

Lattice Study of Trapped Fermions at Unitarity

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In collaboration with:

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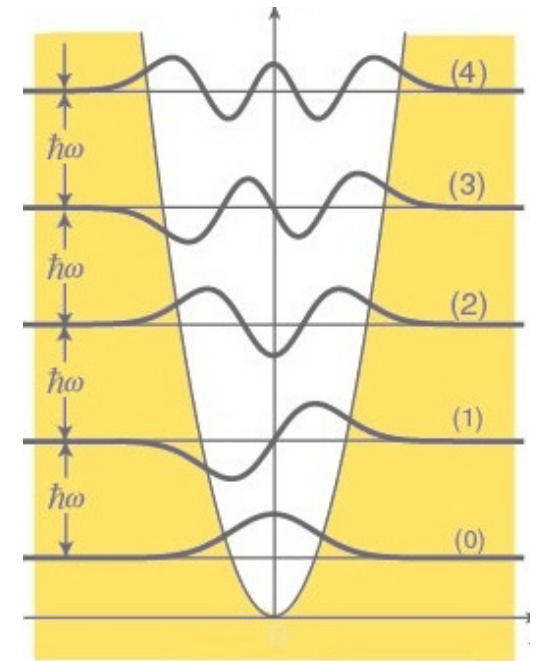


Our Goal

- High-precision studies of large, strongly interacting systems of fermions
 - Unitary fermions - ideal laboratory
 - Universal physics
 - Relevant for all systems with $r_0 \ll n^{-1/3} \ll a$
 - Nuclear matter
 - Cold atom experiments
 - Harmonic trap
- Up to N=20 fermions, <1% errors

Our Goal

- High-precision studies of large, strongly interacting systems of fermions
 - Unitary fermions - ideal laboratory
 - Harmonic trap
 - Control systematics
 - Small S/N problem
 - Trap confinement enhances overlap with ground state
- Up to $N=20$ fermions, $<1\%$ errors



Unitary Fermions

- Vary strength of interaction between fermions
➔ crossover from BEC state to BCS state
- Crossover point called unitarity

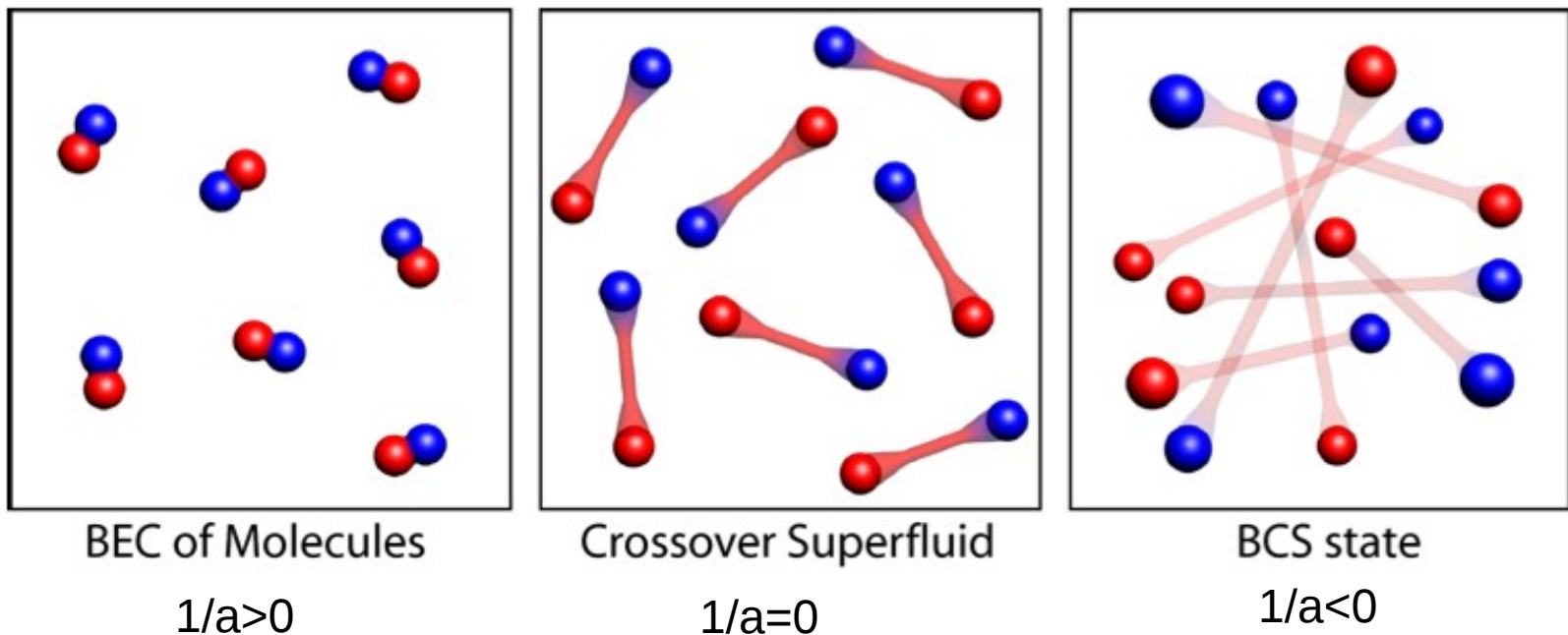


Figure: W. Ketterle and M. Zwierlein, [arXiv:0801.2500](https://arxiv.org/abs/0801.2500)

Unitary Fermions

→ $p \cot \delta = 0$ → no intrinsic scale except the density

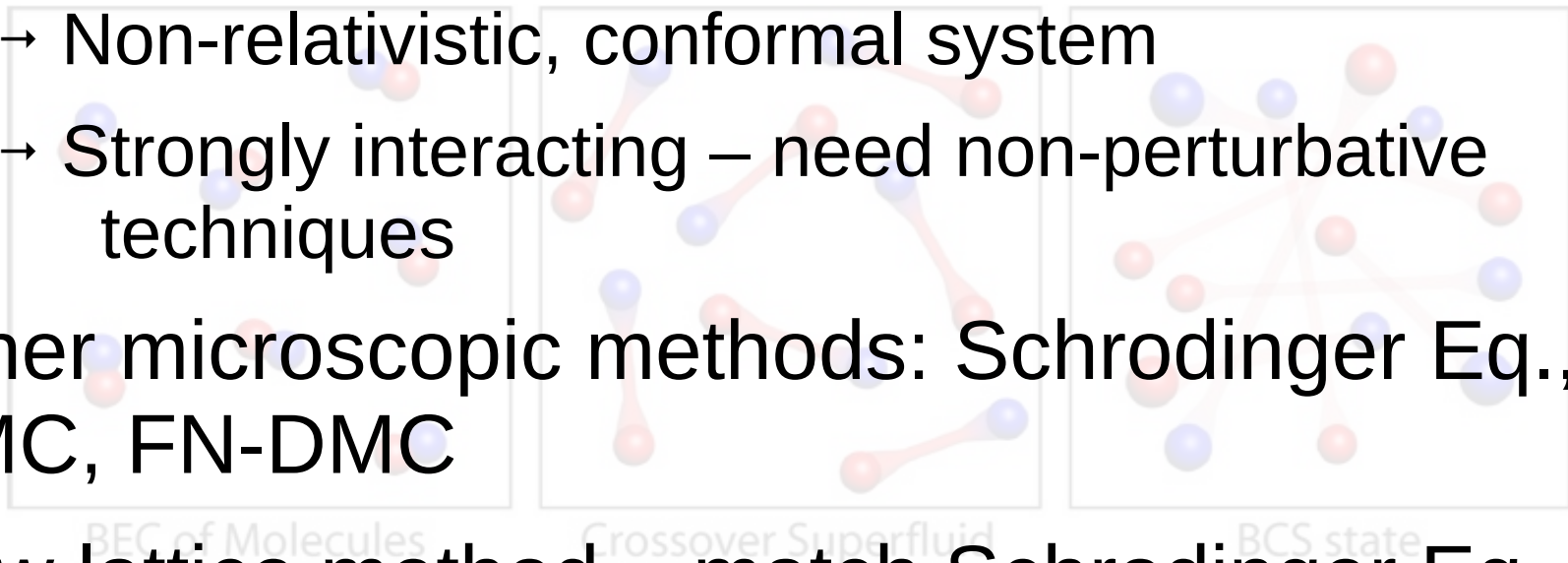
→ Displays universal features

→ Non-relativistic, conformal system

→ Strongly interacting – need non-perturbative techniques

• Other microscopic methods: Schrodinger Eq., GFMC, FN-DMC

• New lattice method – match Schrodinger Eq. to within 1%



Universal Quantities

Bertsch parameter

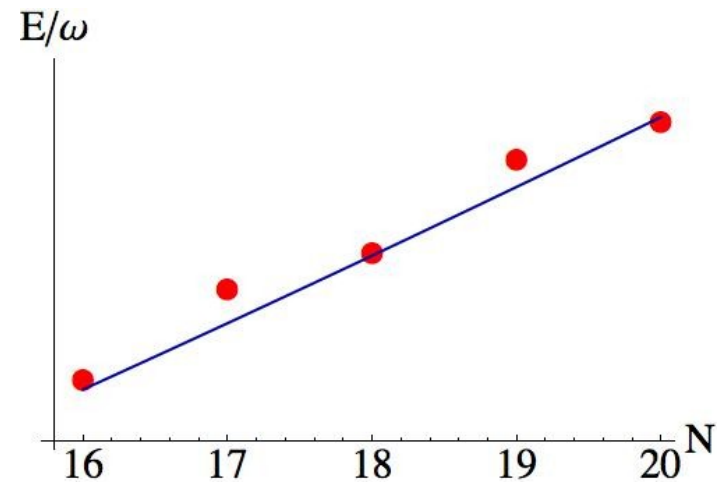
- $\mu_{int}(n) = \xi \mu_{free}(n)$
- Unitary fermions in an SHO:

$$E_{int}(N, \omega) = \sqrt{\xi} E_{free}(N, \omega)$$

Pairing gap

- Energy cost to break fermion pairs
- Even-odd splitting:

$$\Delta(N) = E(N) - \frac{1}{2}(E(N-1) + E(N+1))$$



Universal Quantities

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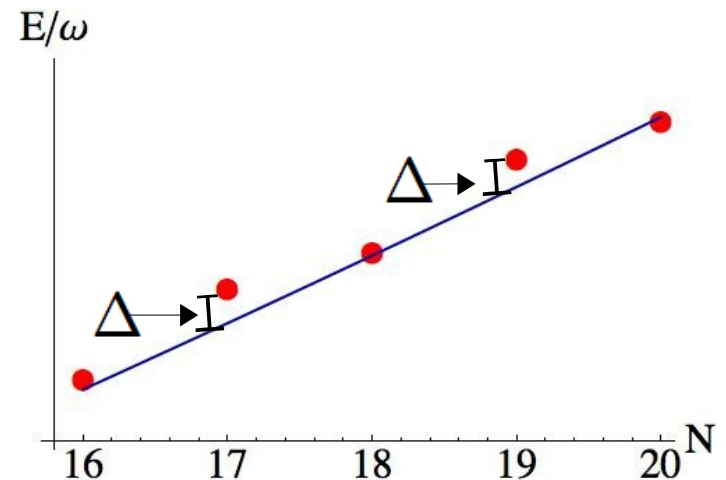
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Lattice Theory

