

# The chiral and angular momentum content of the $\rho$ -meson

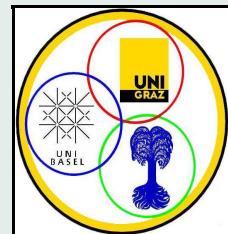
Markus Limmer

L. Ya. Glozman, C. B. Lang

Karl-Franzens-Universität, Graz

LATTICE 2010, June 18

[LYG,CBL,ML PRL103(2009)121601, FewBodySyst47(2010)91]



# Introduction

- Composition of hadronic states in QFT is a subtle issue
- In non-relativistic approaches the concept of a wave function and a complete basis of states is well defined
- BUT: no suitable way of describing a single hadron beyond the ground state in QFT (scattering state)
- In LQCD studies of hadron limited to spectroscopy, axial couplings, form factors, etc.
- More demanding topic: composition of a hadron  $\Rightarrow$  understanding hadron structure in an *ab initio* QCD calculation
- Many Fock components present in a hadron, but only a few are dominant
- Best suited tool: variational method [Michael 1985; Lüscher,Wolff 1990]

# Chiral classification and basis transformation

- Two possible local interpolators with the  $\rho$  quantum numbers  $I, J^{PC} = 1, 1^{--}$  [Cohen, Ji 1997; LYG 2004, 2007]

vector current:  $O_\rho^V = \bar{q} \gamma_i \boldsymbol{\tau} q \rightsquigarrow (0, 1) \oplus (1, 0)$

pseudo-tensor operator:  $O_\rho^T = \bar{q} \sigma^{0i} \boldsymbol{\tau} q \rightsquigarrow (1/2, 1/2)_b$

- Both are a complete set  $\Rightarrow$  signal for chiral symmetry breaking if both couple
- Transformation to angular momentum basis [Nefediev, LYG 2007]

$$\begin{aligned}|(0, 1) \oplus (1, 0); 1\ 1^{--}\rangle &= \sqrt{\frac{2}{3}} |1; {}^3S_1\rangle + \sqrt{\frac{1}{3}} |1; {}^3D_1\rangle \\|(1/2, 1/2)_b; 1\ 1^{--}\rangle &= \sqrt{\frac{1}{3}} |1; {}^3S_1\rangle - \sqrt{\frac{2}{3}} |1; {}^3D_1\rangle\end{aligned}$$

# Variational method

- Representations  $(0, 1) \oplus (1, 0)$  and  $(1/2, 1/2)_b$  form a complete (and orthogonal w.r.t. chiral group) basis  $\Rightarrow$  usage of variational method
- Possibility to study mixing of the two representations
- Cross-correlation matrix

$$C_{ij}(t) = \langle O_i(t) O_j^\dagger(0) \rangle = \sum_{n=1}^N a_i^{(n)} a_j^{(n)*} \exp(-E^{(n)} t)$$

- Overlap of  $O_i$  with physical state  $|n\rangle$  described by  $a_i^{(n)} = \langle 0 | O_i | n \rangle$
- Solve GEVP and get eigenvectors  $u_j^{(n)}$
- Ratios of couplings of different lattice operators to physical state  $|n\rangle$  via

$$\frac{a_i^{(n)}}{a_j^{(n)}} = \frac{C_{ik}(t) u_k^{(n)}}{C_{jk}(t) u_k^{(n)}}$$

# Smearing and the resolution scale

- Hadron decomposition depends on resolution scale
- Point-like interpolators  $\Rightarrow$  resolution scale given by lattice spacing  $a$
- Interested in low (IR) resolution scales (i. e., large  $a$ )  $\Rightarrow$  huge lattice artifacts
- Gauge-invariant smearing of quark fields to size  $R$  such that  $R/a \gg 1$
- Substitute local interpolator by interpolator (with same quantum numbers) using smeared quark fields
- Want Gaussian profile of smeared fields  $\Rightarrow$  Jacobi smearing [Güsken et al. 1989; Best et al. 1997]

$$S = MS_0, \quad M = \sum_{n=0}^N \kappa^n H^n$$

- Hopping term  $H = \sum_{k=1}^3 [U_k(x, t)\delta_{x+\hat{k}, y} + U_k^\dagger(x - \hat{k}, t)\delta_{x-\hat{k}, y}]$

