Toward the Nearly Conformal Composite Higgs Mechanism

Julius Kuti

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With Lattice Higgs Collaboration members: Z. Fodor, K. Holland, D. Nogradi, C. Schroeder

Outline

1. Nf=12

- detailed chiral analysis with estimates of serious limitations

2. Nf=2 sextet chiral symmetry breaking

- first steps toward complete chiral analysis (preliminary)
- this model can be settled!

Talk is based on published results last year:

- L. Topology and higher dimensional representations. Published in JHEP 0908:084,2009. e-Print: arXiv:0905.3586 [hep-lat]
- 2. Nearly conformal gauge theories in finite volume. Phys.Lett.B681:353-361,2009. e-Print: arXiv:0907.4562 [hep-lat]
- 3. Chiral properties of SU(3) sextet fermions e-Print: arXiv:0908.2466 [hep-lat]
- 4. Chiral symmetry breaking in nearly conformal gauge theories e-Print: arXiv:0911.2463 [hep-lat] PoS Lattice 2009
- 5. Calculating the running coupling in strong electroweak models e-Print: arXiv:0911.2934 [hep-lat]

and unpublished new sextet analysis soon to be published

Conformal windows in theory space



We only run with N=3 colors

We are supported by the Wuppertal hardware/software infrastructure

Zoltan Fodor Kalman Szabo Sandor Katz

GPU HARDWARE Sandor Katz

GTX 280 Flops: single 1 Tflop, double 80 Gflops Memory 1GB, Bandwidth 141 GBs⁻¹ 230 Watts, \$350

UCSD Tesla cluster ARRA funded by DOE waiting for Fermi cards

USQCD and Teragrid support

Tesla 1060 Flops: single 1 Tflop, double 80 Gflops Memory 4GB, Bandwidth 102 GBs⁻¹ 230 Watts, \$1200

Tesla 1070 Flops: single 4 Tflops, double 320 Gflops Memory 16GB, Bandwidth 408 GBs⁻¹ 900 Watts, \$8000



Evolution of nearly conformal gauge model with volume

What is the finite volume spectrum? How does the running coupling $g^2(L)$ evolve with L?

At small $g^2(L)$ the zero momentum components of the gauge field dominate the dynamics: Born-Oppenheimer approximation

Originally it was applied to pure-gauge system: Luscher, van Baal

SU(3) pure-gauge model: 27 inequivalent vacua

Low excitations of Hamiltonian (Transfer Matrix) scale with will evolve into glueball states for large L $\sim g^{2/3}(L)\,/\,L$

Three scales of dynamics:(1) WF is localized on one vacuum(perturbation theory)(2) tunneling accross vacua on second scale (nonperturbative, calculable)(3) over the barrier: confinement scale(nonperturbative, not calculable)third stage is either broken chiral symmetry or conformal

Quantum vacuum is at minimum of Veff(C) when massless fermions are turned on early work by van Baal, Kripfganz and Michael Fermions develop a gap $\sim \pi/L$ in the spectrum

k=(1,1,1) antiperiodic minimal when l=0 (mod 2π) A=0 ($P_j = 1$)

k=(0,0,0) periodic minimal when $\vec{l} \neq 0$ nontrivial vacua $P_j = \exp(\pm 2\pi i/3)$

Polyakov loop distributions probe the vacua



Check first that simulation reached stage 3 where testing begins in ernest

Nf=12 fundamental representation

(1) Chiral symmetry breaking hypothesis

(2) First crude asymptotic tests include

(a) Goldstone spectrum, parity partner? tantalizing

(b) $M_{\pi} \sim \sqrt{2Bm_q}$ asymptotics difficult

(c) NLO expansion in p regime has question marks

(d) Is F*L large enough? not at Nf=12

(e) Can RMT help? only if F*L is large enough

(3) Use running coupling tests to complement chiral analysis (Kieran Holland's talk)

(4) Check whether conformal picture provides alternative what is it? has to be better then simple power fits of "anomalous dimensions"

Chiral regimes to identify in theory space:



One-loop expansion in our analysis of p-



Leutwyer, Gasser, P. Hasenfratz,

$$M_{\pi}^{2} = M^{2} \left[1 - \frac{M^{2}}{8\pi^{2}N_{f}F^{2}} ln\left(\frac{\Lambda_{3}}{M}\right) \right],$$

$$F_{\pi} = F \left[1 + \frac{N_{f}M^{2}}{16\pi^{2}F^{2}} ln\left(\frac{\Lambda_{4}}{M}\right) \right],$$

