

# Lattice QCD analysis for the instantaneous interquark potential in the generalized Landau gauge

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## Quantum Chromodynamics — $SU(N_c)$ gauge theory

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{2} \text{Tr} G_{\mu\nu} G^{\mu\nu} + \bar{q} (i\gamma_\mu D^\mu - m) q$$

For actual calculations  $\Rightarrow$  **fix** the gauge

- Landau gauge
- Coulomb gauge
- Maximally Abelian gauge
- ...

For example

physical phenomena  $\Rightarrow$  gauge **invariant**

- confinement
- chiral symmetry breaking
- hadron mass
- ...

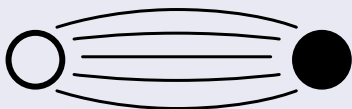
But physical pictures depend on gauges.

# Confinement in the Various Gauge

## Confinement in the Maximally Abelian gauge

**Dual-superconductor picture** — Nambu, 't Hooft, Mandelstam

- QCD vacuum  $\Rightarrow$  Dual superconductor  $\Rightarrow$  Dual Meissner effect



squeezing of the color flux tube



linear potential between quarks

## Confinement in the Landau gauge

**Kugo-Ojima criterion** — T.Kugo, and I.Ojima, Suppl.Prog.Theor. Phys. **66**, 1-130(1979).

If ghost propagator  $D_G(p^2)$  becomes more singular than free propagator

$$\lim_{p^2 \rightarrow 0} p^2 D_G(p^2) = \infty$$

then

- BRS singlet  $\Rightarrow$  color singlet

# Confinement in the Coulomb Gauge

## — Gribov-Zwanziger's Scenario

— V.Gribov, Nucl. Phys. **B139**, 1(1978); D.Zwanzgier, Nucl.Phys. **B518**, 237(1998).

### QCD Hamiltonian in the Coulomb gauge

$$H = \frac{1}{2} \int d^3x \left( \vec{E}^{a,tr} \cdot \vec{E}^{a,tr} + \vec{B}^a \cdot \vec{B}^a \right) + \frac{1}{2} \int d^3x d^3y \rho^a(x) K^{ab}(x, y) \rho^b(y)$$

$K \equiv [M^{-1}(-\nabla^2)M^{-1}]_{x,y}^{ab}$  : instantaneous Coulomb propagator

$M$  : Faddeev-Popov operator

$\rho$  : color charge density

### Coulomb energy part — $\rho K \rho$

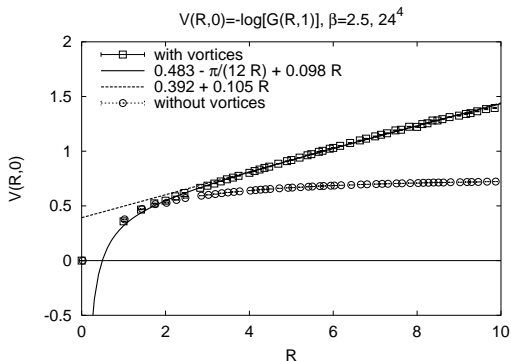
- Gribov horizon  $\Leftarrow$  zero eigenvalues of  $M$
- near Gribov horizon  $\Rightarrow$  **strongly enhancement** of the Coulomb energy
- $\Rightarrow$  Confinement force

# Instantaneous Potential in the Coulomb Gauge

numerical result of the Coulomb energy

## Instantaneous potential

$$V(R, 0) = -\log[\text{Tr}\langle U_4(0, 0)U_4^\dagger(R, 0)\rangle]$$



J.Greenstie, S.Olejnik, PRD.67, 094503(2003).