# Simulation details

- Most details presented at talk of G. P. Engel (Monday, June 14) and in [Gattringer et al. 2009; Engel et al. 2010]
- Two mass-degenerate Chirally Improved fermions [Gattringer 2001; Gattringer,Hip,CBL 2001] and Lüscher-Weisz gauge action [Lüscher,Weisz 1985]
- Three different  $16^3 \times 32$  lattices
- Jacobi smeared sources/sinks  $\Rightarrow$  narrow ( $n$ ) and wide ( $w$ ) profile with widths of  $\approx 0.34$  fm and  $\approx 0.67$  fm

set	$\beta_{LW}$	$m_0$	#conf	$a$ [fm]	$m_\pi$ [MeV]	$m_\rho$ [MeV]	$m_{\rho'}$ [MeV]
A	4.70	-0.050	200	0.150(1)	525(7)	914(10)	1881(63)
B	4.65	-0.060	200	0.150(1)	470(4)	872(11)	1905(71)
C	4.58	-0.077	200	0.144(1)	322(5)	791(15)	1799(72)

# Interpolators

$$O_p^V = \bar{u}_p \gamma_i d_p$$

$$O_n^V = \bar{u}_n \gamma_i d_n$$

$$O_w^V = \bar{u}_w \gamma_i d_w$$

$$O_p^T = \bar{u}_p \gamma_i \gamma_t d_p$$

$$O_n^T = \bar{u}_n \gamma_i \gamma_t d_n$$

$$O_w^T = \bar{u}_w \gamma_i \gamma_t d_w$$

$\gamma_i$  ... matrix in spatial direction

$\gamma_t$  ... matrix in (Euclidean) time direction

# Interpolators

$$O_p^V = \bar{u}_p \gamma_i d_p \rightsquigarrow |(0, 1) \oplus (1, 0); 1\ 1^{--}\rangle$$

$$O_n^V = \bar{u}_n \gamma_i d_n \rightsquigarrow |(0, 1) \oplus (1, 0); 1\ 1^{--}\rangle$$

$$O_w^V = \bar{u}_w \gamma_i d_w \rightsquigarrow |(0, 1) \oplus (1, 0); 1\ 1^{--}\rangle$$

$$O_p^T = \bar{u}_p \gamma_i \gamma_t d_p \rightsquigarrow |(1/2, 1/2)_b; 1\ 1^{--}\rangle$$

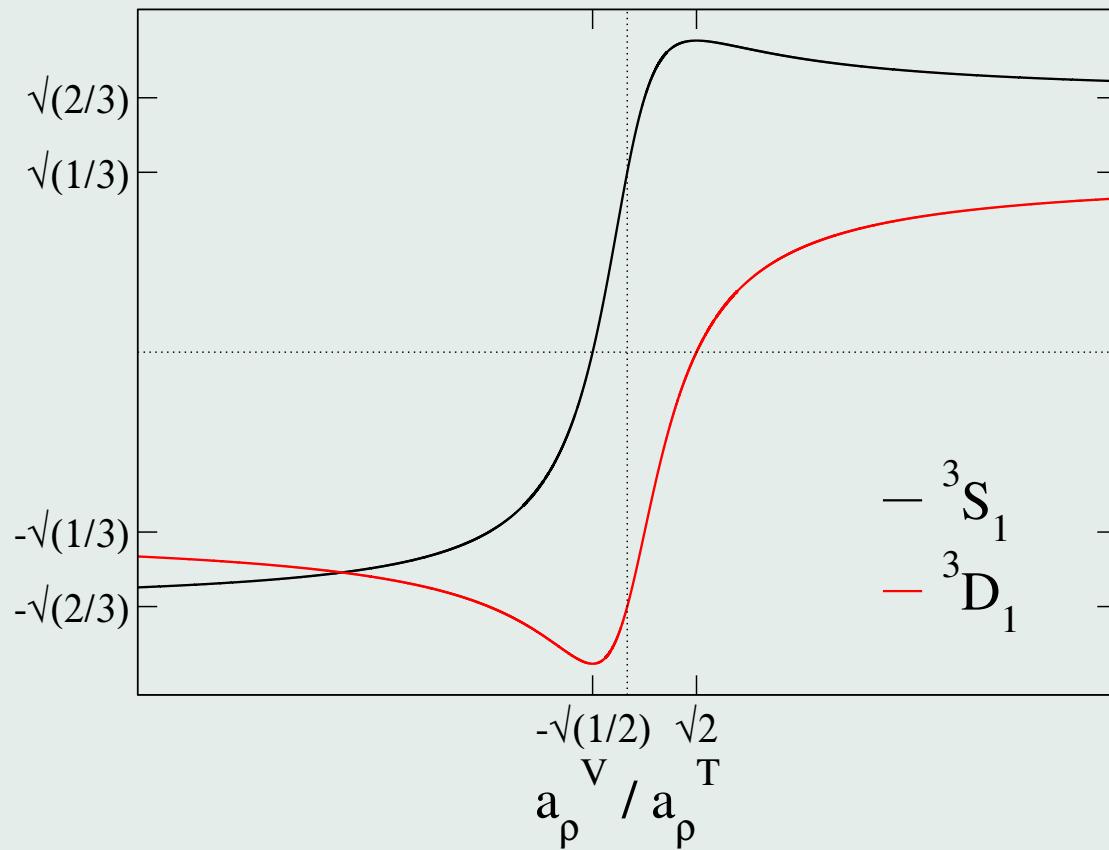
$$O_n^T = \bar{u}_n \gamma_i \gamma_t d_n \rightsquigarrow |(1/2, 1/2)_b; 1\ 1^{--}\rangle$$

