

Conformal window on the lattice

Luigi Del Debbio

University of Edinburgh

Lattice 2010



Outline

1 Introduction

2 Technicolor models

3 Numerical tools

4 Results

5 Outlook

Introduction

Infrared fixed points

[Banks & Zaks 82]

$$\beta(\bar{g}) = \mu \frac{d}{d\mu} \bar{g}(\mu) = -\beta_0 \bar{g}^3 - \beta_1 \bar{g}^5 + O(\bar{g}^7)$$

$\beta_0 > 0 \implies$ asymptotic freedom, $n_f < n_f^{\text{af}}$

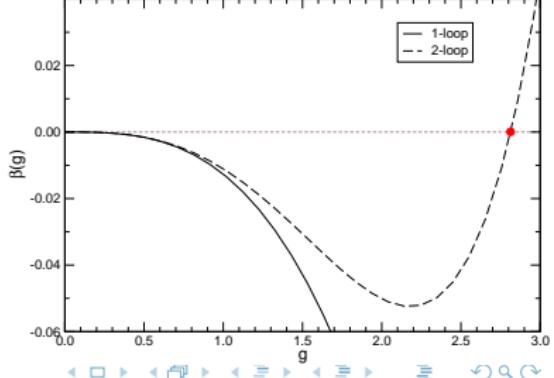
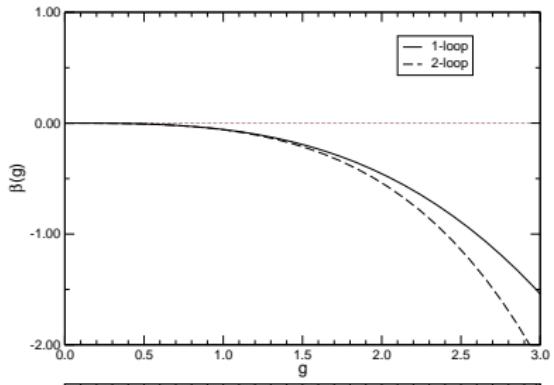
$\beta(\bar{g}) < 0 \implies$ confining theories

$\beta(\bar{g}^*) = 0 \implies$ IR fixed point

conformal window: IR zero of the β function

$$n_f^c < n_f < n_f^{\text{af}}$$

	$(4\pi)^2 \beta_0$	$(4\pi)^4 \beta_1$
SU(3), $n_f = 3$ fund	9	64
SU(3), $n_f = 12$ fund	3	-50
SU(2), $n_f = 2$ adj	2	-40



Introduction

Scheme-dependence

- existence of the IRFP is scheme-independent

$$\beta(\bar{g}) = \mu \frac{d}{d\mu} \bar{g}(\mu)$$

change of scheme: $\bar{g}' = \Phi(\bar{g})$, $\Phi'(\bar{g}) > 0$

$$\boxed{\beta'(\bar{g}') = \frac{\partial}{\partial \bar{g}} \Phi(\bar{g}) \beta(\bar{g})}$$

universality of the first two coefficients

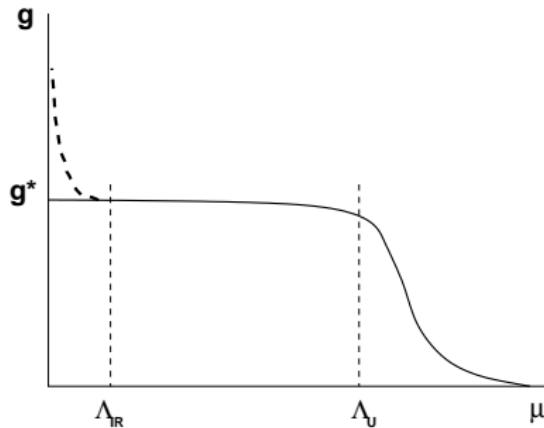
- β -function away from IRFP is not universal!

- $\bar{m}' = \bar{m} F(\bar{g})$

$$\boxed{\gamma' = \gamma + \beta \frac{\partial}{\partial \bar{g}} \log F}$$

Introduction

The scales of an IRFP

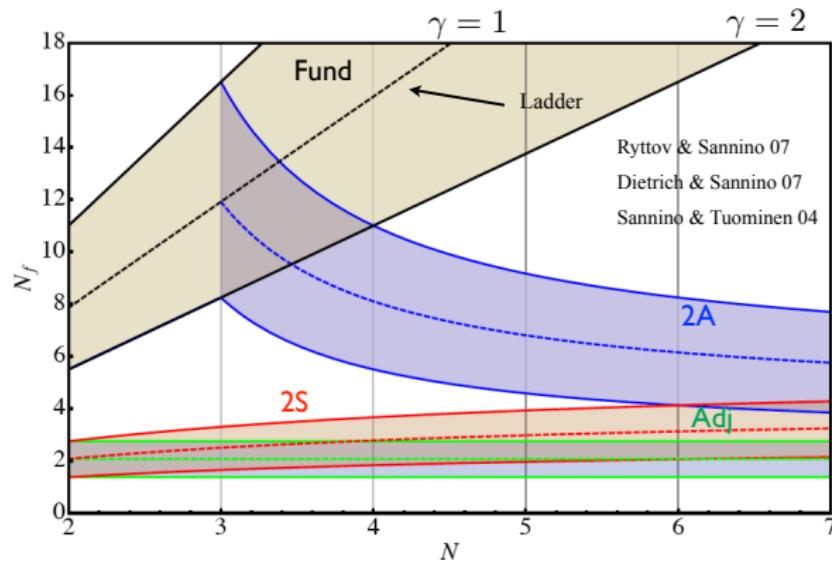


- asymptotic freedom at high-energy scale Λ_{UV}
- scale invariance at large distances \implies no massive spectrum
- large-distance dynamics: anomalous dimensions at the IRFP
- nonperturbative IRFP: strong coupling, large anomalous dimensions
- scale invariance explicitly broken by a mass term/finite volume

Introduction

The conformal window

SU(N) Phase Diagram

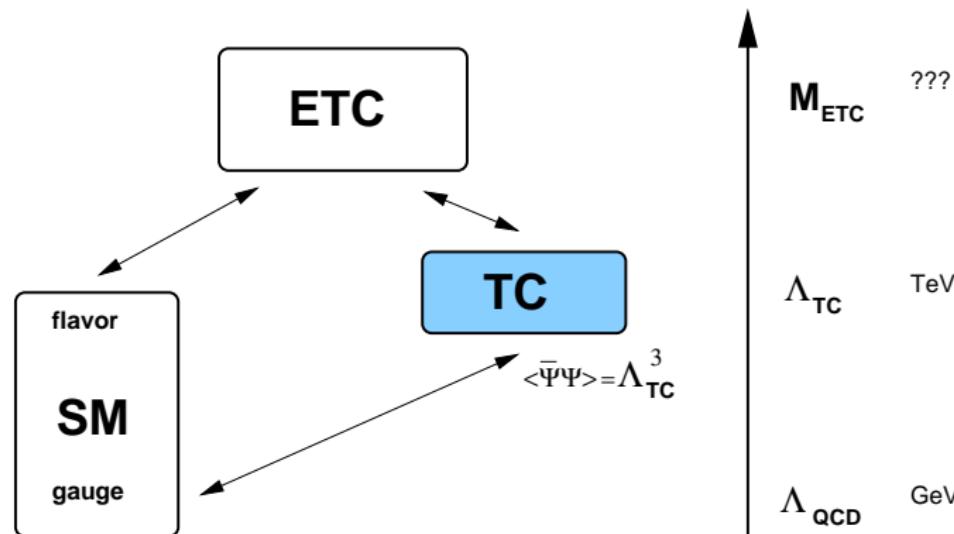


the edge of the window is relevant for phenomenology

Technicolor models

Relevance for pheno

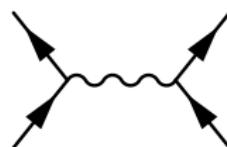
[talk by Chivukula]



Technicolor models

Flavor constraints for TC

- ETC: mass term for the SM fermions, FCNC [talk by Chivukula]



$$\longrightarrow \Delta\mathcal{L} \propto \frac{1}{M_{\text{ETC}}^2} \langle \bar{\Psi}\Psi \rangle_{\text{ETC}} \bar{\psi}\psi, \frac{1}{M_{\text{ETC}}^2} \bar{\psi}\psi \bar{\psi}\psi$$

