

Hadron Form Factors at Large Transfer Momentum

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(with Huey-Wen Lin)

[arXiv:1005.0799](https://arxiv.org/abs/1005.0799)

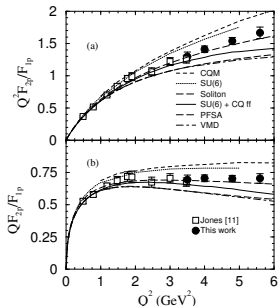
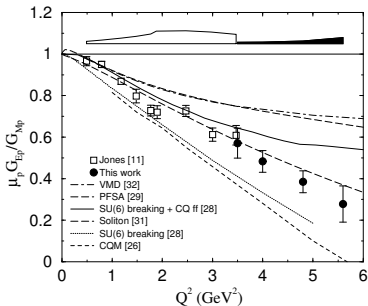
XXVIII International Conference on Lattice Field Theory

Outline

- 1 Experimental Motivation
- 2 Methodology
- 3 Nucleon Form Factors
- 4 Transverse Distributions

Form Factors at High Momentum Transfer

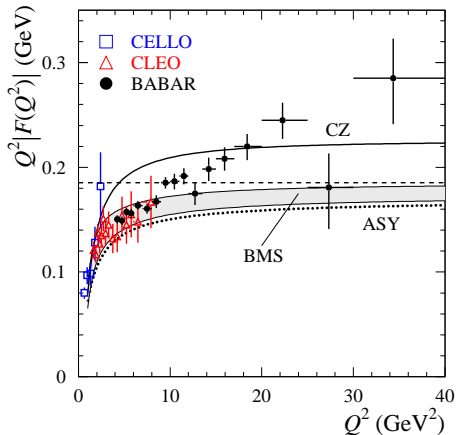
- Experimental results for neutrons available only up to $\approx 4 \text{ GeV}^2$
- JLab 12-GeV upgrade should shed some light on things
- Does G_E^p cross zero around 7 GeV^2 ?



Isn't Perturbative QCD Good Enough?

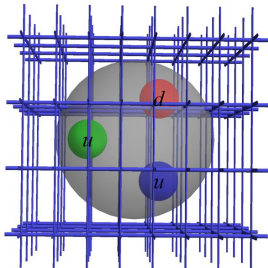
Recent BaBar results for $\gamma^*\gamma \rightarrow \pi^0$

- Disagreement with PQCD over a wide range:
 $4 \text{ GeV}^2 < Q^2 < 40 \text{ GeV}^2$
- An opportunity for lattice QCD

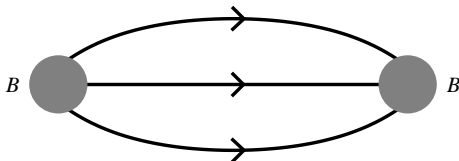


Anisotropic 2+1-Flavor Clover Ensembles

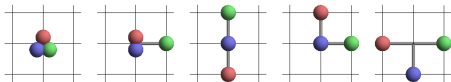
- Anisotropy improves resolution of excited states
- Clover fermion action
- Quenched Ensembles
 - $16^3 \times 64, \xi = 3$
 - $a_t^{-1} \approx 6.0 \text{ GeV}$ ($a_s \approx 0.10 \text{ fm}$)
 - Wilson gauge action
 - Valence $M_\pi \in \{480, 720, 1080\} \text{ MeV}$
- 2+1-Flavor Ensembles
 - $16^3 \times 128, \xi = 3.5$
 - $a_t^{-1} \approx 5.6 \text{ GeV}$ ($a_s \approx 0.12 \text{ fm}$)
 - RG-improved gauge action
 - $M_\pi \in \{450, 580, 875\} \text{ MeV}$



Baryon Correlators



- $B(x) = \epsilon^{abc} [q_1^{aT}(x) C \gamma_5 q_2^b(x)] q_1^c(x)$
- Use various Gaussian smearings at source and sink
 - Quenched: $\sigma \in \{0.5, 2.5, 4.5\}$
 - 2+1-flavor: $\sigma \in \{0.5, 1.5, 2.5, 3.5, 4.5\}$
- More sophisticated operators may be used