Note Nf scaling of pion mass! warning: 2-loop ~ Nf^2 (Bijnens)

$$M_{\pi}(L_s,\eta) = M_{\pi} \left[1 + \frac{1}{2N_f} \frac{M^2}{16\pi^2 F^2} \cdot \widetilde{g}_1(\lambda,\eta) \right], \quad \lambda = ML_s$$
$$F_{\pi}(L_s,\eta) = F_{\pi} \left[1 - \frac{N_f}{2} \frac{M^2}{16\pi^2 F^2} \cdot \widetilde{g}_1(\lambda,\eta) \right],$$

We use staggered action with stout smearing Taste breaking can be included in staggered perturbation theory! structure changing as Nf grows If F*L is not large enough, everything is beginning to break

Nf=12 runs are away from crossover region !



Nf=12 new NLO beta=2.2 chiral analysis in p-regime:



- Consistent with chiral symmetry breaking
- Can this be improved? Hard ...
- What would be the conformal phase analysis?





Slow emergence of asymptotic $M_{\pi} \sim \sqrt{2Bm_q}$ behavior





Nf=12 mq=0.02 rho-A1 splitting pulled out from single correlator with two parity partners

Nf=12 mq=0.015 rho-A1 splitting pulled out from single correlator with two parity partners



$$E_l = \frac{1}{2\theta} l(l+2)$$
 with $l = 0, 1, 2, ...$ rotator spectrum for SU(2)

with $\theta = F^2 L_s^3 (1 + \frac{C(N_f = 2)}{F^2 L_s^2} + O(1/F^4 L_s^4))$ (P. Hasenfratz and F. Niedermayer) there is overall factor $\frac{N_f^2 - 1}{N_f}$ for arbitrary N_f **expansion in 1/F²L²**

 $C(N_f = 2) = 0.45$ expected to grow with N_f

At $FL_s = 0.8$ the correction is 70% and grows with N_f

When expansion collapses in δ – regime, the p-regime analysis needs more scrutiny

Cross checks from several running coupling schemes is important

The Nf=2 sextet model simulations

- (1) Similar check list and warnings as Nf=12 case
- (2) However, Nf=2 makes it easier
- (3) Analysis is preliminary!
- (4) So far quite consistent with chiral symmetry breaking
- (5) Unlike Nf=12 case, no intrinsic barriers with NNLO convergence, F*L squeeze, ...

(6) We are running our couplings (Kieran Holland's talk)

The Nf=2 sextet model simulations



NLO chiral fitting procedure of the Nf=2 sextet model simulations



δ

$$M_{\pi}^{2} = M^{2} \left[1 - \frac{M^{2}}{8\pi^{2}N_{f}F^{2}} ln\left(\frac{\Lambda_{3}}{M}\right) \right], \qquad M^{2} = 2Bm_{q}$$

$$F_{\pi} = F \left[1 + \frac{N_{f}M^{2}}{16\pi^{2}F^{2}} ln\left(\frac{\Lambda_{4}}{M}\right) \right],$$

$$M_{\pi}(L_{s}, \eta) = M_{\pi} \left[1 + \frac{1}{2N_{f}} \frac{M_{\pi}^{2}}{16\pi^{2}F_{\pi}^{2}} \cdot \widetilde{g}_{1}(\lambda, \eta) \right],$$
input
$$F_{\pi}(L_{s}, \eta) = F_{\pi} \left[1 - \frac{N_{f}}{2} \frac{M_{\pi}^{2}}{16\pi^{2}F_{\pi}^{2}} \cdot \widetilde{g}_{1}(\lambda, \eta) \right],$$

$$\lambda = M_{\pi}L, \qquad \eta = \frac{L_T}{L_S}$$

 $\tilde{g}(\lambda,\eta)$ is a shape-dependent expansion in terms of infinite series of Bessel functions Next-to-leading-order fitting procedure of the chiral analysis



Our simulations of the model are not accurate for NNLO analysis with taste breaking











Conclusions and Outlook

- Nf=12 is consistent with chiral symmetry breaking
 to overcome limitations would require big resources
 worth it?
- 2. Nf=2 sextet consistent with chiral symmetry breaking
 easier to control F*L and Nf expansions within our resources
 full chiral analysis can be done
- 3. Important to complement with running couplings