

On the universal $O(N)$ scaling behavior of (2+1)-flavor QCD

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for RBC-Bielefeld-GSI Collaboration:

S. Ejiri, F. Karsch, E. Laermann, C. Miao, S. Mukherjee,
P. Petreczky, C.S., W. Söldner, W. Unger

based on:

S. Ejiri *et al.*, PRD 80 (2009) 094505.

- chiral symmetry of 2-flavor QCD

$$SU_L(2) \times SU_R(2) \simeq O(4)$$

- hence, if m_s is large in (2+1)-flavor QCD:

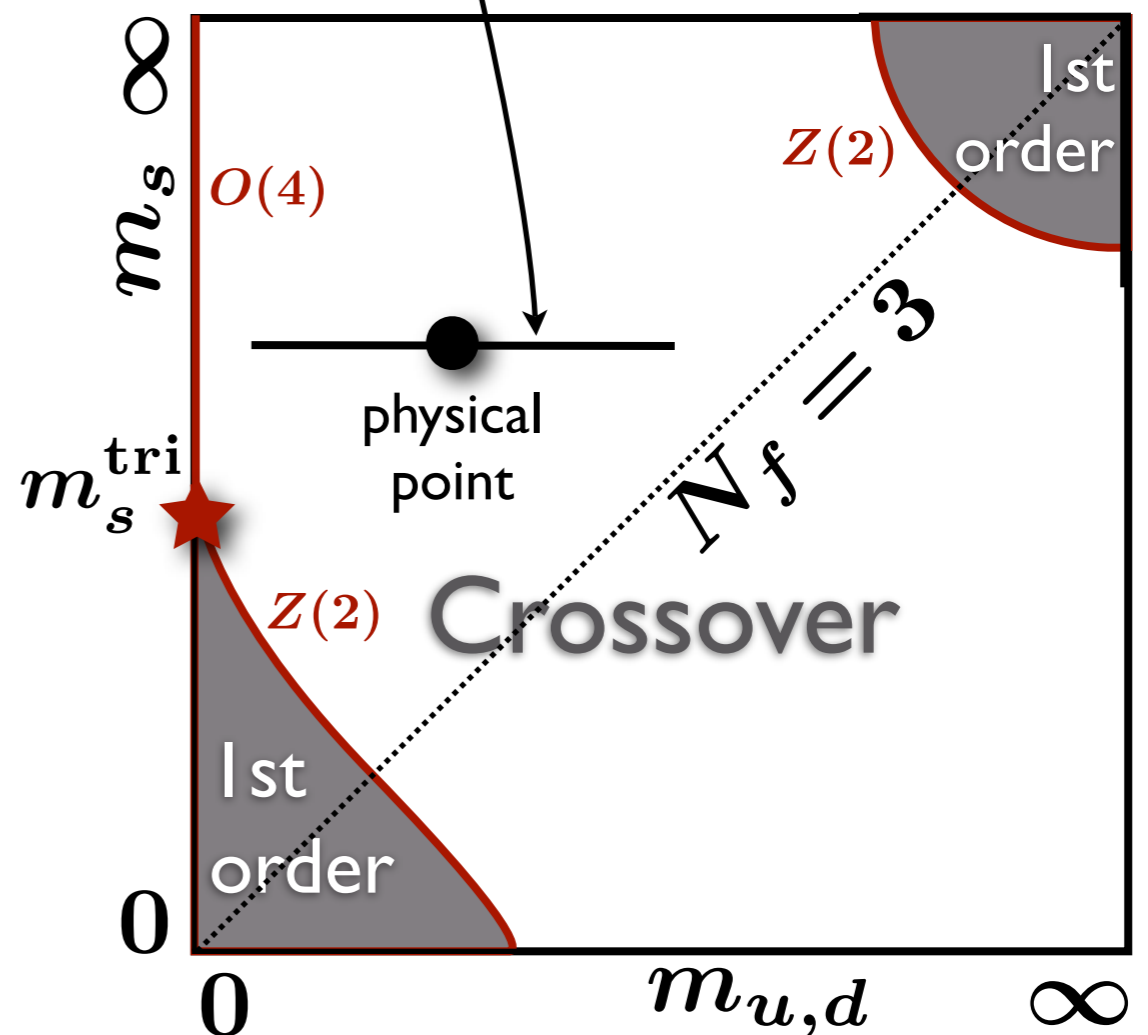
expect universal behavior as of 3d- $O(4)$ spins in the vicinity of T_c and the chiral limit

- so far no clear evidence from simulations
- staggered fermions preserve a flavor non-diagonal $U(1)$ -part of chiral symmetry even at $a > 0$

→ look for $O(2)$ -critical behavior

Simulations with improved staggered fermions (p4fat3)

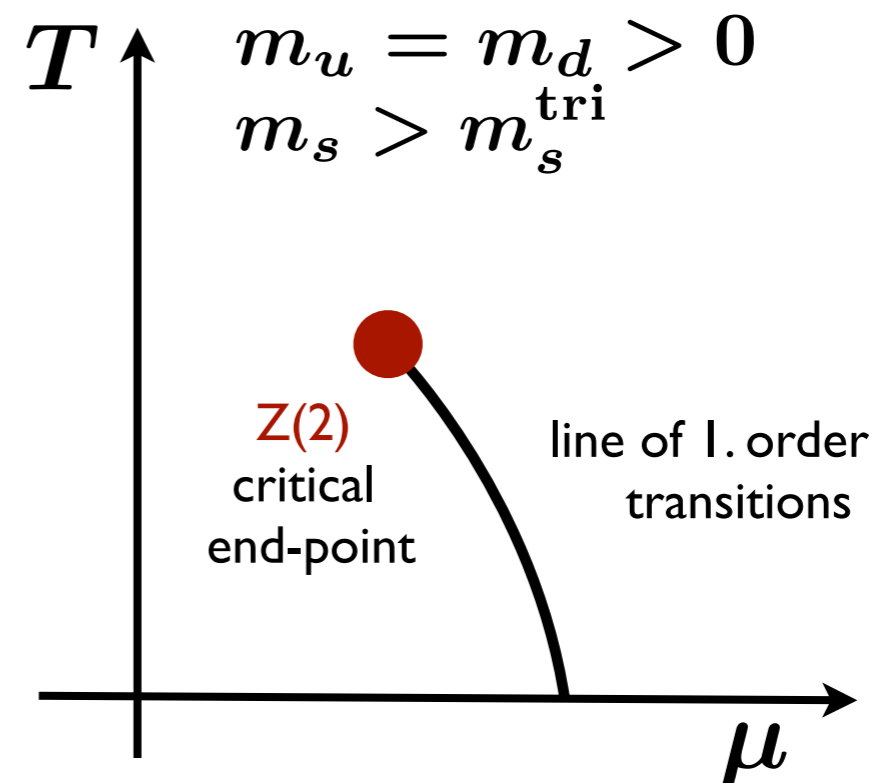
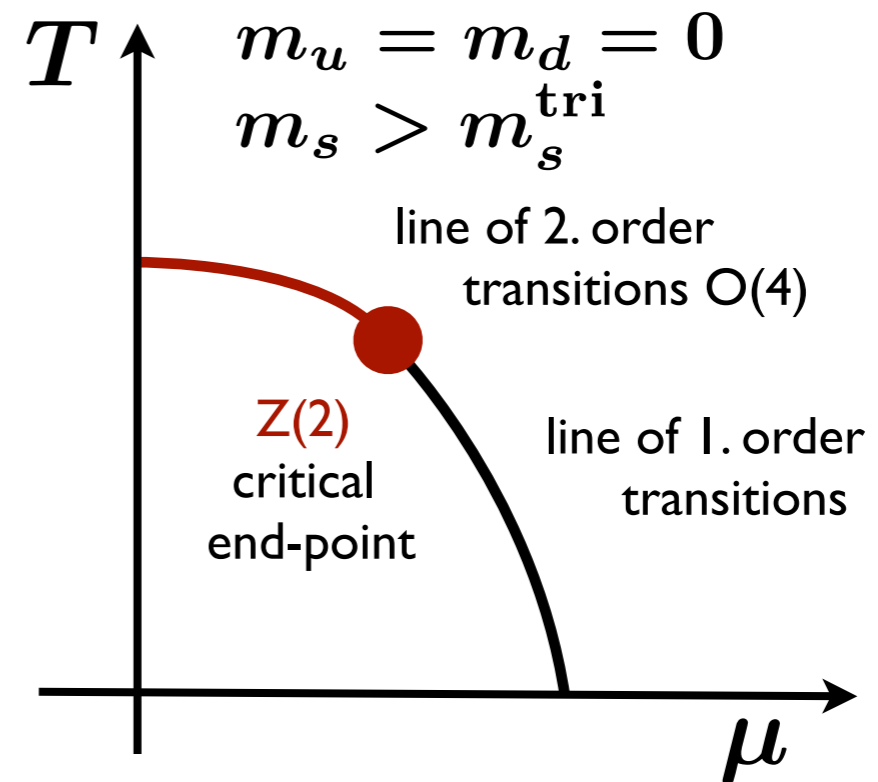
range of simulations ($N_\tau = 4$) $m_q = (2/5 - 1/80)m_s$	m_l/m_s	m_π
	1/80	75 MeV
	1/40	105 MeV
	1/20	150 MeV



- situation at nonzero chemical potential is very unclear
- direct simulations MC simulations are prohibited by the **sign-problem**

→ use Taylor expansion approach

expected (T, μ) -phase diagrams:



- ★ Introduction
- ★ Scaling of chiral condensate (The magnetic EoS)
 - introduce some analogies between spin models and QCD
 - fit p4fat3-data ($N_\tau = 4, 8$) to magnetic scaling function f_G
 - determine important non universal constants of QCD
- ★ Scaling of chiral and mixed susceptibilities
 - fit p4fa3-data to scaling functions f_χ and f'_G
 - more sensitivity to the universality class
 - predictions on the critical line
- ★ Summary

- Thermodynamics in the vicinity of a critical point:

free energy density:

$$-\frac{1}{V} \ln Z = f_s(t, h) + f_r(T, V, H)$$

(singular part) (regular part)

where

$$t = \frac{1}{t_0} \frac{T - T_c}{T_c}$$

(reduced temperature)

$$h = \frac{H}{h_0}$$

(external field)