Variational Method

Given a two-point correlator

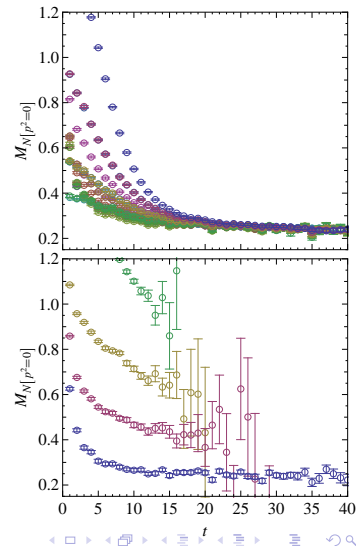
$$C_{ij}(t) = \langle O_j(t) O_i(0) \rangle = \sum_{n=0}^{\infty} Z_{in}^\dagger Z_{jn} e^{-E_n t},$$

take the generalized eigenvalues with respect to intermediate time t_0

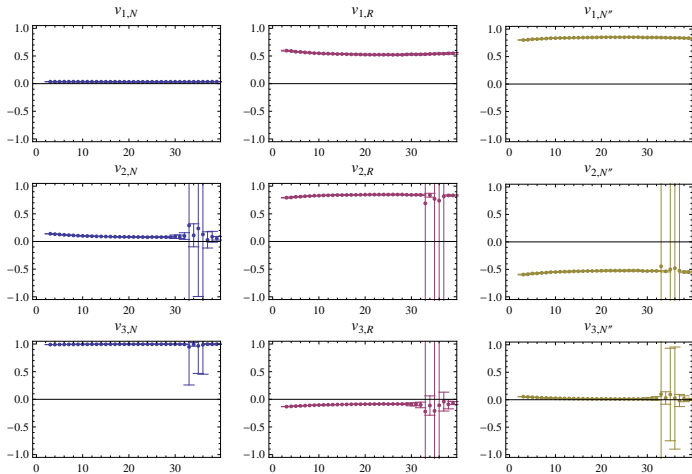
$$C(t_0)^{-1/2} C(t) C(t_0)^{-1/2} \psi = \lambda(t, t_0) \psi.$$

For sufficiently large $t - t_0$,

$$\lambda_n(t, t_0) = e^{-(t-t_0)E_n}.$$



Variational Overlaps

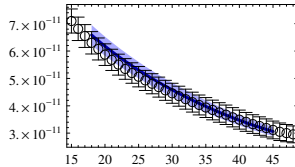
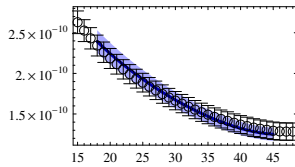
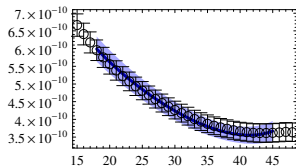
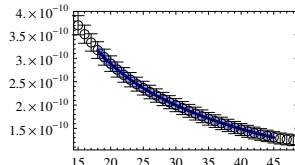
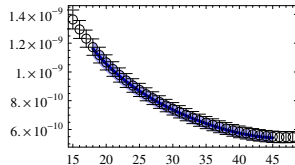
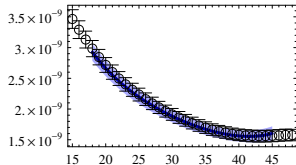
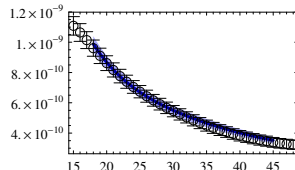
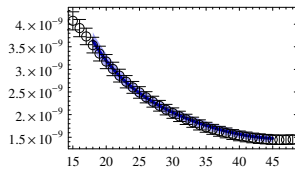
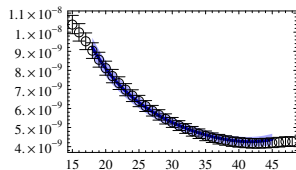


Extract from Three-Point Functions

$$\Gamma_{\mu,AB}^{(3),T}(t_j, t, t_f, \vec{p}_i, \vec{p}_f) = Z_V \sum_n \sum_{n'} f_{n,n'}(p_f, p_i; E_f, E_i; t_j, t, t_f) \\ \times \sum_{s,s'} T_{\alpha\beta} u_{n'}(\vec{p}_f, s')_{\beta} \langle N_{n'}(\vec{p}_f, s') | V_{\mu} | N_n(\vec{p}_i, s) \rangle \bar{u}_n(\vec{p}_i, s)_{\alpha}$$

$$f = \frac{Z_{Bn'}^* Z_{An}}{4E_f E_i} \sqrt{\frac{2E_f}{E_f + M_i}} e^{-E_f(t_f-t)} e^{-E_i(t-t_j)}$$