- running of $\langle \bar{\Psi}\Psi \rangle$: **(near) conformal IR behaviour**

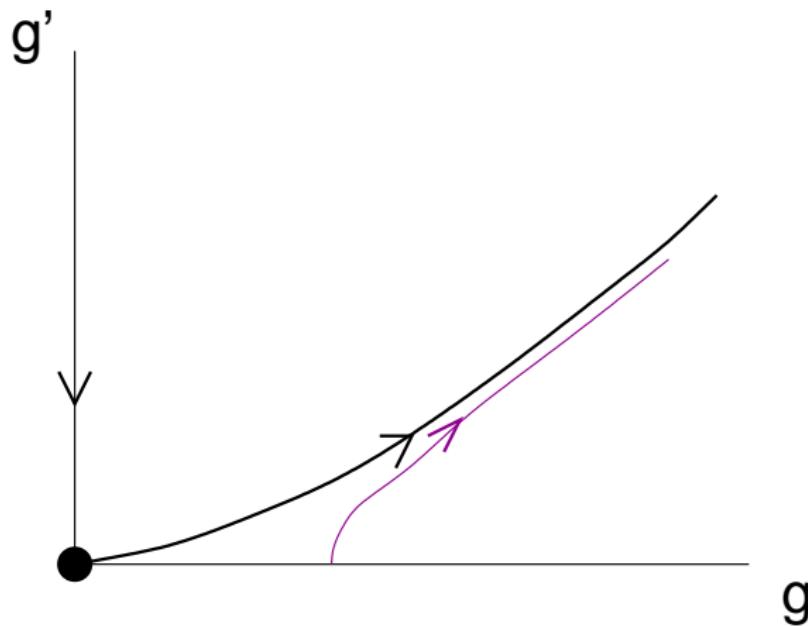
$$\boxed{\langle \bar{\Psi}\Psi \rangle_{\text{ETC}} = \langle \bar{\Psi}\Psi \rangle_{\text{TC}} \exp \left(\int_{\Lambda_{\text{TC}}}^{M_{\text{ETC}}} \frac{d\mu}{\mu} \gamma(\mu) \right)}$$

[holdom, lane, appelquist, luty, sannino, chivukula,...]

Lattice formulation

RG flow for QCD-like theories

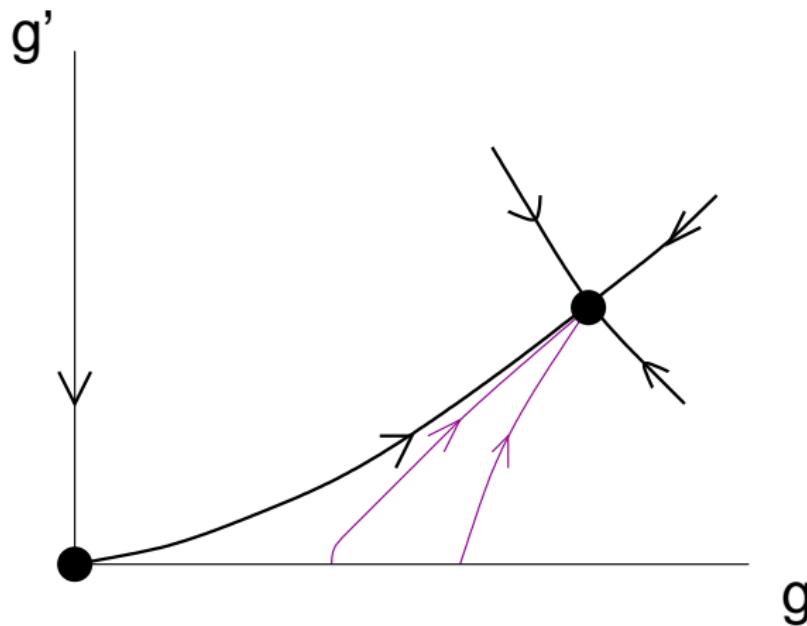
YM theory - g lattice bare coupling, g' space of irrelevant couplings



Lattice formulation

RG flow for conformal theories

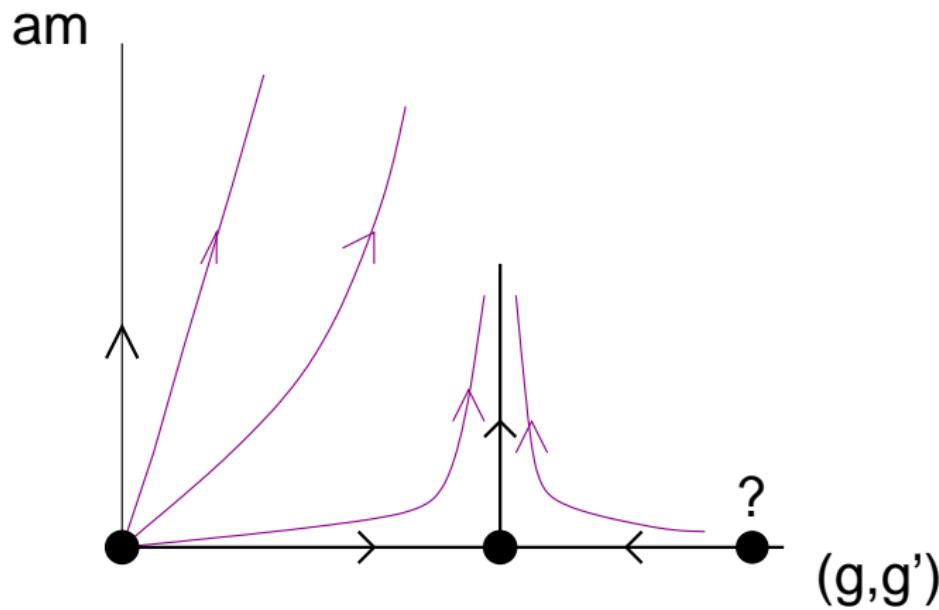
flow in the $m = 0$ plane - assuming only the mass is relevant in the IR



Lattice formulation

RG flow for conformal theories

flow in the presence of a mass deformation



non-trivial UV fixed point? [Kaplan et al 09, Hasenfratz 10, talk by Pallante]

Numerical tools

Scaling laws

Using the linearized RGEs in the vicinity of the IRFP [DeGrand & Hasenfratz 09, Fodor et al 09, Idd et al 10]

- solution of RGEs for the two-point correlators

$$C_H(t; g, \hat{m}, a) = m^{2\Delta_H/y_m} \mathcal{F}(t \hat{m}^{1/y_m}) \sim \exp(-M_H t)$$

- scaling with the quark mass

$$M_H \sim \xi^{-1} \sim m^{1/y_m}$$

- finite-size scaling

$$M = L^{-1} F(x), \quad x = L^{y_m} m = (L/\xi)^{y_m}$$

- scaling is characterized by the critical exponent $y_m = 1 + \gamma$

asymptotic behaviour: $m \rightarrow 0, L \rightarrow \infty$ (beware of $m \rightarrow 0$ at fixed L !!)

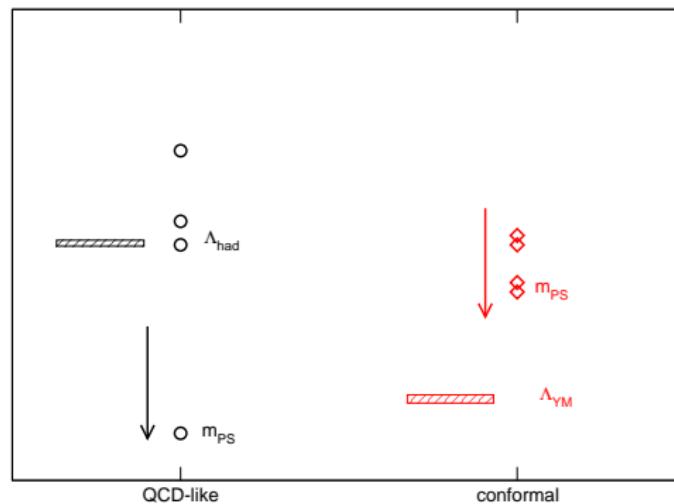
Numerical tools

Conformal spectrum

In the presence of a mass deformation (and infinite volume)

- no parametric separation between the PS and the other channels
- all masses/decay constants scale like m^{1/y_m} (but $M \sim F(0)/L$!!)

→ very different spectrum from a QCD-like theory [Miransky 94]



Matching two lattices of size L/a and $L/(ab)$

$$O_k^{(n)}(\beta) = O_k^{(n-1)}(\beta')$$

$$m = 0$$

$$\begin{aligned}\xi_{\text{lat}}(\beta') &= \xi_{\text{lat}}(\beta)/b \\ a(\beta') &= ba(\beta)\end{aligned}$$

Define bare step scaling function

$$s_0(\beta, b) = \beta - \beta'$$

fixed point

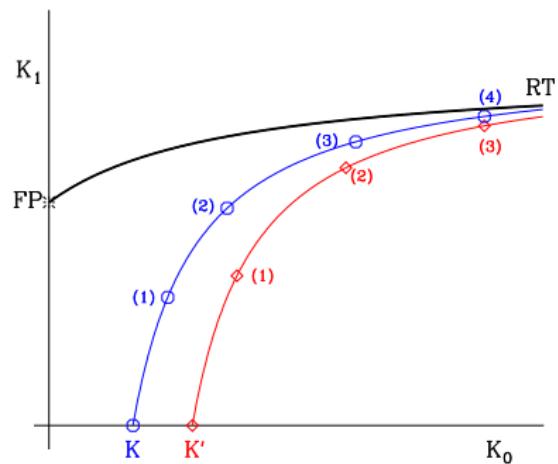
$$s_0(\beta^*, b) = 0$$

Matching m at fixed β

$$\xi_{\text{lat}}(m) = b\xi_{\text{lat}}(m')$$

$$\frac{m'}{m} = b^{y_m}$$

Systematics in matching at strong coupling?



