

# The Aoki phase revisited

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# Summary

- ◉ Motivation
- ◉ Simulations of the Aoki phase
  - > First round: HMC algorithm
  - > Second round: MFA algorithm
  - > Third round: HMC + tiny external source
- ◉ Results and conclusions

(Notation)  $\longrightarrow \langle (i\bar{\psi}\gamma_5\psi)^q \rangle = \left\langle \left( \frac{1}{V} \sum_x i\bar{\psi}\gamma_5\psi \right)^q \right\rangle$

# Motivation

## Our work on the Aoki phase

- > V. Azcoiti, G. Di Carlo, E. Follana and A. Vaquero, Phys. Rev. D79, 014509 (2009).
- > S. Sharpe, Phys. Rev. D79, 054503 (2009).

**Standard wisdom**

$$\langle (i\bar{\psi}\gamma_5\psi)^{2n} \rangle = 0$$

$$\langle (i\bar{\psi}\gamma_5\tau_3\psi)^2 \rangle \neq 0$$

**Our claim**

$$\langle (i\bar{\psi}\gamma_5\psi)^{2n} \rangle \neq 0$$

$$\langle (i\bar{\psi}\gamma_5\tau_3\psi)^2 \rangle \neq 0$$

# Motivation

- Two possible scenarios
  - > New, unexpected Aoki-like phases appear
  - > Infinite tower of sum-rules

$$\langle (i\bar{\psi}\gamma_5\psi)^{2n} \rangle = 0 \quad \forall n \in \mathbb{N} \quad ???$$

- Simulation needed to choose an scenario

# Motivation

- Rich numerical work inside the Aoki phase for QCD
  - > S. Aoki
  - > S. Aoki, A. Gocksch
  - > S. Aoki, T. Kaneda, A. Ukawa, T. Umemura
  - > E.-M. Ilgenfritz, W. Kerler, M. Müller-Preussker, A. Sternbeck, H. Stüben
- Most of them performed with an external source  $i\bar{\psi}\gamma_5\tau_3\psi$

# Motivation

- Why use the external source approach
  - > Standard way of analyzing SSB
  - > Helps convergence of solver (regularizes the small eigenvalues)
- Problems of the external source method
  - > The term  $i\bar{\psi}\gamma_5\psi$  leads to the **sign problem**
  - > We need another approach to verify our claims

# Motivation

- ⊙ Problems of the  $i\bar{\psi}\gamma_5\tau_3\psi$  external source
  - > Usually selects a vacuum PRD79, 014509 (2009)
  - > First thermodynamic limit, then zero field limit
    - Too many simulations
    - Systematic errors in extrapolations
- ⊙ Solution  $\longrightarrow$  P.D.F. formalism
  - > Reconstructs the complete P.D.F. for the observables. See Phys. Lett. B354, 111, (1995)

# Motivation

- Quantities to be measured

- $i\bar{\psi}_u\gamma_5\psi_u$  Parity o.p. 1 flavour  $\langle (i\bar{\psi}_u\gamma_5\psi_u)^2 \rangle$
- $i\bar{\psi}\gamma_5\psi$  Parity o.p. 2 flavour  $\langle (i\bar{\psi}\gamma_5\psi)^2 \rangle$
- $i\bar{\psi}\gamma_5\tau_3\psi$  Aoki o.p.  $\langle (i\bar{\psi}\gamma_5\tau_3\psi)^2 \rangle$

- We always measure the second moment



# First round: Hmc algorithm

- The Aoki phase features small eigenvalues
  - > We need a good solver → GCR + SAP
    - See M. Lüscher, Comput. Phys. Commun. 156 (2004) 209
  - > We need a small stepsize → Autocorrelation
  - > Eigenvalue-crossing is not allowed
- HMC is not ergodic inside the Aoki phase (in the absence of an external source)

# First round: Hmc algorithm

- Results in a  $4^4$  lattice (2 Flavours)

- > Aoki point  $\beta=2.0$   $\kappa=0.25$
- > Physical point  $\beta=3.0$   $\kappa=0.22$

