# EOS in 2+1 flavor QCD with improved Wilson quarks by the fixed scale approach

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### QCD Equation of State on the lattice

Most studies done with staggerd-type quarks

- less computational costs
- a part of chiral sym. preserved ...

 $\rightarrow$  N<sub>f</sub>=2+1, almost physical quark mass,  $\mu \neq 0$ 

■ 4th-root trick to remove unphysical "tastes"
 → non-locality "universality is not guaranteed"

It is important to cross-check with theoretically sound lattice quarks like Wilson-type quarks

Our aim is to investigate

QCD Thermodynamics with Wilson-type quarks

WHOT-QCD Collaboration

### Fixed scale approach to study QCD thermodynamics

Temperature  $T = 1/(N_t a)$  is varied by  $N_t$  at fixed a



Advantages

- Line of Constant Physics
- T=0 subtraction for renorm.
- larger 1/a in whole T region

#### Disadvantages

- T resolution by integer N<sub>t</sub>
- Statistics in lower T region
- coding for odd  $\ensuremath{\mathsf{N}}_t$

#### T-integration method to calculate the EOS

#### We propose a new method ("T-integration method") to calculate the EOS at fixed scales

T.Umeda et al. (WHOT-QCD), Phys.Rev.D79 (2009) 051501(R)

Our method is based on the trace anomaly (interaction measure),

$$\frac{\epsilon - 3p}{T^4} = \left(\frac{N_t^3}{N_s^3}\right) a \frac{d\beta}{da} \left\langle \frac{dS}{d\beta} \right\rangle_{sub}$$

and the thermodynamic relation.

$$\frac{\epsilon - 3p}{T^4} = T \frac{\partial (p/T^4)}{\partial T}$$
$$\implies \frac{p}{T^4} = \int_0^T dT' \ \frac{\epsilon - 3p}{T'^5}$$

## Test in quenched QCD



[\*] G. Boyd et al., NPB469, 419 (1996)

- Our results are roughly consistent with previous results.
- Our results deviate from the fixed N<sub>t</sub>=8 results [\*] at higher T (aT~0.3 or higher)
- Trace anomaly is sensitive to spatial volume at lower T (below T<sub>c</sub>).
   V > (2fm)<sup>3</sup> is ncessarry.

#### Lattice setup

■ T=0 simulation: on 28<sup>3</sup> x 56 by CP-PACS/JLQCD Phys. Rev. D78 (2008) 011502

- RG-improved Iwasaki glue + NP-improved Wilson quarks

- 
$$\beta = 2.05$$
,  $\kappa_{ud} = 0.1356$ ,  $\kappa_s = 0.1351$ 

- V~(2 fm)<sup>3</sup>, a=0.07 fm, ( $m_{\pi} \sim 634 \text{MeV}, \ \frac{m_{\pi}}{m_{\rho}} = 0.63, \ \frac{m_{\eta_s s}}{m_{\phi}} = 0.74$ )

- configurations available on the ILDG/JLDG

T>0 simulations: on  $32^3 \times N_t$  (N<sub>t</sub>=4, 6, ..., 14, 16) lattices

RHMC algorithm, same parameters as T=0 simulation



#### Formulation for $N_f=2+1$ improved Wilson quarks

$$\begin{split} S &= S_g + S_q \\ S_g &= -\beta \left\{ \sum_{x,\mu>\nu} c_0 W_{\mu\nu}^{1\times1}(x) + \sum_{x,\mu,\nu} c_1 W_{\mu\nu}^{1\times2}(x) \right\} \\ S_q &= \sum_{f=u,d,s} \sum_{x,y} \bar{q}_x^f D_{x,y} q_y^f \\ D_{x,y} &= \delta_{x,y} - \kappa_f \sum_{\mu} \{(1 - \gamma_{\mu}) U_{x,\mu} \delta_{x+\bar{\mu},y} + (1 + \gamma_{\mu}) U_{x-\bar{\mu},\mu}^{\dagger} \delta_{x-\bar{\mu},y} \} - \delta_{x,y} c_{SW} \kappa_f \sum_{\mu>\nu} \sigma_{\mu\nu} F_{\mu\nu} \\ c_{SW}(\beta) &= 1 + 0.113g^2 + 0.0209g^4 + 0.0049g^6 \\ Phys. Rev. D73, 034501 \\ CP-PACS/JLQCD \\ \hline \frac{\epsilon - 3p}{T^4} &= \frac{N_t^3}{N_s^3} \left( a \frac{\partial \beta}{\partial a} \left\langle \frac{\partial S}{\partial \beta} \right\rangle_{sub} + a \frac{\partial \kappa_{ud}}{\partial a} \left\langle \frac{\partial S}{\partial \kappa_{ud}} \right\rangle_{sub} + a \frac{\partial \kappa_s}{\partial a} \left\langle \frac{\partial S}{\partial \kappa_s} \right\rangle_{sub} \right) \\ \hline \left\langle \frac{\partial S}{\partial \beta} \right\rangle &= N_s^3 N_t \left( - \left\langle \sum_{x,\mu>\nu} c_0 W_{\mu\nu}^{1\times1}(x) + \sum_{x,\mu,\nu} c_1 W_{\mu\nu}^{1\times2}(x) \right\rangle + N_f \frac{\partial c_{SW}}{\partial \beta} \kappa_f \left\langle \sum_{x,\mu>\nu} Tr^{(c,s)} \sigma_{\mu\nu} F_{\mu\nu} (D^{-1})_{x,x} \right\rangle \right) \\ \hline \left\langle \frac{\partial S}{\partial \kappa_f} \right\rangle &= N_f N_s^3 N_t \left( \left\langle \sum_{x,\mu} Tr^{(c,s)} \{(1 - \gamma_{\mu}) U_{x,\mu} (D^{-1})_{x+\bar{\mu},x} + (1 + \gamma_{\mu}) U_{x-\bar{\mu},\mu}^{\dagger} (D^{-1})_{x-\bar{\mu},x} \} \right) \\ + c_{SW} \left\langle \sum_{x,\mu>\nu} Tr^{(c,s)} \sigma_{\mu\nu} F_{\mu\nu} (D^{-1})_{x,x} \right\rangle \right) \end{split}$$

