# Two-photon decay of $\pi^0$ from two flavor lattice QCD

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# Motivation

• Why we study  $\pi^0 \rightarrow \gamma^{(*)}\gamma^{(*)}$  process ? Lattice analysis of on- and off-shell photon decay amplitude Anomaly effect plays an important role  $\rightarrow$  resent measurement [Jefferson Lab (2009)]

Non-perturbative study of  $\gamma^* \gamma \rightarrow \pi^0$  transition form factor Comparison with pQCD analysis in deep Euclidean region  $\rightarrow$  CELLO, CLEO and BaBar (e<sup>+</sup>e<sup>-</sup> $\rightarrow \gamma^* \rightarrow$  hadron(+ $\gamma$ ))

Applicable to the calculation of NLO hadronic contribution to muon anomalous g-2

 $\rightarrow$  precision test of the SM, constraint on BSM

# In this talk

Analysis of lattice data in CHPT

Fit the lattice results of  $\pi^0 \rightarrow \gamma \gamma$  transition form factor in low energy

 $\rightarrow$  an analysis of chiral behavior

Comparison with CHPT and VMD

 $\rightarrow$  correspondence to LECs and vector meson mass

Pion loop correction to the  $\pi^0 \rightarrow \gamma \gamma$  amplitude

# CHPT with anomalous term

 Wess-Zumino effective action 12  $\mathcal{L}_{WZW} = -\frac{N_c}{96\pi^2 f_0} \varepsilon^{\mu\nu\alpha\beta} F_{\mu\nu} F_{\alpha\beta} \pi^0$ 11 Wess, Zumino (1971), Witten (1983) <sup>0</sup>→γ/ Decay Width (eV) induced from axial-anomaly. At  $N_c = 3$  and in the massless limit:  $\lim_{m_q \to 0} f_{\pi^0 \gamma \gamma}(0,0) = \frac{1}{4\pi^2 f_0}$ Decay width: 7 LO:  $\Gamma^{\text{LO}}_{\pi^0 \to \gamma\gamma} = 7.73(4) \,\text{eV}$ **NLO:**  $\Gamma^{\text{LO}}_{\pi^0 \to \gamma\gamma} = 7.54(4) \,\text{eV}$ Include SU(2) breaking  $\Gamma_{\pi^0 \to \gamma\gamma}^{\pi^0 - \eta} = 7.93(12) \,\text{eV} \,[\text{Ioffe, Oganesian} \,(2007)]$ CHTP + I/N<sub>c</sub> expansion  $\Gamma_{\pi^0 \to \gamma\gamma}^{1/N_c} = 8.10(8) \text{ eV} [\text{Goity, Bernstein, Hostein, (2002)}]$ 

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### Lattice set-up

- Dynamical overlap fermion (Nf=2) Exact chiral symmetry on the lattice
  - Dynamical configuration generated by JLQCD/TWQCD
  - $16^3 \times 32$ , 1/a=1.67 GeV, Iwasaki gauge action
  - $m_{g} = 0.015, 0.025, 0.035, 0.050$
  - $Z_v = 1.38$  : non-perturbative value (RI-MOM)

# $\tau^{0} \rightarrow \gamma \gamma$ transition form factor on lattice

Definition

 $\int d^4x e^{ip_2 x} e^{ip_1 y} \langle \pi^0 | V_{\nu}^{\text{EM}}(x) V_{\mu}^{\text{EM}}(y) | 0 \rangle = \varepsilon_{\mu\nu\alpha\beta} p_1^{\alpha} p_2^{\beta} f_{\pi^0 \gamma^* \gamma^*}(p_1, p_2)$ 

#### 3-point function,V-V-A (Nf=2)

 $\left\langle \nabla A^3 V_{\nu}^{\mathrm{EM}} V_{\mu}^{\mathrm{EM}} \right\rangle_c (P_1, P_2) = \sum_{x,y} \left\langle \nabla A^3(x) V_{\nu}^{\mathrm{EM}}(y) V_{\mu}^{\mathrm{EM}}(0) \right\rangle_c e^{-iQx} e^{-iP_2y}$ 

Fourier transformed to 4D Euclidean momentum:  $P^2 = -p^2$ (fixed  $P_2 = (0, 0, 0, 2\pi n/N_t)$ )



