

# Mass Anomalous Dimension and Running of the Coupling in $SU(2)$ with Six Fundamental Fermions

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# Motivation

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- Most work on (near-)conformal theories have concentrated on  $SU(3)$  with fundamental fermions or those with higher representations
- $SU(2)$  with fundamental fermions requires fewer fermions to be conformal than  $SU(3)$
- Combined with the small size of the representation this can have a small naïve  $S$  parameter
- Six fermions places us between a number of predictions for the lower end of the conformal window:
  - Loss of Asymptotic Freedom: 11
  - Ladder Approximation: 7.86
  - All Orders Beta Function: 7.33 / 5.5
  - UV / IR degrees of Freedom: 4.74

# Walking Technicolor

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- The original Technicolor theories have a tension between larger fermion masses and small FCNCs
- Walking Technicolor circumvents this by having the underlying gauge theory run slowly between the TC and ETC scales
- This results in a large value for  $\langle \bar{\Psi}\Psi \rangle_{ETC}$  avoiding the tension

$$\langle \bar{\Psi}\Psi \rangle_{ETC} \sim \langle \bar{\Psi}\Psi \rangle_{ETC} \left( \frac{M_{ETC}}{\Lambda_{TC}} \right)^\gamma$$

- A viable theory must not only run slowly but must have a large anomalous dimension

# Schrödinger Functional

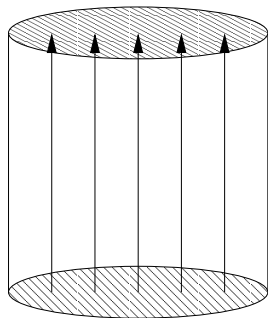
- The SF formalism replaces the (anti-)periodic boundary conditions of standard lattice formulations with Dirichlet boundaries in the time-like direction

$$U(x, k)|_{t=0} = \exp[\eta\tau_3 a/iL],$$

$$U(x, k)|_{t=L} = \exp[(\pi - \eta)\tau_3 a/iL]$$

$$P_+ \psi = 0, \quad \bar{\psi} P_- = 0 \quad \text{at } t = 0,$$

$$P_- \psi = 0, \quad \bar{\psi} P_+ = 0 \quad \text{at } t = L$$



with

$$P_{\pm} = 1/2(1 \pm \gamma_0)$$

- The periodicity is maintained in the space-like directions
- Classically this is equivalent to a background Chromoelectric Flux

## SF Coupling

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- We can use the effective action  $S$  of the theory to define a coupling
- We measure the expectation of derivative of  $S$  as this a better observable

$$\bar{g}^2 = k \left\langle \frac{\partial S}{\partial \eta} \right\rangle^{-1}$$

- Where  $k$  is chosen so  $\bar{g}^2 = g_0^2$  to first order in perturbation theory

$$k = -24 \frac{L^2}{a^2} \sin \left( \frac{a^2}{L^2} (\pi - 2\eta) \right)$$

- $\bar{g}^2$  is deined on the scale of the Lattice size
- We can probe different scales by varying the size of the lattice

## Step Scaling

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- We can define a step scaling function  $\Sigma$  that measures the changing in the coupling after a scaling of size  $s$

$$\Sigma(u, s, a/L) = \bar{g}^2(g_0, sL/a) |_{\bar{g}^2(g_0, L/a)=u}$$

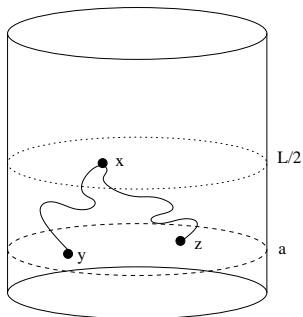
- By simulating at multiple values of  $L$  we can calculate  $\Sigma$  at different values of  $a/L$  allowing us to form a continuum extrapolation

$$\sigma(u, s) = \lim_{a/L \rightarrow 0} \Sigma(u, s, a/L)$$

- The continuum step scaling function can be directly related to the  $\beta$  function

$$-2 \log s = \int_u^{\sigma(u,s)} \frac{dx}{\sqrt{x} \beta(\sqrt{x})}$$

# Mass



- We use the PCAC relation to define the physical quark mass and set it to zero by tuning the bare quark mass

$$m_{PCAC}(x) = \frac{\frac{1}{2}(\partial_0 + \partial_0^*)f_A(x)}{2f_P(x)}$$

- The Schrödinger Functional Boundary conditions suppress fermionic zero modes allowing us to simulate at  $m_q = 0$
- All correlators can be measured from either boundary giving a non-trivial check of the mass