$$O_w^T = \bar{u}_w \gamma_i \gamma_t d_w \rightsquigarrow |(1/2, 1/2)_b; 1\ 1^{--}\rangle$$

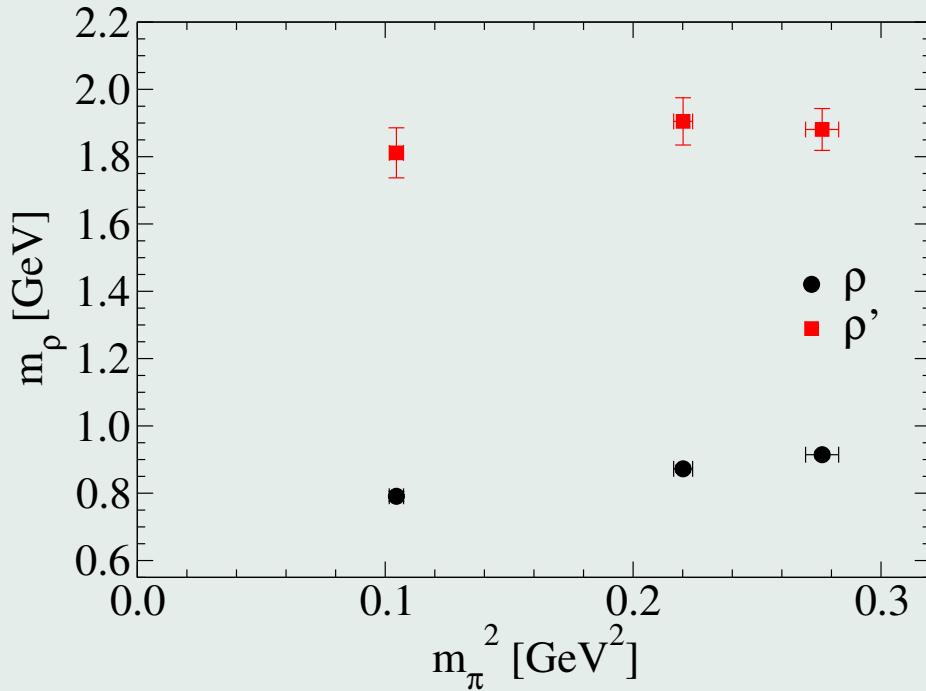
$\gamma_i$  ... matrix in spatial direction

