Condensate enhancement for mass generation in SU(3) gauge theory

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The Lattice Strong Dynamics collaboration

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Our first thoughts



How do the properties of gauge theory change with N_f, N_c and R ?

- * Chiral condensate enhancement $\frac{\langle \overline{\psi}\psi \rangle}{F^3}$
- * S parameter
- ***** Particle spectrum
- ***** Dirac operator eigenvalue spectrum

How should we proceed?



Choices

ETC is complicated enough:

stay with fundamental reps

Start from something we know:

lattice QCD -- SU(3) color then move to SU(2) color

\mathbb{H} Move slowly away from QCD and not too close to N_{fc}:

first do 6 flavors then move to 10 flavors

Chiral and flavor symmetries are crucial:

use DWF

To be able to observe enhancement:

use large cutoff (small a)

To be able to make direct comparisons:

Do a 2-flavor simulation at the same cutoff

Higher demands

□ Computing cost increases as N_f^{3/2}

□ The lattice must have cutoff much larger than the confinement scale to take advantage of slower running. Larger lattice needed as we approach the IRFP.

We do not know the answer

Simulations

- ★ Lattice Volume is 32³ x 64
- ***** Iwasaki gauge action with DWF at $L_s = 16$
- ✤ Input fermion masses m_f = 0.005 to 0.03
- ★ m_{res} ~ 3x10⁻⁵ (2f), 8x10⁻⁴ (6f), 2x10⁻³ (10f)
- \star M_{π} L > 4
- **CPS:** HMC, multi-level simplectic integrator,

mass preconditioning, chronological inversion

Approximately matched lattice spacings :

 $\beta = 2.76$ (2f) $\leftarrow \rightarrow \beta = 2.1$ (6f) $\leftarrow \rightarrow \beta = 1.95$ (10f)

★ Goal: ~1,000 configurations per point

Scale matching 2f, 6f



Reasonable distance from cutoff with $\rm M_{\rho} \sim$ cutoff / 5

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Chiral perturbation theory

$$M_m^2 = \frac{2m\langle \bar{\psi}\psi \rangle}{F^2} \left\{ 1 + zm \left[\alpha_M + \frac{1}{N_f} \log(zm) \right] \right\}$$
$$F_m = F \left\{ 1 + zm \left[\alpha_F - \frac{N_f}{2} \log(zm) \right] \right\}$$
$$\bar{\psi}\psi\rangle_m = \langle \bar{\psi}\psi \rangle \left\{ 1 + zm \left[\alpha_C - \frac{N_f^2 - 1}{N_f} \log(zm) \right] \right\}$$

 \bigstar Log coefficients of $F_m, \langle \bar{\psi}\psi\rangle_m \sim N_f$

★ $\alpha_C \sim 1/a^2 \rightarrow \langle \bar{\psi}\psi \rangle_m$ difficult to measure ★ Instead measure (GMOR) $\frac{M_m^2}{2mF_m} \rightarrow \frac{\langle \bar{\psi}\psi \rangle}{F^3}$ at $m \rightarrow 0$

6-flavor over 2-flavor enhancement



xPT fits and bound





10 and 6 flavor over 2-flavor enhancement



Scale matching 2f, 6f, 10f



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Salient features: Topology

At small lattice spacing the barriers between TC sectors are large

At small m DWF HMC encounter barriers in changing global topology Q

At large volume Q is irrelevant but for us it is a finite size effect

For $0.01 \le m$, Q evolves sufficiently: for m = 0.005 it does not

300 million core hours on LLNL BG/L



Near term LSD plans

- * 10-flavors SU(3)_c fundamental at the same lattice spacing
- ***** Measure enhancement at 2, 6 and 10 flavors
- ***** Measure particle spectrum at 2, 6 and 10 flavors
- ***** Measure S at 2, 6 and 10 flavors
- ***** Measure Dirac eigenvalues at 2, 6 and 10 flavors
- ✤ Fundamental SU(2)_c

Conclusions

- **\mathbb{H}** 2 and 6 flavors SU(3)_c fundamental at same lattice spacing
- # 6 flavors condensate enhancement larger than 50% Excluded no-enhancement at 73% confidence level
- **# 10f very preliminary data indicate larger enhancement.**

Backup slides

Topology (preliminary)

flavors, beta = 2.10 topological charge



Topology (preliminary)

6 flavors, beta = 2.10 topological charge



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Topology (preliminary)

2 flavors, beta = 2.70 topological charge



Domain Wall Fermions (DWF)



- Restoration of Lorentz symmetry : lim(a=0)
- Restoration of chiral symmetry : lim(L_s= infinity)
- The two limits are decoupled !!!
- Computing is only linear in L_s !!!
- Other fermions need $2^6 2^{10}$ more computing for a <= a/2.

Small worlds inside small worlds Technicolor, take 2



The virtues of TC/ETC

- Dynamical explanation of EWS breaking
- ✓ Asymptotically free:
 - no unnatural fine tuning needed
 - no hierarchy problem (breaking scale naturally much smaller than cutoff) it is not trivial
- ✓ ETC provides insights to flavor physics

The problems of TC/ETC

- o Flavor changing neutral currents (ETC)
- o Precision electroweak measurements (TC)
- o Large top quark mass

Precision EW constraints



The **S** parameter of Peskin & Takeuchi assumes a scaled version of QCD with N_f and N_c

27

all: M., = 117 GeV

all: M_u = 340 GeV

all: M., = 1000 GeV

1.00

1.25

0.50 0.75

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Flavor changing neutral currents

- > Fermion masses need new interactions at scale $M >> \Lambda_{TC}$
- > At scales well below $M \to \bar{\Psi} \Psi \bar{T} T \to m_f \sim \frac{<\bar{T}T>}{M^2}$
- > But also have $\bar{\Psi}\Psi\bar{\Psi}\Psi$
- > Flavor changing neutral currents: no known suppression mechanism
- Must keep the scale M very high ~ 1,000 TeV
- But then the quark and lepton masses become too small

Not so fast

Scaling QCD with the number of flavors and colors is not correct.

QCD with many light flavors is a very different theory than QCD with 2 light flavors

$$L\frac{\partial}{\partial L}g(L) = \beta(g) \stackrel{g \to 0}{\sim} b_0 g^3 + b_1 g^5 + b_2 g^7 + \cdots$$
$$b_0 = -\frac{1}{(4\pi)^2} \left(\frac{11}{3}N_c - \frac{2}{3}N_f\right), \qquad b_1 = -\frac{1}{(4\pi)^4} \left[\frac{34}{3}N_c^2 - \left(\frac{13}{3}N_c - \frac{1}{N_c}\right)N_f\right].$$



Walking

Low end of conformal window

$$8 \le N_{fc} \le 12$$

T. Appelquist, G. Fleming, E. Neil, Phys. Rev. Lett. 100, 171607 (2008), hep-lat/0901.3766

Possible effects of walking and the Lattice

$$\label{eq:main_alpha} \Box \quad \frac{m_{a1}^2 - m_{\rho}^2}{m_{a1}^2} \quad \sim {\rm S} \quad {\rm may} \ {\rm be \ smaller}$$

Coupling stays strong at larger scales: could enhance condensate relative to Λ_{TC}

□ Need a true first principles calculation => Need the Lattice.

LHC TeV physics