- linearly rising potential
- 2-3 times larger than physical string tension
- Zwanziger's inequality

D.Zwanziger, PRL.90, 102001(2003).

$$V_{\text{phys}}(R) \leq V_{\text{Coul}}(R)$$

## The role of gluons is changed according to gauges

- Landau gauge
  - $A_\mu^a \Rightarrow$  equal roles
- Coulomb gauge
  - spatial component  $\mathbf{A} = (A_1, A_2, A_3) \Rightarrow$  dynamical variable
  - temporal component  $A_0 \Rightarrow$  potential
- Maximally Abelian gauge
  - diagonal part  $A_\mu^3, A_\mu^8 \Rightarrow$  dominant
  - off-diagonal part  $\Rightarrow$  inactive — K.Amemiya, and H.Suganuma, PRD.60, 114509 (1999)

# Gauge and Properties of Gluons

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In this talk, we discuss the gluon properties from **Landau** gauge toward **Coulomb** gauge.

# Landau Gauge and Coulomb Gauge

## Landau gauge

$$\partial_\mu A_\mu = 0$$

global form (in the Euclidean metric)

$$\text{minimize } R \equiv \int d^4x \text{Tr} \{A_\mu(x)A_\mu(x)\}$$

- Lorentz covariance is kept
- gauge field fluctuation  $\Rightarrow$  **suppressed**

## Coulomb gauge

$$\nabla \cdot \mathbf{A} = 0$$

- Lorentz covariance is **broken**  
 $A_\mu \rightarrow A_0$  and  $\mathbf{A}$
- compatible with **canonical quantization**



## Similarity of gauge fixing condition

- Landau gauge

$$\partial_1 A_1 + \partial_2 A_2 + \partial_3 A_3 + \partial_4 A_4 = 0$$

- Coulomb gauge

$$\partial_1 A_1 + \partial_2 A_2 + \partial_3 A_3 = 0$$

# Generalization of the Landau Gauge

## generalized Landau gauge ( $\lambda$ -gauge)

— cf. C. Bernard, *et al.*, Nucl.Phys.B (Proc.Suppl.) 17 (1990) 593; 20 (1991) 410.

$$\partial_i A_i + \lambda \partial_4 A_4 = 0$$

- $\lambda = 0 \Rightarrow$  Coulomb gauge
- $\lambda = 1 \Rightarrow$  Landau gauge
  
- by varying  $\lambda$ -parameter :  $\lambda = 1 \rightarrow 0$
- we can analyze continuous change of gluon properties
- from Landau gauge  $\Rightarrow$  Coulomb gauge

## on lattice QCD

maximize

$$R_\lambda[U] \equiv \sum_x \sum_{i=1}^3 \text{Re Tr } U_i(x) + \lambda \sum_x \text{Re Tr } U_4(x)$$

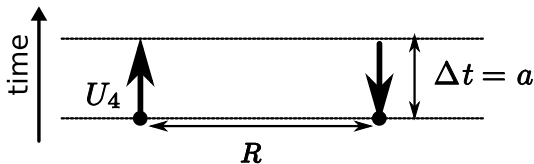
using a gauge transformation of link-variables

# “Instantaneous” Interquark Potential

## Definition of the “instantaneous” potential

$$V(R) \equiv -\frac{1}{a} \log \text{Tr} \langle U_4(\vec{x}, a) U_4^\dagger(\vec{y}, a) \rangle, \quad R = |\vec{x} - \vec{y}|$$

$a$ : lattice spacing

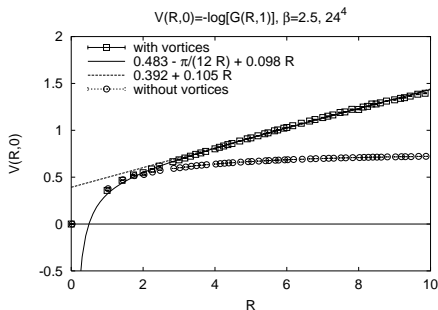


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⇐ Coulomb gauge result

J.Greensite, S.Olejnik, PRD.67, 094503(2003).