$\gamma_t$  ... matrix in (Euclidean) time direction

# Coupling “strength”

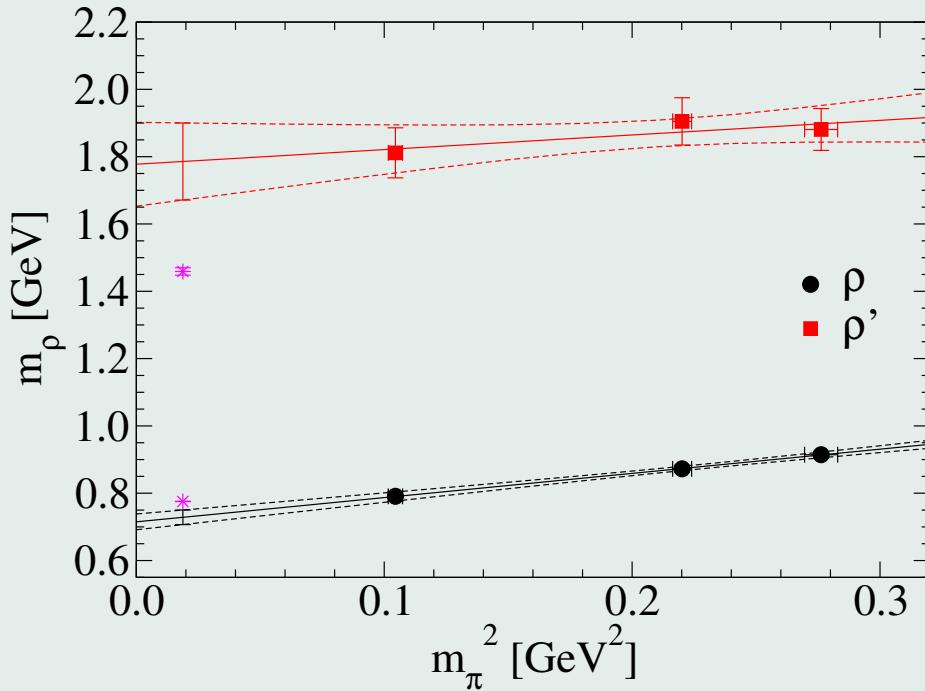


# Vector meson mass



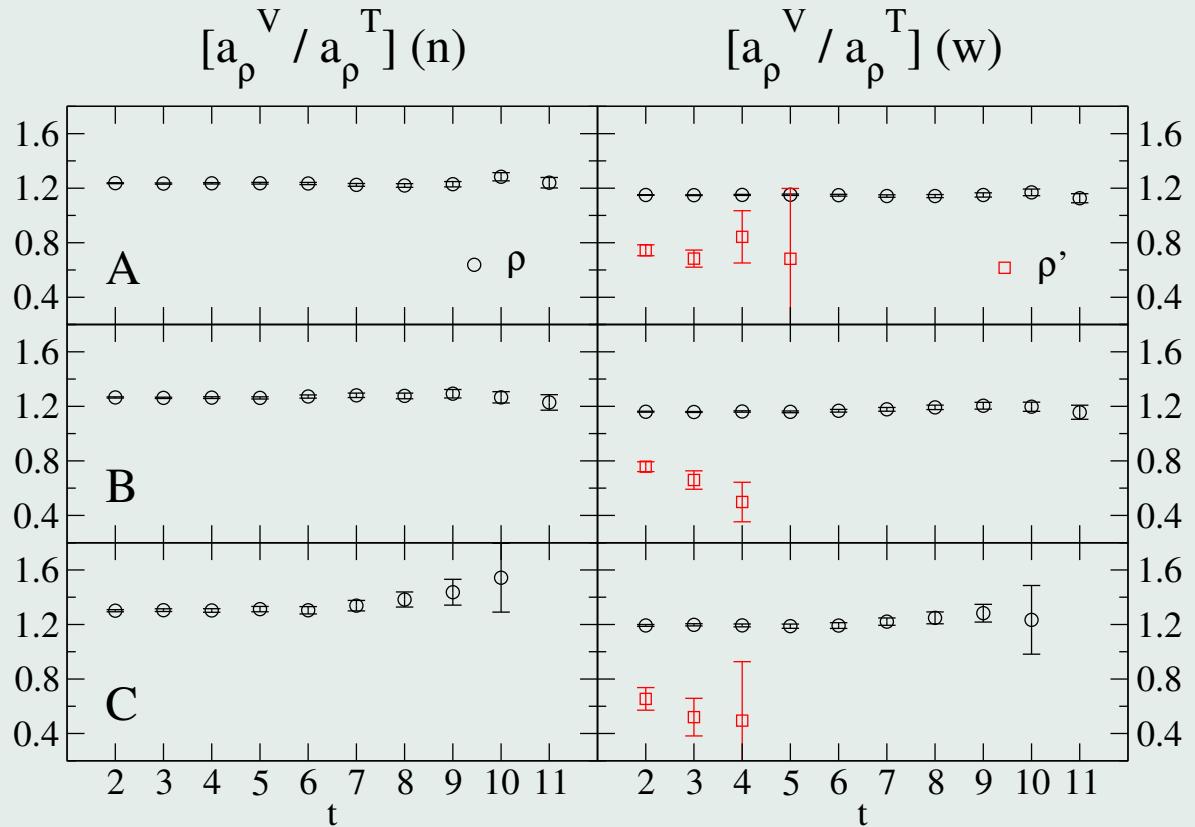
- Masses extracted from  $4 \times 4$  matrix, including  $O_{n/w}^V, O_{n/w}^T$