QCD:

$H \sim m_q$
(quark mass)
our choice:
 $H = m_l/m_s$

assume:

$$f_s(t, h) = b^{-d} f_s(b^{y_t} t, b^{y_h} h)$$

choose: $b = h^{-1/y_h}$



$$f_s(t, h) = h_0 h^{1+1/\delta} f_M(z)$$

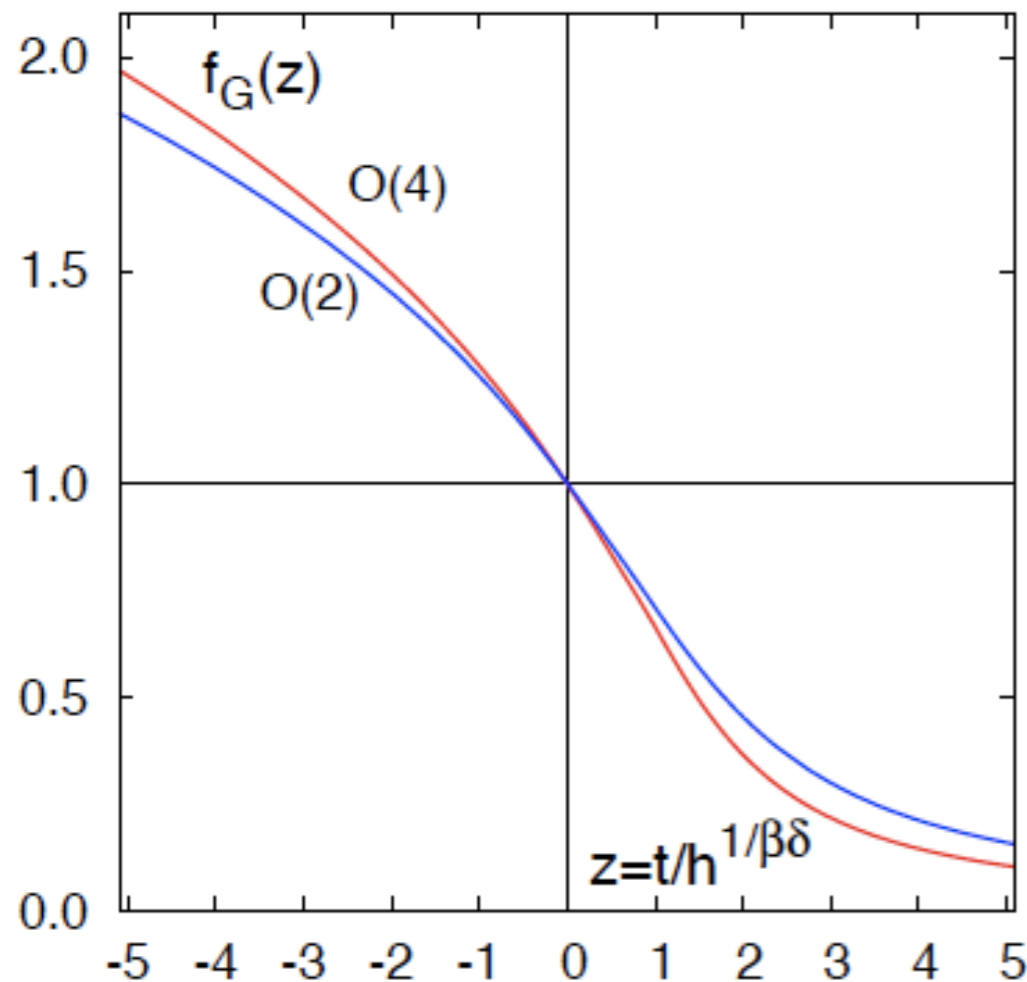
(“magnetic version” of the free energy density)

$$z = t/h^{1/\beta\delta}$$

(scaling variable)

- order parameter (magnetization):

$$M = -\frac{\partial f_s(t, h)}{\partial H} = \frac{1}{h_0} \frac{\partial f_s(t, h)}{\partial h} \equiv h^{1/\delta} f_G(z)$$



universal scaling function

$$f_G(z) = - \left[\left(1 + \frac{1}{\delta} \right) f_M(z) - \frac{z}{\beta\delta} f'_M(z) \right]$$

- scaling variable:

$$z = t/h^{1/\beta\delta}$$

- scaling function and critical exponents are known to high precision in condensed matter literature [e.g. Engels *et al.*]

- scaling function includes **Goldstone effect** in the limit of $z \rightarrow -\infty$

$$z \rightarrow -\infty : \quad h \rightarrow 0, t < 0 \quad M \sim (-t)^\beta + c(t)\sqrt{h}$$

- **lattice action**: improved staggered fermions (p4fat3), (2+1)-flavor
- **algorithm**: exact RHMC
- **strange quark mass**: fixed to physical values, $N_\tau = 4 : am_s = 0.065$
 $N_\tau = 8 : am_s = 0.024$
 $\rightarrow m_{\bar{s}s} \simeq 669 \text{ MeV}$
- **light quark mass**: $m_q/m_s = 1/80 : \rightarrow m_\pi = 75 \text{ MeV}$
 $m_q/m_s = 1/20 : \rightarrow m_\pi = 150 \text{ MeV}$
- **statistics** (measurements separated by 10 trajectories):

lattice dim.	m_q/m_s	statistics	lattice dim.	m_q/m_s	statistics
$32^3 \times 4 \blacktriangle$	1/80	$\mathcal{O}(20000)$			
$32^3 \times 4 \blacktriangle$	1/40	$\mathcal{O}(20000)$			
$16^3 \times 4 \blacktriangle$	1/40	$\mathcal{O}(30000)$	$32^3 \times 8$	1/40	just started
$16^3 \times 4 \blacktriangle$	1/20	$\mathcal{O}(40000)$	$32^3 \times 8 \blacklozenge$	1/20	$\mathcal{O}(20000)$
$16^3 \times 4 \blacktriangle$	1/10	$\mathcal{O}(40000)$	$32^3 \times 8 \star$	1/10	$\mathcal{O}(30000)$
$16^3 \times 4 \blacktriangle$	1/5	$\mathcal{O}(40000)$	$32^3 \times 8 \star$	1/5	$\mathcal{O}(30000)$
$16^3 \times 4 \blacktriangle$	2/5	$\mathcal{O}(40000)$			
$\beta = 3.2800, \dots$		3.3300	$\beta = 3.4800, \dots$		3.5400

\blacktriangle S. Ejiri *et al.* [**RBC-Bielefeld-GSI**], PRD 80 (2009) 094505.

\blacklozenge M. Cheng *et al.* [**RBC-Bielefeld-GSI**], PRD 81 (2010) 054504.