Simultaneous Fitting



Dirac and Pauli Form Factors

- General Form:

$$\langle N_2 | V_\mu | N_1 \rangle(q) = \bar{u}_{N_2}(p') \left[F_1(q^2) \left(\gamma_\mu - \frac{q_\mu}{q^2} \not{q} \right) + \sigma_{\mu\nu} q_\nu \frac{F_2(q^2)}{M_{N_1} + M_{N_2}} \right] u_{N_1}(p) e^{-iq \cdot x}$$

- Ground-State Only:

$$\langle N | V_\mu | N \rangle(q) = \bar{u}_N(p') \left[F_1(q^2) \gamma_\mu + \sigma_{\mu\nu} q_\nu \frac{F_2(q^2)}{2M_N} \right] u_N(p) e^{-iq \cdot x}$$

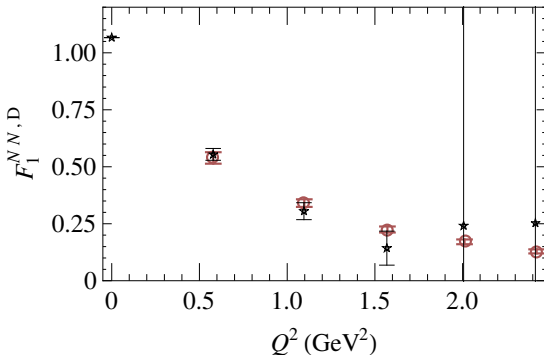
- Sachs Form Factors

$$G_E(Q^2) = F_1(Q^2) - \frac{Q^2}{4M_N^2} F_2(Q^2) \quad (1)$$

$$G_M(Q^2) = F_1(Q^2) + F_2(Q^2) \quad (2)$$

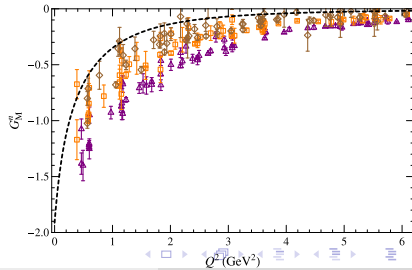
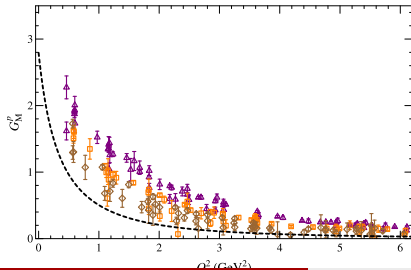
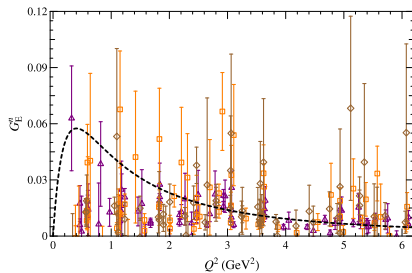
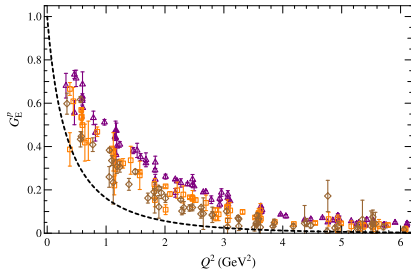
Comparison with Ratio Method

$$\frac{C_{SS}^{(3)}(t_i, t, t_f; p_i, p_f)}{C_{SS}^{(2)}(t_i, t_f; p_f)} \left(\frac{C_{ps}^{(2)}(t, t_f, p_i) C_{SS}^{(2)}(t_i, t, p_f) C_{ps}^{(2)}(t_i, t_f, p_f)}{C_{ps}^{(2)}(t, t_f, p_f) C_{SS}^{(2)}(t_i, t, p_i) C_{ps}^{(2)}(t_i, t_f, p_i)} \right)^{\frac{1}{2}}$$