- Non-relativistic fermions interacting via point interactions
- No fermion determinant \longrightarrow quenched = exact
- Highly improved interactions give $p \cot \delta = 0$ to arbitrary order in p^2
- Harmonic trap easily implemented:

$$\mathcal{T} = \mathcal{D}^{-1/2}(1 - \mathcal{V})\mathcal{D}^{-1/2}$$

$$\mathcal{V} = V_{int} + V_{SHO}$$

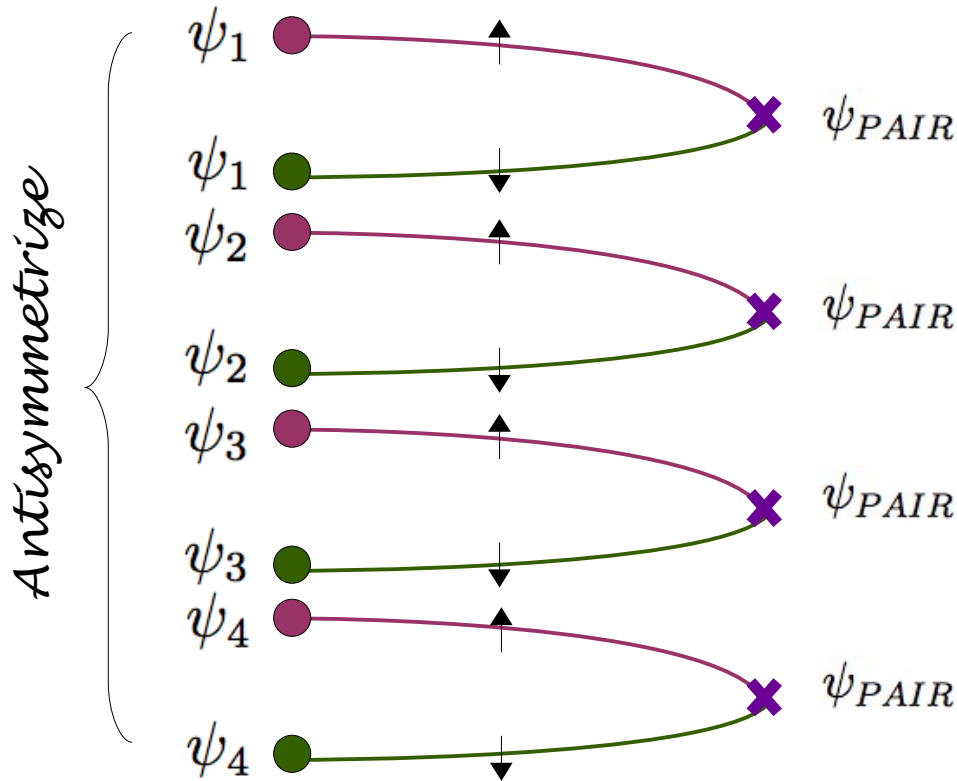
$$V_{SHO} = 1 - e^{-\frac{1}{2}m\omega^2 \sum_{i=1}^3 (L_i/2 - x_i)^2}$$

* Details introduced in talk by M. Endres

N-body Correlators

- Slater determinant of single-particle SHO states

N-body Correlators

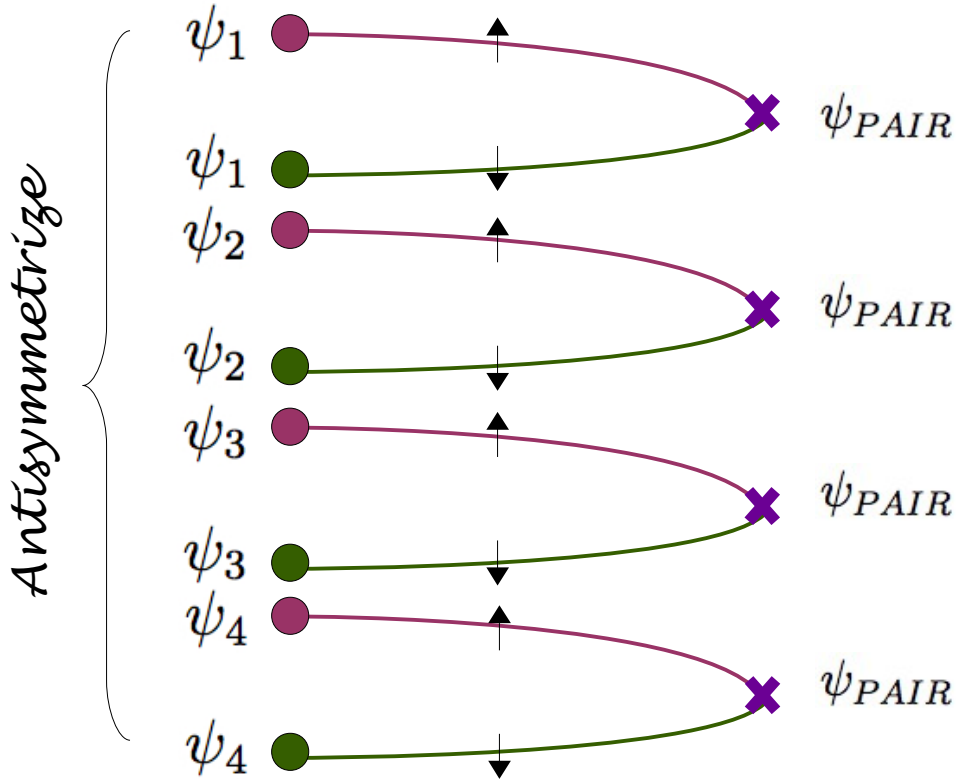


- Slater determinant of single-particle SHO states
- Include pair correlations*

$$\psi_{PAIR} \propto \frac{e^{-(x^2+y^2)/(2L_0^2)}}{|x-y|}$$

*J. Carlson, et al, Phys. Rev. Lett. **91** (2003)

N-body Correlators

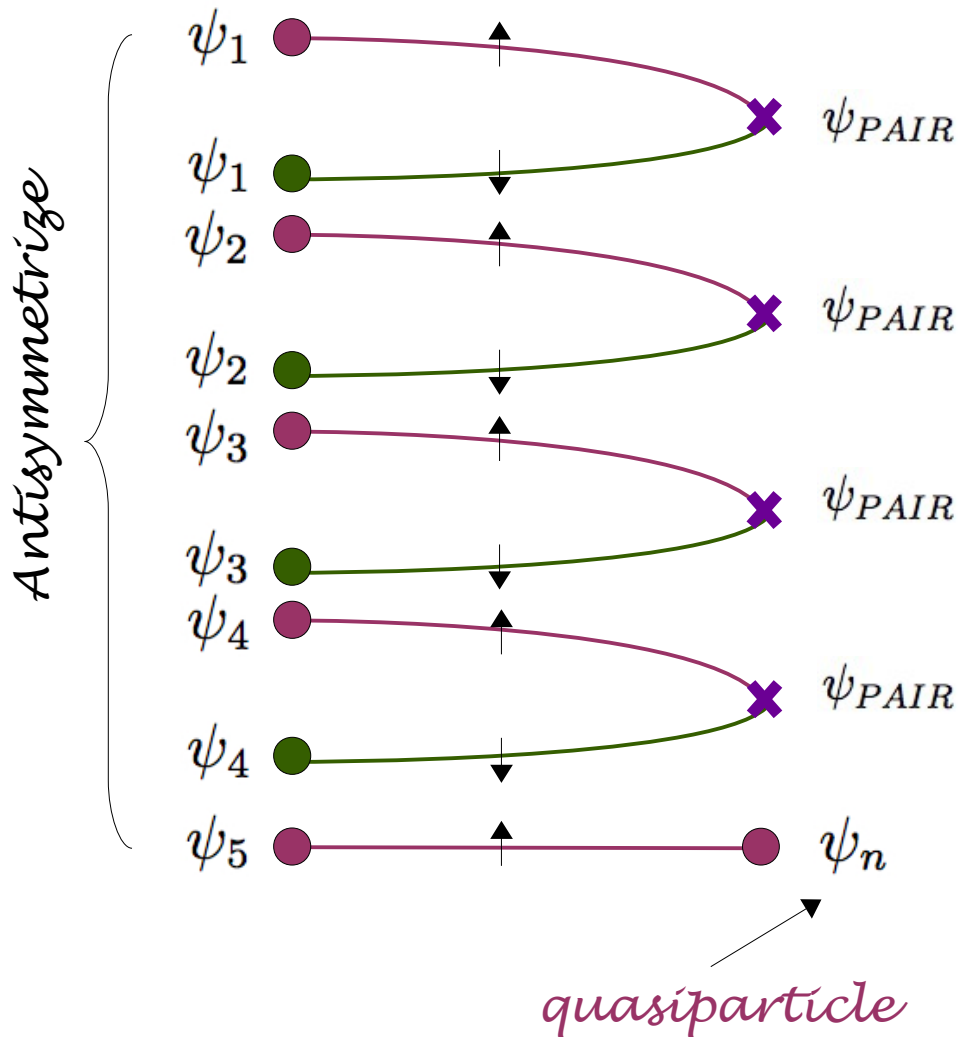


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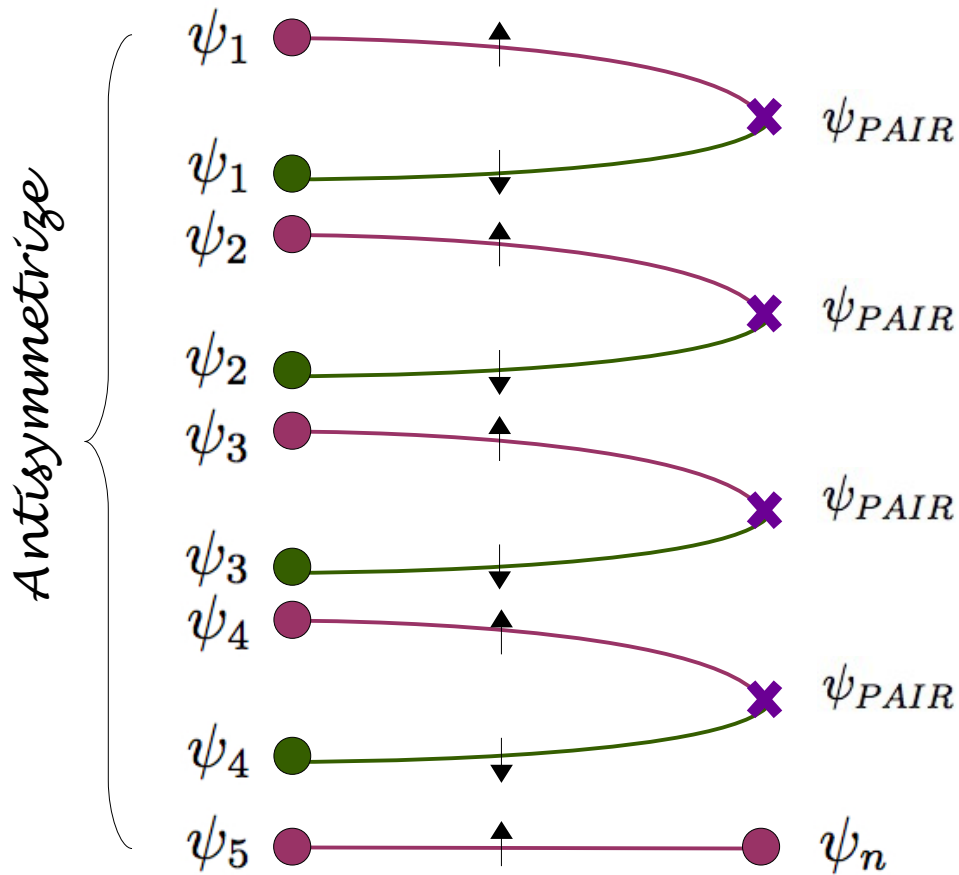
$$C_N(\tau) = \begin{vmatrix} \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_1 \psi_1 \rangle & \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_1 \psi_2 \rangle & \cdots & \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_1 \psi_{N/2} \rangle \\ \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_2 \psi_1 \rangle & \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_2 \psi_2 \rangle & \cdots & \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_2 \psi_{N/2} \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_{N/2} \psi_1 \rangle & \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_{N/2} \psi_2 \rangle & \cdots & \langle \psi_{PAIR} | (\mathcal{T})^\tau | \psi_{N/2} \psi_{N/2} \rangle \end{vmatrix}$$

N-body Correlators



- Slater determinant of single-particle SHO states
 - Include pair correlations
 - For odd N , add single particle state at sink
- Replace n th row in slater matrix

N-body Correlators



- Slater determinant of single-particle SHO states
- Include pair correlations
- For odd N , add single particle state at sink
- For SHO, either is sufficient

