# Search for the IR fixed point in the Twisted Polyakov Loop scheme (II)

# Etsuko Itou (Osaka University) arXiv:0910.4196 and Work in progress

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Numerical simulation was carried out on the vector supercomputer NEC SX-8 in YITP, Kyoto University and RCNP, Osaka University SR and BlueGene in KEK

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## Introduction

#### large flavor SU(3) gauge theories:



Gardi and Karliner Nucl.Phys.B529:383-423,1998. Ncw~12 from resummation method of 3 terms of the beta fn

Miransky and Yamawaki Phys.Rev.D55:5051-5066,1997 Ncw=11.9 based on the same idea of Gardi et al.

Ryttov and Sannino Phys.Rev.D78:065001,2008 Ncw=8.25 from calclulation of the anomalous dimension

## Previous and Recent studies in lattice QCD with Nf=12

➢Iwasaki et al: Phys.Rev.D76:034504

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Deuzeman, Lombordo and Pallante arXiv:0904.4662 [hep-ph] there is a conformal phase...

- ➢Fodor, Holland, Kuti, Nogragi, and Schroeder arXiv:0809.4890, 0911:2463 [hep-lat]
- Nf=12 is below the conformal window. (stout-smeared staggered fermion)



The difficulty to study low energy behavior comes from a large discretization error .

#### Summary of our work

Measure the renormalized coupling in the Twisted Polyakov loop scheme (TPL). The existence of the fixed point does not depend on the renormalization scheme.

>Estimate the  $O(a^2)$  discretization error.

Study the growth rate of the renormalized coupling.

We focus on the step scaling at each coupling rather than the beta function, since the latter suffers from accumulation of errors at low energy region.



When the coupling stops running, the ratio (sigma/u) becomes "1".

# Twisted Polyakov loop scheme

Twisted Polyakov loop (TPL) scheme de Divitiis et al.:NPB422:382

TPL coupling: 
$$g_{TP}^2 = \frac{1}{k} \frac{\langle \sum_{y,z} P_1(y,z,L/2a) P_1(0,0,0)^* \rangle}{\langle \sum_{x,y} P_3(x,y,L/2a) P_3(0,0,0)^* \rangle}$$

no O(a/L) error

At tree level,  $g_{TP}^2$  is proportional to bare coupling.  $g_{TP}^2|_{\text{tree}} = g_0^2$  $k = \frac{1}{24\pi^2} \sum_{n=-\infty}^{\infty} \frac{(-1)^n}{n^2 + (1/3)^2} = 0.0318$ 

#### Remarks on TPL coupling

>If the theory is conformal, the correlation length becomes infinity. The divergence part of the OPE for Ploop correlator depends on the conformal dimension of each operator. It cancels out between numerator and denominator. The coefficient of OPE depends on boundary contribution and gives a nontrivial constant. The ratio of Polyakov loop gives a nontrivial constant ( $\neq$ 1).

However, even if there is no conformal field theory, the IR limit of TPL coupling goes to constant.  $g_{TP}^2 \rightarrow \frac{1}{k} \sim 32$ 

 $1/L \rightarrow 0$ 

This is a lattice artifact of the TPL coupling.

## Simulation: Nf=12

### Nf=12 QCD

Lattice size for step scaling

s = 1 : L = 4, 5, 6, 8s = 2 : L = 8, 10, 12, 16

Parameter  $4.5 \le \beta \le 20$ 

Interpolation of beta at fixed L/a

 $f(\beta) = \frac{C_1}{\beta} + \frac{C_2}{\beta^2} + \frac{C_3}{\beta^3} + \frac{C_4}{\beta^4}$ 

Interpolation to L/a=5

At high energy, the data of the large lattice size are larger than that of the small lattice .

This behavior is same as one of the SF scheme.

- Wilson action + staggered fermion
- Hybrid Monte Carlo algorithm (exact algorithm)
- > 50,000 ~ 100,000 Trj. for each parameter



#### In low beta region (4.5 < beta < 6.0)



### Continuum extrapolation for each step scaling



Lattice size for step scaling

s = 1 : L = 4, 5, 6, 8s = 2 : L = 8, 10, 12, 16

From u= 0.471 to 2.36, we measure the scaling function. The region of these data values corresponds to 5< beta < 18.

We take the continuum limit using linear extraplation of a<sup>2</sup>.

At the strong coupling region, there is a large scaling violation.



#### The scaling function for each step scaling



#### The scaling function for each step scaling





### <u>Summary</u>

>In low beta, the TPL coupling's behavior is different from SF scheme.

> In our analyses, we estimated the discretization errors in  $O(a^2)$ .

 $\succ$ We study the growth rate of the renormalized coupling.

> The TPL coupling in the case of Nf=12 does not show a signal of IR fixed point.

>Now, we are working in progress to study the lower energy region and have to do some improvement to reduce the systematic error.

#### **Future directions**

Numerical measurement of the anomalous dimension at the IR point the composite operator of fermions is interesting (Del Debbio et al.)

- Study Nf dependence (for example Nf=8 or 16) to study arbitrary Nf, we need Overlap or Domain-wall fermion
- The other gauge group and representation of fermions many people are studying...(Sannino, Yamada, Ohki...)

SU(2) fund.Nf=8 case (Twisted Wilson loop sheme) H.Ohki's talk (17 June pm5:40-) Back up slides

$$\langle O^{i}(x)O^{j}(\mathbf{0})\rangle \sim \sum_{k} \frac{C^{ijk}}{|x|^{\Delta_{i}+\Delta_{j}-\Delta_{k}}}O^{k}$$

#### Step scaling procedure



1. Tune the value of beta (bare coupling) for a small lattice size, which gives the renormalized coupling "u".

2. Carry out the simulation for the twice size of lattice.

3. Take the continuum limit (energy scale  $\mu = 1/2L_0$  )

Notation:

input renormalized coupling :  $u = g^2(1/L)$ output renormalized coupling :  $\sigma(u) = g^2(1/sL)$ 

#### Comparison the continuum extrapolations