Numerical tools

Schrödinger functional

Schrödinger functional [ALPHA collaboration]

$$Z[\eta] = e^{-\Gamma[\eta]} = \int_{L \times L^3} \mathcal{D}U \mathcal{D}\psi \mathcal{D}\bar{\psi} \exp(-S[U, \psi, \bar{\psi}])$$

Dirichlet boundary conditions at $t = 0, L$, dependent on η

$$\frac{k}{\bar{g}^2(L)} = \left. \frac{\partial \Gamma}{\partial \eta} \right|_{\eta=0}, \quad \bar{g}^2 = g_0^2 + O(g_0^4)$$

Lattice step scaling function

$$\Sigma(u, a/L) = \bar{g}^2(bL) \Big|_{\bar{g}^2(L)=u, m=0}$$

Step scaling function

$$\sigma(u) = \lim_{a \rightarrow 0} \Sigma(u, a/L)$$

fixed point:

$$\sigma(u^*) = u^*$$

Numerical tools

Schrödinger functional

Definition of the renormalized mass

$$\partial_\mu (A_R)_\mu = 2\bar{m}P_R$$

$$\begin{aligned}(A_R)_\mu(x) &= Z_A \bar{\psi}(x) \gamma_\mu \gamma_5 \psi(x) \\ (P_R)(x) &= Z_P \bar{\psi}(x) \gamma_5 \psi(x)\end{aligned}$$

Lattice step scaling function

$$\begin{aligned}Z_P(g_0, L/a) &= c \frac{\sqrt{3f_1}}{f_P(L/2)} \\ \Sigma_P(u, a/L) &= \frac{Z_P(g_0, bL/a)}{Z_P(g_0, L/a)}, \quad \bar{g}^2(L) = u\end{aligned}$$

Step scaling function

$$\sigma_P(u) = \lim_{a \rightarrow 0} \Sigma_P(u, a/L)$$

Numerical tools

Potential schemes

[Bilgici et al 09, Fodor et al 09, talk by Etou]

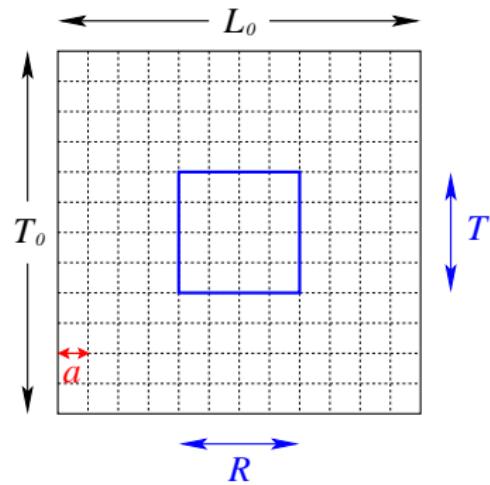
$$g_w^2(L_0, R/L_0, a/L_0) = \frac{1}{k} (R/a)^2 \chi(R/a, L_0/a)$$

step scaling function

$$\begin{aligned}\Sigma_w(u, R/L_0, a/L_0) &= g_w^2(bL_0, R/L_0, a/L_0), \\ u &= g_w^2(L_0, R/L_0, a/L_0)\end{aligned}$$

$$\sigma_w(u) = \lim_{a \rightarrow 0} \Sigma_w(u, R/L_0, a/L_0)$$

$O(a)$ improved



Numerical tools

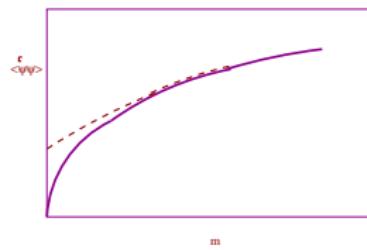
Phase diagram

[Deuzeman et al 09]

ANALOGY WITH STRONGLY COUPLED QED



- At the critical point g_c
 $\langle\bar{\psi}\psi\rangle = am^{1/\delta}$
- In the symmetric phase
Conformality perturbatively broken by Coulomb forces:
 $\langle\bar{\psi}\psi\rangle = am^d$, $d = \frac{3-\gamma}{1+\gamma}$
- Only in the free limit:
 $d = 1$.
- d measured in QED
Kogut et al. 1985
 $1/\delta < d < 1$



Lattice strategy:

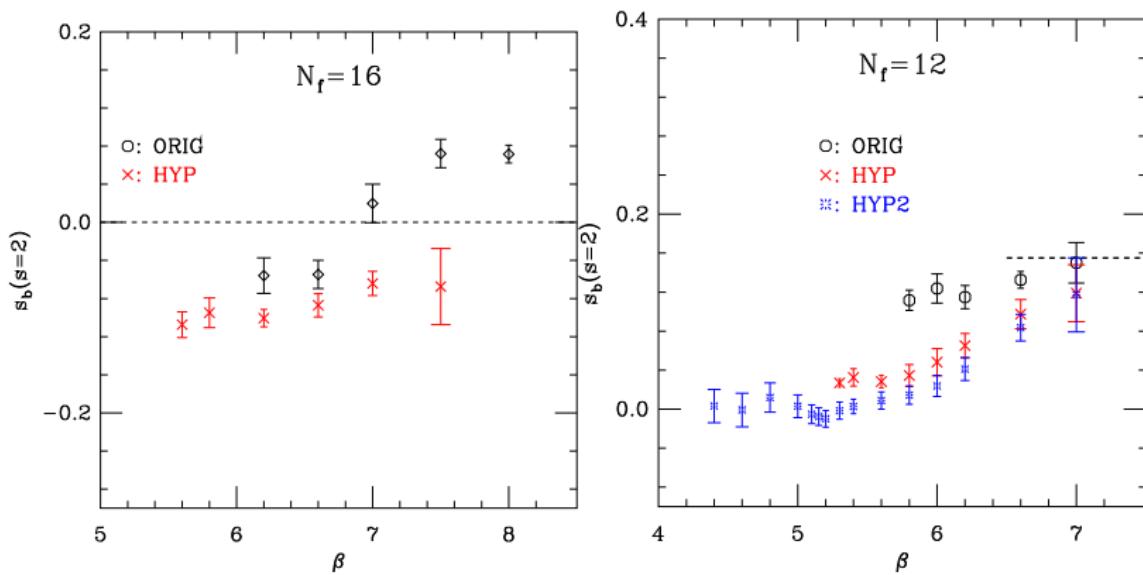
-Contrast χ^2 of different fits -

- Symmetric Phase $\langle\bar{\psi}\psi\rangle = am^d$
- Broken Phase $\langle\bar{\psi}\psi\rangle = am + b$

Results

SU(3) fundamental

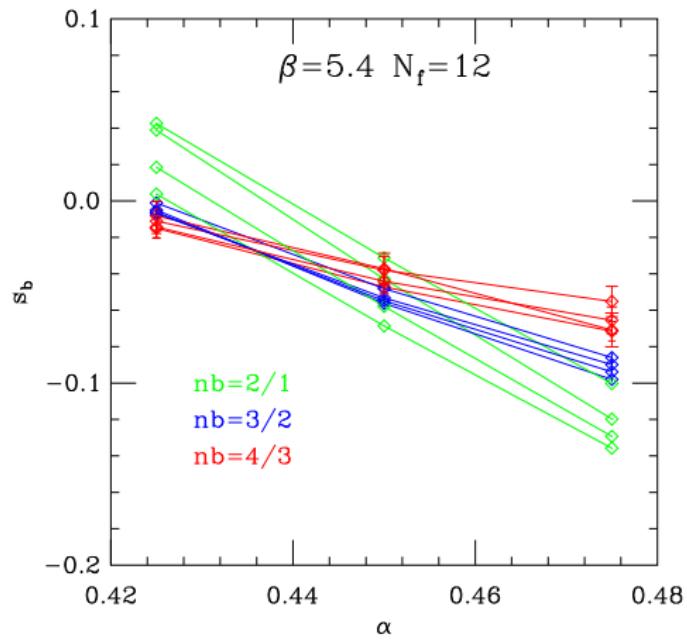
$n_f = 16, 12$ staggered, MCRG [Hasenfratz 10]