Simulation	$\langle (i\bar{\psi}_u \gamma_5 \psi_u)^2 \rangle$	$\langle (i\bar{\psi} \gamma_5 \psi)^2 \rangle$	$\langle (i\bar{\psi} \gamma_5 \tau_3 \psi)^2 \rangle$
Aoki Cold [0]	$(1.92 \pm 0.03) \times 10^{-2}$	$(2.68 \pm 0.12) \times 10^{-2}$	$(5.00 \pm 0.09) \times 10^{-2}$
Aoki Hot [0]	$(1.87 \pm 0.03) \times 10^{-2}$	$(2.70 \pm 0.05) \times 10^{-2}$	$(4.79 \pm 0.08) \times 10^{-2}$
Aoki Hot [1]	$(6.52 \pm 0.72) \times 10^{-3}$	$(-4.49 \pm 0.50) \times 10^{-2}$	$(7.10 \pm 0.30) \times 10^{-2}$
Physical [0]	$(2.100 \pm 0.002) \times 10^{-3}$	$(4.151 \pm 0.003) \times 10^{-3}$	$(4.246 \pm 0.005) \times 10^{-3}$

- > Statistics  $\sim 10000$  trajectories per point

# Second round: MFA

- MFA relies on quenched generation weighted by the determinant
  - Phys. Rev. Lett. 65, 2239, (1990)
    - > Eigenvalue crossing allowed
    - > Determinant fluctuations suppressed
    - > Poor sampling
- We need huge configuration ensembles to obtain meaningful measurements

# Second round: MFA

- Results in a  $4^4$  and a  $6^4$  lattice

Simulation	$\langle (i\bar{\psi}_u \gamma_5 \psi_u)^2 \rangle$	$\langle (i\bar{\psi} \gamma_5 \psi)^2 \rangle$	$\langle (i\bar{\psi} \gamma_5 \tau_3 \psi)^2 \rangle$
MFA $4^4$ [All]	$(2.25 \pm 0.16) \times 10^{-2}$	$(-1.9 \pm 2.8) \times 10^{-2}$	$(1.1 \pm 0.4) \times 10^{-1}$
MFA $4^4$ [0]	$(2.57 \pm 0.46) \times 10^{-2}$	$(-0.5 \pm 3.4) \times 10^{-2}$	$(1.0 \pm 0.5) \times 10^{-1}$
Hmc+Weights	$(1.75 \pm 0.17) \times 10^{-2}$	$(1.73 \pm 0.15) \times 10^{-2}$	$(5.3 \pm 0.6) \times 10^{-2}$
MFA $6^4$ [All]	$(4.6 \pm 1.0) \times 10^{-2}$	$(-1.3 \pm 4.0) \times 10^{-1}$	$(3.1 \pm 4.0) \times 10^{-1}$
MFA $6^4$ [0]	$(4.6 \pm 0.7) \times 10^{-2}$	$(6.4 \pm 0.5) \times 10^{-2}$	$(1.2 \pm 0.3) \times 10^{-1}$

- > Weights  $4^4$  [0] ~ 86.8% [1] ~ 13.2% Error 7.1%
- > Statistics  $4^4$  ~ 2140000 configurations  
 $6^4$  ~ 50000 configurations

# Third round: Hmc + tiny external source

- If the source is small enough...
  - > ...regularizes the inverse of the dirac operator, so it can not diverge
  - > ...this enhances eigenvalue crossing
  - > ...if the external field is small enough (of order  $1/V$ ), it does not select a vacuum
- HMC becomes ergodic again

# Results and Conclusions

- First time the Aoki phase without external source is simulated with dynamical fermions
  - > Errors too large at this moment, but...
  - > ...data consistent among different methods
  - > Need to perform simulations at larger volumes

Simulation	$\langle (i\bar{\psi}_u \gamma_5 \psi_u)^2 \rangle$	$\langle (i\bar{\psi} \gamma_5 \psi)^2 \rangle$	$\langle (i\bar{\psi} \gamma_5 \tau_3 \psi)^2 \rangle$
Aoki Cold [0]	$(1.92 \pm 0.03) \times 10^{-2}$	$(2.68 \pm 0.12) \times 10^{-2}$	$(5.00 \pm 0.09) \times 10^{-2}$
MFA 4 <sup>4</sup> [0]	$(2.6 \pm 0.5) \times 10^{-2}$	$(-0.5 \pm 3.4) \times 10^{-2}$	$(1.0 \pm 0.5) \times 10^{-1}$
MFA 6 <sup>4</sup> [0]	$(4.6 \pm 0.7) \times 10^{-2}$	$(6.4 \pm 0.5) \times 10^{-2}$	$(1.2 \pm 0.3) \times 10^{-1}$