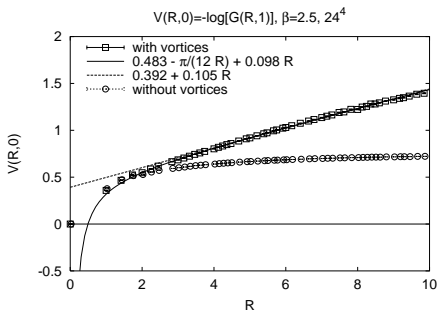
Coulomb + linear potential  
(2-3 times larger than physical  
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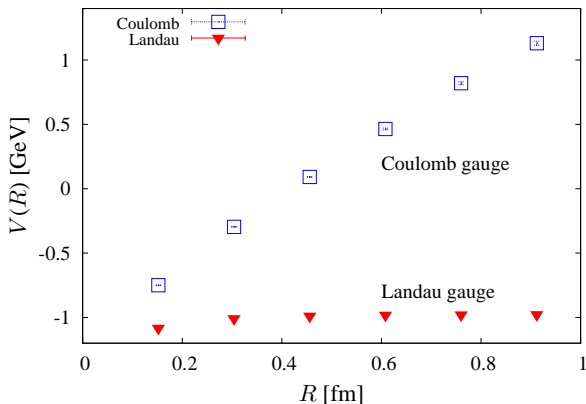
⇒ Analyze **gauge dependence** of  $V(R)$  in the generalized Landau gauge.

# Lattice QCD Calculation

lattice QCD result ( $\beta = 5.8$ ,  $a = 0.152$  fm,  $16^4$  lattice)

“instantaneous” potential

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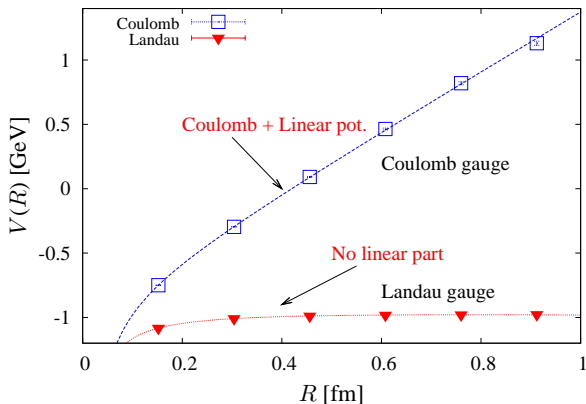
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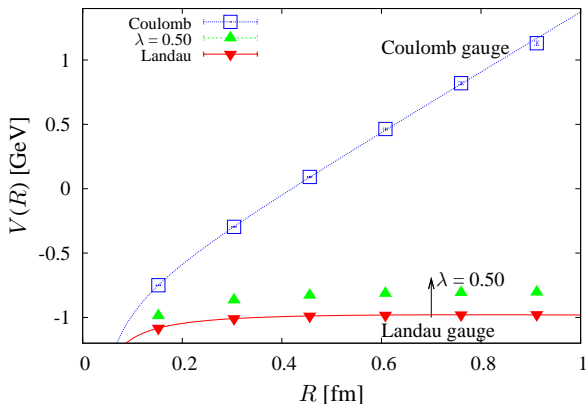
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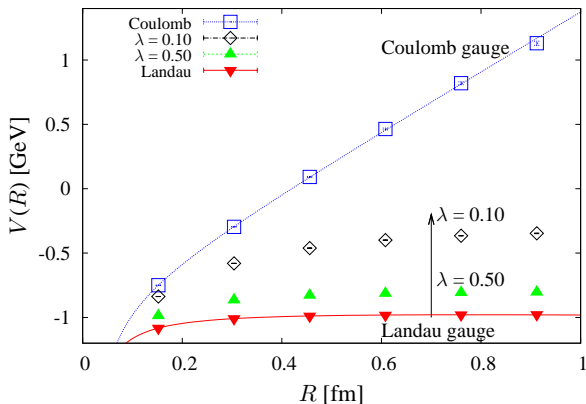