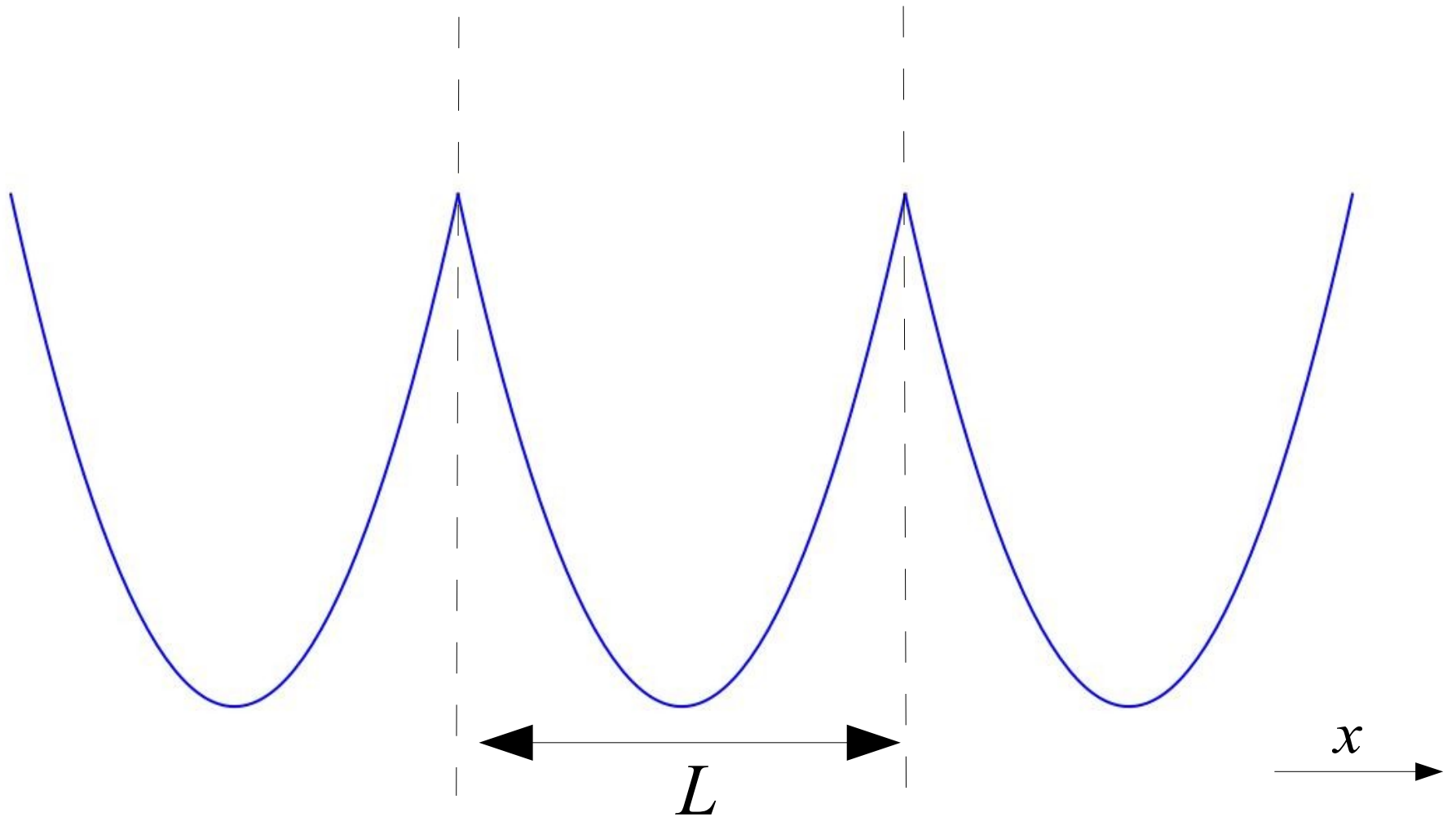
Systematic Errors

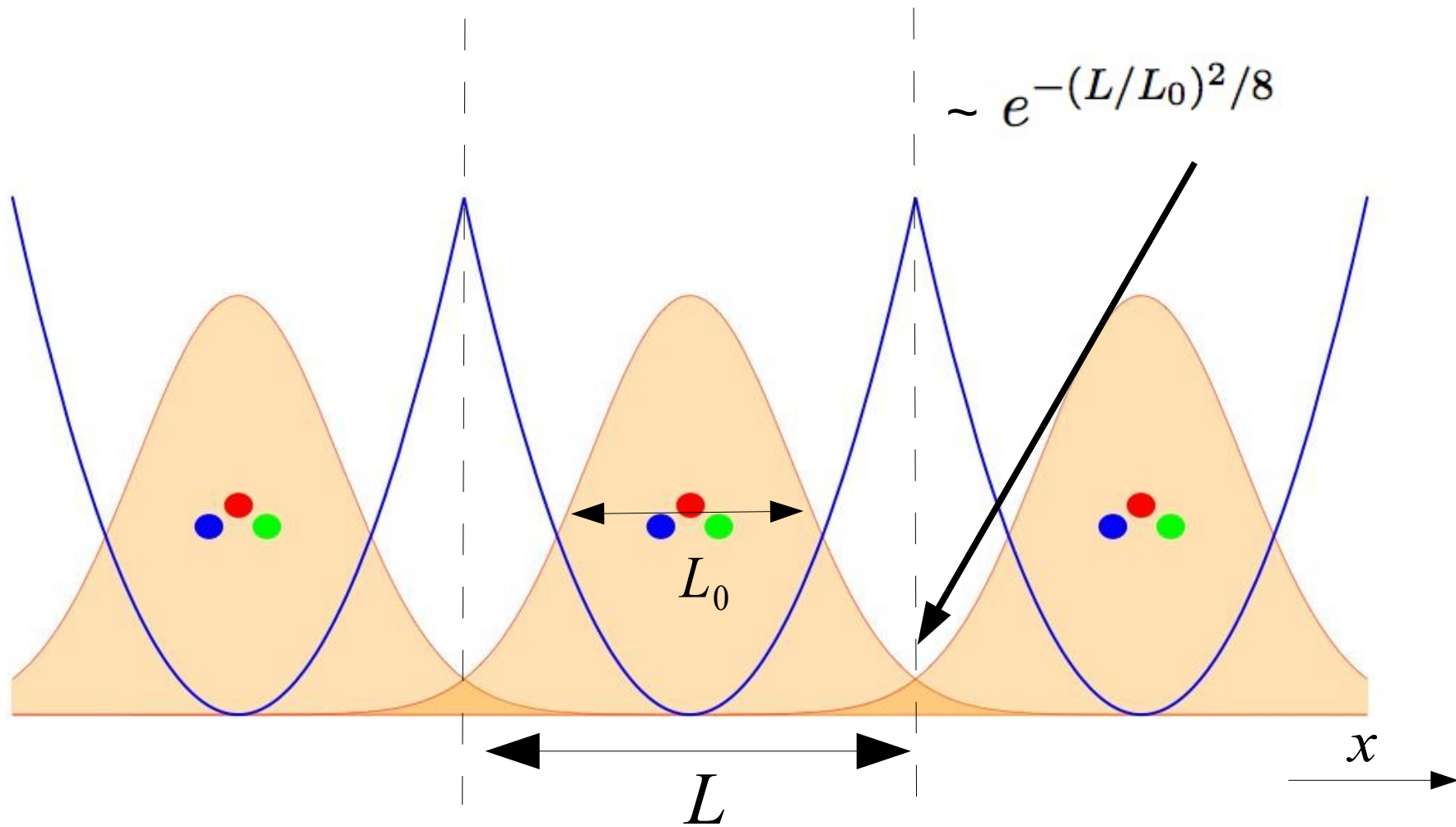
- Tunable scales set finite volume, finite lattice spacing (b_t , b_s) effects
 - ωb_t → temporal discretization error
 - b_s/L_0 → spatial discretization error
 - L_0/L → finite volume error
- Temporal discretization error
 - Exponential form of SHO offers some improvement
 - Ensure small b_t errors by choosing small ω

$$L_0 = (m\omega)^{-1/2}$$

Spatial Errors

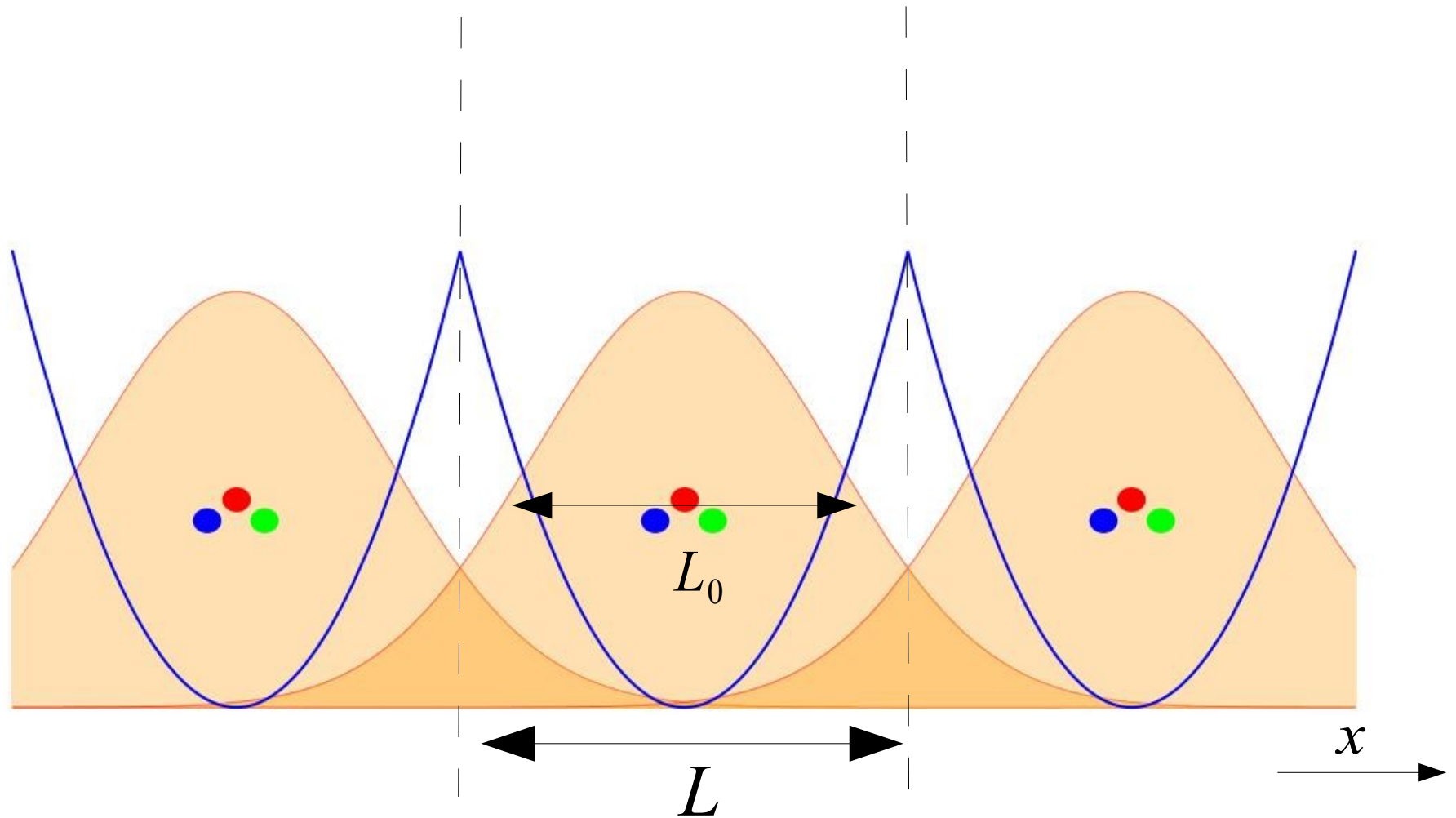
Position space potential: PeriodicBC



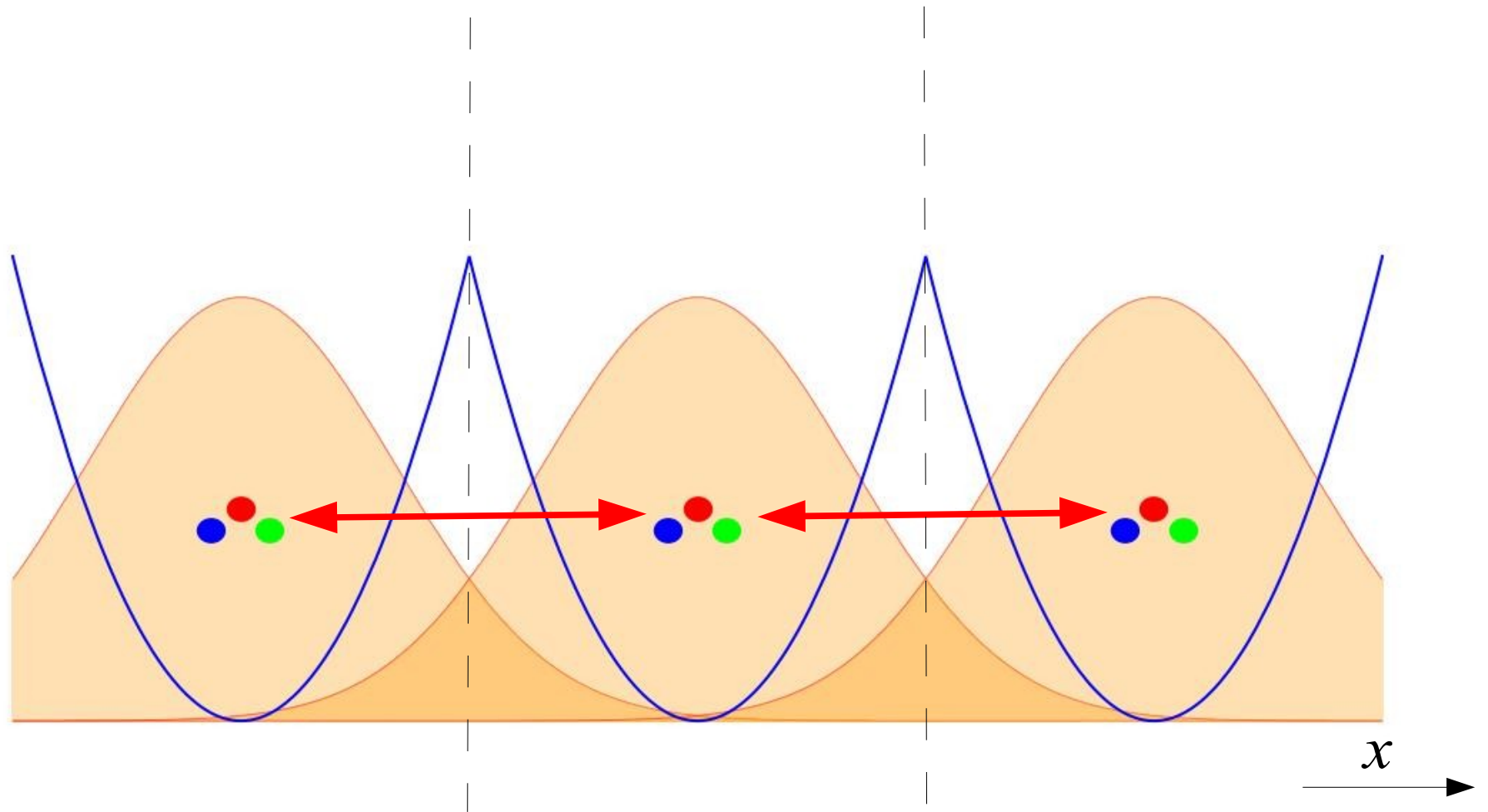


(ω fixed)

