Pion and
Kaon EM matrix elements

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## Matrix element of the electromagnetic operator between Kaon and pion states

Baum, Itzhak ${ }^{1}$ Lubicz, Vittorio ${ }^{2}$ Martinelli, Guido ${ }^{3}$ Simula, Silvano ${ }^{2}$

${ }^{1}$ Rome University "La Sapienza"
${ }^{2}$ University of Rome III and INFN - Roma Tre
${ }^{3}$ University of Rome "La Sapienza" and INFN Rome

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## Outline

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## Kaon rare semileptonic decays as new physics probes

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Rare Kaon decays have not been detected yet:

$$
B R\left(K_{L} \rightarrow \pi^{0} \ell^{ \pm} \ell^{\mp}\right)_{\exp }<6.6 \cdot 10^{-10}
$$

In the SM they are estimated to be

$$
\begin{gathered}
B R(K \rightarrow \pi e e)_{S M} \sim 1.5 \cdot 10^{-12} \\
B R(K \rightarrow \pi \mu \mu)_{S M} \sim 3 \cdot 10^{-10}
\end{gathered}
$$

New physics can be the leading contribution, mediated through the Electro-magnetic and Chromo-magnetic operators:

$$
\begin{aligned}
& \mathcal{Q}_{E M}^{+}=\bar{s} F_{\mu \nu} \sigma^{\mu \nu} d \\
& \mathcal{Q}_{C M}^{+}=\bar{s} G_{\mu \nu} \sigma^{\mu \nu} G_{\mu \nu} d
\end{aligned}
$$



■ Sensitive to hadronic matrix elements.

## Previous work

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First lattice calculation [Becirevic et al. 2001] of the EM form factor

$$
f_{T}\left(q^{2}=0\right)=0.77 \pm 0.06 \pm 0.03
$$

With the slope in $q^{2}$

$$
\lambda=1.21 \pm 0.05 \mathrm{GeV}^{-2}
$$

- Quenched ( $n_{f}=0$ )
- High pion masses $\left(530<m_{\pi}<800 \mathrm{MeV}\right)$
- One lattice size ( $\left.a^{-1}=2.7(1) \mathrm{GeV}\right)$


## Lattice details

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- ETMC lattice QCD simulations [ETMC 0701012, 0911.5061]
- Dynamical flavors: $n_{f}=2$

■ Pion mass range: $270<m_{\pi}<600 \mathrm{MeV}$

- Lattice sizes: $24^{3} \times 48$ and $32^{3} \times 64$

■ Lattice step sizes: $a=0.068,0.085,0.10 \mathrm{fm}$
■ Action is Symanzik tree-level improved with maximally twisted-mass Wilson fermions

■ Non perturbative renormalization in the $\mathrm{RI} / \mathrm{MOM}$ scheme [ETMC 1004.1115]

- 3-point correlators with all-to-all stochastic propagator calculation, increase accuracy
■ Breit momentum frame: $\vec{p}_{K}=\vec{p}, \vec{p}_{\pi}=-\vec{p}$


## Electromagnetic form factor calculation

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$$
\mathcal{Q}_{E M}=\bar{s} \sigma^{\mu \nu} d
$$

The EM form factor is acquired from the EM matrix element by [Becirevic et al. 2001]:

$$
\left\langle\frac{\pi^{0}}{\sqrt{2}}\right| \mathcal{Q}_{E M}\left|K^{0}\right\rangle=i\left(p_{K}^{\mu} p_{\pi}^{\nu}-p_{K}^{\nu} p_{\pi}^{\mu}\right) \frac{\sqrt{2} f_{T}}{m_{K}+m_{\pi}}
$$