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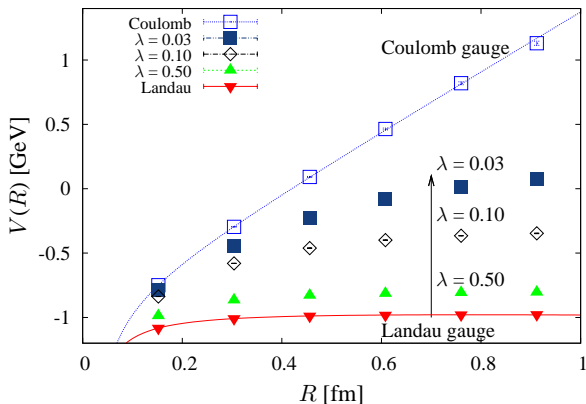
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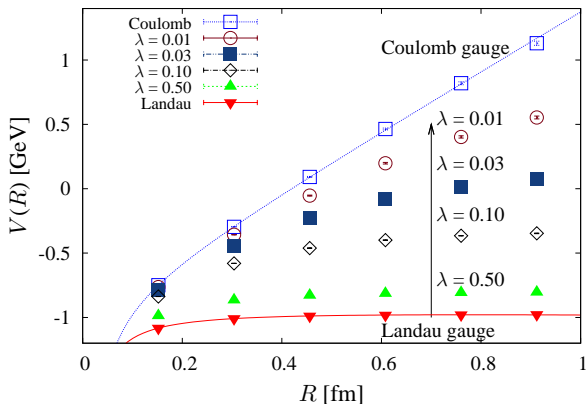
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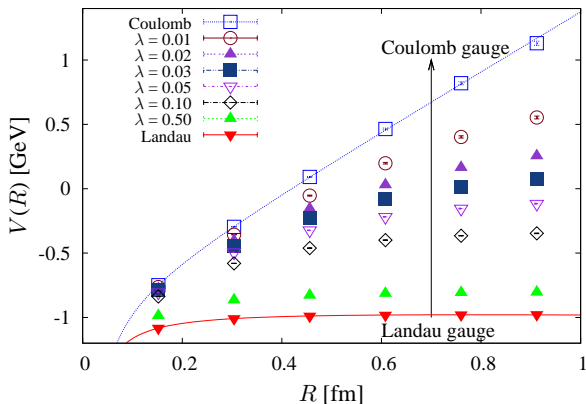
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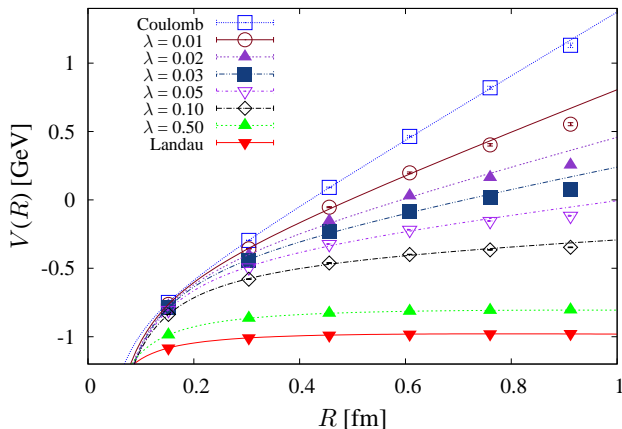


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Coulomb + Linear pot.  
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- Landau gauge ( $\lambda = 1$ )  
No linear part  
 $\sigma \sim 0$  [GeV/fm]
- large **gap** between Landau and Coulomb
- $\lambda = 1 \rightarrow 0$   
potential is changing **continuously**

# Instantaneous Potential Analysis

## Coulomb plus linear functional form

$$V(R) = -\frac{A}{R} + \sigma R + \text{Const.}$$

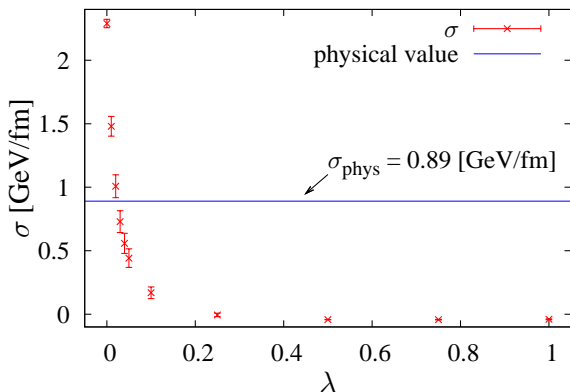


- Landau and Coulomb gauge is connected **continuously**
- Coulomb + linear pot. works well for  $R < 0.8$  fm