# Mass Renormalisation

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- The correlators  $f_P, f_A$  are the pseudoscalar and axial currents defined by:

$$f_P(x) = -1/12 \int d^3y d^3z \langle \bar{\psi}(x) \gamma_5 \tau^a \psi(x) \bar{\zeta}(y) \gamma_5 \tau^a \zeta(z) \rangle,$$

$$f_A(x) = -1/12 \int d^3y d^3z \langle \bar{\psi}(x) \gamma_0 \gamma_5 \tau^a \psi(x) \bar{\zeta}(y) \gamma_5 \tau^a \zeta(z) \rangle$$

- From these we can calculate the pseudoscalar renormalisation constant  $Z_P$  allowing us to calculate the running of the quark mass

$$Z_P = \sqrt{3f_1/f_P(L/2)}$$

$$f_1 = \frac{-1}{12L^6} \int d^3u d^3v d^3y d^3z \langle \bar{\zeta}'(u) \gamma_5 \tau^a \zeta'(v) \bar{\zeta}(y) \gamma_5 \tau^a \zeta(z) \rangle$$

- $Z_P$  requires boundary fields set to Identity and must be calculated on different lattices from  $\bar{g}^2$



# Mass Scaling

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- We can define a step scaling function for the mass in the same way as for the coupling

$$\Sigma_P(u, s, a/L) = \frac{Z_P(g_0, sL/a)}{Z_P(g_0, L/a)} \Big|_{\bar{g}^2(g_0, L/a)=u}$$

- Again we use the raw data to construct a continuum limit

$$\sigma_P(u, s) = \lim_{a \rightarrow 0} \Sigma_P(u, s, a/L)$$

- We can use this to form an estimator for  $\gamma^*$

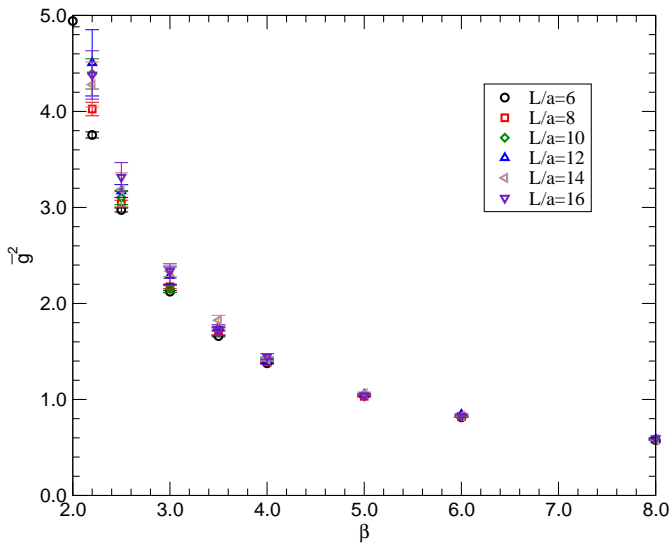
$$\hat{\gamma}(u) = -\frac{\log |\sigma_P(u, s)|}{\log |s|}$$

# Simulation Details

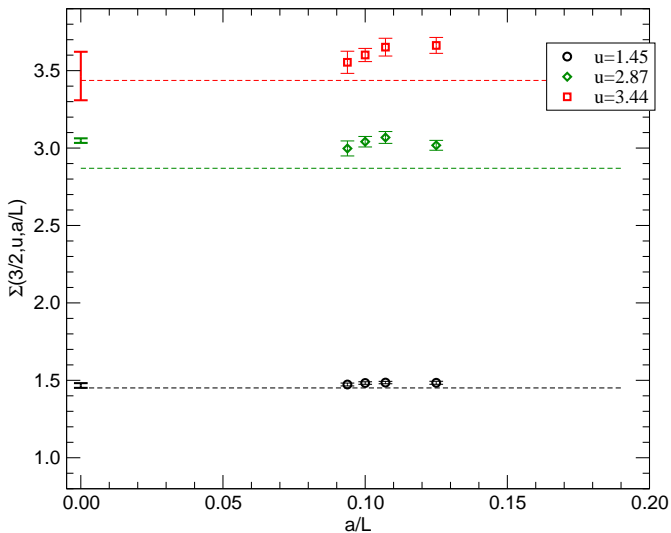
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- We use unimproved Wilson fermions and an RHMC with 4 pseudofermions
- $\kappa_c$  is determined from lattices with unit boundary conditions
- We simulate on lattices of size  $L = 6, 8, 10, 12, 14, 16$  for a range of bare couplings between  $\beta = 2$  and  $\beta = 8$
- We interpolate in  $L$  to produce results for  $L = 9, 9\frac{1}{3}, 10\frac{2}{3}, 15$
- This gives us multiple pairs of lattices with  $s = 3/2$  to allow for a continuum limit