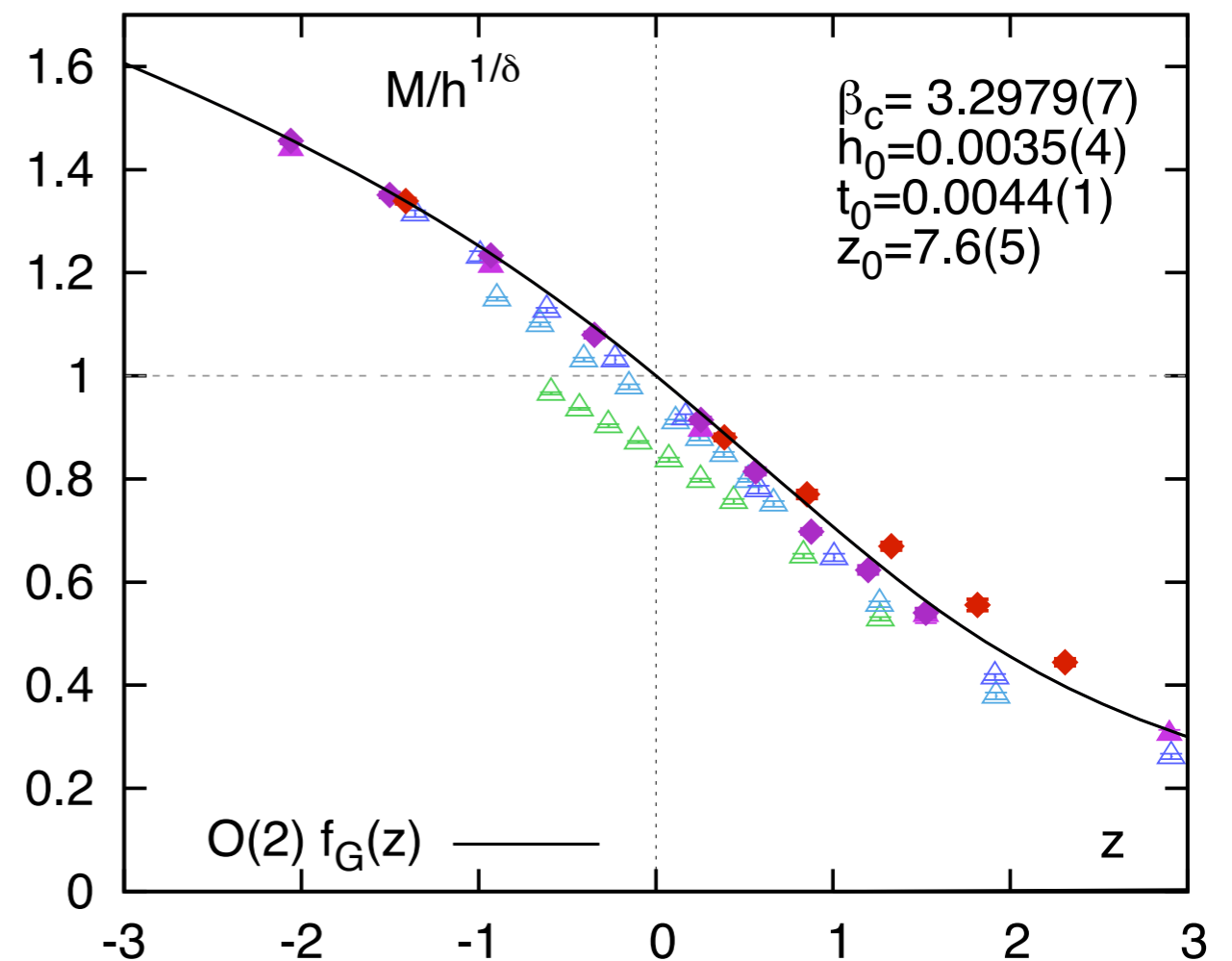
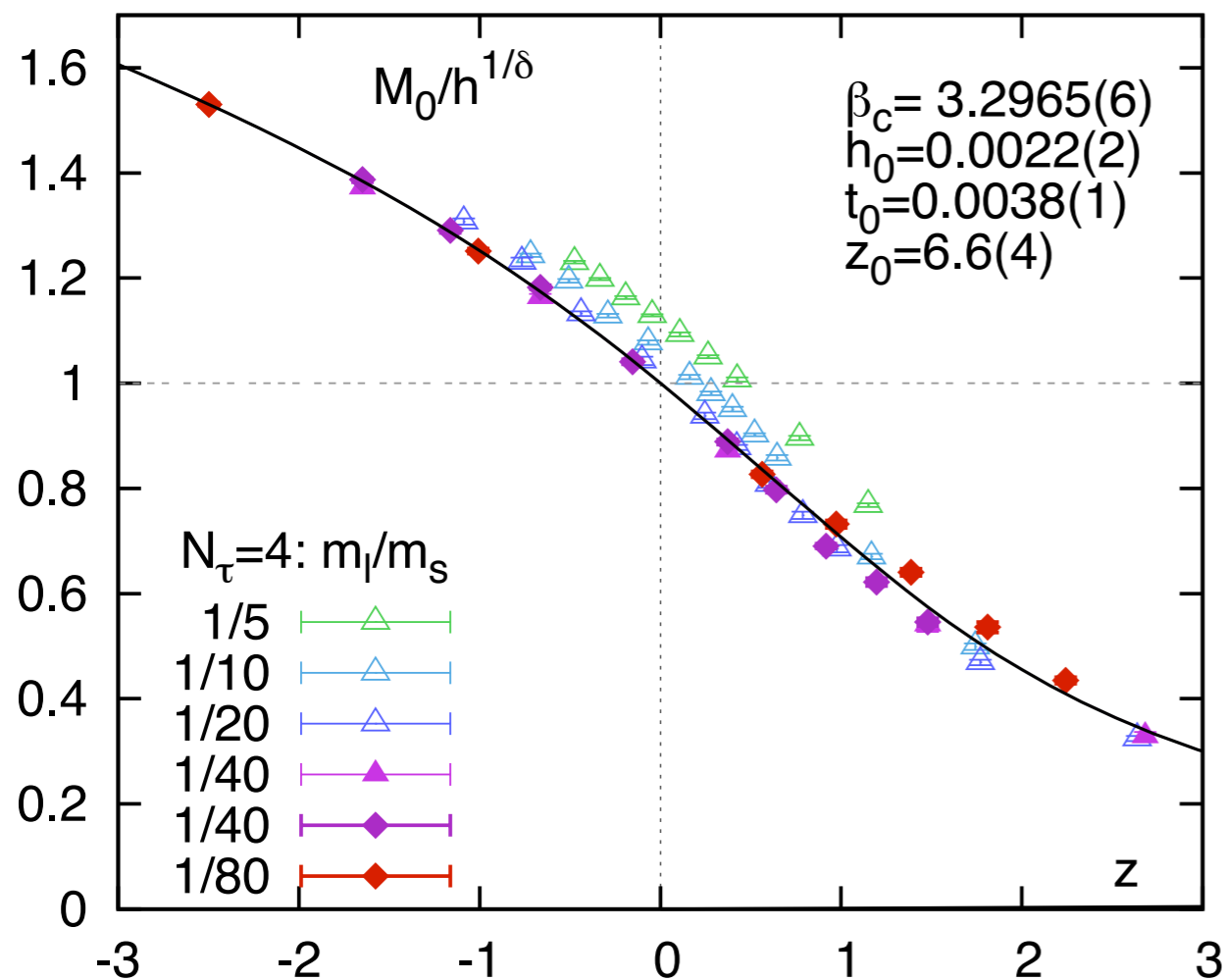
\star A. Bazavov *et al.* [**HotQCD**], PRD 80 (2009) 014504.

- **two order parameter:**

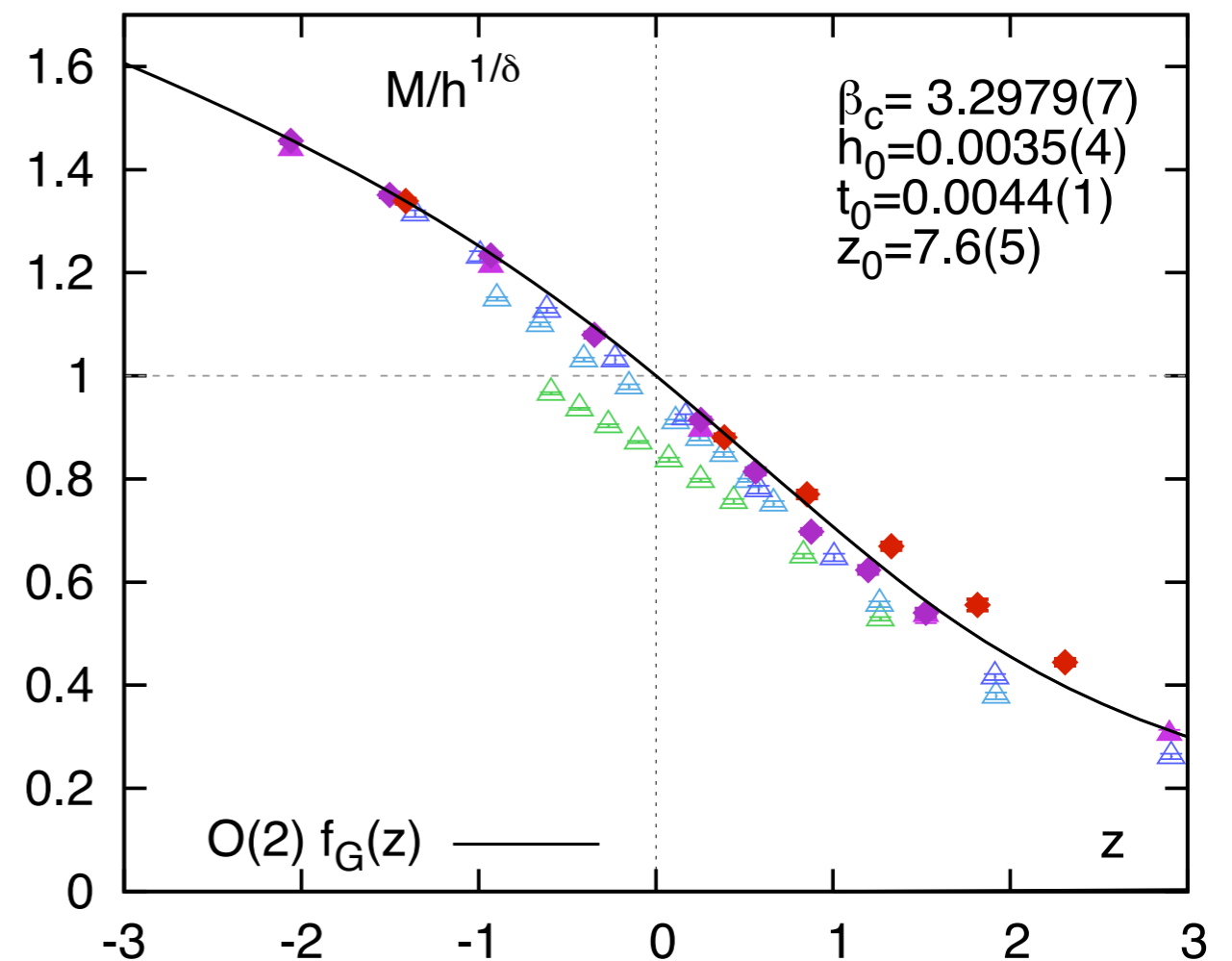
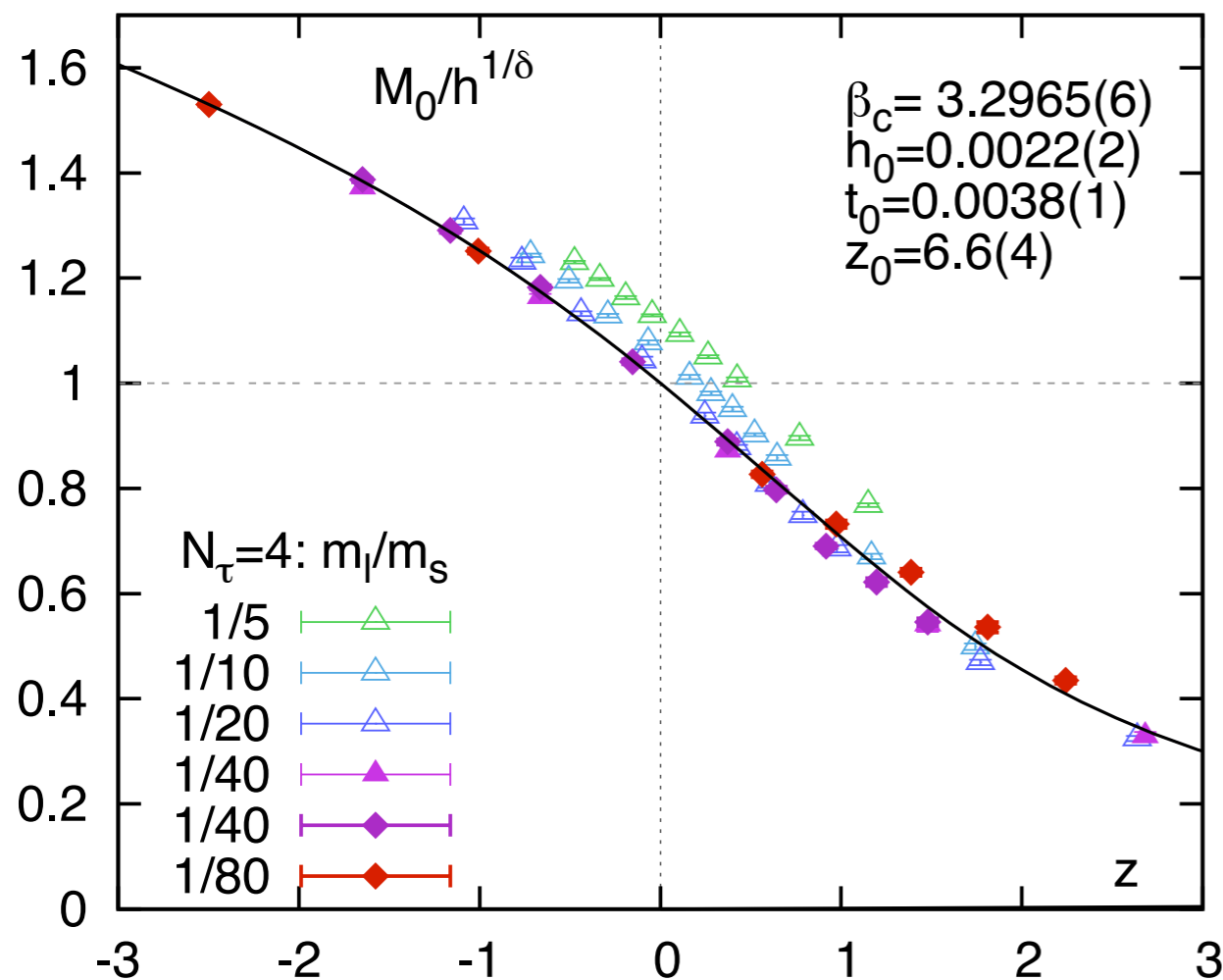
$$\left. \begin{aligned} M_0 &= m_s \langle \bar{\psi}\psi \rangle_l / T^4 \\ M &= m_s \left(\langle \bar{\psi}\psi \rangle_l - \frac{m_l}{m_s} \langle \bar{\psi}\psi \rangle_s \right) / T^4 \end{aligned} \right\} = h^{1/\delta} f_G(z)$$