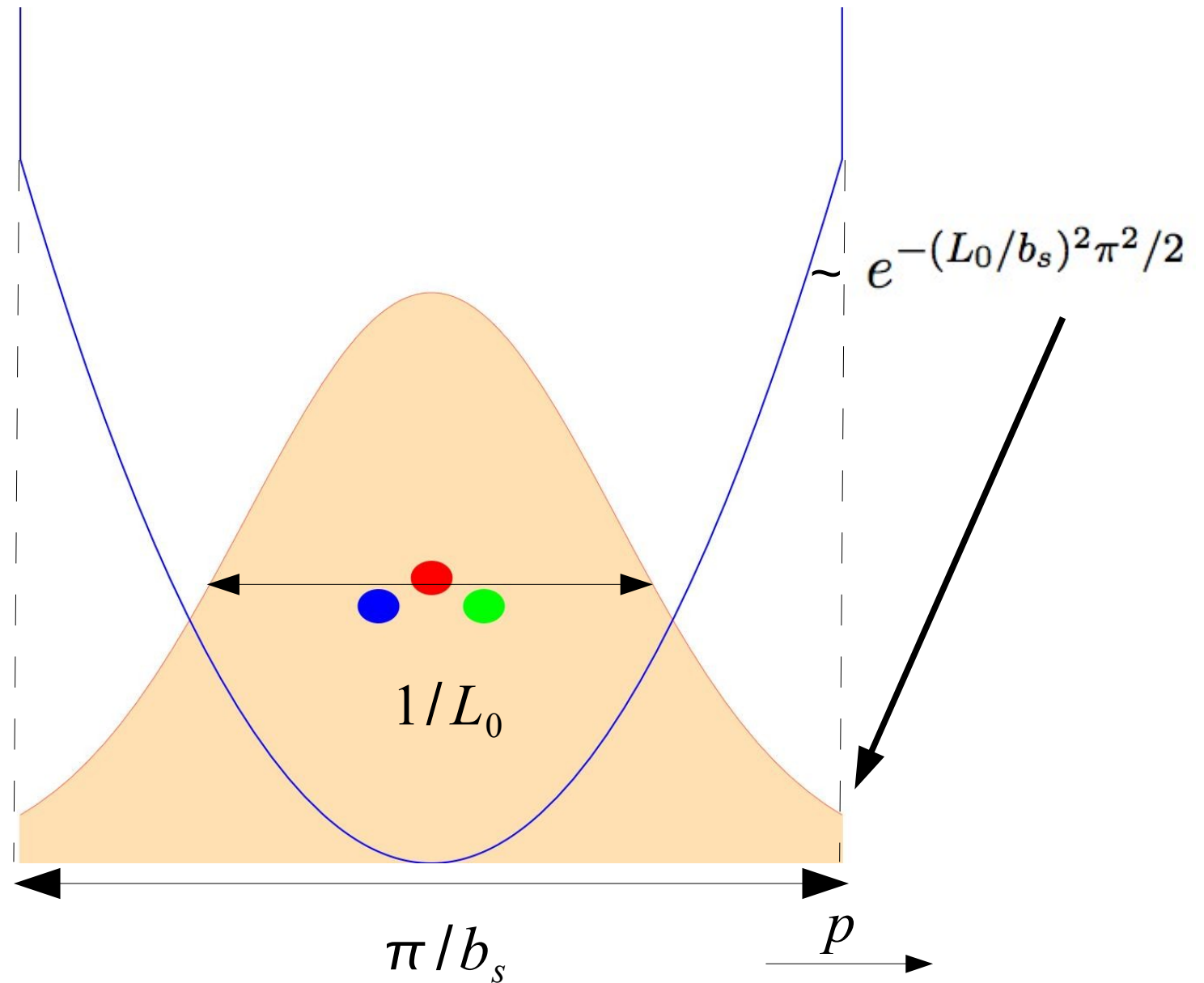
Increase L_0 , keep L fixed



Interactions with image charges lower energy

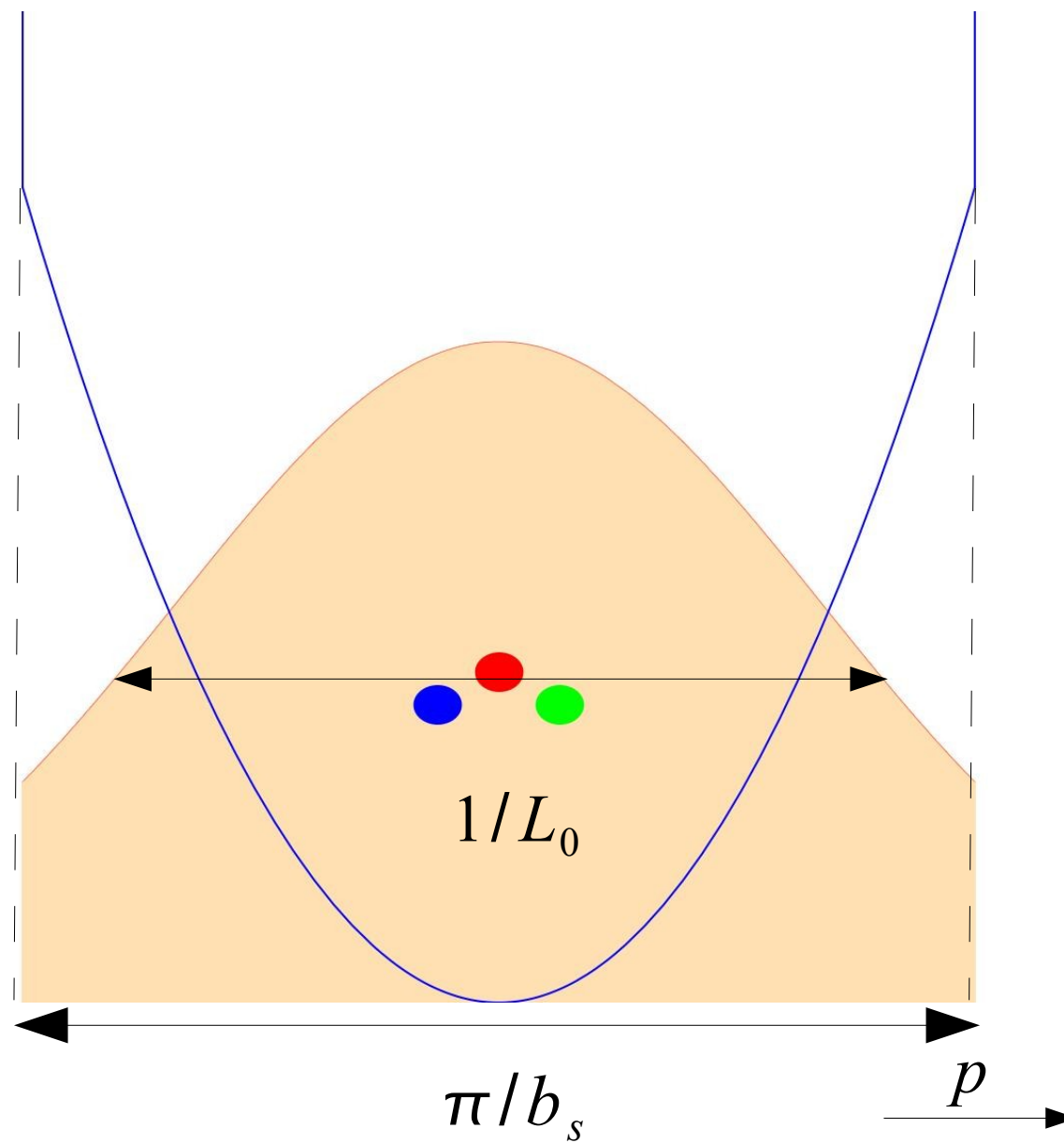


Momentum Space: Hard Cutoff

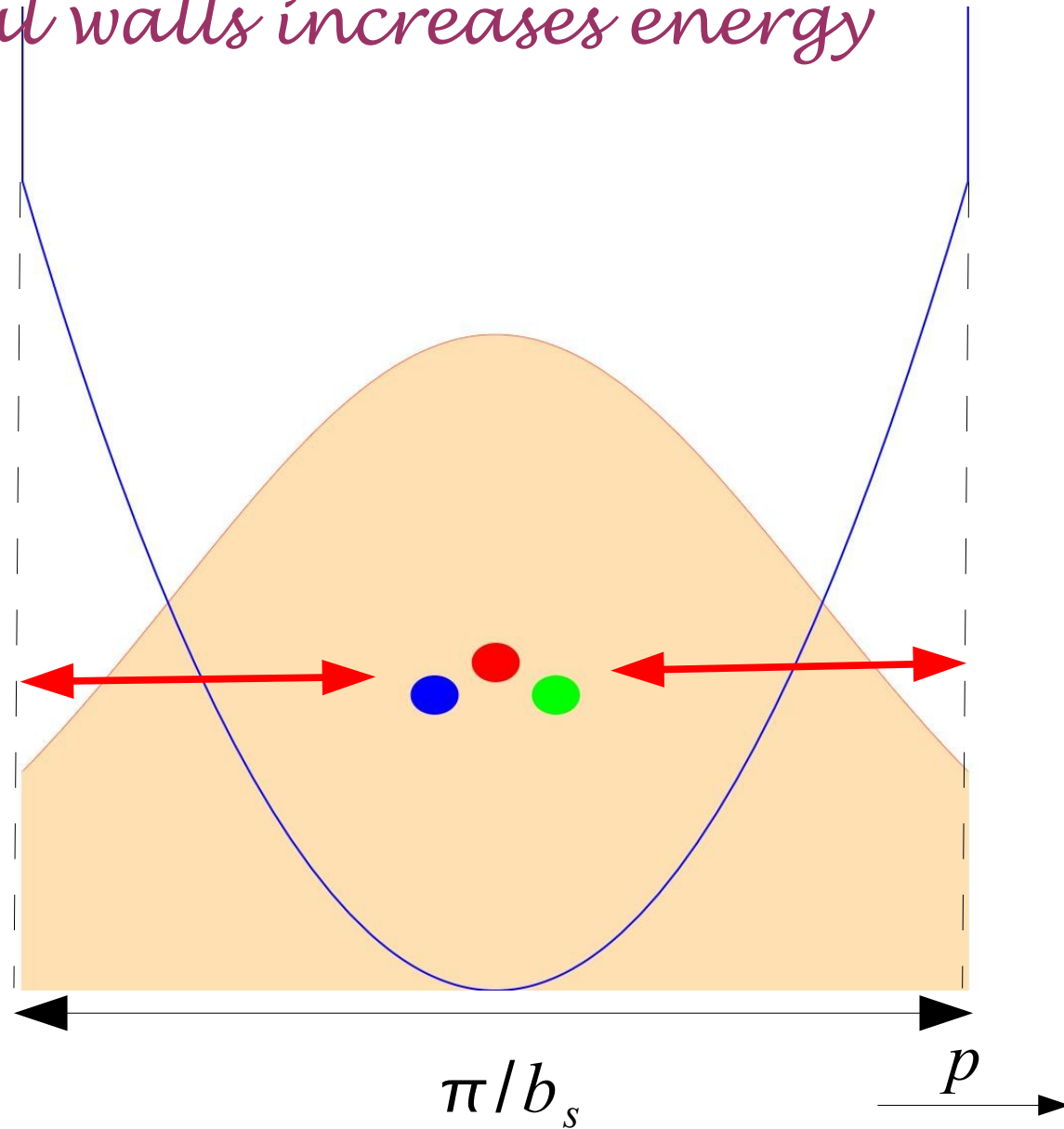


Reduce L_0

(ω fixed)