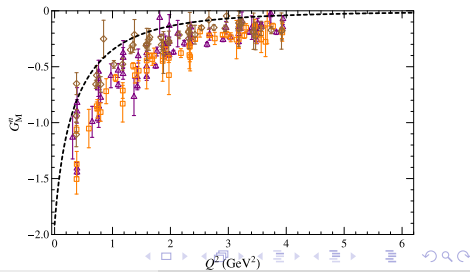
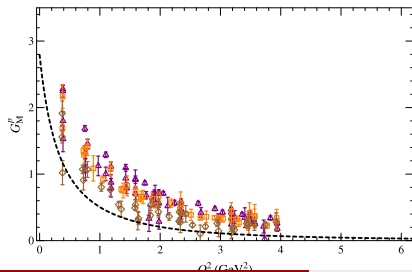
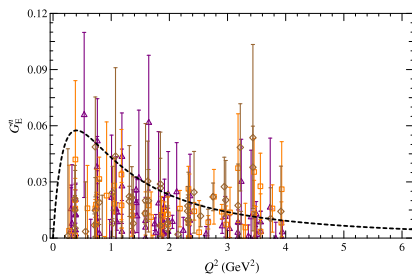
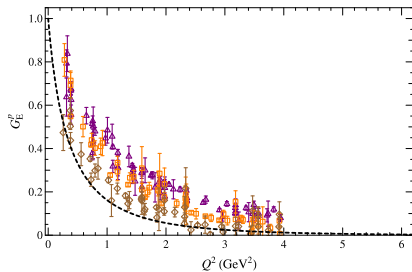
Sachs Form Factors

Quenched Results

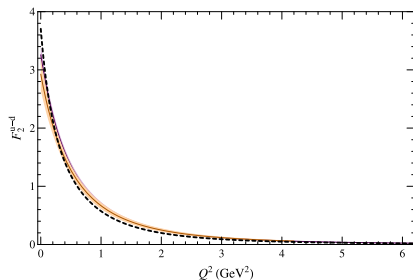
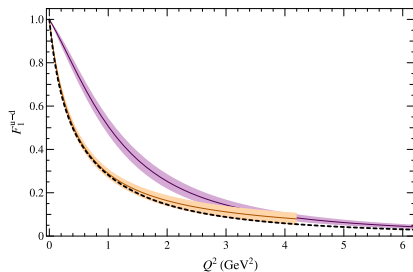


Sachs Form Factors

Dynamical Results



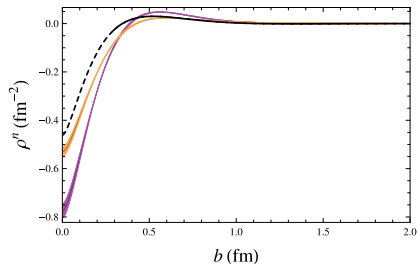
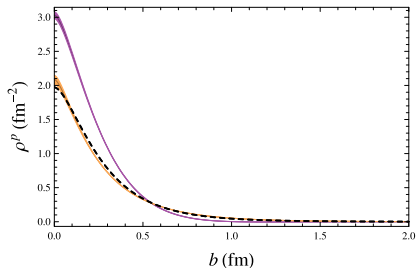
Isovector Form Factors



Transverse Charge Distribution

To avoid relativistic distortions, we calculate densities in a plane transverse to a direction of infinite boost:

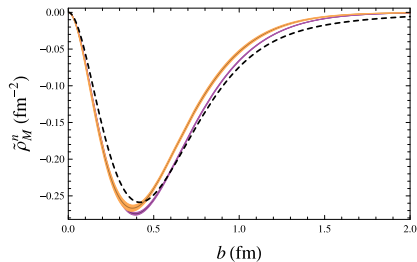
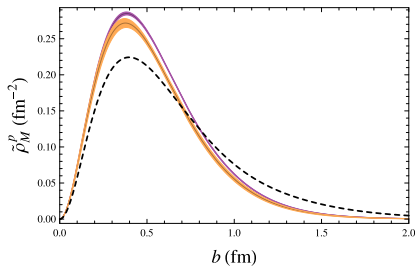
$$\rho(b) = \int_0^\infty \frac{Q dQ}{2\pi} J_0(bQ) F_1(Q^2)$$