# String Tension in the Generalized Landau Gauge

## String tension of the instantaneous potential

- Coulomb gauge ( $\lambda = 0$ )  $\Rightarrow \sigma \sim 2.5\sigma_{\text{phys}}$  over confinement
- Landau gauge ( $\lambda = 1$ )  $\Rightarrow \sigma \sim 0$  no linear part

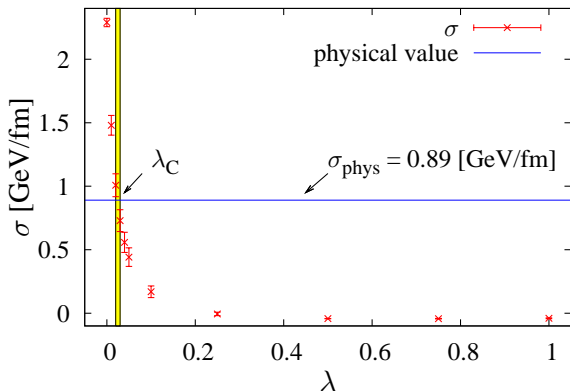


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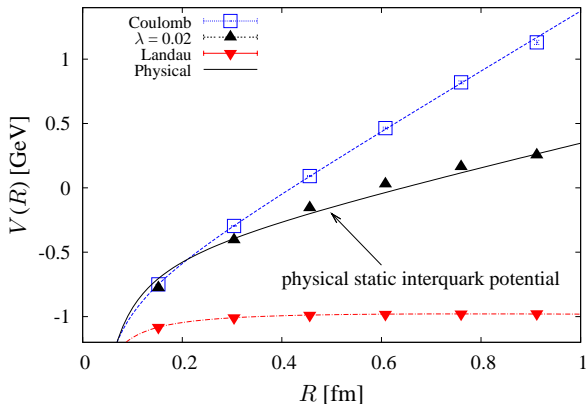


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- $\lambda_C \sim 0.02$   
 $\sigma(\lambda) \simeq \sigma_{\text{phys}}$

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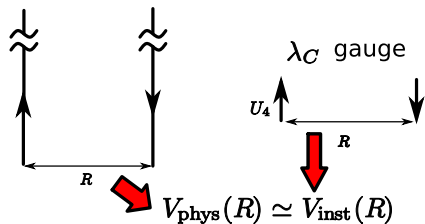
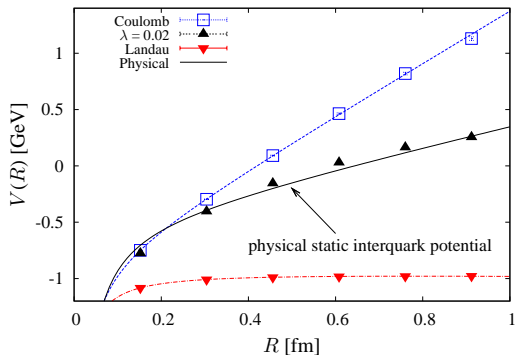
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- $\lambda_C \sim 0.02$   
 $\sigma(\lambda) \simeq \sigma_{\text{phys}}$
- in  $\lambda_C$  gauge  
instantaneous potential reproduces the physical interquark potential



# $\lambda_C$ gauge

in  $\lambda_C$  gauge

instantaneous potential reproduces the physical interquark potential

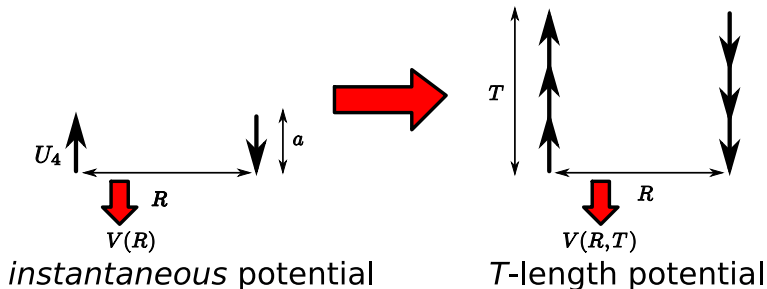


useful for modeling QCD ?

