Wilson}$$

$$a = 0.091 \text{ fm}, \quad L = 2.9 \text{ fm}$$

$$m_\pi = 410 \text{ MeV}$$

Our results :

$$g_{\rho\pi\pi} = 5.24(51)$$

$$M_R = 0.4064(46)$$

$$\Gamma = \frac{g_{\rho\pi\pi}^2}{6\pi} \frac{k_\rho^3}{m_\rho^2} \\ = 113(22) \text{ MeV}$$

exp. :

$$g_{\rho\pi\pi} = 5.98(2)$$

$$\Gamma = 149.1(8) \text{ MeV}$$

: slightly smaller than exp.

possible reason :

(1) mass dependence

(2) $O(a^2)$ error

: Future improvement

Trial calcs :

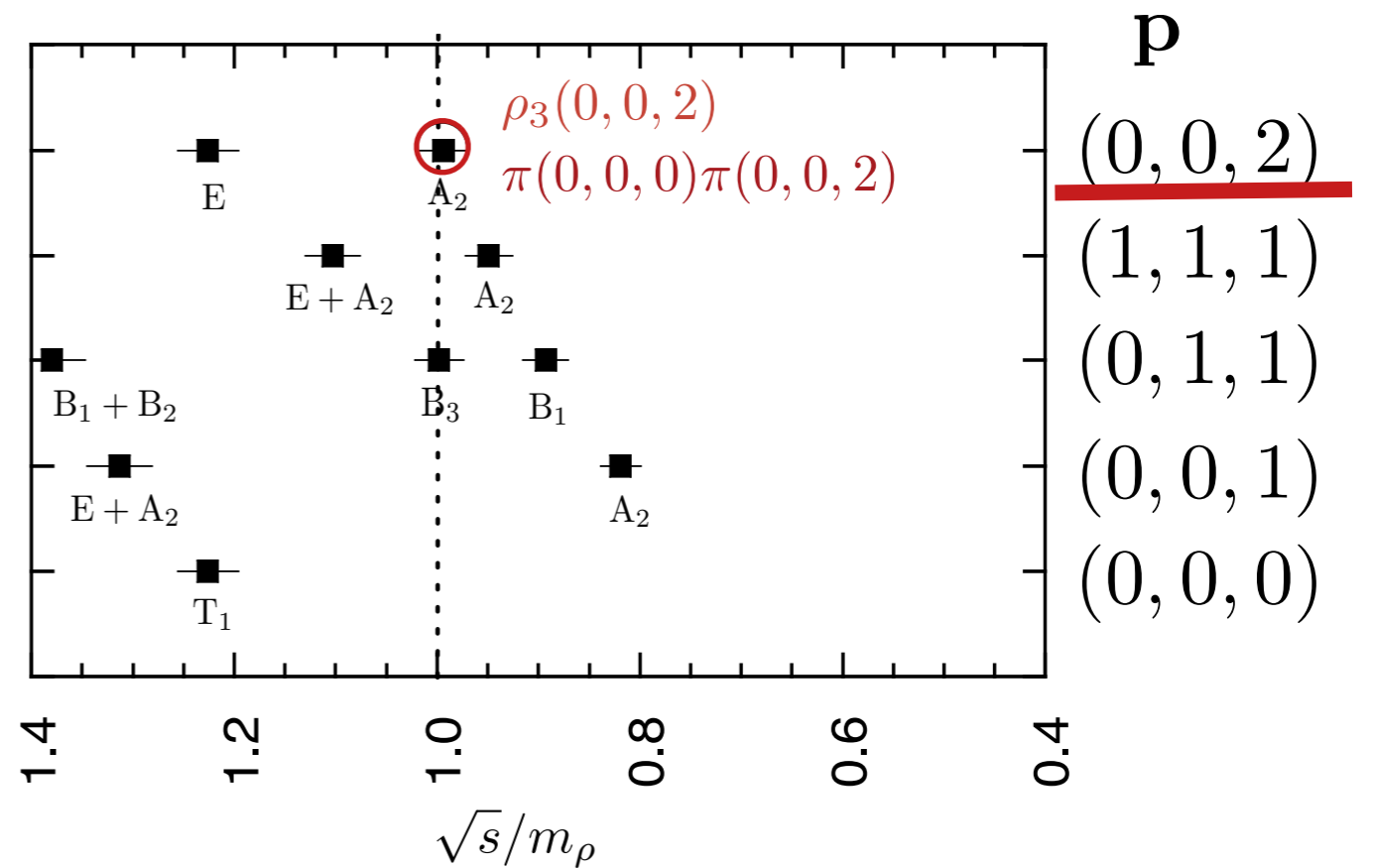