(subtracted condensate to remove UV-div. $\sim m_l/a^2$)

- **three fit parameter:** critical temperature T_c (critical coupling β_c), normalization constants t_0, h_0



- $O(2)$ slightly preferred, however, re-parametrization $z \rightarrow 1.2z$ moves $O(2)$ onto $O(4)$
 \rightarrow scaling functions almost indistinguishable, we can not discriminate between $O(2)$ and $O(4)$
- $z_0 = t_0/h_0^{1/\beta\delta}$ is independent under re-scaling (t_0, h_0 not)
- $z_0(m_s, a^2)$ might be a QCD invariant, which only depend on strange quark mass and lattice artifacts

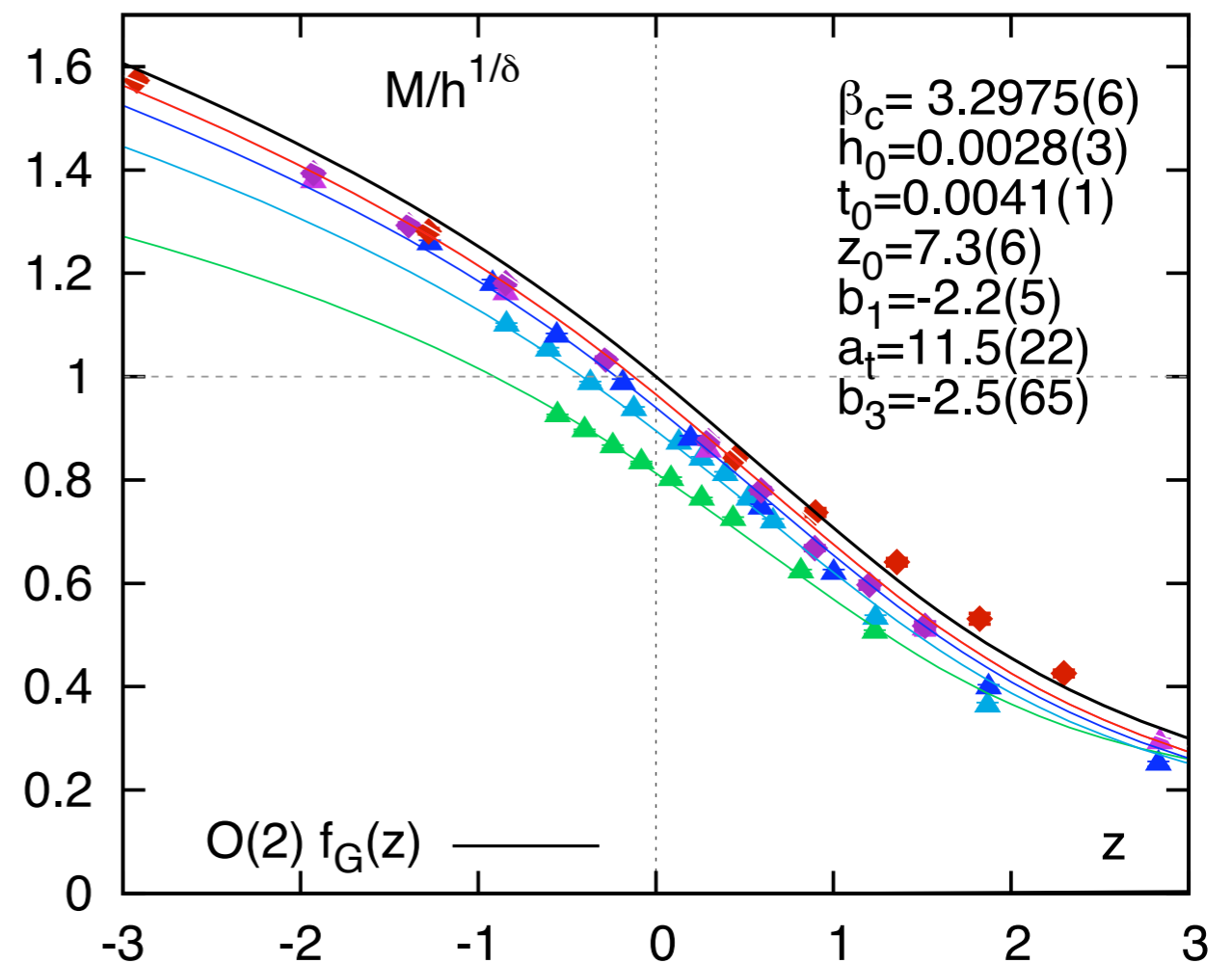
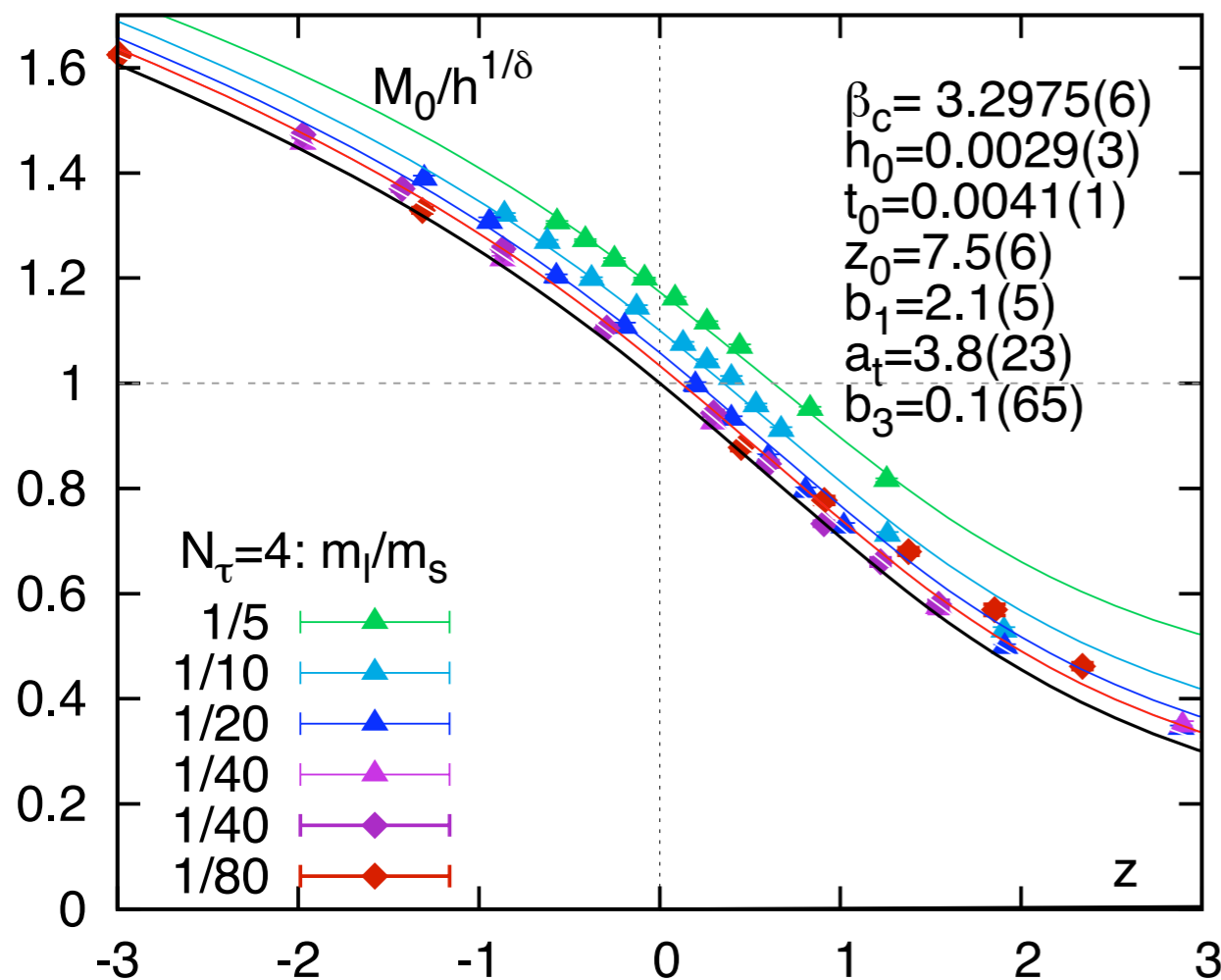


Deviations from scaling (Nt=4)