# T-length Potential

## Instantaneous potential

$$V(R) \equiv -\frac{1}{a} \log \text{Tr} \langle U_4(\vec{x}, a) U_4^\dagger(\vec{y}, a) \rangle, \quad R = |\vec{x} - \vec{y}|$$



## T-length potential

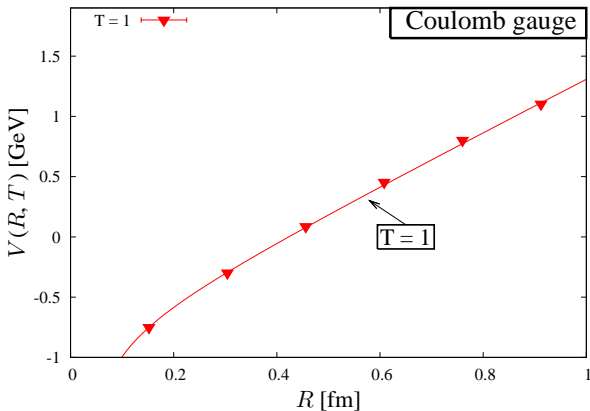
$$V(R, T) \equiv -\frac{1}{T} \log \text{Tr} \langle L(\vec{x}, T) L^\dagger(\vec{y}, T) \rangle, \quad R = |\vec{x} - \vec{y}|$$

Polyakov line:  $L(\vec{x}, T) \equiv U_4(\vec{x}, a) U_4(\vec{x}, 2a) \cdots U_4(\vec{x}, T)$

# T-length Potential in the Coulomb Gauge

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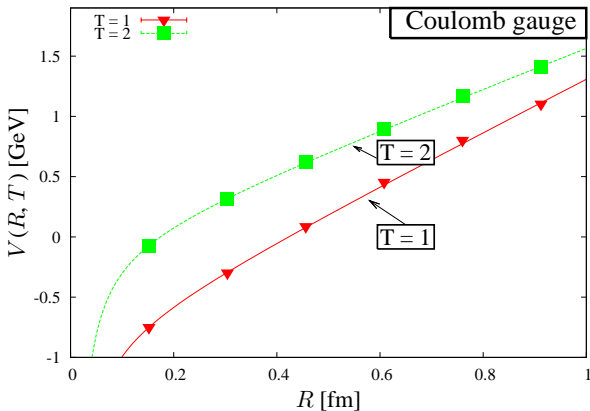
- Coulomb + linear

for SU(2) Coulomb gauge J.Greensite, S.Olejnik,  
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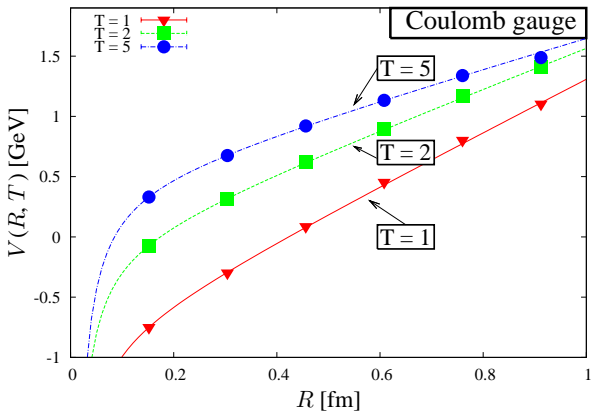
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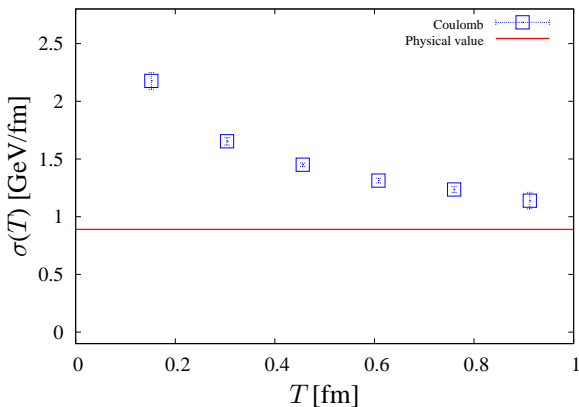
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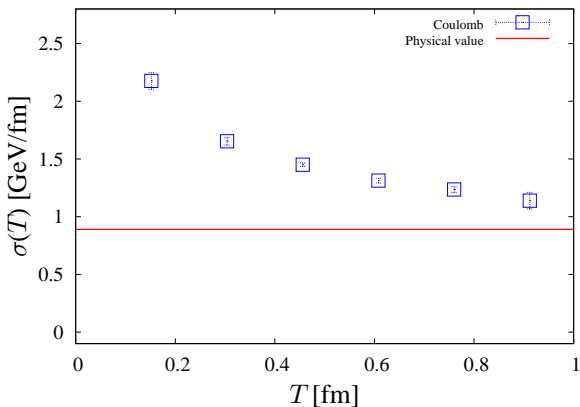
- Coulomb + linear
- string tension is changed