Results

SU(3) fundamental

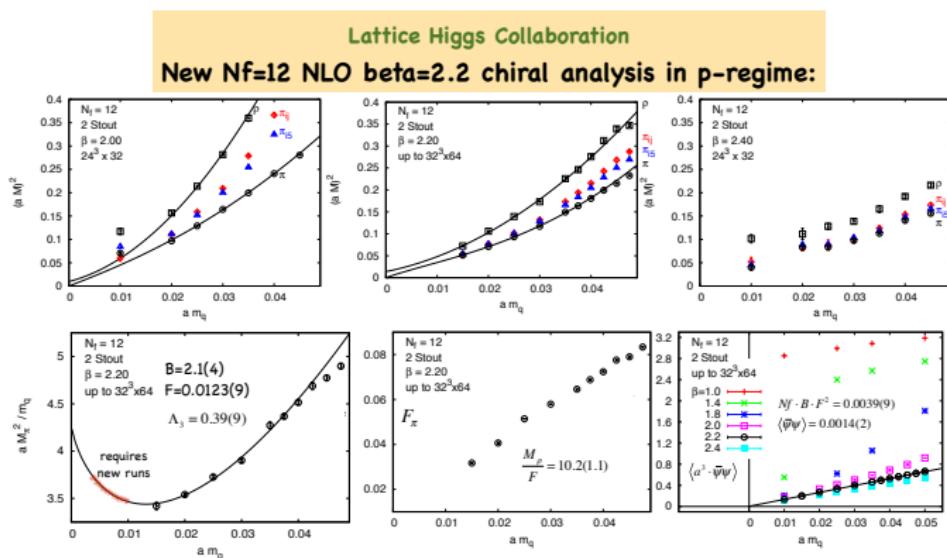
$n_f = 12$ staggered, MCRG [Hasenfratz 10]



Results

SU(3) fundamental

$n_f = 12$ staggered, spectrum [talks by Kuti]



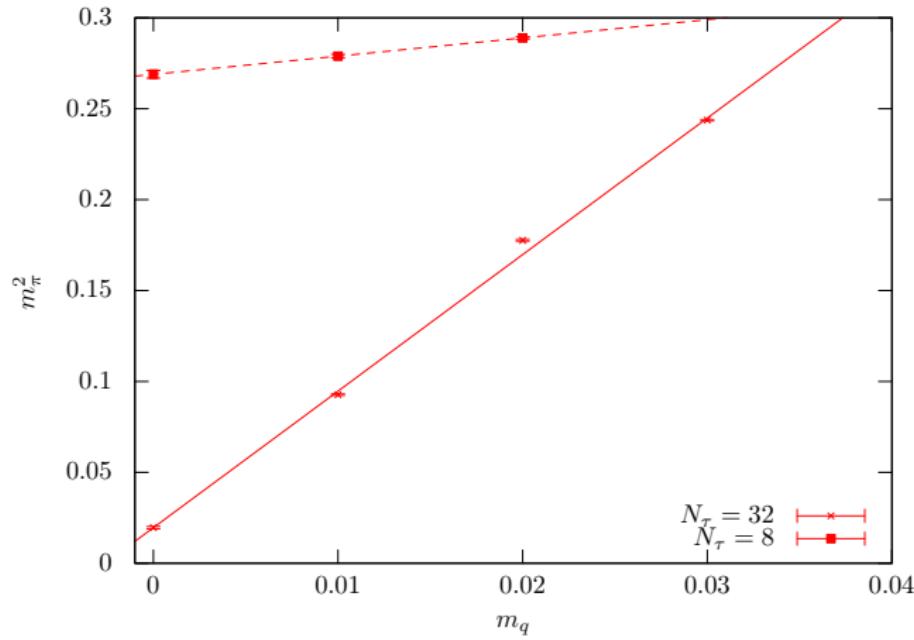
Consistent with chiral symmetry breaking

+ eigenvalues distribution in agreement with RMT

Results

SU(3) fundamental

$n_f = 12$ staggered, spectrum [Jin & Mawhinney 09/10]



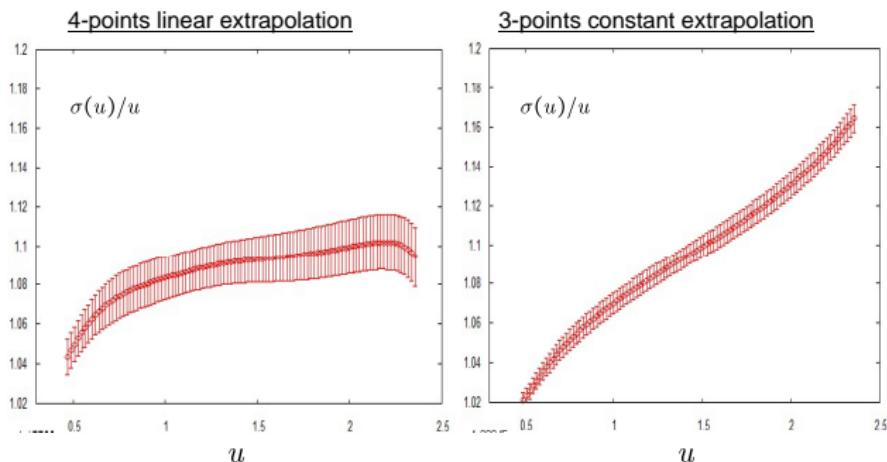
data consistent with broken chiral symmetry

Results

SU(3) fundamental

$n_f = 12$ staggered, TPL scheme [talk by Itou]

There is a large systematic error.



There is no signal of the fixed point.

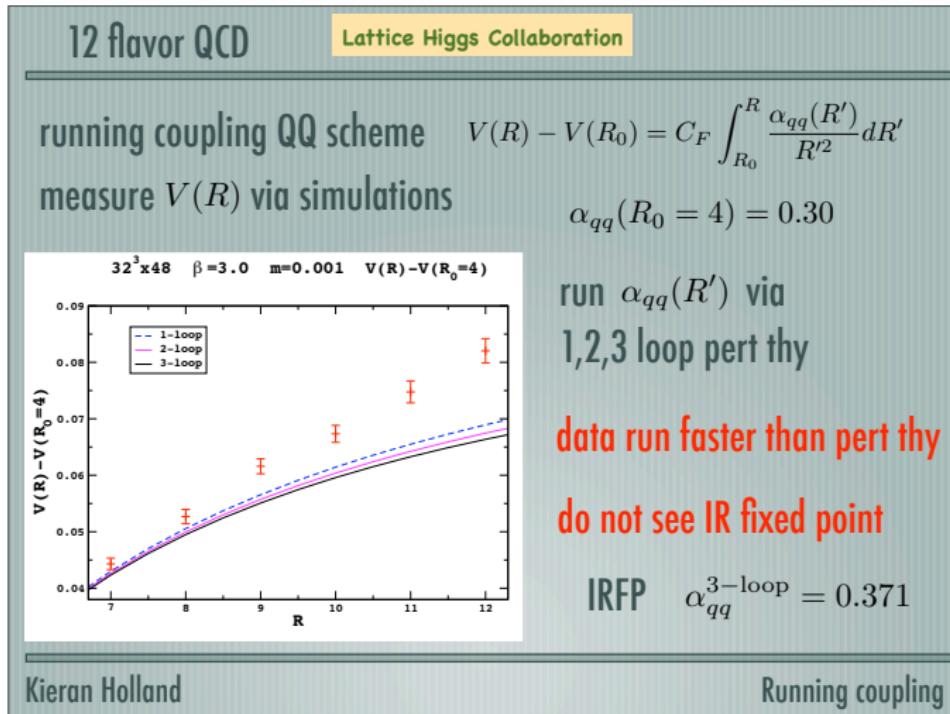
We should be study the large u region ($\beta < 5$) to find the fixed point
and

have to do some improvements (larger lattice size simulation).

Results

SU(3) fundamental

$n_f = 12$ staggered, QQ scheme [talk by Holland]



Results

SU(3) fundamental

$n_f = 12$ staggered, phase diagram [poster by Lombardo, talk by Pallante]

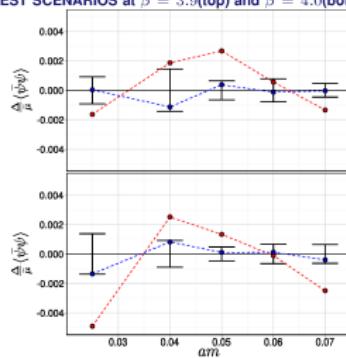


CHIRALLY SYMMETRIC PHASE : ANALOGIES WITH LATTICE QED



SIDE-BY-SIDE COMPARISON OF THE TWO SIMPLEST SCENARIOS at $\beta = 3.9$ (top) and $\beta = 4.0$ (bottom)

- $Y = bX^d$
 $(\chi^2/ndf) = 0.4-0.3$,
Chiral Symmetry :
Consistent with data
- $Y = bX + K$
 $(\chi^2/ndf) = 3.9-4.0$,
Chiral Symmetry Breaking:
Excluded within 2σ