- mass range $m_l/m_s < 1/20$ is well described by scaling function
- deviations from scaling substantial for $m_l/m_s > 1/20$
- include regular part into the fit:

$$M = h^{1/\delta} f_G(z) + a_t(T - T_c)H + b_1H + b_3H^3$$

→ results for β_c, t_0, h_0 are recovered within errors



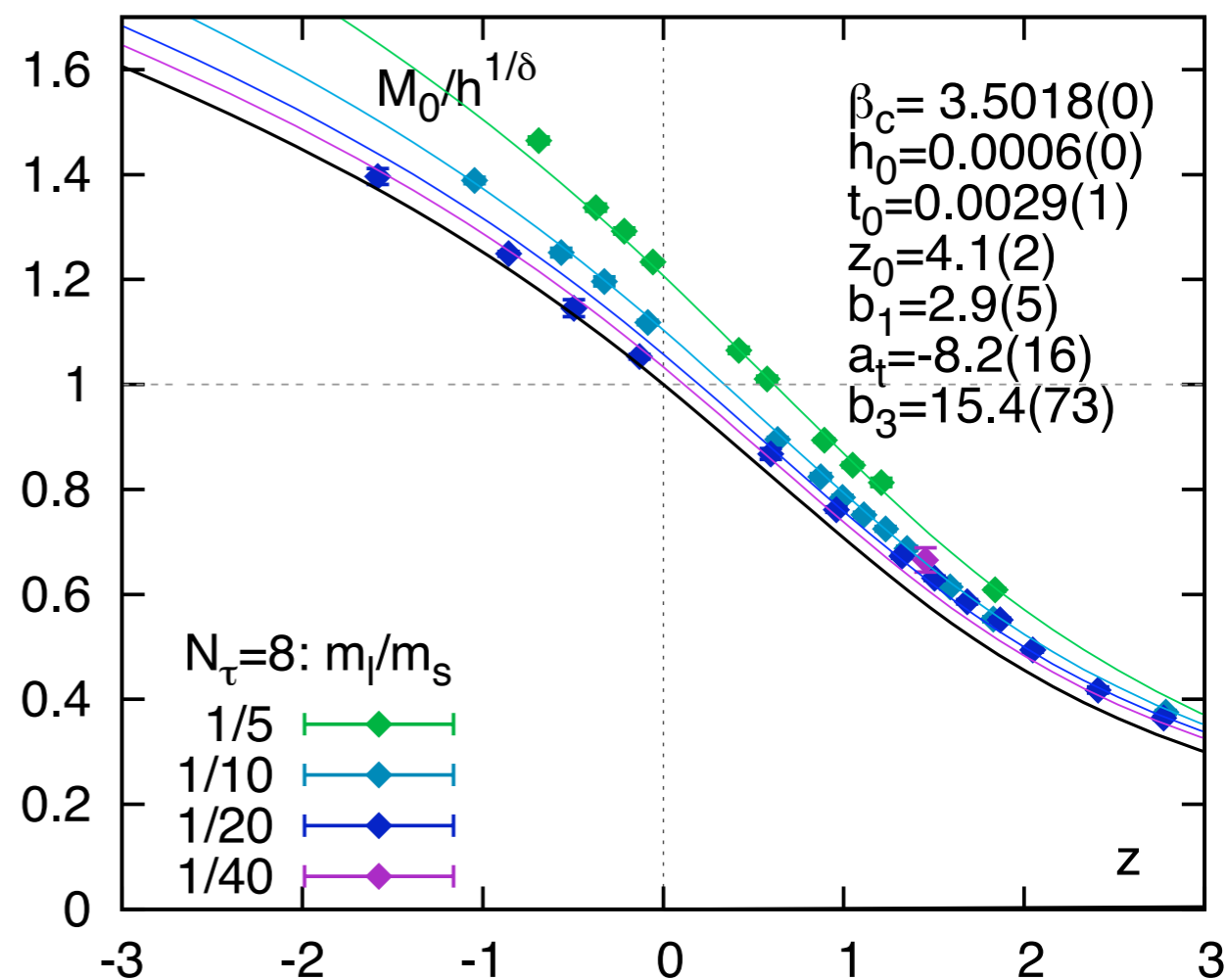
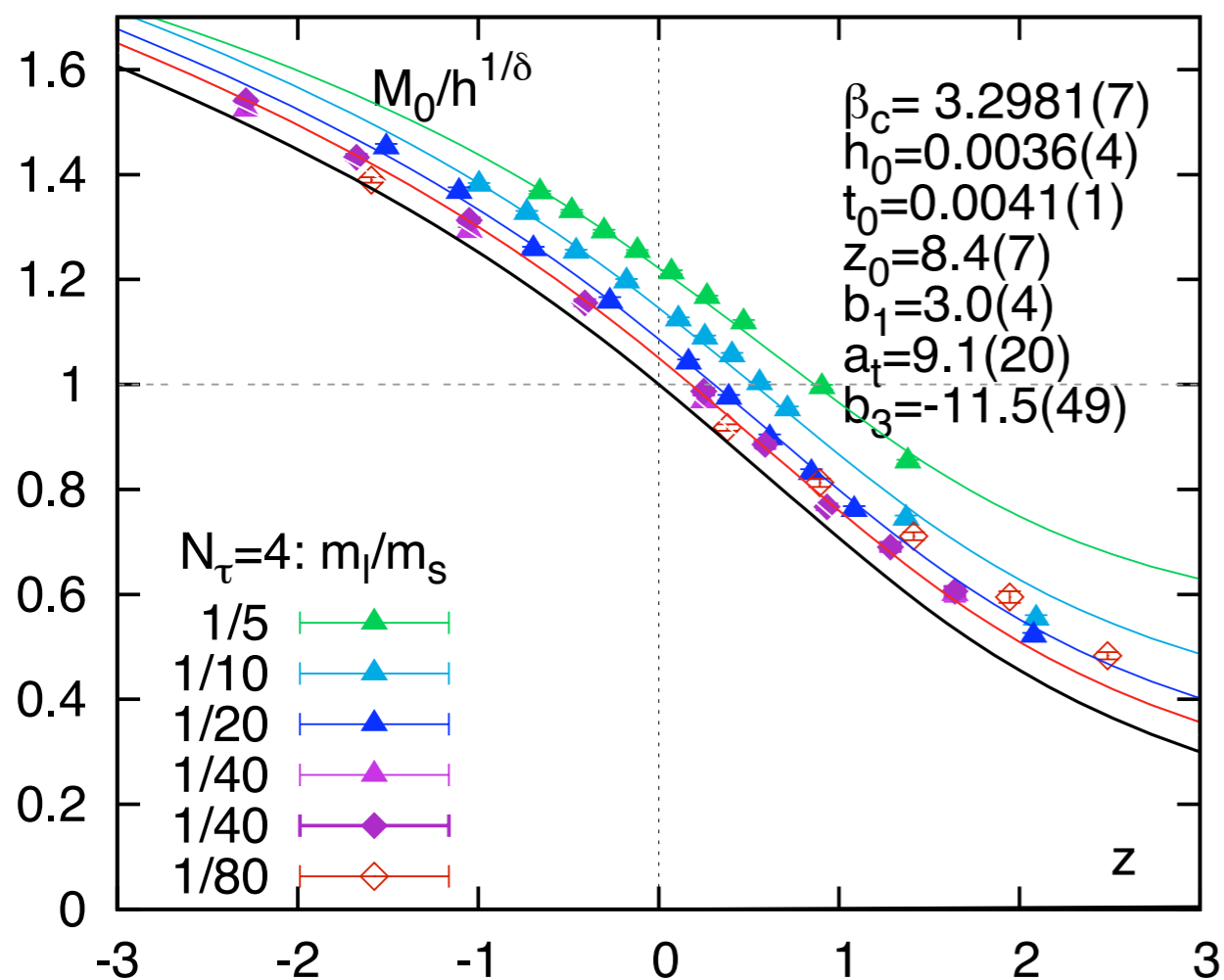
Magnetic EoS in QCD ($N_t=8$)

- $N_\tau = 8$: fit w/o scaling violations not possible yet
- fit for $\beta_c, t_h, h_0, a_t, h_1, h_3$ (range $m_l/m_s \geq 1/40$) works reasonably well \rightarrow **assume z_0 to be stable/reliable**

• cutoff dependence:

N_τ	4	8
z_0	7.5(9)	4.3(5)