Transverse Magnetization Distribution

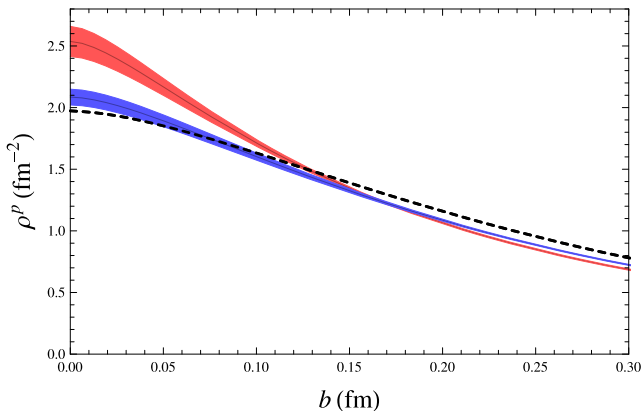
The transverse magnetization is defined with respect to a B -field in the \hat{x} direction:

$$\tilde{\rho}_M(b) = b \sin^2 \phi \int_0^\infty \frac{Q^2 dQ}{2\pi} J_1(bQ) F_2(Q^2)$$



How High Q^2 Is High Enough?

Consider what we would get for the density in the nucleon core if we only had data up to a typical Q^2 of 2 GeV^2



Summary

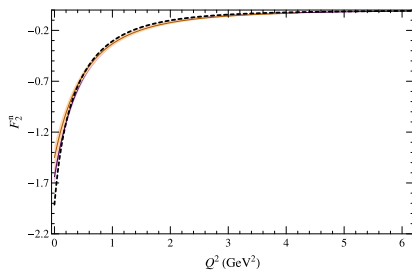
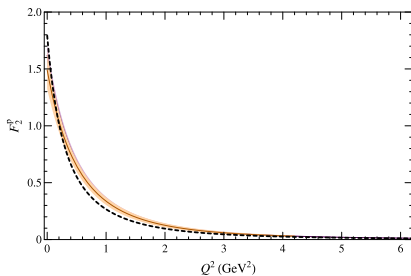
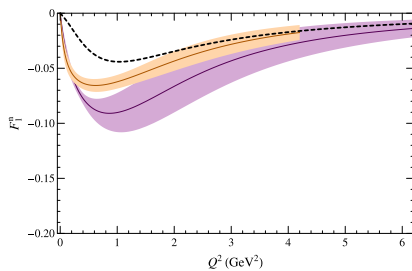
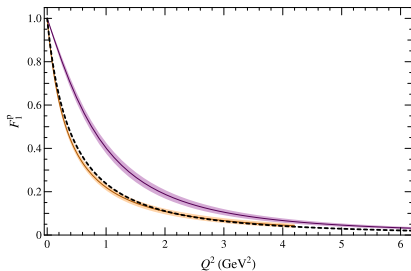
- Conclusions

- Method allows extension to higher momentum transfer
- Will extend reach to match future experiments
- Essential for accurately determining densities in the hadronic core

- Future Work

- Introduce more sophisticated operators
- Nucleon axial and pion form factors (see H-W Lin's poster)

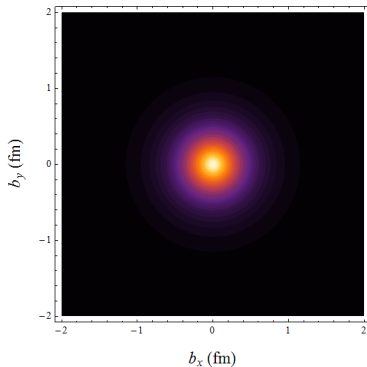
Quenched-Dynamical Fit Comparisons



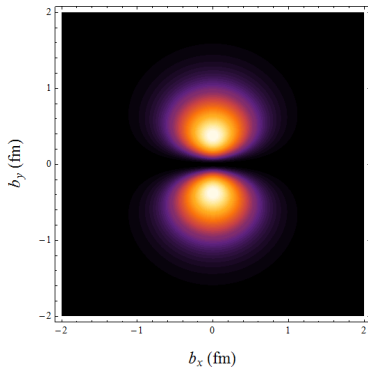
Two-Dimensional Densities

Proton

$$F_1(\vec{b})$$



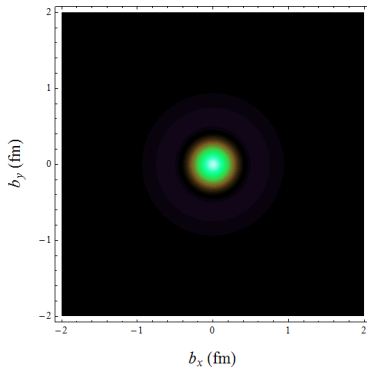
$$F_2(\vec{b})$$



Two-Dimensional Densities

Neutron

$$F_1(\vec{b})$$



$$F_2(\vec{b})$$