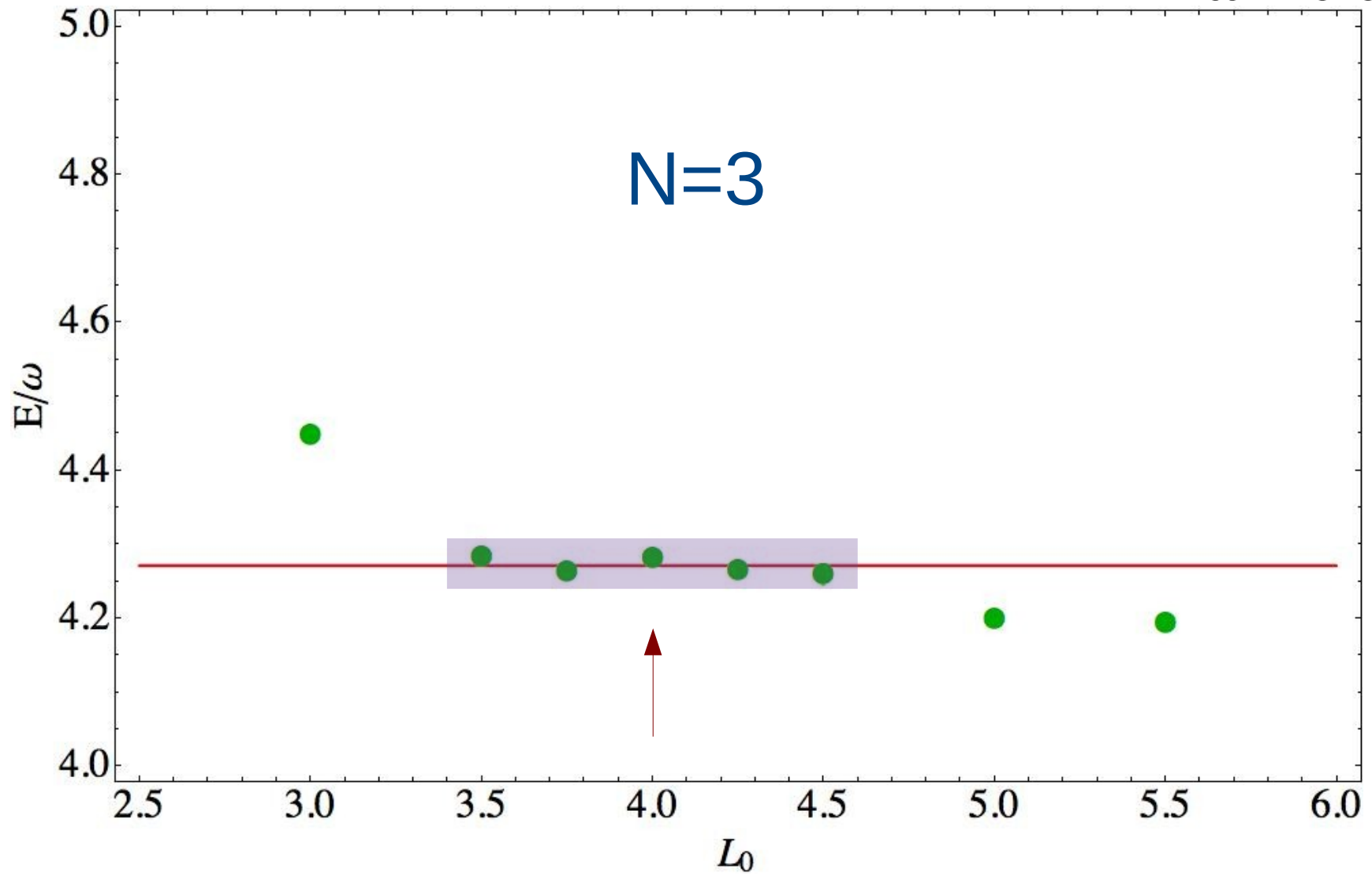
Reduce L_0 - more sensitivity to infinite potential walls increases energy



Spatial Errors

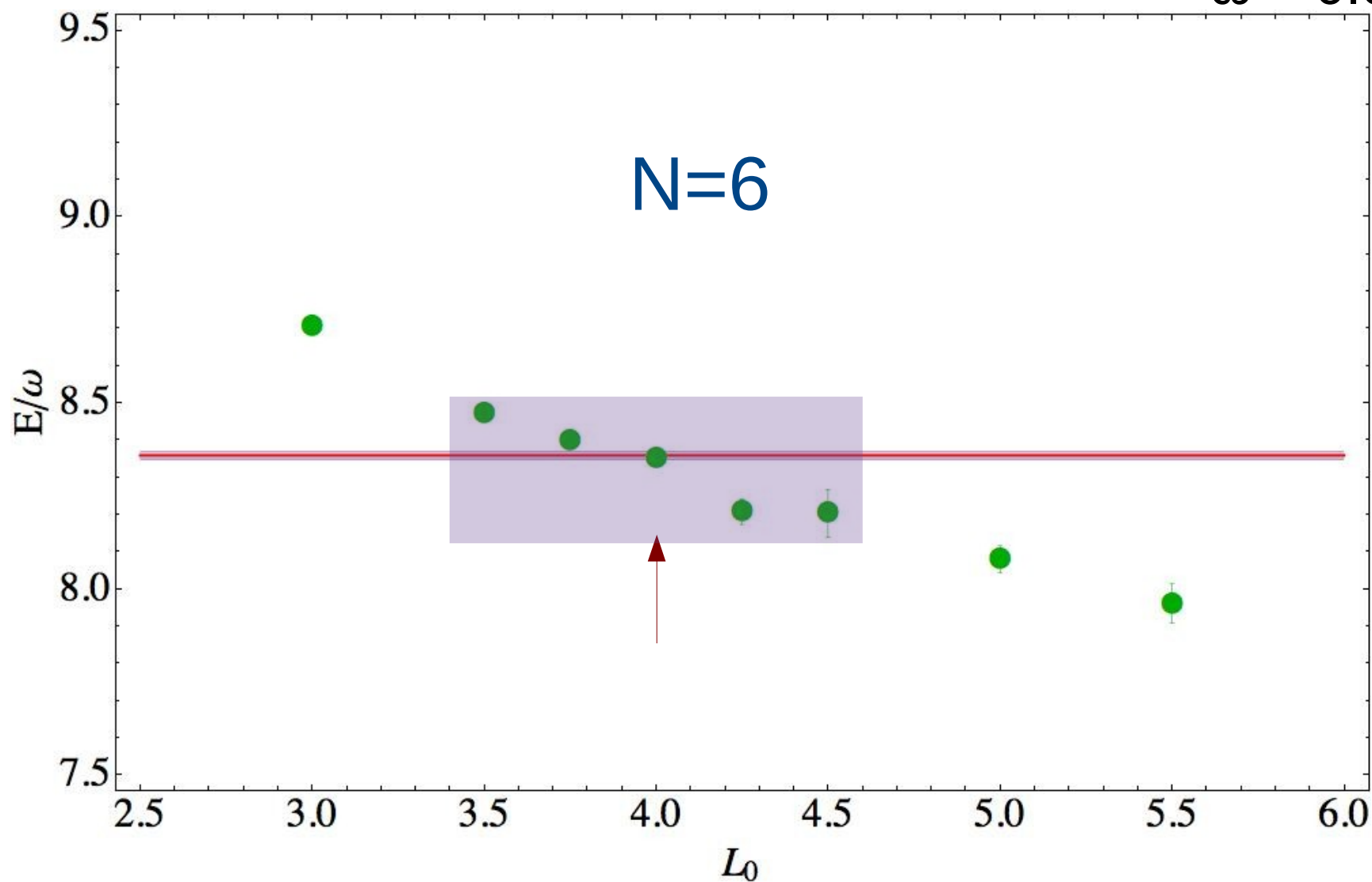
- Both finite volume and spatial discretization errors affected by changing L_0
 - Finite volume errors push energy down for large L_0
 - Discretization errors push energy up for small L_0
- Performed tests at various values of L_0 to choose ideal value