\rightarrow **further studies are needed to control continuum limit**

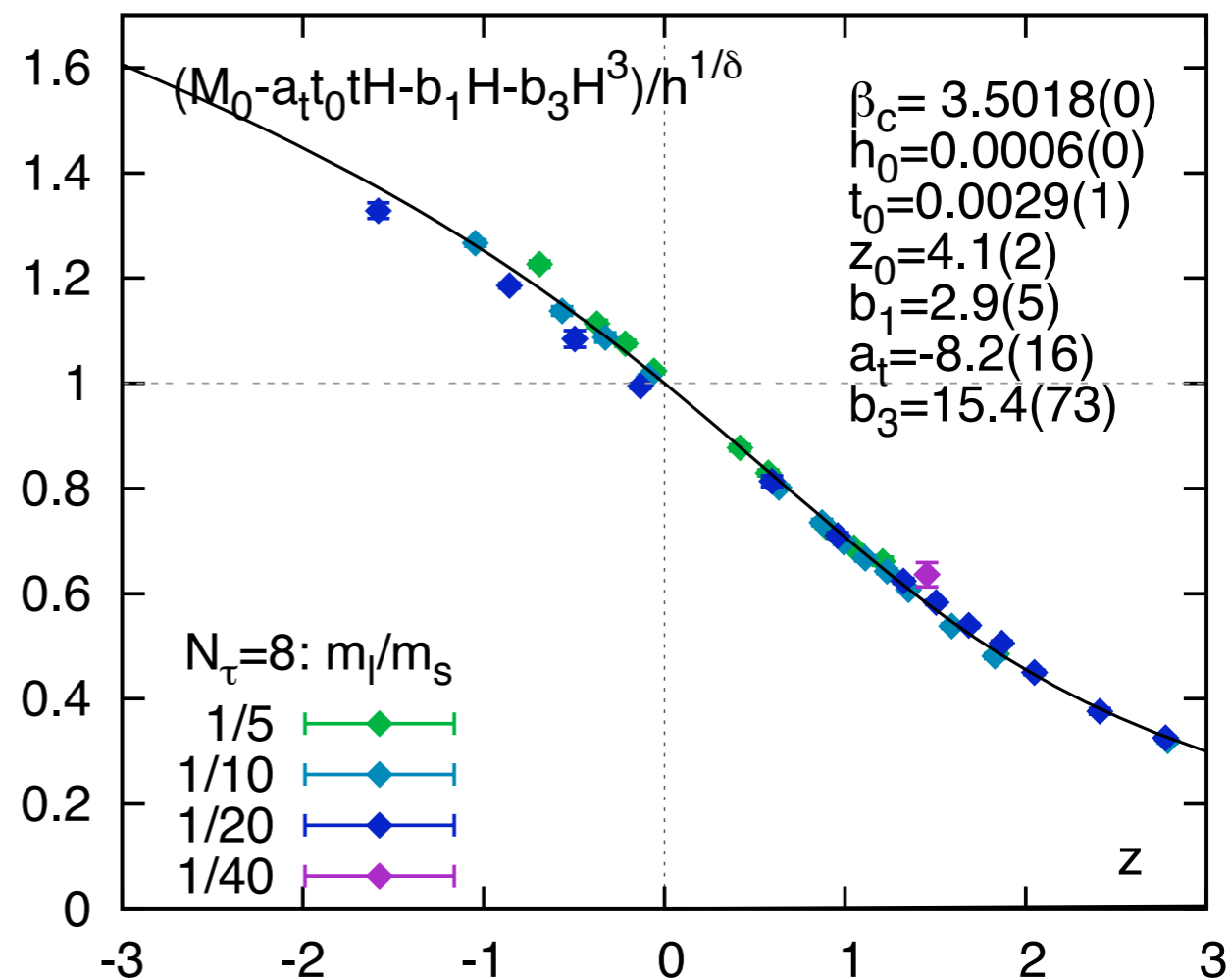
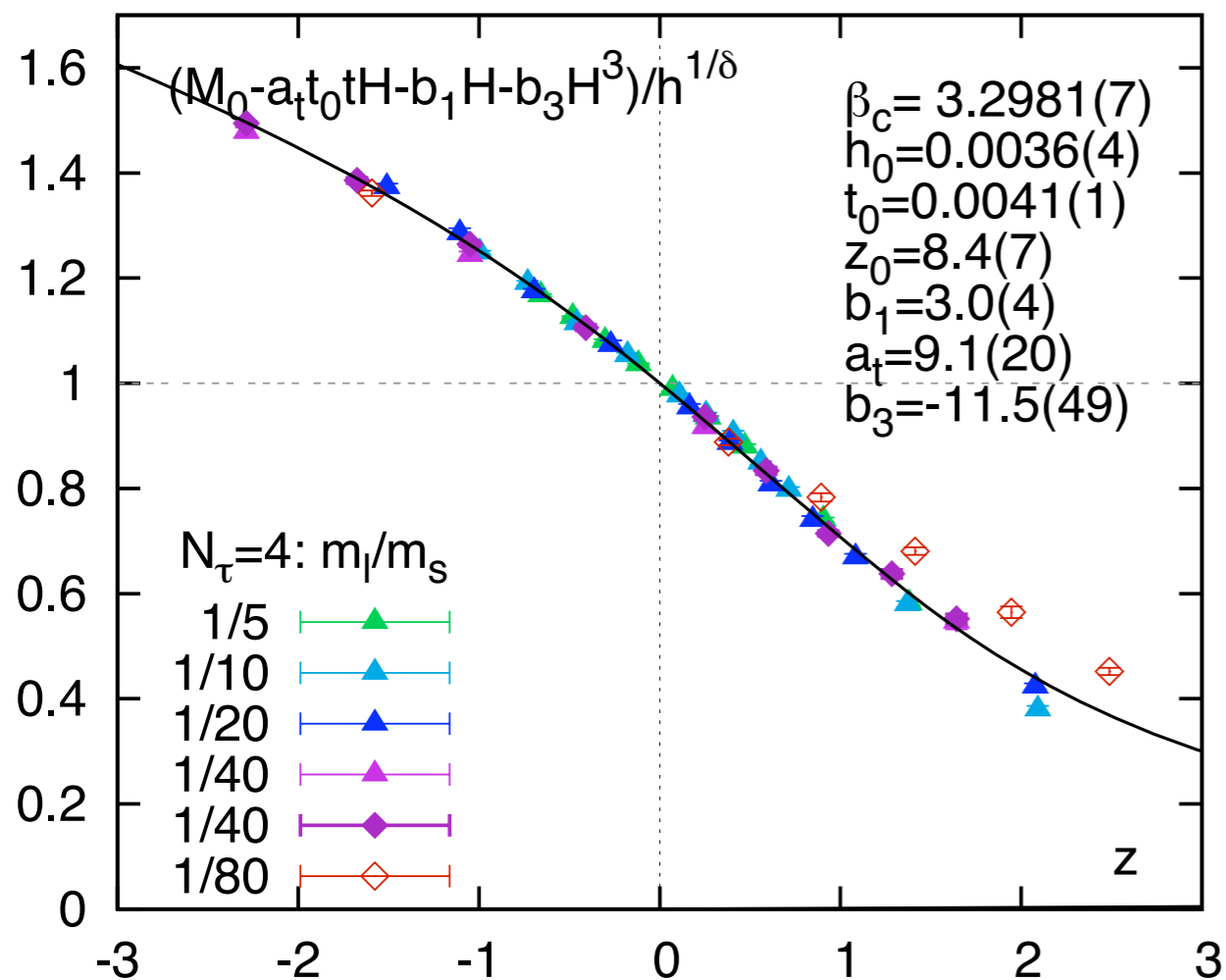


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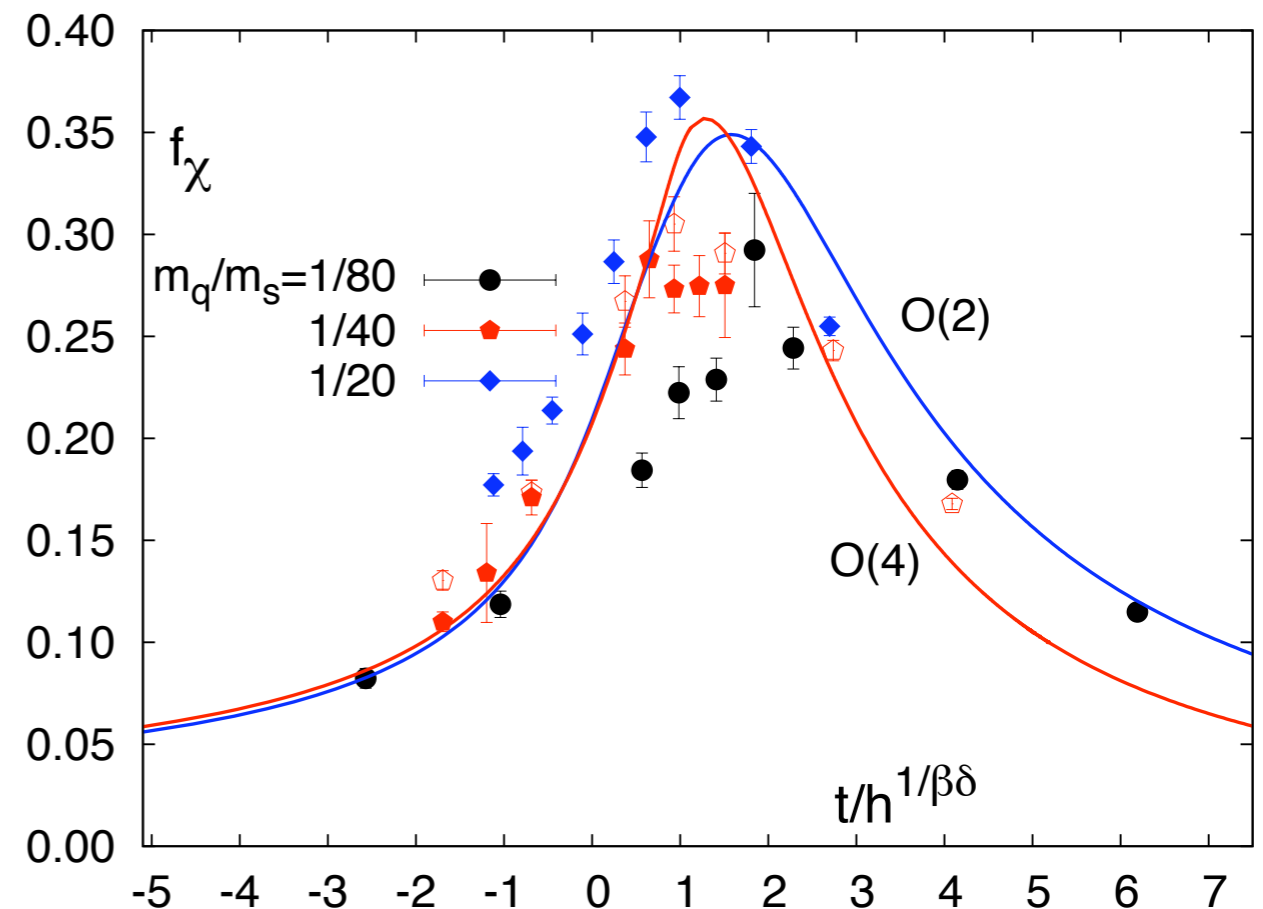
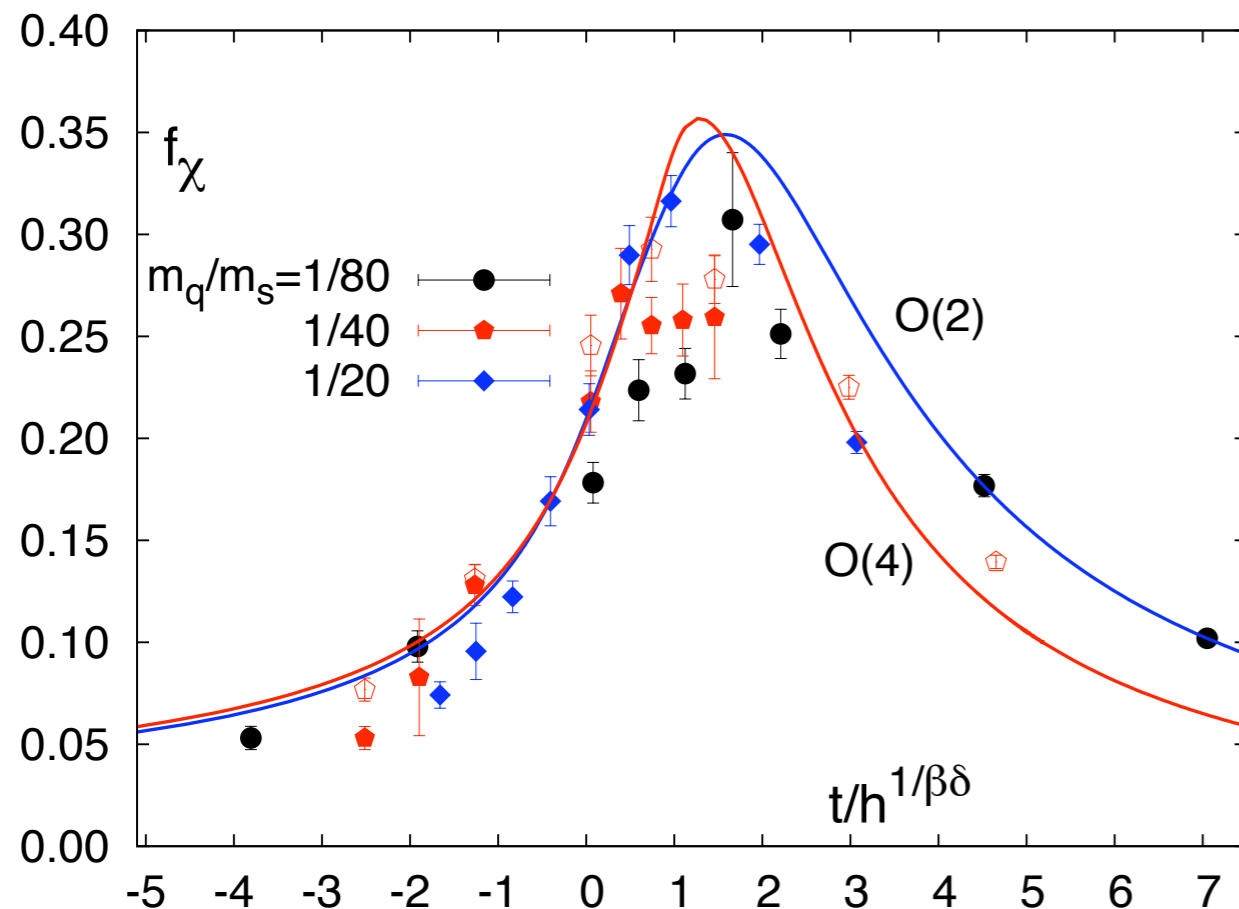
\rightarrow **further studies are needed to control continuum limit**