$N_f = 2 + 1$, Imp. Wilson

$\beta = 1.9$, $32^3 \times 64$

$a = 0.091$ fm , $L = 2.9$ fm

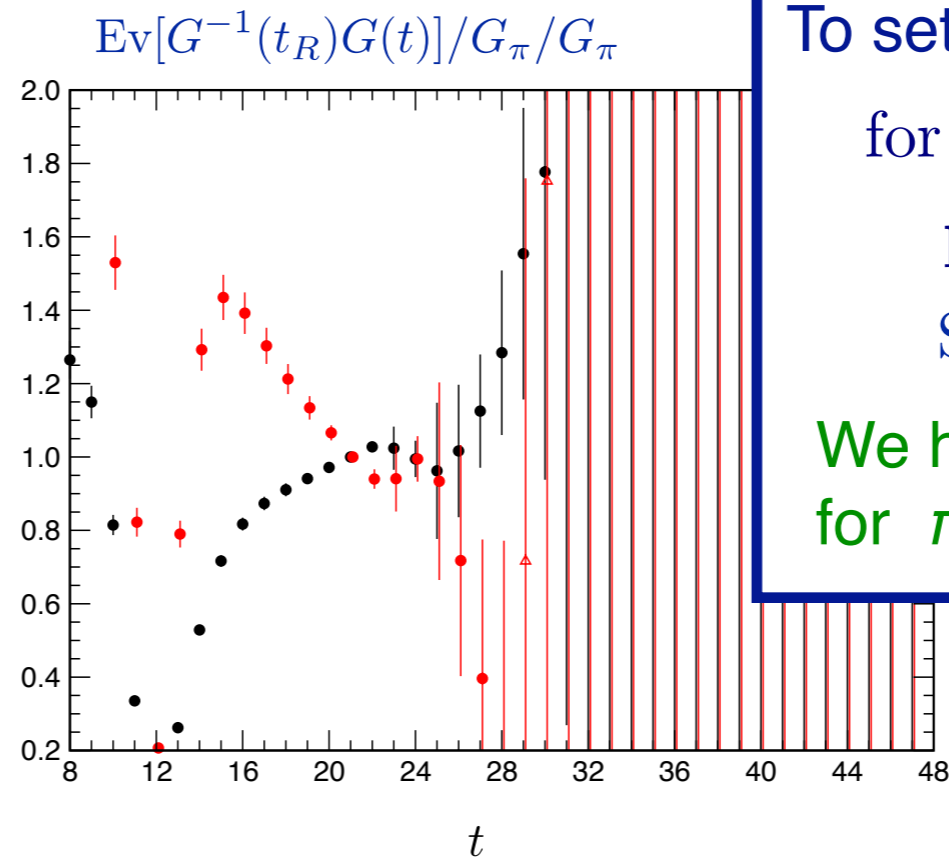
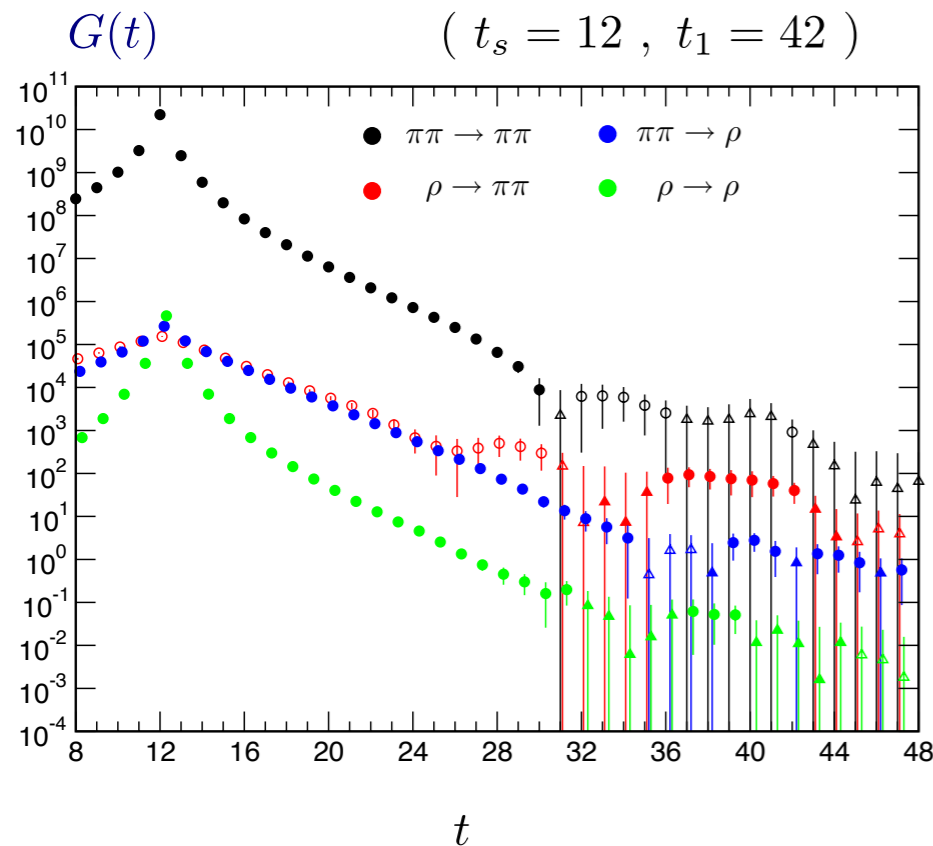
$m_\pi = 300$ MeV

generated by PACS-CS col.
PRD79(2009)034503.



Results of #c. = 92

(70 day by 1.4Tf machine)



To set $\sqrt{s} \sim m_\rho$
for $m_\pi \rightarrow 0$
 $\mathbf{p} \rightarrow$ larger
Stat. Error. \rightarrow large

We have to improve method
for π with high momentum.