# $\pi^0 \rightarrow \gamma \gamma$ transition form factor on lattice

#### Extraction of transition form factor

Using conserved axial current and local vector current

 $\rightarrow$  no contact term

□ In the low energy [~  $Q^2 \ll (10m_{\pi})^2 \sim (m_{\pi'})^2$ ], where pion pole is dominant  $\langle \nabla AV_{\nu}V_{\mu}\rangle(P_1, P_2) = \frac{f_{\pi}m_{\pi}^2}{Q^2 + m_{\pi}^2} \varepsilon_{\mu\nu\alpha\beta} P_1^{\alpha} P_2^{\beta} f_{\pi^0\gamma^*\gamma^*}(P_1, P_2) + \cdots$  $\Rightarrow \langle \nabla AV_{\nu}V_{\mu}\rangle(P_1, P_2)/(\varepsilon_{\mu\nu\alpha\beta}P_{1\,\alpha}P_{2\,\beta}) = -\frac{f_{\pi}m_{\pi}^2}{Q^2 + m_{\pi}^2} f_{\pi^0\gamma^*\gamma^*}^{\text{lat}}(P_1, P_2)$ 

#### Momentum assignment

To avoid the lattice artifact due to Lorentz violation,

e.g.  $\varepsilon_{\mu\nu\alpha\beta}P_{1\,\alpha}P_{2\,\beta}(P_{1\,\mu}^2 + P_{2\,\nu}^2)a^2$ , we take a reduced momentum assignment:  $P_1 = \left(0, 0, \frac{2\pi n_3}{L_3}, 0\right), P_2 = \left(0, 0, 0, \frac{2\pi n_4}{L_4}\right)$  at  $\mu = 2, \nu = 1$  $P_1 = \left(0, \frac{2\pi n_2}{L_2}, 0, 0\right), P_2 = \left(0, 0, 0, \frac{2\pi n_4}{L_4}\right)$  at  $\mu = 3, \nu = 1$ 

# Fit function to lattice results

### Based on VMD



### CHPT combined with VMD

Low energy : chiral loops, High energy : tail of vector meson pole LECs corresponds to vector mass and coupling



# NLO in CHPT

$$\begin{aligned} \mathbf{Chiral\ expansion\ at\ O(p^{6})} & \text{Borasoy, Nissler (2004)} \\ \mathcal{L}_{WZW} \epsilon^{\mu\nu\alpha\beta} d^{4}x &= -\frac{\sqrt{2}}{8\pi^{2}f_{0}^{3}} \Big( \operatorname{tr} \partial_{\mu}\pi[\pi, [\pi, v_{\alpha}\partial_{\nu}v_{\beta}]] \\ &- \operatorname{tr} \partial_{\mu}\pi[\pi, \partial_{\nu}v_{\alpha}][\pi, v_{\beta}] - 2i\operatorname{tr} \partial_{\mu}\pi\partial_{\nu}\pi\partial_{\alpha}v_{\beta} + \cdots \Big) \epsilon^{\mu\nu\alpha\beta} d^{4}x \end{aligned}$$

Contribution to transition from factor



## New fit ansatz - based on CHPT

Re-summation of bubble diagram

$$\frac{1}{f_{\pi}^{2}} \longrightarrow \frac{P^{2}}{P^{2} + M^{2}} = \frac{I}{f_{\pi}^{2}} + 2\frac{I}{f_{\pi}^{2}} + \left(2\frac{I}{f_{\pi}^{2}}\right)^{2} + \dots + [\text{counter term}]$$

$$+ \left[\frac{I(P^2;m^2)}{f_\pi^2} + \frac{64\pi^2}{3}P^2w_1\right] \left[1 - 2\left(\frac{I(P^2;m^2)}{f_\pi^2} + \frac{64\pi^2}{3}P^2w_1\right)\right]^{-1}$$

then we have

$$\begin{split} \Lambda_{\pi^{0}\gamma^{*}\gamma^{*}} &= 1 - \frac{2}{f_{\pi}^{2}} \Delta_{\pi} - \frac{256\pi^{2}}{3} m_{\pi}^{2} w_{3} + \frac{1}{2} \Big[ \frac{P_{1}^{2}}{P_{1}^{2} + M^{2}(P_{1}^{2}, m_{\pi}^{2})} + \frac{P_{2}^{2}}{P_{2}^{2} + M^{2}(P_{2}^{2}, m_{\pi}^{2})} \Big],\\ M(P^{2}, m^{2}) &= \Big[ -\frac{I(P^{2}, m^{2})}{2P^{2} f_{\pi}^{2}} - \frac{32\pi^{2}}{3} w_{1} \Big]^{-1/2} \end{split}$$