# Vector meson mass

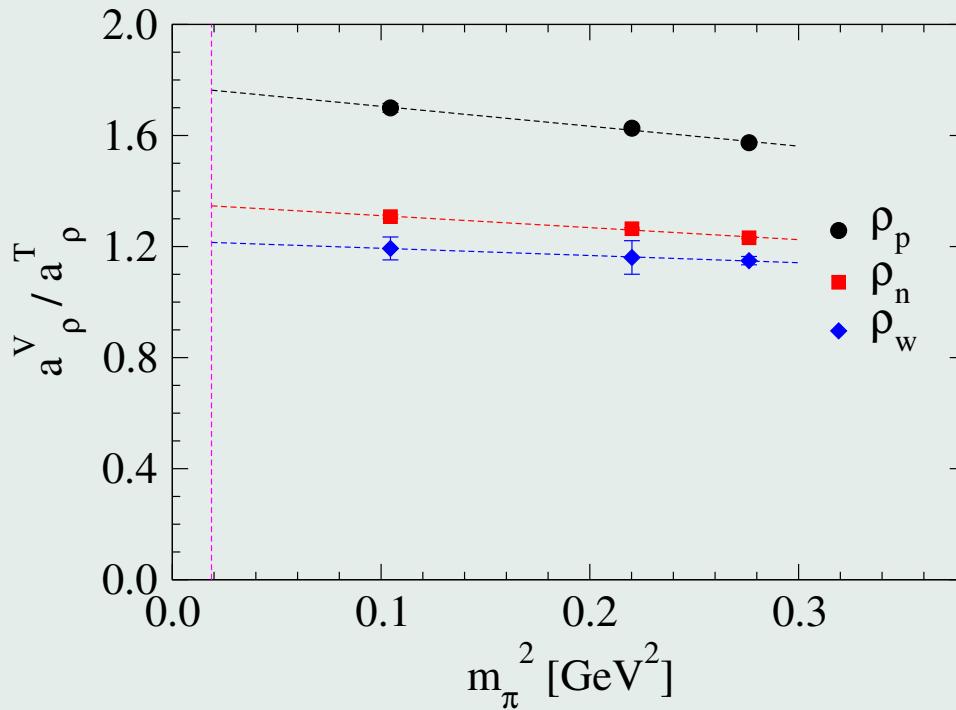


- Masses extracted from  $4 \times 4$  matrix, including  $O_{n/w}^V, O_{n/w}^T$
- “Chiral extrapolation” linear in  $m_\pi^2$