To obtain the matrix elements from the 3-point correlators, we look at the lattice times far from the pion and Kaon sources

$$
C_{3}^{K \pi} \rightarrow \frac{\sqrt{Z_{K} Z_{\pi}}}{4 E_{K} E_{\pi}}\left\langle\pi^{0}\right| \mathcal{Q}_{E M}\left|K^{0}\right\rangle e^{-E_{K} t_{x}-E_{\pi}\left(t_{y}-t_{x}\right)}
$$

and use the ratio

$$
\frac{C_{3}^{K \pi} C_{3}^{\pi K}}{C_{2}^{\pi}\left(t_{y}\right) C_{2}^{K}\left(t_{y}\right)} \rightarrow \frac{\left\langle\pi^{0}\right| \mathcal{Q}_{E M}\left|K^{0}\right\rangle^{2}}{16 E_{K} E_{\pi}}
$$

where $t_{y}$ is a fixed point $t_{y}=T / 2$.

## Calculations

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■ Interpolation in momentum to $q^{2}=0$, assuming pole behaviour:

$$
f_{T}\left(q^{2}\right)=\frac{f_{T}(0)}{1-q^{2} \lambda}
$$

■ Interpolation to physical strange mass:

$$
\left(2 m_{K}^{2}-m_{\pi}^{2}\right)_{L A T T} \rightarrow\left(2 m_{K}^{2}-m_{\pi}^{2}\right)_{P H Y S} \propto\left(m_{s}\right)_{P H Y S}
$$




## Extrapolation in masses

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Extrapolation in $m_{\pi}^{2}$ to physical pion mass $m_{\pi^{0}}=135 \mathrm{MeV}$

- linear $f=A m_{\pi}^{2}+B$

■ quadratic $f=A^{\prime} m_{\pi}^{4}+B^{\prime} m_{\pi}^{2}+C^{\prime}$

- log-linear $f=A^{\prime \prime} m_{\pi}^{2} \ln \left(m_{\pi}^{2}\right)+B^{\prime \prime} m_{\pi}^{2}+C^{\prime \prime}$


- Small finite volume effects


## Results

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Other lattice spacings:





## Results

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We find (Preliminary results, no systematic effects)

$$
\begin{gathered}
f_{T}\left(q^{2}=0\right)=0.430 \pm 0.066^{\text {stat }} \\
\lambda=1.61 \pm 0.41^{\text {stat }} \mathrm{GeV}^{-2} \text { (slope in } q^{2}, \text { pole fit) }
\end{gathered}
$$

To compare with [Becirevic et al. 2001] (linear fit)

$$
\begin{gathered}
f_{T}(0)=0.77 \pm 0.06 \pm 0.03 \\
\lambda=1.21 \pm 0.05 \mathrm{GeV}^{-2}
\end{gathered}
$$

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Confronting results for $m_{K}=m_{\pi}$ for similar lattice sizes $a$


Similar behaviour, difference may be due to:

- Extrapolation from large pion masses

■ Quenching effects

## Work in progress

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- Electro-magnetic operator $\mathcal{Q}_{E M}^{+}=\bar{s} \sigma^{\mu \nu} d$
- Combined fit for all lattice spacings
- Systematic errors analysis (chiral extrapolation, momentum dependence)
■ Chromo-magnetic operator $\mathcal{Q}_{C M}^{+}=\bar{s} G_{\mu \nu} \sigma^{\mu \nu} d$
- No previous lattice calculation
- Matrix elements calculation
- Renormalization

$$
\mathcal{Q}_{C M}^{\text {renorm }}=Z_{C M}\left(\mathcal{Q}_{C M}^{\text {bare }}+\frac{c}{a} \mathcal{Q}_{S}\right)
$$

- additive - subtraction of mixing with scalar operator
- multiplicative - 1-loop lattice perturbation theory $[\mathrm{H}$. Panagopoulos et al.]


## Conclusions

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■ Previously, single calculation (2001) of the EM operator

- Our calculations were performed for a large range of masses and lattice spacings
■ Higher statistical accuracy achieved
- Values at $q^{2}=0$ differ, may be due to either quenching or smaller pion masses
- Slope in $q^{2}$ is consistent with previous result, but with higher preliminary statistical error
■ Chromo-magnetic operator is work-in-progress

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Thank you!

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