$L = 32$
 $\omega = 0.013$



F. Werner and Y. Castin, Phys. Rev. Lett. **97**. 150401 (2006)

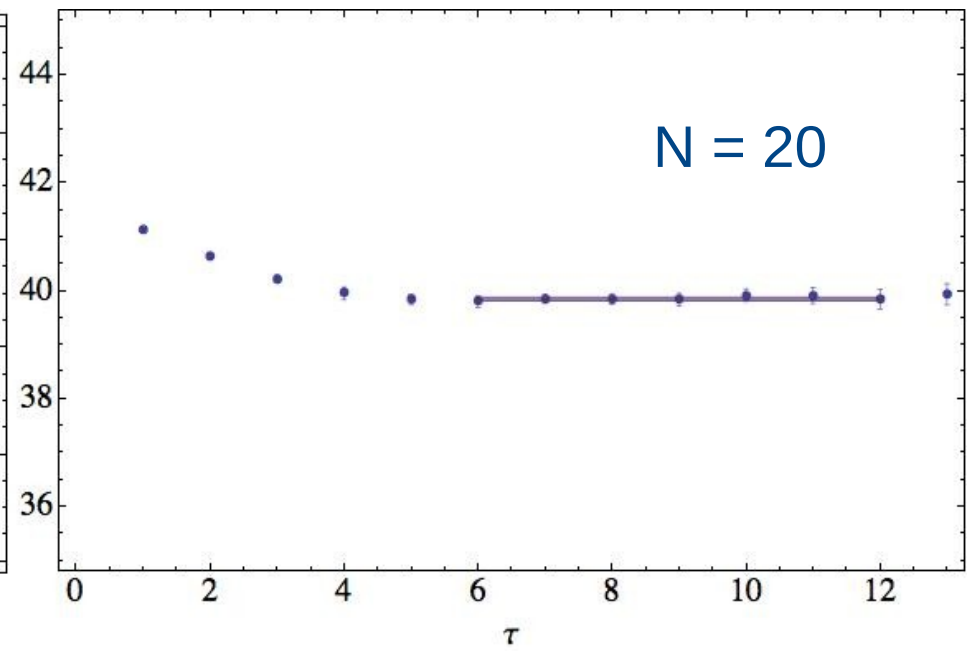
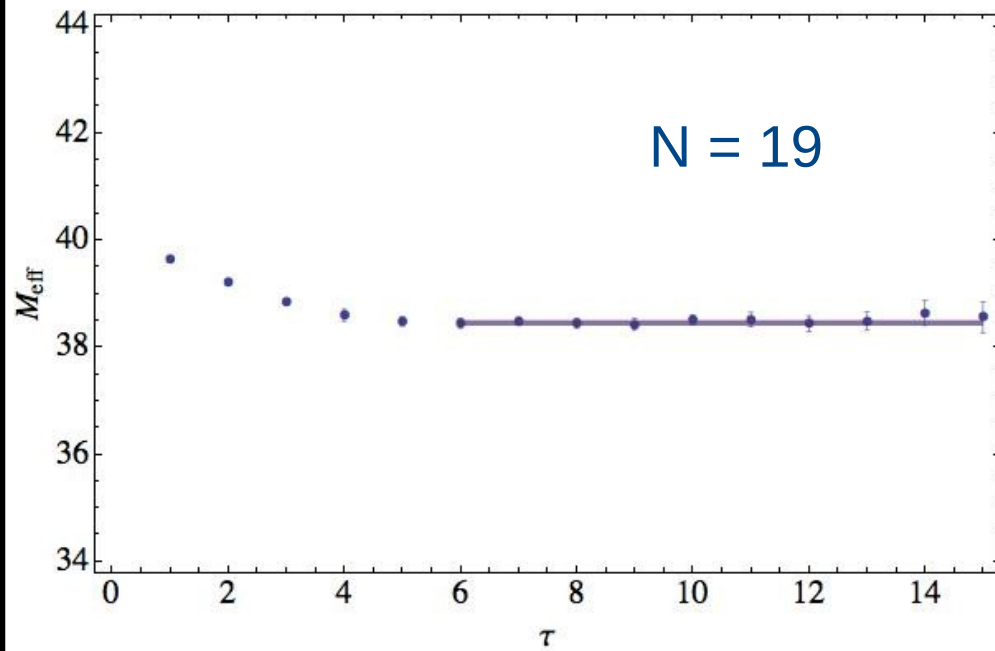
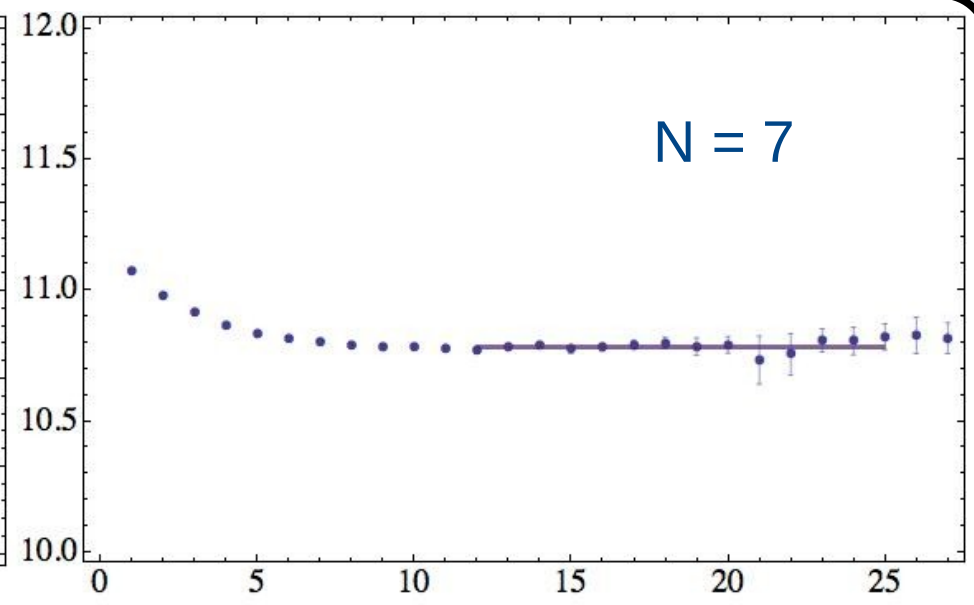
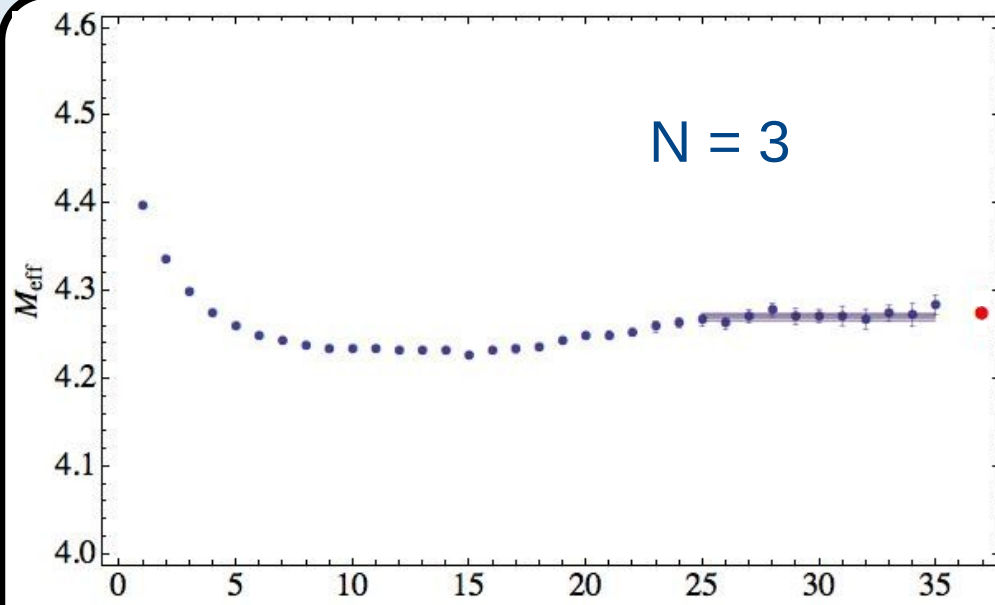
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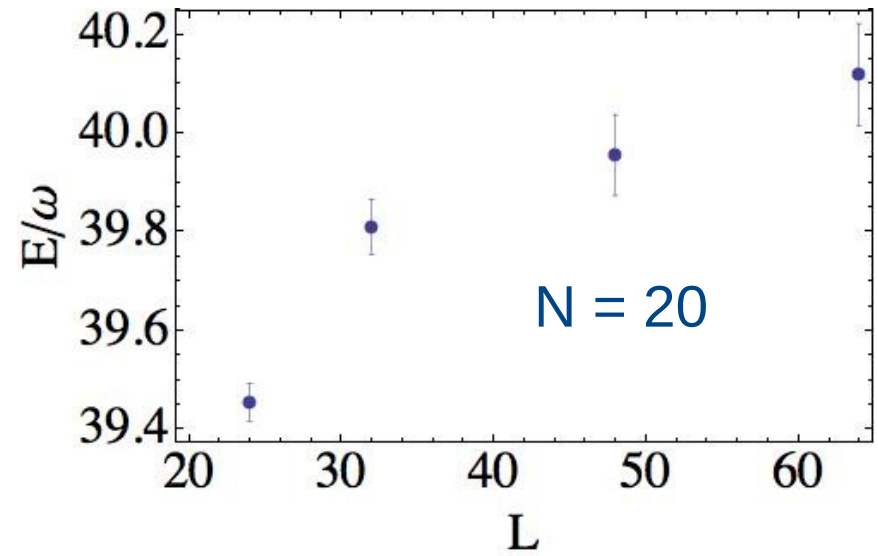
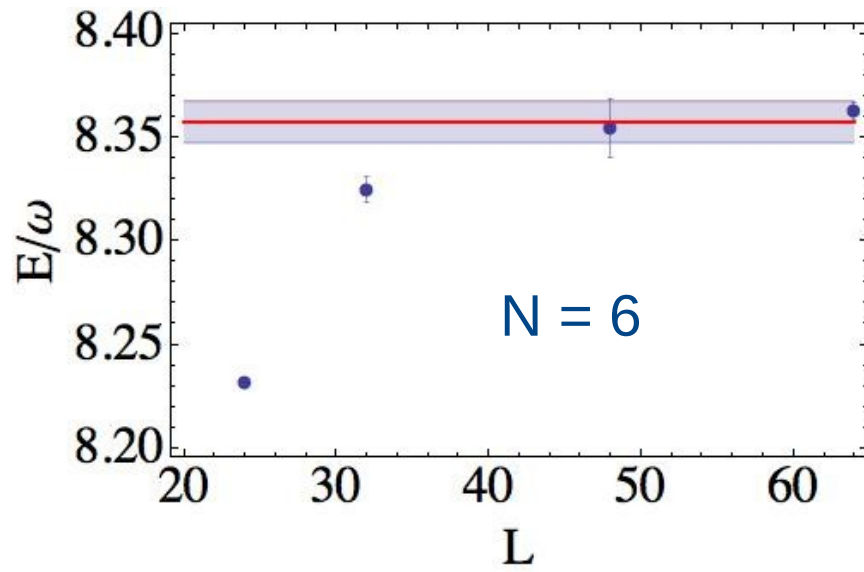
 D. Blume, private communication

Results

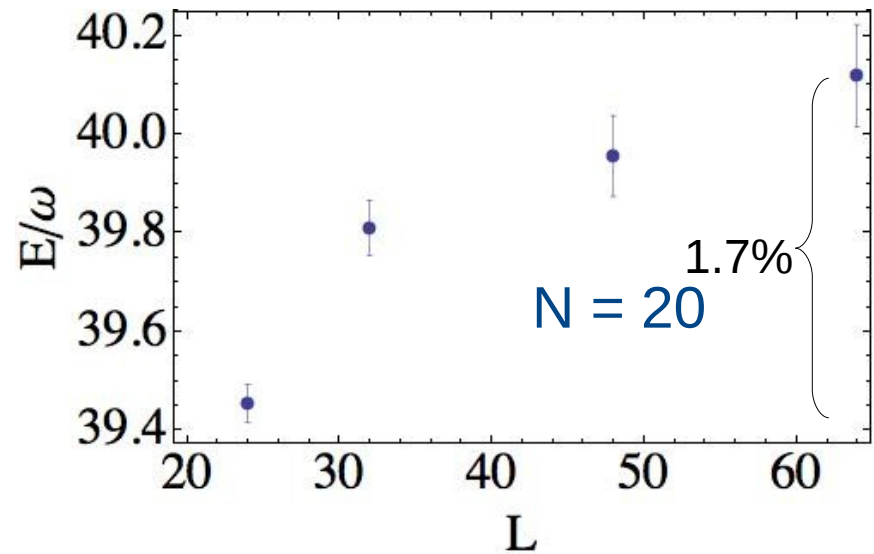
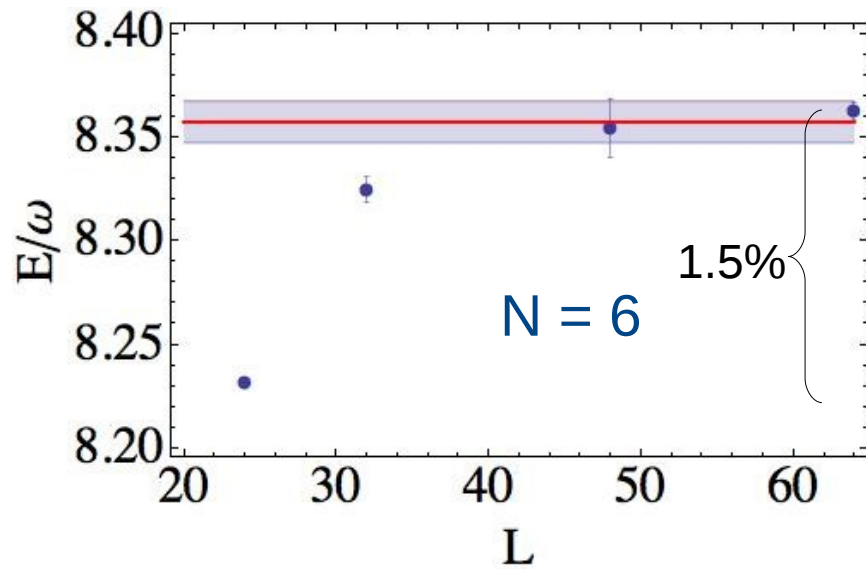
A large, multi-tiered fountain is the central focus, with a prominent vertical jet of water. The fountain is surrounded by a paved walkway where several people are walking. In the background, there are lush green trees and a large, light-colored, dome-shaped structure, possibly a stadium or arena, under a clear blue sky.