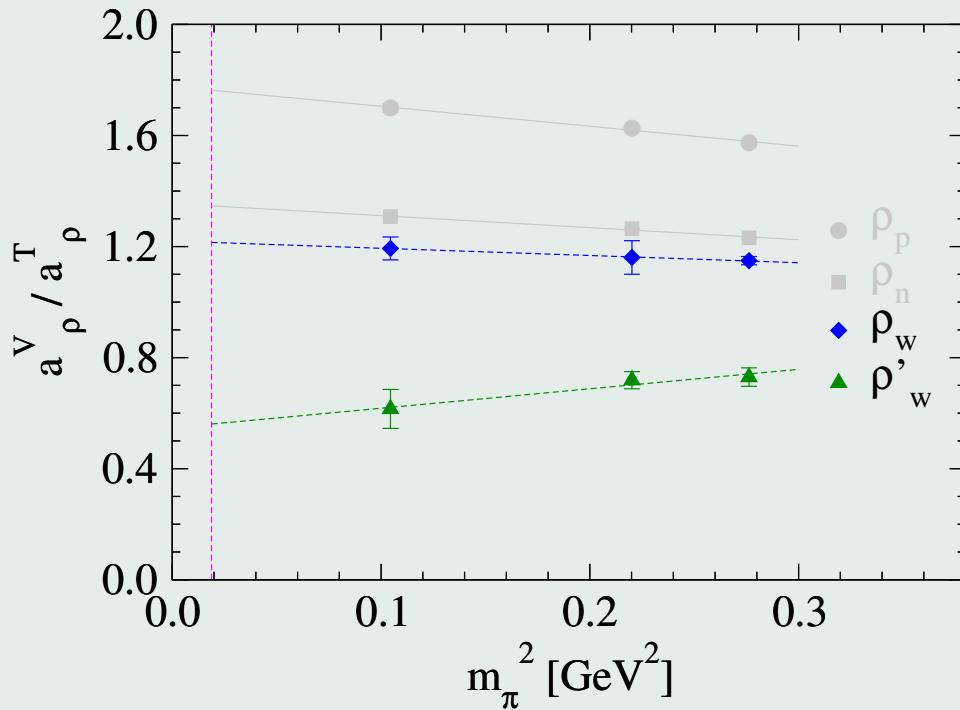
# Stability of the $a$ -ratios



# Vector meson decomposition



# Vector meson decomposition



# Vector meson decomposition

- Can read off fractions of  $|1; {}^3S_1\rangle$  wave and  $|1; {}^3D_1\rangle$  wave from  $a$ -ratios
- Our results:

$$\rho_p: \frac{(0, 1) \oplus (1, 0)}{(1/2, 1/2)_b} = 1.763 \quad \Rightarrow \quad 0.9950|1; {}^3S_1\rangle + 0.09943|1; {}^3D_1\rangle$$

$$\rho_n: \frac{(0, 1) \oplus (1, 0)}{(1/2, 1/2)_b} = 1.347 \quad \Rightarrow \quad 0.9997|1; {}^3S_1\rangle - 0.02328|1; {}^3D_1\rangle$$

$$\rho_w: \frac{(0, 1) \oplus (1, 0)}{(1/2, 1/2)_b} = 1.215 \quad \Rightarrow \quad 0.9973|1; {}^3S_1\rangle - 0.07311|1; {}^3D_1\rangle$$

$$\Rightarrow \rho \approx |1; {}^3S_1\rangle$$

# Vector meson decomposition

- Can read off fractions of  $|1; {}^3S_1\rangle$  wave and  $|1; {}^3D_1\rangle$  wave from  $a$ -ratios
- Our results:

$$\rho_p: \frac{(0, 1) \oplus (1, 0)}{(1/2, 1/2)_b} = 1.763 \quad \Rightarrow \quad 0.9950|1; {}^3S_1\rangle + 0.09943|1; {}^3D_1\rangle$$

$$\rho_n: \frac{(0, 1) \oplus (1, 0)}{(1/2, 1/2)_b} = 1.347 \quad \Rightarrow \quad 0.9997|1; {}^3S_1\rangle - 0.02328|1; {}^3D_1\rangle$$

$$\rho_w: \frac{(0, 1) \oplus (1, 0)}{(1/2, 1/2)_b} = 1.215 \quad \Rightarrow \quad 0.9973|1; {}^3S_1\rangle - 0.07311|1; {}^3D_1\rangle$$

$$\rho'_w: \frac{(0, 1) \oplus (1, 0)}{(1/2, 1/2)_b} = 0.5607 \quad \Rightarrow \quad 0.9029|1; {}^3S_1\rangle - 0.4298|1; {}^3D_1\rangle$$

$$\Rightarrow \rho \approx |1; {}^3S_1\rangle$$

$$\Rightarrow \rho' \not\approx |1; {}^3S_1\rangle$$

# Summary

- Task: hadron composition from an *ab initio* QCD calculation
- Here: vector meson  $\rho$  ( $I, J^{PC} = 1, 1^{--}$ ) was used
- Two different types of interpolators can be used:  $O_\rho^V$ ,  $O_\rho^T$
- Variational method applied
- Used smeared quark fields to vary resolution scale
- Simulation carried out on three  $16^3 \times 32$  lattices with  $a \approx 0.15$  fm and  $m_\pi \approx 320 - 520$  MeV
- Results: ground state mostly  $|1; {}^3S_1\rangle$ , but first excitation has significant admixture of  $|1; {}^3D_1\rangle$  (contradictory to quark model!)

**Thank you!**