Last term is same form as VMD ansatz

## New fit ansatz - based on CHPT

### Fitting at



# New fit ansatz - combined with VMD

### One-loop CHPT + VMD



 $w_3$  : LECs at O(p<sup>6</sup>),  $z_{1,2}$  :VMD coupling

# New fit ansatz - combined with VMD

### • Fit results



Decay width

$$\Gamma_{\pi^0} \gamma \gamma = \frac{m_\pi^3 \alpha_e^2}{64\pi^3 f_\pi^2} \Lambda_{\pi^0 \gamma \gamma}^2$$
  
= 8.15(3) eV : CHPT  
= 7.62(6) eV : CHPT+VMD



# Summary

- CHPT analysis has been done with  $\pi^0$  decay to 2 photon
- Both re-summed CHPT and CHPT + VMD fit function well describes lattice data
- Decay width of  $\pi^0$

 $\Gamma_{\pi^0 \gamma \gamma} = 7.89(3)(\pm 37) \,\mathrm{eV}$ 

which is consistent with experiment, but remains large fit uncertainties in our fitting region

### • Future

To reduce sys. error, need small momentum  $\rightarrow$  large lattice,  $24^3 \times 48$  extend to higher momentum as 3 or 4 GeV<sup>2</sup>

- $\rightarrow$  cross check of CLEO, BaBar result
- → provides suggestive information for NLO hadronic contribution to muon g-2

# Backup

Vector meson mass



Comparison with experiment (preliminary)

Brodsky, Lepage (1980)

$$f_{\pi^0 \gamma \gamma^*}(0, Q^2) = \frac{1}{4\pi^2 f_\pi} \frac{1}{1 + Q^2/(8\pi^2 f_\pi^2)} \to \frac{2f_\pi}{Q^2} \left(Q^2 \gg \Lambda^2\right)$$

- pQCD: division from BaBar at  $p_1^2 > 15 \,\mathrm{GeV}^2$
- Our fit result is
  consistent with both
  CLEO and BaBar, but still
  large error.



# Hadronic contribution: NLO

π, η, . . .

Hadronic light-by-light scattering diagram (3 off-shell, I on-shell)
 Direct calculation: hard task !

Indirect calculation: decomposition into two diagram  $(\pi^0, \eta^0, ...) \rightarrow \gamma^* \gamma^*$ 

π, η, ...

μμ

Pion pole dominance model

Focus on pion decay, other meson decay is sub-leading: M. Knecht, A. Nyffeler (2002)

μμ

$$a_{\mu}^{\text{LbyL}} = \int_{p_1} \int_{p_2} \left[ p_1^2 p_2^2 (p_1 + p_2)^2 \right]^{-1} \left[ ((p + p_1)^2 - m_{\mu}^2) ((p - p_2)^2 - m_{\mu}^2 \right]^{-1} \Pi_{\pi^0 \gamma^* \gamma^*} (p_1, p_2, p) \\ \Pi_{\pi^0 \gamma^* \gamma^*} (p_1, p_2, p) \sim \left\{ \frac{f_{\pi^0 \gamma^* \gamma^*} (q_1^2, (q_1 + q_2)^2) f_{\pi^0 \gamma^* \gamma} (q_2^2, 0)}{p_2^2 - m_{\pi}^2}, \frac{f_{\pi^0 \gamma^* \gamma^*} (q_1^2, q_2^2) f_{\pi^0 \gamma^* \gamma} ((q_1 + q_2)^2, 0)}{(p_1 + p_2)^2 - m_{\pi}^2} \right\}$$

Since we have assumed some resonance models (e.g.VMD) in LbyL diagram, then <u>non-perturbative form</u> plays important role.

π, η, . . .

# Comparison with CHPT

• П<sub>V-A</sub><sup>(0)</sup>



CHPT prediction:  $Q^2 \Pi_{V-A}^{(0)}(Q^2) = \frac{f_{\pi}^2 m_{\pi}^2}{Q^2 + m_{\pi}^2}$ well describe lattice data  $\checkmark$ other lattice artifact including  $\Pi_{V-A}$  term is small ((aQ)<sup>2</sup> < 2)