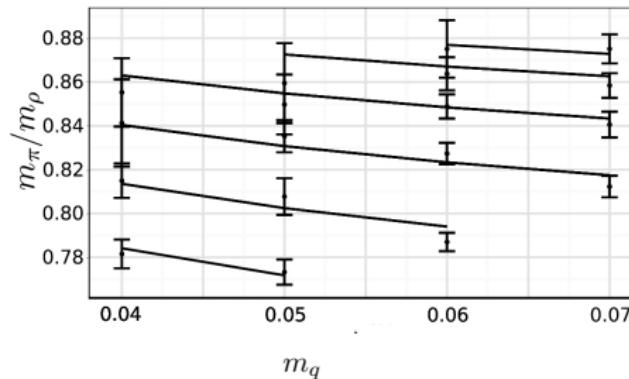


A. Deuzeman, E. Pallante, M.P. Lombardo

Results

SU(3) fundamental

$n_f = 12$ staggered, phase diagram [poster by Lombardo, talk by Pallante]



$g_L = 1.22 - 1.30$ from top to bottom

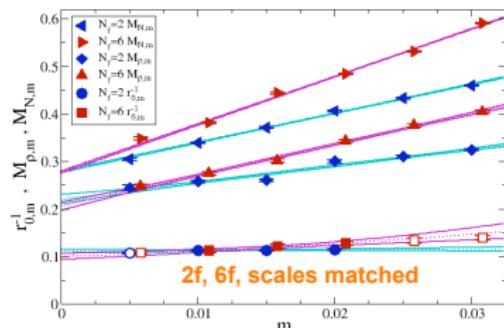
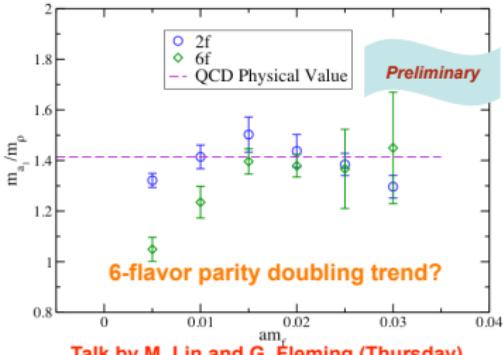
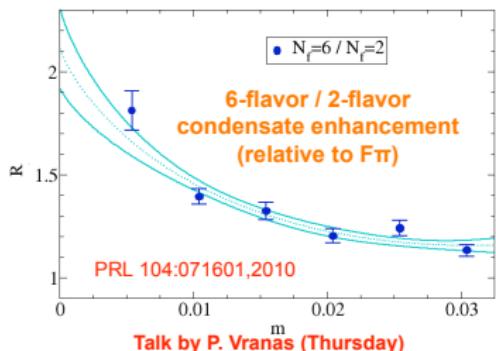
SPECTRUM AND CHIRAL SYMMETRY: SAME TREND AS IN SYMMETRIC QED,
CONSISTENT WITH CHIRAL SYMMETRY

A. Deuzeman, E. Pallante, M.P. Lombardo

Results

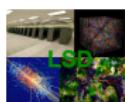
SU(3) fundamental: a trip at the edge of the window

Lattice Strong Dynamics collaboration



S parameter Talk by D. Schaich (Thursday)

- DWF ~ 1,000 configurations
- SU(3c), SU(2c) fundamental
- Small lattice spacing $a \sim N_f$ matched
- $V = 32^3 \times 64 \quad M_\pi L > 4$
- Condensate enhancement
- Spectrum
- S parameter
- Eigenvalue spectrum
- LLNL BG/L ~ 300 million core hours

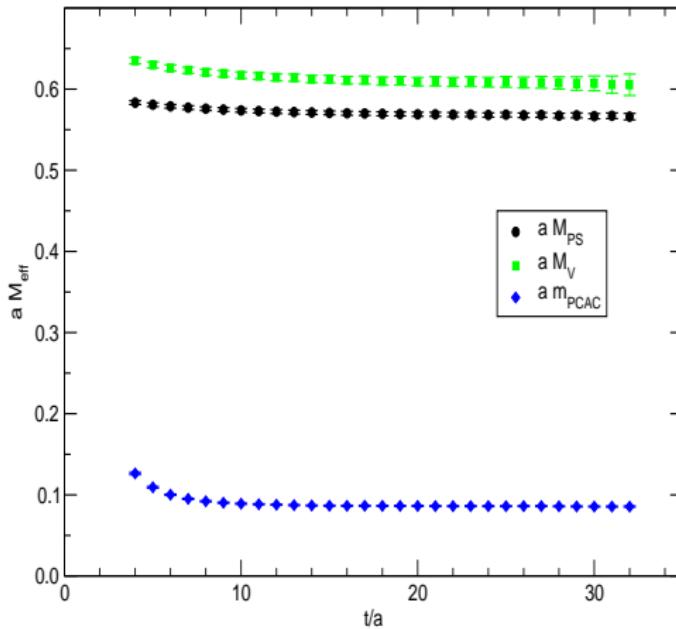


Results

SU(2) adjoint

$n_f = 2$ Wilson, spectrum [Catterall et al 07-09, Idd et al 08-10, Hietanen et al 08]

$\beta = 2.10$, 64×24^3 lattice



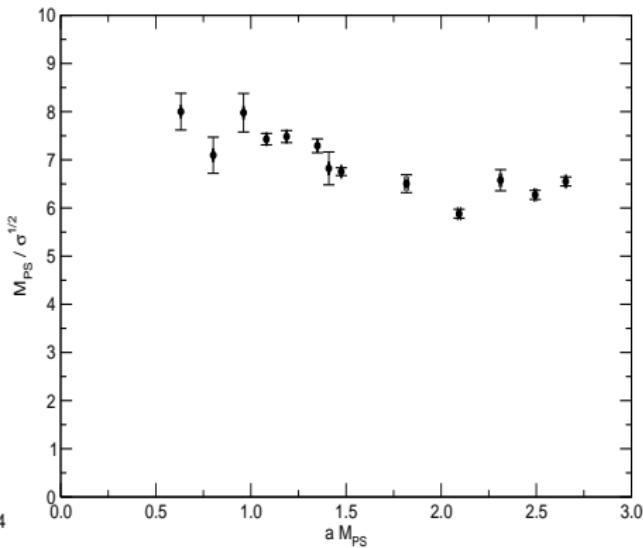
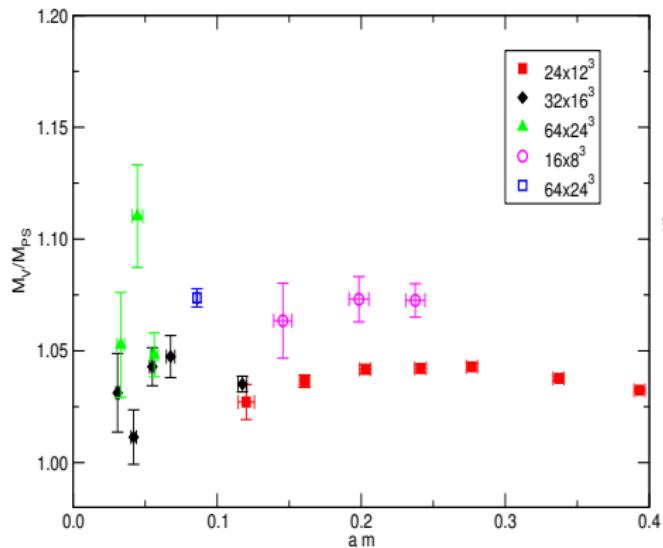
systematic errors on smaller volumes

[talk by Kerrane]

Results

SU(2) adjoint

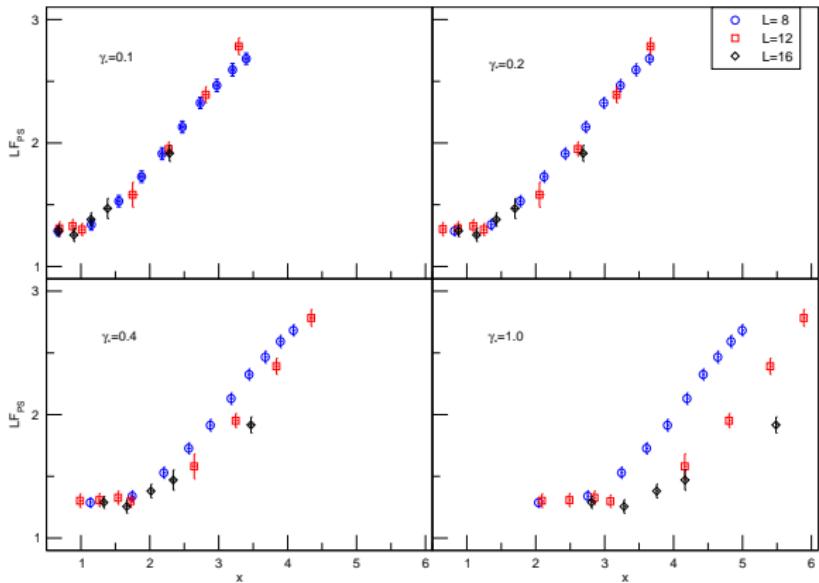
$n_f = 2$ Wilson, spectrum [talks by Pica & Patella]