- scaling function:

$$f_{\chi}(z) = \frac{1}{\delta} \left(f_G(z) - \frac{z}{\beta} f'_G(z) \right)$$

- chiral susceptibilities scale reasonably well
- f_{χ} more sensitive to universality class, however, statistics still not sufficient



- mixed susceptibility:

$$\chi_t \equiv \frac{\partial M}{\partial T} = \frac{1}{t_0 T_c} \frac{\partial M}{\partial t} = \frac{1}{t_0 T_c} h^{(\beta-1)/\beta\delta} f'_G(z)$$

where

$$t = \frac{1}{t_0} \frac{T - T_c}{T_c} \quad h = \frac{H}{h_0}$$

(reduced temperature) (external field)

- introducing **chemical potential**:

$$t = \frac{1}{t_0} \left(\frac{T - T_c}{T_c} + \kappa_\mu \left(\frac{\mu_l}{T} \right)^2 \right)$$

in the chiral limit: μ_l does not break chiral symmetry
 \rightarrow couples only to reduced temperature

- (other) mixed susceptibility:

$$c_2^{\bar{\psi}\psi} \equiv \frac{\partial^2 M}{\partial(\mu_l/T)^2} \Big|_{\mu_l=0} = \frac{2\kappa_\mu}{t_0 T_c} \frac{\partial M}{\partial t} = \frac{2\kappa_\mu}{t_0 T_c} h^{(\beta-1)/\beta\delta} f'_G(z)$$

$$\propto \chi_t$$

- curvature of critical line in the ciral limit:

$$t = 0 \quad \longleftrightarrow \quad \frac{T}{T_c} = 1 - \kappa_\mu \left(\frac{\mu_l}{T} \right)^2$$

t_0, h_0, T_c known from scaling analysis of magnetic EoS

- fit $2\kappa_\mu f'_G(z)$ to χ_t -data (one fit parameter)
- **preliminary** result from fit to O(2) scaling curve:

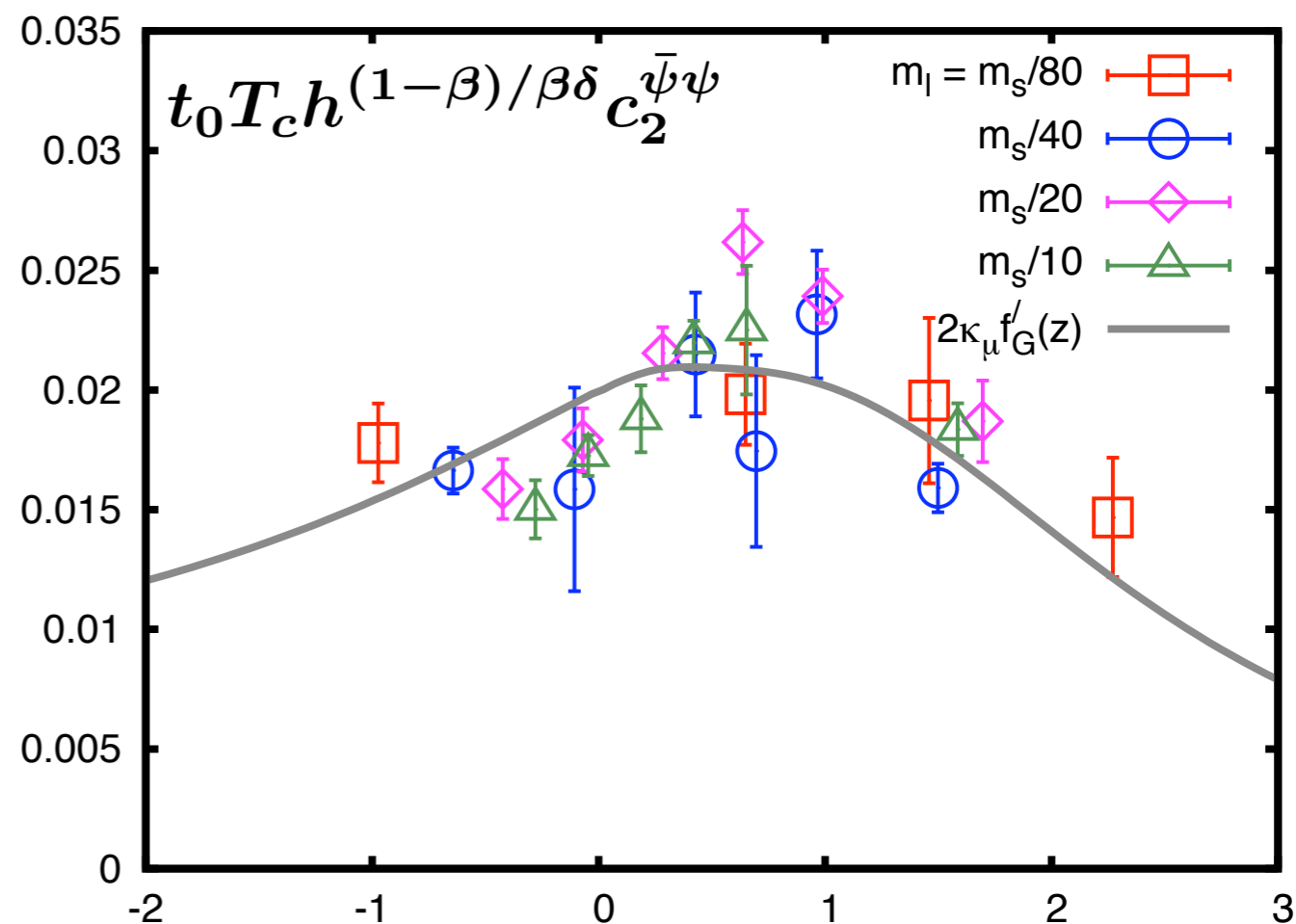
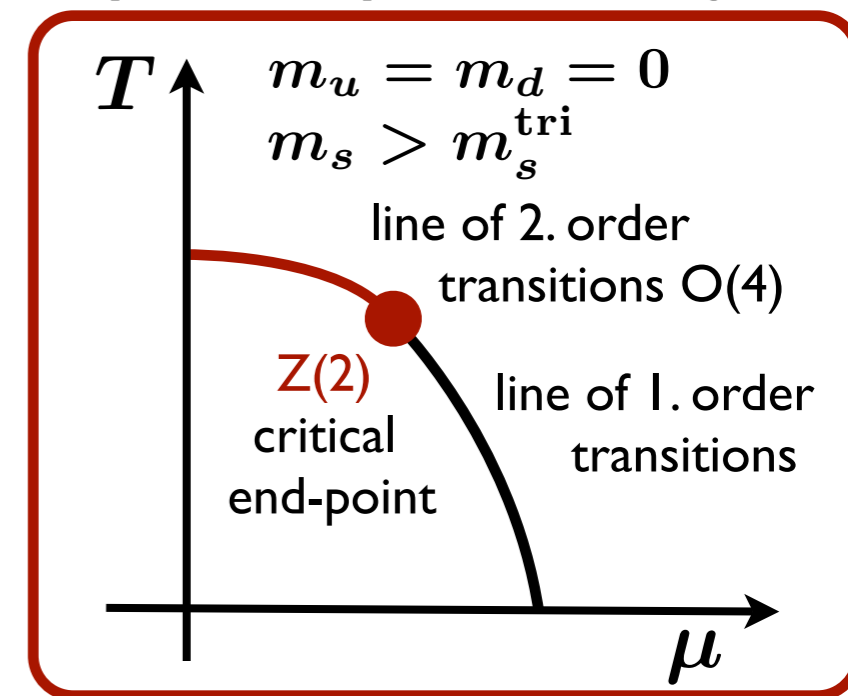
$$\kappa_\mu = 0.035(1)$$