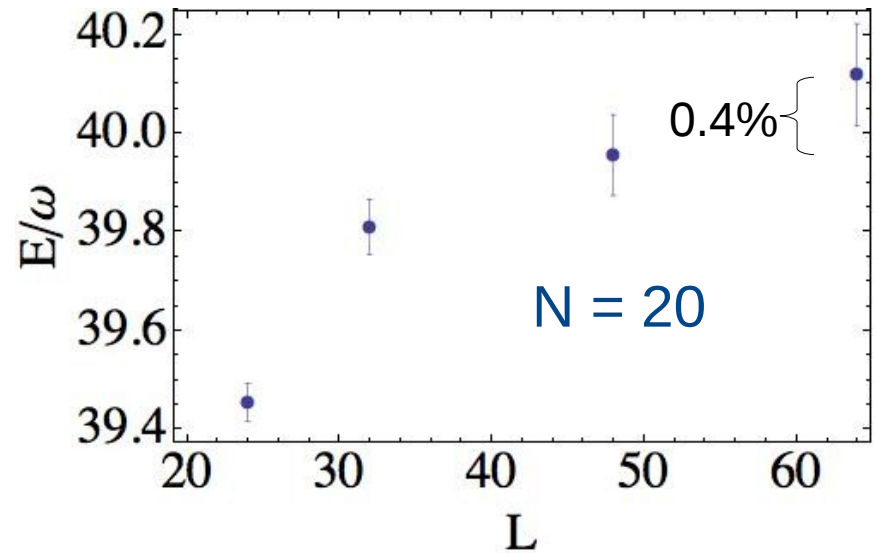
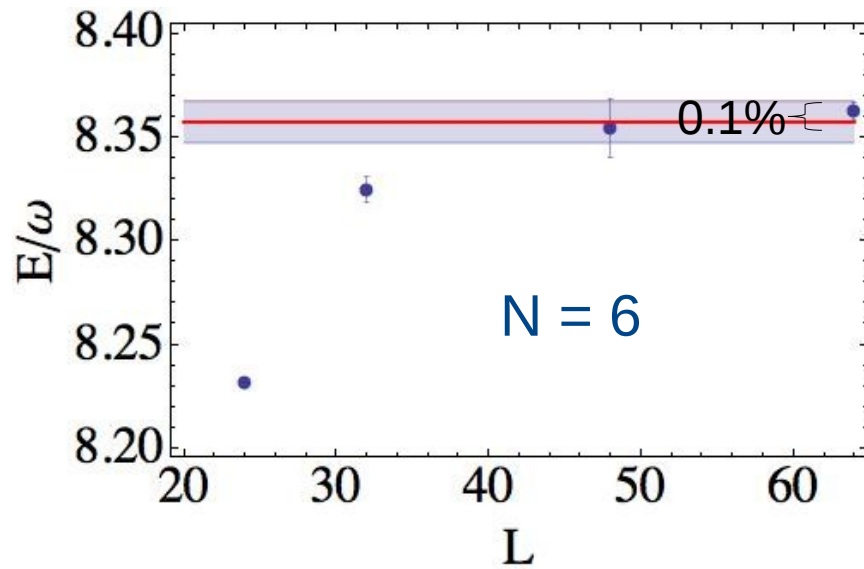
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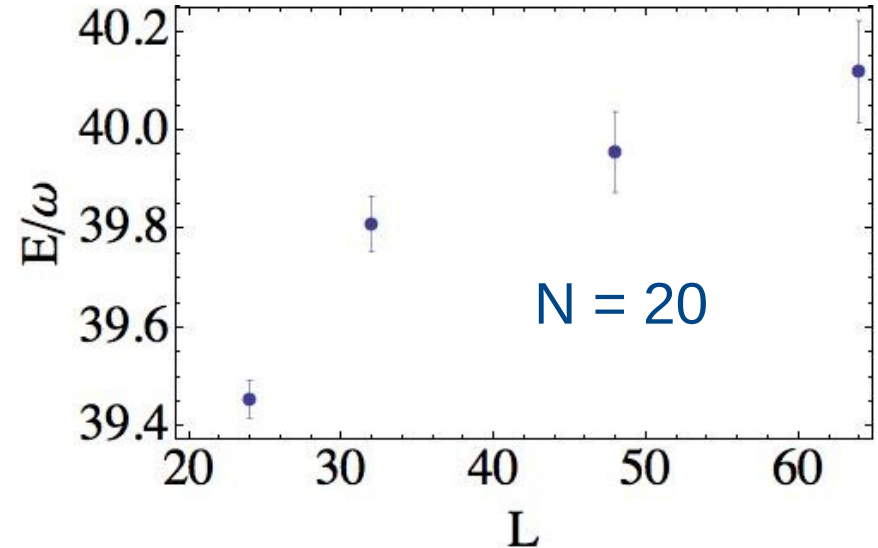
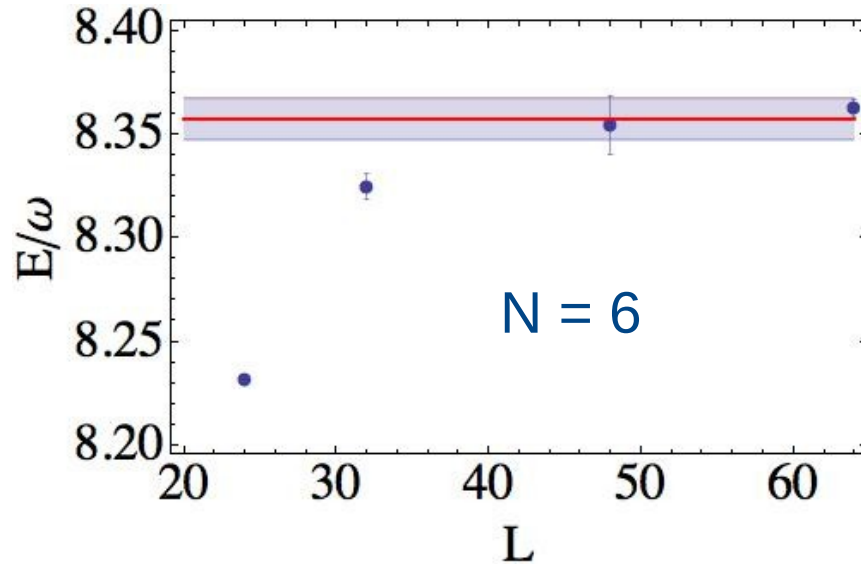
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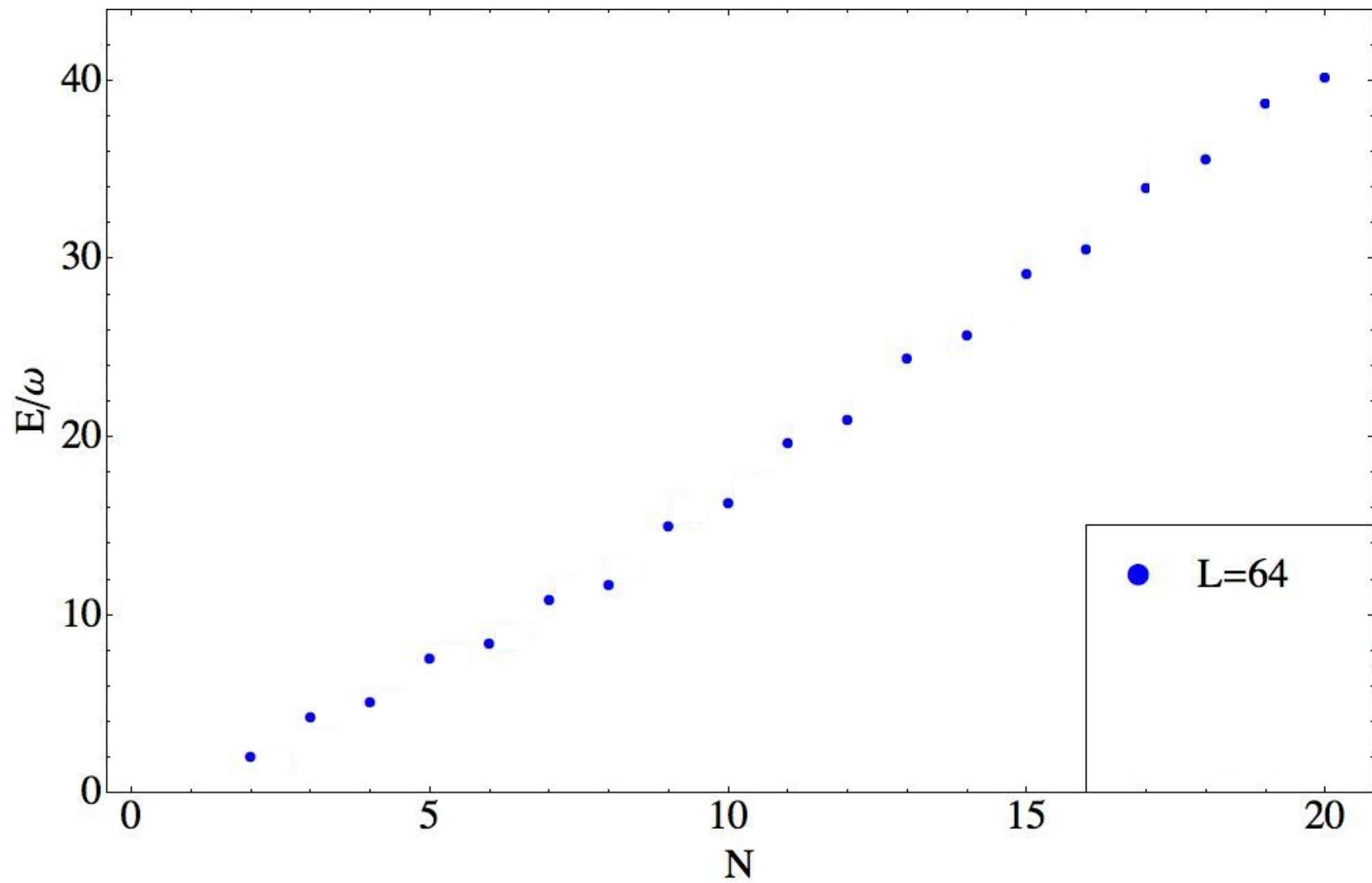
N	This Work	Comparison	% Deviation
3	4.253(2)(4)	4.2727*	0.5
4	5.058(1)(1)	5.028(20)†	0.6
5	7.513(3)(2)	7.457(10)‡	0.8
6	8.338(4)(5)	8.357(10)‡	0.2

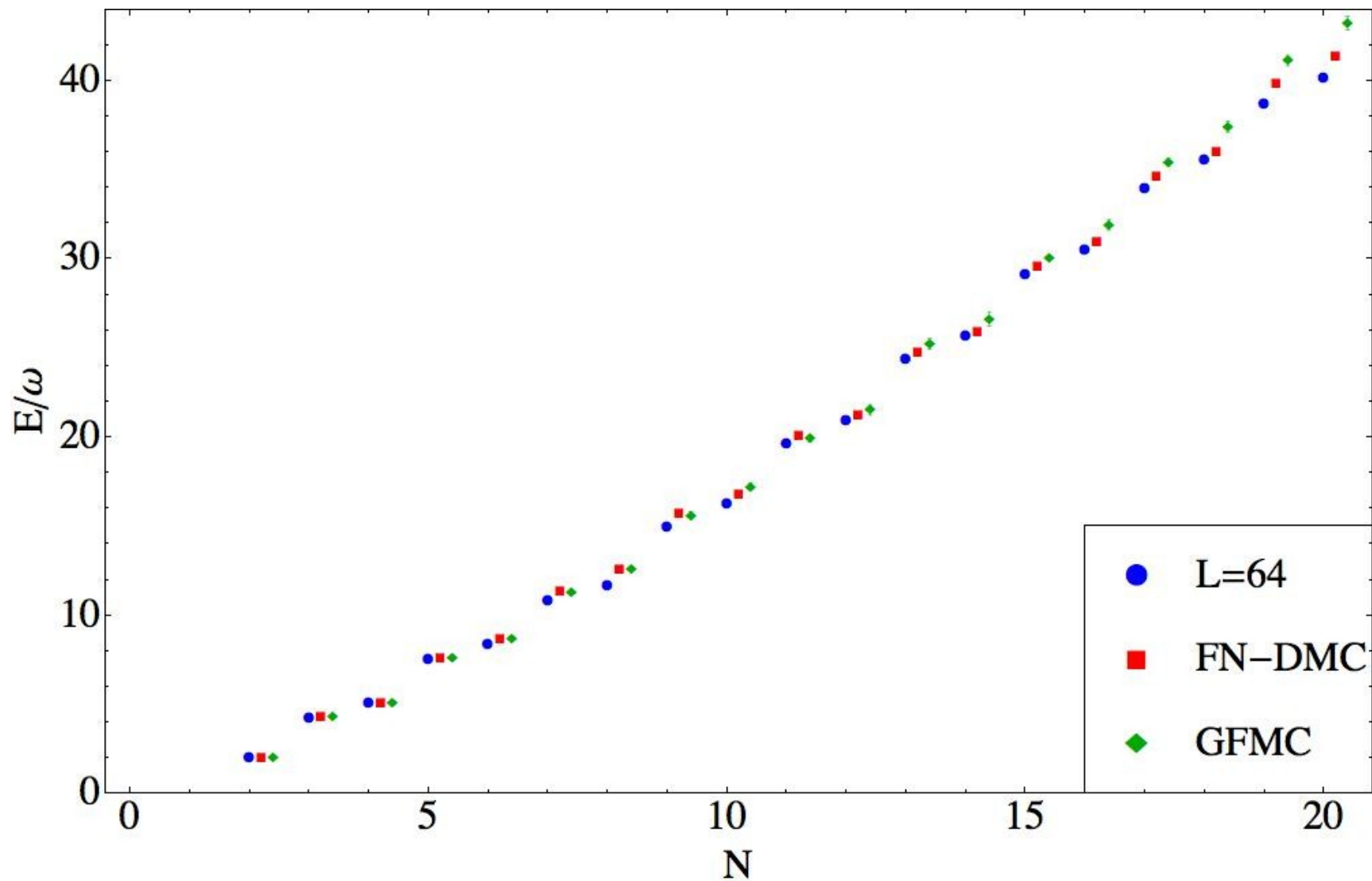
*F. Werner and Y. Castin, Phys. Rev. Lett. **97**. 150401 (2006)

†D. Blume, J. von Stecher, and C. Greene, Phys. Rev. Lett. **99**. 233201 (2007)