Results

SU(2) adjoint

$n_f = 2$ Wilson, FSS spectrum [talk by Pica]

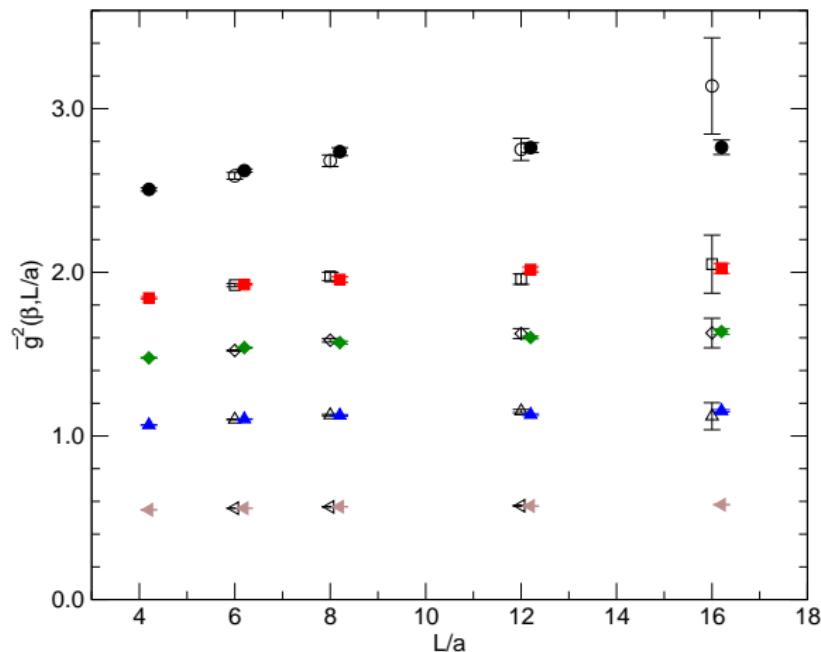


for several γ , $LF_{PS} = F(L^{\gamma_m} m)$; $\gamma < 0.5$ is favoured

Results

SU(2) adjoint

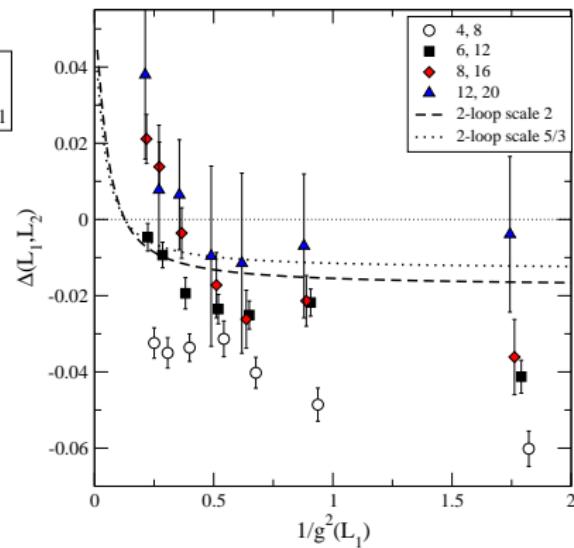
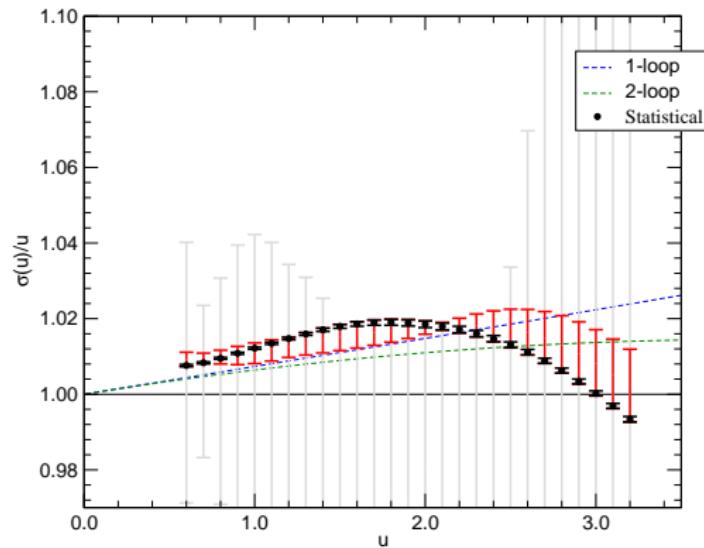
$n_f = 2$ Wilson, Schrödinger functional coupling [Hietanen et al 08, Bursa et al 09]



Results

SU(2) adjoint

$n_f = 2$ Wilson, Schrödinger functional coupling [Hietanen et al 08, Bursa et al 09]

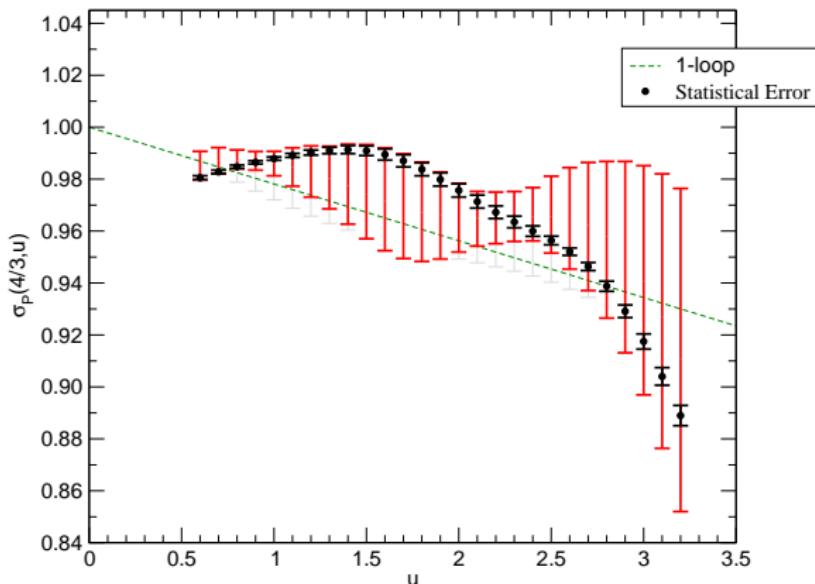


$$\Delta(L_1, L_2) = 1/g^2(L_2) - 1/g^2(L_1)$$

Results

SU(2) adjoint

$n_f = 2$ Wilson, Schrödinger functional mass [Bursa et al 09]

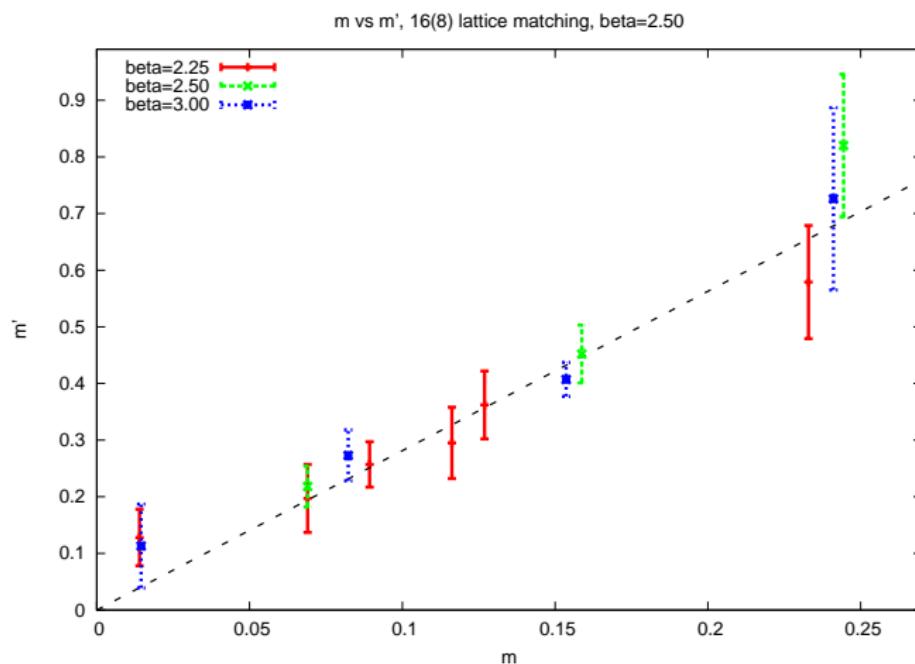


$$\log |\sigma_P(u, b)| = -\gamma \log b \implies 0.05 < \gamma < 0.56$$

Results

SU(2) adjoint

$n_f = 2$ Wilson, MCRG mass matching - PRELIMINARY [talk by Keegan]