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# $T$ -length String Tension in the Generalized Landau Gauge

## $T$ -length potential analysis

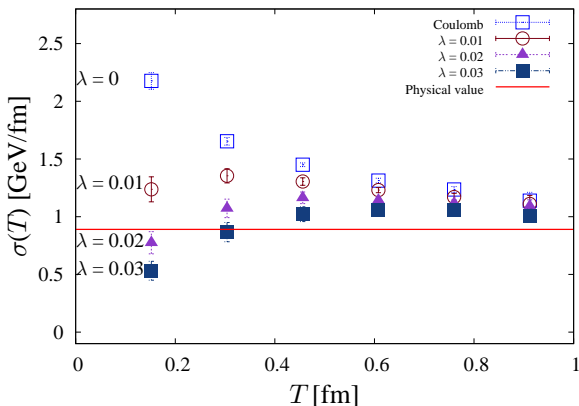
$$V_\lambda(R, T) = -\frac{A_\lambda(T)}{R} + \sigma_\lambda(T)R + \text{Const.}$$



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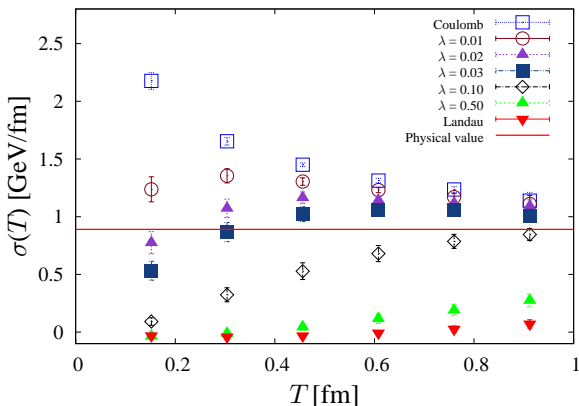
- Near the Coulomb gauge ( $\lambda \sim 0$ )  
 $\sigma_\lambda(T) \rightarrow \sigma_{\text{phys}}$   
 $T \sim 1$  [fm]



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- Near the Coulomb gauge ( $\lambda \sim 0$ )  
 $\sigma_\lambda(T) \rightarrow \sigma_{\text{phys}}$   
 $T \sim 1$  [fm]
- $T \rightarrow \infty$   
 $\sigma_\lambda(T) \rightarrow \sigma_{\text{phys}}$   
for all  $\lambda$  ?

# Summary

We have analyzed the instantaneous potential in the generalized Landau gauge ( $\lambda$ -gauge)

- The potential changes continuously from Landau toward Coulomb gauge.
- Coulomb + linear ( $V(R) = -A/R + \sigma R + \text{Const.}$ )
- String tension is drastically changed near the Coulomb gauge.
- In specified  $\lambda$ -parameter, physical interquark potential is reproduced

We also have analyzed  $T$ -length potential.

- Coulomb + linear potential also works well.
- The asymptotic value seems to coincide with the physical value near the Coulomb gauge.