- **for orientation:** reweighting std. action, $m_l/m_s = 1/27$

$$\kappa_\mu = 0.0288(9)$$

Z. Fodor and S.D. Katz,
JHEP 0404 (2004)

expected phase diagram



★ The magnetic EoS

- EoS consistent with 3d-O(N) scaling already at physical masses
- we find no evidence for nearby first order phase transitions

★ Scaling of chiral susceptibilities

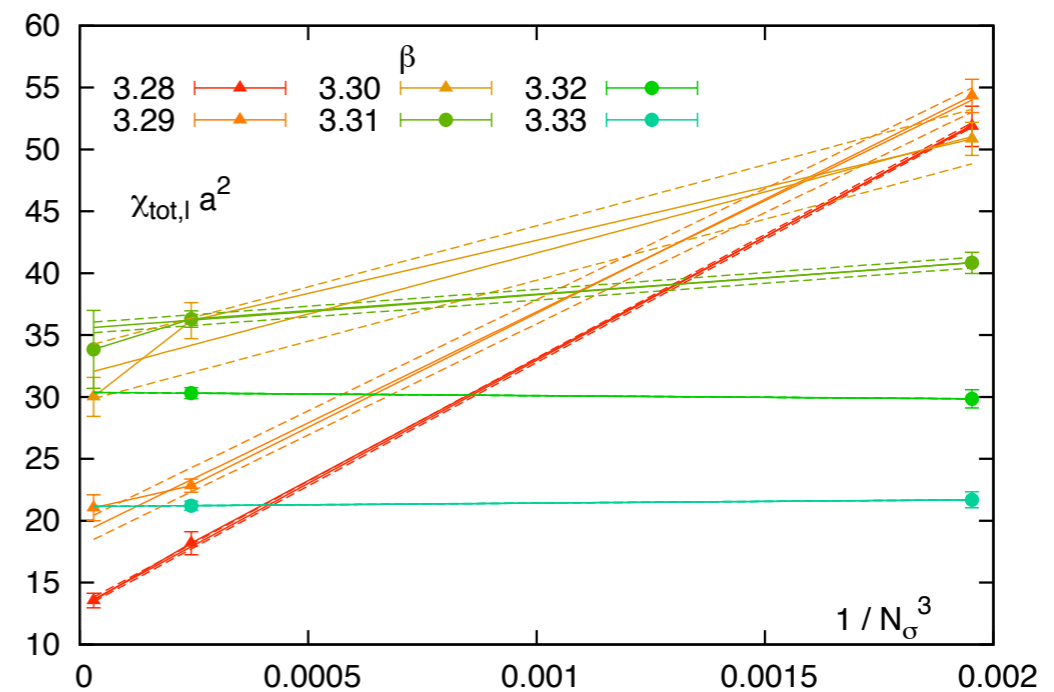
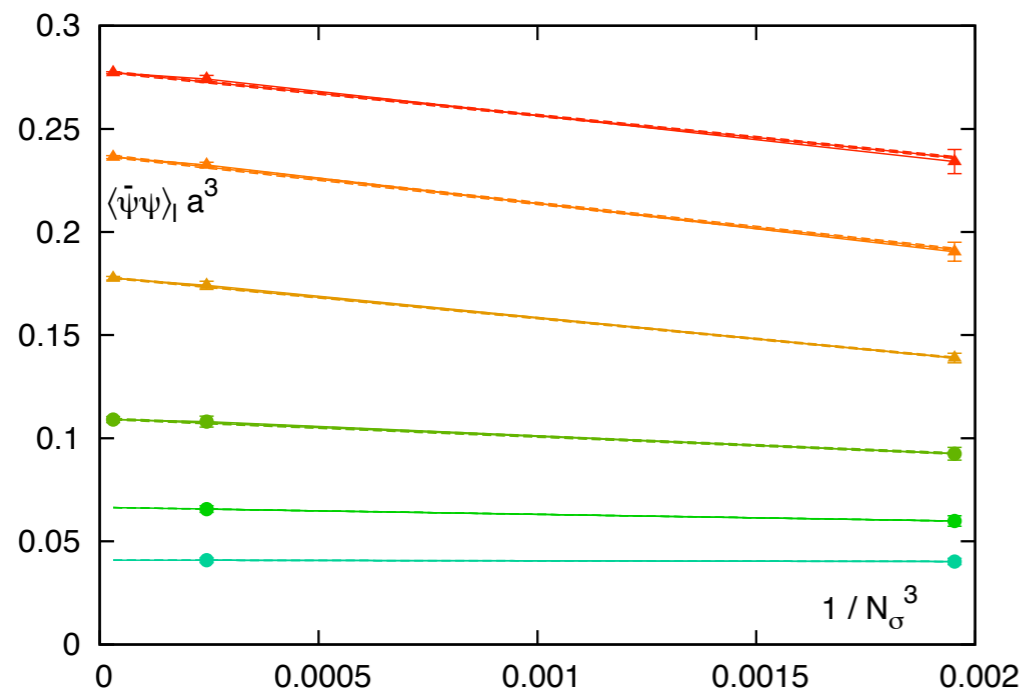
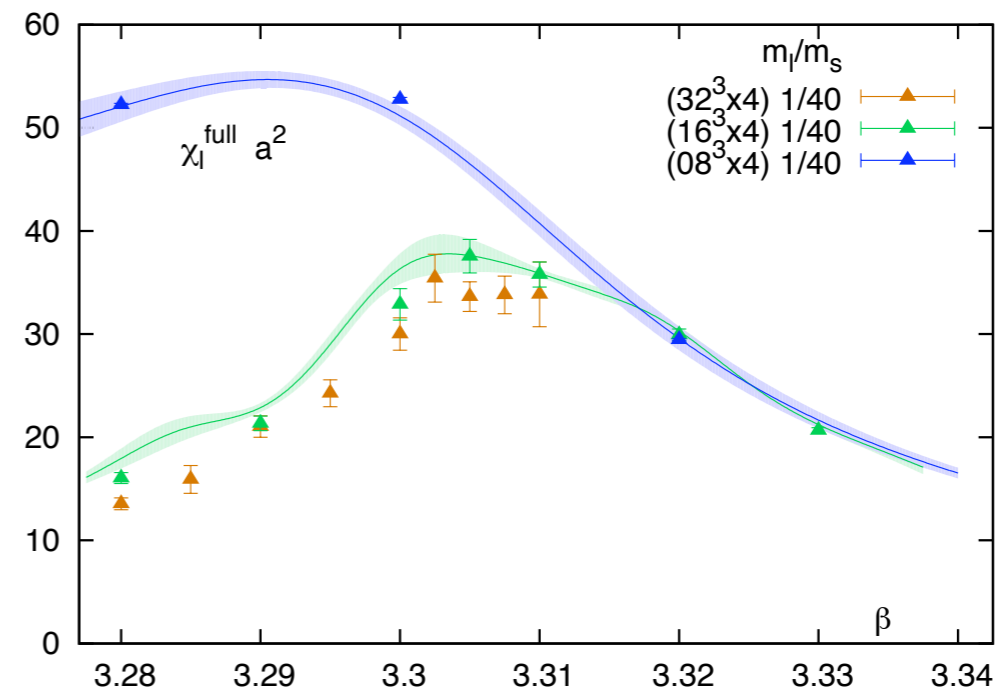
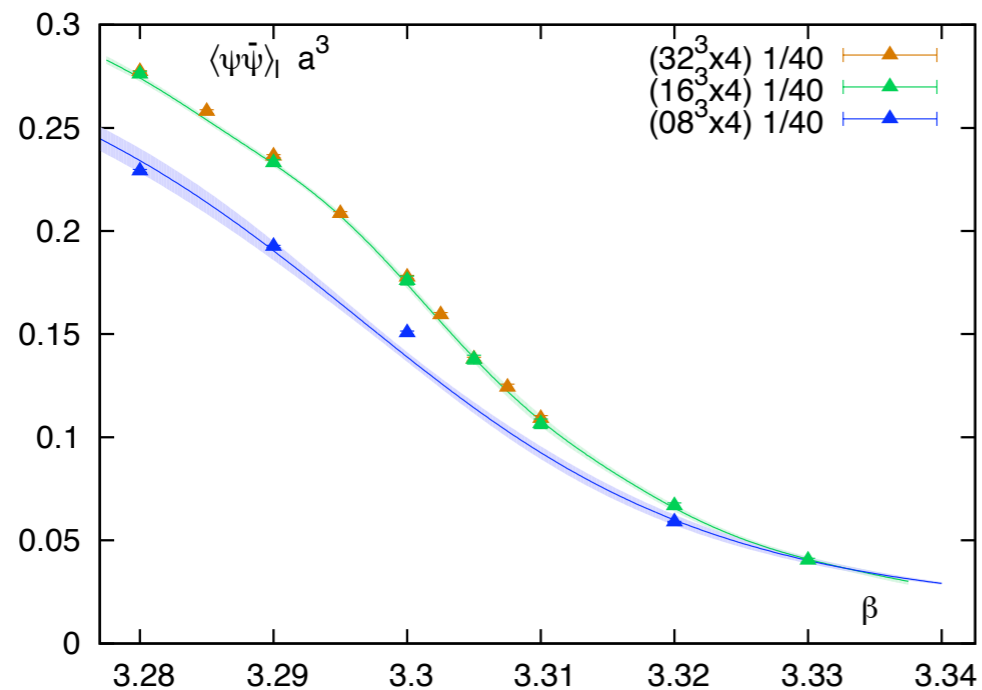
- statistics not (yet) sufficient to discriminate between universality classes

★ Non-zero chemical potentials

- the curvature of the critical line in the chiral limit can be extracted from scaling properties of mixed susceptibilities

→ all this needs to be confirmed in the continuum limit

- no evidence for finite size scaling \rightarrow crossover
- no strong temperature dependence of $1/V$ corrections
- thermodynamic limit well under control (for smallest mass: $m_\pi L \simeq 3$)



- scaling function:

$$f_{\chi}(z) = \frac{1}{\delta} \left(f_G(z) - \frac{z}{\beta} f'_G(z) \right)$$

- full susceptibilities scale reasonably well (after subtraction of regular part)
- f_{χ} more sensitive to universality class, however, statistics still not sufficient