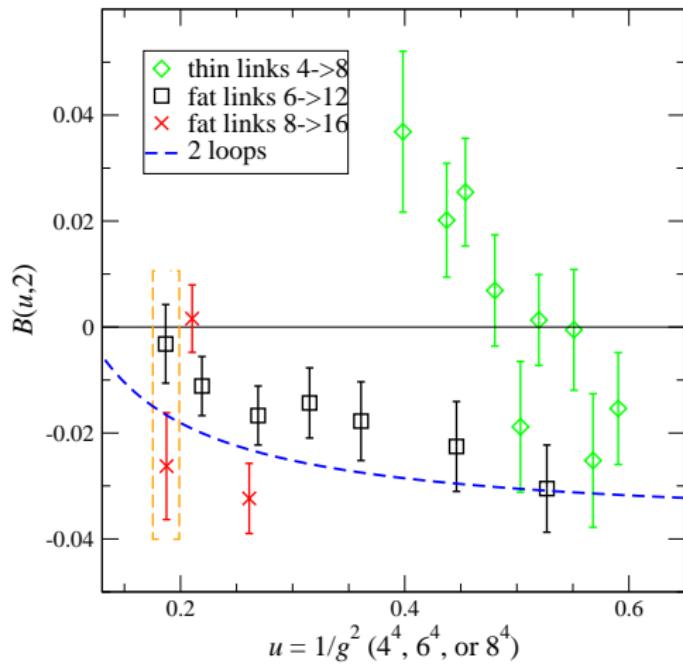


$$m' = b^{y_m} m \implies \gamma = 0.49(13)$$

Results

SU(3) sextet

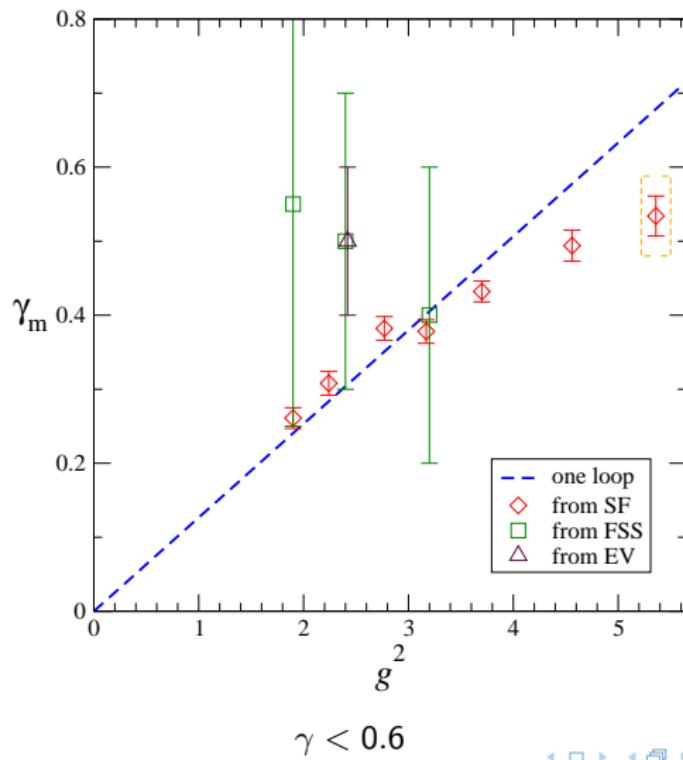
$n_f = 2$ Wilson, SF coupling [talk by Svetitsky]



Results

SU(3) sextet

$n_f = 2$ Wilson, SF mass

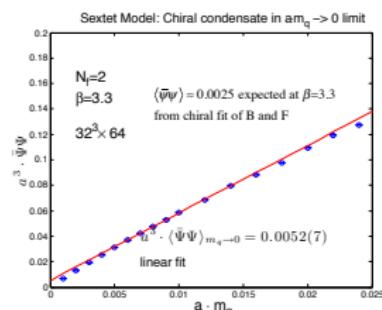
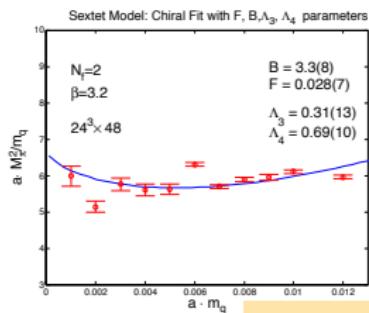
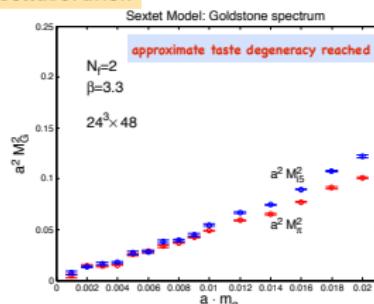
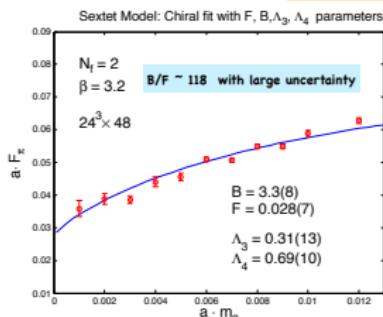


Results

SU(3) sextet

$n_f = 2$ staggered [talk by Kuti]

Lattice Higgs Collaboration



Consistent with chiral symmetry breaking

Results

SU(3) sextet

$n_f = 2$ staggered [Kogut & Sinclair 10]

$N_t = 4, 6, 8$ - separated deconfinement and chiral transitions

β values for these transitions:

N_t	β_d	β_χ
4	5.40(1)	6.3(1)
6	5.54(1)	6.6(1)
8	5.66(1)	6.7(1)

Increase in β_χ from $N_t = 6$ to $N_t = 8$ is very small

need to simulate at larger N_t and check how β_χ changes

Outlook

Towards phenomenology

- tools have been developed and tested - look for better tools
- preliminary results need to be put on solid ground
 - benchmark codes/analysis
 - control systematics: small masses & large volumes
 - improved actions [talks by Karavirta & Mykkänen]
- $SU(2)$ $n_f = 2$ adjoint, consistent results
- more results needed for $SU(3)$ $n_f = 12$ fundamental, $n_f = 2$ sextet
- all methods have systematics that “muddle” the answers
more confident when different approaches yield consistent results
- anomalous dimensions are small - better control of systematics
- deformations away from conformality
 - $n_f < n_f^c$
 - mass deformation
 - 4fermi interactions
- approach the conformal window from the chirally broken phase
LSD collaboration: deviations from QCD
- lattice results to be taken into account for model building
- talk to phenomenologists!

Lattice formulation

What can the lattice provide?

- analytical studies so far (*voodoo QCD*)
 - Schwinger-Dyson
 - AdS/CFT
 - other assumptions on the NP dynamics
- lattice provides first-principles results (w. systematic errors)
- Spectroscopy
 - masses and decay constants
 - finite-size scaling
 - eigenvalue distributions
- MCRG
 - two-lattice matching
- SF/TPL/QQ schemes
 - running coupling
 - running mass
- Phase diagram

Lattice formulation

RG flows

- field theory on a lattice

$$a, \phi(x), \{K_\alpha\}$$

- integrate out UV \longrightarrow blocked lattice

$$a' = ba, \phi'(x')$$

- constant physics

$$\int \mathcal{D}\phi' \exp[-S'(\phi')] = \int \mathcal{D}\phi \exp[-S(\phi)]$$

- flow in the space of couplings

$$S(\phi) = \sum_\alpha K_\alpha O_\alpha(\phi) \mapsto S'(\phi') = \sum_\alpha K'_\alpha O_\alpha(\phi')$$

- RG trajectory

$$\begin{aligned} \{K_\alpha\} &\mapsto \{K'_\alpha\} \mapsto \{K''_\alpha\} \\ \xi_{\text{lat}} &\mapsto \xi_{\text{lat}}/b \dots \end{aligned}$$