‡D. Blume, private communication

PRELIMINARY

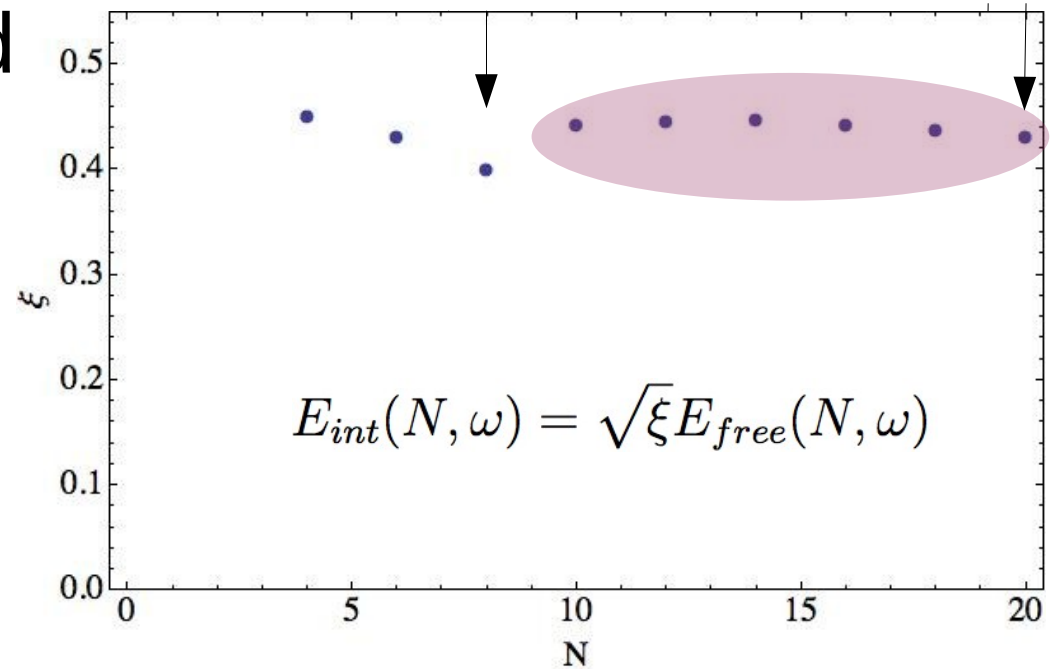




- FN-DMC: D. Blume, J. von Stecher, Chris H. Greene, arXiv:0708.2734
- ◆ GFMC: S. Y. Chang and G. F. Bertsch, arXiv:physics/0703190

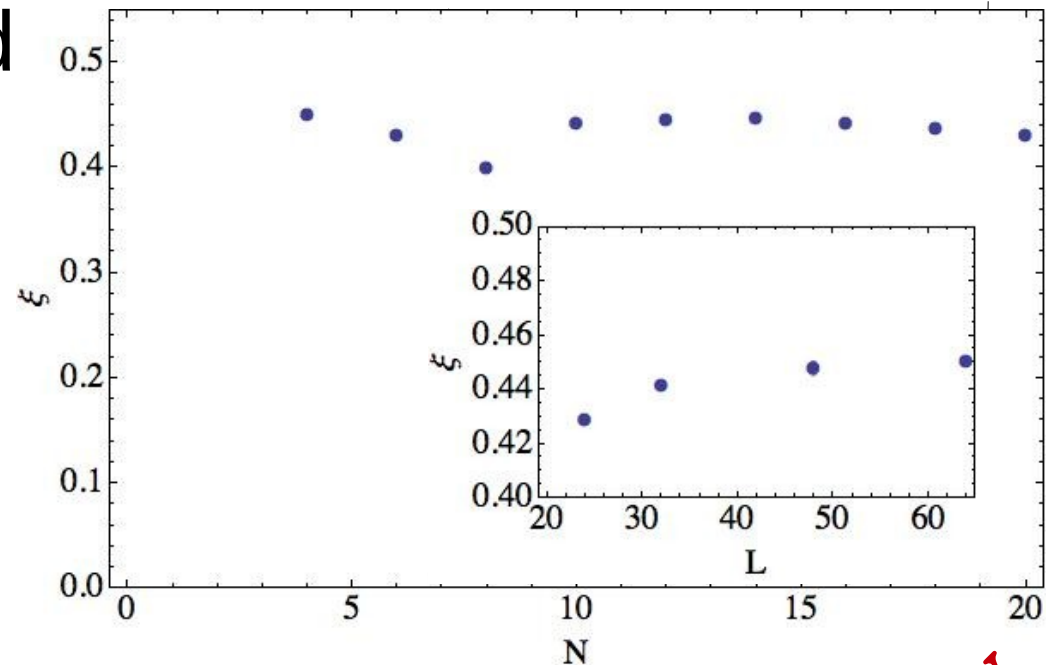
Bertsch Parameter

- Performed correlated fits to second shell
- Clear shell structure
➔ haven't reached thermodynamic limit



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- Finite volume effects smaller than 1%
- Results from SHO higher than from box



This Work: 0.450(1)
FN-DMC ~ 0.465
GFMC ~ 0.500

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GFMC: J. Carlson, et al (2008)
Lattice EFT: D. Lee (2008)

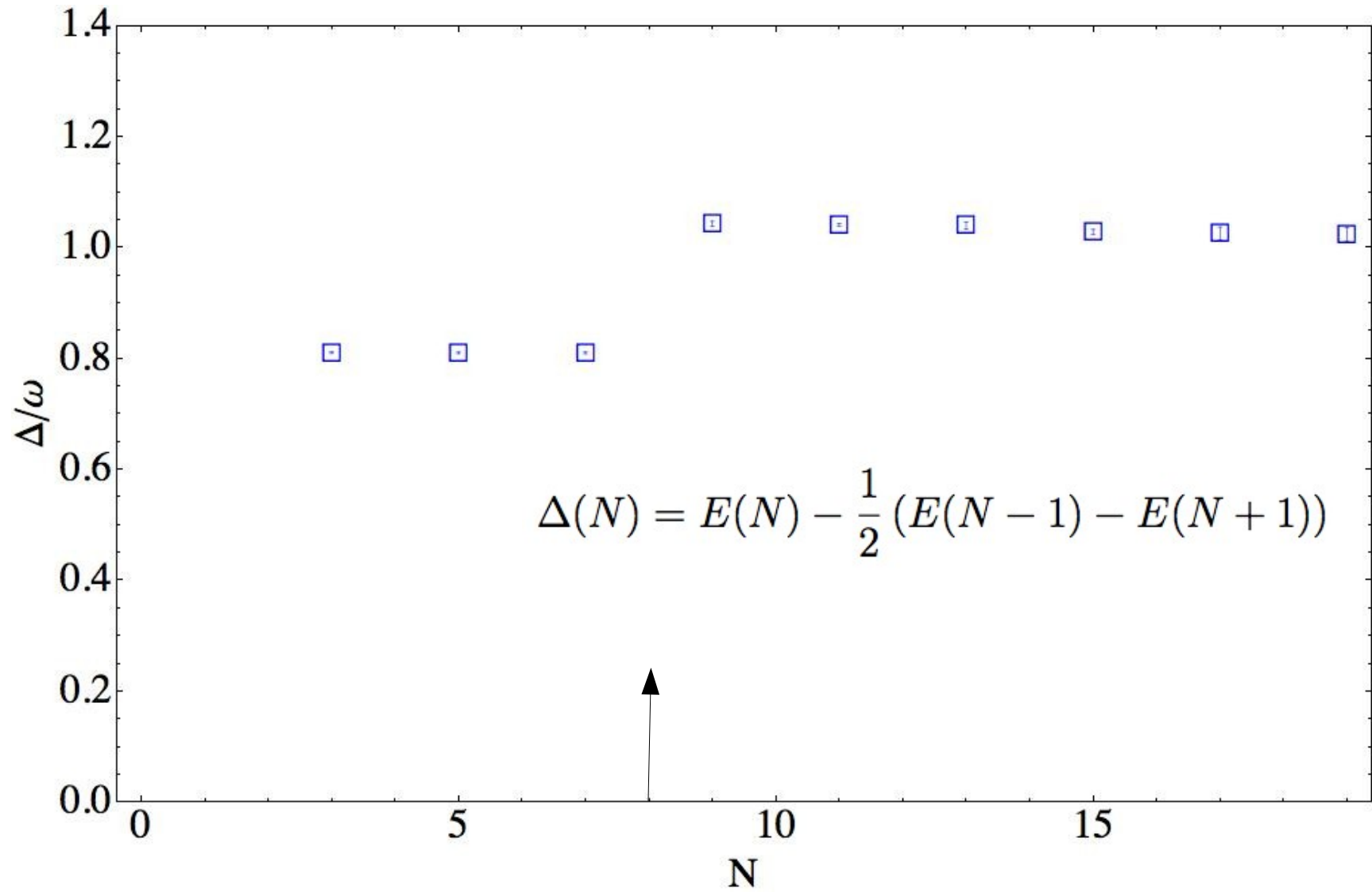
Box:

This Work: 0.412(4)
GFMC: 0.40(1)
Lattice EFT: 0.329(5)

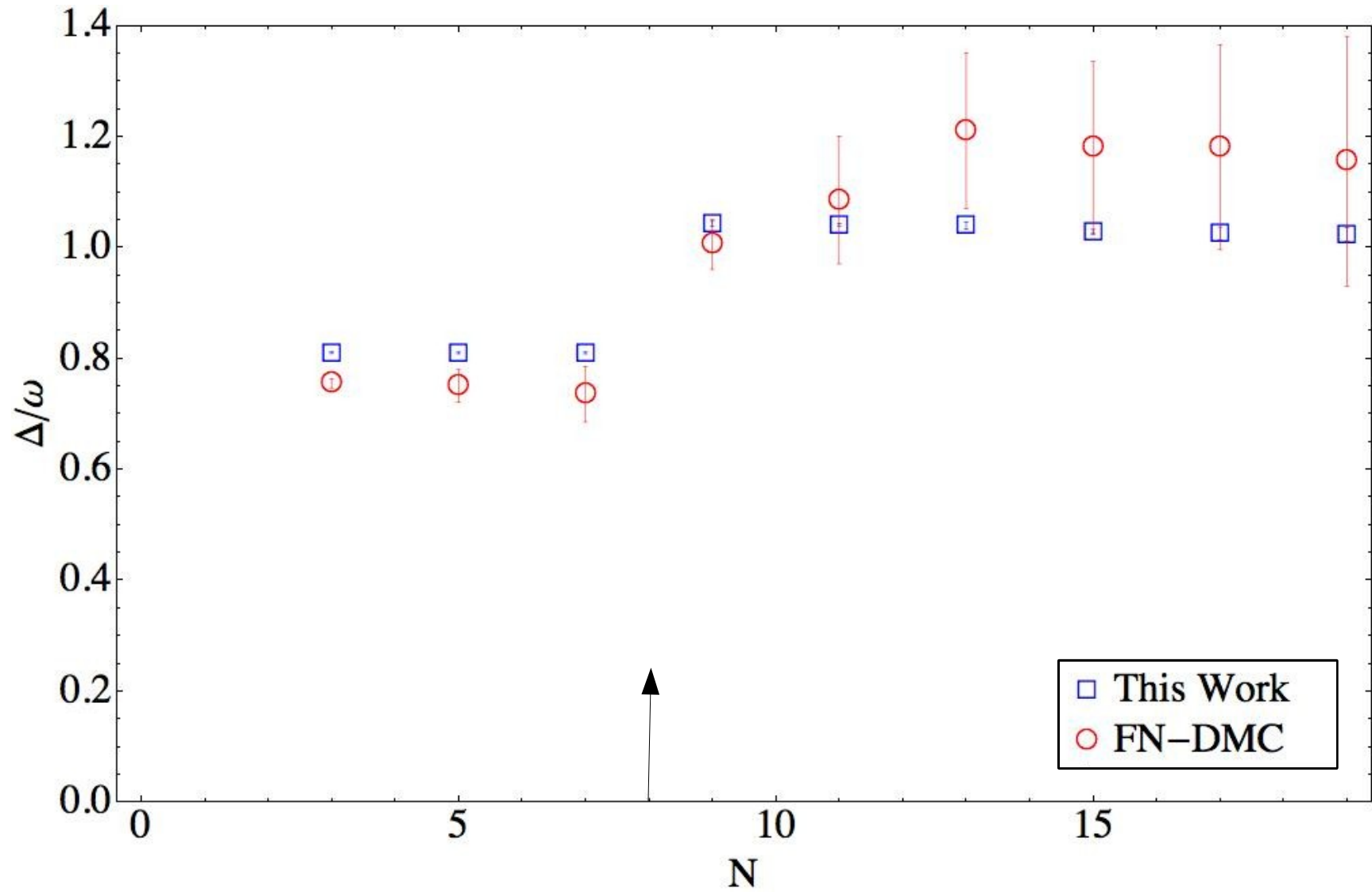
SHO:

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Gap



Gap



SHO Conclusions

- Possible to study large N to high precision
- Results consistent with high precision Schrodinger Eq. solutions for $N \leq 6$ to within 1%
- Better systematics and smaller errors than previous methods

Future directions

- Working to reduce lattice spacing errors further
- Ability to precisely tune $p \cot \delta$ suggests applicability to nuclear systems
- Trap confinement may be useful for studying bound states of hadrons

Acknowledgements

Parts of these calculations were performed on New York Blue (BG/L) at Brookhaven National Laboratory