Noise method ( #noise = 1 for each color & spin indices )

Trace anomaly needs Beta-functions in  $N_f=2+1$  QCD

$$\frac{\epsilon - 3p}{T^4} = \frac{N_t^3}{N_s^3} \left( a \frac{\partial \beta}{\partial a} \left\langle \frac{\partial S}{\partial \beta} \right\rangle_{sub} + a \frac{\partial \kappa_{ud}}{\partial a} \left\langle \frac{\partial S}{\partial \kappa_{ud}} \right\rangle_{sub} + a \frac{\partial \kappa_s}{\partial a} \left\langle \frac{\partial S}{\partial \kappa_s} \right\rangle_{sub} \right)$$

Direct fit method Phys. Rev. D64 (2001) 074510

fit 
$$\beta$$
,  $\kappa_{ud}$ ,  $\kappa_s$  as functions of  $(am_{\rho})$ ,  $\left(\frac{m_{\pi}}{m_{\rho}}\right)$ ,  $\left(\frac{m_{\eta_{ss}}}{m_{\phi}}\right)$ 

$$\begin{pmatrix} \beta \\ \kappa_L \\ \kappa_S \end{pmatrix} = \vec{c}_1 + \vec{c}_2 (m\rho) + \vec{c}_3 (am\rho)^2 + \vec{c}_4 \left(\frac{m\pi}{m\rho}\right) + \vec{c}_5 \left(\frac{m\pi}{m\rho}\right)^2 + \vec{c}_6 (am\rho) \left(\frac{m\pi}{m\rho}\right) \\ = \vec{c}_7 \left(\frac{m\eta_{ss}}{m\phi}\right) + \vec{c}_8 \left(\frac{m\eta_{ss}}{m\phi}\right)^2 + \vec{c}_9 (am\rho) \left(\frac{m\eta_{ss}}{m\phi}\right) + \vec{c}_{10} \left(\frac{m\pi}{m\rho}\right) \left(\frac{m\eta_{ss}}{m\phi}\right) \\ = \frac{\partial X}{\partial (am\rho)} \quad \text{with fixed } \left(\frac{m\pi}{m\rho}\right), \left(\frac{m\eta_{ss}}{m\phi}\right) \quad (X = \beta, \kappa_{ud}, \kappa_s)$$

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Meson spectrum by CP-PACS/JLQCD Phys. Rev. D78 (2008) 011502.

3 ( $\beta$ ) x 5 ( $\kappa_{ud}$ ) x 2 ( $\kappa_s$ ) = 30 data points



### Trace anomaly in $N_f=2+1$ QCD



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#### Trace anomaly in $N_f=2+1$ QCD



#### Quark mass dependence of Trace anomaly



### Equation of State in $N_f=2+1$ QCD



T-integration

 $\frac{p}{T^4} = \int_0^T dT' \frac{\epsilon - 3p}{T'^5}$ 

is performed by the trapezoidal rule.

•  $\varepsilon$  /T<sup>4</sup> is calculated from

$$\frac{\epsilon}{T^4} = \frac{\epsilon - 3p}{T^4} + \frac{3p}{T^4}$$

Large error in whole T region

#### Summary & outlook

We reported on EOS in  $N_f=2+1$  QCD using improve Wilson quarks

- Beta functions More work needed Reweighting method ?
- Equation of state

More statistics needed in the lower temperature region

■ N<sub>f</sub>=2+1 QCD just at the physical point

the physical point (pion mass ~ 140MeV) by PACS-CS

Finite density

We can combine our approach with the Taylor expansion method, to explore EOS at  $\mu \neq 0$ 

# Back-up slides

#### A systemtic error



#### Meson spectrum by CP-PACS/JLQCD Phys. Rev. D78 (2008) 011502.





$$\left(a\frac{\partial\beta}{\partial a}, a\frac{\partial\kappa_{ud}}{\partial a}, a\frac{\partial\kappa_s}{\partial a}\right)_{\text{simulation point}}$$
  
=  $(-0.330(3), 0.00288(5), 0.00247(5)) m_{
ho}$ 

$$= (-0.340(3), 0.00286(5), 0.00242(5)) \quad m_{K^*}$$

$$= (-0.345(3), 0.00285(5), 0.00242(5)) \quad m_{\phi}$$

### Trace anomaly ( $\beta \& \kappa$ derivative part)

$$\frac{\epsilon - 3p}{T^4} = \frac{N_t^3}{N_s^3} \left( a \frac{\partial \beta}{\partial a} \left\langle \frac{\partial S}{\partial \beta} \right\rangle_{sub} + a \frac{\partial \kappa_{ud}}{\partial a} \left\langle \frac{\partial S}{\partial \kappa_{ud}} \right\rangle_{sub} + a \frac{\partial \kappa_s}{\partial a} \left\langle \frac{\partial S}{\partial \kappa_s} \right\rangle_{sub} \right)$$



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