Lattice formulation

RG fixed points

RG transformation

$$K'_i = R_i(K)$$

fixed point

$$K^*_i = R_i(K^*)$$

linearized flow equations in a *neighbourhood* of K^*

$$\delta K'_i = R_{ij} \delta K_j, \quad R_{ij} = \left. \frac{\partial R_i}{\partial K_j} \right|_{K^*}$$

eigenvalues and eigenvectors of R_{ij}

$$u'_i = b^{y_i} u_i$$

y_i critical exponents

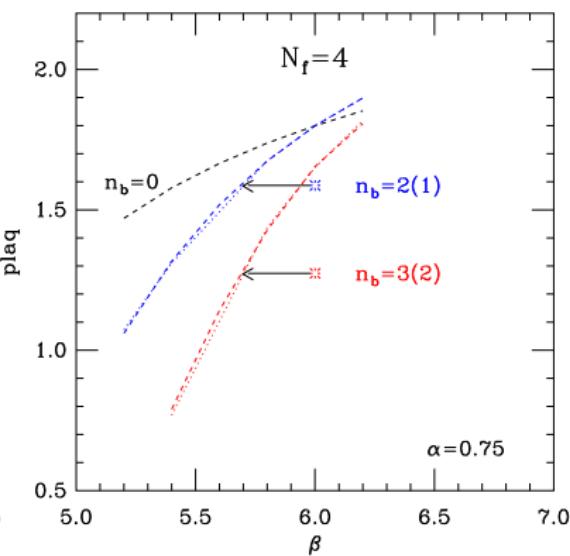
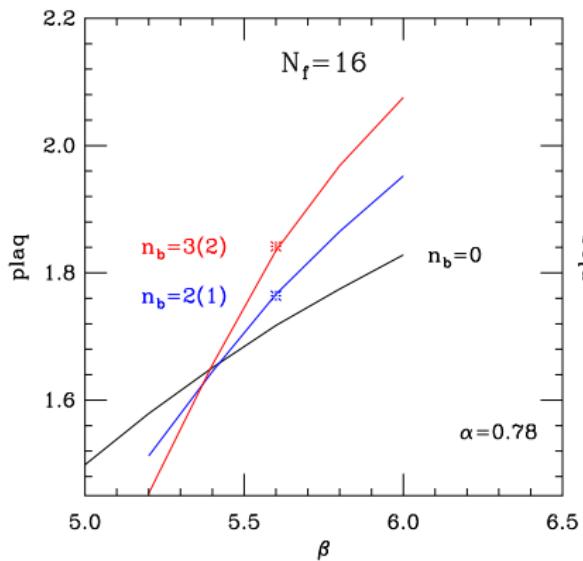
$y_i > 0$: relevant coupling at the fixed point

Results

SU(3) fundamental

(for $n_f < 4$ see other talks!!)

$n_f = 16$ staggered, evidence for IRFP from MCRG [Hasenfratz 09]



Introduction

Infrared fixed points

SU(N) gauge + n_f massless fermions in representation R

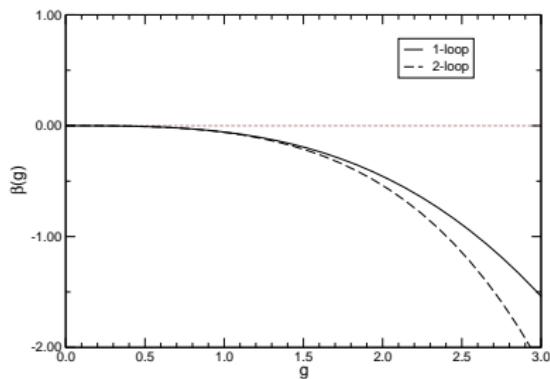
$$\beta(\bar{g}) = \mu \frac{d}{d\mu} \bar{g}(\mu) = -\beta_0 \bar{g}^3 - \beta_1 \bar{g}^5 + O(\bar{g}^7)$$

$$\beta_0 = \frac{1}{(4\pi)^2} \frac{11}{3} C_2(A) \left[1 - \frac{4}{11} x_f \right]$$

$$x_f = \frac{T_R n_f}{C_2(A)}$$

asymptotic freedom: $\beta_0 > 0$

$$x_f < \frac{11}{4} \iff n_f < n_f^{\text{af}}$$



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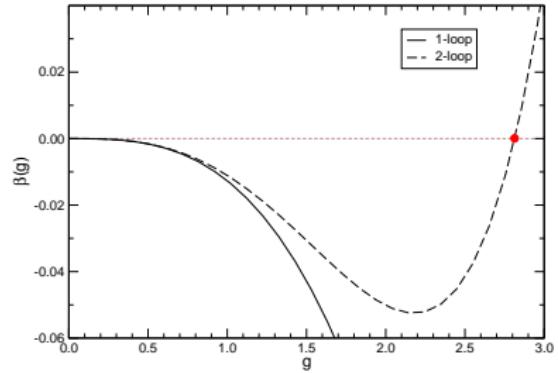
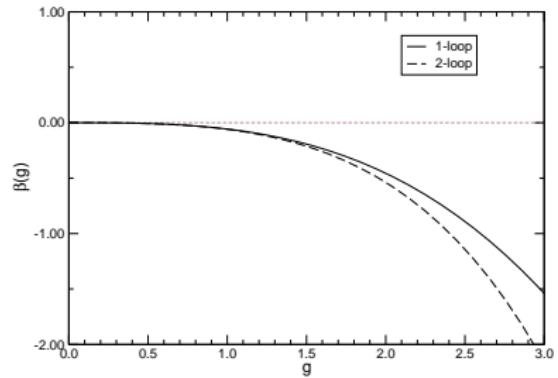
$$\beta(\bar{g}) = \mu \frac{d}{d\mu} \bar{g}(\mu) = -\beta_0 \bar{g}^3 - \beta_1 \bar{g}^5 + O(\bar{g}^7)$$

$$\beta_1 = \frac{1}{(4\pi)^4} \frac{34}{3} C_2(A)^2 [1 - F_1(R)x_f]$$

$$F_1(R) = \frac{1}{34} \left[20 + 12 \frac{C_2(R)}{C_2(A)} \right]$$

$$\beta_1 < 0 \iff 1 < F_1(R)x_f$$

	$(4\pi)^2 \beta_0$	$(4\pi)^4 \beta_1$
SU(3), $n_f = 3$ fund	9	64
SU(3), $n_f = 12$ fund	3	-50
SU(2), $n_f = 2$ adj	2	-40



Introduction

Infrared fixed points

- UV Gaussian fixed point at $g = 0$ if $n_f < n_f^{\text{af}}$
- IR fixed point at g^* if $n_f > n_f^c$

$$g^{*2} = -\frac{\beta_0}{\beta_1} = -(4\pi)^2 \frac{11}{34} \frac{1}{C_2(A)} \frac{1 - \frac{4}{11}x_f}{1 - F_1(R)x_f}$$

- conformal window

$$n_f^c < n_f < n_f^{\text{af}}$$

- for fundamental fermions:

$$x_f^c = \frac{17}{10 + 3\frac{N^2-1}{N^2}}, \quad n_f^c = \frac{34N}{10 + 3\frac{N^2-1}{N^2}}$$

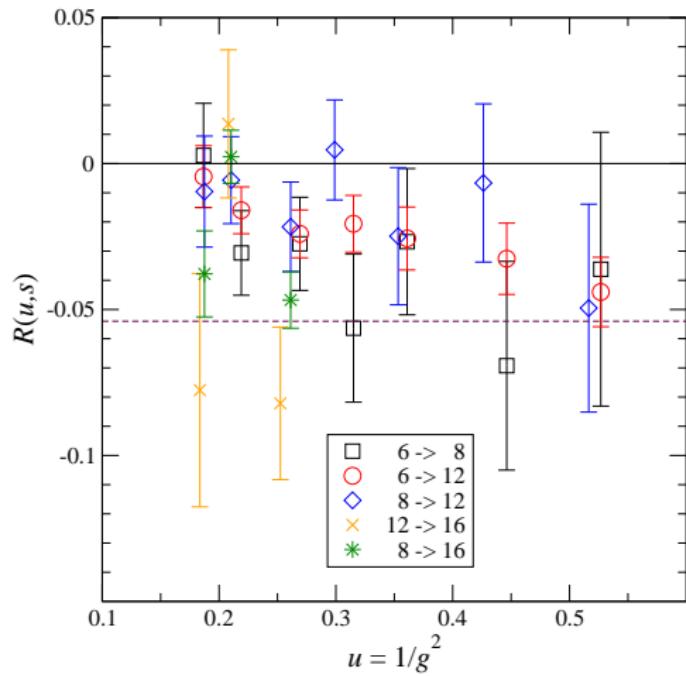
- for adjoint fermions:

$$x_f^c \equiv n_f^c = 34/32$$

Results

SU(3) sextet

$n_f = 2$ Wilson, SF coupling [DeGrand et al 10]



Numerical tools

Schrödinger functional

Schrödinger functional [ALPHA collaboration]

$$Z[\eta] = e^{-\Gamma[\eta]} = \int_{L \times L^3} \mathcal{D}U \mathcal{D}\psi \mathcal{D}\bar{\psi} \exp(-S[U, \psi, \bar{\psi}])$$

Dirichlet boundary conditions at $t = 0, L$, dependent on η

$$\left. \frac{\partial \Gamma}{\partial \eta} \right|_{\eta=0} = \frac{k}{\bar{g}^2}, \quad \bar{g}^2 = g_0^2 + O(g_0^4)$$

Lattice step scaling function:

$$\Sigma(u, a/L) = \bar{g}^2(bL) \Big|_{\bar{g}^2(L)=u, m=0}$$

Step scaling function:

$$\sigma(u) = \lim_{a \rightarrow 0} \Sigma(u, a/L)$$

$$\beta(\sqrt{\sigma(u)}) = \beta(\sqrt{u}) \sqrt{\frac{u}{\sigma(u)}} \sigma'(u)$$