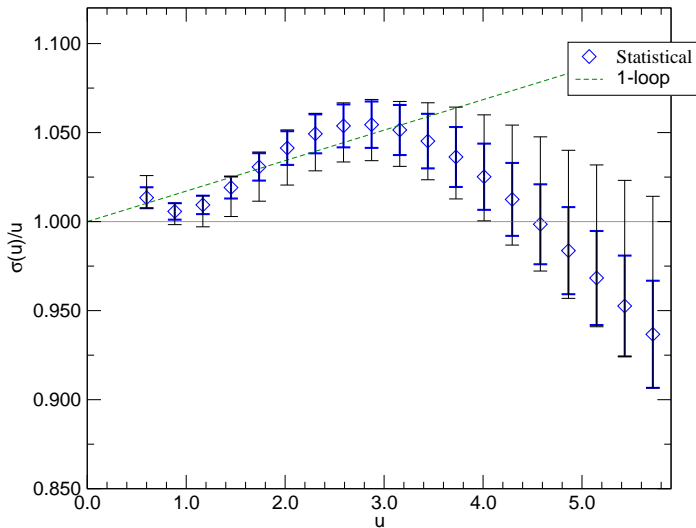
# Coupling Results



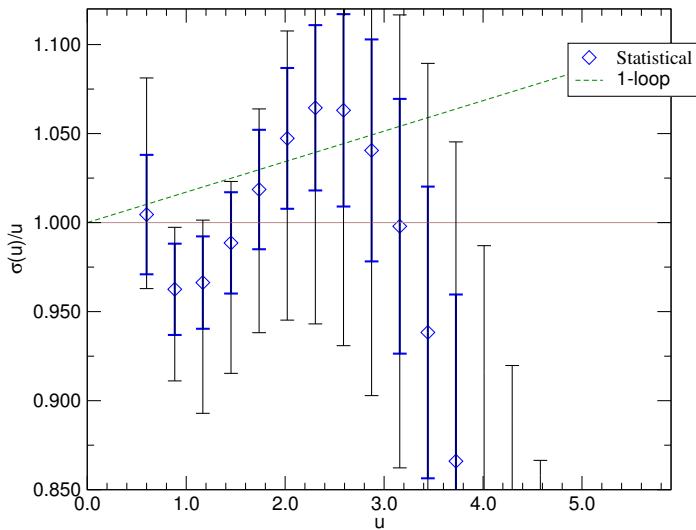
# Continuum Extrapolation



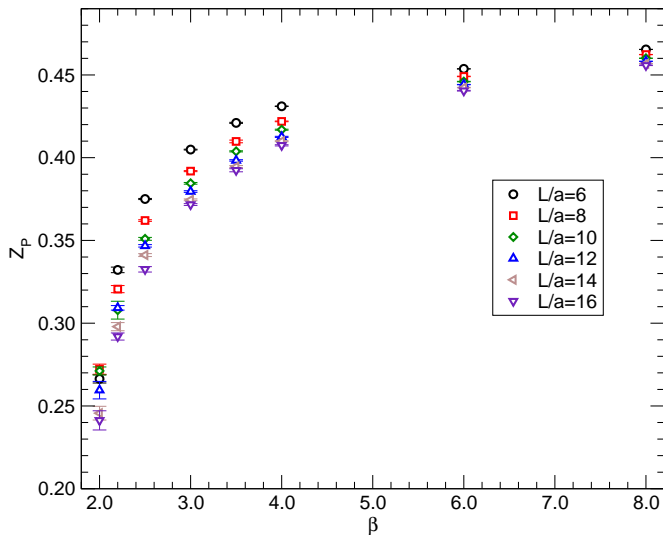
# Constant Extrapolation Scaling



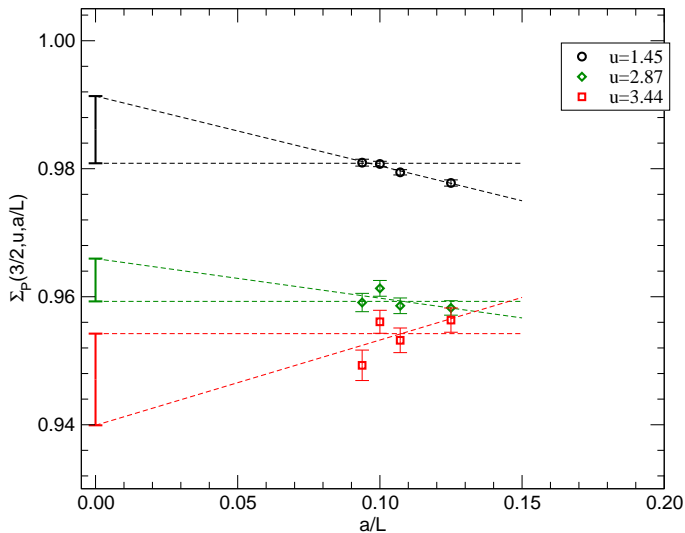
# Linear Extrapolation Scaling



# $Z_P$ Results

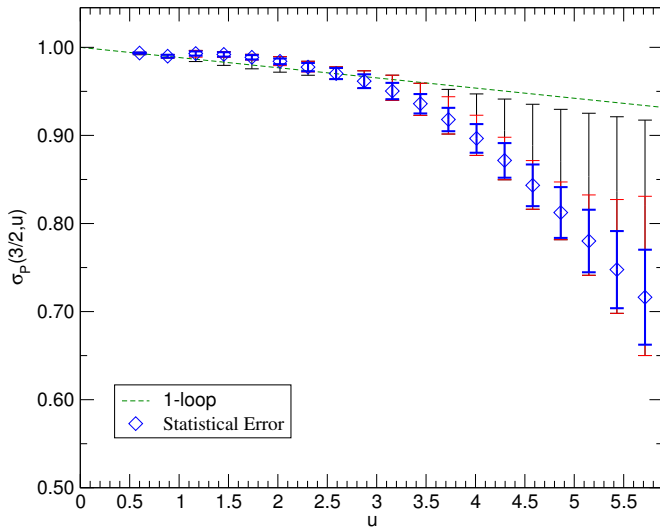


# $Z_P$ Continuum Extrapolation

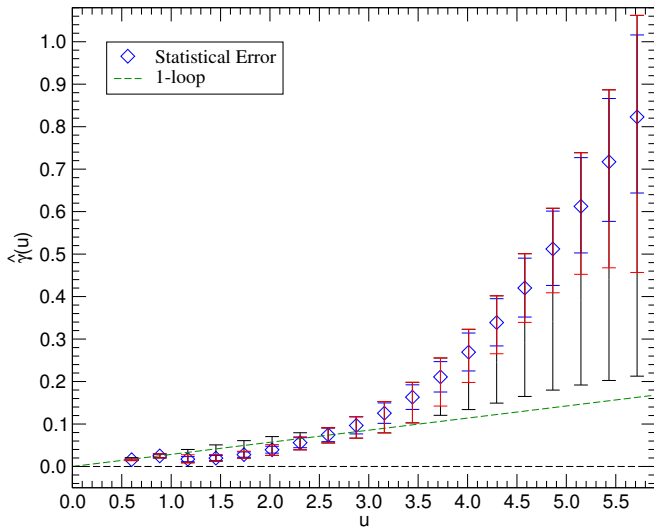




# Mass Scaling



# Mass Anomalous Dimension

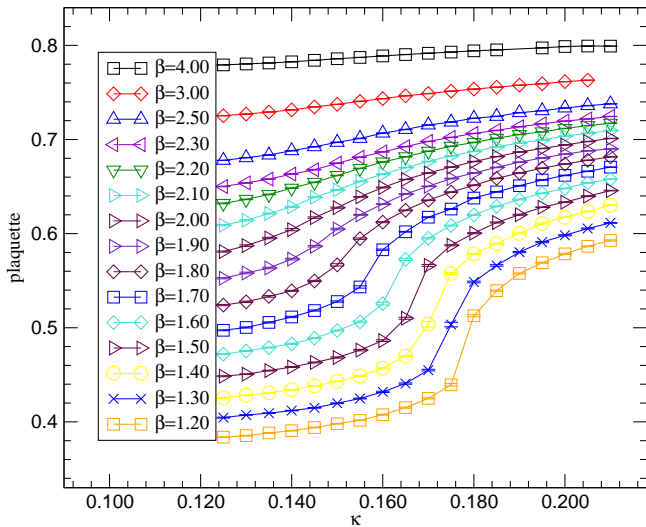


# Summary

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- The coupling runs very slowly across the range investigated
- It is hard to conclusively argue for the existence of an IRFP
- Further simulation is required to pin down the correct continuum extrapolation
- The anomalous dimension is better determined but exact values are limited by the uncertainty in the fixed point
- The data favours a low value of the anomalous dimension except at the extreme range of the coupling

# Thermal Scan



# Total Error

